# Counting instantons in N=1 theories of class S<sub>k</sub>

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[1512.06079 Coman, EP, Taki, Yagi]

[1703.00736 Mitev,EP]

[1712.01288 Bourton, EP]

# Motivation: N=2 exact results

\* Seiberg-Witten theory: effective theory in the IR Classical Integrability

\* Nekrasov: instanton partition function Quantum Integrability

\* Pestun: observables in the UV (path integral on the sphere localizes)

String/M-/F-theory realizations

**\*** Gaiotto: 4D N=2 **class S**: 6D (2,0) on Riemann surface C<sub>g,n</sub>

★ AGT: 4D partition functions = 2D CFT correlators

4D SC Index = 2D correlation function of a TFT

2D/4D relations

## What can we do for N=1 theories?

Superconformal Index

[Romelsberge 2005] [Kinney,Maldacena,Minwalla,Raju 2005]

- *Mathebric Construction of the series of the*
- Witten: IIA/M-theory approach to curves

#### No Localization: No Nekrasov! No Pestun!

# What can we do for N=1 theories?

[Leigh, Strassler 1995] Can construct conformal N=1 theories. 

[Kachru, Silverstein 1998] AdS/CFT natural route to several examples. [Lawrence, Nekrasov, Vafa1998]

6D (1,0) on a Riemann Surface. 

[Gaiotto,Razamat 2015] [Heckman,Vafa....]

Class  $S_k$  ( $S_{\Gamma}$ ): [Gaiotto,Razamat 2015]

2D/4D

relation

#### **\*** Conformal

- **★** Obtained by orbifolding N=2 (inheritance)
- **★** Labeled by punctured Riemann Surface
- **k** Index = 2D correlation function of a TFT

#### Plan

Is there  $AGT_k$ ? 4D partition functions = 2D CFT correlators

- **\*** Introduce N=1 theories in class  $S_k$
- **\*** Spectral curves for N=1 theories in class  $S_k$
- **\*** From the curves: 2D symmetry algebra and representations
- ★ Conformal Blocks → Instanton partition function
- **\*** Instanton partition function from Dp/D(p-4) branes on orbifold



6D (2,0) SCFT on Riemann surface: 4D N=2 theories of class S

Transverse  $C^2/Z_k$  Orbifold the 6D (2,0) SCFT to 6D (1,0) SCFT

6D (1,0) SCFT on Riemann surface: 4D N=1 theories of class Sk

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$	$x^{10}$
N M5-branes			_	_	•	•	_	•	•	•	_
$A_{k-1}$ orbifold	•	•	•	•	_	_	•	_		•	•

#### [Gaiotto,Razamat 2015]

#### [Gaiotto,Razamat 2015]

![](_page_7_Figure_2.jpeg)

 $\epsilon = \Gamma^0 \Gamma^1 \Gamma^2 \Gamma^3 \Gamma^4 \Gamma^5 \epsilon = \Gamma^0 \Gamma^1 \Gamma^2 \Gamma^3 \Gamma^6 \epsilon = \Gamma^4 \Gamma^5 \Gamma^7 \Gamma^8 \epsilon$ 

YM coupling:  $1/g^2$ 

Type IIB

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$
$A_{M-1}$ orbifold	•	•	•	•	•	•				
N D3-branes	_	_			•	•	•	•	•	•
$A_{k-1}$ orbifold	•	•	•	•			•			•

 $\checkmark$  N=1 orbifold daughter of N=4 SYM  $\Gamma = \mathbb{Z}_k \times \mathbb{Z}_M$ 

**Vert** Useful for AdS/CFT (**orbifold inheritance**)  $AdS_5 \times \mathbb{S}^5 / (\mathbb{Z}_k \times \mathbb{Z}_M)$ [Bershadsky, Kakushadze, Vafa 1998]

String theory technics to calculate instantons

[Dorey, Hollowood, Khoze, Mattis,...] [Lerda,...]

#### [Gaiotto,Razamat 2015]

4D field theory point of view

![](_page_9_Figure_3.jpeg)

N=2 class S mother theory

![](_page_9_Figure_5.jpeg)

N=1 class  $S_k$  daughter theory

	$U(1)_t$	$U(1)_{\alpha_c}$	$U(1)_{\beta_{i+1-c}}$	$U(1)_{\gamma_i}$
$V_{(i,c)}$	0	0	0	0
$\Phi_{(i,c)}$	-1	0	-1	+1
$Q_{(i,c-1)}$	+1/2	-1	+1	0
$\widetilde{Q}_{(i,c-1)}$	+1/2	+1	0	-1

![](_page_9_Picture_8.jpeg)

Begin with N=2 class S with SU(kN) factors:

$$W_{\mathcal{S}} = \sum_{c=1}^{M-1} \left( Q_{(c-1)} \Phi_{(c)} \tilde{Q}_{(c-1)} - \tilde{Q}_{(c)} \Phi_{(c)} Q_{(c)} \right)$$

N=2 class S mother theory

![](_page_10_Figure_4.jpeg)

# Curves

[1512.06079 Coman, EP, Taki, Yagi]

#### **Curves from M-theory**

![](_page_12_Figure_1.jpeg)

![](_page_12_Figure_2.jpeg)

The NS5/D4 is the classical configuration.

 $v = x^4 + ix^5$ 

Take in account tension of the branes: include quantum effects.

M-theory: a single MS brane with non trivial topology

$$(t-1)(t-q)v^2 - P_1(t)v + P_2(t) = 0$$

 $(v - m_1)(v - m_2)t^2 + (-(1 + q)v^2 + qMv + u)t + q(v - m_3)(v - m_4) = 0$ coupling constant  $q = e^{2\pi i \tau w}$   $u = tr \phi^2$   $M = m_1 + m_2 + m_3 + m_4$ 

![](_page_13_Figure_0.jpeg)

$$\begin{array}{c} (v^{k} - m_{1}^{k})(v^{k} - m_{2}^{k})t^{2} + P(v)t + q(v^{k} - m_{3}^{k})(v^{k} - m_{4}^{k}) = 0 \\ P(v) = -(1 + q)v^{2k} + u_{k}v^{k} + u_{2k} \\ \text{vevs of gauge invariant operators:} \\ \text{parameterize the Coulomb branch} \\ \langle \operatorname{tr}(\Phi_{(1)} \cdots \Phi_{(k)}) \rangle \sim u_{k} \\ \langle \operatorname{tr}(\Phi_{(1)} \cdots \Phi_{(k)})^{2} \rangle \sim u_{2k} \end{array}$$

 $-m_4$ 

 $-m_3$ 

 $R^4/Z_2$ 

$$\phi_{k\ell}^{(3)}(t) = (-1)^{\ell} \frac{\mathfrak{c}_L^{(\ell,k)} t - \mathfrak{c}_R^{(\ell,k)}}{t^{k\ell}(t-1)} \quad \text{for } \ell = 1$$

The above coefficients can be directly obtained by taking the limit q

$$u_{k\ell}(q=0) \longrightarrow (-1)^{\ell+1} \mathfrak{c}_R^{(\ell,k)} \cdot \mathfrak{c}_R^{(\ell,k)$$

**IOP PATE** The UV curves corresponding to the free trinion and to the SCQCD theory three and respectively four punctured<sup>5</sup> spheres with the punctures at twhile those at t = 1 and t = q are simple punctures  $\bullet$ , see [28].

![](_page_14_Figure_5.jpeg)

nvelation

Figure 2: The UV curves of the trinion and of the  $SCQCD_k$  theories. They are  $SCQCD_k^{tirely}$  the spectral production of the second transformed at t = 0 and t = 0 an at t = 1 and at t = q. then we find

Gaiotto Shifts in x for k = 1. Due to the orbifold relation (2.2), we are allowed to shift the variable  $d_{t} = 0$ k = 1, but not for k > 1. This shift is the consequence of the additional U(1) degrees of freedom that for present for k = 1 but, as we shall see more in detail later, disappear for k > 1. For k = 1, if we get  $k \ge 1$ , if k = 1, if  $k \ge 1$ , if equation  $\sum_{i=0}^{N} x^{i} \phi_{i}$  to  $\sum_{i=0}^{N} x^{i} \phi'_{i}$  by making the transformation  $x \to x - \kappa \phi_{1}$ , then we find B  $\searrow$  (C)  $-\kappa\phi_1$ , then we find

the and the h We remind that  $\phi_0 = 1$  before and after the transformation. It is clear that the shift leaves  $\Omega_2 = d\lambda_{SW} = dx \wedge dt$  unchanged however the structure of the poles of  $\lambda_{SW}$  on the various sheets of the does change, see [28]. If we put the shift parameter  $\kappa$  equal to  $\frac{1}{N}$ , then the coefficient  $\phi'_1$  vanishes curve is known as the Galotto curve. /Let us denote the curve coefficients for the Galotto curve b barameters we shall review later, their expansion around the poles in t gives the charges of the  $\mathbf{W}_N$  algebra. mation Parameteringbine para <sup>5</sup>The UV curves are characterized by the meromorphic differentials  $\phi_s^{(n)}$  that have only poles and no branch cuts. T punctures  $\star$  discussed in [28] will not be relevant for our purposes here. instead of the U(1) instead of the UE

Figure 9.13: Triality, using trivalent diagrams.

ty looks like 19 equation!

# Instantons from the 2D Blocks

 $\mathcal{Z}_{\text{inst}} = \mathcal{B}_{\mathbf{w}}(\mathbf{w}_1, \mathbf{w}_2, \mathbf{w}_3, \mathbf{w}_4 | q)$ 

[1703.00736 Mitev, EP]

![](_page_16_Figure_1.jpeg)

**Figure 2:** The UV curves of the trinion and of the  $SCQCD_k$  spheres. The full punctures are depicted by  $\odot$  and placed at t = 0 at t = 1 and at t = q.

Gaiotto Shifts in x for k = 1. Due to the orbifold relation ( k = 1, but not for k > 1. This shift is the consequence of the present for k = 1 but, as we shall see more in detail later, disappendix equation  $\sum_{i=0}^{N} x^i \phi_i$  to  $\sum_{i=0}^{N} x^i \phi'_i$  by making the transformation xB A  $\begin{pmatrix} Q \\ \phi'_{\ell} = \sum_{j=N-\ell}^{N} {j \choose N-\ell} \phi_{N-j} (-\kappa \phi_1)^{j+i-N} = \sum_{j=N-\ell}^{N} {j \choose N-\ell} (-\kappa \phi_1)^{j+i-N} (-\kappa \phi_1)^{j+i-N} = \sum_{j=N-\ell}^{N} {j \choose N-\ell} (-\kappa \phi_1)^{j+i-N} (-\kappa \phi$ 

 $\Omega_2 = d\lambda_{SW} = dx \wedge dt \text{ unchanged/ however the structure of the p}$ does change, see [28]. If we put the shift parameter  $\kappa$  equal to  $\frac{1}{N}$ , 1/q curve is known as the Galotto curve. Let us denote the curve we shall review/later, their expansion around the poles in t give  $\frac{D}{A}$ 

С

А

<sup>5</sup>The UV curves are characterized by the meromorphic differentials φ<sub>s</sub><sup>(n)</sup> the punctures ★ discussed in [28] will not be relevant for our purposes here.
 Figure 9.13: Triality, using trivalent diagrams.

### From the curves to the 2D CFT

$$\lim_{\epsilon_{1,2}\to 0} \left\langle \left\langle J_{\ell}(t) \right\rangle \right\rangle_n = \phi_{\ell}^{(n)}(t)$$

 $\left<\left< J(t) \right>\right>_n \stackrel{\mathrm{def}}{=} \frac{n\text{-point W-block with insertion of } J(t)}{n\text{-point W-block}}$ 

**\*** The symmetry algebra that underlies the 2D CFT =  $W_{kN}$  algebra

**\*** The reps are **standard** reps of the  $W_{kN}$  algebra

★ Obtain them from the N=2 SU(kN) after replacing:

$$m_{j+Ns}^{\mathrm{SU}(Nk)} \longmapsto m_j \,\mathrm{e}^{\frac{2\pi i}{k}s} \qquad a_{j+Ns}^{\mathrm{SU}(Nk)} \longmapsto a_j \,\mathrm{e}^{\frac{2\pi i}{k}s}$$

[1703.00736 Mitev, EP]

# 2D Conformal Blocks = Instanton P.F.

K We have the reps of the  $W_{kN}$  algebra for  $\varepsilon_{1,2} = 0$  (from the curve)

**\*** Demand: the structure of the multiplet (null states) not change  $\varepsilon_{1,2} \neq 0$ 

**\*** The blocks for  $\varepsilon_{1,2} \neq 0$ : proposal for the instanton partition functions:

$$\mathcal{Z}_{\text{inst}} = \mathcal{B}_{\mathbf{w}}(\mathbf{w}_1, \mathbf{w}_2, \mathbf{w}_3, \mathbf{w}_4 | q)$$

If w and c turn on  $Q \neq 0$  as in Liouville/Toda,

then we obtain them from the N=2 SU(kN) after replacing:

$$m_{j+Ns}^{\mathrm{SU}(Nk)} \longmapsto m_j e^{\frac{2\pi i}{k}s} \qquad a_{j+Ns}^{\mathrm{SU}(Nk)} \longmapsto a_j e^{\frac{2\pi i}{k}s}$$

[1703.00736 Mitev, EP]

# Instantons from D(p-4) branes

[1712.01288 Bourton, EP]

#### **Instantons from D branes**

Tr 
$$\int_{Dp} d^{p+1}x \quad C_{p-3} \wedge F \wedge F$$

#### Instanton on Dp brane = D(p-4) brane

$$\mu^{\mathbf{C}} := [B_1, B_2] + IJ = 0 ,$$
  
$$\mu^{\mathbf{R}} := [B_1, B_1^{\dagger}] + [B_2, B_2^{\dagger}] + II^{\dagger} - J^{\dagger}J = 0$$

The Dp-D(p-4) strings give the NXK I, Jt and the D(p-4)-D(p-4) the KXK B1, B2 auxiliary matrices of the ADHM construction

$$\mathcal{M}_{K\text{-inst}}^{Dp} \simeq \mathcal{M}_{\mathrm{Higgs}}^{D(p-4)}$$

[Witten 1995, Douglas 1995, Dorey 1999, ...]

$$KD(p-4)$$

$$= \{B_1, B_2, I, J \mid \mu^{\mathbf{C}} = 0, \mu^{\mathbf{R}} = 0\} / U(K)$$
$$= \{X_{Dp}^{\perp} = 0, \mathcal{V}_{D(p-4)} = 0\} / U(K)$$

![](_page_21_Figure_0.jpeg)

## Mass deformed N=4 SYM

First we practice without the orbifolds:

\* Obtain the  $\mathcal{Z}_{inst}$  for mass deformed N=4 SYM (N=2\*) from a 2D SCI computation: 2D theory with (4,4) susy

$$D3/D(-1) \xleftarrow{\text{T-duality}} D5/D1$$

N

K

- ✤ The 2D SCI of the gauge theory on the D1 branes
  - = instantons of the 6D (2,0) on  $\mathbb{R}^4 \times T^2$  = M-strings.
- **\*** KK reducing: instantons of mass deformed 5D N=2 MSYM on  $\mathbb{S}^1$  and further KK reduce to mass deformed N=4 SYM in 4D.

### Instantons with an orbifold

- **\*** "Orbifold" the 2D SCI of mass deformed N=4 SYM with one, say the  $Z_M$  orbifold: we get M-strings on a transverse orbifold
  - = instantons of an  $SU(N)^{M}$  quiver when reduce down to 5D/4D

2D theory with (0,4) susy

Further "Orbifold" the 2D SCI with the Z<sub>k</sub> orbifold we get something new that should correspond to instantons of class Sk, the rational, trigonometric and elliptic uplift.
2D theory with (0,2) susy

#### S<sub>k</sub> Instantons

 $\mathcal{Z}_{K\text{-inst}}^{\text{D5}}(a, m, \epsilon_1, \epsilon_2) = \mathcal{Z}_{\text{Higgs}}^{K \text{ D1}}(a, m, \epsilon_1, \epsilon_2) = \mathcal{I}_{2D} = \text{Tr}_{\mathcal{M}_{\text{Higgs}}^{K \text{ D1}}}(-1)^F e^{\epsilon_1 J_L} e^{\epsilon_2 J_R} e^{\xi R} e^{a J_G} e^{m J_F}$ 

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$
N D5	—	—	_	_	_	_	•	•	•	•
$\mathbb{Z}_M$	•	•	•	•	•	•	×	×	×	×
$\mathbb{Z}_k$	•	•	•	•	×	×	•	×	×	•
K D1	•	•	•	•	—		•	•	•	•

![](_page_24_Figure_3.jpeg)

The same answer as in [1703.00736 Mitev, EP] from conformal blocks!

#### **Instantons from D0 branes**

![](_page_25_Figure_1.jpeg)

#### Summary

#### $\mathbf{M}$ We constructed spectral **curves** for N=1 theories in class $S_k$

 $\mathcal{M}$  The curves: 2D symmetry algebra ( $W_{kN}$ ) and representations

#### Conformal Blocks Instanton partition function

#### *Market Series of Contraction From Dp/D(p-4) on orbifold*

## **Future**

Compute one, two instantons with standard QFT techniques Ж

![](_page_27_Picture_2.jpeg)

**C** Go away from the orbifold point [to appear Bourton, EP]

![](_page_27_Picture_4.jpeg)

![](_page_27_Picture_5.jpeg)

Other N=1 theories

- The perturbative part? N=1 partition function on S<sup>4</sup>! \*
- Get the AGT<sub>k</sub> from (1,0) 6D à la Cordova and Jafferis \*

# Thank you!