

The Niels Bohr International Academy

VILLUM FONDEN

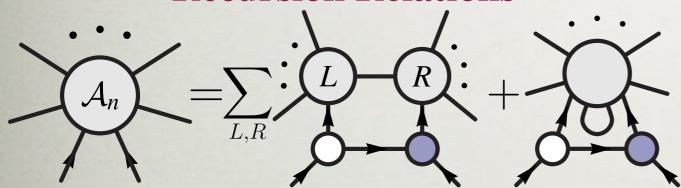




◆ Loop *integrands* are becoming "easy" to construct —(exposing and) preserving much simplicity

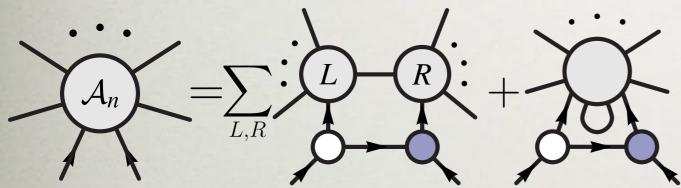
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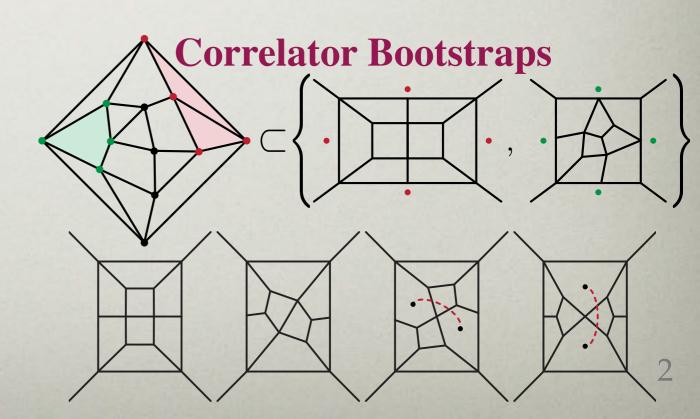
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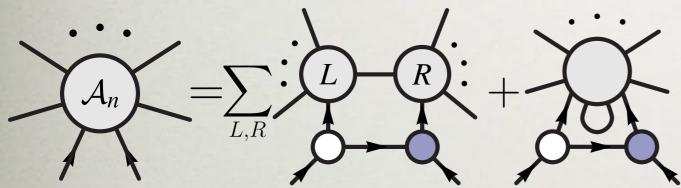
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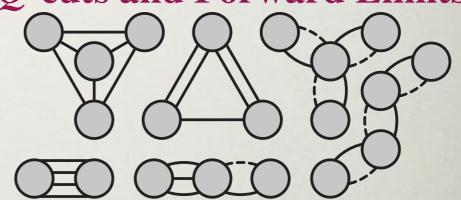


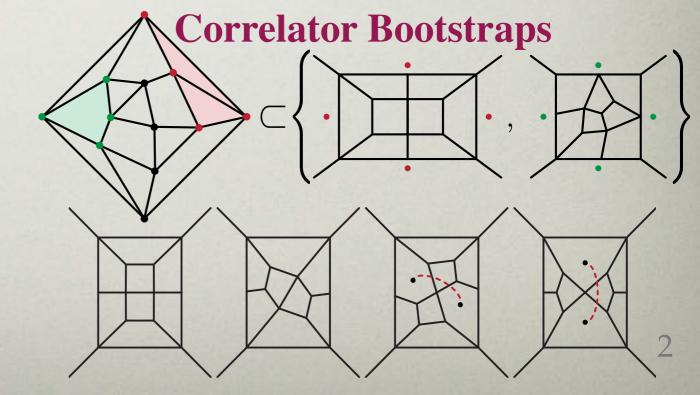
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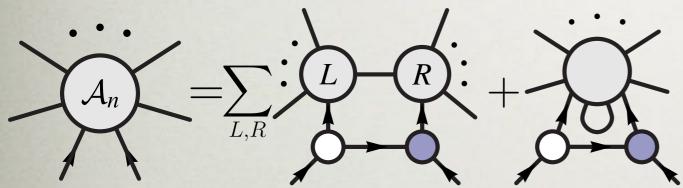
Q-cuts and Forward Limits



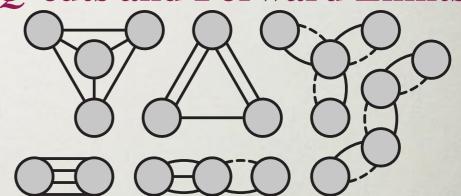


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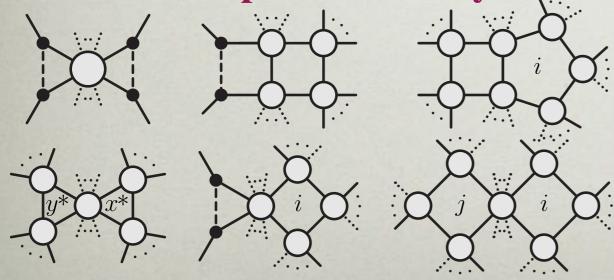
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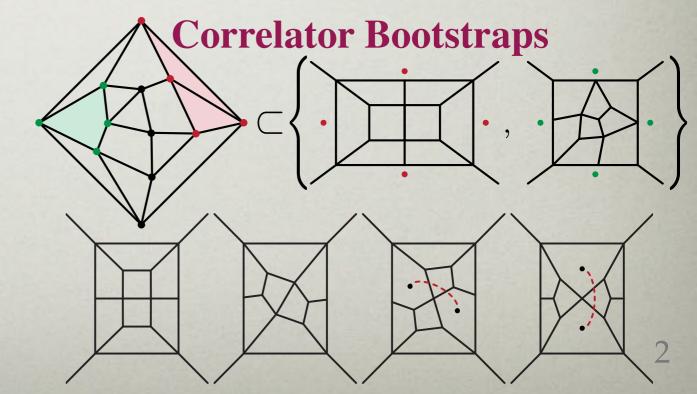


Q-cuts and Forward Limits



Prescriptive Unitarity





◆ Exempli gratia: we now have closed formulae for all amplitude integrands in planar SYM through 3 loops: (ask me about non-planar, non-SUSY!

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$$\mathcal{A}_n^{L=2} = \sum_{\mathcal{L}} f_{\mathcal{L}}$$

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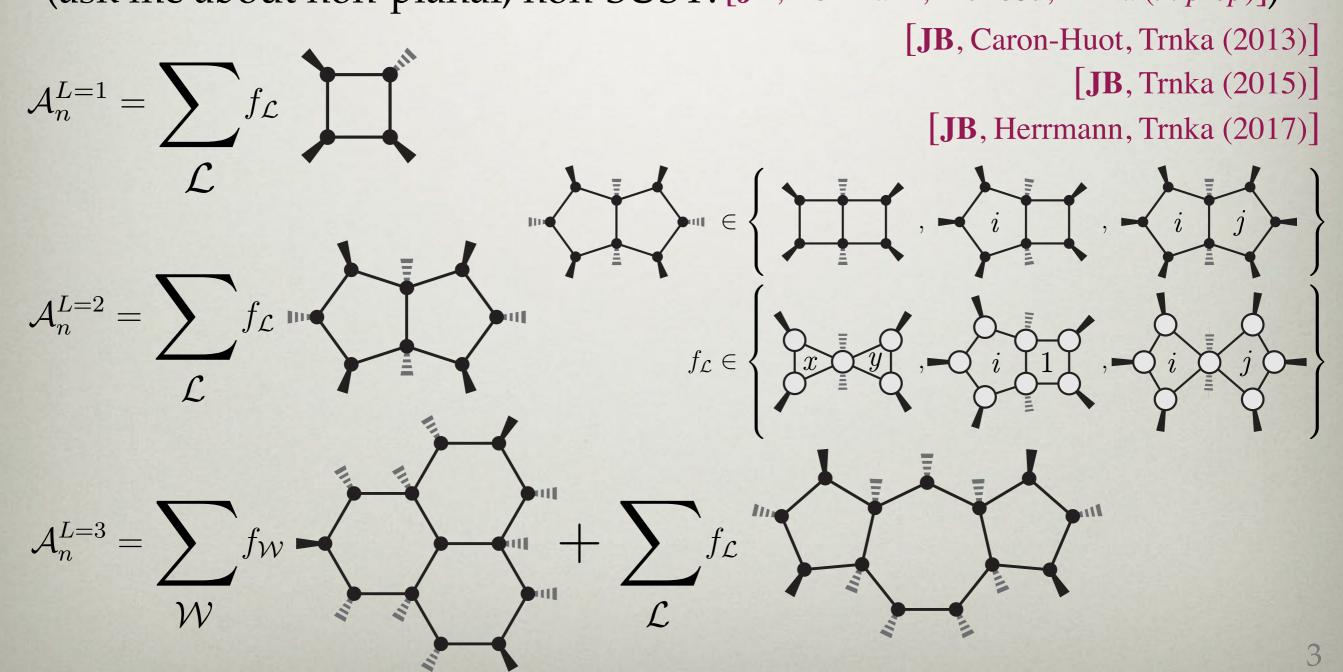
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[JB, Herrmann, Trnka (2017)]

$$\mathcal{A}_n^{L=2} = \sum_{\mathcal{L}} f_{\mathcal{L}}$$
 and $\mathcal{A}_n^{L=3} = \sum_{\mathcal{W}} f_{\mathcal{W}}$ and $\mathcal{L}_n^{L=3} = \sum_{\mathcal{L}} f_{\mathcal{L}}$

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Roadmap: Polylogs to Traintracks

- * Spiritus Movens(/Loop Integration Polemics)
 When has an integrand been integrated?
- **◆ Integrating Loop Integrals Rationally**
 - ► dual-conformal sufficiency [JB, Dixon, Dulat, Panzer (to appear)]
 - momentum twistor reducibility

[JB, McLeod, von Hippel, Wilhelm (2018)]

- Ubiquity of Non-Polylogarithmicity
 - integrals beyond (even elliptic) polylogarithms

[JB, McLeod, Spradlin, von Hippel, Wilhelm (2017)]

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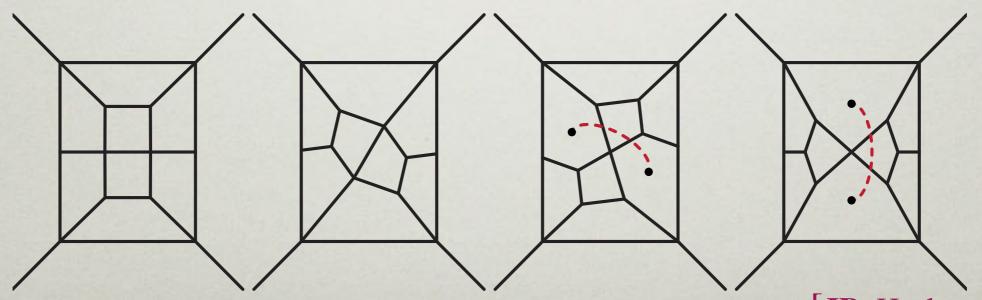
$$\left(\frac{3}{2}\zeta_3 - \pi^2\log(2) + \zeta_2 + \frac{197}{72}\right)$$

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Are these "numbers" MZVs?



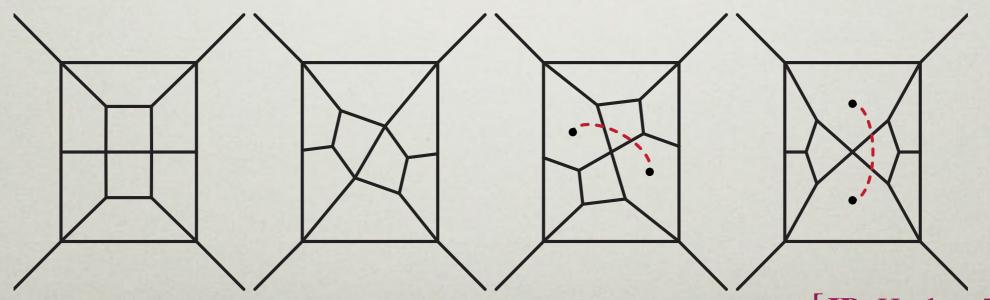
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implications for BES...

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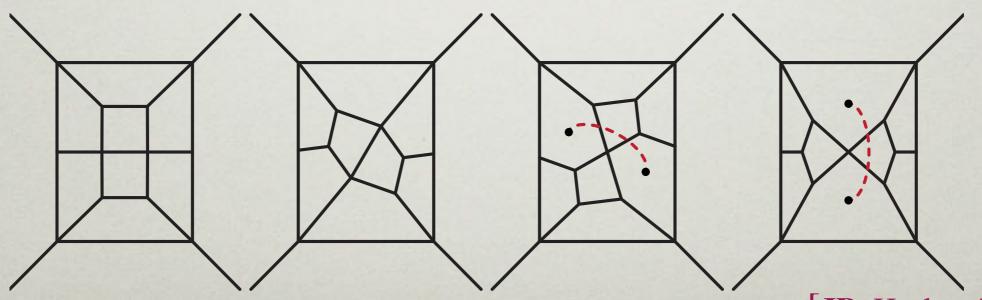
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Are these "numbers" MZVs? YES!

[O. Schnetz (private corr.)]



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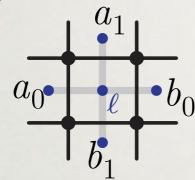
$$\mathcal{A}_n^{L=2,\mathrm{MHV}} = \sum_{a < b < c < d < a} \frac{1}{c} \frac{1}{b}$$

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$$=\frac{(\ell_1,N_1)(\ell_2,N_2)}{(\ell_1,a)(\ell_1,a+1)(\ell_1,b)(\ell_1,b+1)(\ell_1,\ell_2)(\ell_2,c)(\ell_2,c+1)(\ell_2,d)(\ell_2,d+1)}$$

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When the result is a *function*, this is more subtle—depending on various (often valid) criteria

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+ Certifiability

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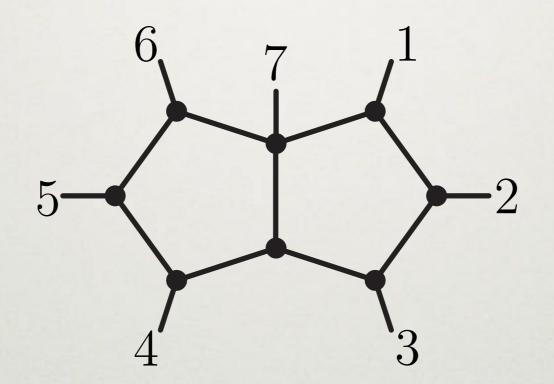
 $= \int_{0}^{1} d^{4}\vec{\alpha} \frac{1}{f_{1}f_{2}g_{2}} = \int_{0}^{1} \frac{ds}{\sqrt{4s^{3} - g_{2}s - g_{3}}} H_{3}(s)$

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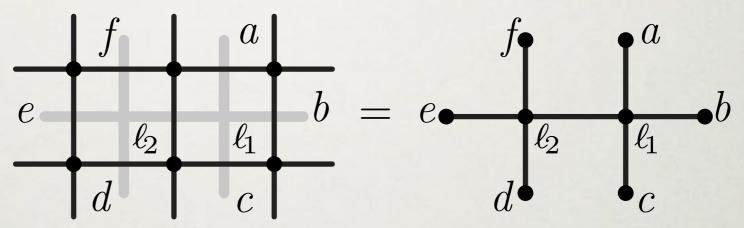
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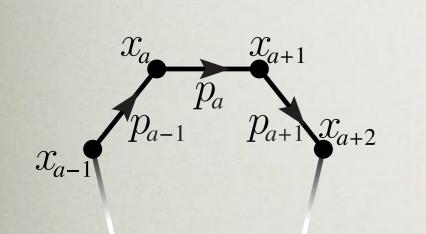
◆ Parameterize kinematic variables using: momentum twistors

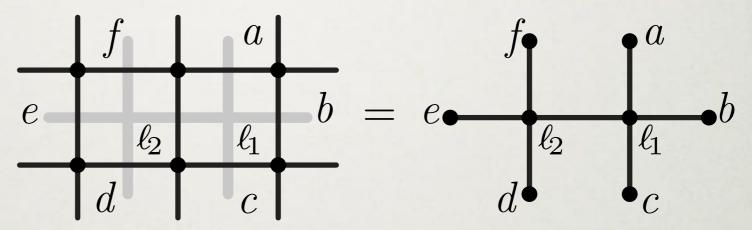
[JB, McLeod, von Hippel, Wilhelm (2018)]

chosen non-redundantly

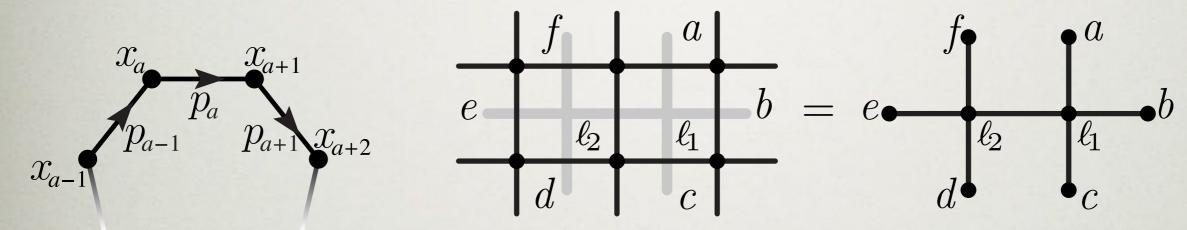
→ Partial fraction to death (e.g. use HyperInt) [Panzer (2014)]





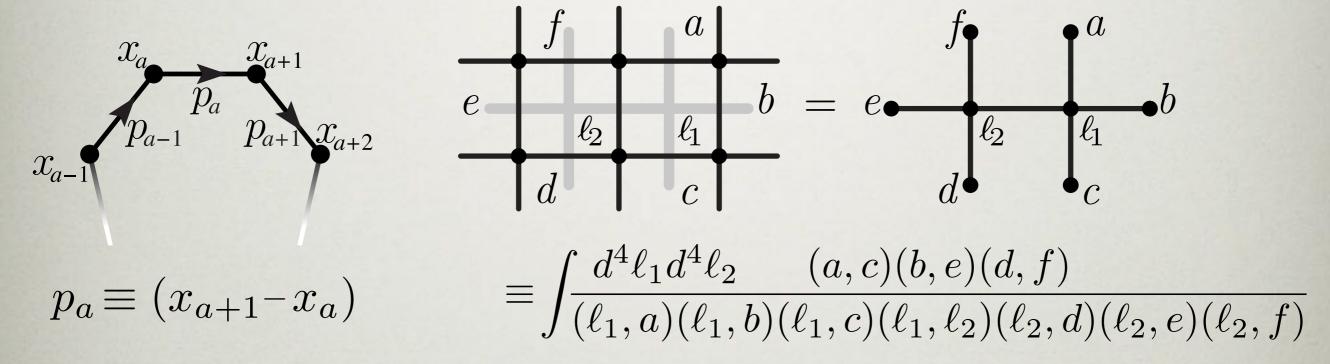


$$p_a \equiv (x_{a+1} - x_a)$$



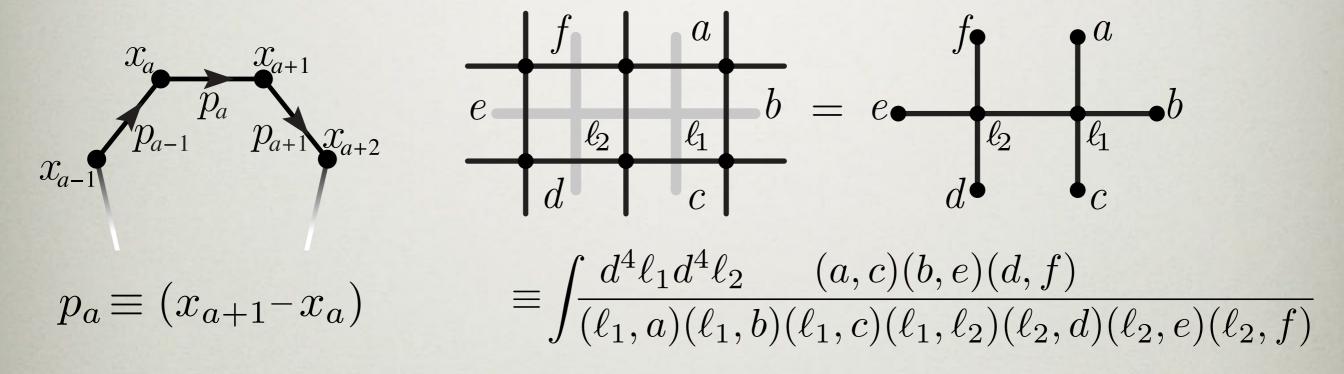
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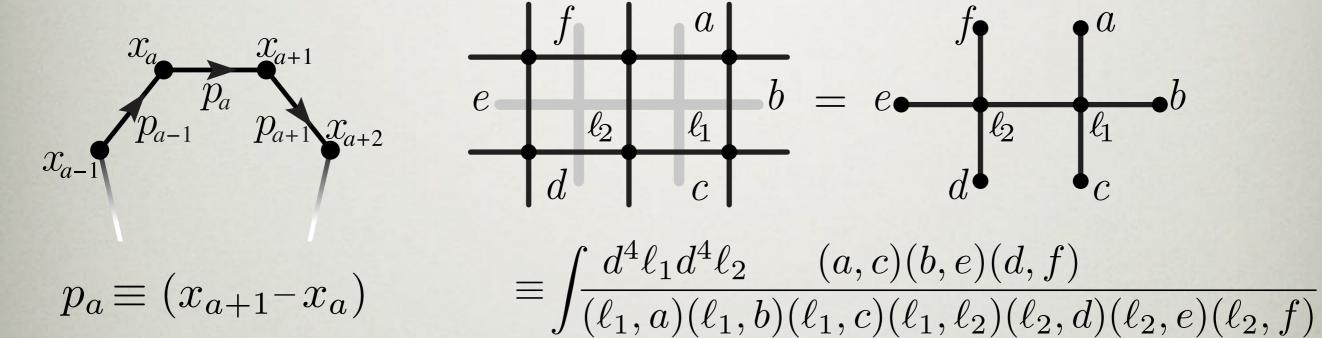
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 \bullet Dual-Conformal Invariance is conformality in x's

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$$(ab;cd) \equiv \frac{(a,b)(c,d)}{(a,c)(b,d)}$$

[Drummond, Henn, Smirnov, Sokatchev; Drummond, Korchemsky, Henn; ...]

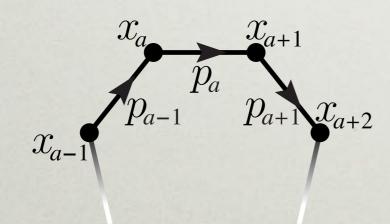
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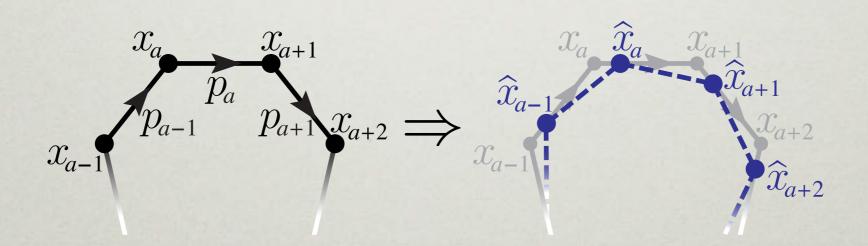
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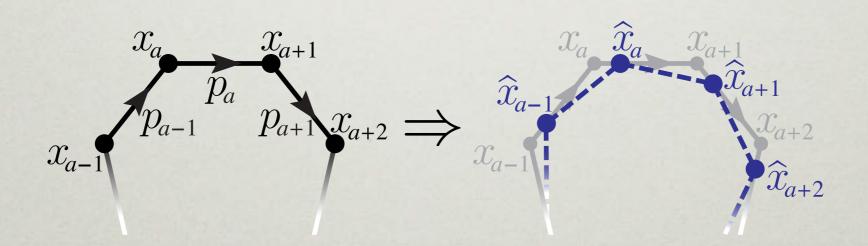
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 Coefficients of each divergence can be obtained as strictly finite (Feynman-) parametric integrals which can always be rendered manifestly DCI

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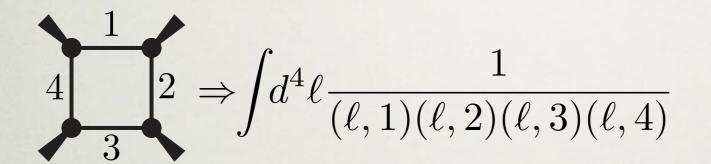
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$$\mathcal{F} \equiv \alpha_1 \alpha_2(1,2) + \alpha_2 \alpha_3(2,3) + \alpha_1 \alpha_3(1,3) + \alpha_1 \alpha_4(1,4) + \alpha_2 \alpha_4(2,4) + \alpha_3 \alpha_4(3,4)$$

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$$\begin{array}{c}
1 \\
4 \\
\hline
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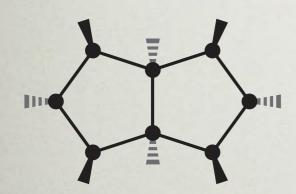
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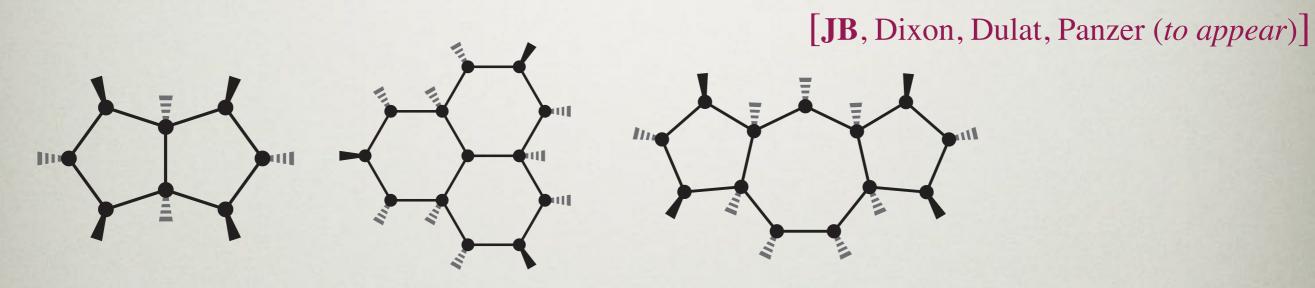
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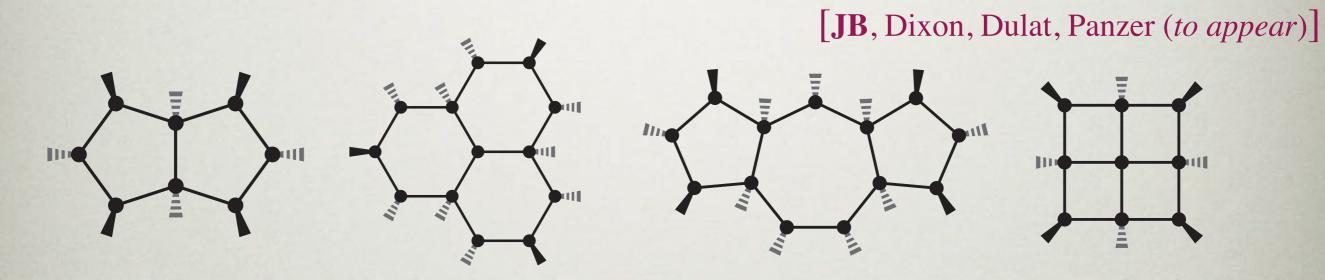
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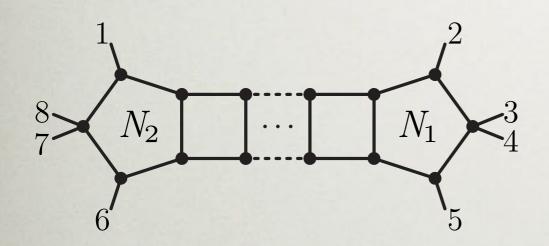
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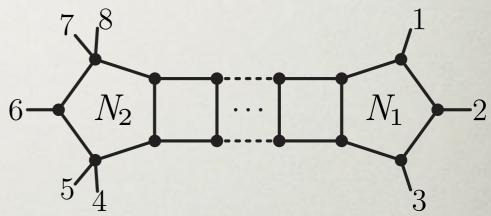


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$$\begin{array}{c}
 & 1 \\
 & N_{2} \\
 & N_{2} \\
 & N_{2} \\
 & N_{1}
\end{array}$$

$$\begin{array}{c}
 & 3 \\
 & 5 \\
 & M_{1}
\end{array}$$

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 & 3 \\
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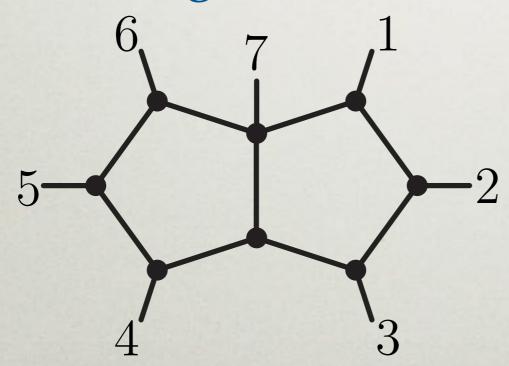
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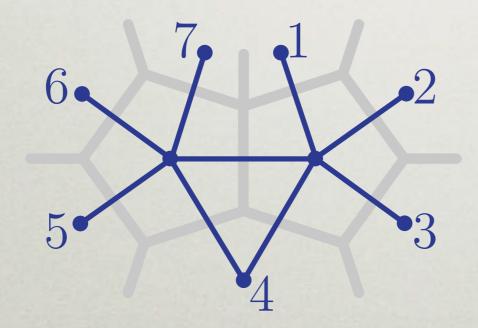
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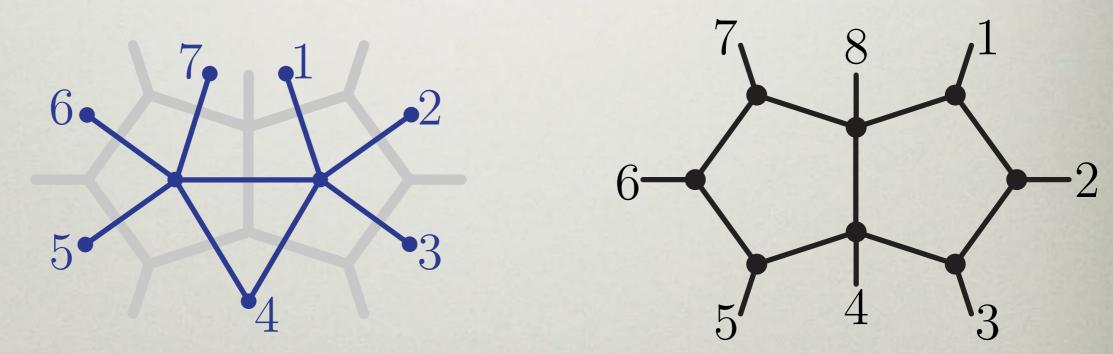
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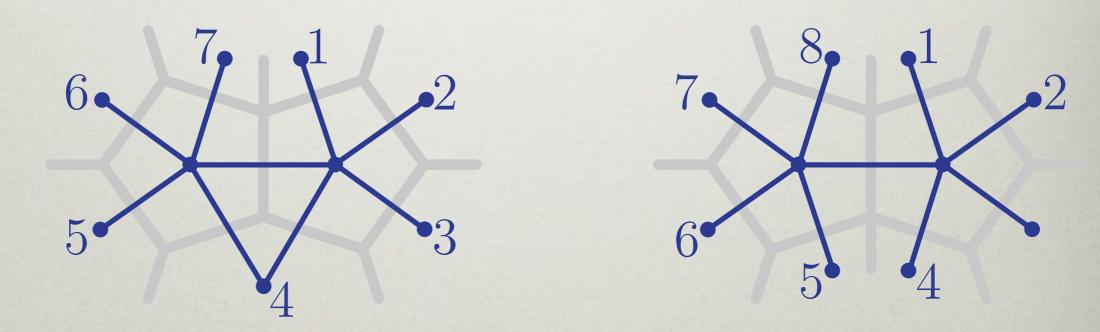
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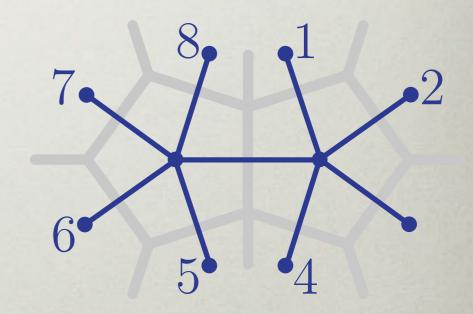
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rescaling-independent cross ratios: n(n-5)/2

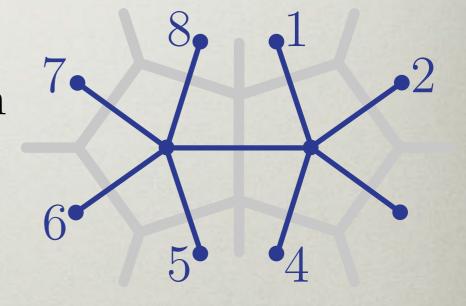
actually independent cross ratios: 3n-15

- ◆ Although a good start, we haven't yet eliminated all conformal redundancies—just the rescalings
 - —which is to say that parity-even cross-ratios are:
 - too great in number
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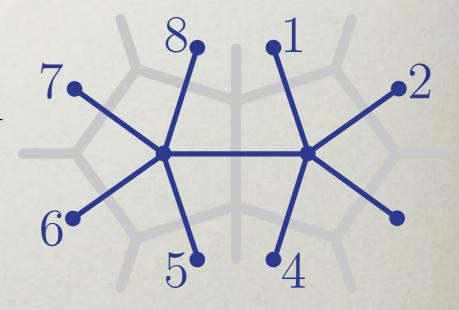
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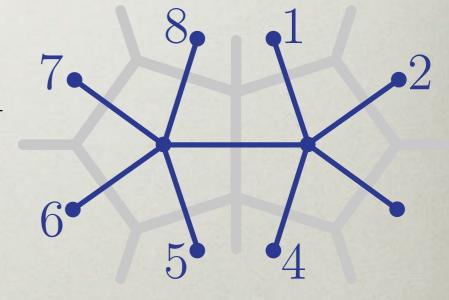
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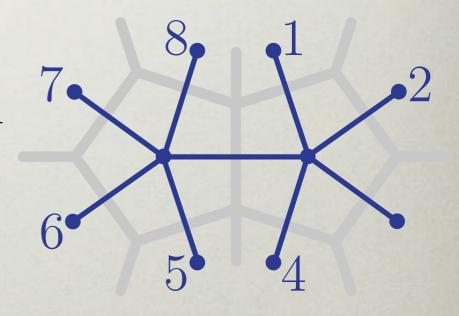
$$\sqrt{(1-u-v-w)^2-4uvw}$$

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- satisfy (complex) algebraic relations

$$\sqrt{(1-u-v-w)^2-4uvw}$$

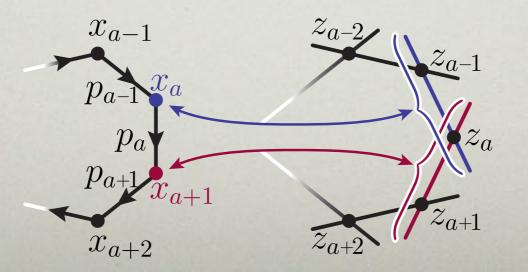


rescaling-independent cross ratios: n(n-5)/2

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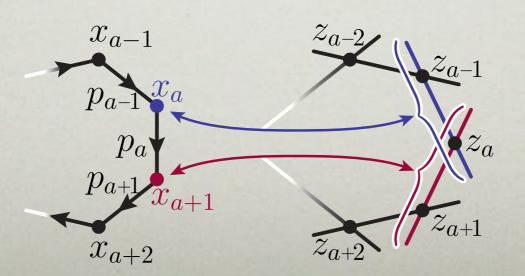
Momentum-Twistor Magic

Unsurprisingly (to most of us), momentum twistors are (closer to) the right kind of conformal variables



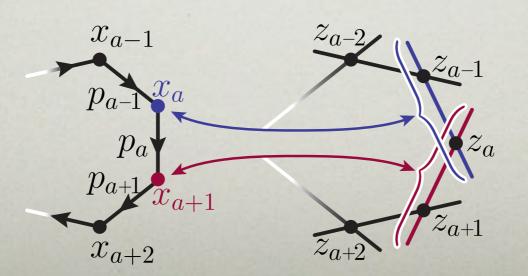
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- Unsurprisingly (to most of us), momentum twistors are (closer to) the right kind of conformal variables
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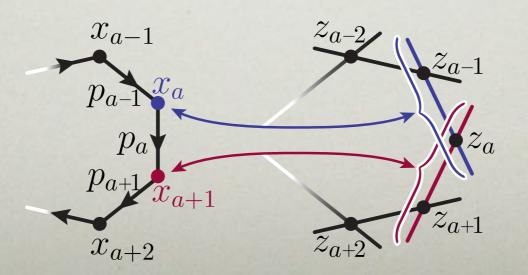


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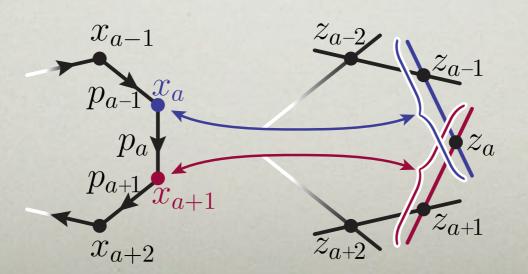
- manifest the rank of the Gramian
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 - manifest the rank of the Gramian
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 - rationalize all 6x6 Gram determinants



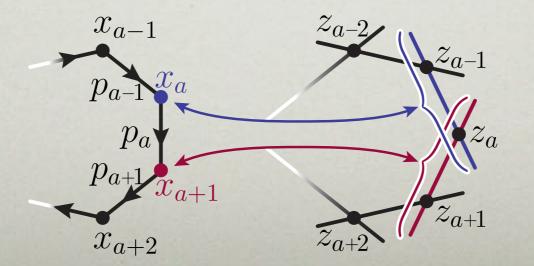
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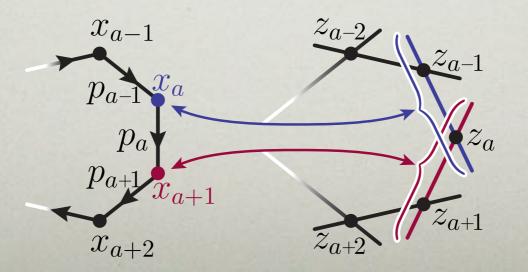
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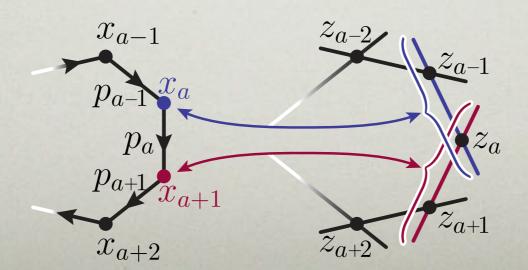


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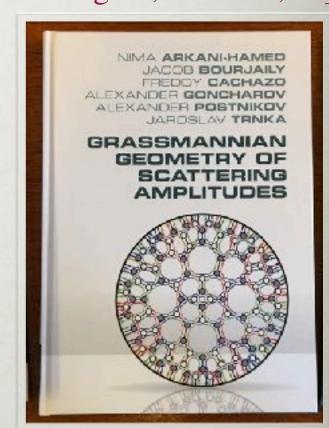






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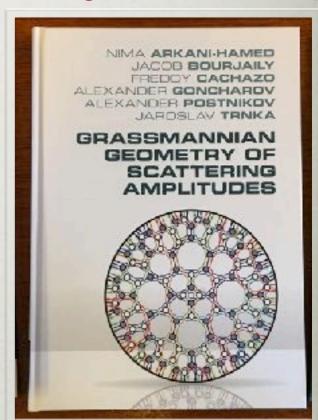


$$(1\ 0\ 0\ e_4\ 0\ -1\ -1-e_2(1+e_3)\ -1-(e_2+e_5)(1+e_3))$$
 $(1\ 1\ 0\ 0\ 0\ 0\ -e_2e_3\ -e_3(e_2+e_5))$
 $(1\ 1\ 0\ 0\ 0\ 1\ 1+e_1e_4\ 1\ 1\ e_1\ e_1\ e_1$

JB, McLeod, von Hippel, Wilhelm (2018)

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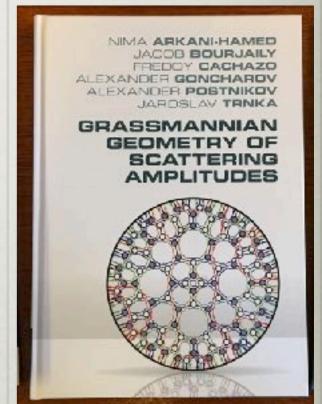


$$Z_{\rm B}^{(8)} \equiv \begin{pmatrix} 1 & 0 & 0 & e_4 & 0 & -1 & -1 - e_2 (1 \\ 1 & 1 & 0 & 0 & 0 & 0 & -e_2 e_3 \\ 0 & 0 & 1 & 1 + e_1 e_4 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 & 1 + e_1 & e_1 \end{pmatrix}$$

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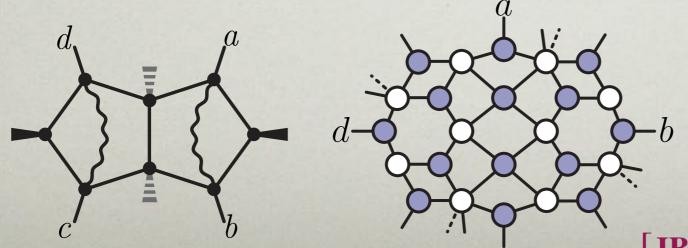


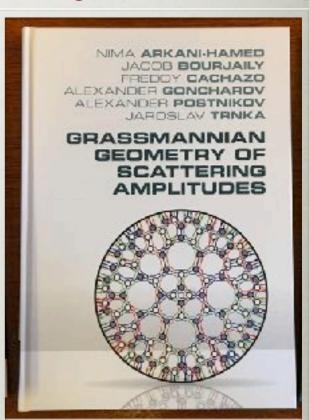
$$Z_{\rm B}^{(8)} \equiv \begin{pmatrix} 1 & 0 & 0 & e_4 & 0 & -1 & -1 - e_2(1 + e_3) & -1 - (e_2 + e_5)(1 + e_3) \\ 1 & 1 & 0 & 0 & 0 & 0 & -e_2 e_3 & -e_3(e_2 + e_5) \\ 0 & 0 & 1 & 1 + e_1 e_4 & 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 1 + e_1 & e_1 & e_1 \end{pmatrix}$$

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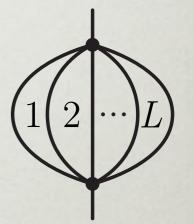




◆ Despite their ubiquity at low multiplicity and low loop orders, iterated polylogarithms are far from the only class of integrals that are needed in QFT

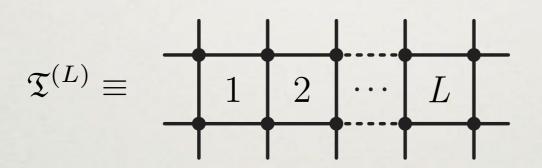
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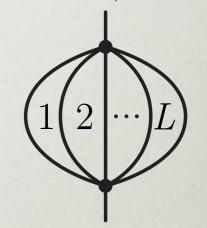
[Bloch, Kerr, Vanhove; Broadhurst;...]



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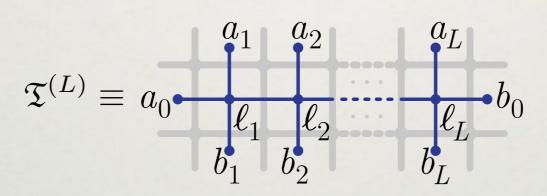
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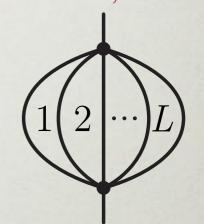




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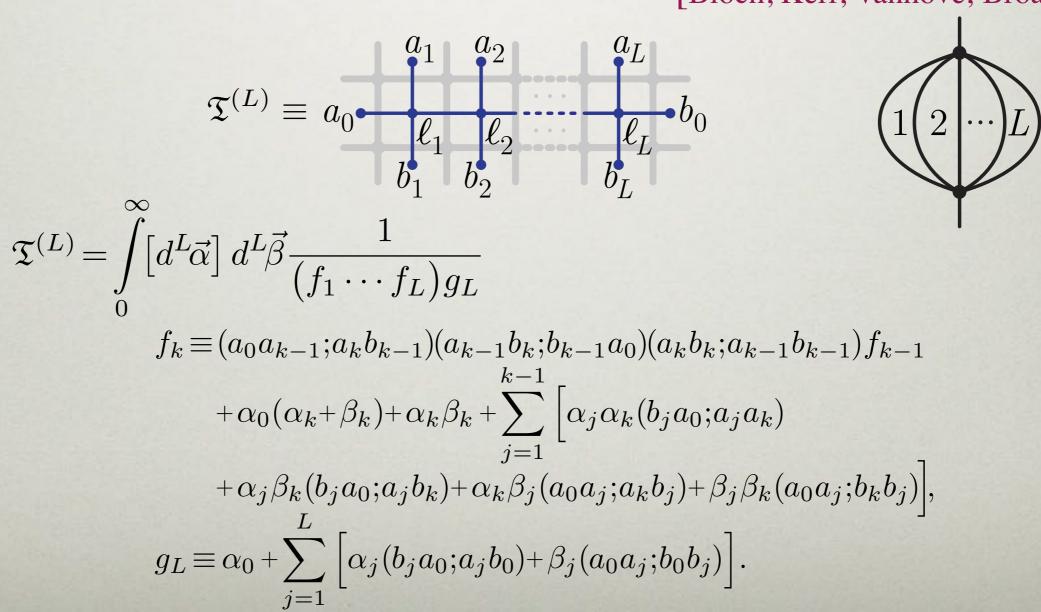




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[JB, He, McLeod, von Hippel, Wilhelm (2018)] Bloch, Kerr, Vanhove; Broadhurst;...]

$$\mathfrak{T}^{(L)} \equiv a_0 \qquad a_1 \qquad a_2 \qquad a_L \\ b_1 \qquad b_2 \qquad b_L \qquad b_0 \\ \mathfrak{T}^{(L)} = \int_0^\infty \left[d^L \vec{\alpha} \right] d^L \vec{\beta} \frac{1}{(f_1 \cdots f_L)g_L} = \int_{\sqrt{4x^3 - g_2(\vec{z})x - g_3(\vec{z})}}^{dx \, d^{L-2}\vec{z}} G'(x, \vec{z}),$$

$$f_k \equiv (a_0 a_{k-1}; a_k b_{k-1})(a_{k-1} b_k; b_{k-1} a_0)(a_k b_k; a_{k-1} b_{k-1}) f_{k-1} \\ + \alpha_0 (\alpha_k + \beta_k) + \alpha_k \beta_k + \sum_{j=1}^{k-1} \left[\alpha_j \alpha_k (b_j a_0; a_j a_k) \right. \\ + \alpha_j \beta_k (b_j a_0; a_j b_k) + \alpha_k \beta_j (a_0 a_j; a_k b_j) + \beta_j \beta_k (a_0 a_j; b_k b_j) \right],$$

$$g_L \equiv \alpha_0 + \sum_{j=1}^{L} \left[\alpha_j (b_j a_0; a_j b_0) + \beta_j (a_0 a_j; b_0 b_j) \right].$$

 Despite their ubiquity at low multiplicity and low loop orders, iterated polylogarithms are far from the only class of integrals that are needed in QFT

$$\mathcal{A}(\overbrace{\varphi_{12},\ldots,\varphi_{12}}^{L+1},\varphi_{13},\overbrace{\varphi_{34},\ldots,\varphi_{34}}^{L+1},\varphi_{24},\varphi_{24}))$$
 [Bloch, Kerr, Vanhove; Broadhurst;...]
$$\mathcal{I}^{(L)} \equiv a_0 \underbrace{\begin{array}{c} a_1 & a_2 \\ b_1 & b_2 \end{array}}_{b_1} \underbrace{\begin{array}{c} a_L \\ b_L \end{array}}_{b_L} b_0$$
 [Bloch, Kerr, Vanhove; Broadhurst;...]
$$\mathcal{I}^{(L)} = \int_0^\infty [d^L \vec{\alpha}] \ d^L \vec{\beta} \frac{1}{(f_1 \cdots f_L)g_L} = \int_0^\infty \underbrace{\begin{array}{c} dx \ d^{L-2}\vec{z} \\ \sqrt{4x^3 - g_2(\vec{z})x - g_3(\vec{z})}} G'(x, \vec{z}),$$

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$$\mathcal{A}(\overbrace{\varphi_{12},\ldots,\varphi_{12}},\varphi_{13},\varphi_{13},\overbrace{\varphi_{34},\ldots,\varphi_{34}},\varphi_{24},\varphi_{24}) \qquad [\text{Bloch, Kerr, Vanhove; Broadhurst;}...]$$

$$\mathcal{I}^{(L)} = \underbrace{\int_{0}^{24} [d^L \vec{\alpha}] \ d^L \vec{\beta} \frac{1}{(f_1 \cdots f_L)g_L}}_{24} = \underbrace{\int_{0}^{24} \frac{dx \ d^{L-2} \vec{z}}{\sqrt{4x^3 - g_2(\vec{z})x - g_3(\vec{z})}} G'(x,\vec{z}),$$

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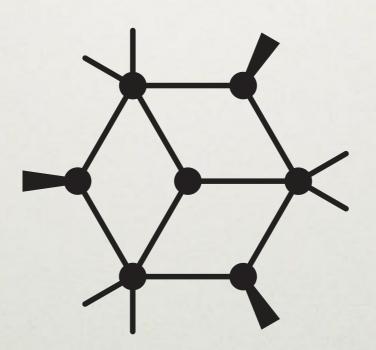
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$$16$$

Further Novelties Found...

◆ It turns out that traintracks do not saturate functional complexity at fixed loop-order...



[JB, McLeod, von Hippel, Wilhelm (in progress)]

Questions?