

Bootstrapping $4d \mathcal{N} = 2$ conformal theories

Madalena Lemos



IGST 2018
21 August 2018, Copenhagen University

together with C. Beem, M. Cornagliotto, P. Liendo, W. Peelaers, L. Rastelli,
V. Schomerus, B. van Rees

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 2$ theories: growing list of theories

The Superconformal Bootstrap Program

What is the space of consistent 4d SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories

The Superconformal Bootstrap Program

What is the space of consistent 4d SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories
 - many lacking a Lagrangian description,
 - isolated SCFTs

The Superconformal Bootstrap Program

What is the space of consistent 4d SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories
many lacking a Lagrangian description,
isolated SCFTs

Can we bootstrap specific theories?

The Superconformal Bootstrap Program

What is the space of consistent 4d SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories
 - many lacking a Lagrangian description,
 - isolated SCFTs

Can we bootstrap specific theories?

- “Simplest” $\mathcal{N} = 2$ Argyres-Douglas theory?

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_n(x_n) \rangle$

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_n(x_n) \rangle$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_k \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} \mathcal{O}_k(0)$$

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_n(x_n) \rangle$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k \text{ prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_n(x_n) \rangle$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k \text{ prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

→ Finite radius of convergence

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_n(x_n) \rangle$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k \text{ prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

- Finite radius of convergence
- n -point function by recursive use of the OPE until
 $\langle \mathbb{1} \rangle = 1$

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

Operator Product Expansion

$$\mathcal{O}_1(x) \mathcal{O}_2(0) = \sum_{k \text{ prim.}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

- Finite radius of convergence
- n -point function by recursive use of the OPE until
 $\langle \mathbb{1} \rangle = 1$

Conformal Bootstrap

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

Operator Product Expansion

$$\mathcal{O}_1(x) \mathcal{O}_2(0) = \sum_{k \text{ prim.}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

- Finite radius of convergence
- n -point function by recursive use of the OPE until
 $\langle \mathbb{1} \rangle = 1$

CFT data strongly constrained

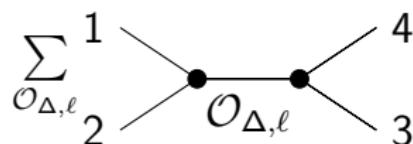
- ▶ Unitarity
- ▶ Associativity of the operator product algebra

$$(\mathcal{O}_1 \mathcal{O}_2) \mathcal{O}_3 = \mathcal{O}_1 (\mathcal{O}_2 \mathcal{O}_3)$$

Conformal Bootstrap

Crossing Symmetry

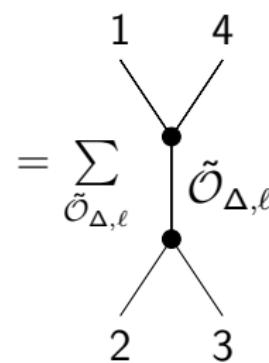
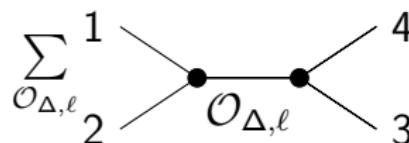
$$\langle (\mathcal{O}_1(x_1) \mathcal{O}_2(x_2)) \mathcal{O}_3(x_3) \mathcal{O}_4(x_4) \rangle =$$



Conformal Bootstrap

Crossing Symmetry

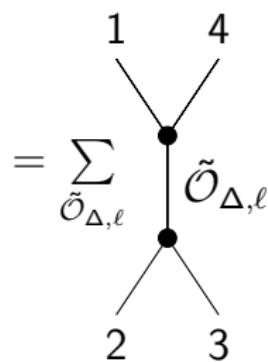
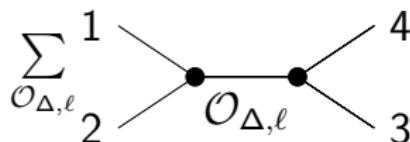
$$\langle \mathcal{O}_1(x_1) (\mathcal{O}_2(x_2) \mathcal{O}_3(x_3)) \mathcal{O}_4(x_4) \rangle =$$



Conformal Bootstrap

Crossing Symmetry

$$\langle (\mathcal{O}_1(x_1) \mathcal{O}_2(x_2)) \mathcal{O}_3(x_3) \mathcal{O}_4(x_4) \rangle =$$



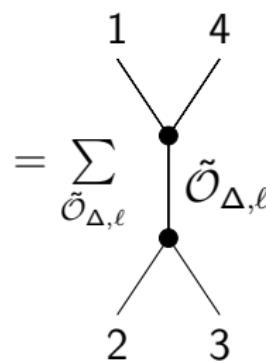
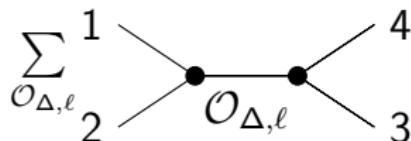
$$\frac{1}{x_{12}^{2\Delta_{\mathcal{O}}} x_{34}^{2\Delta_{\mathcal{O}}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}_3 \mathcal{O}_4 \mathcal{O}_{\Delta,\ell}} g_{\Delta,\ell}(z, \bar{z}) =$$

where $\Delta_{\mathcal{O}_i} = \Delta_{\mathcal{O}}$, $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2} = z\bar{z}$, $v = \frac{x_{23}^2 x_{14}^2}{x_{13}^2 x_{24}^2} = (1-z)(1-\bar{z})$

Conformal Bootstrap

Crossing Symmetry

$$\langle \mathcal{O}_1(x_1) (\mathcal{O}_2(x_2) \mathcal{O}_3(x_3)) \mathcal{O}_4(x_4) \rangle =$$

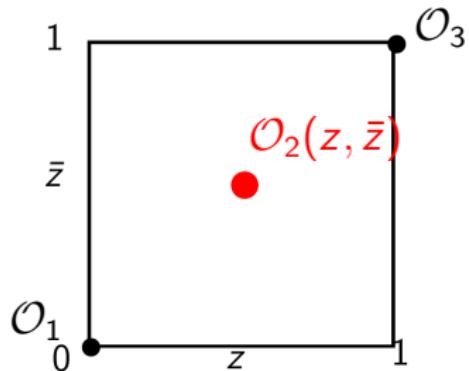


$$\frac{1}{x_{12}^{2\Delta_{\mathcal{O}}} x_{34}^{2\Delta_{\mathcal{O}}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}_3 \mathcal{O}_4 \mathcal{O}_{\Delta,\ell}} g_{\Delta,\ell}(z, \bar{z}) = \\ \frac{1}{x_{14}^{2\Delta_{\mathcal{O}}} x_{23}^{2\Delta_{\mathcal{O}}}} \sum_{\tilde{\mathcal{O}}_{\Delta,\ell}} \lambda_{\mathcal{O}_1 \mathcal{O}_4 \tilde{\mathcal{O}}_{\Delta,\ell}} \lambda_{\mathcal{O}_2 \mathcal{O}_3 \tilde{\mathcal{O}}_{\Delta,\ell}} g_{\Delta,\ell}(1-z, 1-\bar{z})$$

where $\Delta_{\mathcal{O}_i} = \Delta_{\mathcal{O}}$, $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2} = z\bar{z}$, $v = \frac{x_{23}^2 x_{14}^2}{x_{13}^2 x_{24}^2} = (1-z)(1-\bar{z})$

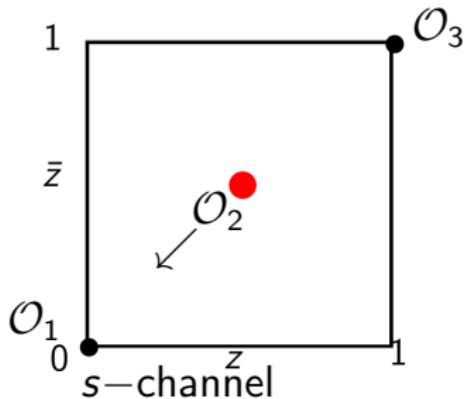
The crossing equation

$$\mathcal{O}_4(\infty)$$



The crossing equation

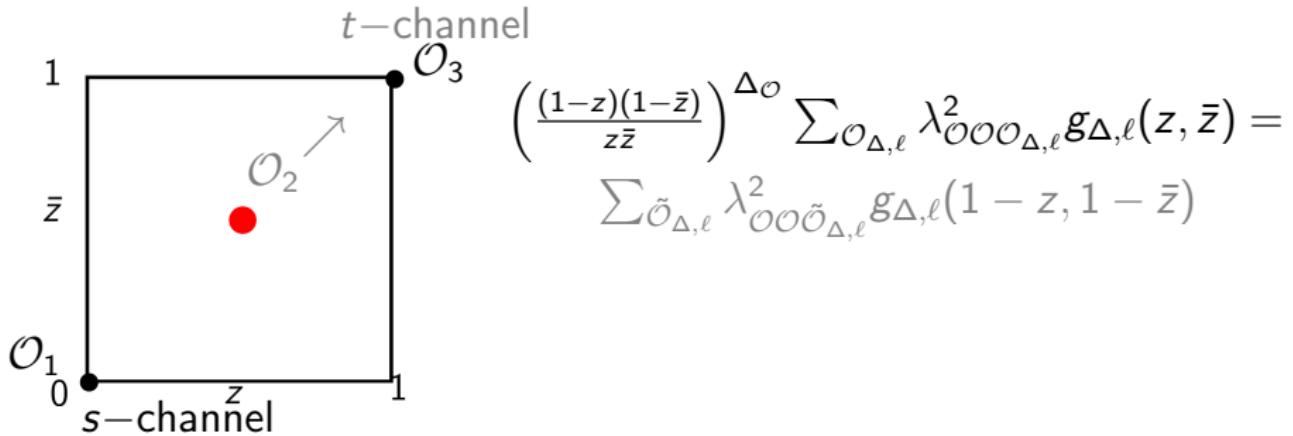
$$\mathcal{O}_4(\infty)$$



$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}} \right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) =$$

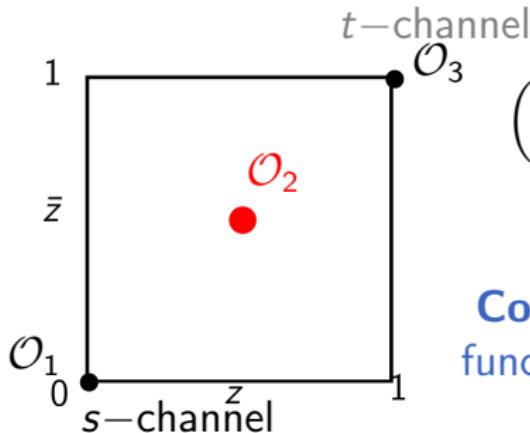
The crossing equation

$$\mathcal{O}_4(\infty)$$



The crossing equation

$$\mathcal{O}_4(\infty)$$

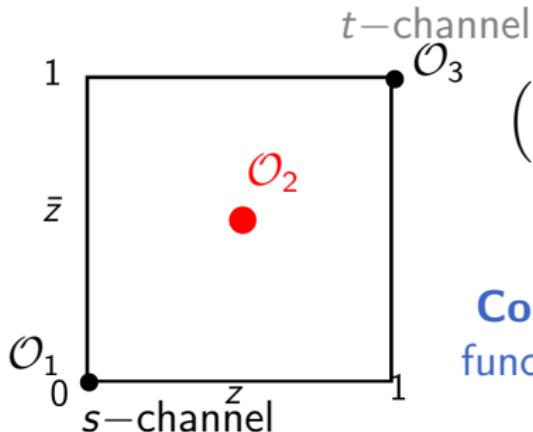


$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}} \right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) = \sum_{\tilde{\mathcal{O}}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\tilde{\mathcal{O}}_{\Delta,\ell}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$$

Consistent CFT \Rightarrow All four point functions are **crossing symmetric**

The crossing equation

$$\mathcal{O}_4(\infty)$$



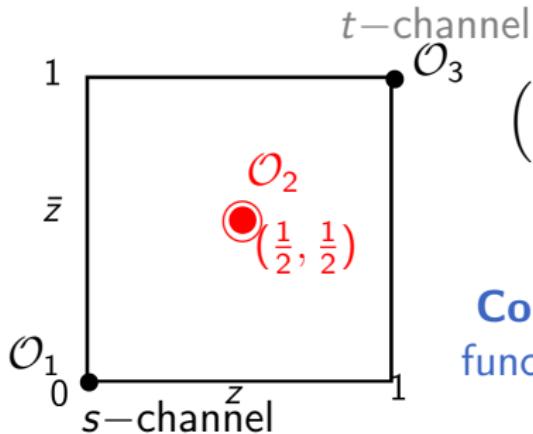
$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}} \right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) = \sum_{\tilde{\mathcal{O}}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\tilde{\mathcal{O}}_{\Delta,\ell}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$$

Consistent CFT \Rightarrow All four point functions are crossing symmetric

Our tools

The crossing equation

$$\mathcal{O}_4(\infty)$$



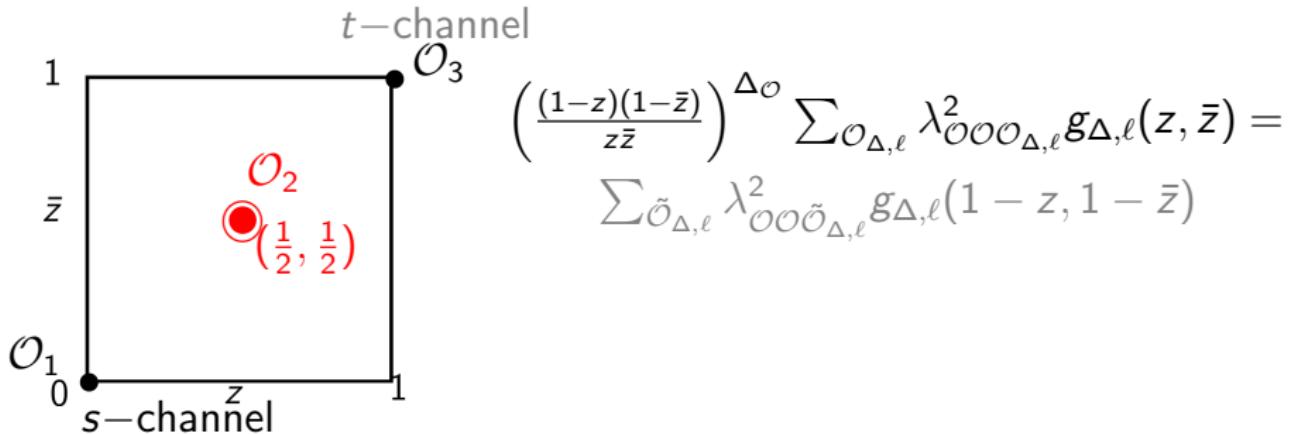
$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}} \right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) = \sum_{\tilde{\mathcal{O}}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\tilde{\mathcal{O}}_{\Delta,\ell}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$$

Consistent CFT \Rightarrow All four point functions are **crossing symmetric**

Our tools

The crossing equation

$$\mathcal{O}_4(\infty)$$

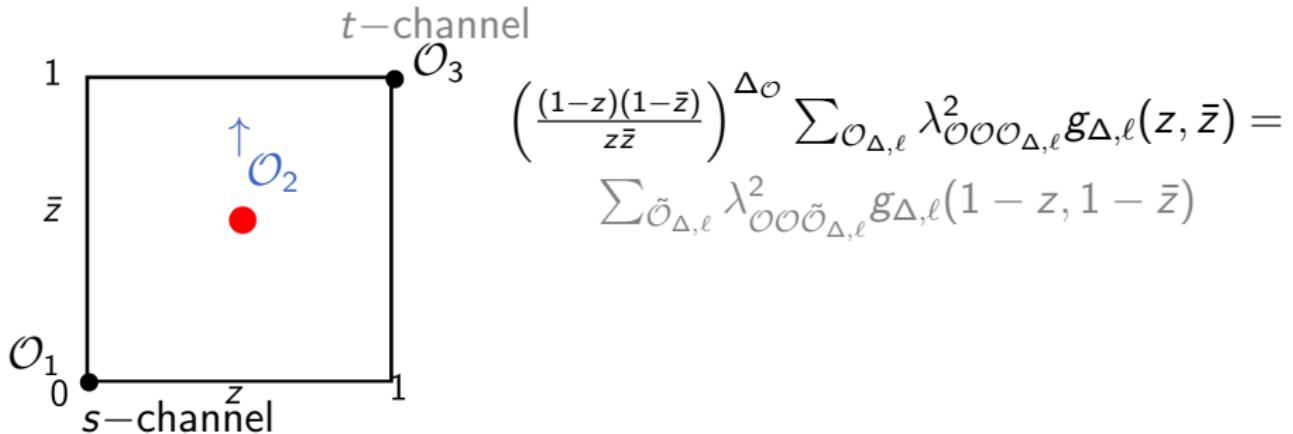


Our tools

- ▶ Numerical bootstrap [Rattazzi Rychkov Tonni Vichi]

The crossing equation

$$\mathcal{O}_4(\infty)$$



Our tools

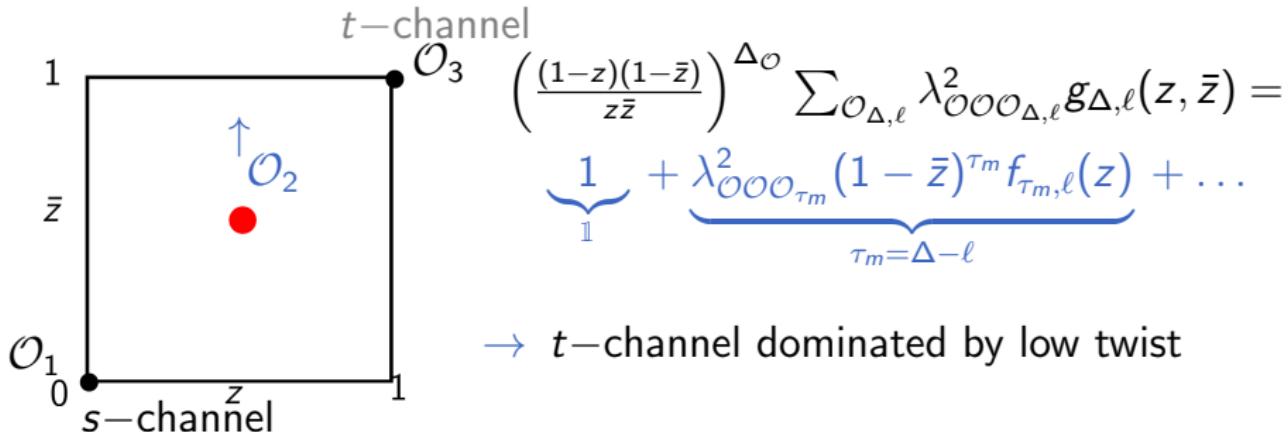
► Numerical bootstrap [Rattazzi Rychkov Tonni Vichi]

► Lightcone bootstrap [see Alday's talk]

[Alday Maldacena, Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]

The crossing equation

$$\mathcal{O}_4(\infty)$$



Our tools

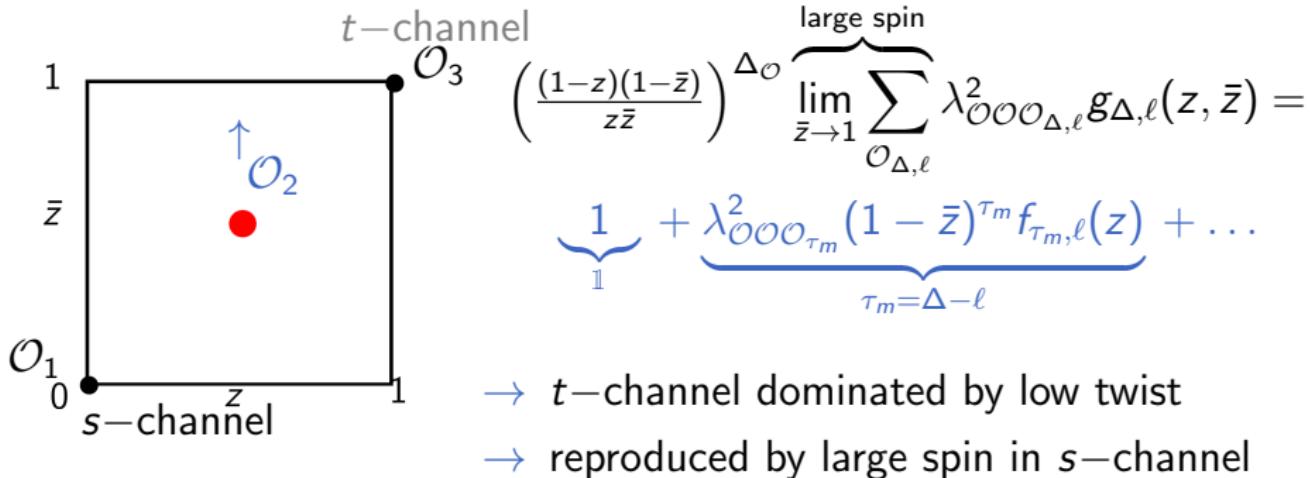
► Numerical bootstrap [Rattazzi Rychkov Tonni Vichi]

► Lightcone bootstrap [see Alday's talk]

[Alday Maldacena, Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]

The crossing equation

$$\mathcal{O}_4(\infty)$$



Our tools

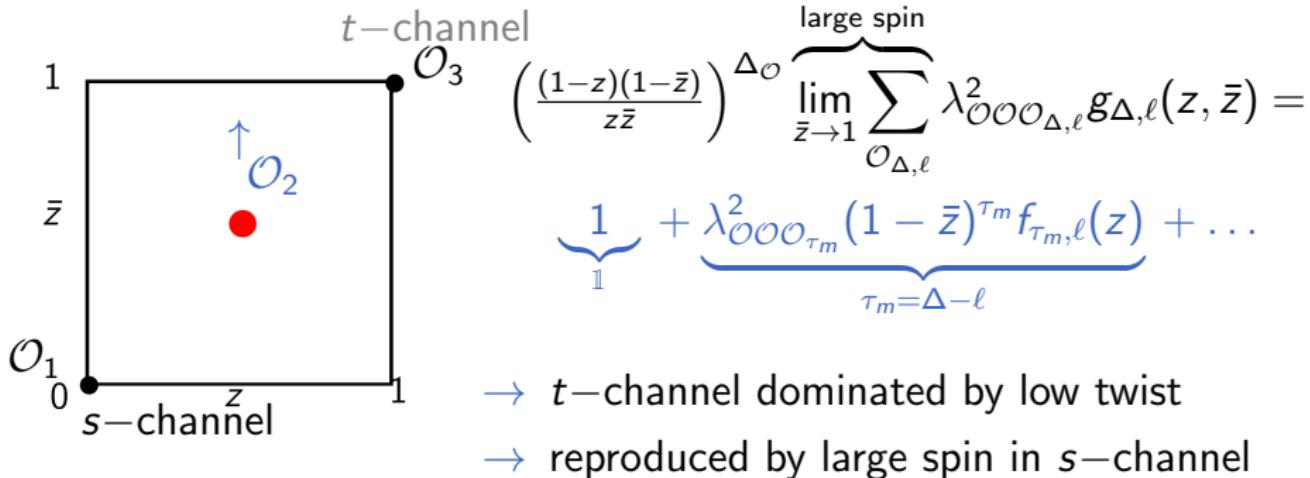
► Numerical bootstrap [Rattazzi Rychkov Tonni Vichi]

► Lightcone bootstrap [see Alday's talk]

[Alday Maldacena, Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]

The crossing equation

$$\mathcal{O}_4(\infty)$$



Our tools

► Numerical bootstrap [Rattazzi Rychkov Tonni Vichi]

► Lightcone bootstrap [see Alday's talk]

[Alday Maldacena, Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]

↪ Lorentzian inversion formula of [Caron-Huot, see his talk]

Numerical bootstrap review

Numerical bootstrap review

- ▶ Solving crossing equations \Rightarrow constraining space of solutions
 - ↪ How large can an OPE coefficient be?

Numerical bootstrap review

- ▶ Solving crossing equations \Rightarrow constraining space of solutions
 - ↪ How large can an OPE coefficient be?

Sum rule: identical scalars \mathcal{O}

$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}} \right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) = \\ \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$$

Numerical bootstrap review

- ▶ Solving crossing equations \Rightarrow constraining space of solutions
 - ↪ How large can an OPE coefficient be?

Sum rule: identical scalars \mathcal{O}

$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}} \right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) = \\ \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$$

→ Identity operator $\lambda_{\mathcal{O}\mathcal{O}1} = 1$

$$u = z\bar{z}, v = (1-z)(1-\bar{z})$$

Numerical bootstrap review

- ▶ Solving crossing equations \Rightarrow constraining space of solutions
 - ↪ How large can an OPE coefficient be?

Sum rule: identical scalars \mathcal{O}

$$\left(\frac{(1-z)(1-\bar{z})}{z\bar{z}}\right)^{\Delta_{\mathcal{O}}} \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(z, \bar{z}) = \\ \sum_{\mathcal{O}_{\Delta,\ell}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$$

→ Identity operator $\lambda_{\mathcal{O}\mathcal{O}\mathbb{1}} = 1$

$$u = z\bar{z}, v = (1-z)(1-\bar{z})$$

$$\sum_{\substack{\mathcal{O}_{\Delta,\ell} \in \mathcal{O}\mathcal{O} \\ \mathcal{O}_{\Delta,\ell} \neq \mathbb{1}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta,\ell}}^2 \underbrace{\frac{u^{\Delta_{\mathcal{O}}} g_{\Delta,\ell}(v, u) - v^{\Delta_{\mathcal{O}}} g_{\Delta,\ell}(u, v)}{v^{\Delta_{\mathcal{O}}} - u^{\Delta_{\mathcal{O}}}}}_{F_{\Delta,\ell}(u,v)} = 1$$

Numerical bootstrap review

Sum rule

$$\sum_{\mathcal{O}_{\Delta,\ell} \in \mathcal{OO}, \mathcal{O}_{\Delta\ell} \neq \mathbb{1}} \lambda_{\mathcal{OO}\mathcal{O}_{\Delta,\ell}}^2 \quad F_{\Delta,\ell} = \quad 1$$

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 F_{\Delta, \ell} = 1$$

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that
 - ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that

- ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$
- ↪ $\psi \cdot F_{\Delta, \ell}(z, \bar{z}) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that
 - ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$
 - ↪ $\psi \cdot F_{\Delta, \ell}(z, \bar{z}) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum
 - ↪ Minimize $\psi \cdot 1$

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that
 - ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$
 - ↪ $\psi \cdot F_{\Delta, \ell}(z, \bar{z}) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum
 - ↪ Minimize $\psi \cdot 1$
- ▶ $\boxed{\lambda_{\phi\phi\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \leq \psi \cdot 1}$

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that
 - ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$
 - ↪ $\psi \cdot F_{\Delta, \ell}(z, \bar{z}) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum
 - ↪ Minimize $\psi \cdot 1$
- ▶ $\boxed{\lambda_{\phi\phi\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \leq \psi \cdot 1}$
- ▶ Truncate $\psi = \sum_{m,n}^{m,n \leq \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$

Numerical bootstrap review

Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that
 - ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$
 - ↪ $\psi \cdot F_{\Delta, \ell}(z, \bar{z}) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum
 - ↪ Minimize $\psi \cdot 1$
- ▶
$$\boxed{\lambda_{\phi\phi\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \leq \psi \cdot 1}$$
- ▶ Truncate $\psi = \sum_{m,n}^{m,n \leq \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$
 - ↪ Increase $\Lambda \Rightarrow$ bounds get stronger

Numerical bootstrap review

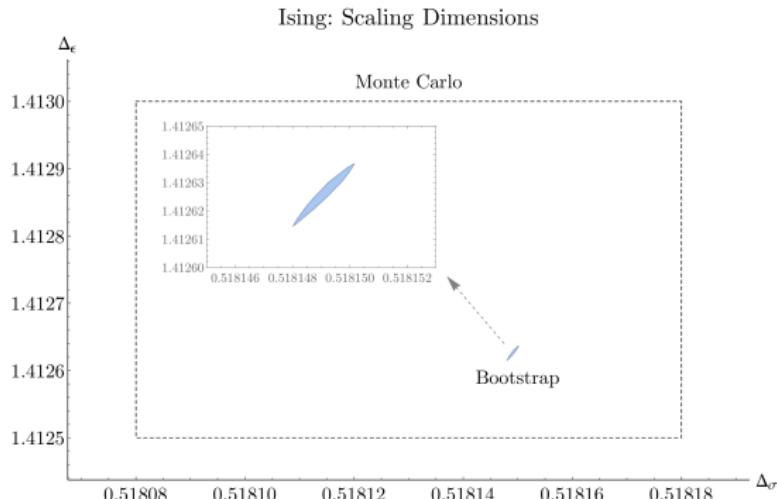
Sum rule

$$\lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \psi \cdot F_{\Delta_\star, \ell_\star} + \sum_{\substack{\mathcal{O}_{\Delta, \ell} \in \mathcal{OO}, \mathcal{O}_{\Delta \ell} \neq \mathbb{1} \\ \mathcal{O}_{\Delta, \ell} \neq \mathcal{O}_{\Delta_\star, \ell_\star}}} \lambda_{\mathcal{O}\mathcal{O}\mathcal{O}_{\Delta, \ell}}^2 \psi \cdot F_{\Delta, \ell} = \psi \cdot 1$$

- ▶ Find Functional ψ such that
 - ↪ $\psi \cdot F_{\Delta_\star, \ell_\star}(z, \bar{z}) = 1$
 - ↪ $\psi \cdot F_{\Delta, \ell}(z, \bar{z}) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum
 - ↪ Minimize $\psi \cdot 1$
- ▶
$$\boxed{\lambda_{\phi\phi\mathcal{O}_{\Delta_\star, \ell_\star}}^2 \leq \psi \cdot 1}$$
- ▶ Truncate $\psi = \sum_{m,n}^{m,n \leq \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$
 - ↪ Increase $\Lambda \Rightarrow$ bounds get stronger
 - ↪ Always true bounds

3d Ising Model

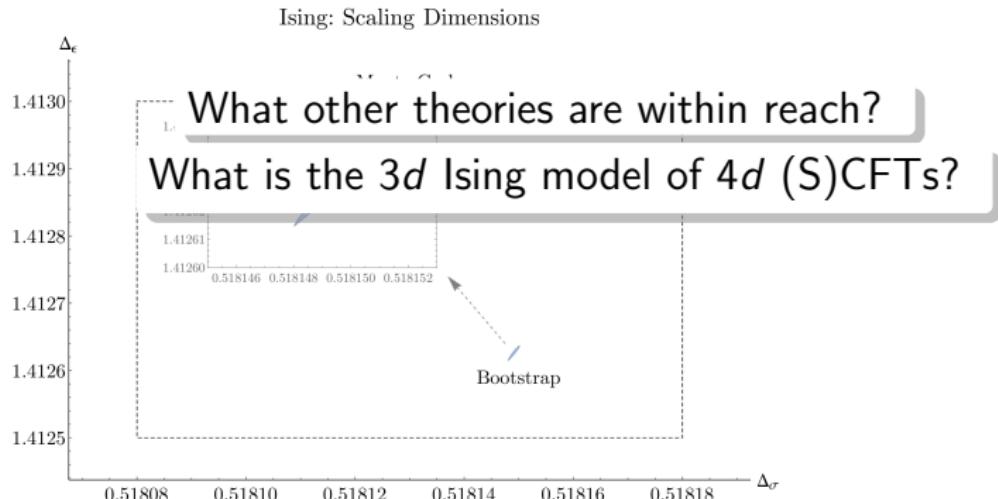
[Poland Simmons-Duffin Kos, Simmons-Duffin, Poland Simmons-Duffin Kos Vichi]



One \mathbb{Z}_2 -even, one \mathbb{Z}_2 -odd relevant scalar operator

3d Ising Model

[Poland Simmons-Duffin Kos, Simmons-Duffin, Poland Simmons-Duffin Kos Vichi]



One \mathbb{Z}_2 -even, one \mathbb{Z}_2 -odd relevant scalar operator

The Superconformal Bootstrap

Conformal field theory defined by

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

Conformal field theory defined by

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families

The Superconformal Bootstrap

Conformal field theory defined by

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families
- ▶ Finite re-organization of an infinite amount of data

The Superconformal Bootstrap

Conformal field theory defined by

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations? [see Rastelli's talk]

The Superconformal Bootstrap

Conformal field theory defined by

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

- ▶ Conformal families \leadsto Superconformal families
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations? [see Rastelli's talk]

- Yes, for $4d \mathcal{N} \geq 2$ [Beem ML Liendo Peelaers Rastelli van Rees]
(and also $6d \mathcal{N} = (2,0)$ and $2d \mathcal{N} = (0,4)$ [Beem Rastelli van Rees])

The Superconformal Bootstrap

Conformal field theory defined by

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{\lambda_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

- ▶ Conformal families \leadsto Superconformal families
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations? [see Rastelli's talk]

- Yes, for $4d \mathcal{N} \geq 2$ [Beem ML Liendo Peelaers Rastelli van Rees]
(and also $6d \mathcal{N} = (2,0)$ and $2d \mathcal{N} = (0,4)$ [Beem Rastelli van Rees])
- Subsector $\mathcal{N} \geq 2$ SCFTs captured by $2d$ chiral algebra

A solvable subsector

$4d \mathcal{N} = 2$ SCFTs $\rightarrow 2d$ chiral algebra [see Rastelli's talk]

A solvable subsector

4d $\mathcal{N} = 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ $SU(2)_R$ current \mapsto 2d stress tensor $T(z)$

A solvable subsector

4d $\mathcal{N} = 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ $\underbrace{SU(2)_R}_{\in \text{Super-stress tensor multiplet}}$ current \mapsto 2d stress tensor $T(z)$

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A trivial statement in $2d$

- (super-)stress tensor four-point function fixed in terms of c_{2d}

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A trivial statement in $2d$

- (super-)stress tensor four-point function fixed in terms of c_{2d} ($\langle TT \rangle \propto c$)

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A trivial statement in $2d$

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)

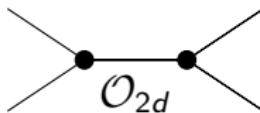
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2$$


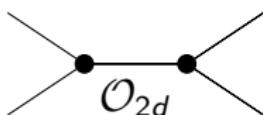
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2$$


- $\lambda_{\mathcal{O}_{2d}}^2$

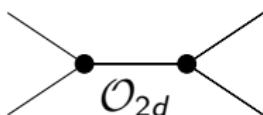
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2$$


$$\lambda_{\mathcal{O}_{2d}}^2 \leadsto \lambda_{\mathcal{O}_{4d}}^2$$

assumptions: interacting theory, unique stress tensor

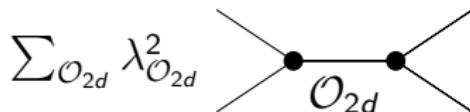
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:



$$\lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 \geq 0$$

4d unitarity

assumptions: interacting theory, unique stress tensor

A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra [see Rastelli's talk]

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:



$$\lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 \geq 0 \Rightarrow \text{New unitarity bounds}$$

4d unitarity

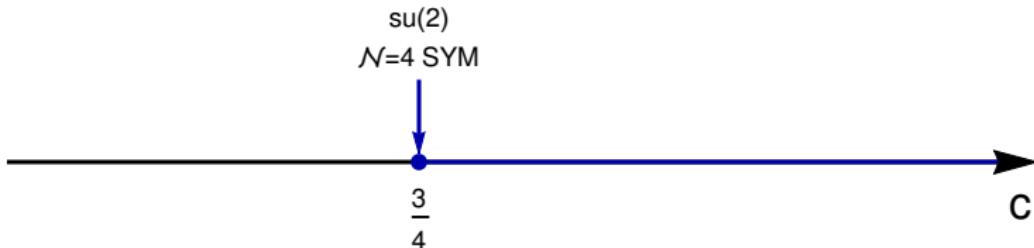
assumptions: interacting theory, unique stress tensor

Landscape of $4d \mathcal{N} \geq 2$ SCFTs

From $2d$ (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

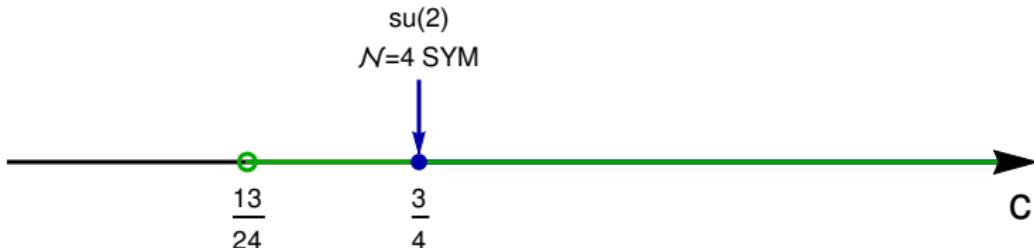
From 2d (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

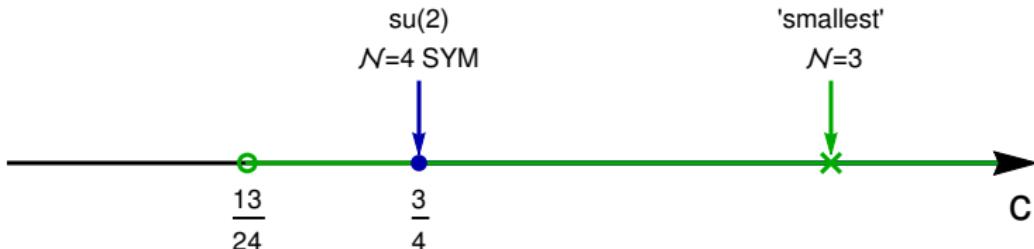
From 2d (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

From 2d (super-)stress tensor four-point function

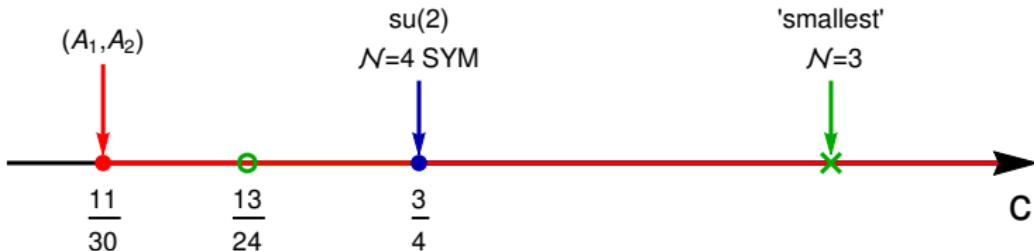
(assumptions: interacting theory, unique stress tensor)

→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator

→ 4d $\mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

From 2d (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

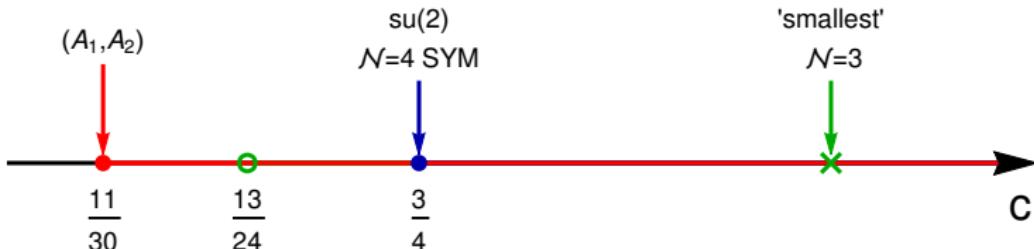
→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator

→ 4d $\mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]

↪ Saturated by the (A_1, A_2) Argyres-Douglas theory



Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description
- Strongly coupled isolated SCFT – no marginal deformations

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)]$ = Lee-Yang minimal model
[Beem Rastelli]

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)]$ = Lee-Yang minimal model
[Beem Rastelli]

Our tools beyond protected subsector

- ▶ Numerical bootstrap
 - [Rattazzi Rychkov Tonni Vichi]

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)]$ = Lee-Yang minimal model
[Beem Rastelli]

Our tools beyond protected subsector

- ▶ Numerical bootstrap
 - [Rattazzi Rychkov Tonni Vichi]
- ▶ Lightcone bootstrap

[Alday Maldacena, Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]

- Lorentzian inversion formula of [Caron-Huot]

The “simplest” Argyres-Douglas theory

How can we approach it?

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ

$$\Delta_\phi = \frac{6}{5}$$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ $(\mathcal{Q}_\alpha^I \phi = 0)$

$$\Delta_\phi = \frac{6}{5}$$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ $(\mathcal{Q}_\alpha^I \phi = 0)$

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ $(\mathcal{Q}_\alpha^I \phi = 0)$
$$\Delta_\phi = \frac{6}{5}$$
 $U(1)_r$ charge $r = \Delta_\phi$
- ▶ Study $\langle \phi(x_1) \phi(x_2) \bar{\phi}(x_3) \bar{\phi}(x_4) \rangle$

The “simplest” Argyres-Douglas theory

How can we approach it?

- Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

- Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)
$$\boxed{\Delta_\phi = \frac{6}{5}}$$
$$U(1)_r \text{ charge } r = \Delta_\phi$$

- ▶ Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

- ▶ Two OPE channels:

$$\hookrightarrow \phi\phi \sim \phi^2 + \dots$$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)
$$\boxed{\Delta_\phi = \frac{6}{5}}$$
$$U(1)_r \text{ charge } r = \Delta_\phi$$
- ▶ Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$
- ▶ Two OPE channels:
 - ↪ $\phi\phi \sim \phi^2 + \dots$
 - ↪ $\phi\bar{\phi} \sim \text{Identity} + \text{Super-stress tensor} + \dots$

The “simplest” Argyres-Douglas theory

How can we approach it?

- Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

- Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$
- Two OPE channels:
 - $\leftrightarrow \phi\phi \sim \phi^2 + \dots$
 - $\leftrightarrow \phi\bar{\phi} \sim \text{Identity} + \text{Super-stress tensor} + \dots$
- Conformal blocks \leadsto superconformal blocks

The “simplest” Argyres-Douglas theory

How can we approach it?

- Known: $4d \mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

- Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

- Two OPE channels:

$$\hookrightarrow \phi\phi \sim \phi^2 + \dots$$

$$\hookrightarrow \phi\bar{\phi} \sim \text{Identity} + \text{Super-stress tensor} + \dots$$

- Conformal blocks \rightsquigarrow superconformal blocks

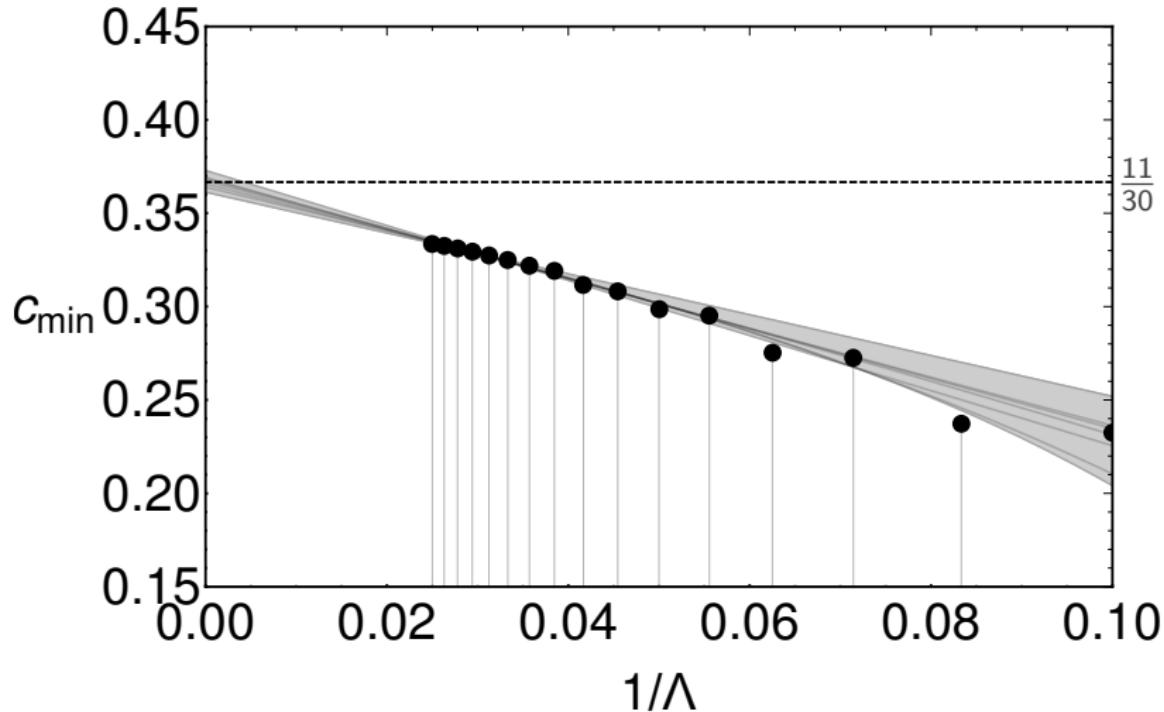
(only in $\phi\bar{\phi}$ channel) [Fitzpatrick Kaplan Khandker Li Poland Simmons-Duffin]

Minimum allowed central charge

Does $\langle \phi\bar{\phi}\bar{\phi}\bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?

Minimum allowed central charge

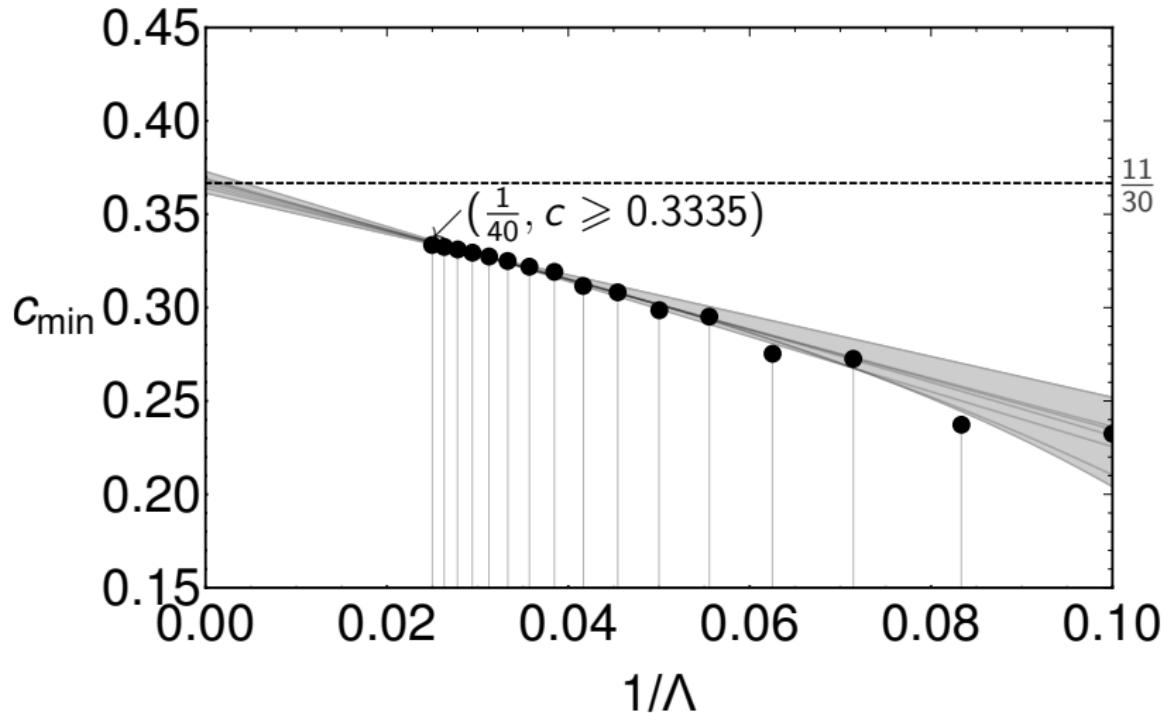
Does $\langle \phi \bar{\phi} \bar{\phi} \bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?



[Cornagliotto ML Liendo]

Minimum allowed central charge

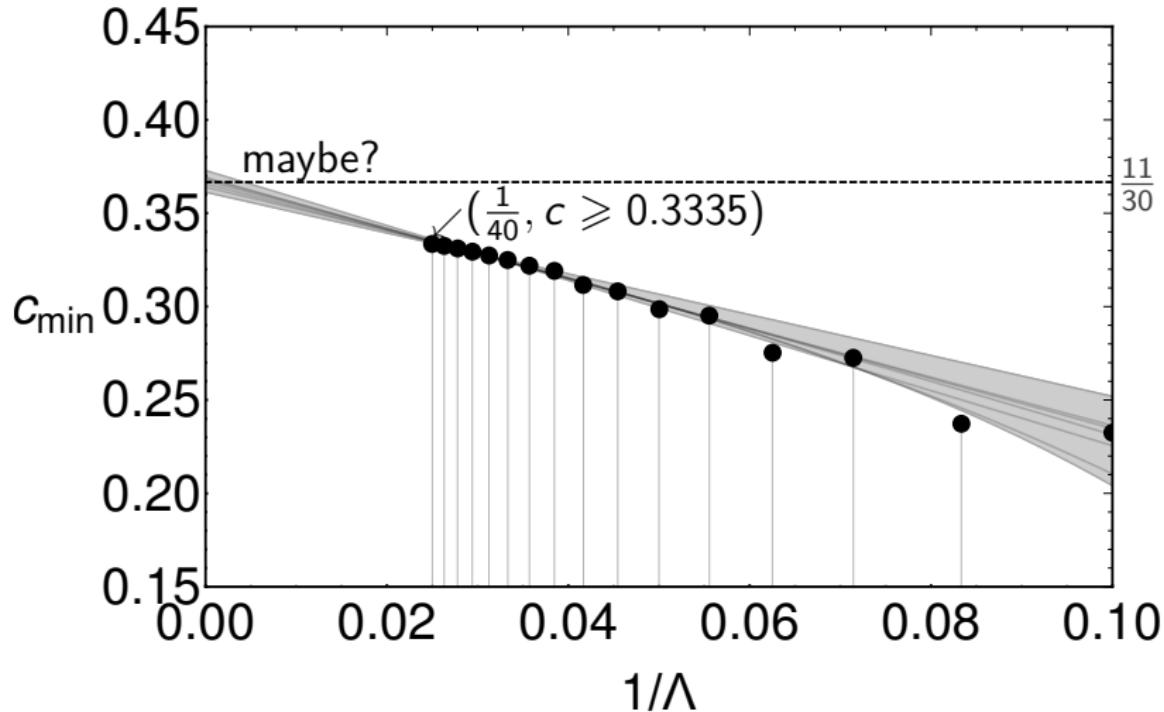
Does $\langle \phi \bar{\phi} \bar{\phi} \bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?



[Cornagliotto ML Liendo]

Minimum allowed central charge

Does $\langle \phi \bar{\phi} \bar{\phi} \bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?



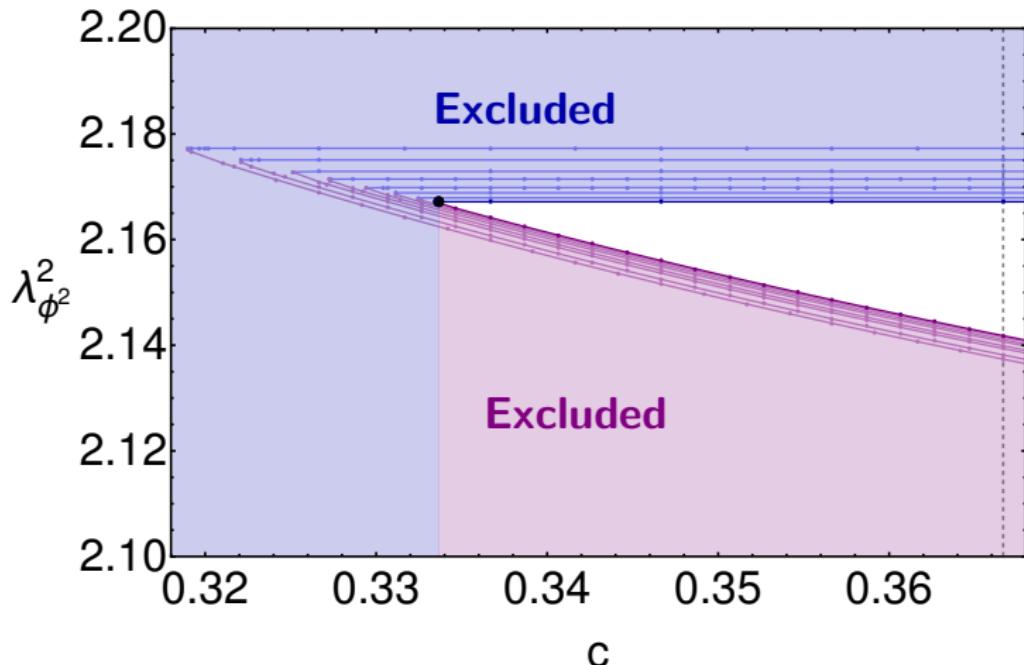
[Cornagliotto ML Liendo]

Bounding OPE coefficients

$$\phi\phi \sim \underbrace{\lambda_{\phi^2}^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$

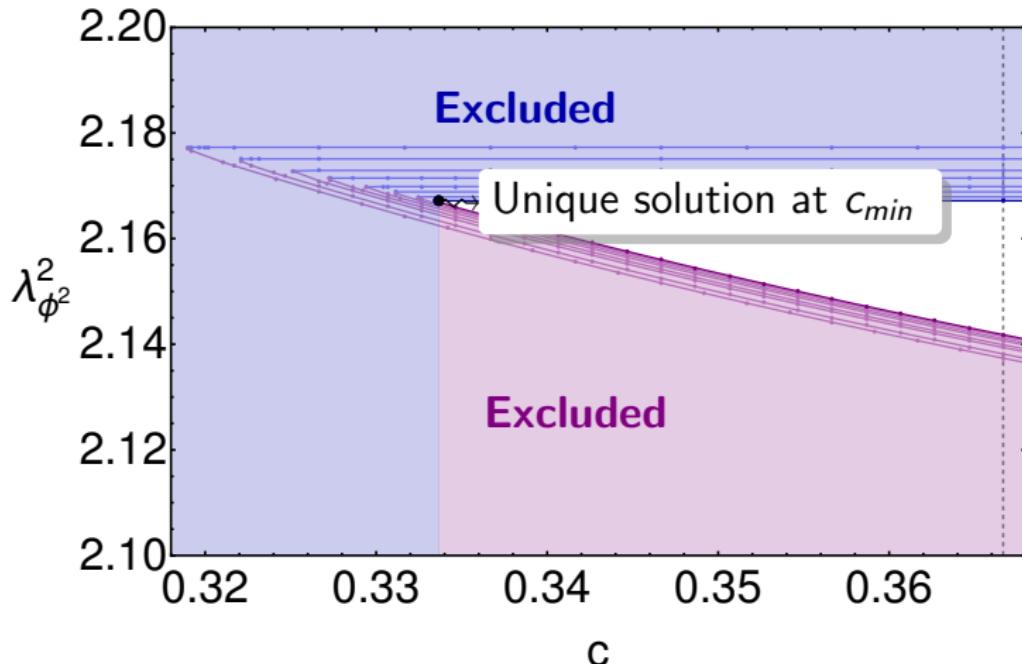
Bounding OPE coefficients

$$\phi\phi \sim \underbrace{\lambda_{\phi^2}^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$



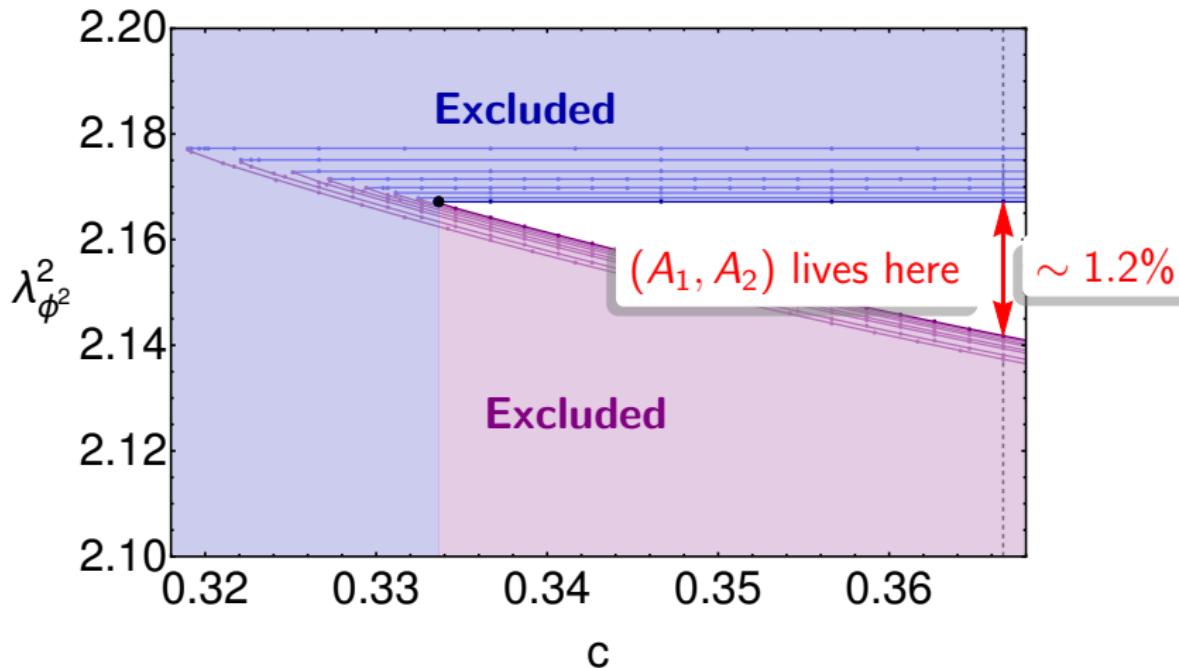
Bounding OPE coefficients

$$\phi\phi \sim \underbrace{\lambda_{\phi^2}^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$



Bounding OPE coefficients

$$\phi\phi \sim \underbrace{\lambda_{\phi^2}^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$

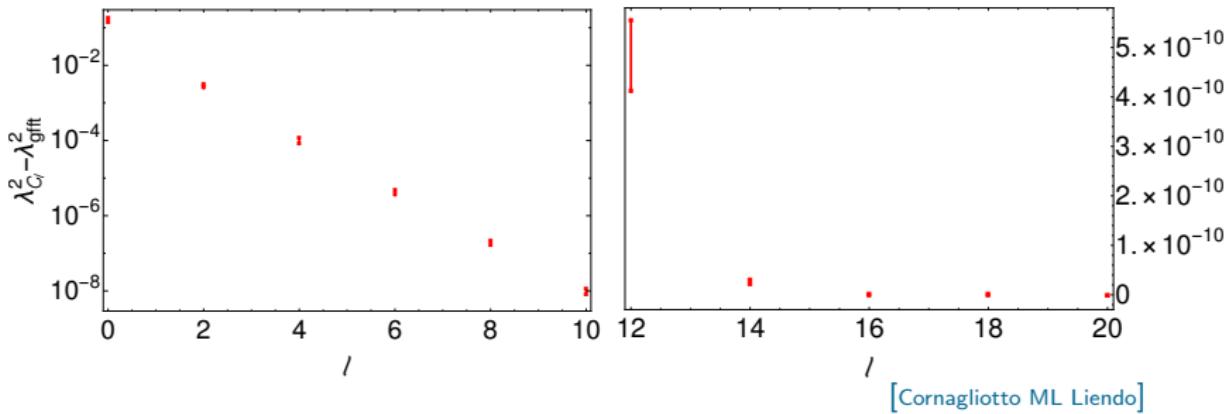


Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

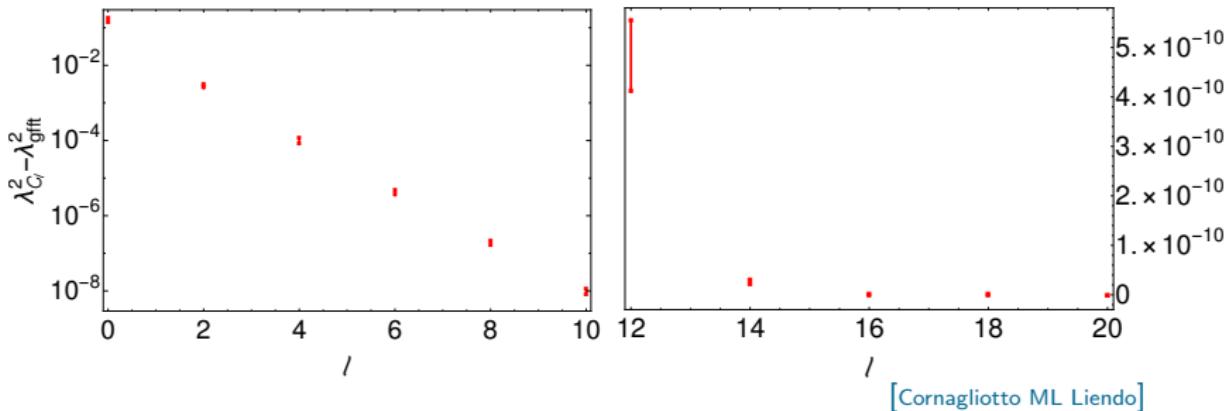
Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 + \underbrace{\lambda_{C_\ell}^2}_{\Delta=2\Delta_\phi} + \underbrace{\lambda_{C_{\ell>0}}^2}_{\Delta=2\Delta_\phi+\ell} + \dots$$

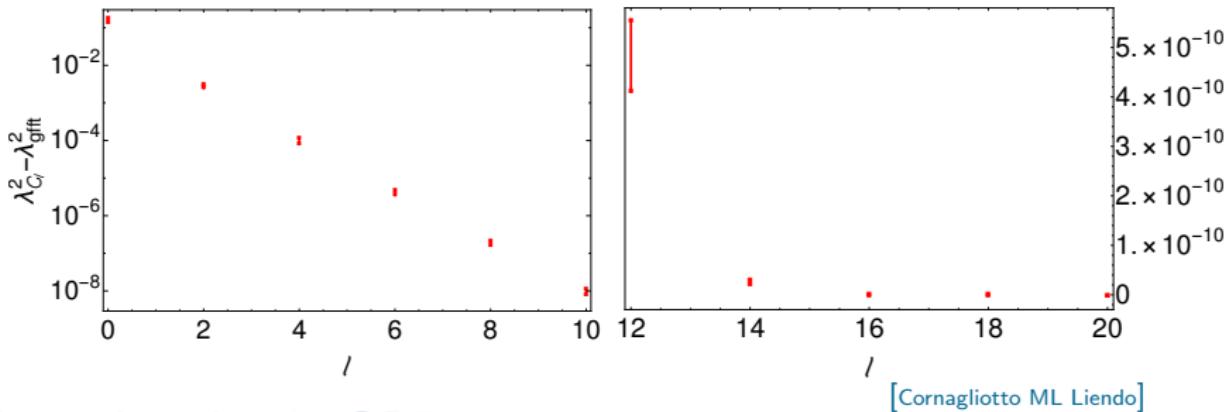


Inverting the $\phi\phi$ OPE

→ Same as bosonic inversion, valid for $\ell > 1$

Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

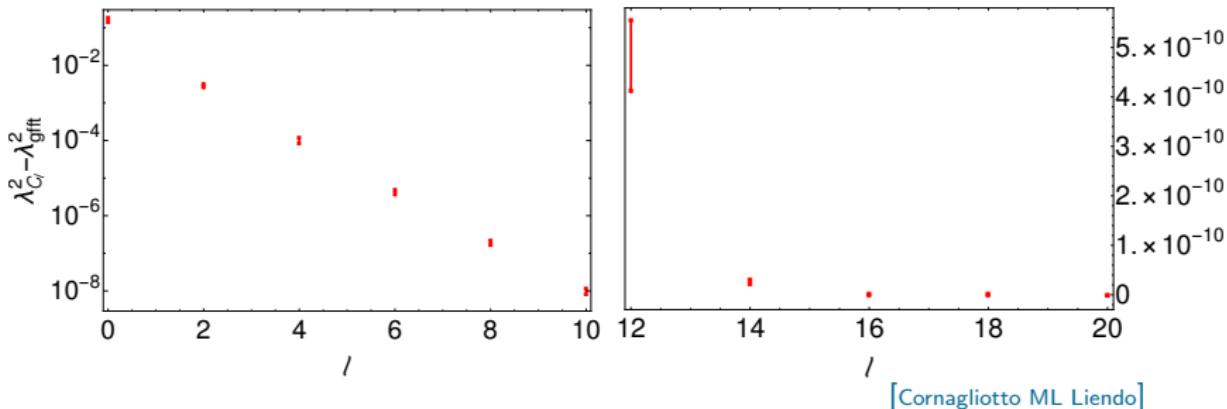


Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE

Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



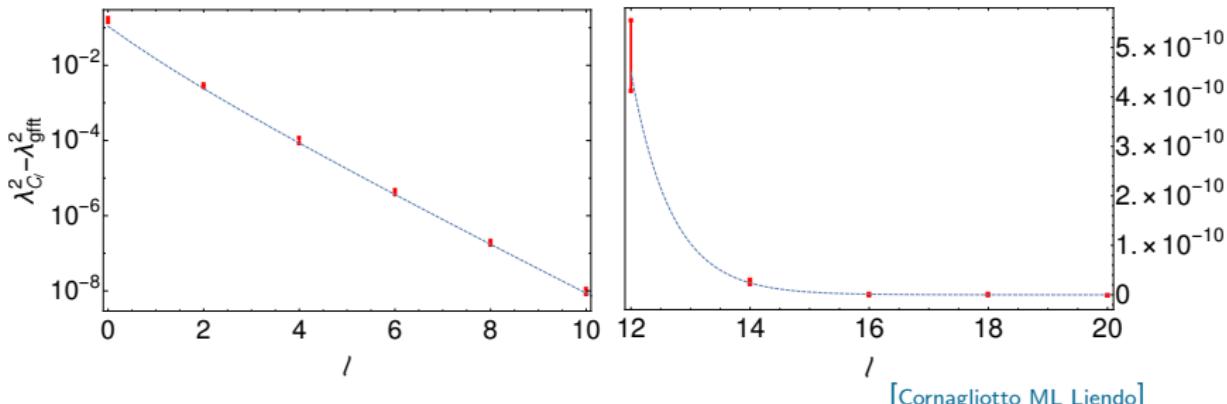
[Cornagliotto ML Liendo]

Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$

Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



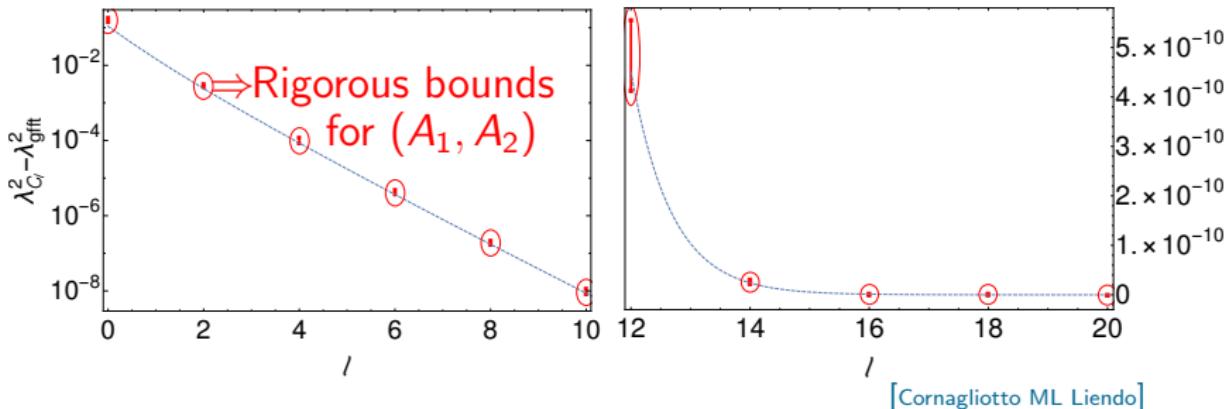
[Cornagliotto ML Liendo]

Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$
- Get s -channel ($\phi\phi$) large spin

Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

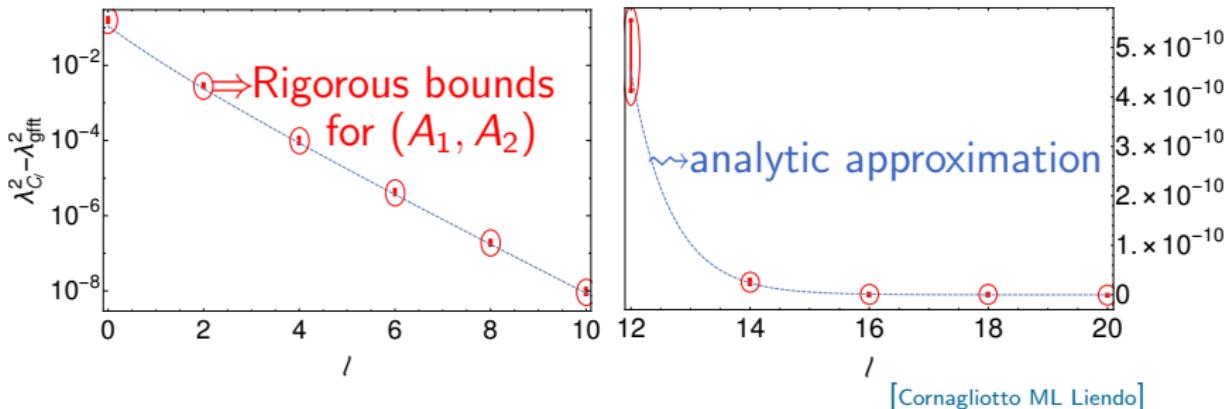


Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$
- Get s -channel ($\phi\phi$) large spin

Lorentzian inversion formula

$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$
- Get s -channel ($\phi\phi$) large spin

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

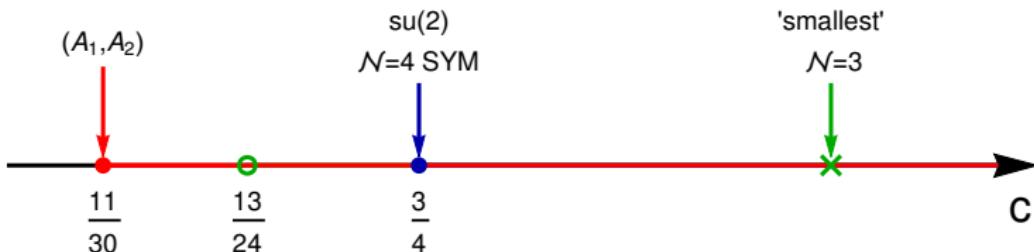
Landscape of $4d \mathcal{N} \geq 2$ SCFTs

Projection of space of SCFTs to an axis

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ $4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

→ $4d \mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



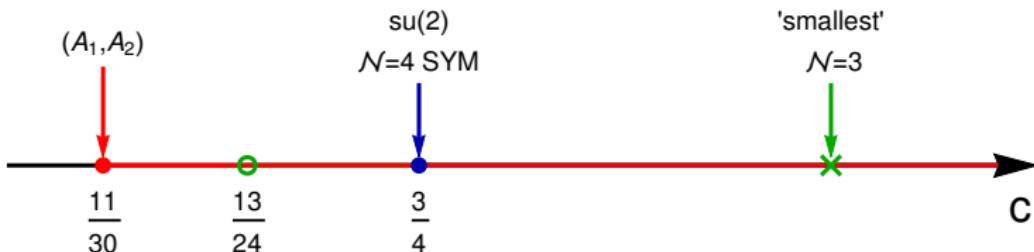
Landscape of $4d \mathcal{N} \geq 2$ SCFTs

Projection of space of SCFTs to an axis

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ $4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

→ $4d \mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



Finer view of the space of theories:

⇒ Organize theories by flavor symmetry

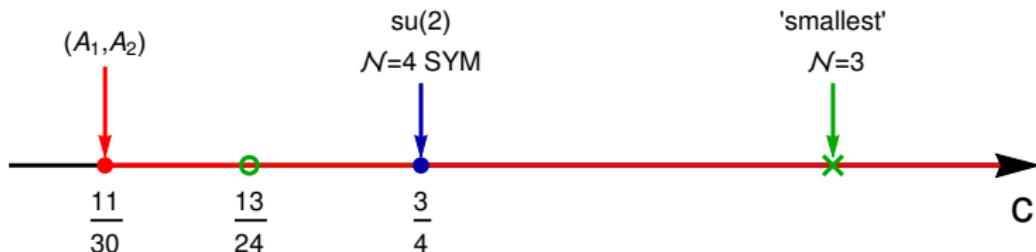
Landscape of $4d \mathcal{N} \geq 2$ SCFTs

Projection of space of SCFTs to an axis

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ $4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

→ $4d \mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



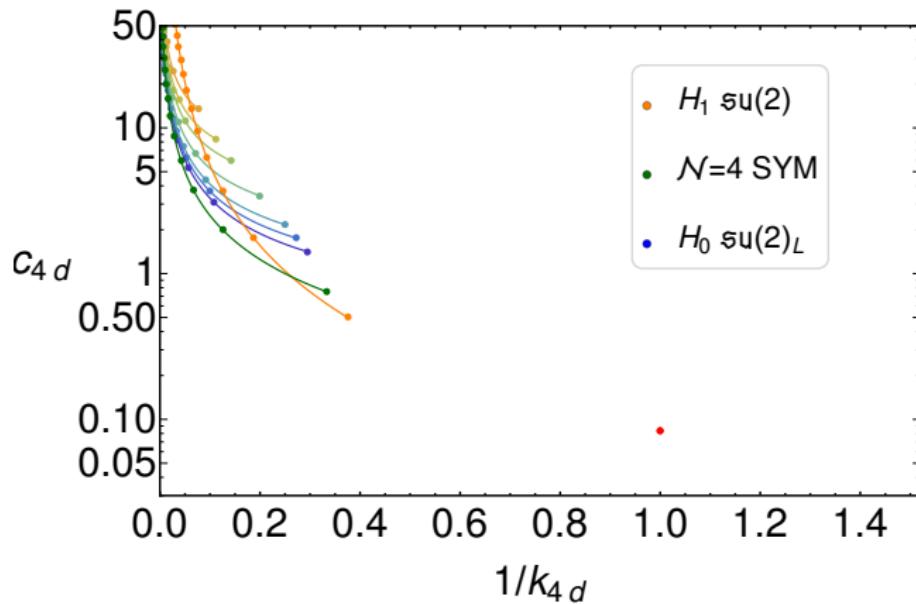
Finer view of the space of theories:

⇒ Organize theories by flavor symmetry

$$\langle TT \rangle \propto c, \quad \langle JJ \rangle \propto k$$

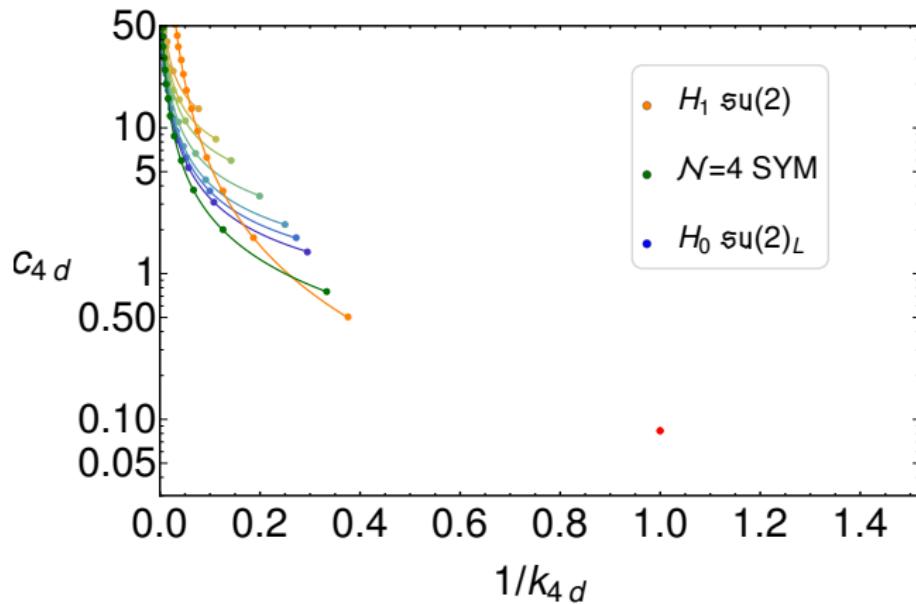
$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

- ▶ 4d Flavor current supermultiplet



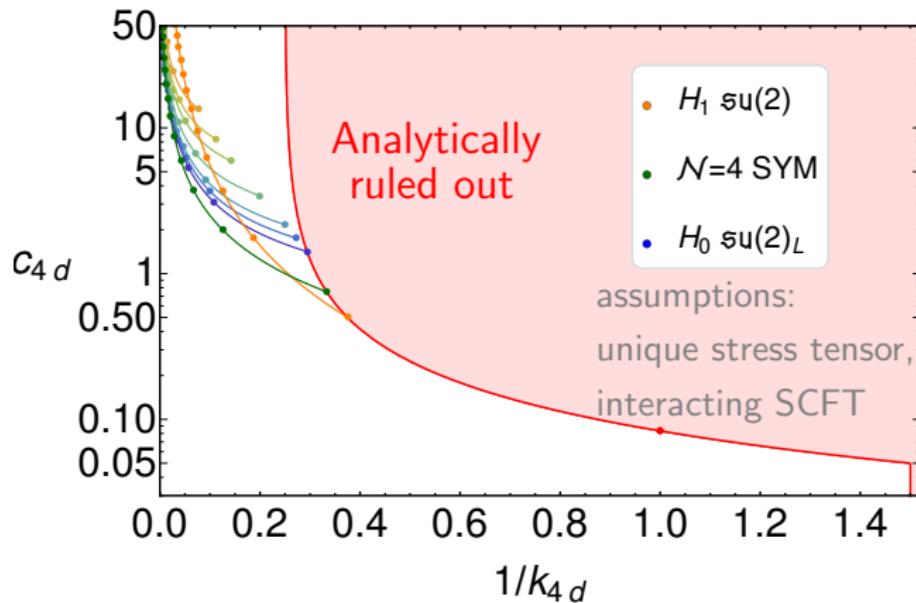
$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d}$



$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

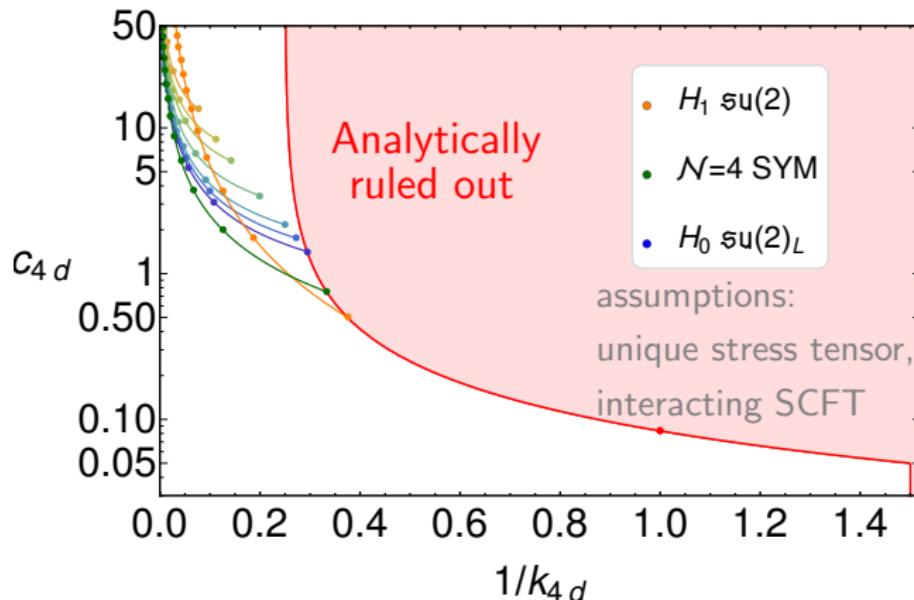
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \lambda_{4d}^2 \geqslant 0$



[Beem ML Liendo Peelaers Rastelli van Rees]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

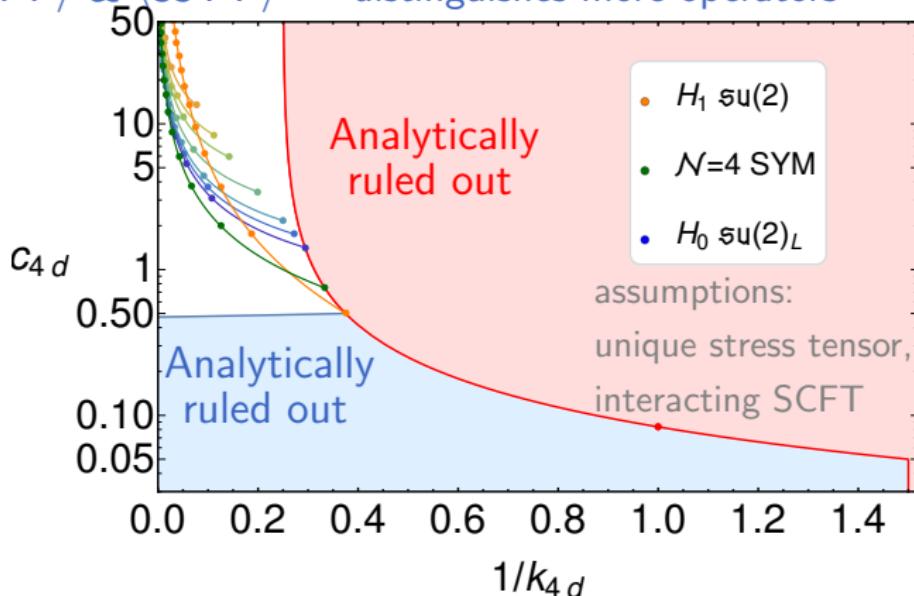
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum \lambda_{4d}^2 \geqslant 0$



[Beem ML Liendo Peelaers Rastelli van Rees]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

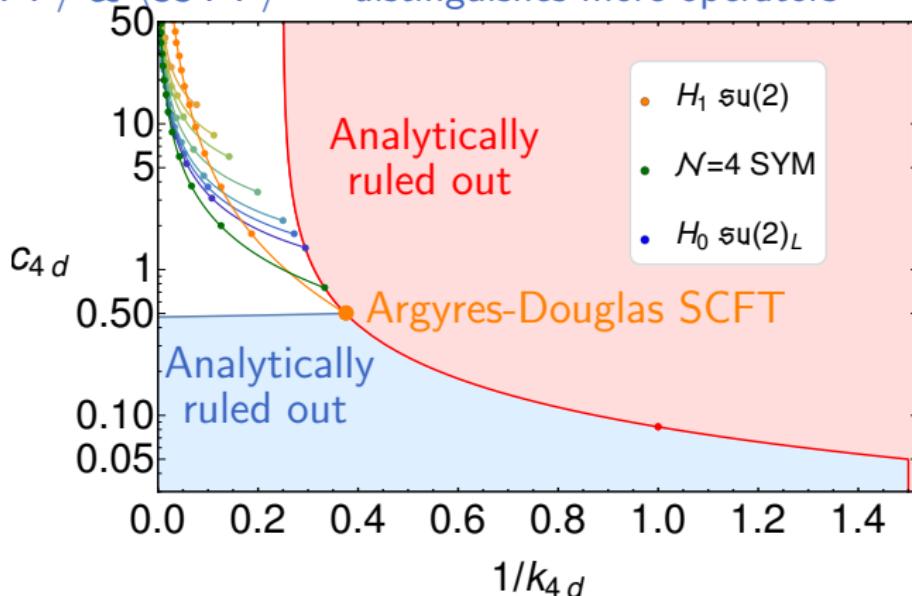
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum \lambda_{4d}^2 \geq 0$
- ▶ $\langle TTTT \rangle$ & $\langle JJTT \rangle \rightsquigarrow$ distinguishes more operators



[Beem ML Liendo Peelaers Rastelli van Rees, ML Liendo]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

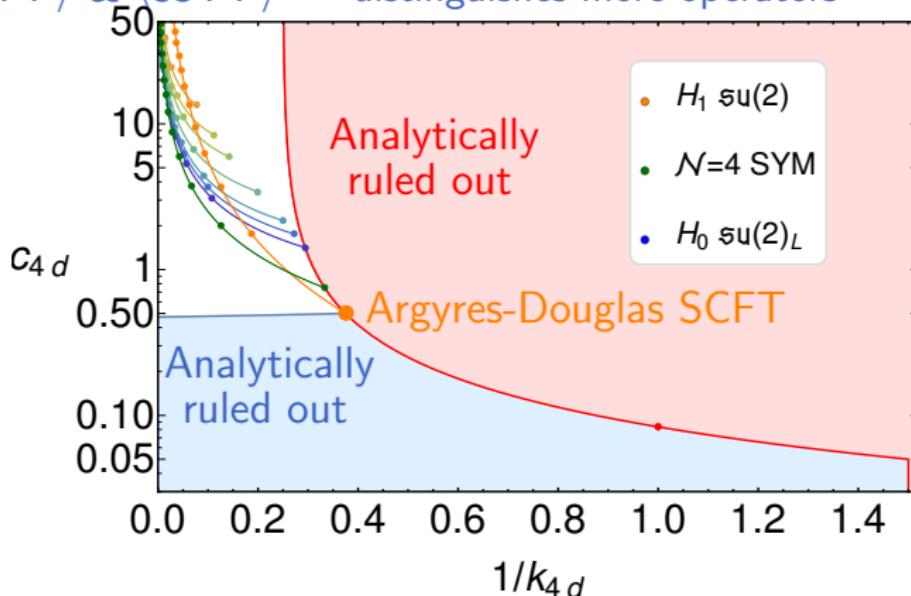
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum \lambda_{4d}^2 \geq 0$
- ▶ $\langle TTTT \rangle$ & $\langle JJTT \rangle \rightsquigarrow$ distinguishes more operators



[Beem ML Liendo Peelaers Rastelli van Rees, ML Liendo]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum \lambda_{4d}^2 \geq 0$
- ▶ $\langle TTTT \rangle$ & $\langle JJTT \rangle \rightsquigarrow$ distinguishes more operators



Only for $su(2)$, $su(3)$, $so(8)$, g_2 , f_4 , e_6 , e_7 , e_8

[Beem ML Liendo Peelaers Rastelli van Rees, ML Liendo]

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

→ Mixed system: stress tensor & flavor current multiplets

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?
- Two-dimensional long blocks [Cornagliotto ML Schomerus]
needed for $c > \frac{13}{24}$ for $\mathcal{N} = 3$ SCFTs

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?
- Two-dimensional long blocks [Cornagliotto ML Schomerus]
needed for $c > \frac{13}{24}$ for $\mathcal{N} = 3$ SCFTs
- Weight-shifting operators? [Karateev Kravchuk Simmons-Duffin]
- Calogero-Sutherland approach? [see Schomerus' talk]

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other isolated $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?
- Two-dimensional long blocks [Cornagliotto ML Schomerus]
needed for $c > \frac{13}{24}$ for $\mathcal{N} = 3$ SCFTs
- Weight-shifting operators? [Karateev Kravchuk Simmons-Duffin]
- Calogero-Sutherland approach? [see Schomerus' talk]

What is the “smallest” $\mathcal{N} = 3$ SCFT?

Thank you!

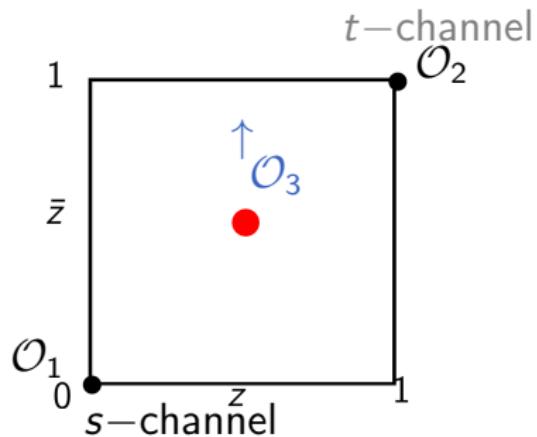
Backup slides

Outline

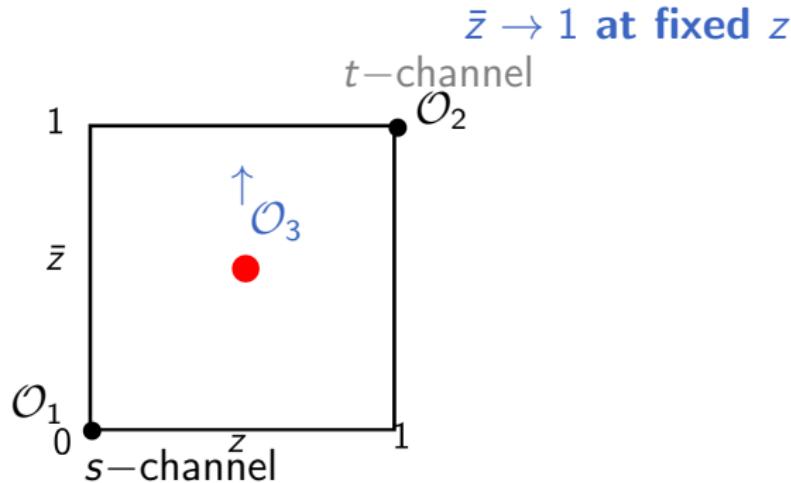
Inversion formula

- ⑤ Lorentzian inversion formula for (A_1, A_2)
- ⑥ A solvable subsector
- ⑦ Constraining the space of 4d $\mathcal{N} = 2$ SCFTs

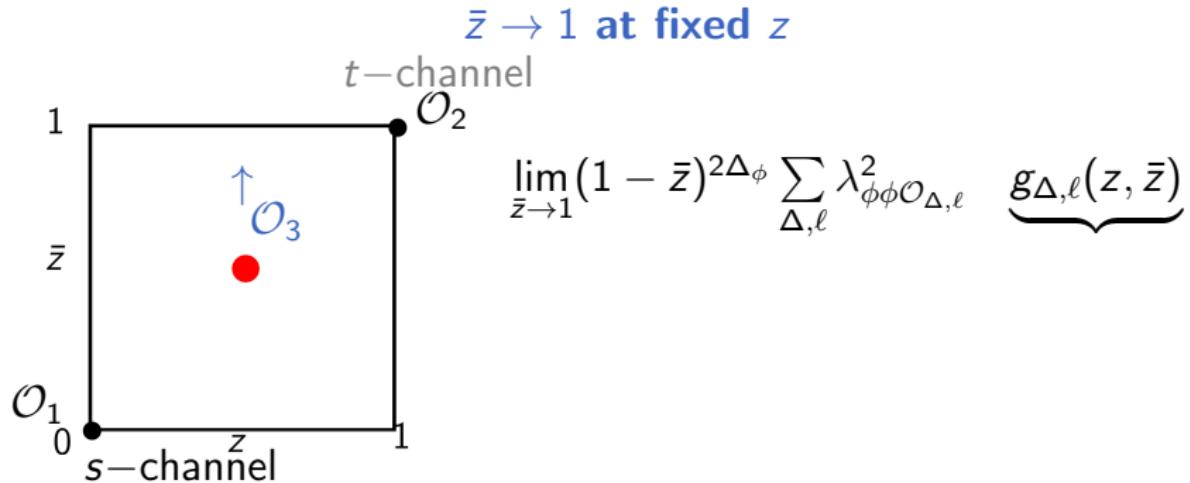
The lightcone bootstrap



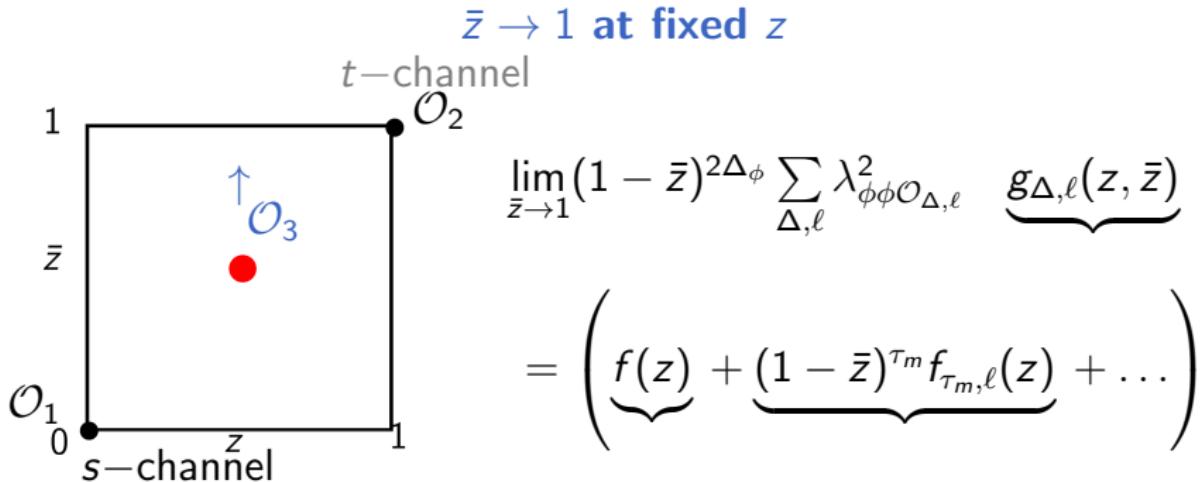
The lightcone bootstrap



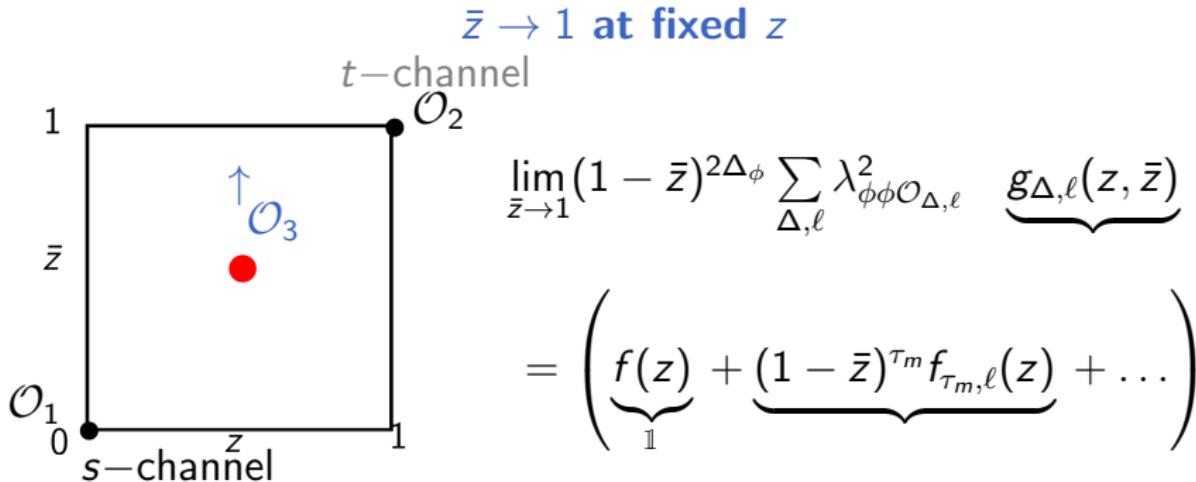
The lightcone bootstrap



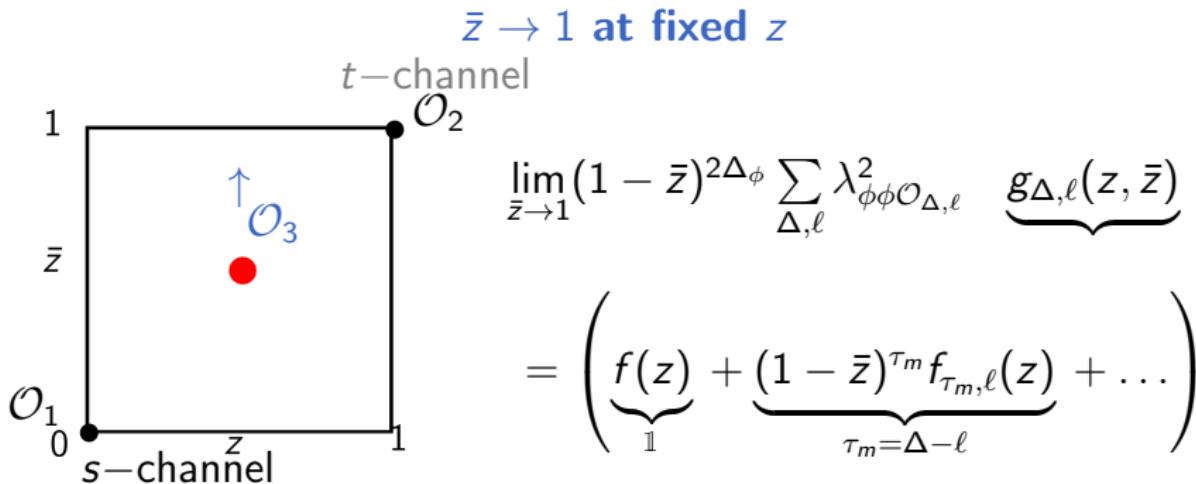
The lightcone bootstrap



The lightcone bootstrap

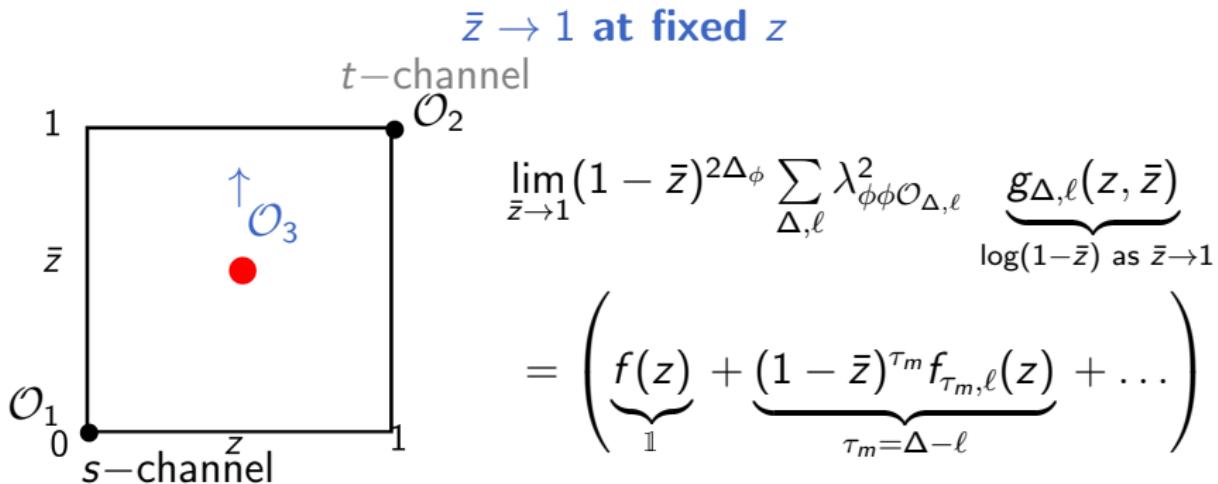


The lightcone bootstrap



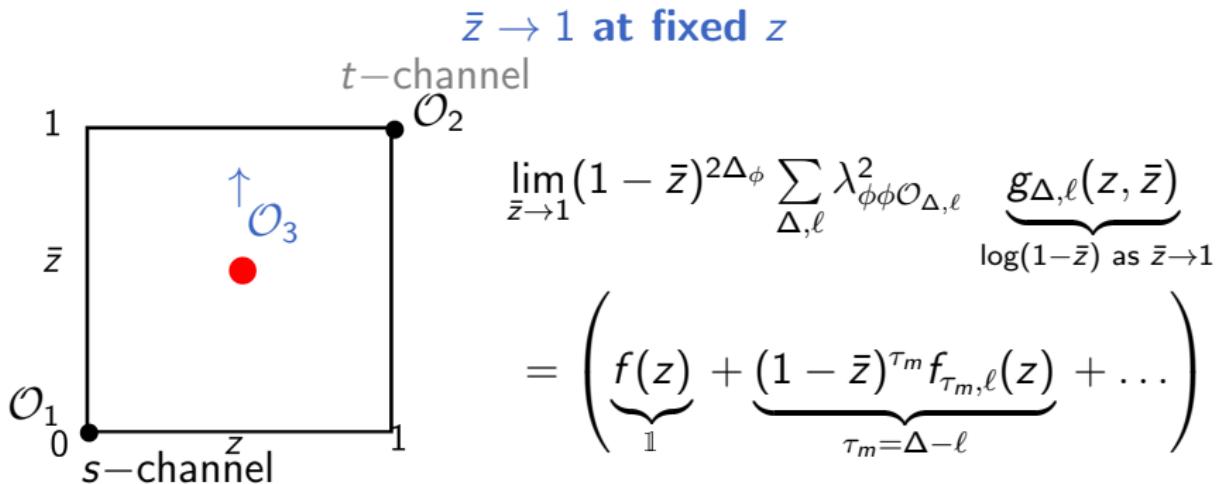
→ t -channel dominated by lowest twist τ_m operators

The lightcone bootstrap



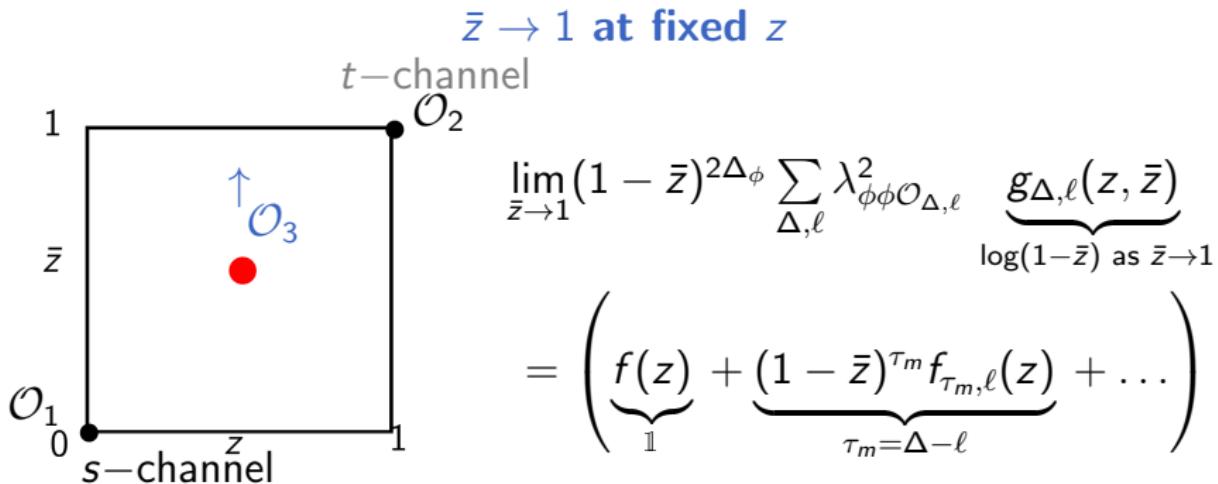
→ t -channel dominated by lowest twist τ_m operators

The lightcone bootstrap



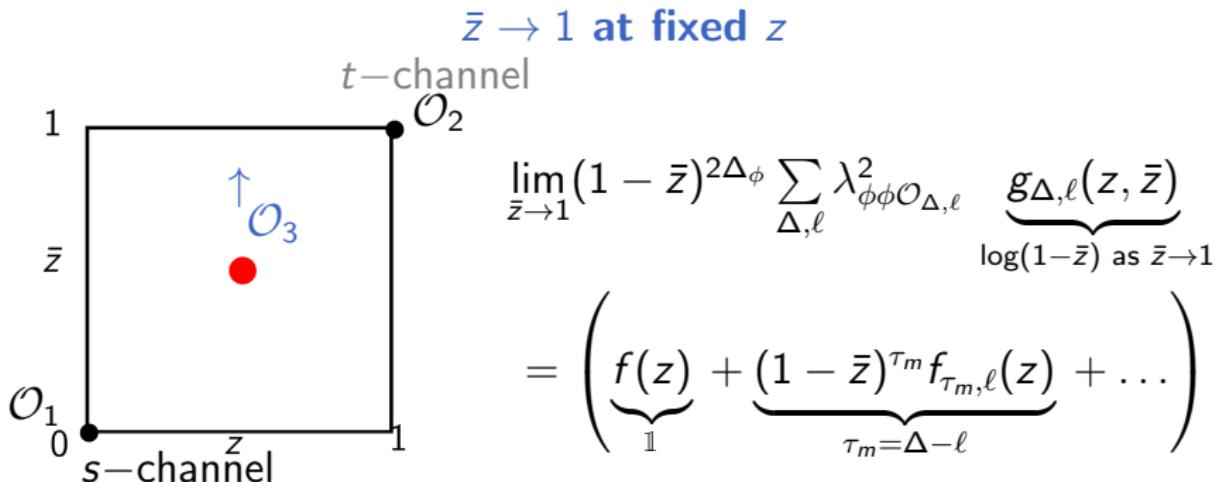
- t -channel dominated by lowest twist τ_m operators
- behavior reproduced by infinite sum over s -channel spins

The lightcone bootstrap



- t -channel dominated by lowest twist τ_m operators
- behavior reproduced by infinite sum over s -channel spins
- ↔ Large spin spectrum of CFT

The lightcone bootstrap



- t -channel dominated by lowest twist τ_m operators
- behavior reproduced by infinite sum over s -channel spins
- ↔ Large spin spectrum of CFT

$$\mathbb{1} \Rightarrow \Delta \rightarrow 2\Delta_\phi + 2n + \ell \quad (\phi \square^n \partial_{\mu_1} \dots \partial_{\mu_\ell} \phi)$$

A Lorentzian inversion formula

Large spin perturbation theory

→ Very successful for $3d$ Ising model

[Alday Zhiboedov, Simmons-Duffin]

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for $3d$ Ising model
[Alday Zhiboedov, Simmons-Duffin]
- down to spin two!

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for $3d$ Ising model
[Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for $3d$ Ising model
[Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula
- $c(\Delta, \ell)$ with poles where operators are, residues $\sim \lambda_{\Delta, \ell}^2$

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for $3d$ Ising model
 - [Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula
- $c(\Delta, \ell)$ with poles where operators are, residues $\sim \lambda_{\Delta, \ell}^2$
- Need to know full correlation function to get full spectrum

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for 3d Ising model
 - [Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula
- $c(\Delta, \ell)$ with poles where operators are, residues $\sim \lambda_{\Delta, \ell}^2$
- Need to know full correlation function to get full spectrum
- only makes sense for integer ℓ

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for 3d Ising model
 - [Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula
- $c(\Delta, \ell)$ with poles where operators are, residues $\sim \lambda_{\Delta, \ell}^2$
- Need to know full correlation function to get full spectrum
- only makes sense for integer ℓ
- [Caron-Huot] Inversion formula *analytic* in spin for $\ell > 1$

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for 3d Ising model
 - [Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula
- $c(\Delta, \ell)$ with poles where operators are, residues $\sim \lambda_{\Delta, \ell}^2$
- Need to know full correlation function to get full spectrum
- only makes sense for integer ℓ
- [Caron-Huot] Inversion formula *analytic* in spin for $\ell > 1$
- Operators organize in trajectories

A Lorentzian inversion formula

Large spin perturbation theory

- Very successful for 3d Ising model
 - [Alday Zhiboedov, Simmons-Duffin]
- down to spin two!
- Invert s –channel OPE: Euclidean inversion formula
- $c(\Delta, \ell)$ with poles where operators are, residues $\sim \lambda_{\Delta, \ell}^2$
- Need to know full correlation function to get full spectrum
- only makes sense for integer ℓ
- [Caron-Huot] Inversion formula *analytic* in spin
for $\ell > 1$
- Operators organize in trajectories
- large ℓ dominated by low t –channel twists

Outline

Inversion formula

- ⑤ Lorentzian inversion formula for (A_1, A_2)
- ⑥ A solvable subsector
- ⑦ Constraining the space of 4d $\mathcal{N} = 2$ SCFTs

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Feed in low twist in t -channel ($\bar{\phi}\phi$)

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Feed in low twist in t -channel ($\bar{\phi}\phi$)
 - $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

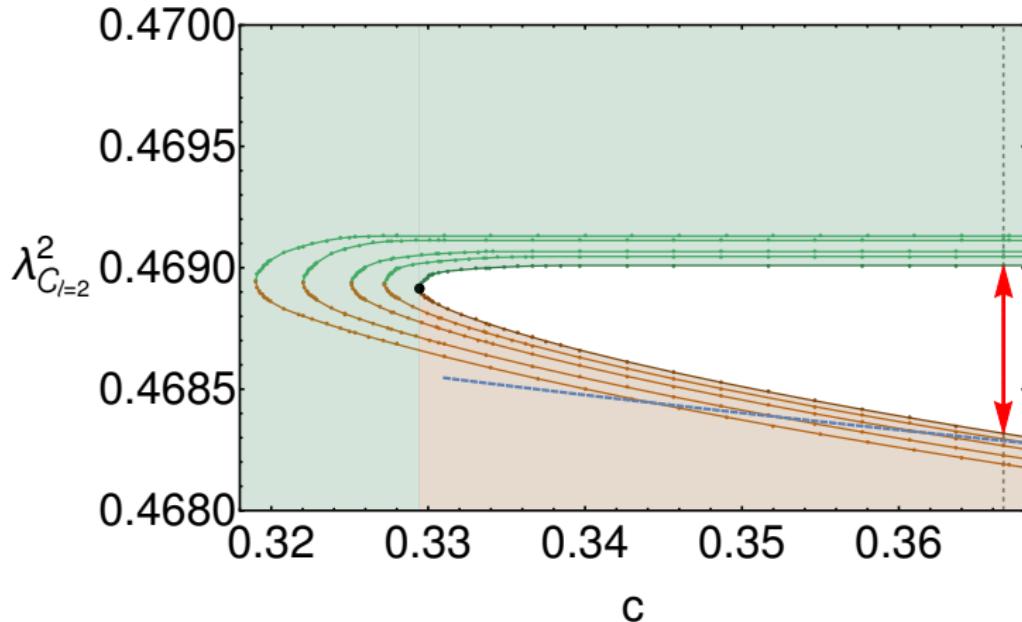
- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Feed in low twist in t -channel ($\bar{\phi}\phi$)
 - $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$
- and in u -channel ($\phi\phi$)
 - $\phi\phi \sim \phi^2 + \dots$

Bounding OPE coefficients

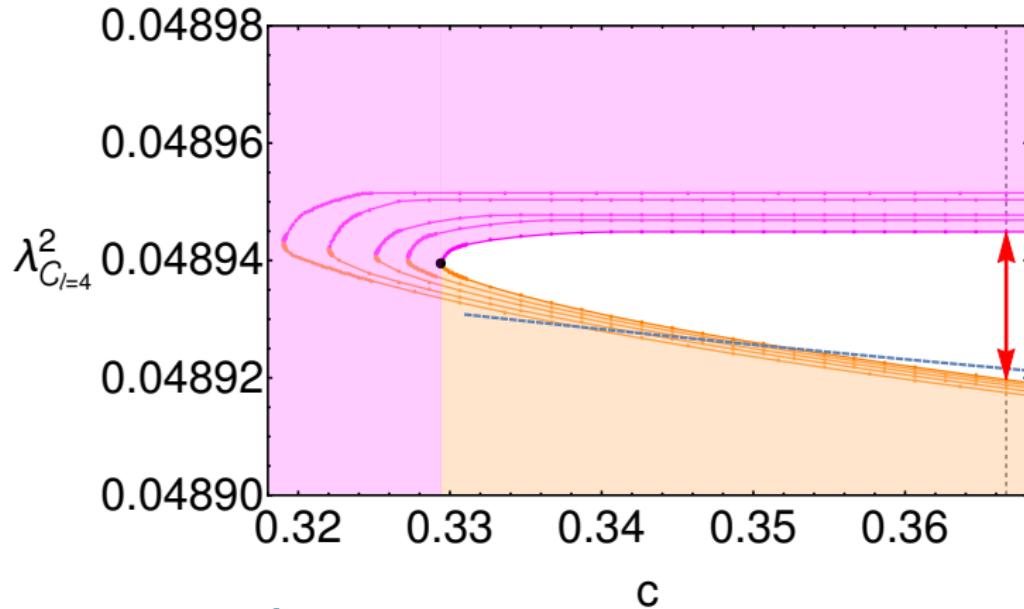
$$\phi\phi \sim \lambda_{\phi^2}^2 \underbrace{\frac{\phi^2}{\Delta=2\Delta_\phi}} + \lambda_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



[Cornagliotto ML Liendo]

Bounding OPE coefficients

$$\phi\phi \sim \lambda_{\phi^2}^2 + \underbrace{\lambda_{C_\ell}^2}_{\Delta=2\Delta_\phi} + \underbrace{\lambda_{C_{\ell>0}}^2}_{\Delta=2\Delta_\phi+\ell} + \dots$$



[Cornagliotto ML Liendo]

A Lorentzian inversion formula

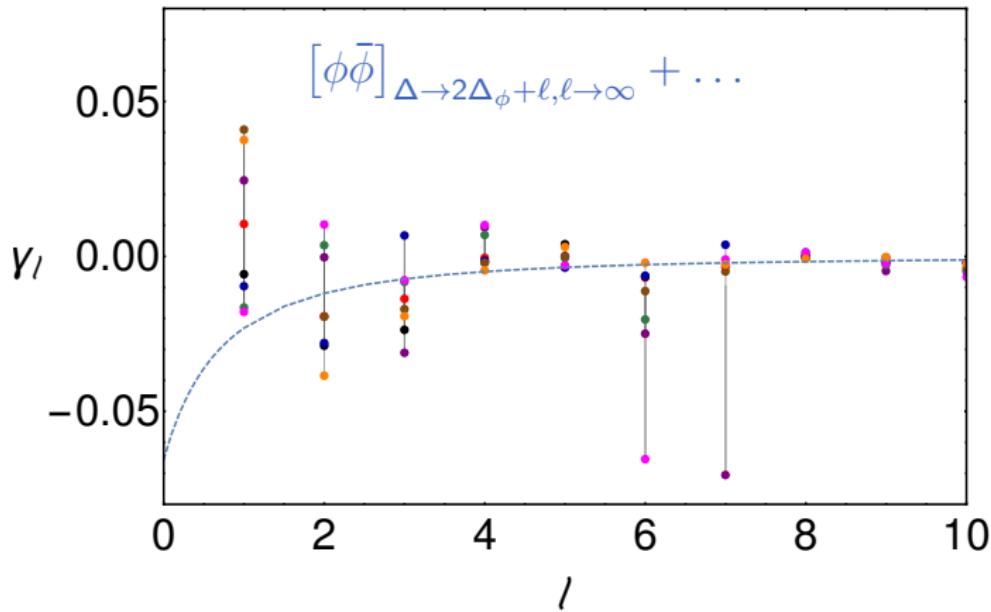
Inverting the $\phi\bar{\phi}$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$

A Lorentzian inversion formula

Inverting the $\phi\bar{\phi}$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$



Outline

Inversion formula

- ⑤ Lorentzian inversion formula for (A_1, A_2)
- ⑥ A solvable subsector
- ⑦ Constraining the space of 4d $\mathcal{N} = 2$ SCFTs

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1, j_2)}, \quad \}$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Claim

→ Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$,

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Claim

→ Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Claim

→ Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$

$$\langle \mathcal{O}_1^{l_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{l_n}(z_n, \bar{z}_n) \rangle$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Claim

- Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$
- Restrict to operators with $\Delta = 2R + j_1 + j_2$

$$\langle \mathcal{O}_1^{I_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{I_n}(z_n, \bar{z}_n) \rangle$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Claim

- Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$
- Restrict to operators with $\Delta = 2R + j_1 + j_2$

$$u_{I_1}(\bar{z}_1) \dots u_{I_n}(\bar{z}_n) \langle \mathcal{O}_1^{I_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{I_n}(z_n, \bar{z}_n) \rangle$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1 j_2), \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f}\}$$

Claim

- Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$
- Restrict to operators with $\Delta = 2R + j_1 + j_2$

$$u_{I_1}(\bar{z}_1) \dots u_{I_n}(\bar{z}_n) \langle \mathcal{O}_1^{I_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{I_n}(z_n, \bar{z}_n) \rangle = f(z_i)$$

- Meromorphic!

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent \mathbb{Q}

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{} \times \underbrace{\bar{\mathfrak{sl}}_2}_{}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$
- twisted translations $u_I(\bar{z})$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes
- $$\Delta = 2R + j_1 + j_2$$
- ▶ On plane
$$\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$$
- twisted translations $u_I(\bar{z})$
- ↪ diagonal subalgebra $\bar{\mathfrak{sl}}_2 \times \mathfrak{su}(2)_R$ is \mathbb{Q} exact

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim Q + S$
- Cohomology at the origin \Rightarrow non-empty classes
- $$\Delta = 2R + j_1 + j_2$$
- ▶ On plane
$$\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$$
- twisted translations $u_I(\bar{z})$
- ↪ diagonal subalgebra $\bar{\mathfrak{sl}}_2 \times \mathfrak{su}(2)_R$ is \mathbb{Q} exact
- ↪ anti-holomorphic dependence drops out

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q'$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}'$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

→ $q(z, \bar{z})\tilde{q}(0) \sim$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

$$\rightarrow q(z, \bar{z}) \tilde{q}(0) \sim \bar{z} \tilde{Q}^*(z, \bar{z}) \tilde{Q}(0) \sim$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

$$\rightarrow q(z, \bar{z}) \tilde{q}(0) \sim \bar{z} \tilde{Q}^*(z, \bar{z}) \tilde{Q}(0) \sim \frac{\bar{z}}{z \bar{z}}$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z})$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

$$\rightarrow q(z, \bar{z}) \tilde{q}(0) \sim \bar{z} \tilde{Q}^*(z, \bar{z}) \tilde{Q}(0) \sim \frac{\bar{z}}{z \bar{z}} = \frac{1}{z}$$

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

→ Stress tensor $T_{\mu\nu}$

$4d \ \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

→ Stress tensor $T_{\mu\nu}$ \rightsquigarrow superdescendant

$4d \ \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- Stress tensor supermultiplet

$4d \ \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- Stress tensor supermultiplet

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

- ↪ Global \mathfrak{sl}_2 enhances to Virasoro

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

- ↪ Global \mathfrak{sl}_2 enhances to Virasoro
- ↪ $c_{2d} = -12c_{4d}$

$4d \ \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

→ Theory with flavor symmetry

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Theory with flavor symmetry
- Multiplet containing flavor current

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Theory with flavor symmetry
- Multiplet containing flavor current
- ↪ Affine Kac Moody current algebra

$$J^a(z) J^b(0) \sim -\frac{k_{4d}/2\delta^{ab}}{z^2} + i f^{abc} \frac{J^c(0)}{z} + \dots ,$$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Theory with flavor symmetry
- Multiplet containing flavor current
- ↪ Affine Kac Moody current algebra

$$J^a(z)J^b(0) \sim -\frac{k_{4d}/2\delta^{ab}}{z^2} + if^{abc}\frac{J^c(0)}{z} + \dots,$$

↪ $k_{2d} = -\frac{k_{4d}}{2}$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow chiral algebra

Which operators are in the cohomology?

- Theory with flavor symmetry
- Multiplet containing flavor current
- ↪ Affine Kac Moody current algebra

$$J^a(z) J^b(0) \sim -\frac{k_{4d}/2 \delta^{ab}}{z^2} + i f^{abc} \frac{J^c(0)}{z} + \dots ,$$

- ↪ $k_{2d} = -\frac{k_{4d}}{2}$
- ...

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow Chiral algebra

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow Chiral algebra

→ Cohomology classes \Rightarrow Operators in chiral algebra

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow Chiral algebra

- Cohomology classes \Rightarrow Operators in chiral algebra
- conformal weight $h = R + j_1 + j_2 \geqslant 0$

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow Chiral algebra

- Cohomology classes \Rightarrow Operators in chiral algebra
- conformal weight $h = R + j_1 + j_2 \geqslant 0$
- Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow Chiral algebra

- Cohomology classes \Rightarrow Operators in chiral algebra
- conformal weight $h = R + j_1 + j_2 \geq 0$
- Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
- Very specific non-unitary chiral algebra constrained by unitarity of $4d$ theory

$4d \mathcal{N} \geqslant 2$ SCFT \longrightarrow Chiral algebra

- Cohomology classes \Rightarrow Operators in chiral algebra
- conformal weight $h = R + j_1 + j_2 \geqslant 0$
- Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
- Very specific non-unitary chiral algebra constrained by unitarity of $4d$ theory
 - some operators acquire negative norms

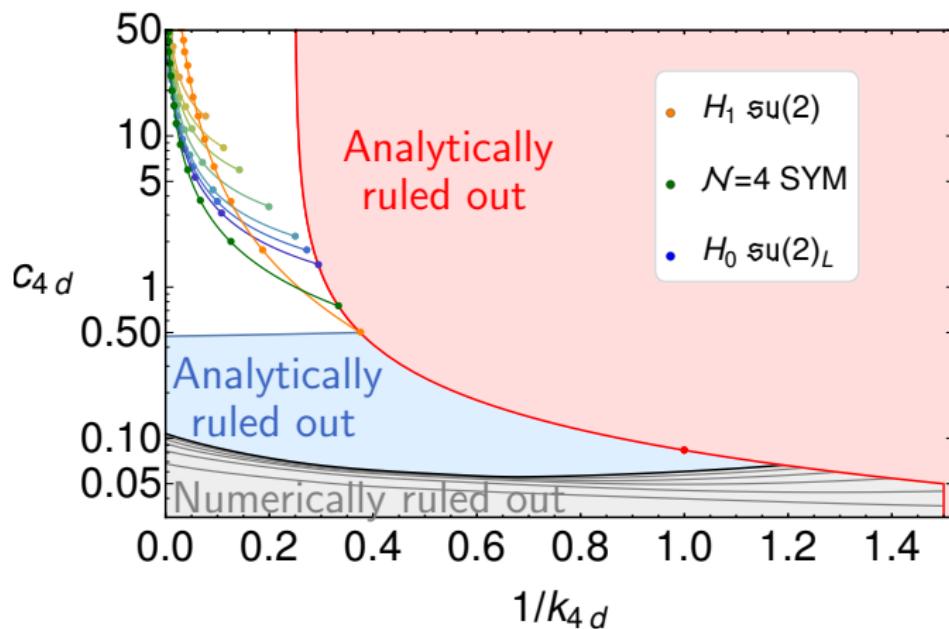
Outline

Inversion formula

- ⑤ Lorentzian inversion formula for (A_1, A_2)
- ⑥ A solvable subsector
- ⑦ Constraining the space of 4d $\mathcal{N} = 2$ SCFTs

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

$su(2)$ flavor symmetry

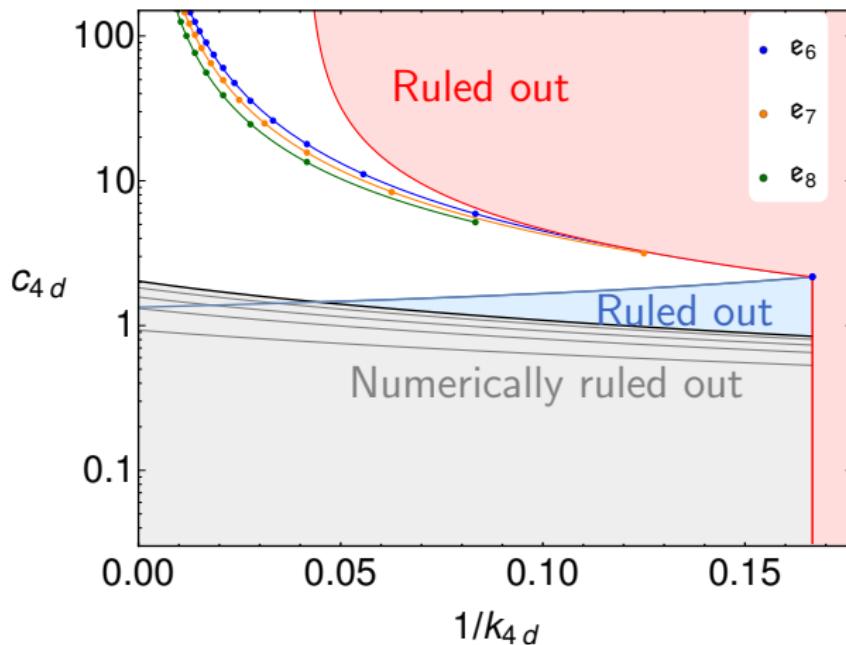


[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]

[Beem, ML, Liendo, Rastelli, van Rees]

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

e_6 flavor symmetry

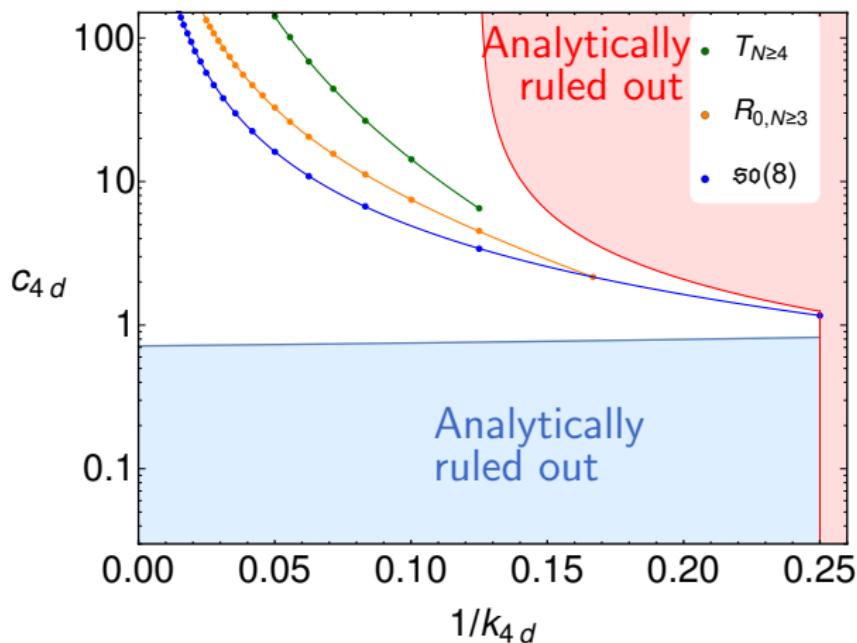


[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]

[Beem, ML, Liendo, Rastelli, van Rees]

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

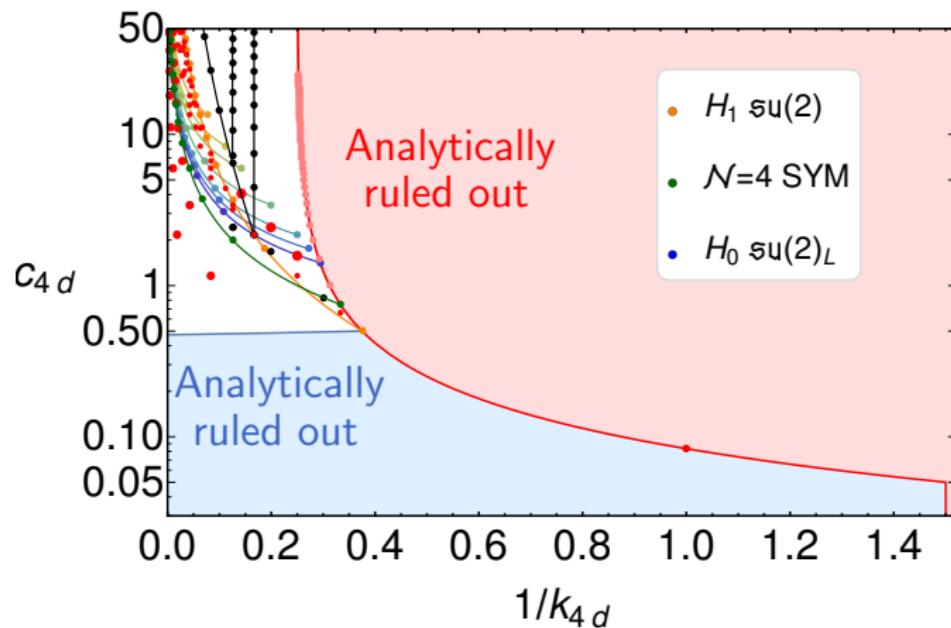
$su(4)$ flavor symmetry



[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

$su(2)$ flavor symmetry



[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]