

# Cosmic Inflation Tutorial

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Niels Bohr Institute  
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## Cosmic Inflation:

- Concrete technical idea with great phenomenology
- Couched in the language of “solving tuning problems”
- Big open questions

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# Conclusions

- A) Inflation is a technical tool for connecting cosmological observables with high energy physics. Impressive successes.
- B) However, without a meta theory about how inflation started (and how it “competes” with other scenarios, such as the Standard Big Bang) big questions are unresolved.
- C) Complex sociology as different individuals choose to give A) and B) different weight.

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Priors!!

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# OUTLINE

1. Big Bang basics
2. Inflation basics
3. Eternal inflation
4. An alternative to Eternal Inflation
5. Further thoughts

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# Einstein equation + Homogeneity & Isotropy →

## Friedmann Eqn.


$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_k + \rho_r + \rho_m + \rho_{DE})$$

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Hubble parameter  
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"Scale factor"

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The diagram shows the Friedmann equation enclosed in a red rectangular box. Above the box, the text  $\propto a^{-2}$  has a blue arrow pointing down to the  $\rho_k$  term in the equation. Below the box, the word "Curvature" has a red arrow pointing up to the  $\left( \frac{\dot{a}}{a} \right)^2$  term. Further down, the text "Scale factor" has a red arrow pointing up to the  $a$  in the denominator of the fraction inside the curvature term.

Curvature

"Scale factor"

## Friedmann Eqn.

$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_k + \rho_r + \rho_m + \rho_{DE})$$

$\propto a^{-2}$

$\propto a^{-4}$

Curvature

Relativistic Matter

“Scale factor”



## Friedmann Eqn.

$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_k + \rho_r + \rho_m + \rho_{DE})$$

Diagram illustrating the Friedmann Equation and the scaling of its components:

- $\propto a^{-2}$  (Curvature term,  $\rho_k$ )
- $\propto a^{-4}$  (Relativistic Matter term,  $\rho_r$ )
- $\propto a^{-3}$  (Non-relativistic Matter term,  $\rho_m$ )

Labels for the terms in the equation:

- Curvature (points to  $\rho_k$ )
- Relativistic Matter (points to  $\rho_r$ )
- Non-relativistic Matter (points to  $\rho_m$ )
- "Scale factor" (points to  $a$  in the denominator of  $\dot{a}/a$ )

# Friedmann Eqn.

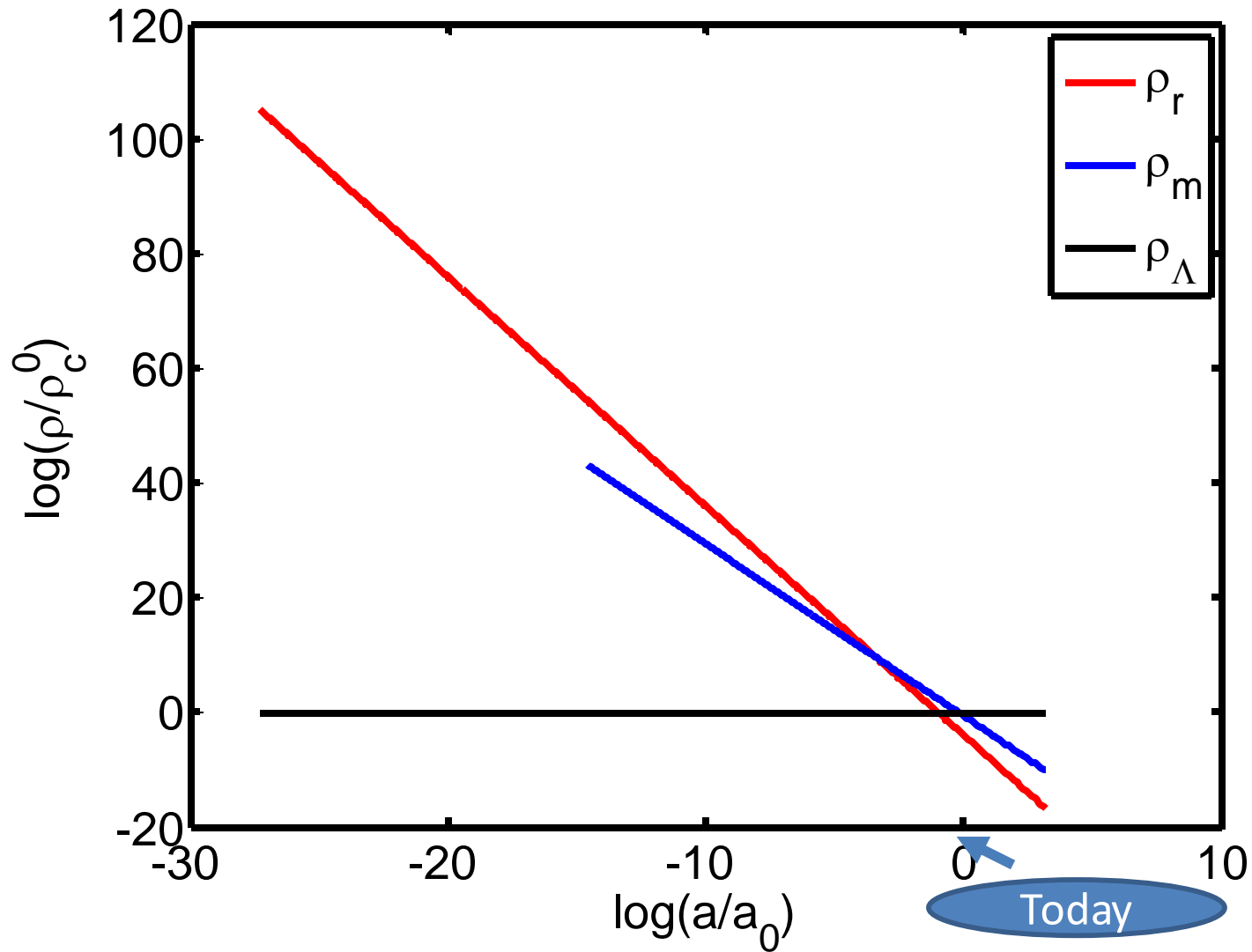
The diagram illustrates the Friedmann equation,  $H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3}G(\rho_k + \rho_r + \rho_m + \rho_{DE})$ , enclosed in a red box. Annotations include:

- Curvature:** An orange arrow points from the label to the  $\left(\frac{\dot{a}}{a}\right)^2$  term.
- “Scale factor”:** An orange arrow points from the label to the  $\dot{a}$  in the numerator of the curvature term.
- Relativistic Matter:** An orange arrow points from the label to the  $\rho_r$  term.
- Non-relativistic Matter:** An orange arrow points from the label to the  $\rho_m$  term.
- Dark Energy:** An orange arrow points from the label to the  $\rho_{DE}$  term.

Scaling relations for the density terms are shown with blue arrows:

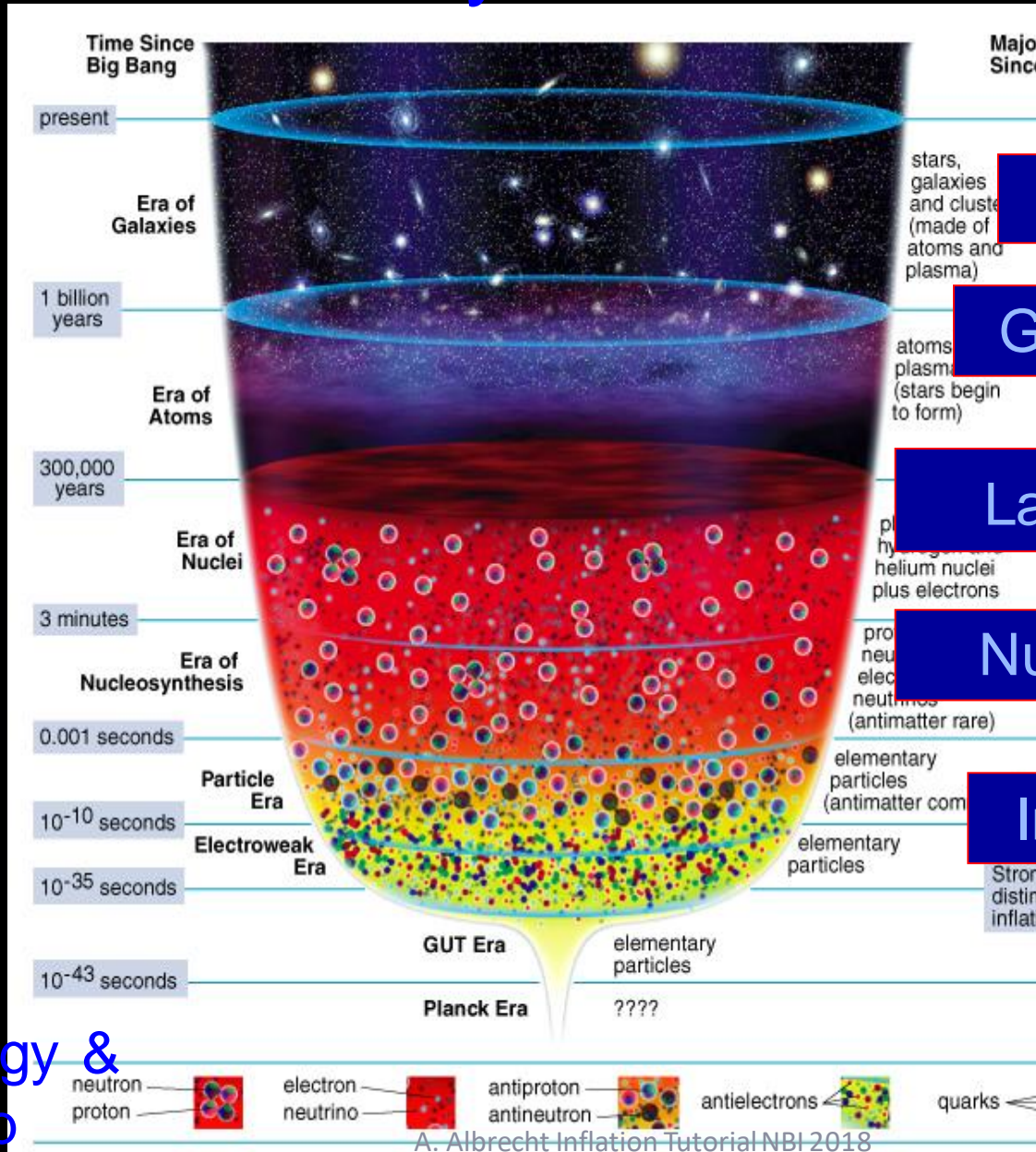
- $\rho_k \propto a^{-2}$
- $\rho_r \propto a^{-4}$
- $\rho_m \propto a^{-3}$
- $\rho_{DE} \propto a^{\approx 0}$

# Evolution of Cosmic Matter



Time

# The History of the Universe



Today

Dark Energy

Galaxy Formation

Last Scattering

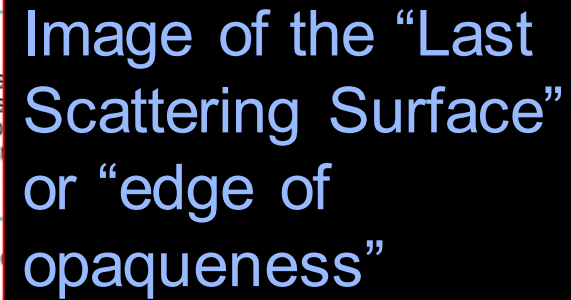
Nuclear & HEP

Inflation?

Extra Dimensions?

High Energy & Temp

# The History of the Universe





Time

# The History of the Universe

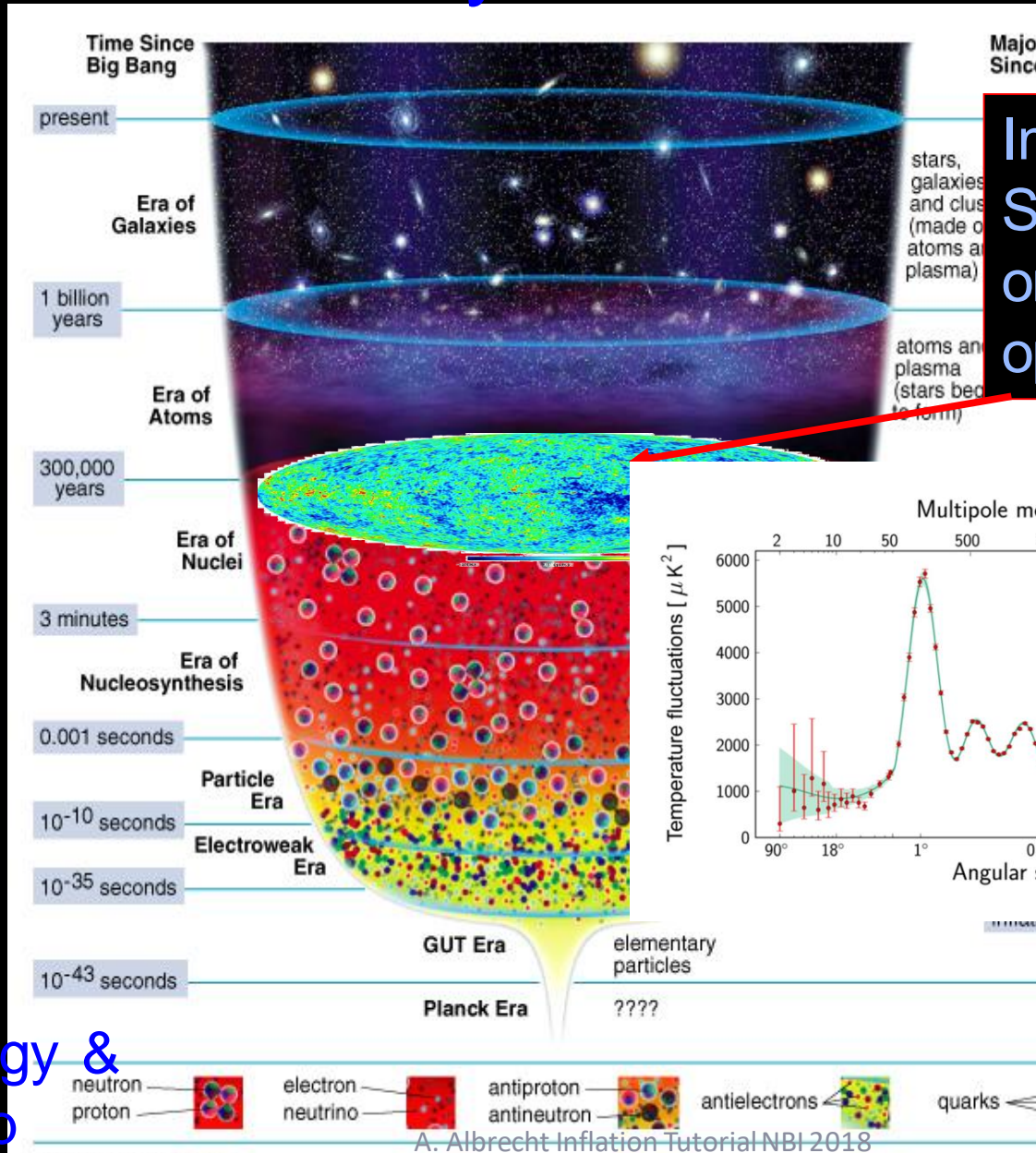
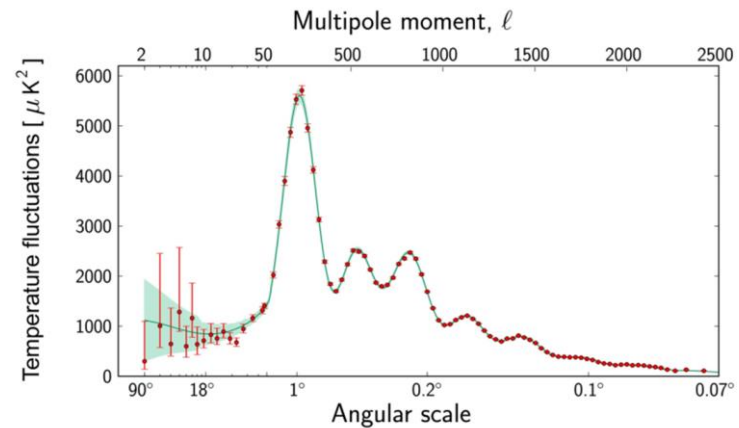
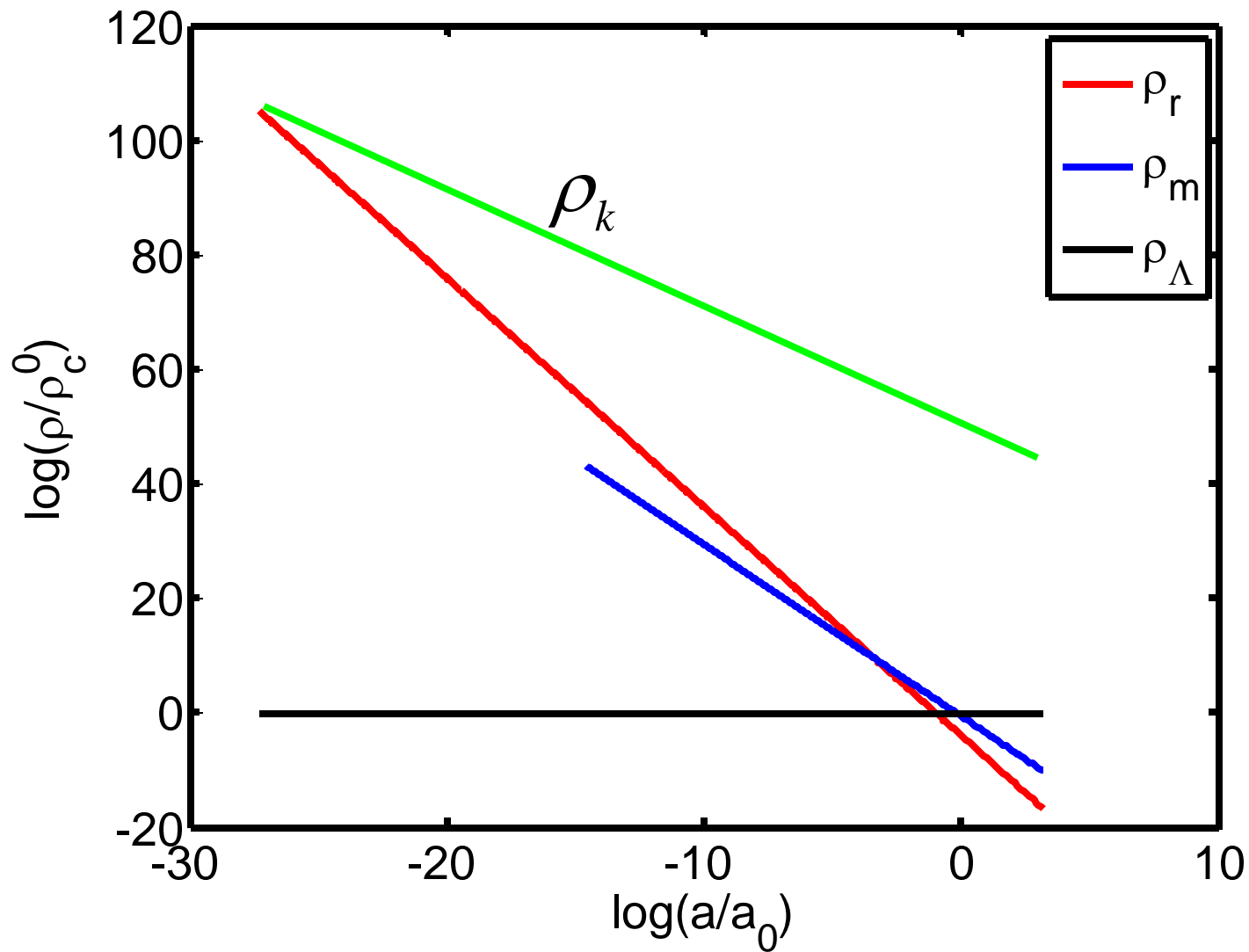


Image of the “Last Scattering Surface” or “edge of opaqueness”

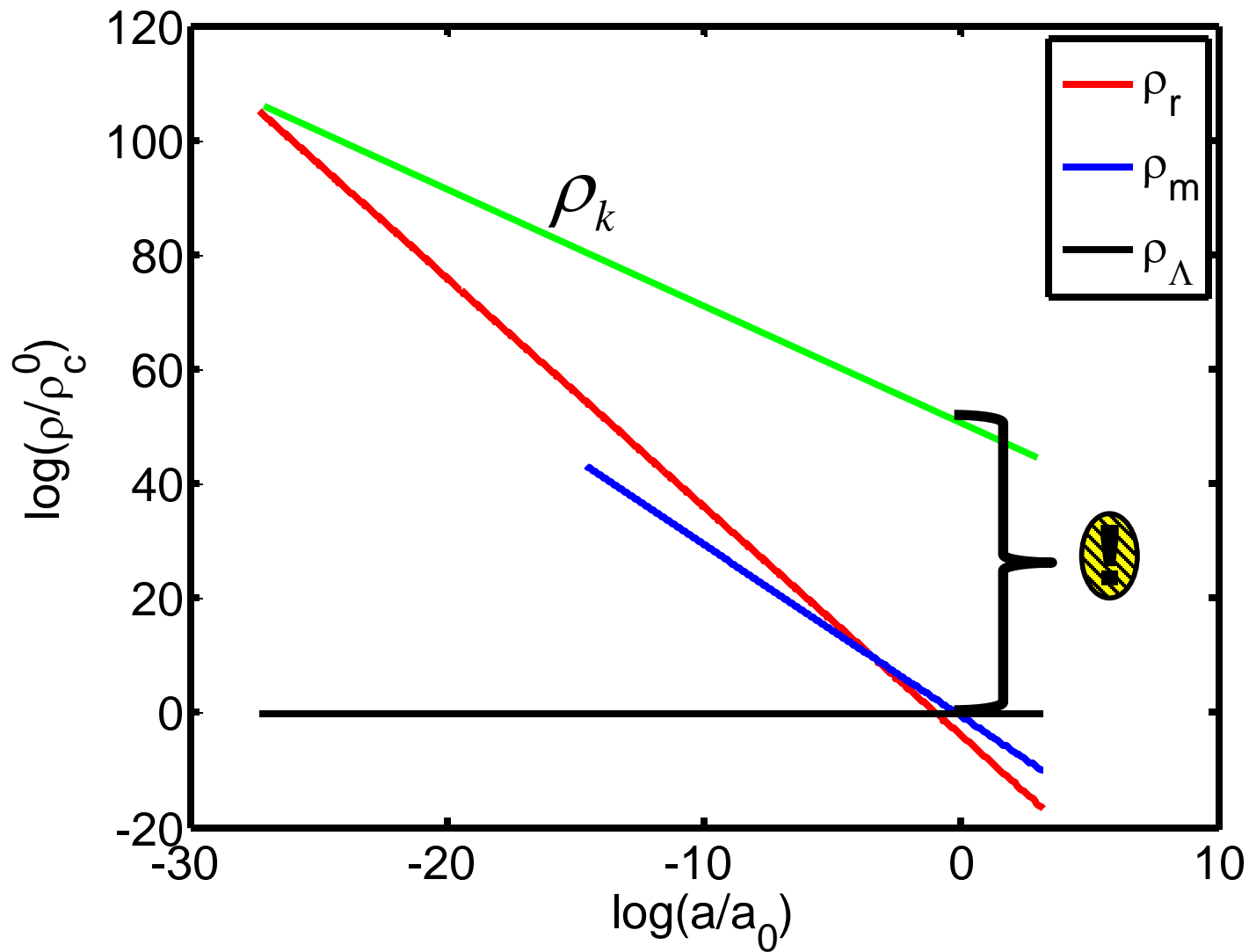


High Energy & Temp

## The curvature feature/“problem”

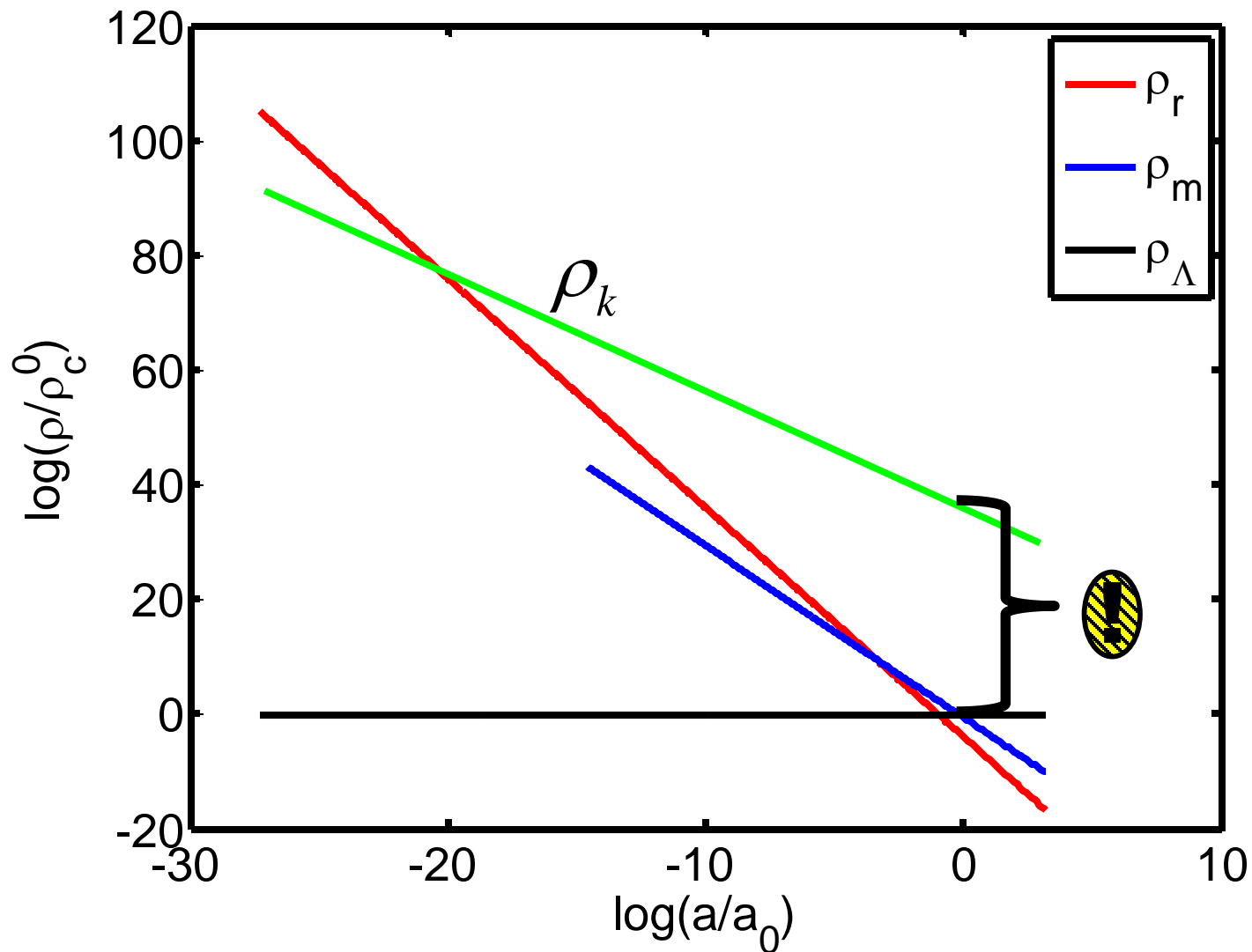


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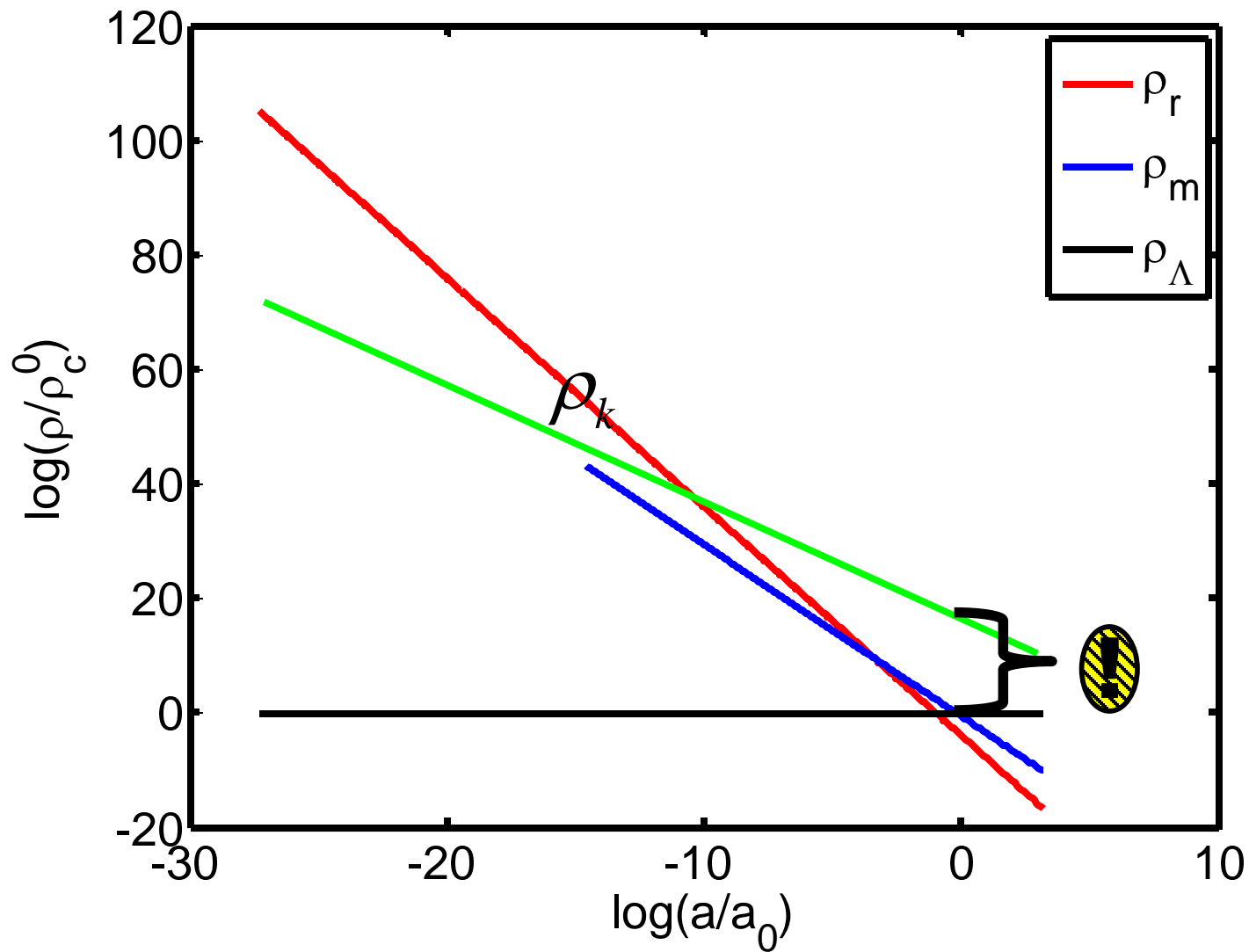




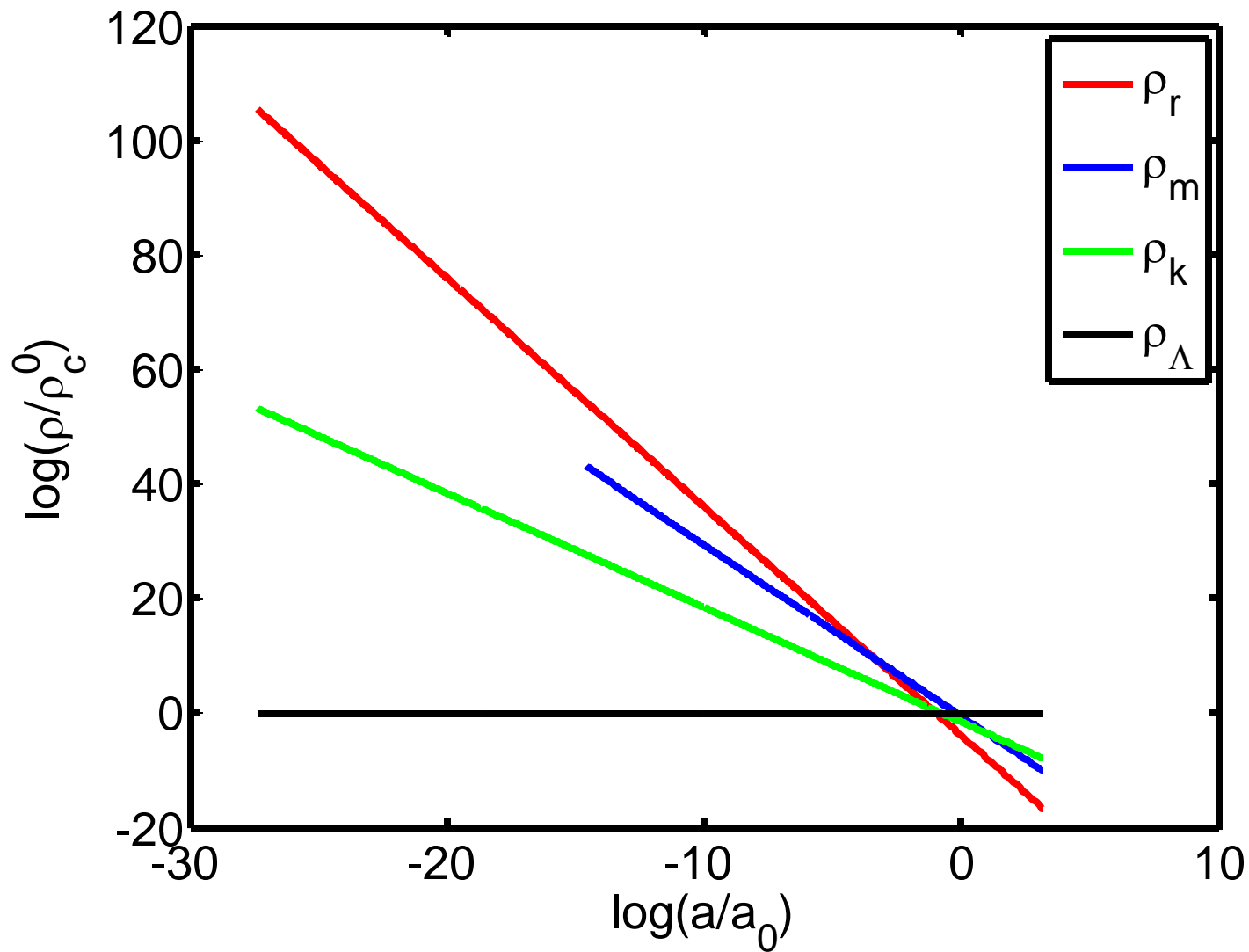
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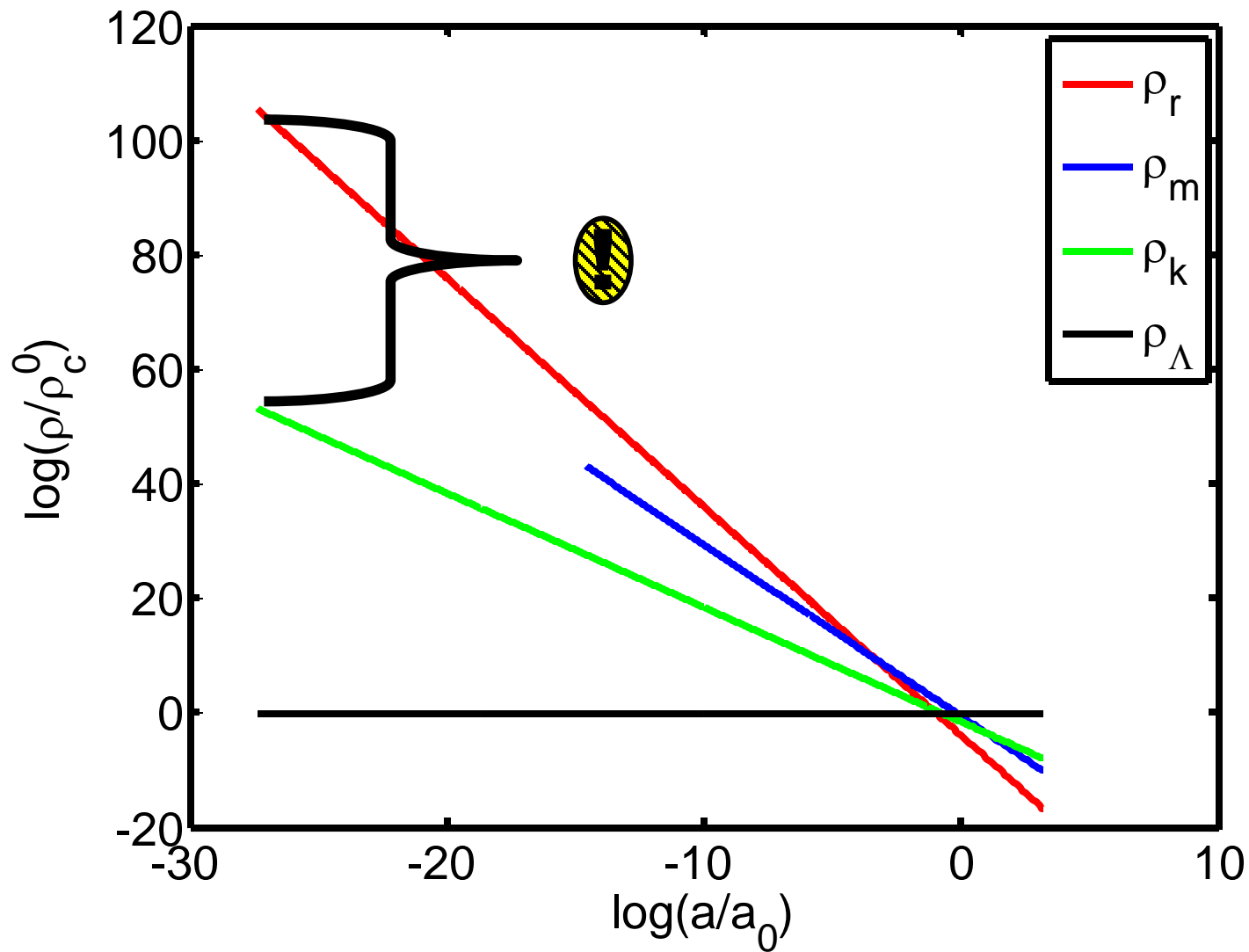
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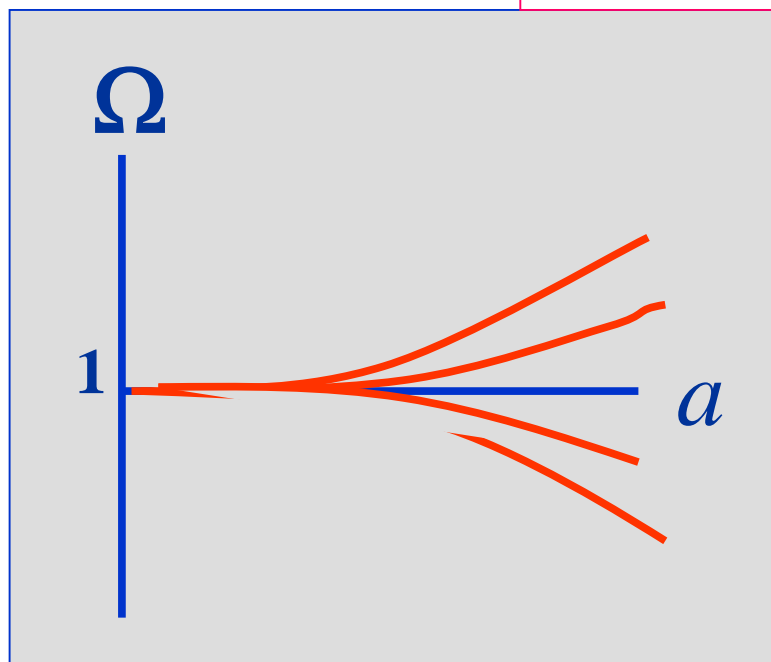
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In the SBB, flatness is an “unstable fixed point”:



At  $T = 10^{16} \text{ GeV}$

or  $\frac{a}{a_0} = 10^{-28}$

The “GUT scale”

$$\Omega \equiv \frac{\rho}{\rho_c}$$

$$\rho_c \equiv \frac{3H^2}{8\pi}$$

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2 = \frac{8\pi}{3} \rho - \frac{k}{a^2}$$

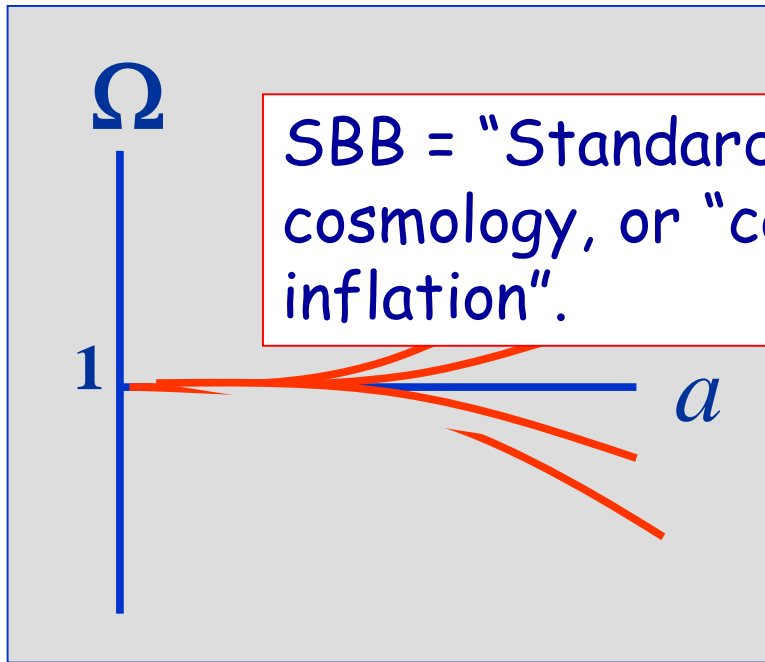
$$\propto a^{-3}, a^{-4}$$

Dominates  
with time

Require  $\rho = \rho_c$  to 55 decimal places to get  $\rho \approx \rho_c$  today

In the SBB, flatness is an “unstable fixed point”:

i



SBB = “Standard Big Bang” cosmology, or “cosmology without inflation”.

$$\rho_c \equiv \frac{3H^2}{8\pi}$$

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2 = \frac{8\pi}{3} \rho - \frac{k}{a^2}$$

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Require  $\rho = \rho_c$  to 55 decimal places to get  $\rho \approx \rho_c$  today

## Gravitational instability: The Jeans Length

$$R_{Jeans} \equiv \lambda_J \equiv c_s \left( \frac{\pi c^2}{G \rho} \right)^{1/2}$$

Sound speed

Average energy  
density

- Overdense regions of size  $> R_{Jeans}$  collapse under their own weight.

If the size is  $< R_{Jeans}$  they just oscillate

## SBB Homogeneity:

On very large scales the Universe is highly homogeneous, despite the fact that gravity will clump matter on scales greater than  $R_{\text{Jeans}}$

At the GUT epoch the observed Universe consisted of  $10^{79} R_{\text{Jeans}}$  sized regions.

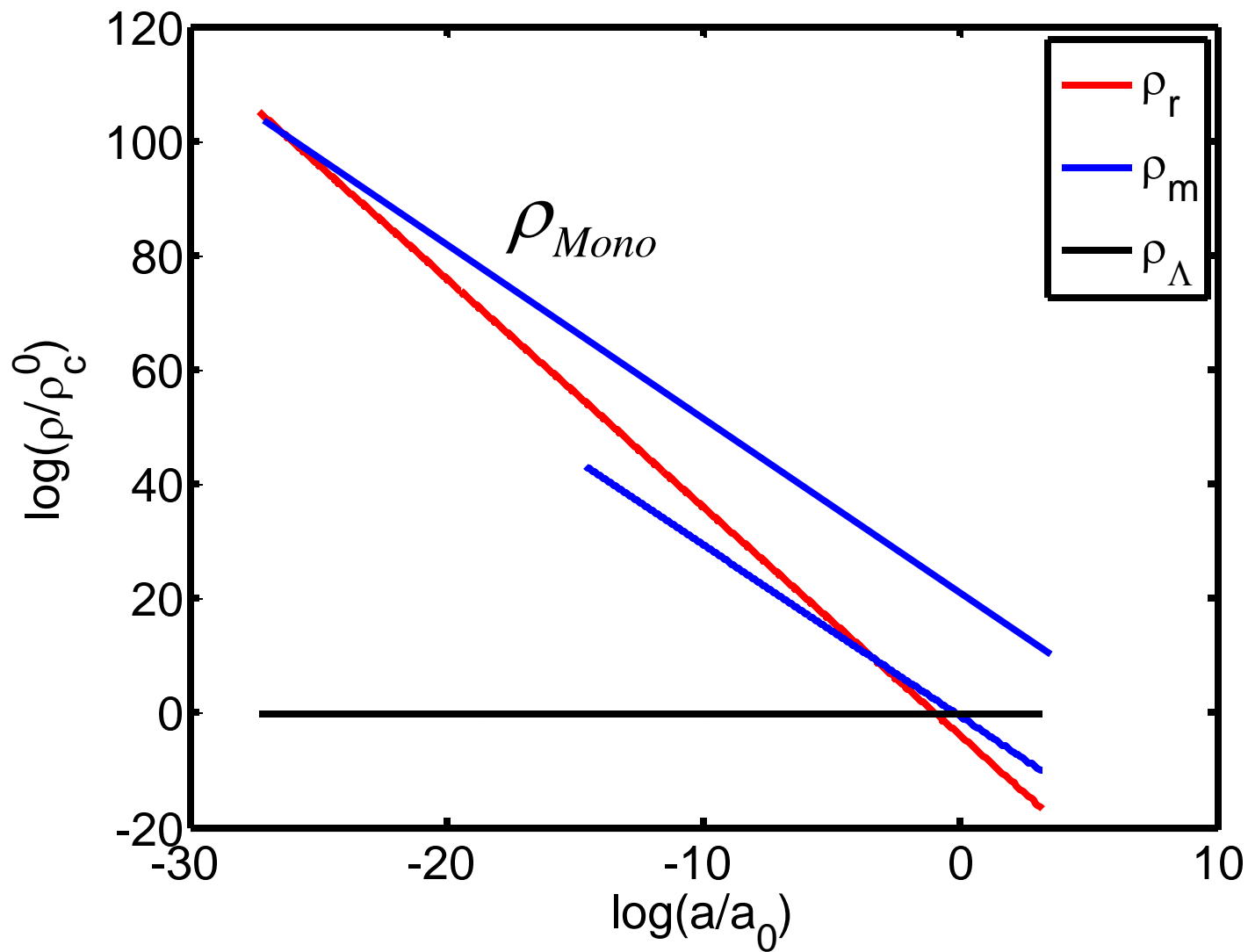
→ The Universe was very smooth to start with.

NB: Flatness & Homogeneity → SBB Universe starts in highly *unstable* state.

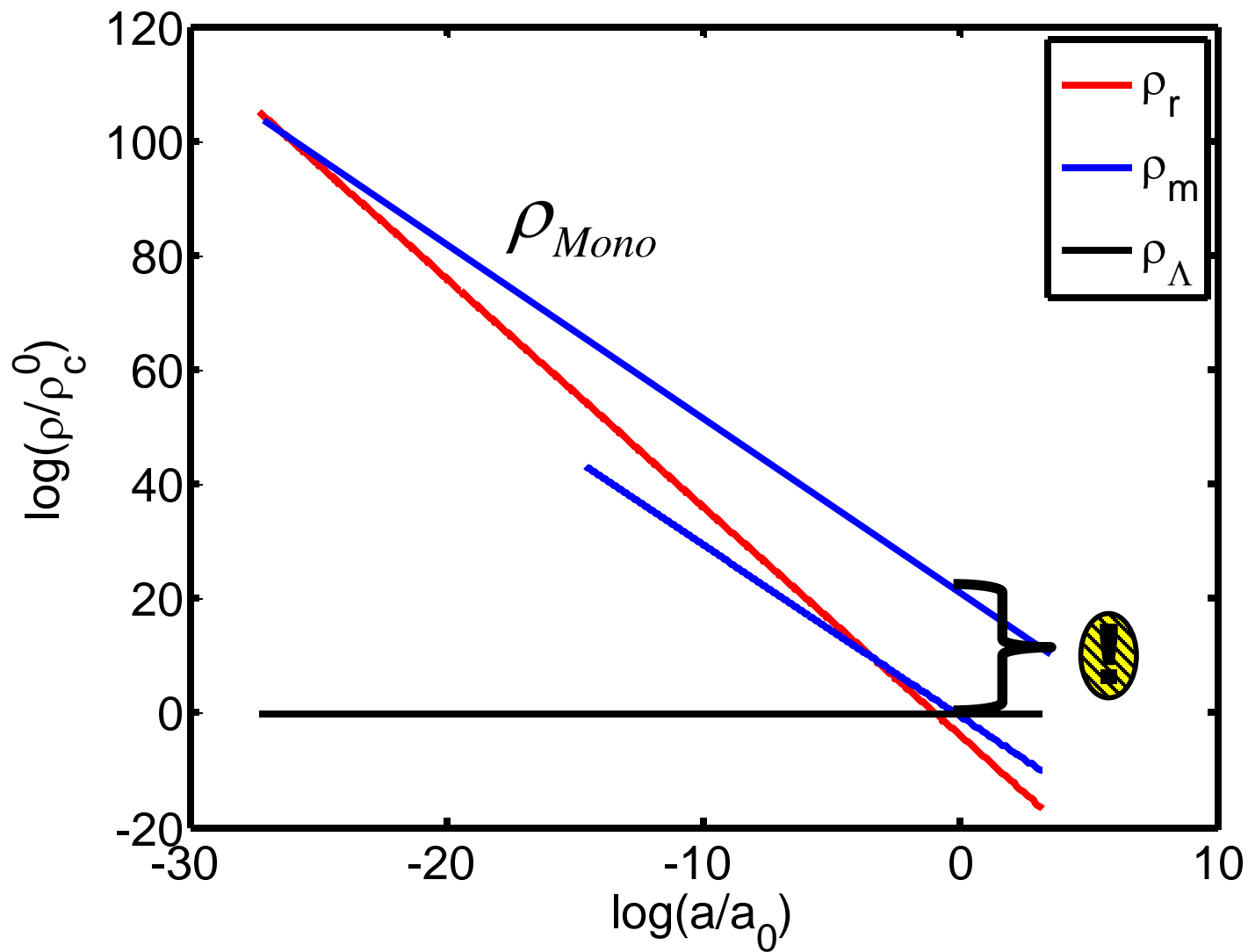
$$S_{\text{Univ}} \approx 10^{-35} S_{\text{bh-Max}} = 10^{-35} 4\pi M_{\text{Univ}}^2$$



## The monopole “problem”



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## SBB Monopoles

- A GUT phase transition (or any other process) that injects stable non-relativistic matter into the universe at early times (deep in radiation era, ie  $T_i = 10^{16}$  GeV) will *\*ruin\** cosmology:

$$\left. \frac{\rho_M}{\rho_{Normal}} \right|_T = \frac{\rho_M(T_i) \left( \frac{T}{T_i} \right)^3}{\rho_{Normal}(T_i) \left( \frac{T}{T_i} \right)^4} = \left( \frac{\rho_M(T_i)}{\rho_{Normal}(T_i)} \right) \times \left( \frac{T_i}{T} \right) \longrightarrow " \infty "$$

Monopole dominated Universe



## SBB Horizon

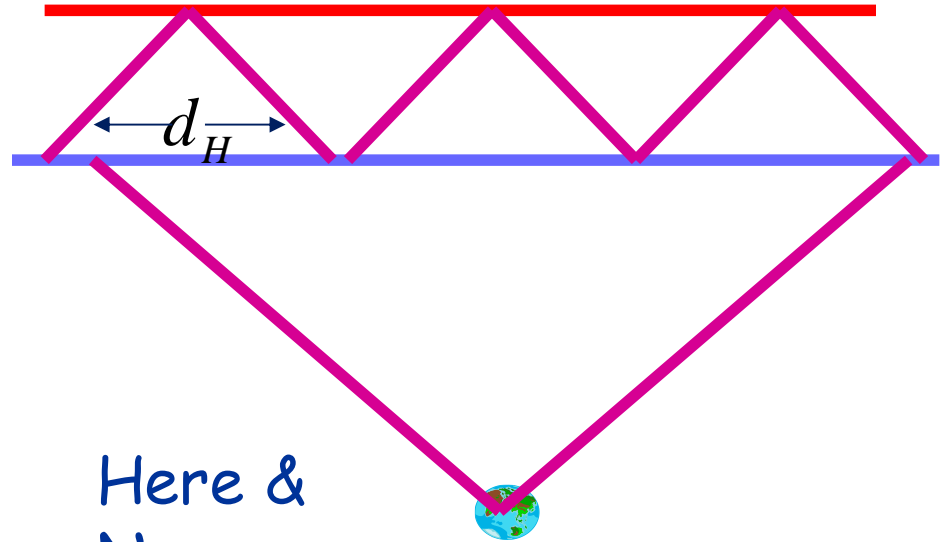
$t=0$

$10^{80}$  causally disconnected regions  
at the GUT epoch

Horizon: The distance light has  
traveled since the big bang:

$$d_H = a(t) \int_0^t \frac{dt'}{a(t')}$$

Here &  
Now



The flatness, homogeneity & horizon features become “problems” if one feels one must explain initial conditions.

Basically, the SBB says the universe must start in a highly balanced (or “fine tuned”) state, like a pencil on its point.

Must/can one explain this?

Inflation says “yes”

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## Friedmann Eqn.

The diagram shows the Friedmann equation enclosed in a red box:

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3} G (\rho_k + \rho_r + \rho_m + \rho_{DE})$$

Annotations above the box indicate the scaling of each density term with the scale factor  $a$ :

- $\propto a^{-2}$  points to  $\rho_k$  (Curvature)
- $\propto a^{-4}$  points to  $\rho_r$  (Relativistic Matter)
- $\propto a^{-3}$  points to  $\rho_m$  (Non-relativistic Matter)
- $\propto a^{\approx 0}$  points to  $\rho_{DE}$  (Dark Energy)

Annotations below the box identify the physical components:

- Curvature (points to  $\rho_k$ )
- Relativistic Matter (points to  $\rho_r$ )
- Non-relativistic Matter (points to  $\rho_m$ )
- Dark Energy (points to  $\rho_{DE}$ )



# Now add cosmic inflation

Friedmann Eqn.

The diagram shows the Friedmann equation with annotations for different energy components and their scaling with the scale factor  $a$ :

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3} G (\rho_I + \rho_k + \rho_r + \rho_m + \rho_{DE})$$

Annotations and scaling relations:

- $\rho_I$  (Inflaton):  $\propto a^{\approx 0}$  (circled in red)
- $\rho_k$  (Curvature):  $\propto a^{-2}$
- $\rho_r$  (Relativistic Matter):  $\propto a^{-4}$
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Labels for the energy components:

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$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_I + \rho_k + \rho_r + \rho_m + \rho_{DE})$$

The diagram illustrates the Friedmann equation with the following annotations:

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$$H_I = \frac{\dot{a}}{a} \approx \text{const} \rightarrow a \approx e^{H_I t}$$

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Friedmann Eqn.

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$\propto a^{\approx 0}$

Inflaton

Curvature

Relativistic Matter

Non-relativistic Matter

Dark Energy

# Now add cosmic inflation

Fried

Just  
suppose!

$$\propto a^{\approx 0}$$

$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_I + \rho_k + \rho_r + \rho_m + \rho_{DE})$$

Inflaton

Curvature

Relativistic Matter

Non-relativistic  
Matter

Dark  
Energy

# The inflaton:

~Homogeneous scalar field  $\phi$  obeying

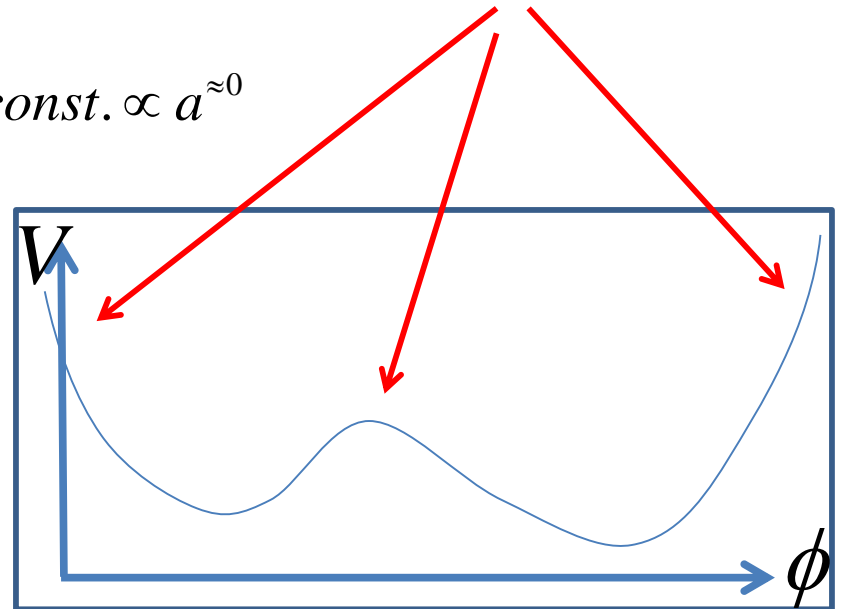
$$\ddot{\phi} + 3H\dot{\phi} = -\Gamma_{\phi}\dot{\phi} - V'(\phi)$$

*Cosmic damping*

*Coupling to ordinary matter*

All potentials have a “low roll” (overdamped) regime where

$$\rho_I = \frac{1}{2}\dot{\phi}^2 + V(\phi) \approx V(\phi) \approx \text{const.} \propto a^{\approx 0}$$



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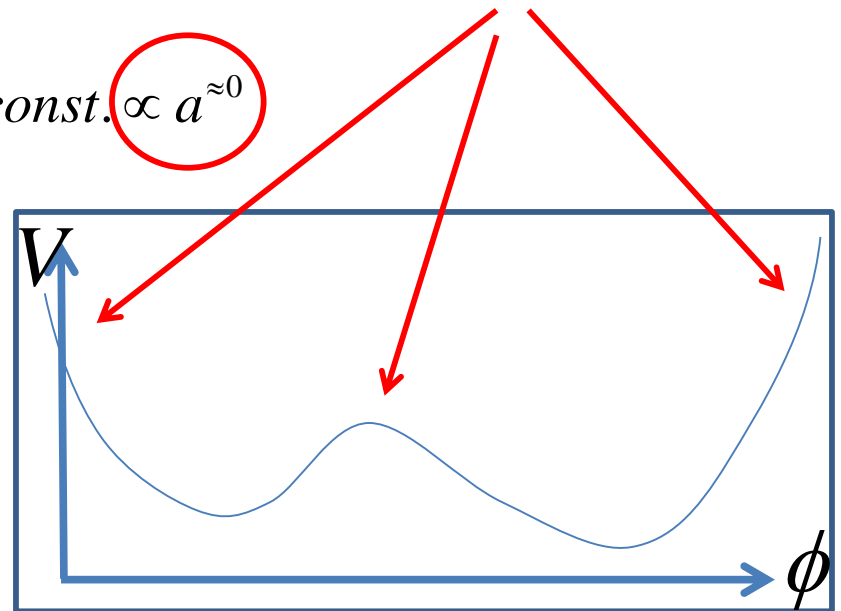
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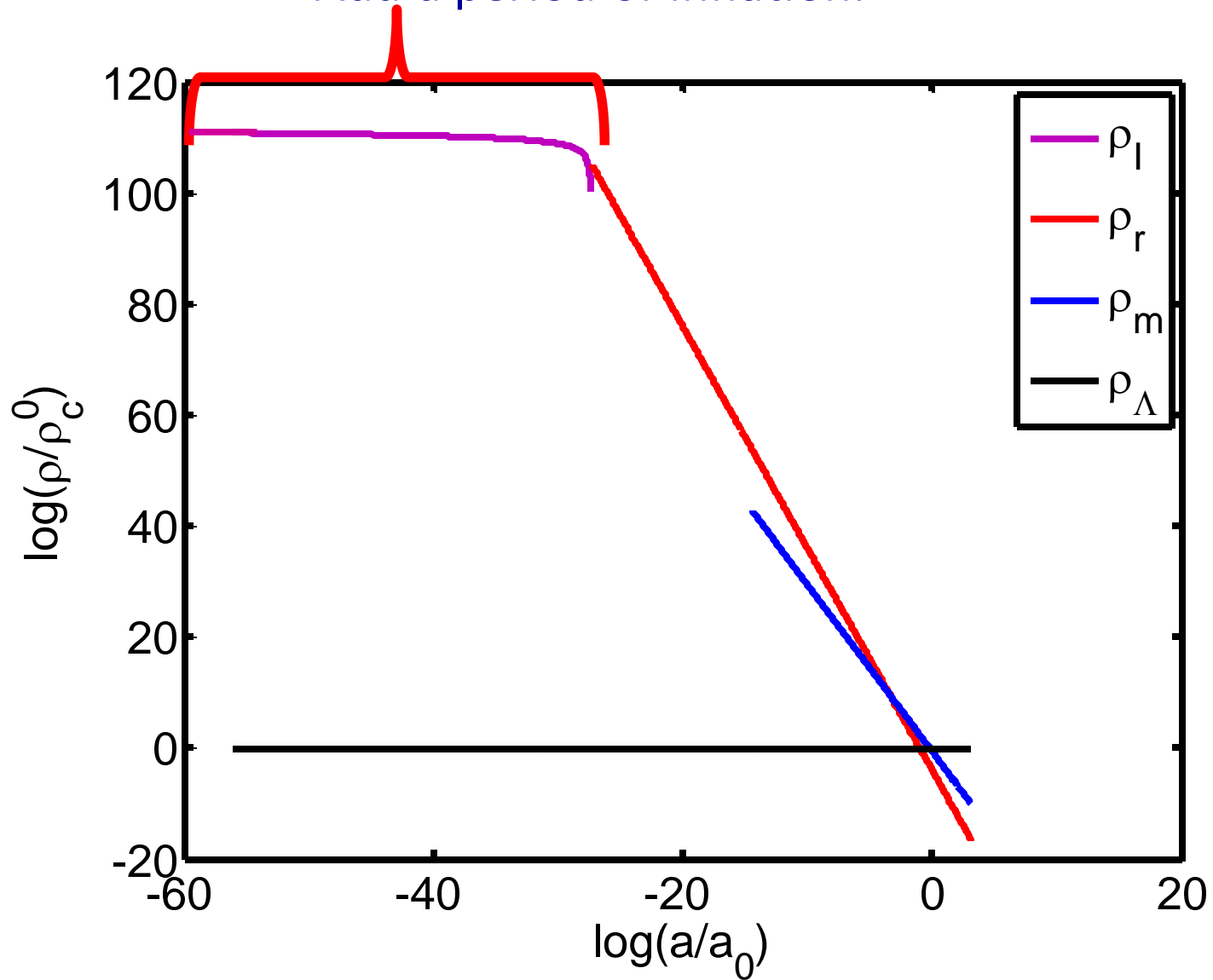
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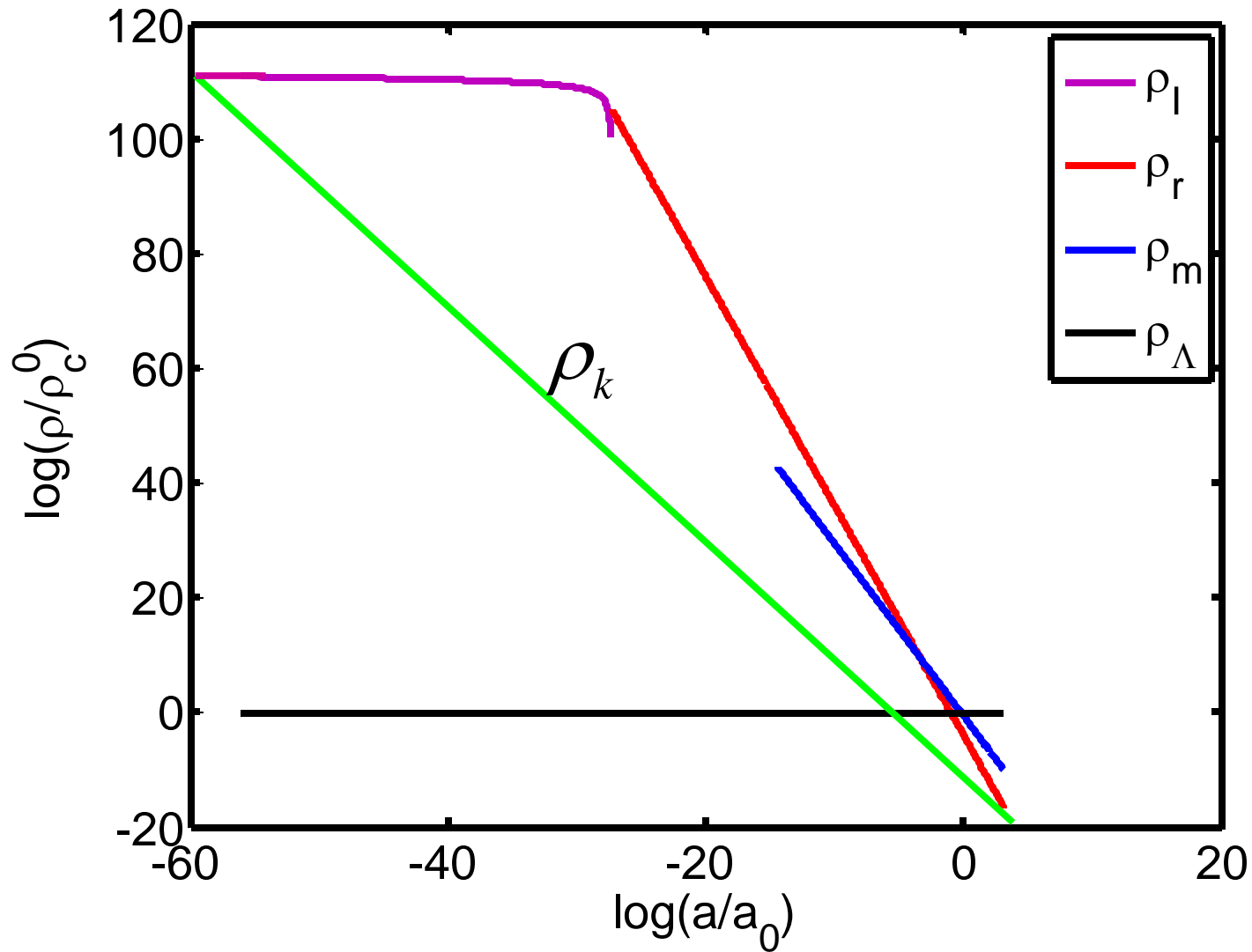
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Add a period of Inflation:

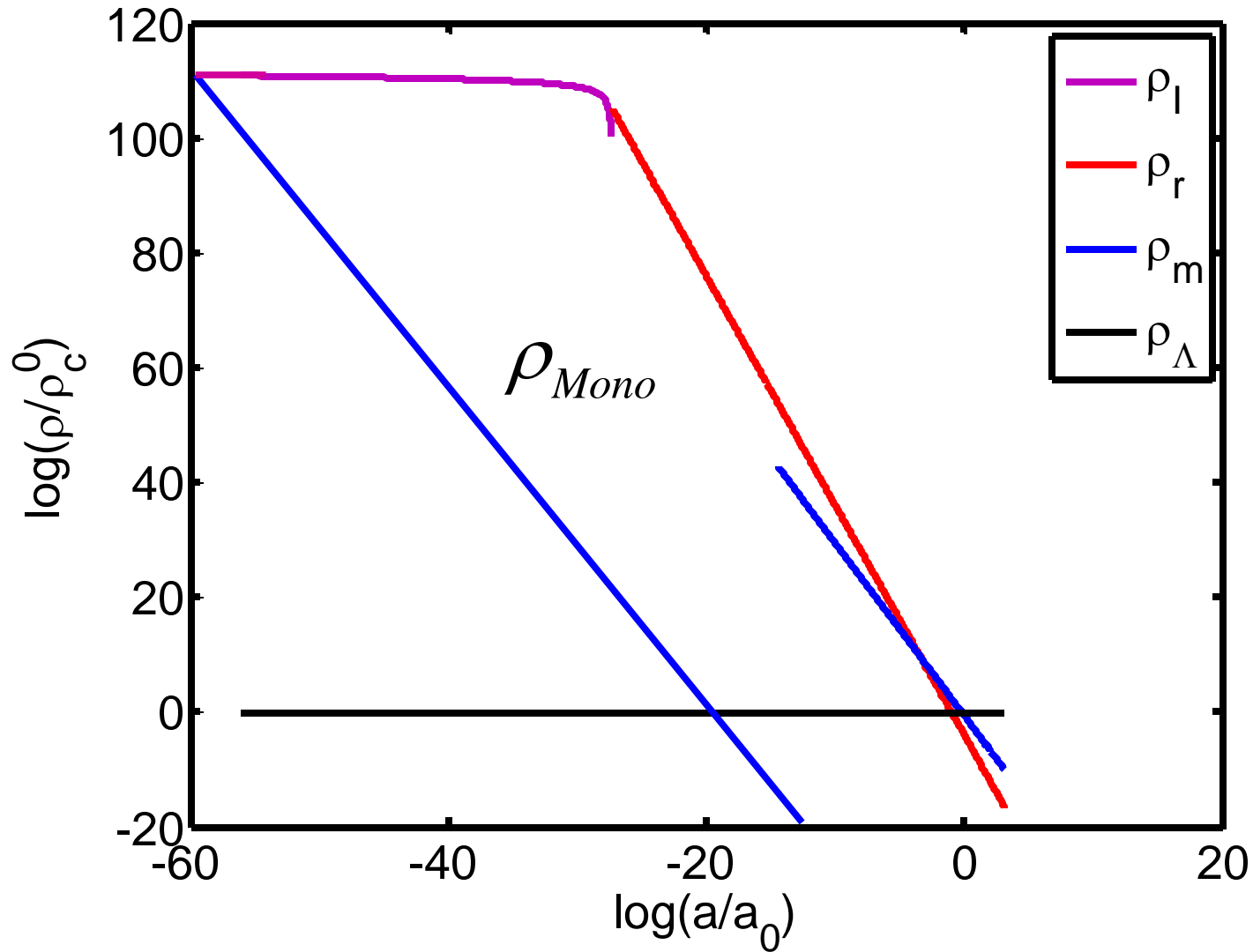


With inflation, initially large curvature is OK:

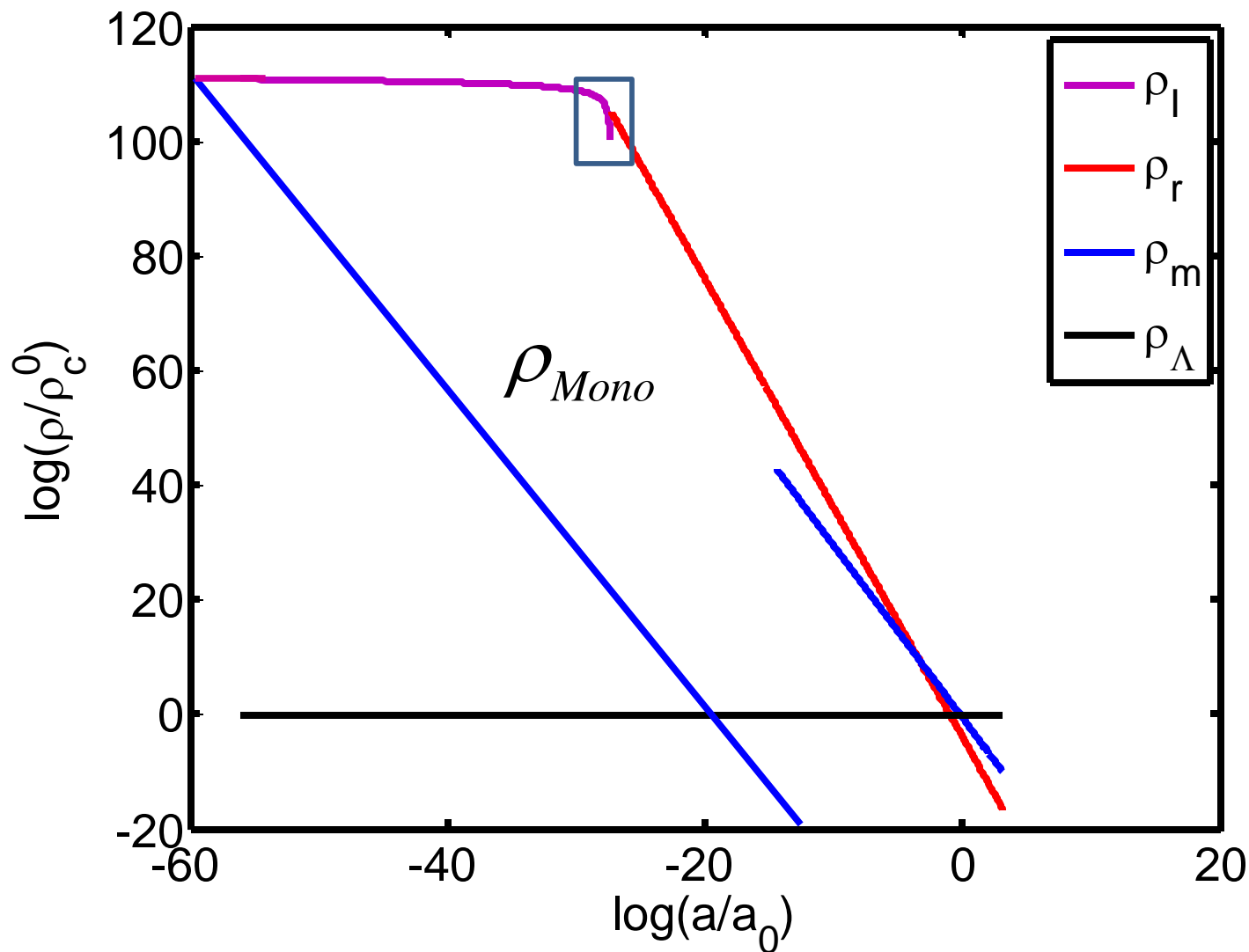




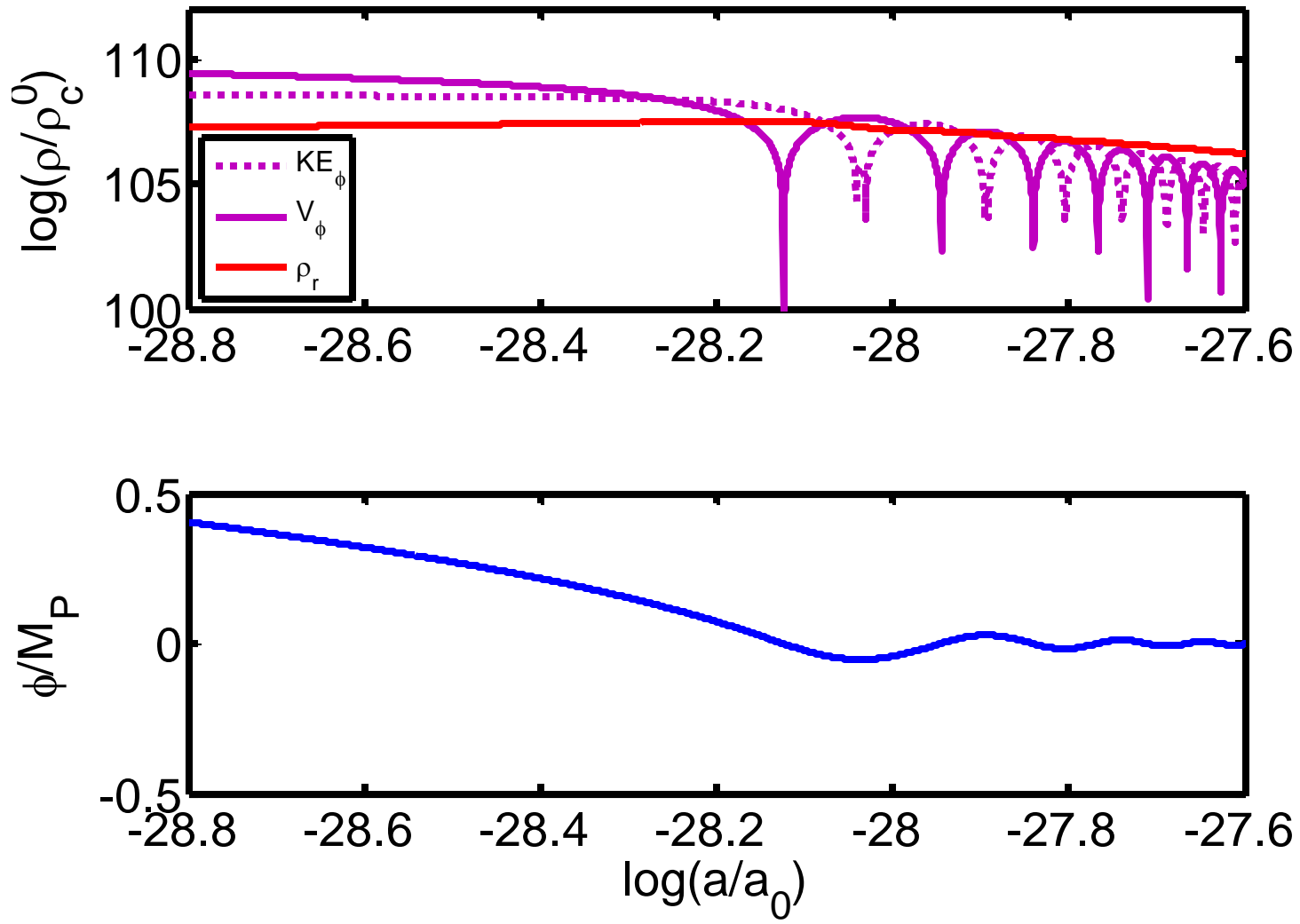
With inflation, early production of large amounts of non-relativistic matter (monopoles) is ok :



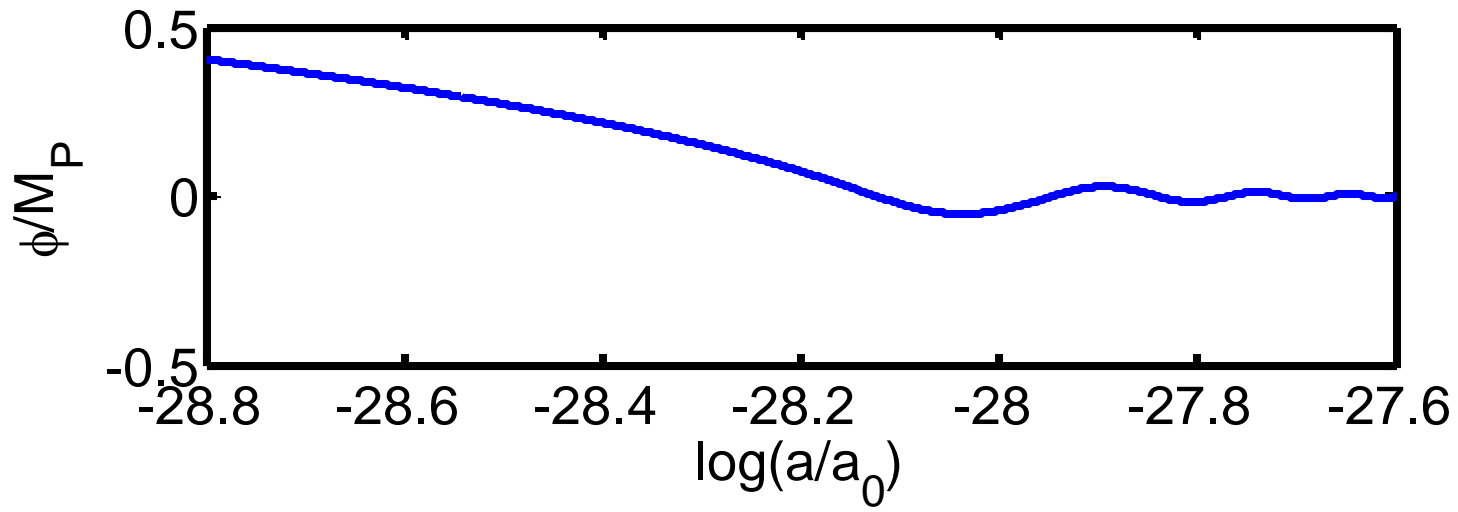
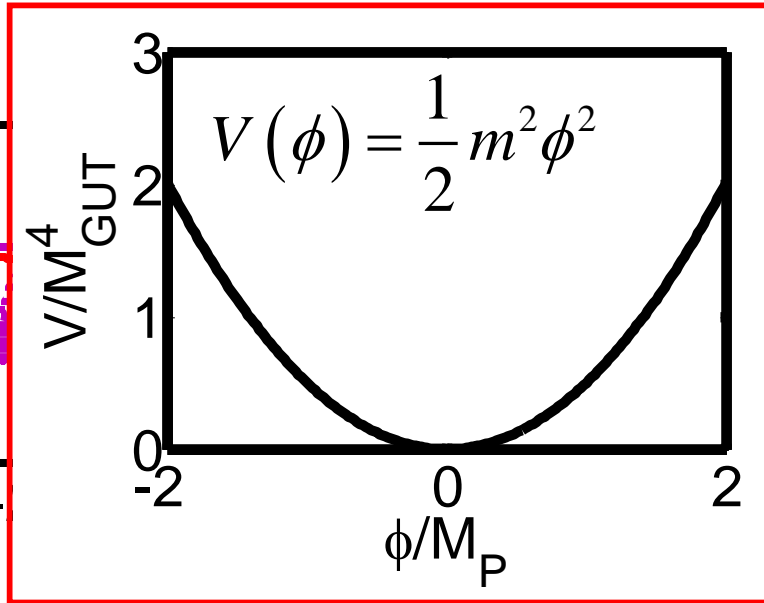
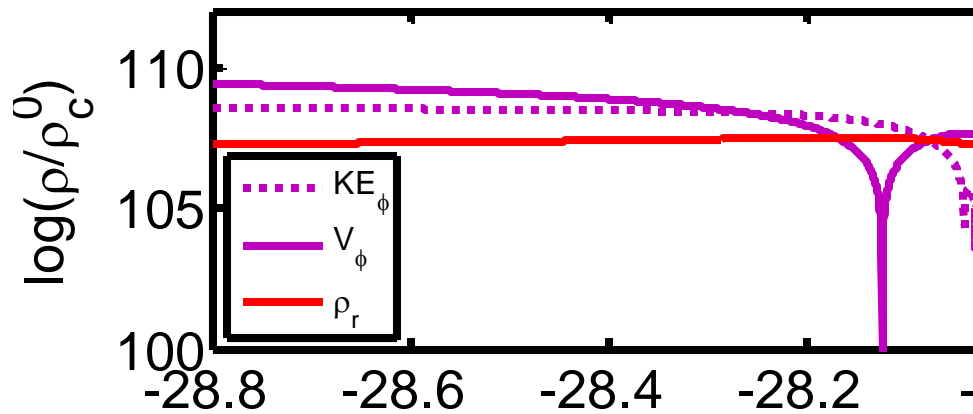
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## Inflation detail:

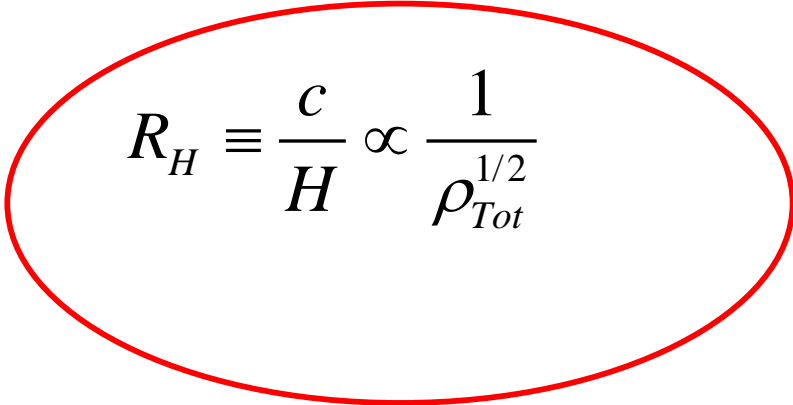


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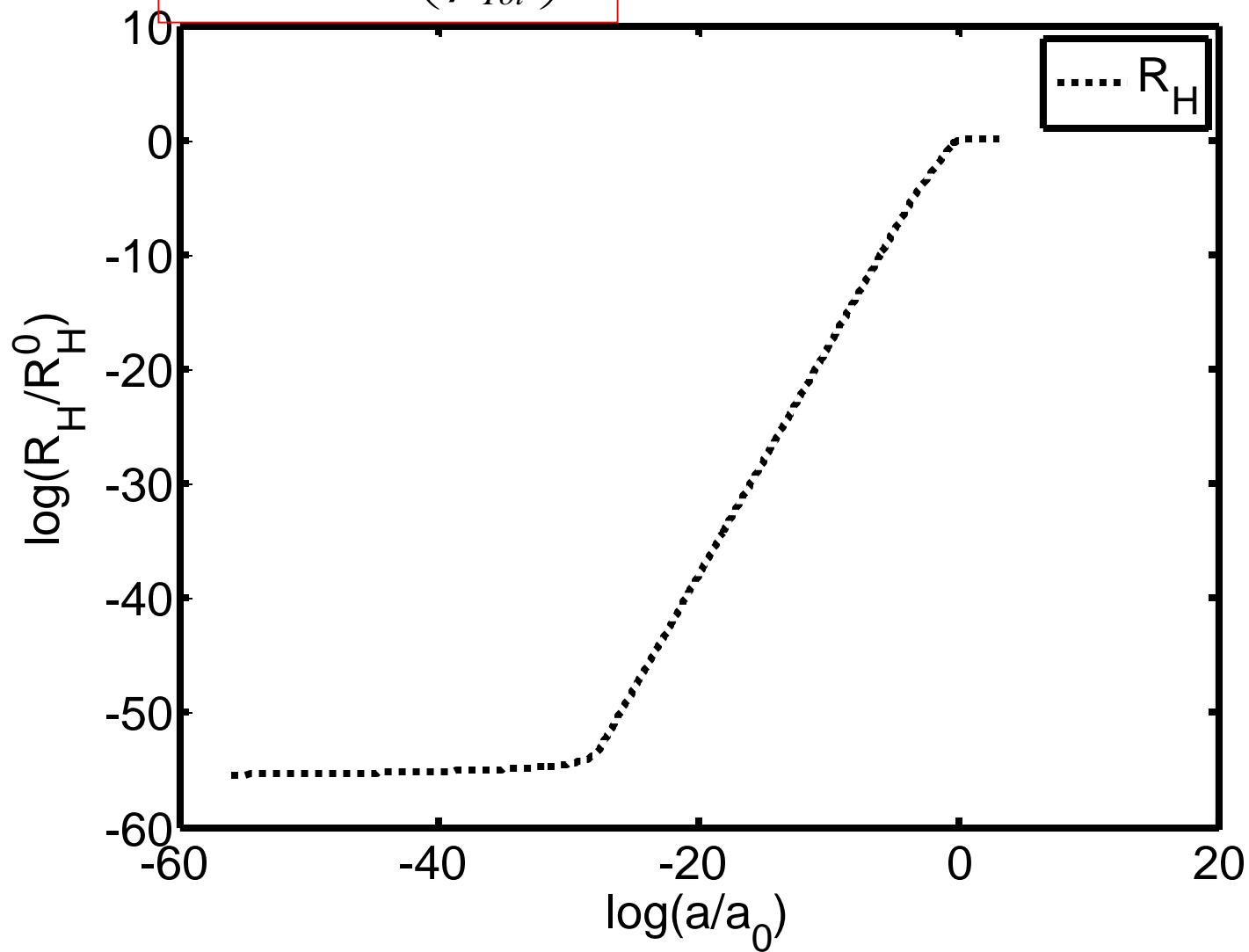


# Hubble Length

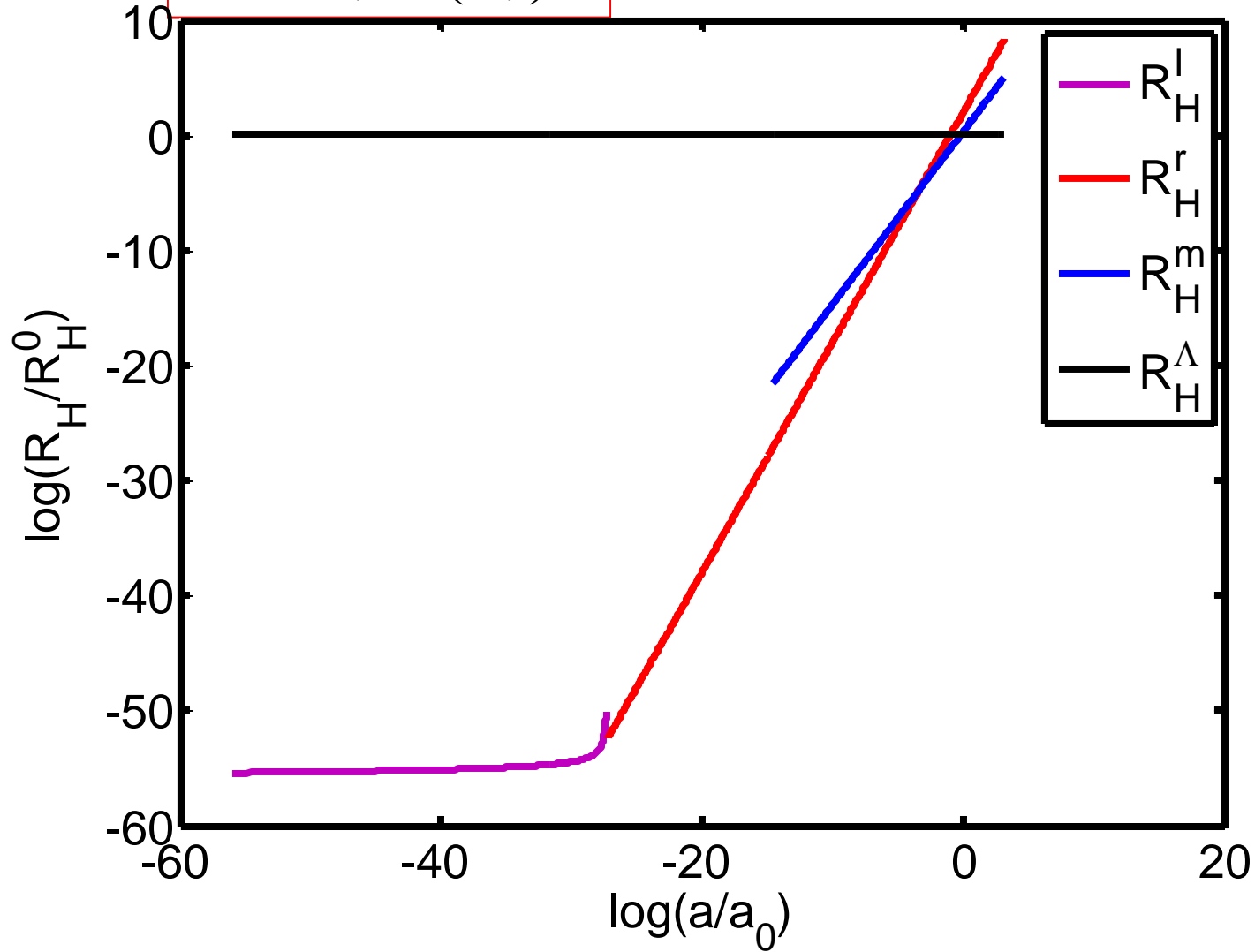
$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_I + \rho_k + \rho_r + \rho_m + \rho_{DE}) \equiv \frac{8\pi}{3} G \rho_{Tot}$$


$$R_H \equiv \frac{c}{H} \propto \frac{1}{\rho_{Tot}^{1/2}}$$

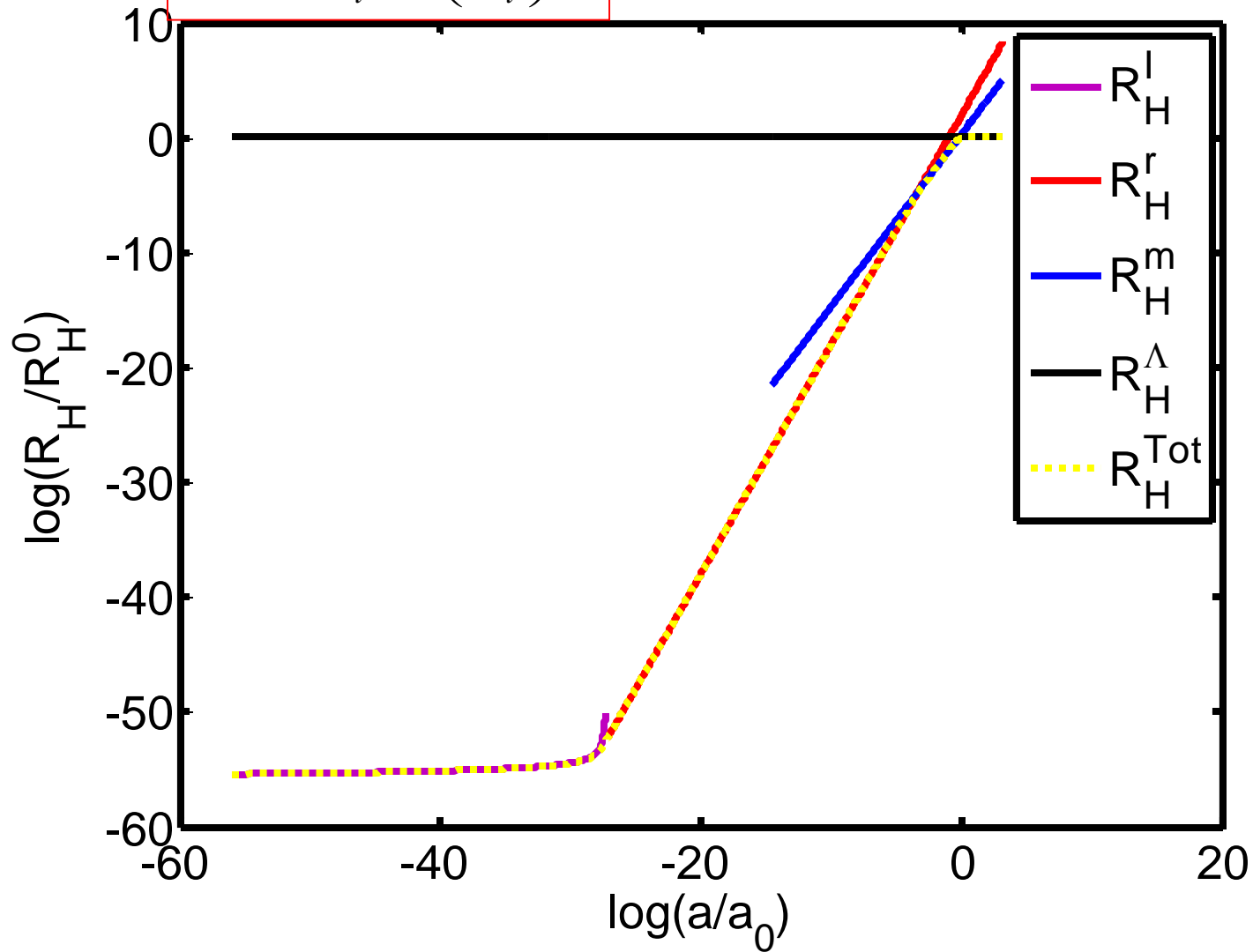
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$$R_H^i \equiv \frac{c}{H_i} \propto \left( \frac{1}{\rho_i} \right)^{1/2}$$

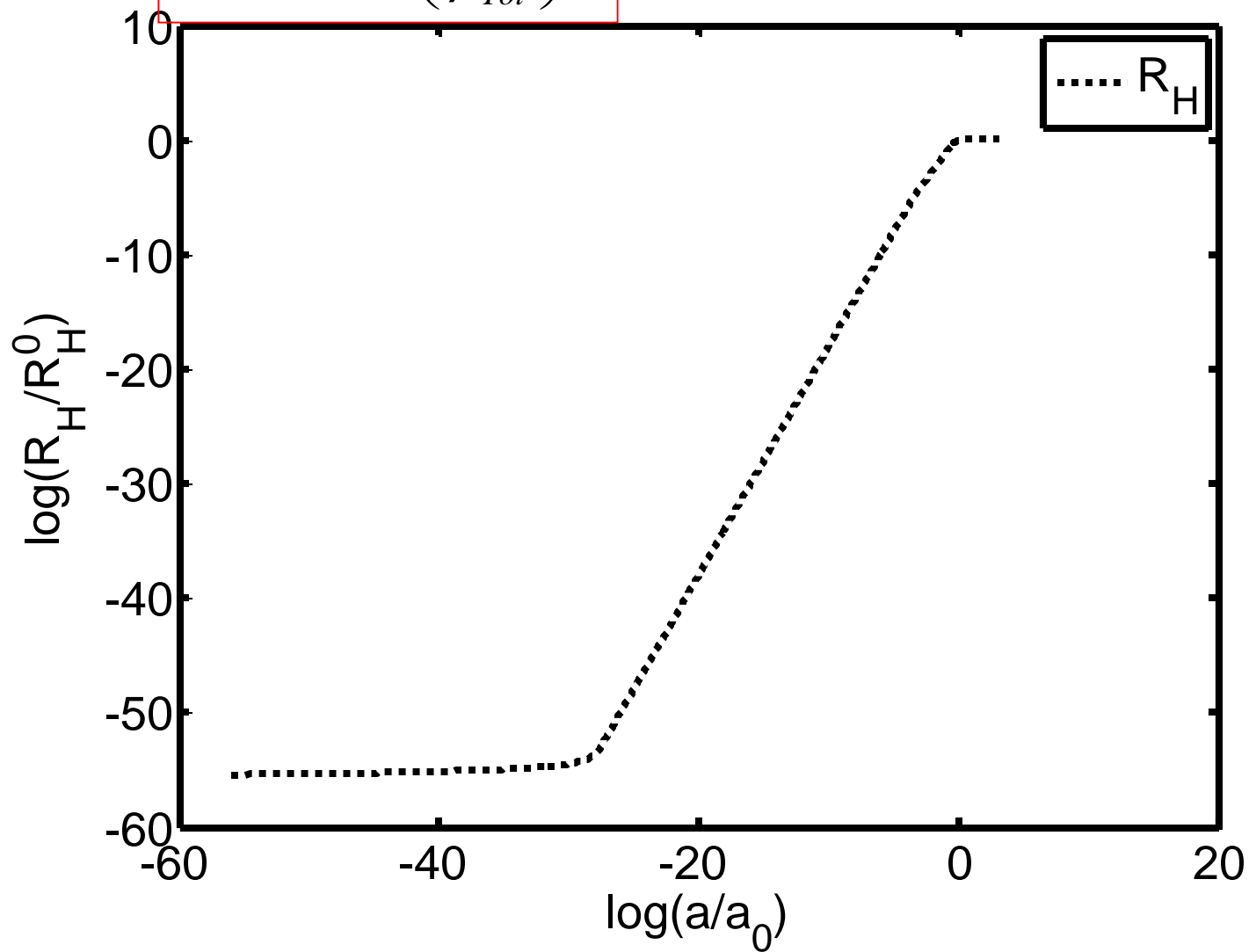


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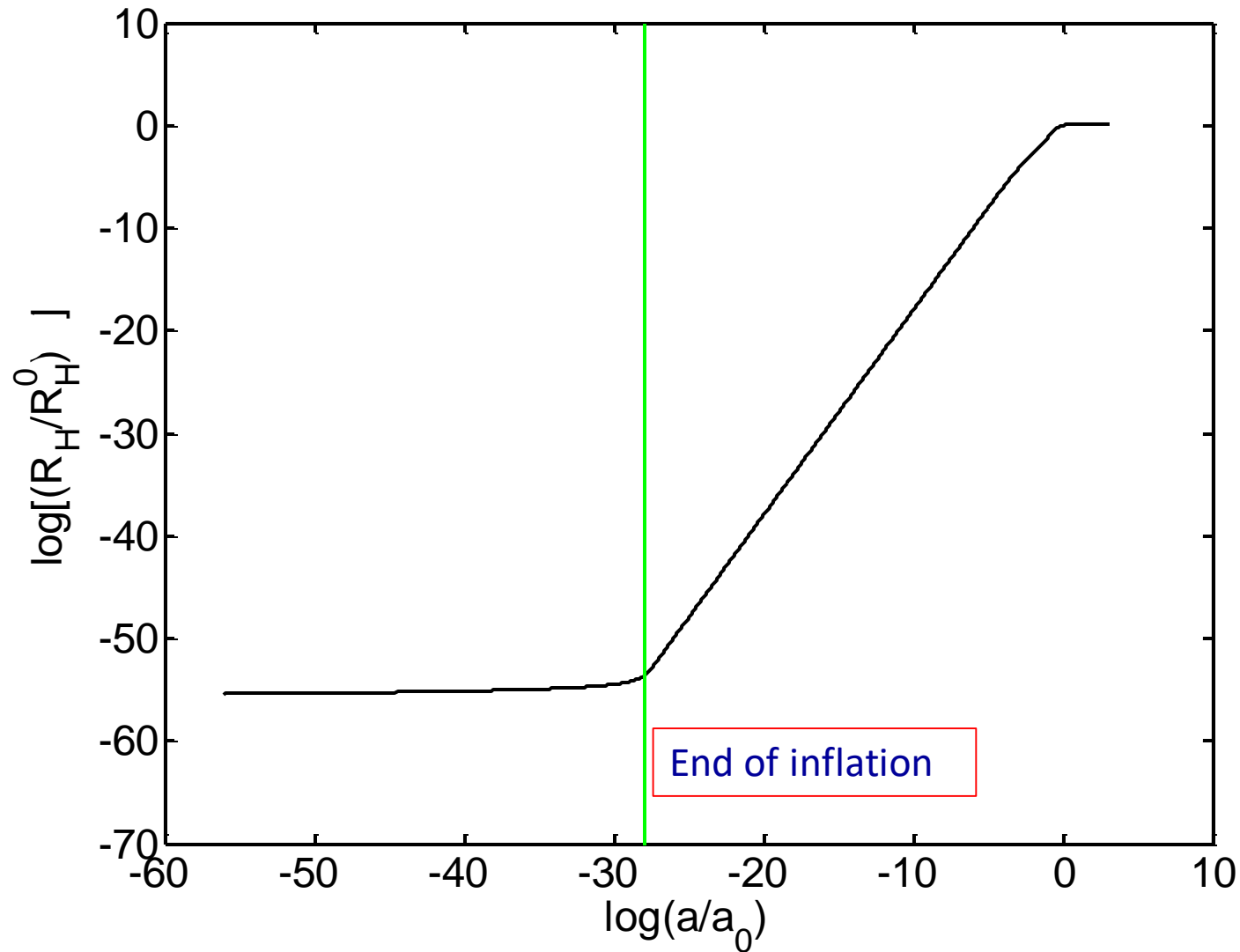




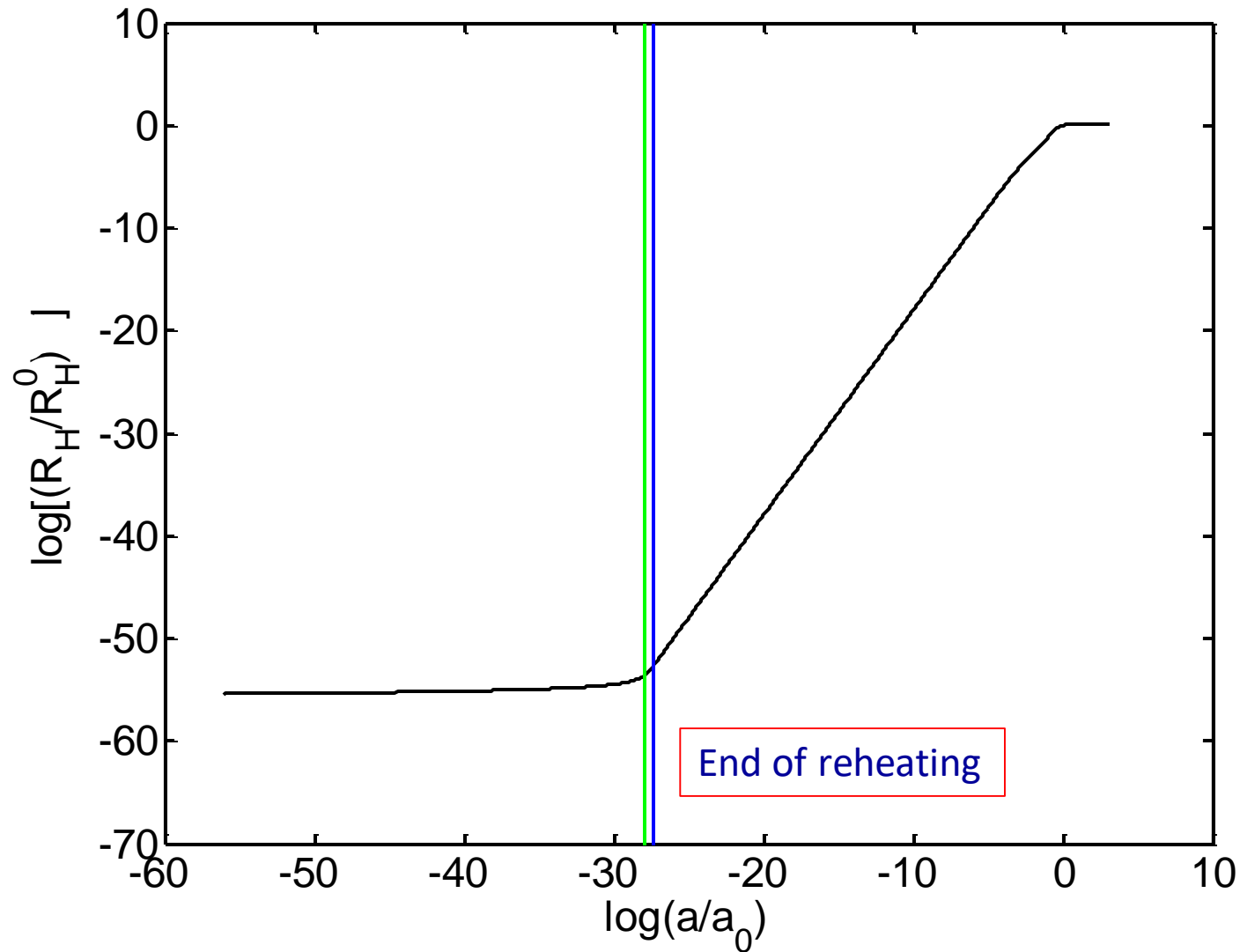
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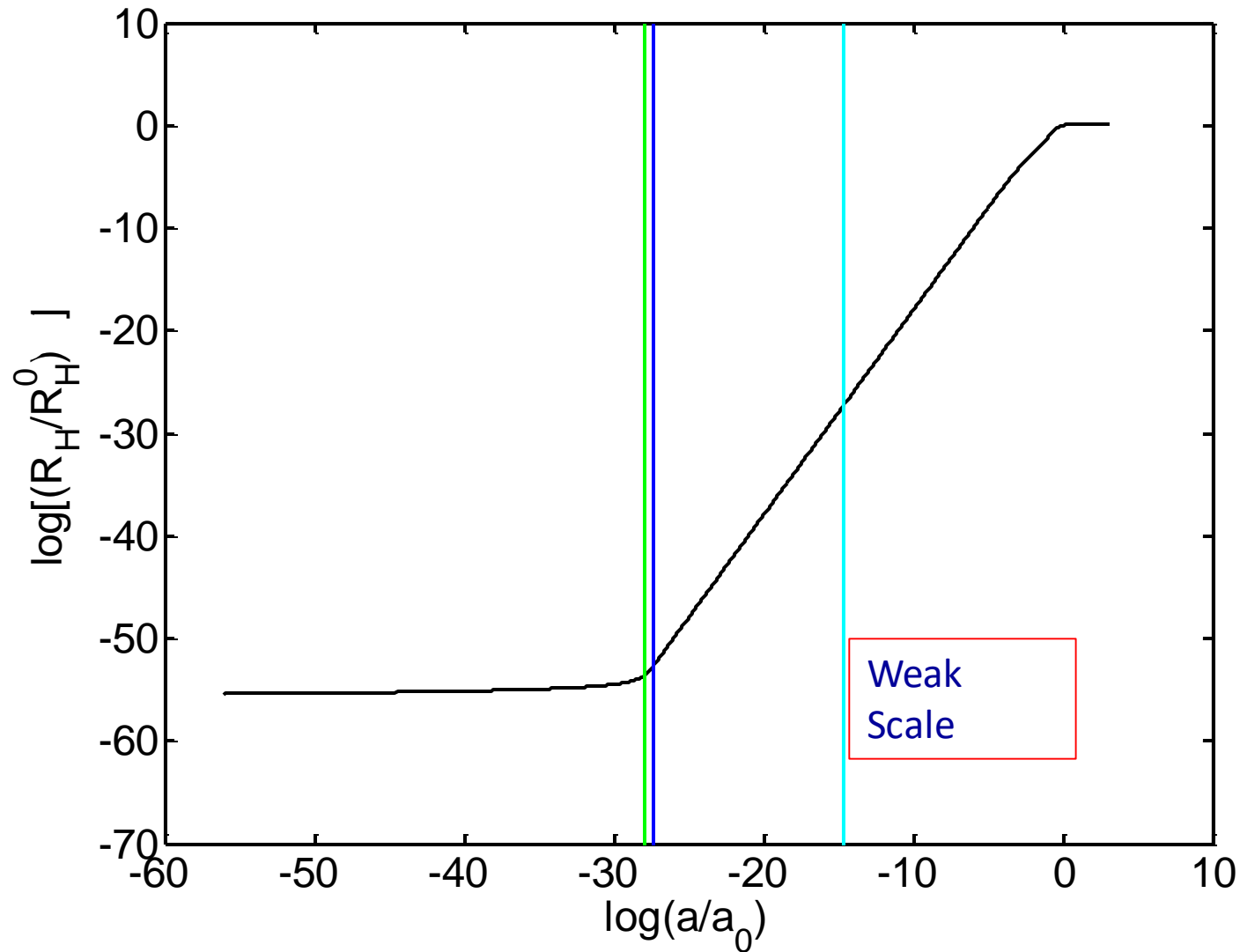
# Evolution of Cosmic Length



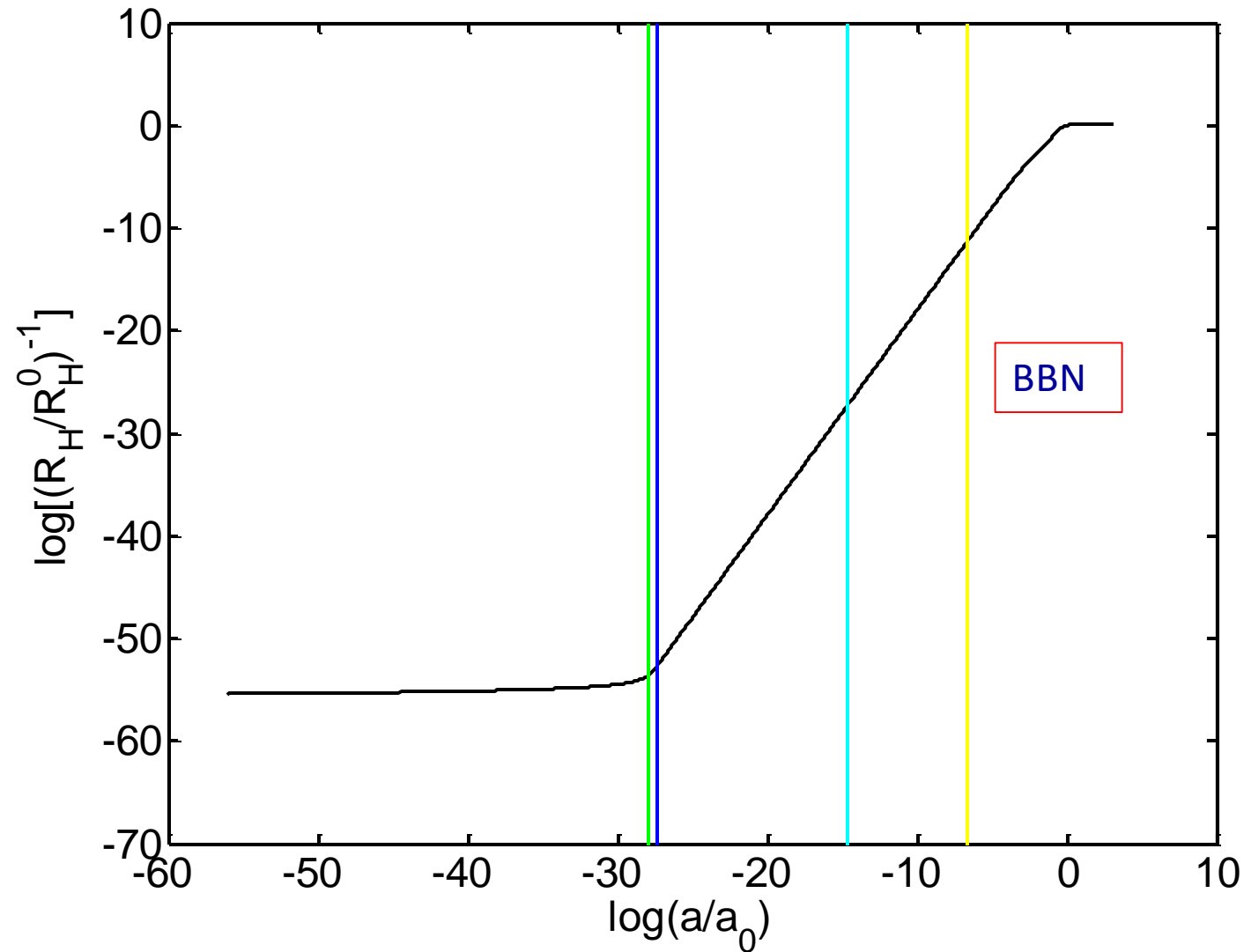
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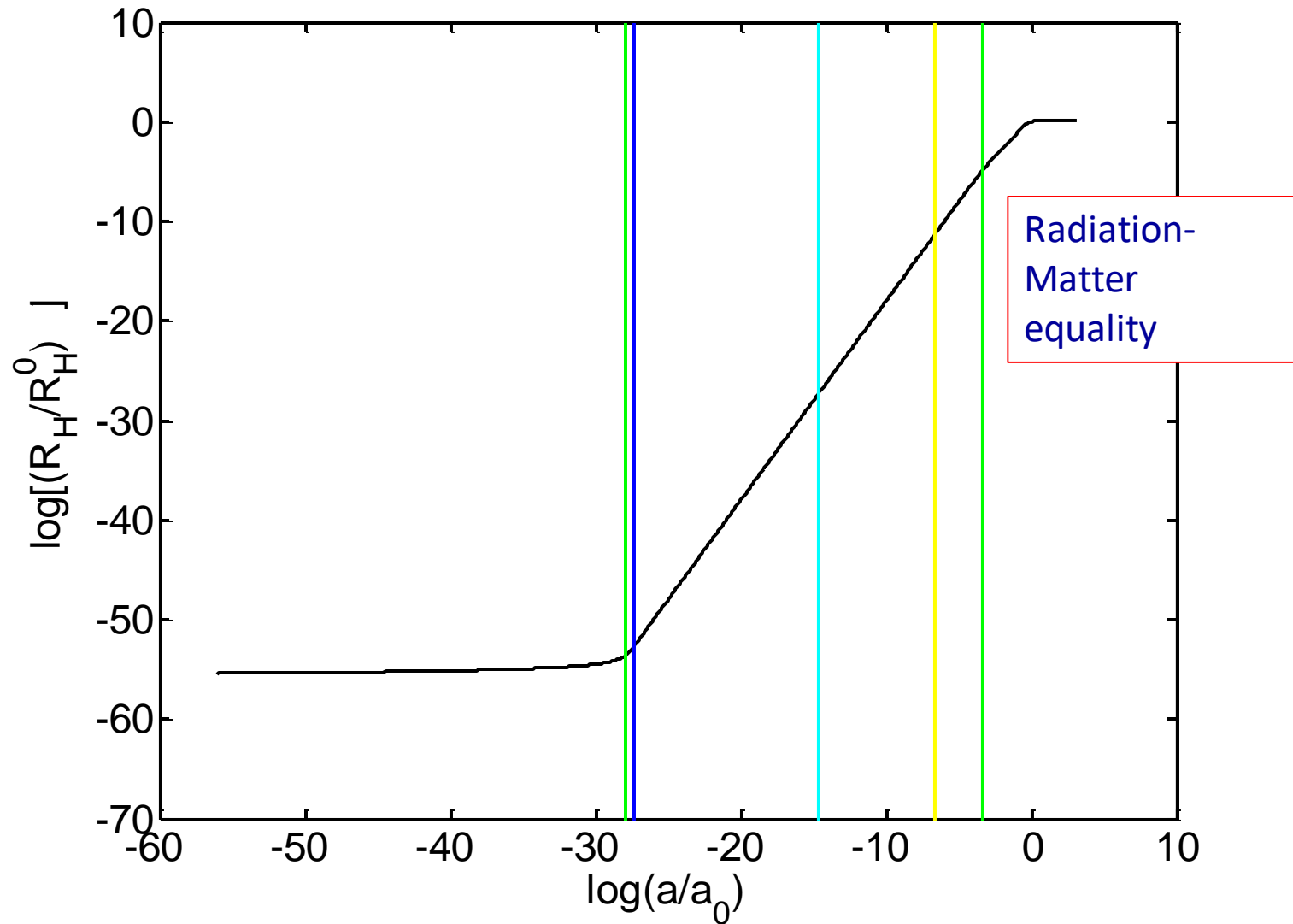
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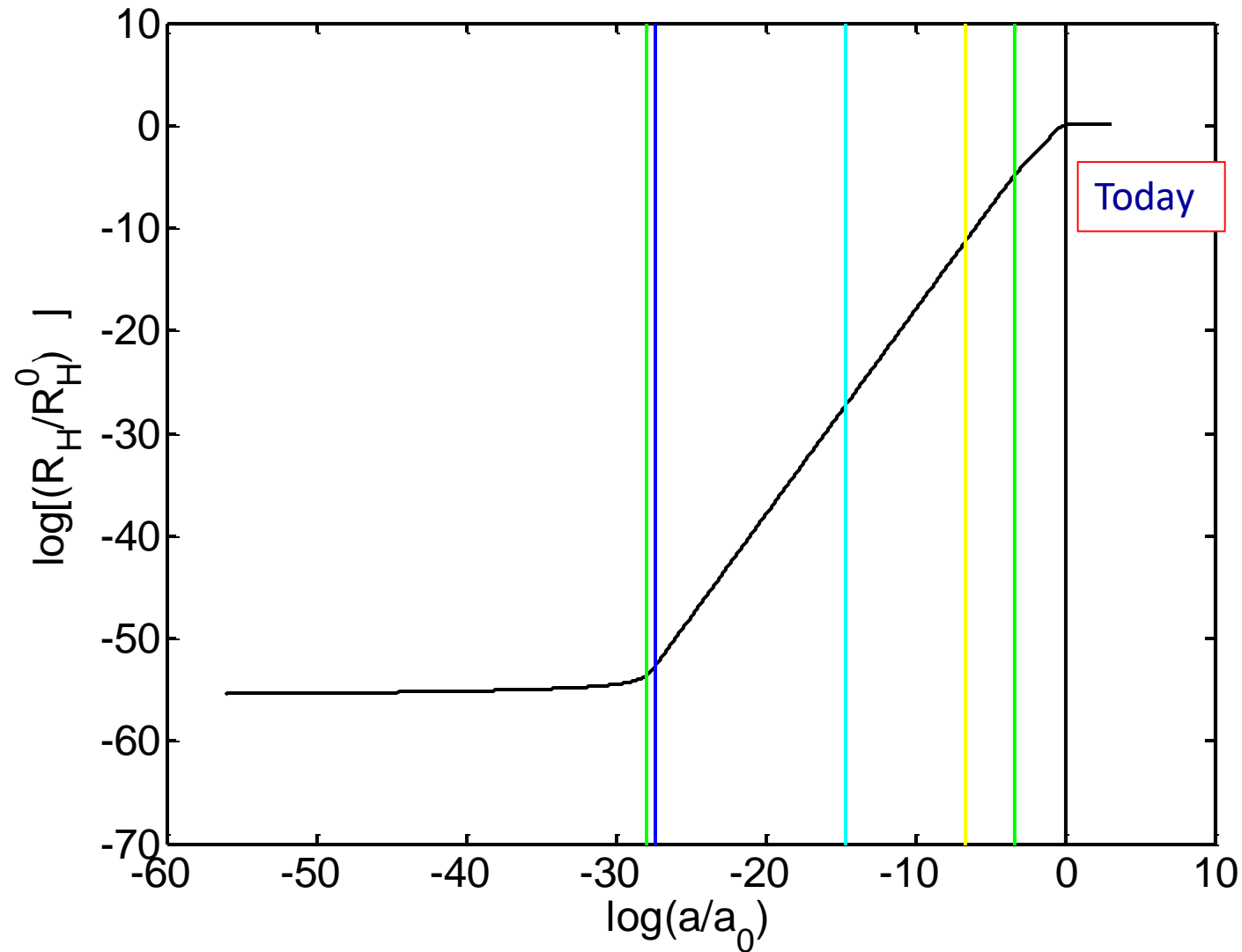
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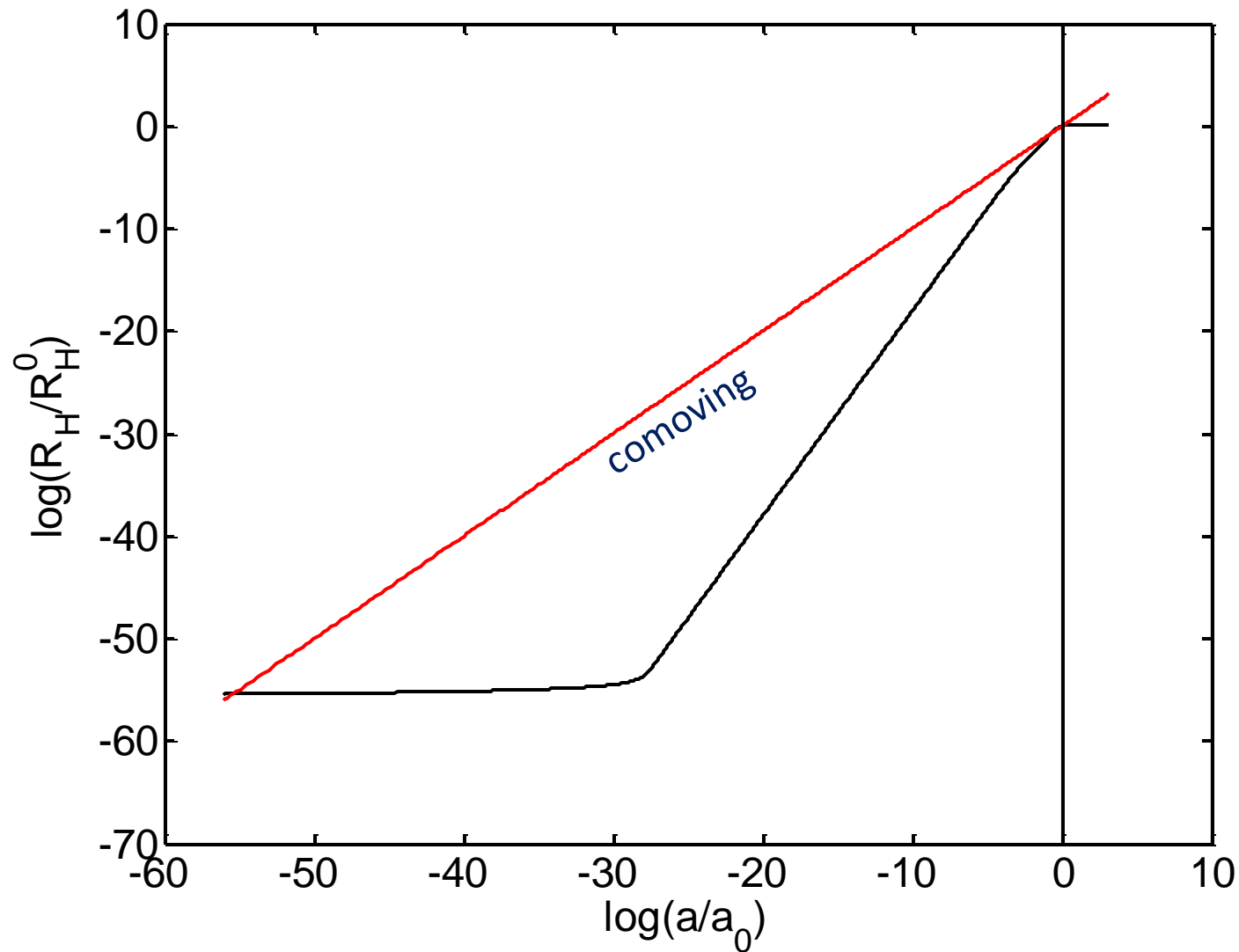
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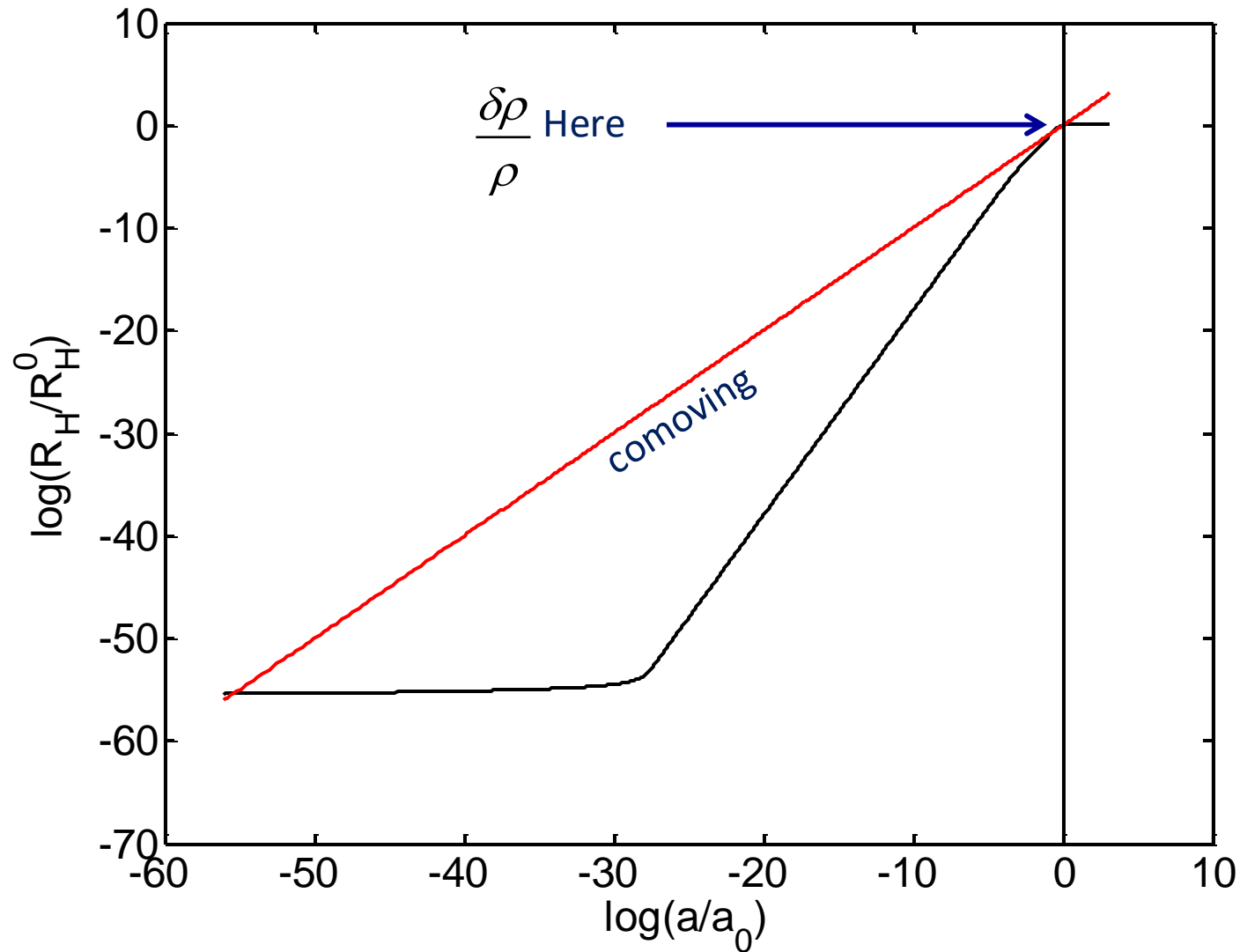


# Perturbations from inflation

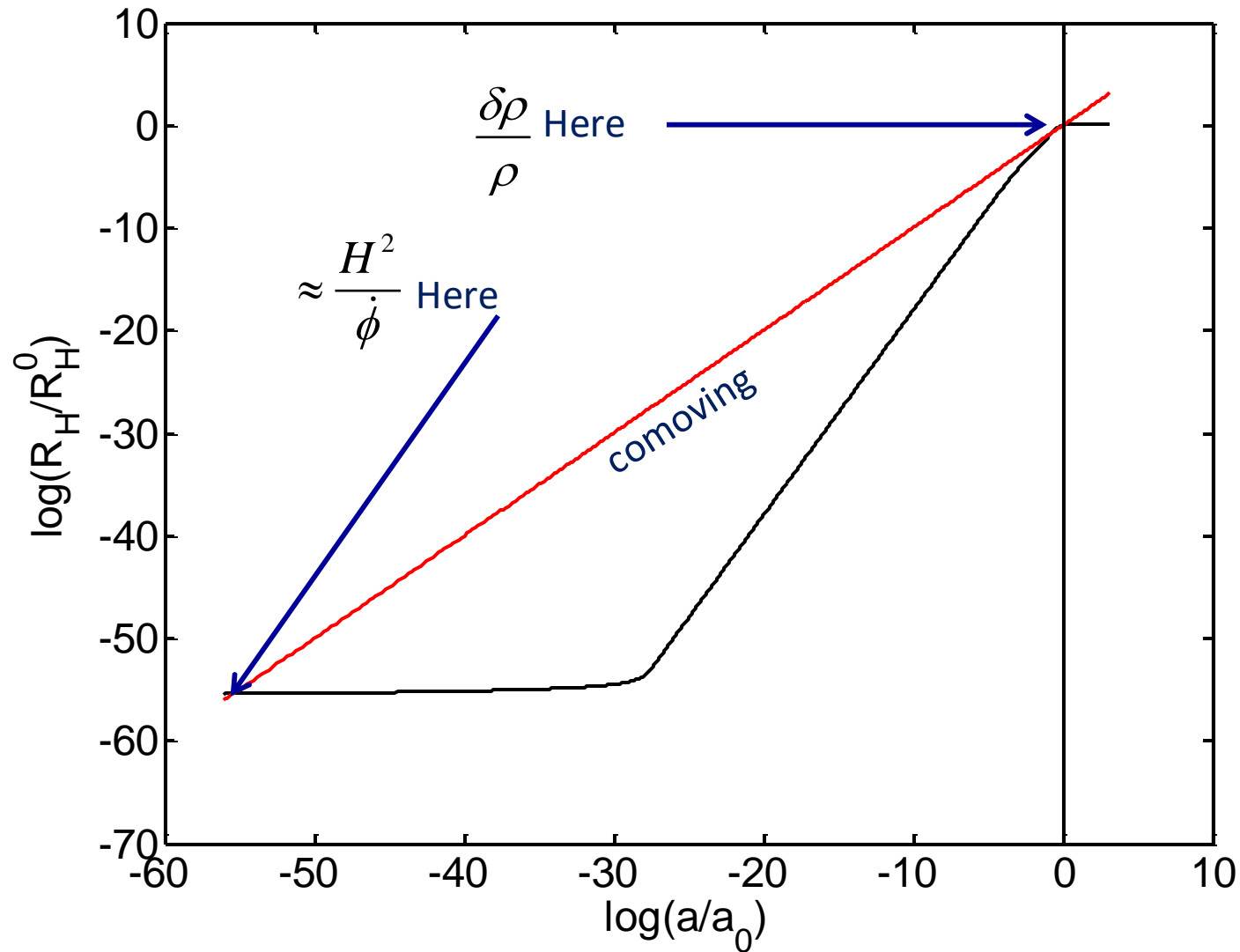


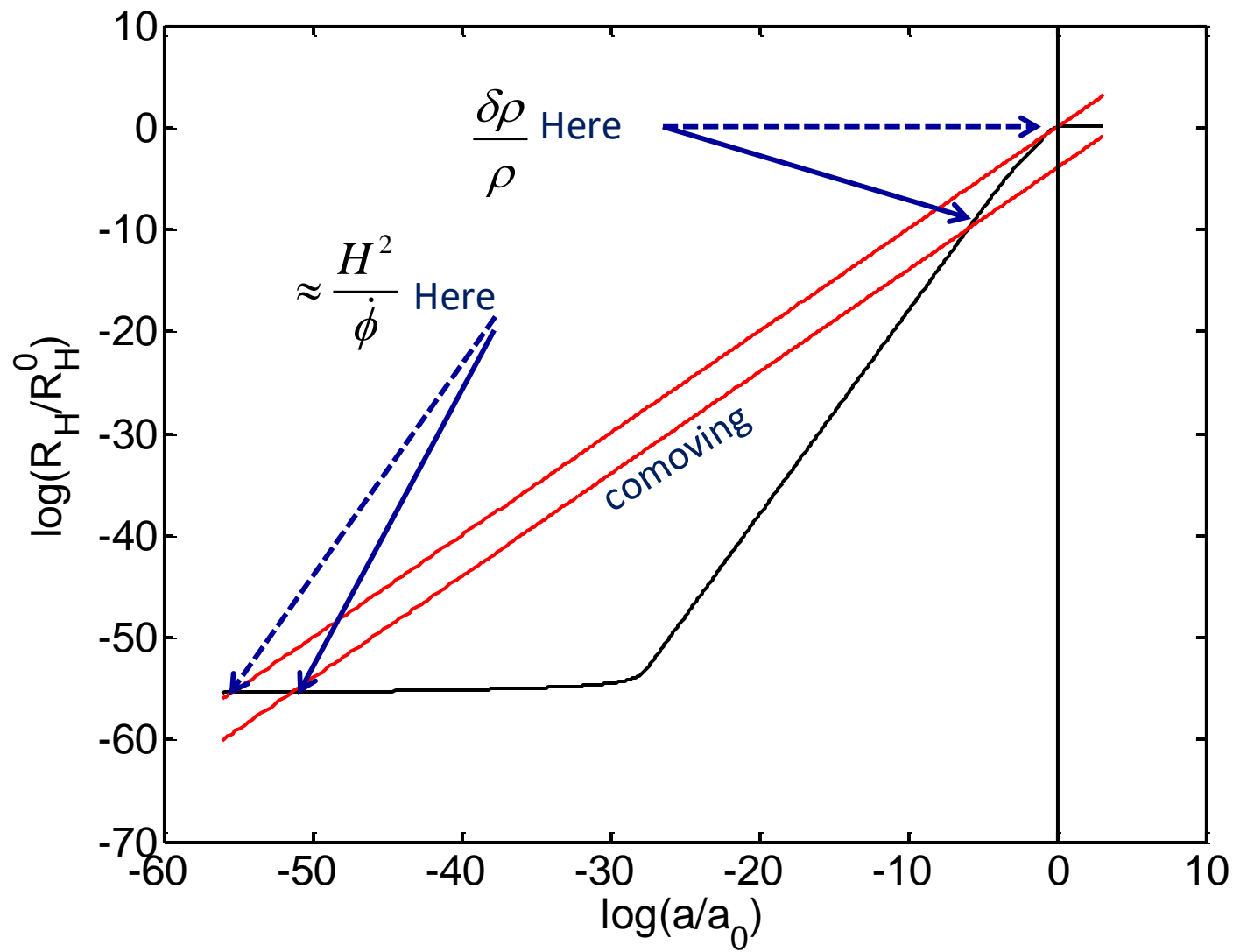


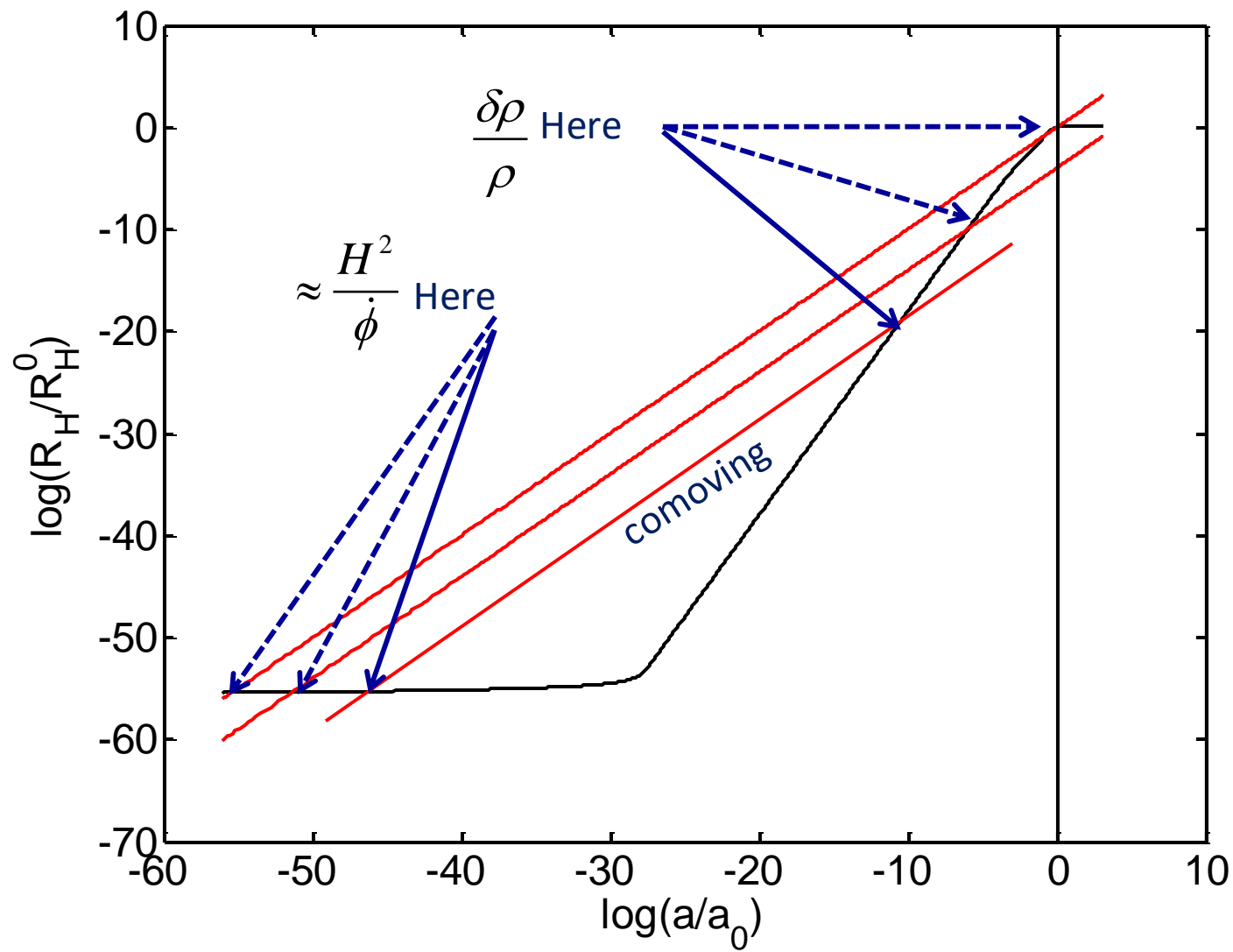
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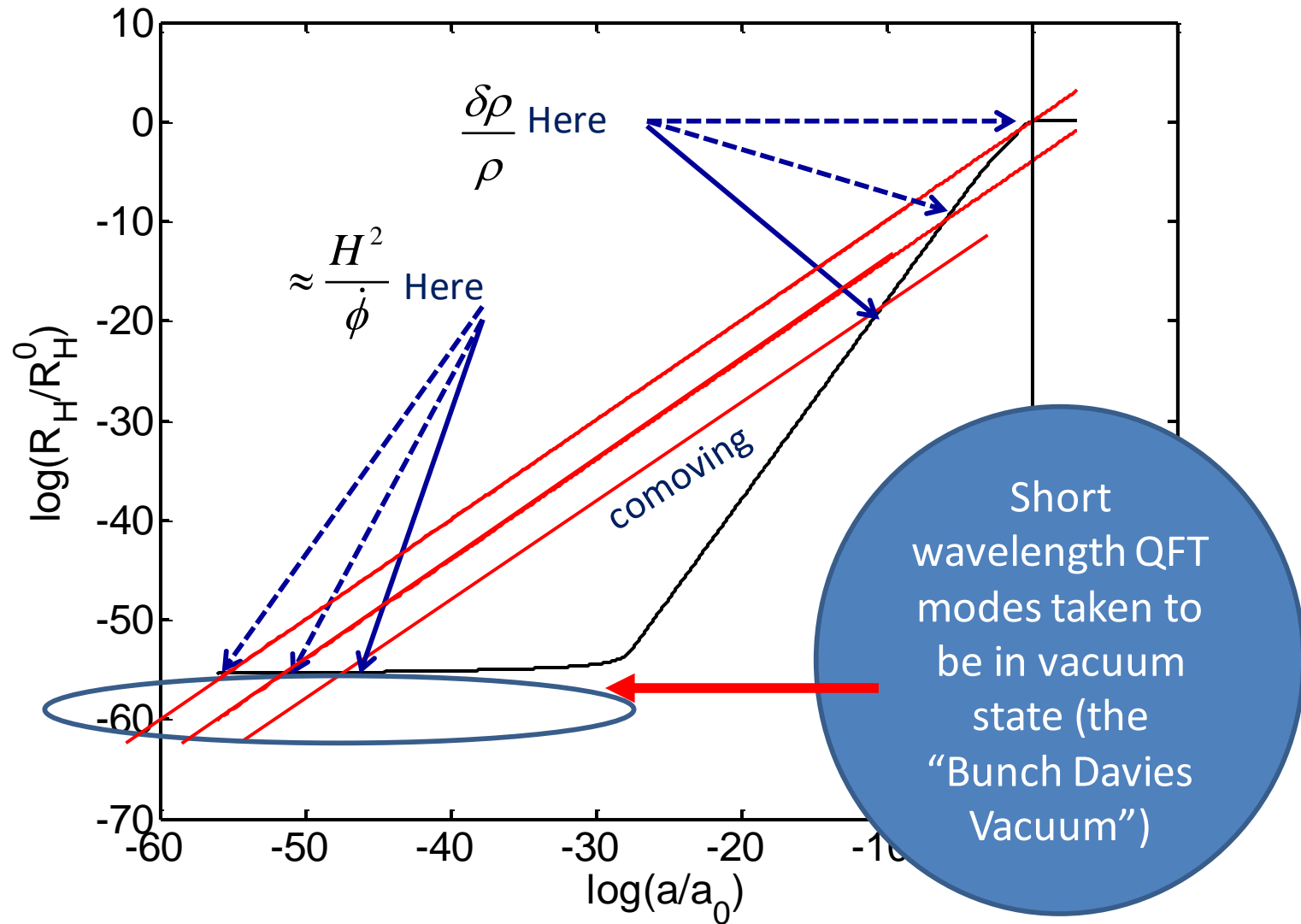


# Perturbations from inflation

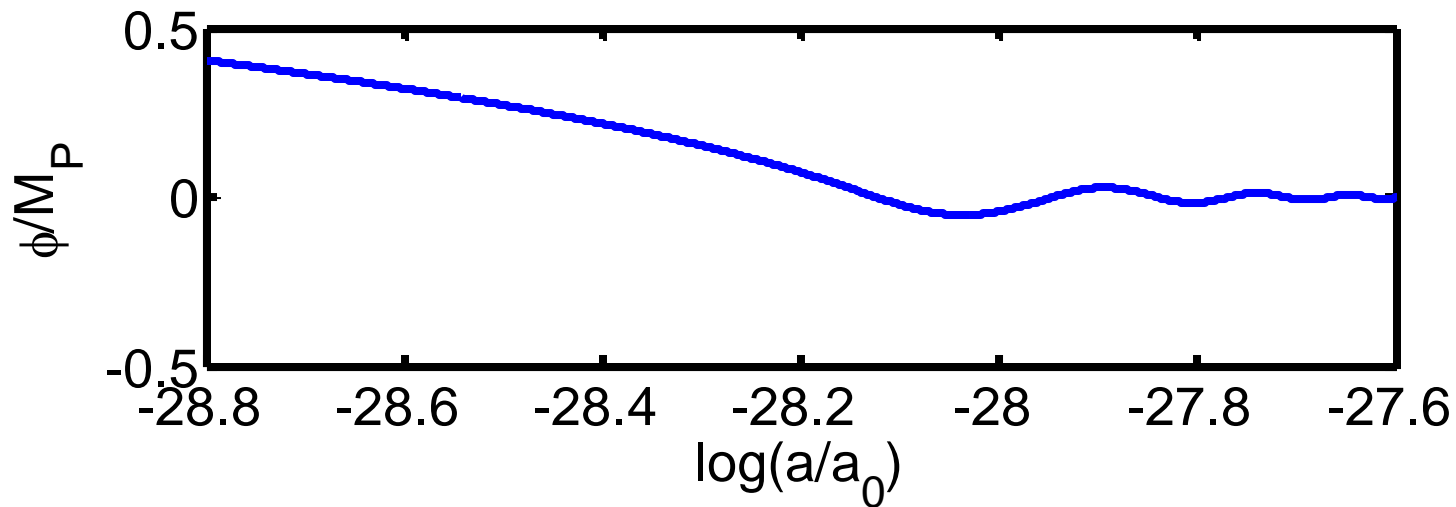
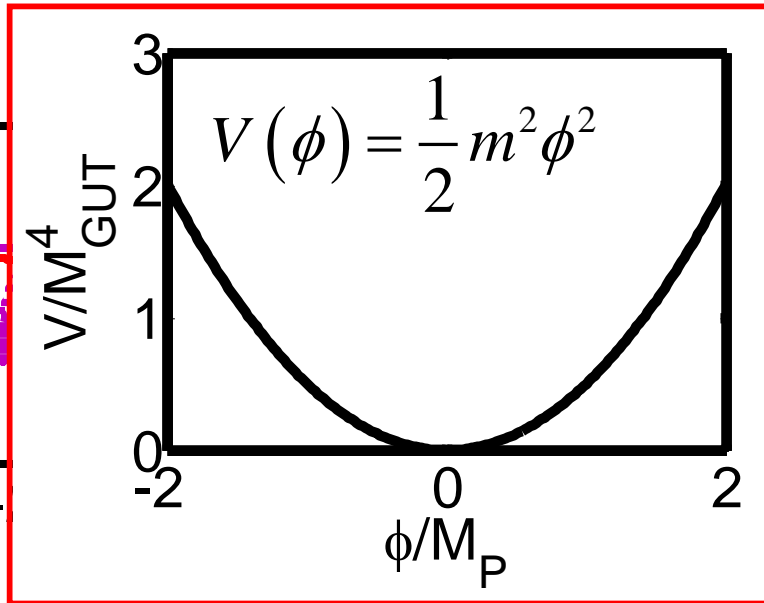
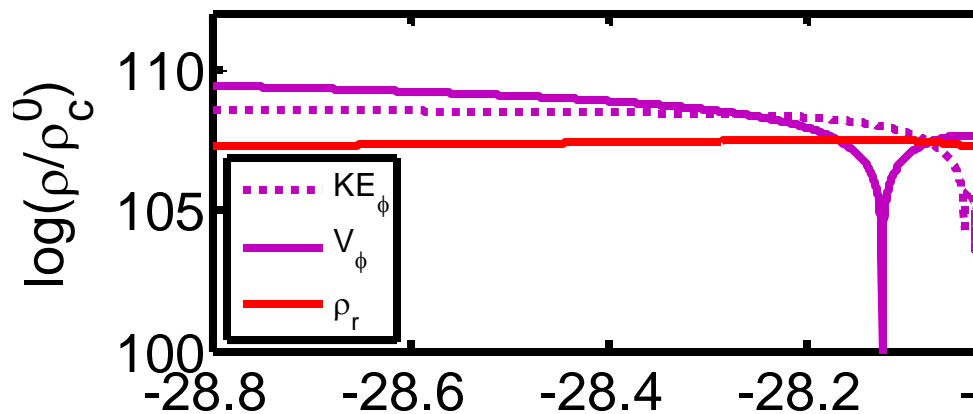




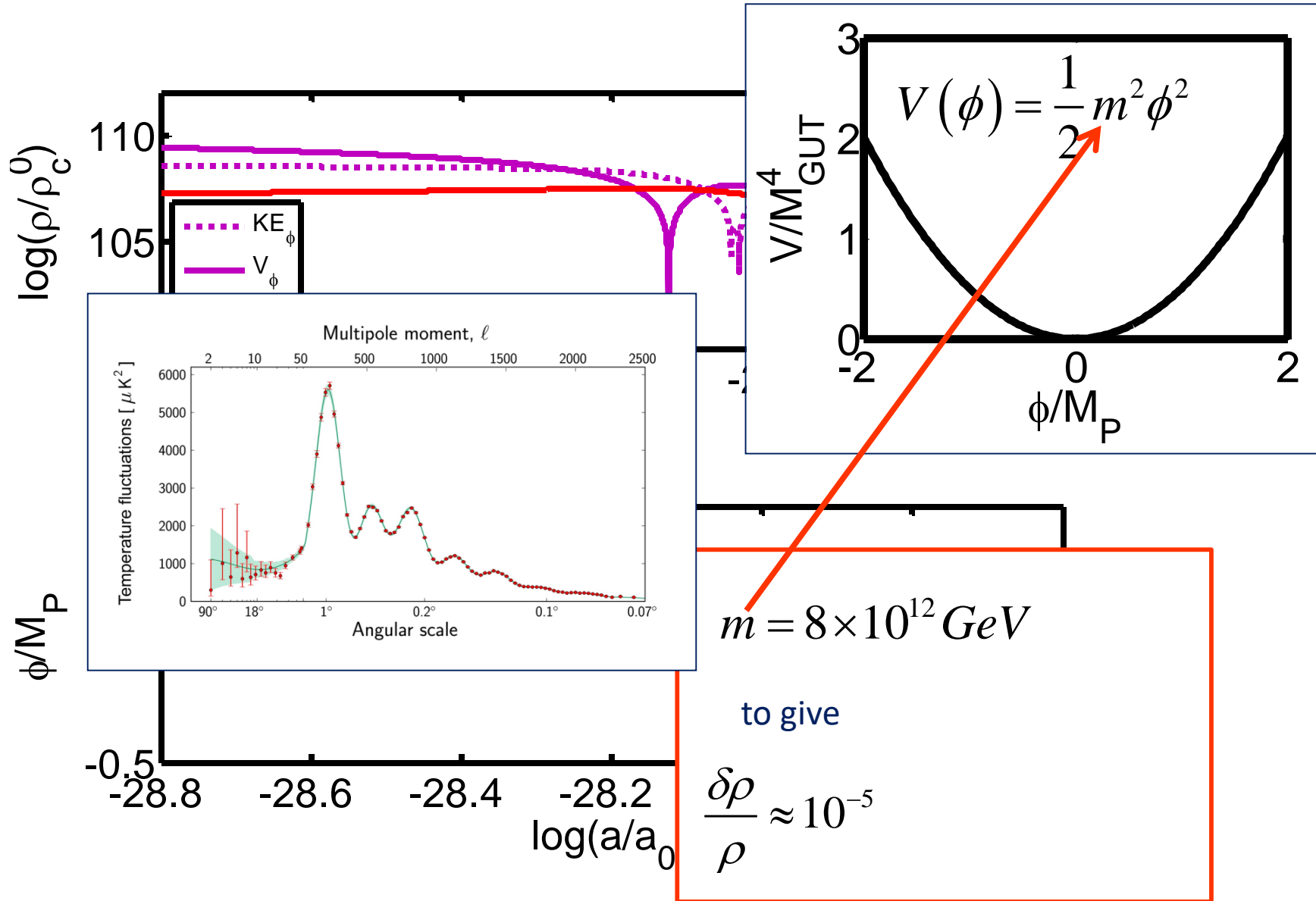




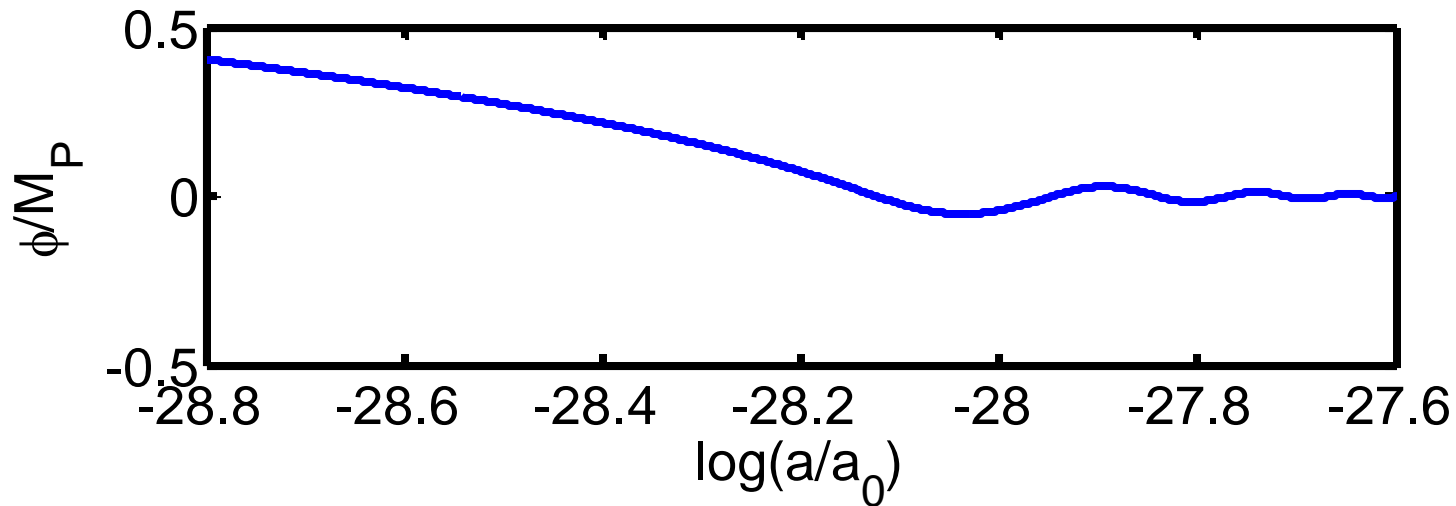
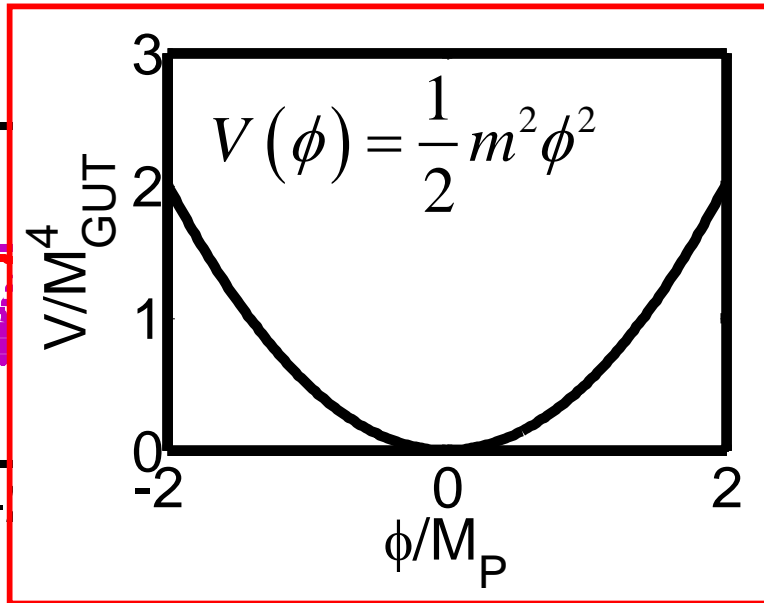
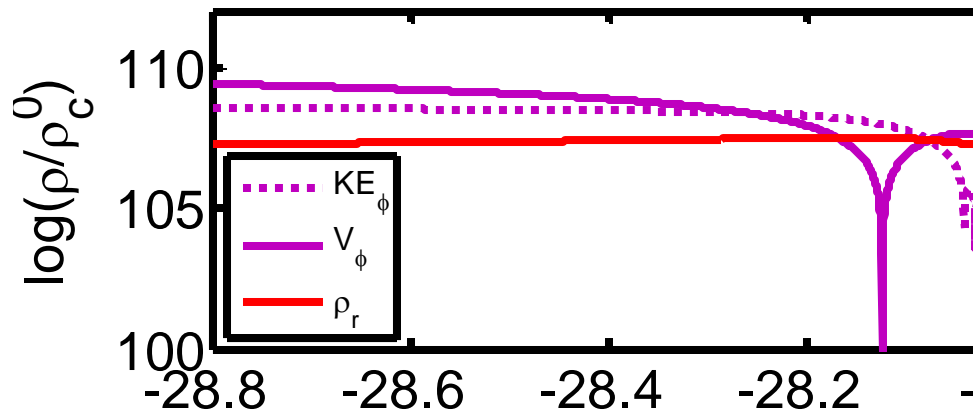
## Inflation detail:



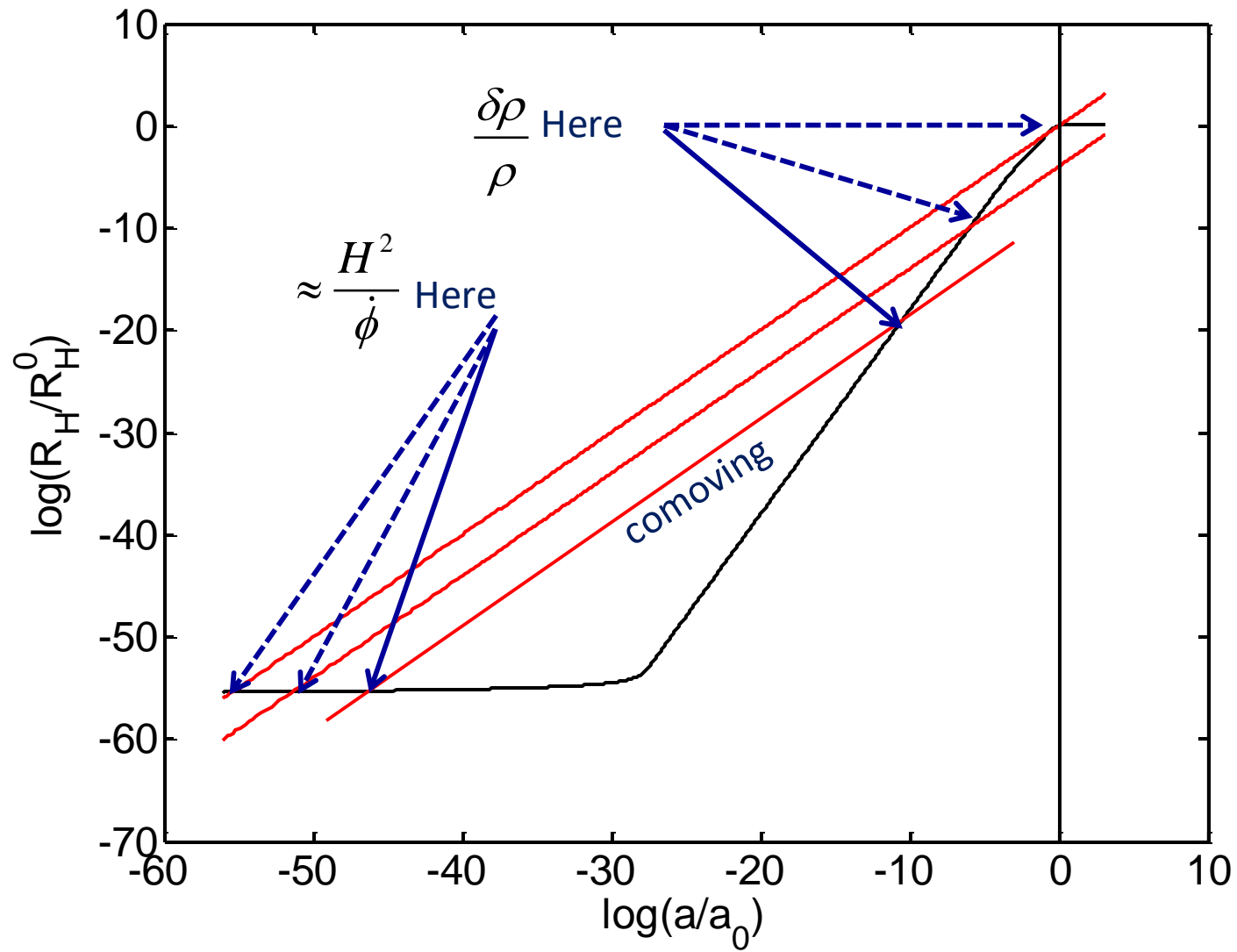
# Inflation detail:



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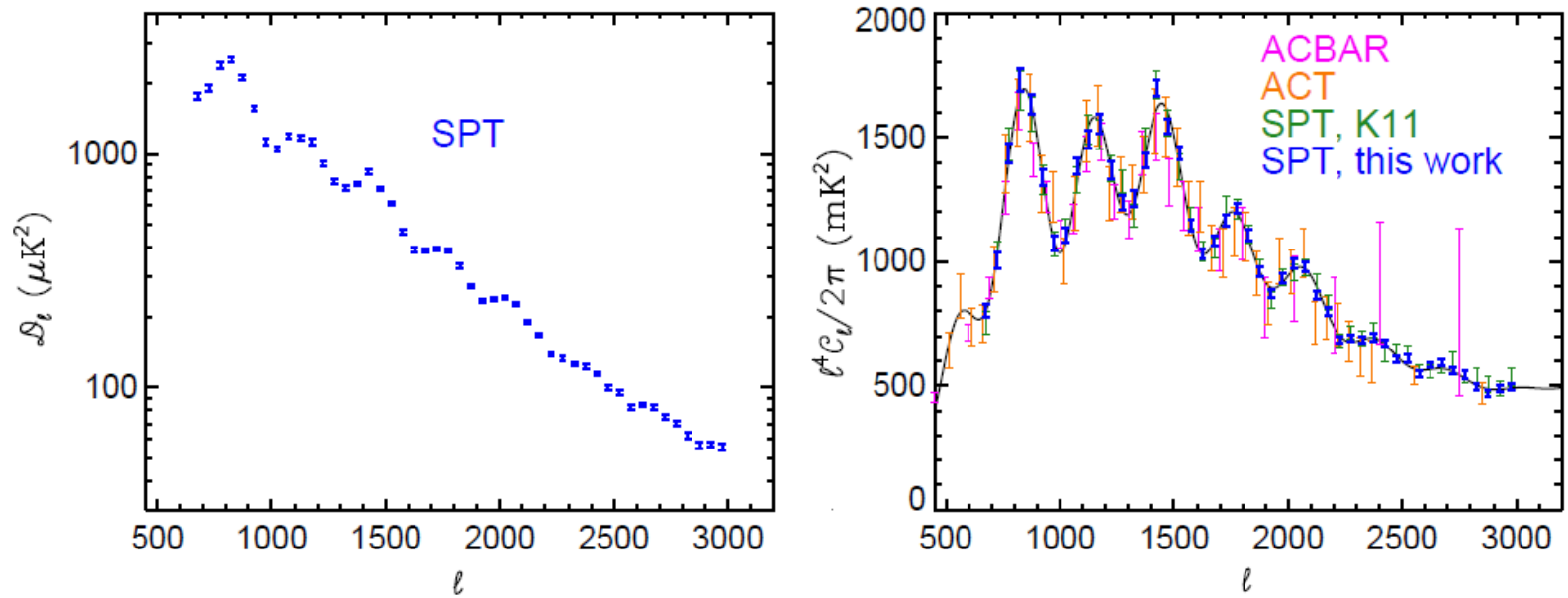
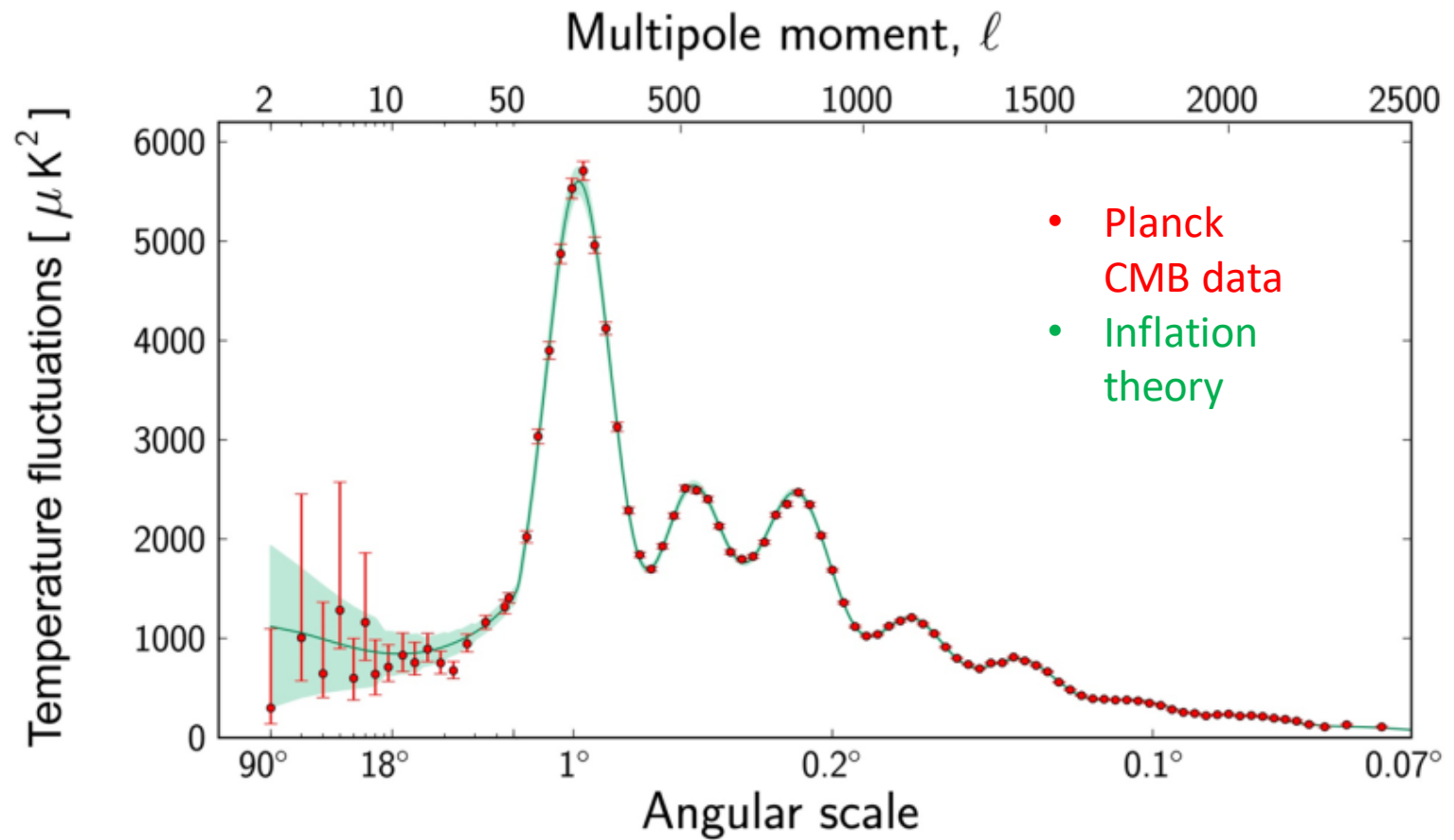


FIG. 3.— **Left panel:** The SPT power spectrum. The leftmost peak at  $\ell \sim 800$  is the third acoustic peak. **Right panel:** A comparison of the new SPT bandpowers with other recent measurements of the CMB damping tail from ACBAR (Reichardt et al. 2009), ACT (Das et al. 2011b), and SPT (K11). Note that the point source masking threshold differs between these experiments which can affect the power at the highest multipoles. In order to highlight the acoustic peak structure of the damping tail, we plot the bandpowers in the right panel as  $\ell^4 C_\ell / (2\pi)$ , as opposed to  $D_\ell = \ell(\ell + 1)C_\ell / (2\pi)$  in the left panel. The solid line shows the theory spectrum for the  $\Lambda$ CDM model + foregrounds that provides the best fit to the SPT+WMAP7 data. The bandpower errors shown in these plots contain sample and noise variance terms only; they do not include beam or calibration uncertainties.

## A MEASUREMENT OF THE COSMIC MICROWAVE BACKGROUND DAMPING TAIL FROM THE 2500-SQUARE-DEGREE SPT-SZ SURVEY

K. T. STORY,<sup>1,2</sup> C. L. REICHARDT,<sup>3</sup> Z. HOU,<sup>4</sup> R. KEISLER,<sup>1,2</sup> K. A. AIRD,<sup>5</sup> B. A. BENSON,<sup>1,6</sup> L. E. BLEEM,<sup>1,2</sup>  
J. E. CARLSTROM,<sup>1,2,6,7,8</sup> C. L. CHANG,<sup>1,6,8</sup> H.-M. CHO,<sup>9</sup> T. M. CRAWFORD,<sup>1,7</sup> A. T. CRITES,<sup>1,7</sup> T. DE HAAN,<sup>10</sup>  
M. A. DOBBS,<sup>10</sup> J. DUDLEY,<sup>10</sup> B. FOLLIN,<sup>4</sup> E. M. GEORGE,<sup>3</sup> N. W. HALVERSON,<sup>11</sup> G. P. HOLDER,<sup>10</sup> W. L. HOLZAPFEL,<sup>3</sup>  
S. HOOVER,<sup>1,2</sup> J. D. HRUBES,<sup>5</sup> M. JOY,<sup>12</sup> L. KNOX,<sup>4</sup> A. T. LEE,<sup>3,13</sup> E. M. LEITCH,<sup>1,7</sup> M. LUEKER,<sup>14</sup> D. LUONG-VAN,<sup>5</sup>  
J. J. McMAHON,<sup>15</sup> J. MEHL,<sup>8,1</sup> S. S. MEYER,<sup>1,2,6,7</sup> M. MILLEA,<sup>4</sup> J. J. MOHR,<sup>16,17,18</sup> T. E. MONTROY,<sup>19</sup> S. PADIN,<sup>1,7,14</sup>  
T. PLAGGE,<sup>1,7</sup> C. PRYKE,<sup>20</sup> J. E. RUHL,<sup>19</sup> J. T. SAYRE,<sup>19</sup> K. K. SCHAFER,<sup>1,6,21</sup> L. SHAW,<sup>10</sup> E. SHIROKOFF,<sup>3</sup>  
H. G. SPIELER,<sup>13</sup> Z. STANISZEWSKI,<sup>19</sup> A. A. STARK,<sup>22</sup> A. VAN ENGELEN,<sup>10</sup> K. VANDERLINDE,<sup>10</sup> J. D. VIEIRA,<sup>14</sup>  
R. WILLIAMSON,<sup>1,7</sup> AND O. ZAHN<sup>23</sup>

Submitted to ApJ



## The Basic Tools of Inflation:

Consider a scalar field with:

$$\mathcal{L}(\varphi) = \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - V(\varphi)$$

⇒ If  $V(\varphi) \gg$  all space and time derivative (squared) terms

⇒ Then

$$T_\mu^\nu \approx \begin{pmatrix} V(\varphi) & 0 & 0 & 0 \\ 0 & -V(\varphi) & 0 & 0 \\ 0 & 0 & -V(\varphi) & 0 \\ 0 & 0 & 0 & -V(\varphi) \end{pmatrix}$$

⇒ Which implies  $p = -\rho$   $w = -1$  ⇒  $\frac{d\rho}{da} \approx 0$

⇒  $a \sim e^{Ht}$  } Inflation

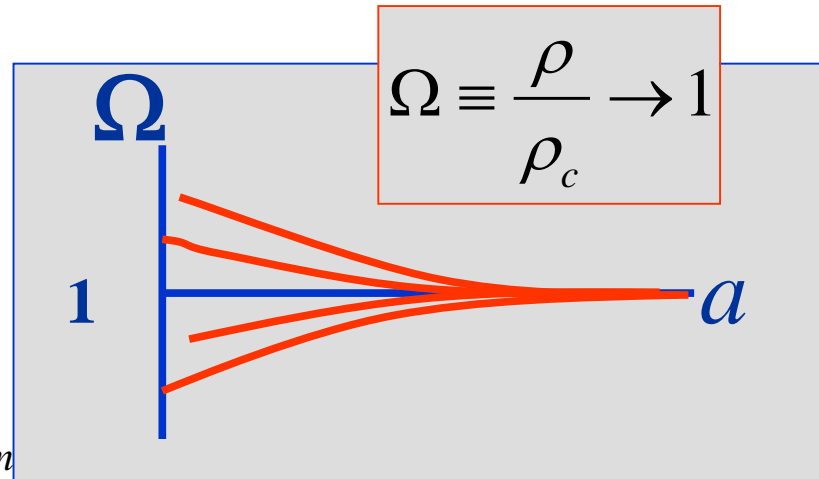
# A period of early inflation gives:

## Flatness:

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2 = \frac{8\pi}{3}\rho - \frac{k}{a^2}$$

$$\rho_\phi = \text{const.}$$

Dominates over  $\rho_r$  &  $\rho_m$   
during inflation



## Homogeneity

At horizon crossing:

$$\left.\frac{\delta\rho}{\rho}\right|_H \equiv \delta_H \approx \left(\frac{H^2}{\dot{\phi}}\right) \approx \left(\frac{HV'}{2\pi\dot{\phi}^2}\right) \approx \text{const}$$

Evaluate when  $k=H$  during inflation

A suitably adjusted potential will give :

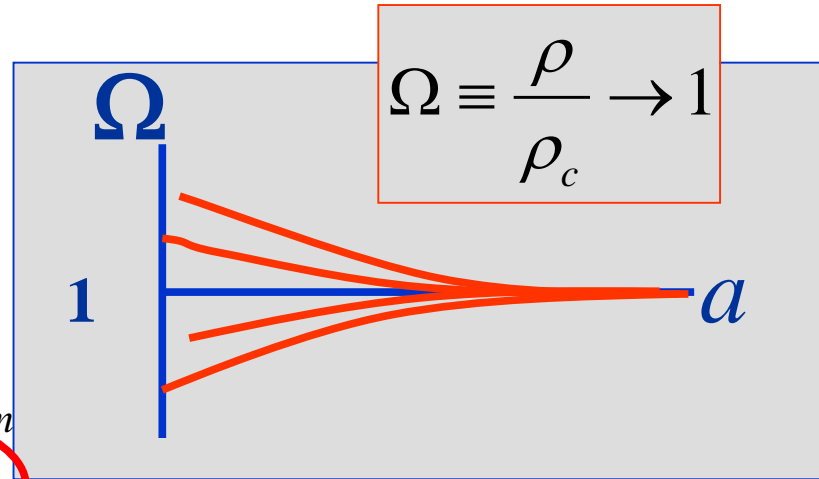
$$\delta_H \approx 10^{-5}$$

# A period of early inflation gives:

## Monopoles:

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2 = \frac{8\pi}{3} \rho - \frac{k}{a^2}$$

$\rho_\phi$  Dominates over  $\rho_r$  &  $\rho_m$  during inflation (and  $\rho_M$ )

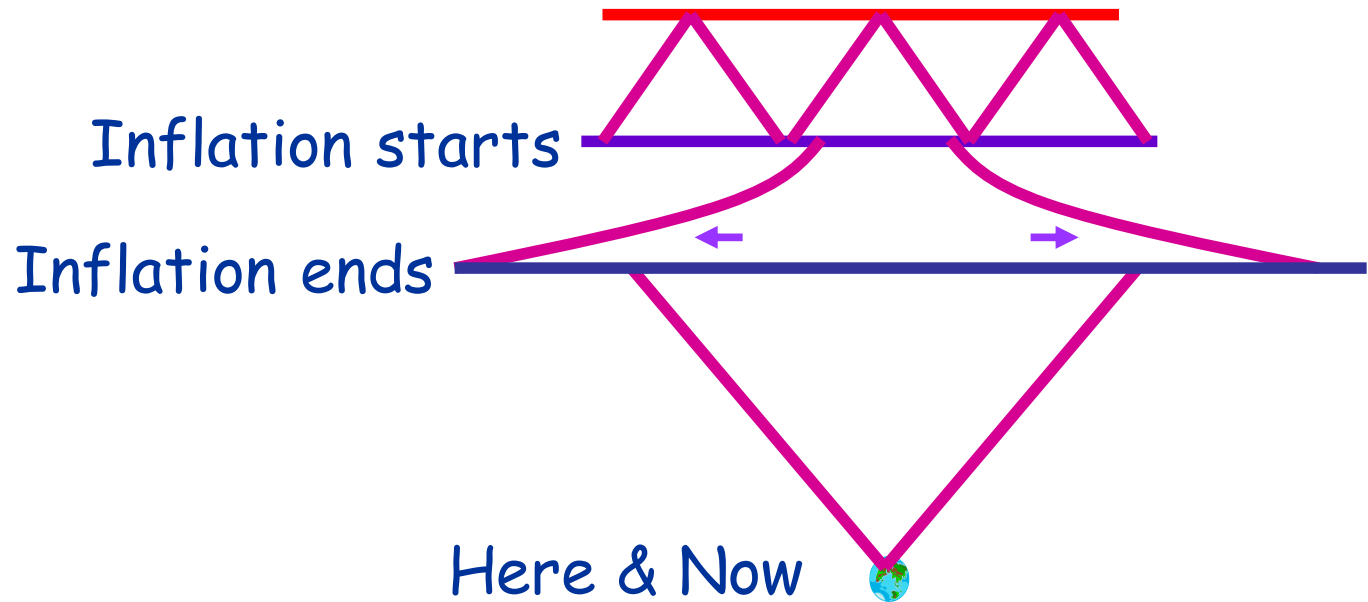


$$\left. \frac{\rho_M}{\rho_\phi} \right|_a = \frac{\rho_M(a_1) \left( \frac{a_1}{a} \right)^3}{\rho_\phi(a) = \text{Const}} = \left( \frac{\rho_M(a_1)}{\rho_\phi(a_1)} \right) \times e^{-3Ht} \longrightarrow 0$$

$$a \sim e^{Ht}$$

Monopoles are erased

## Inflation Horizon:



# I) Inflation in the era of WMAP

## I.0 What is Cosmic Inflation?

### I.1 Successes

## Inflation and the arrow of time

Introduction

Arrow of time basics

Inflation and the arrow of time

Implications

Can the Universe Afford Inflation?



## ➤ Inflation:

- An early period of nearly exponential expansion set up the “initial” conditions for the standard big bang

## ➤ Predictions:

- $\Omega_{\text{total}}=1$  (to one part in 100,000 as measured)
- Characteristic oscillations in the CMB power
- Nearly scale invariant perturbation spectrum
- Characteristic Gravity wave, CMB Polarization etc
- etc

# WMAP

- $\Omega_{\text{total}}=1$

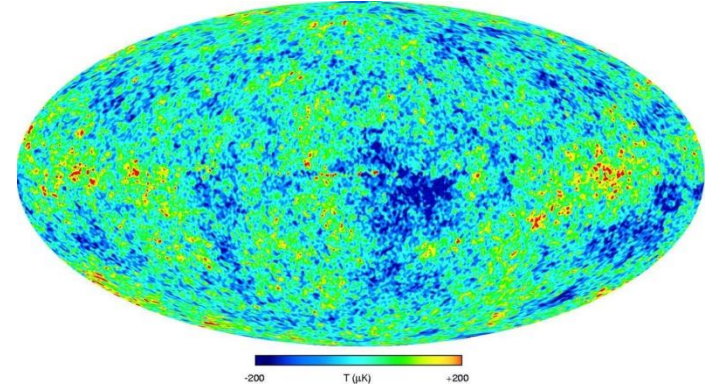


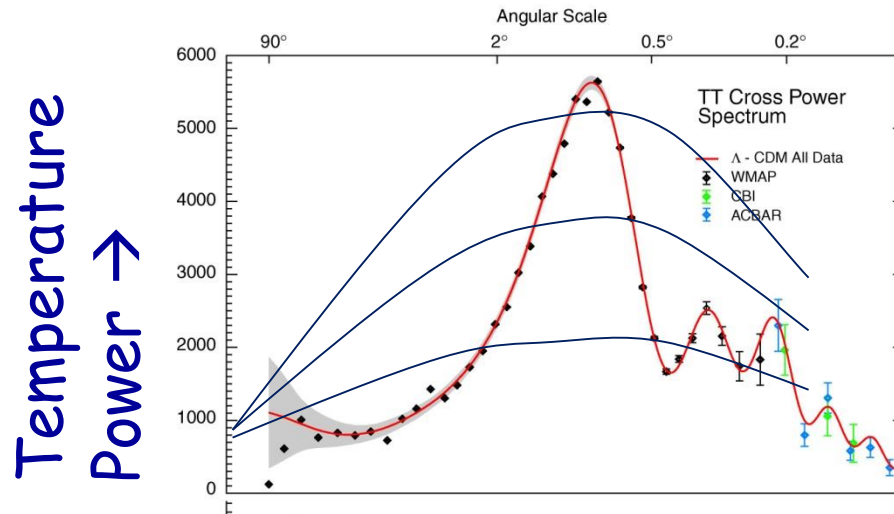
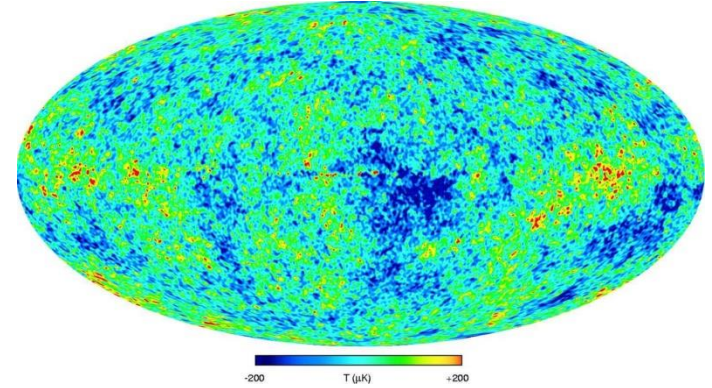
Table 3. “Best” Cosmological Parameters

Description	Symbol	Value	+ uncertainty	– uncertainty
Total density	$\Omega_{tot}$	1.02	0.02	0.02
Equation of state of quintessence	$w$	$< -0.78$	95% CL	—
Dark energy density	$\Omega_{\Lambda}$	0.73	0.04	0.04
Baryon density	$\Omega_b h^2$	0.0224	0.0009	0.0009
Baryon density	$\Omega_b$	0.044	0.004	0.004
Baryon density (cm <sup>-3</sup> )	$n_b$	$2.5 \times 10^{-7}$	$0.1 \times 10^{-7}$	$0.1 \times 10^{-7}$
Matter density	$\Omega_m h^2$	0.135	0.008	0.009
Matter density	$\Omega_m$	0.27	0.04	0.04
Light neutrino density	$\Omega_{\nu} h^2$	$< 0.0076$	95% CL	—

Bennett et al Feb 11 '03

# WMAP

• Characteristic oscillations in the CMB power



— Inflation  
— "Active" models

← Angular scale

Adapted from  
Bennett et al Feb 11 '03

# WMAP

- Nearly scale invariant perturbation spectrum

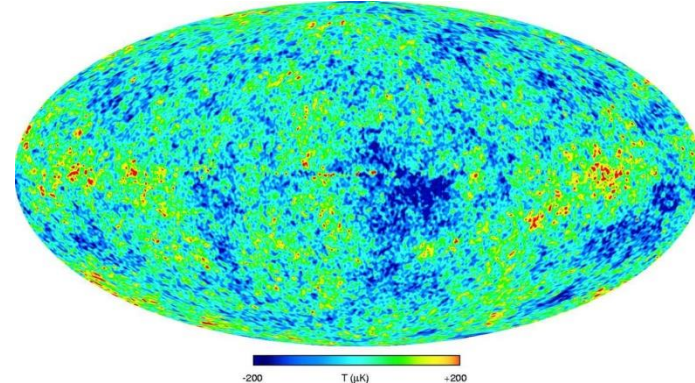


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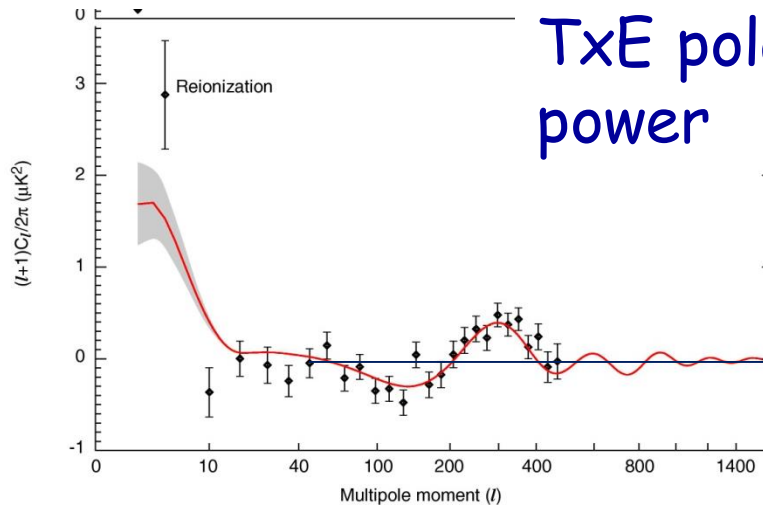
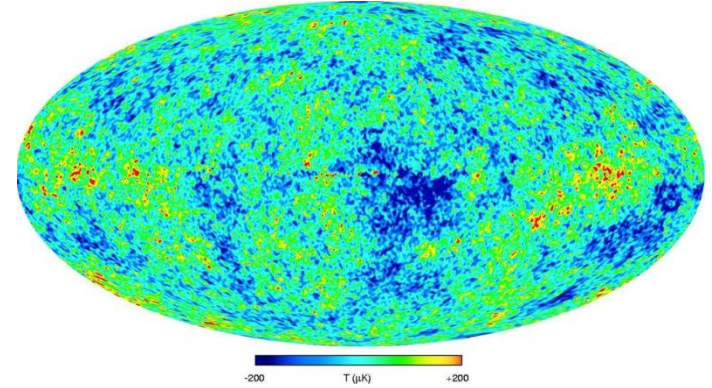
Description	Symbol	Value	+ uncertainty	– uncertainty
Power spectrum normalization (at $k_0 = 0.05 \text{ Mpc}^{-1}$ ) <sup>c</sup>	$A$	0.833	0.086	0.083
Scalar spectral index (at $k_0 = 0.05 \text{ Mpc}^{-1}$ ) <sup>c</sup>	$n_s$	0.93	0.03	0.03
Running index slope (at $k_0 = 0.05 \text{ Mpc}^{-1}$ ) <sup>c</sup>	$dn_s/d\ln k$	–0.031	0.016	0.018
Tensor-to-scalar ratio (at $k_0 = 0.002 \text{ Mpc}^{-1}$ )	$r$	$< 0.71$	95% CL	—
Redshift of decoupling	$z_{dec}$	1089	1	1
1 $\sigma$ uncertainty in $z_{dec}$	$\Delta z_{dec}$	195	2	2
Hubble constant	$h$	0.71	0.04	0.03

$$\left. \frac{\delta \rho}{\rho}(k) \right|_{H=k} = A k^{1-n_s}$$

Bennett et al Feb 11 '03

# WMAP

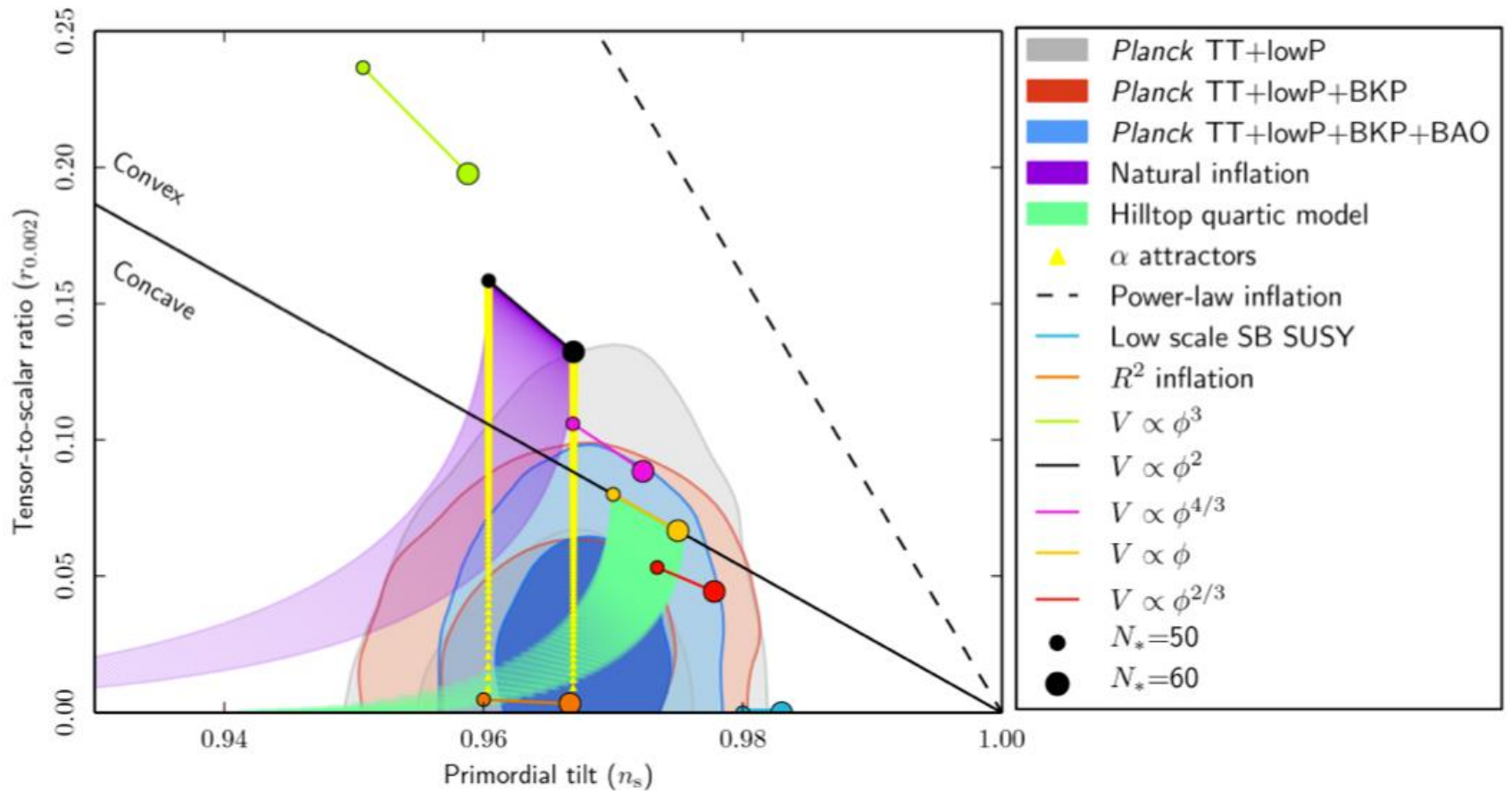
- Characteristic Gravity wave, CMB Polarization etc



TxE polarization power

— Inflation  
— "Active" models

Bennett et al Feb 11 '03



[https://www.aanda.org/articles/aa/full\\_html/2016/10/aa25898-15/aa25898-15.html](https://www.aanda.org/articles/aa/full_html/2016/10/aa25898-15/aa25898-15.html)

fig 55

# What happened before inflation?

## What happened before inflation?

Related to some more contentious topics:

- How well does inflation “solve” tuning problems?
- Eternal inflation



# OUTLINE

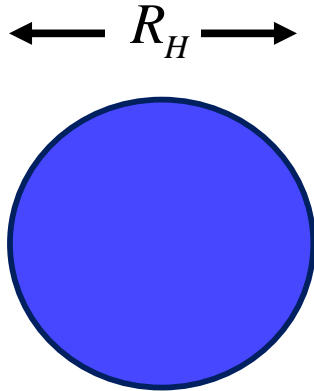
1. Big Bang basics
2. Inflation basics
3. Eternal inflation
4. An alternative to Eternal Inflation
5. Further thoughts

# OUTLINE

1. Big Bang basics
2. Inflation basics
3. Eternal inflation
4. An alternative to Eternal Inflation
5. Further thoughts

- Eternal inflation

# Quantum fluctuations during slow roll:



A region of one field coherence length (  $= R_H$  ) gets a new quantum contribution to the field value from an uncorrelated comoving mode of size  $\Delta\phi = H$  in a time  $\Delta t = H^{-1}$  leading to a (random) quantum rate of change:

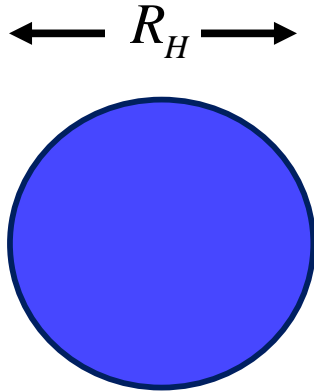
$$\frac{\Delta\phi}{\Delta t} \equiv \dot{\phi}_Q = H^2$$

Thus

$$\frac{\dot{\phi}_Q}{\dot{\phi}} = \frac{H^2}{\dot{\phi}}$$

measures the importance of quantum fluctuations in the field evolution

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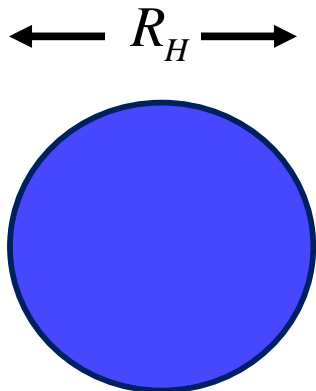
$$\frac{\Delta\phi}{\Delta t} \equiv \dot{\phi}_Q = H^2$$

Thus

$$\frac{\dot{\phi}_Q}{\dot{\phi}} = \frac{H^2}{\dot{\phi}} \quad \left( \approx \frac{\delta\rho}{\rho} \approx 10^{-5} \right)$$

measures the importance of quantum fluctuations in the field evolution

# Quantum fluctuations during slow roll:



For realistic perturbations the evolution is very classical

A region of one field coherence length ( $= R_H$ ) gets a new quantum contribution to the field value from an uncorrelated comoving mode of size  $\Delta\phi = H$  in a time  $\Delta t = H^{-1}$  leading to a (random) quantum rate of change:

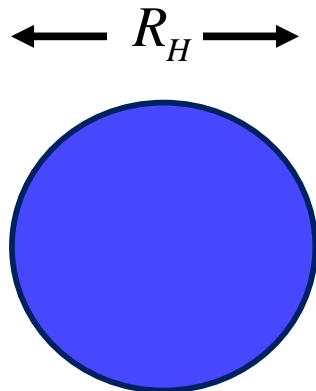
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measures the importance of quantum fluctuations in the field evolution

# Quantum fluctuations during slow roll:



For realistic perturbations the evolution is very classical

(But not as classical as most classical things we know!)

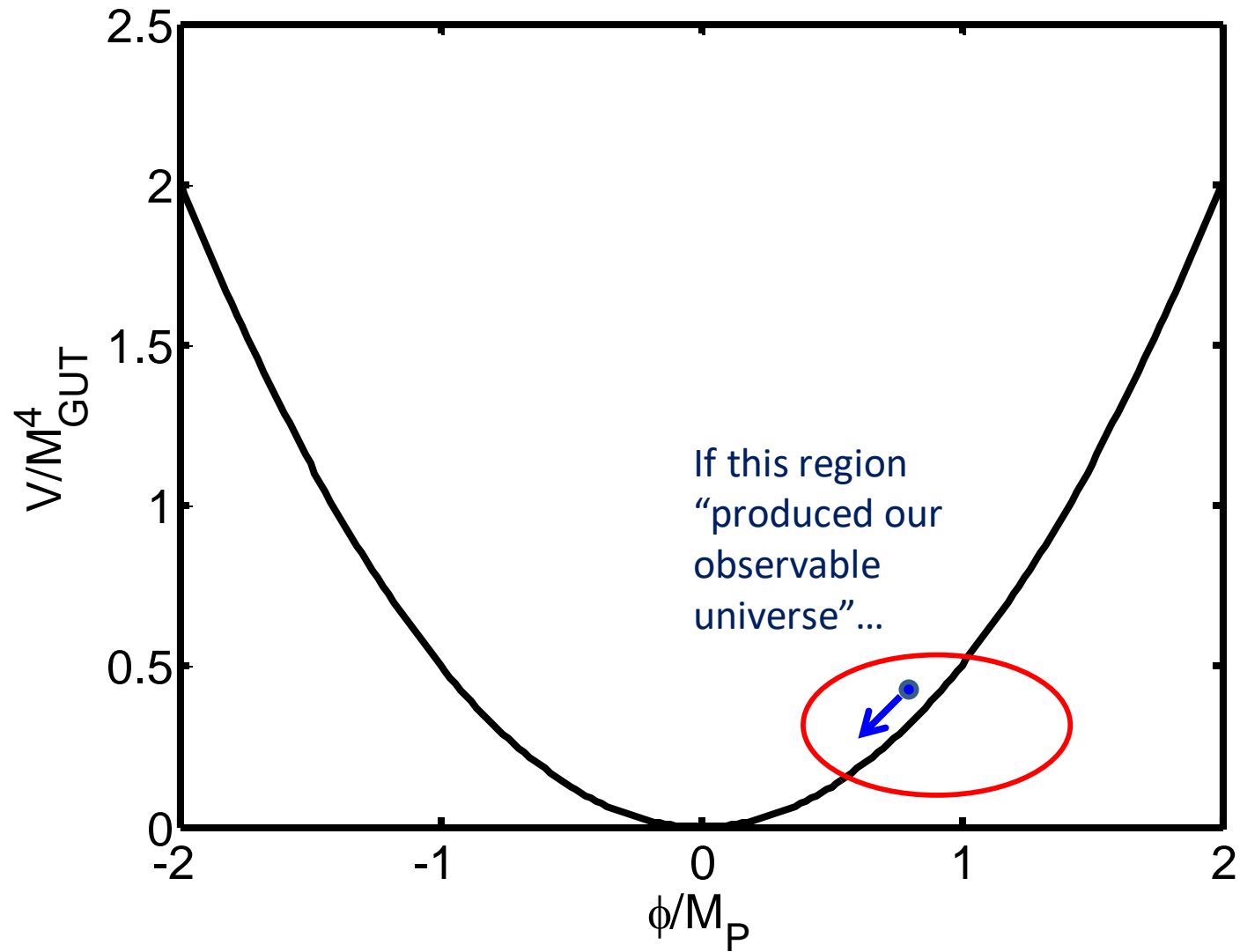
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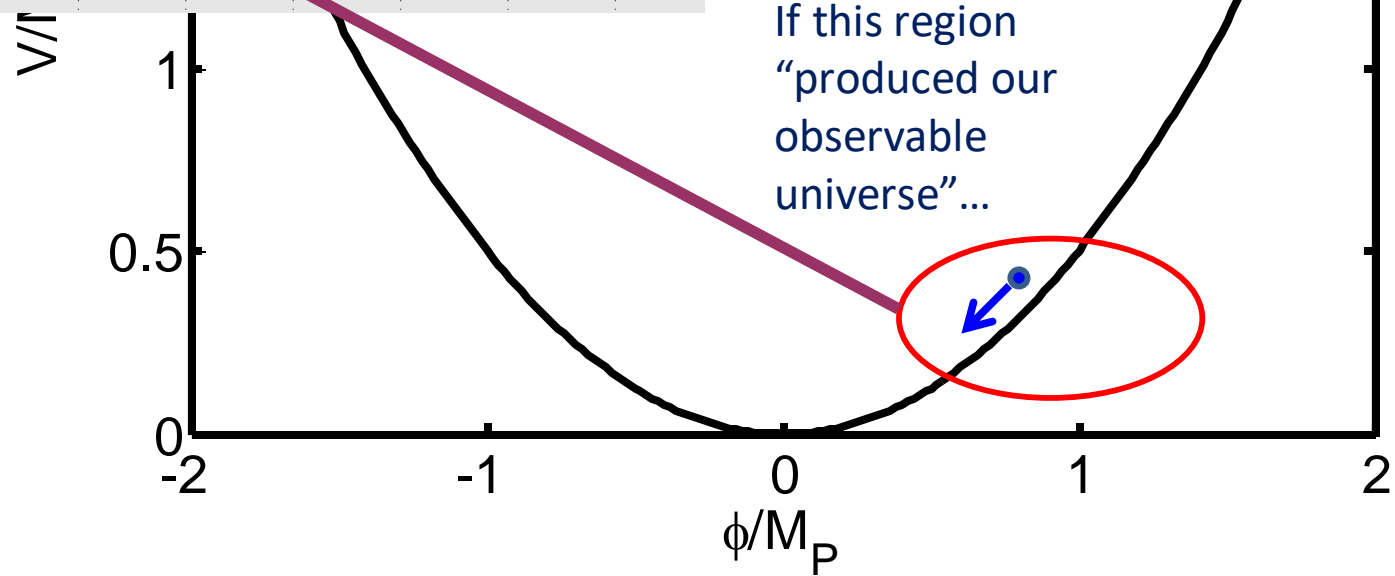
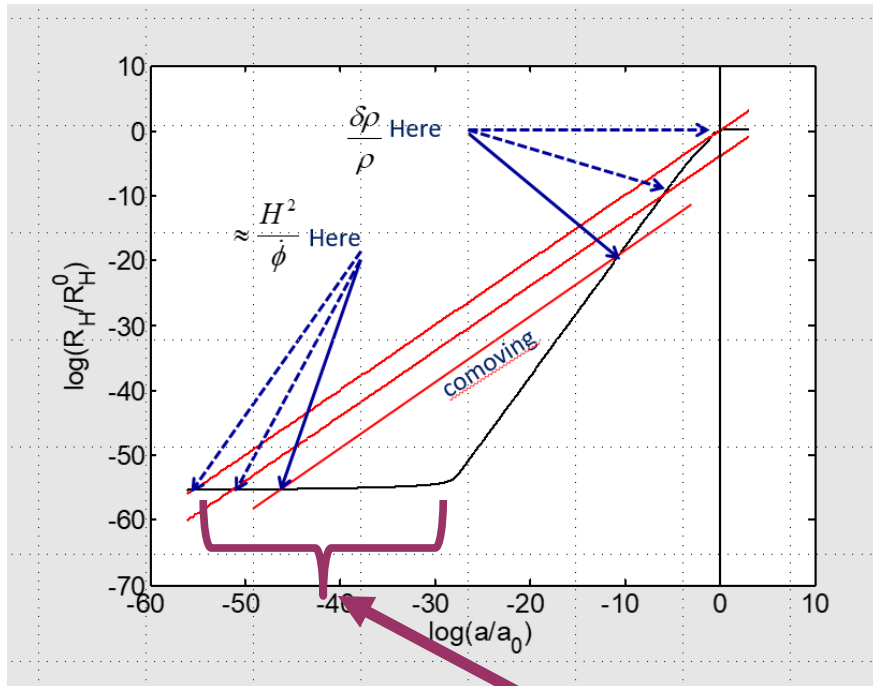
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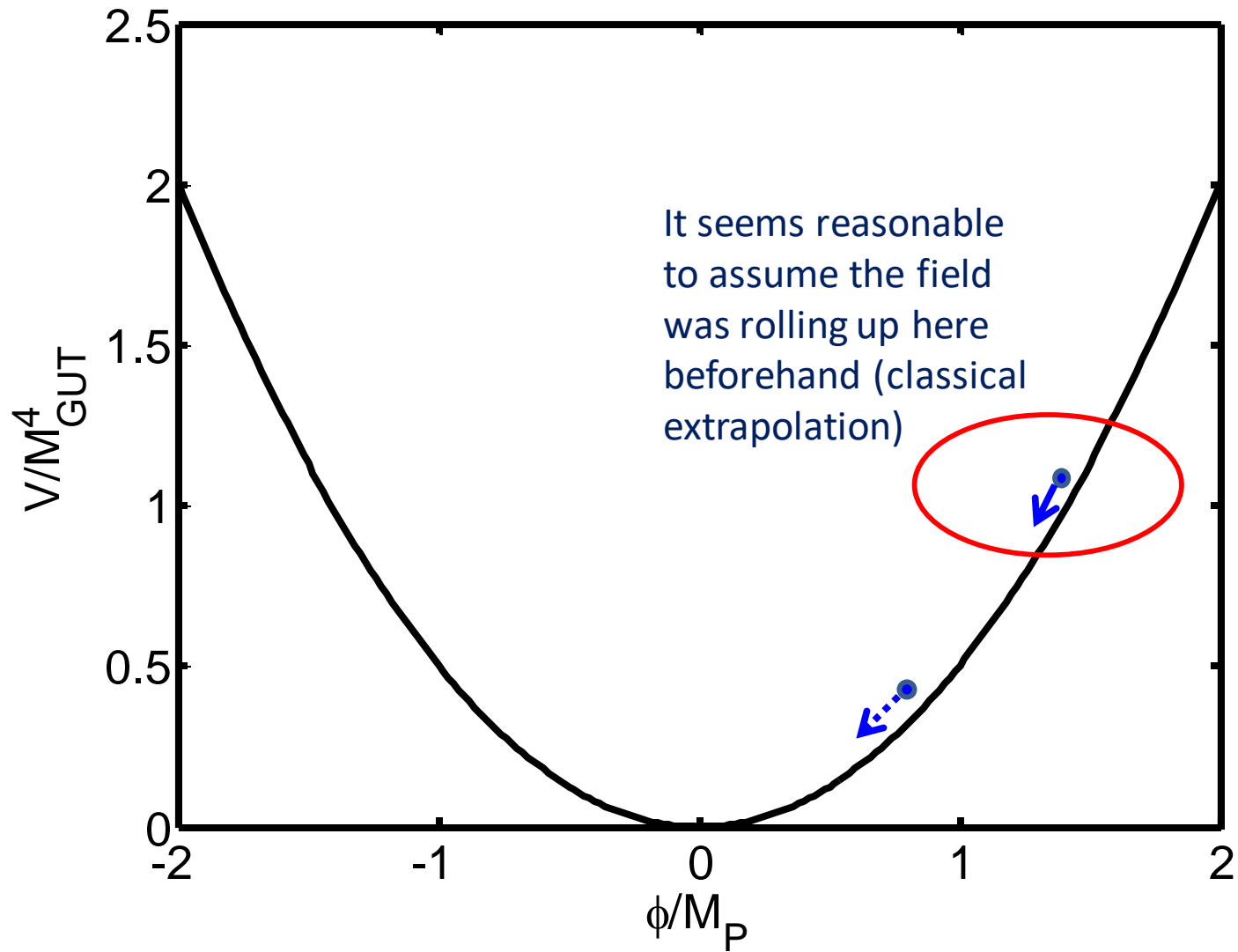
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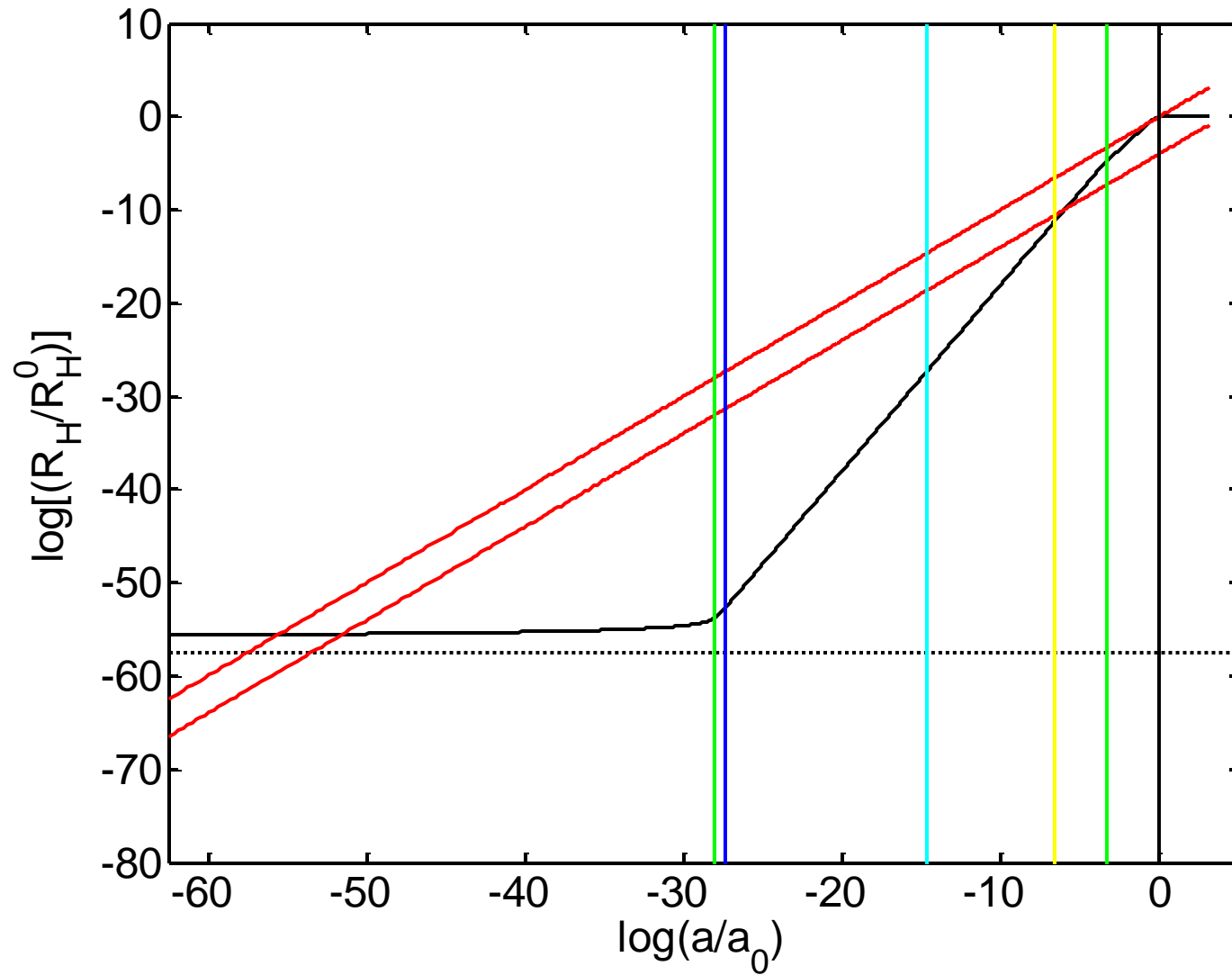




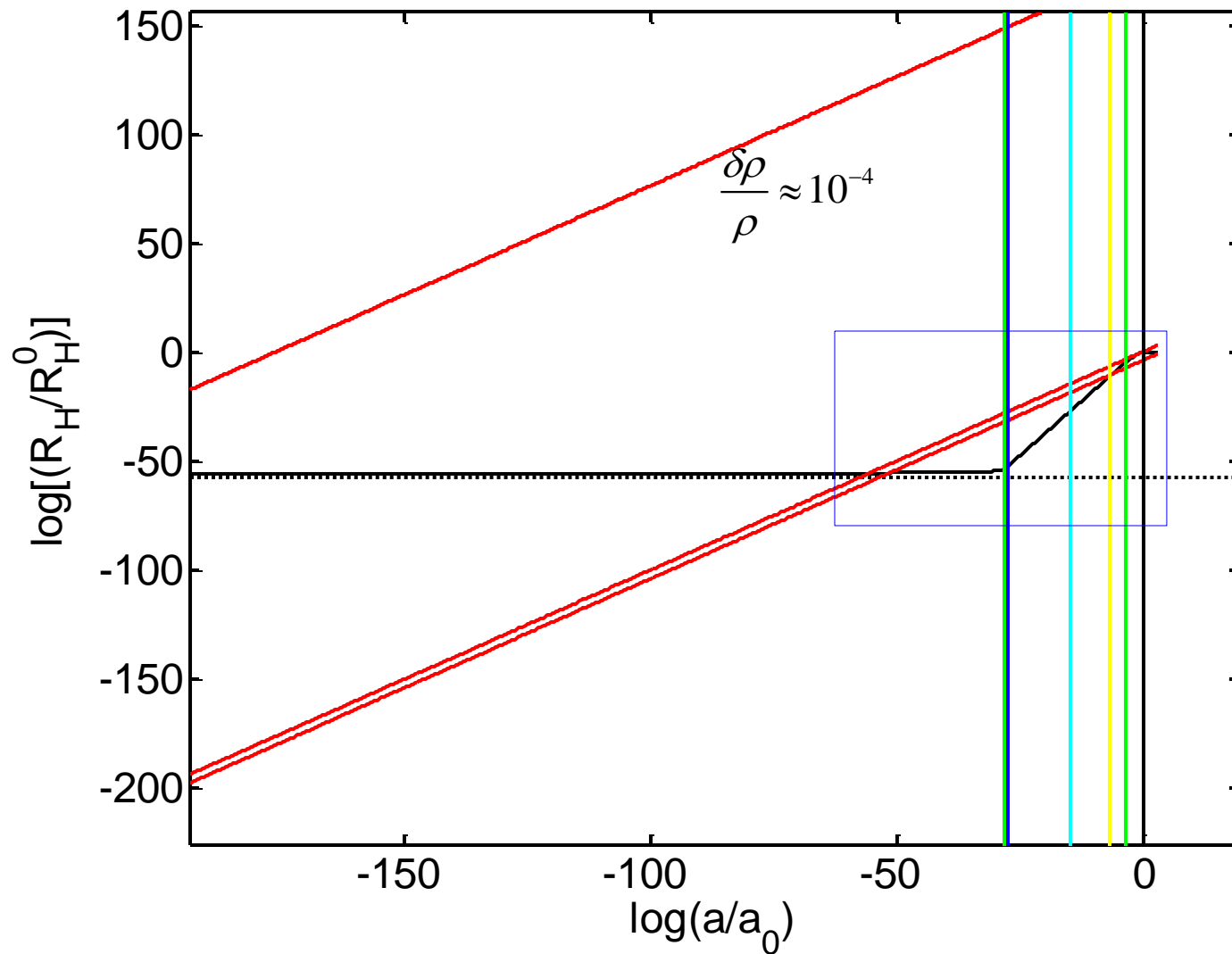




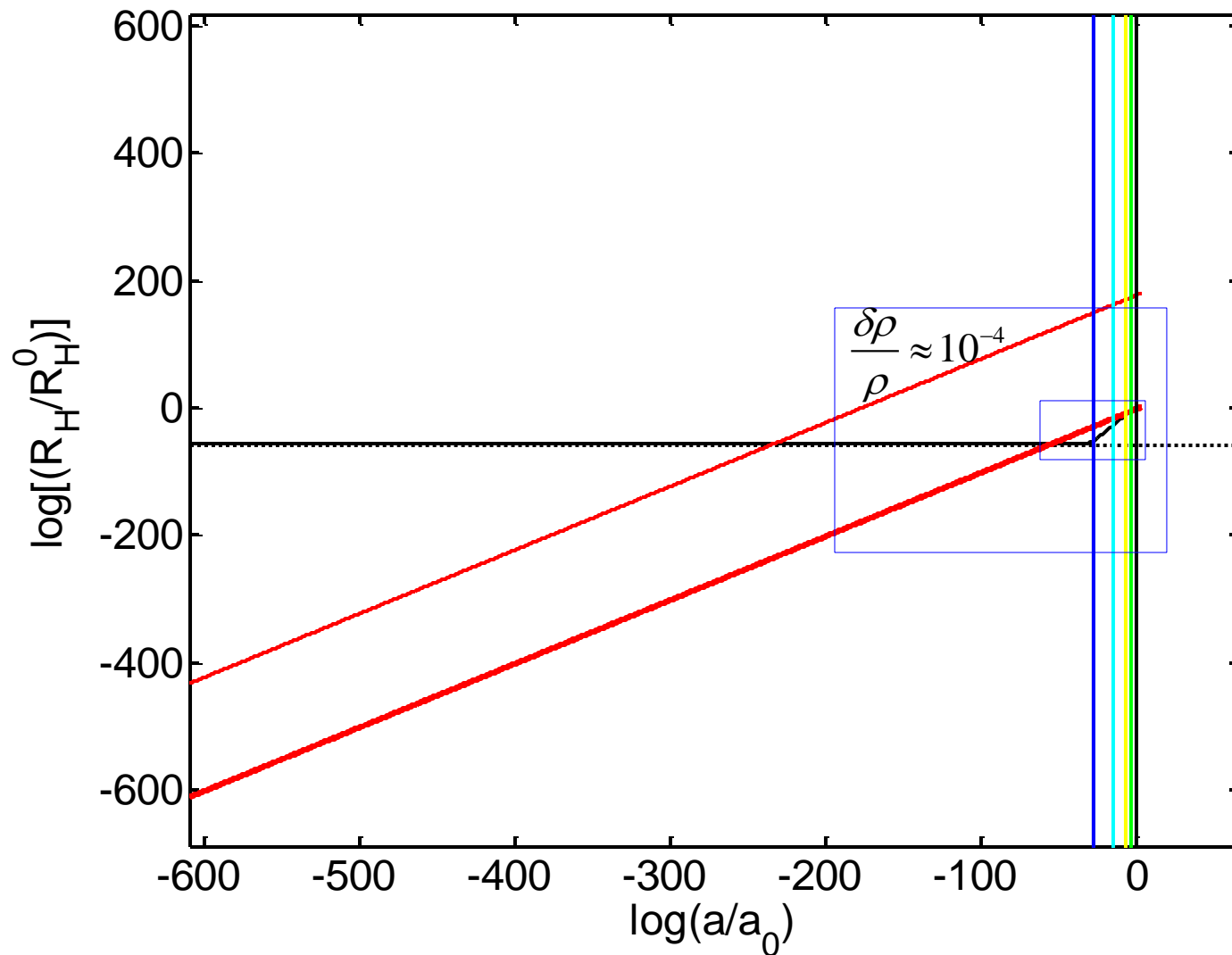
# Evolution of Cosmic Length



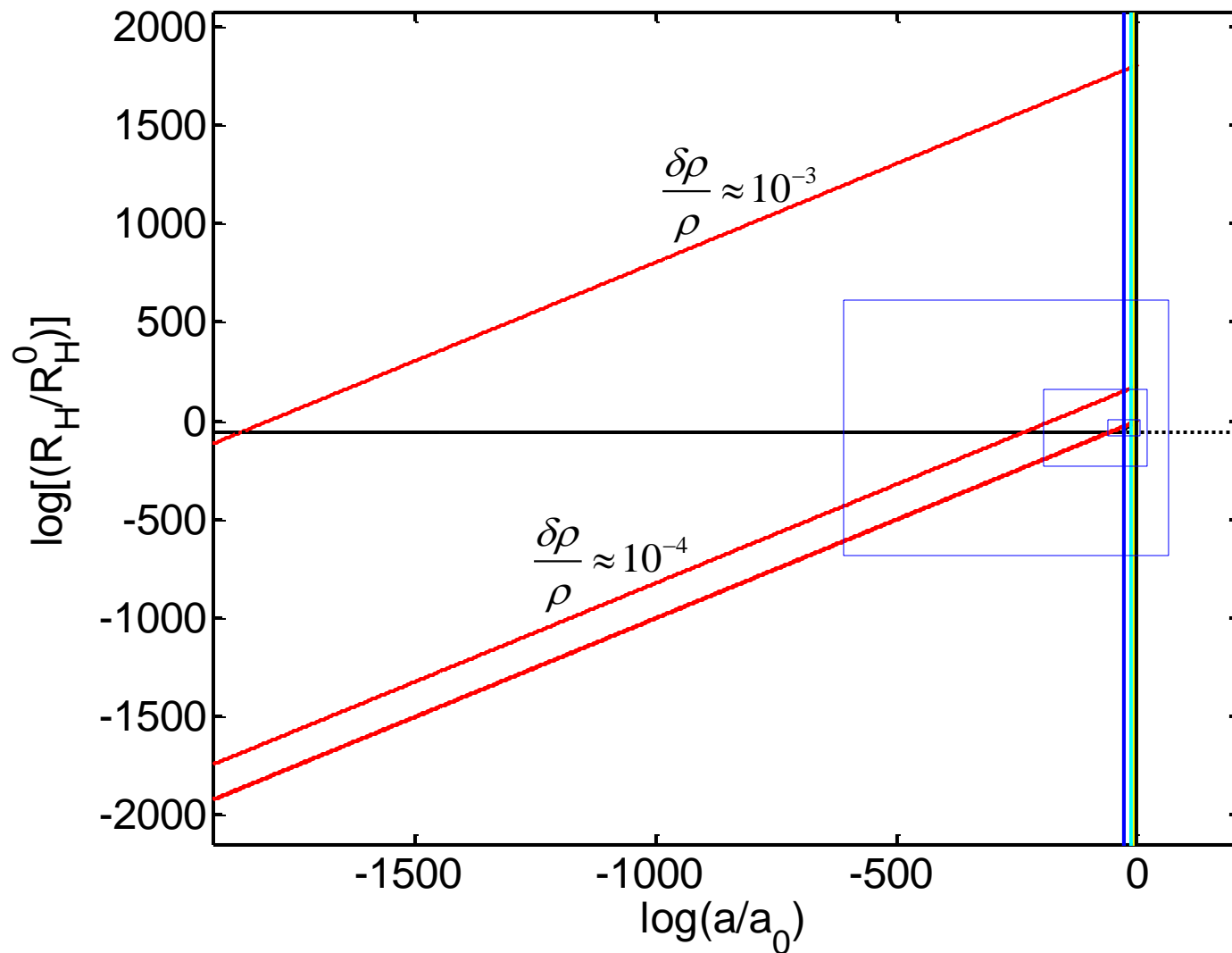
# Evolution of Cosmic Length (zooming out)



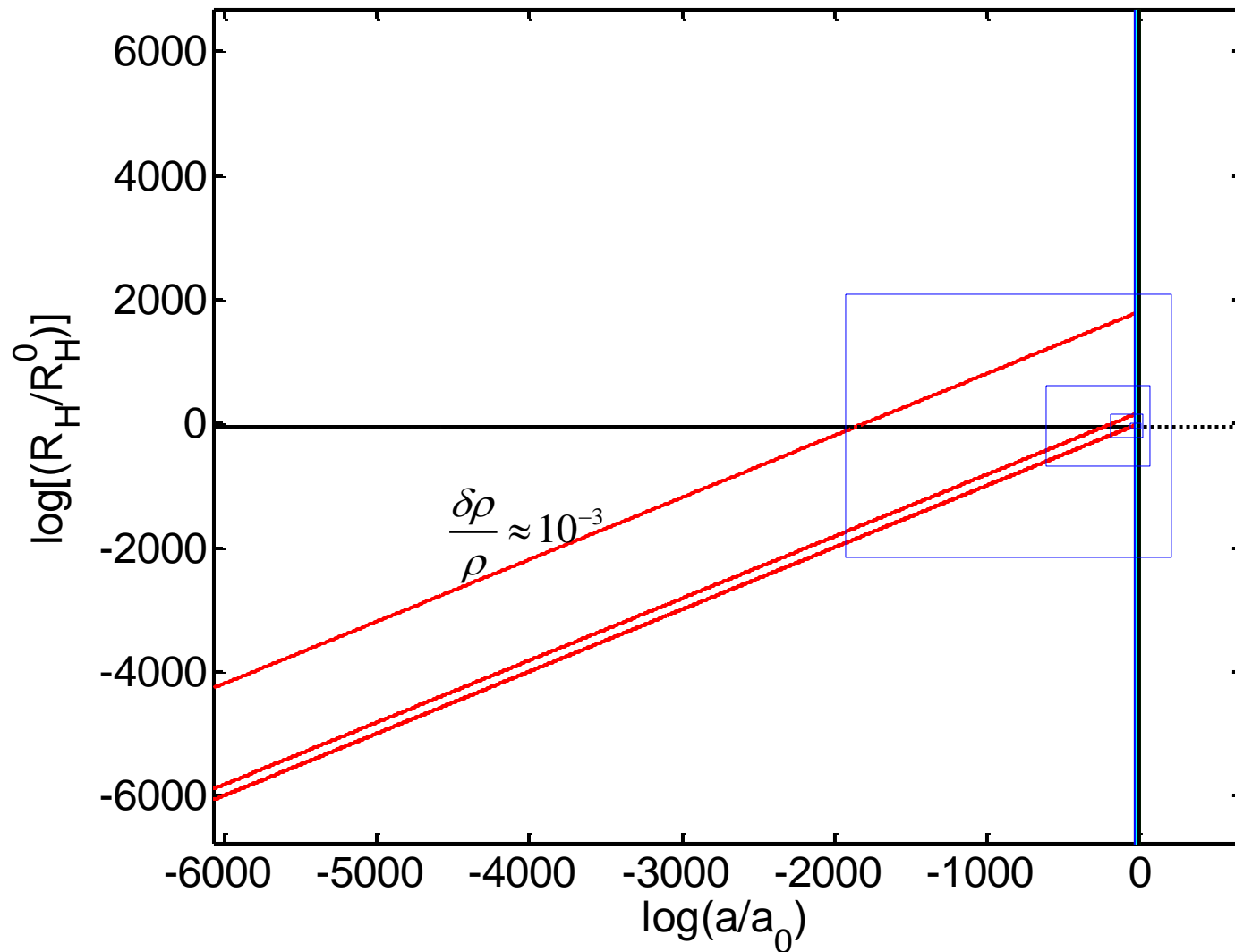
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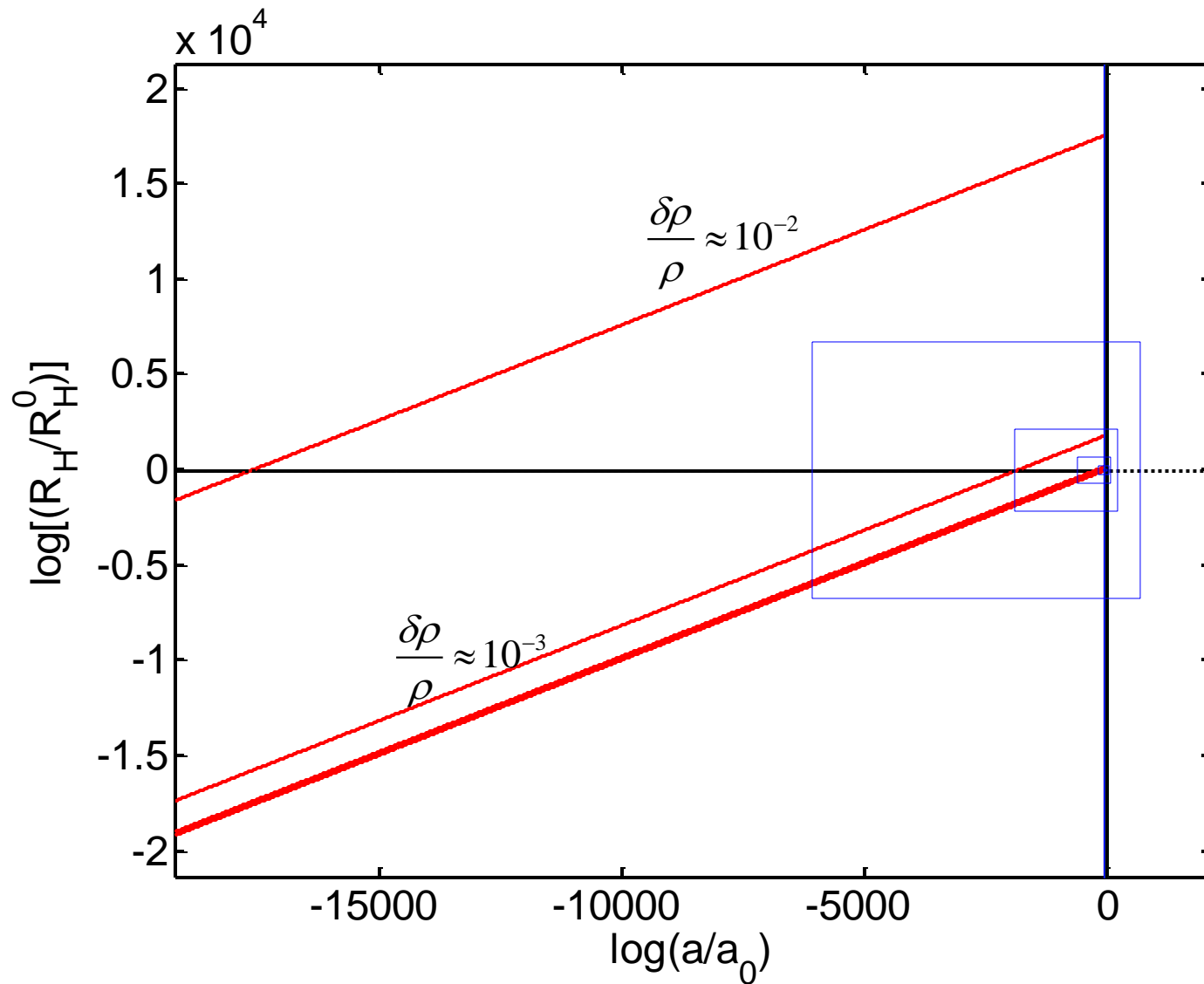
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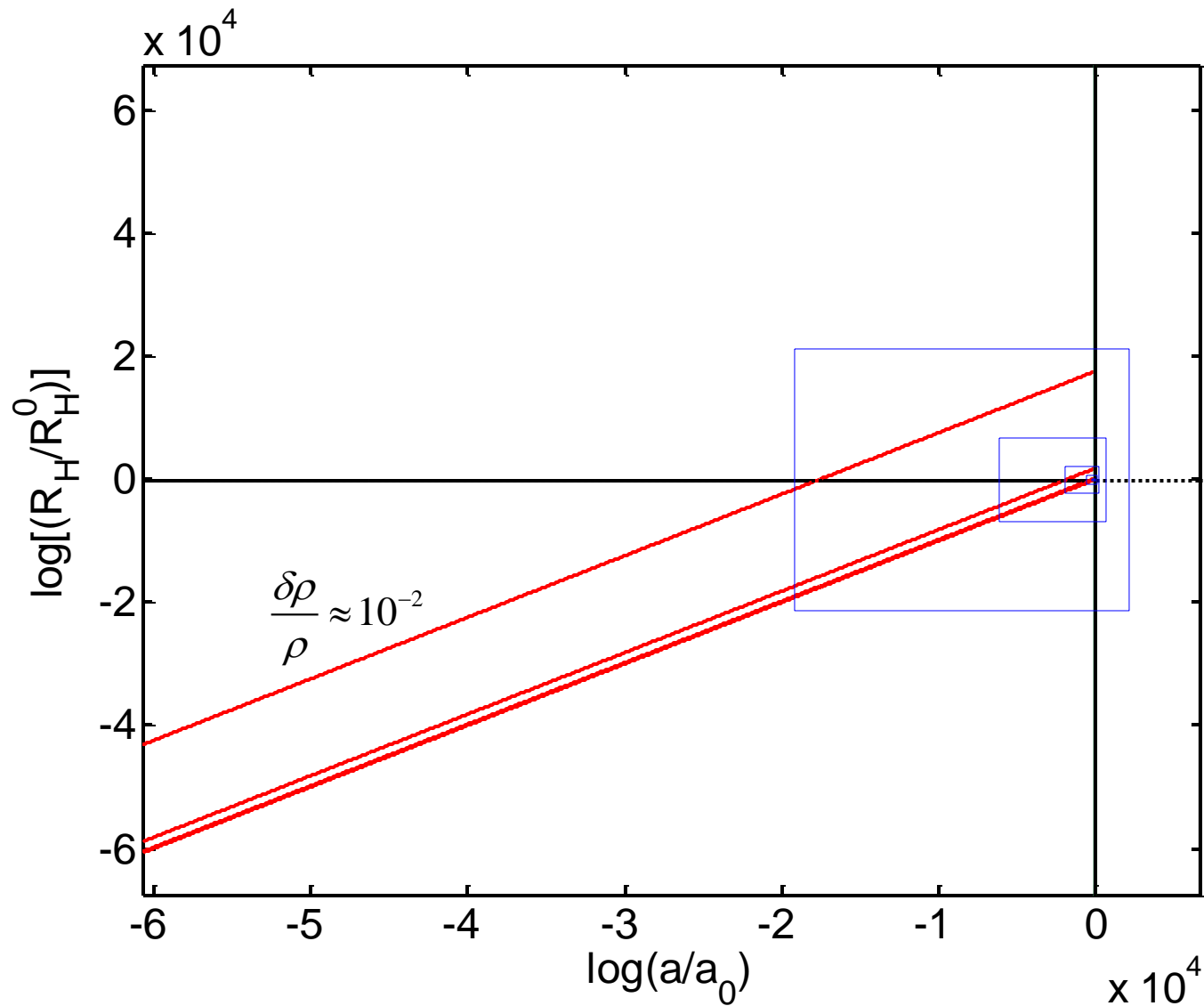


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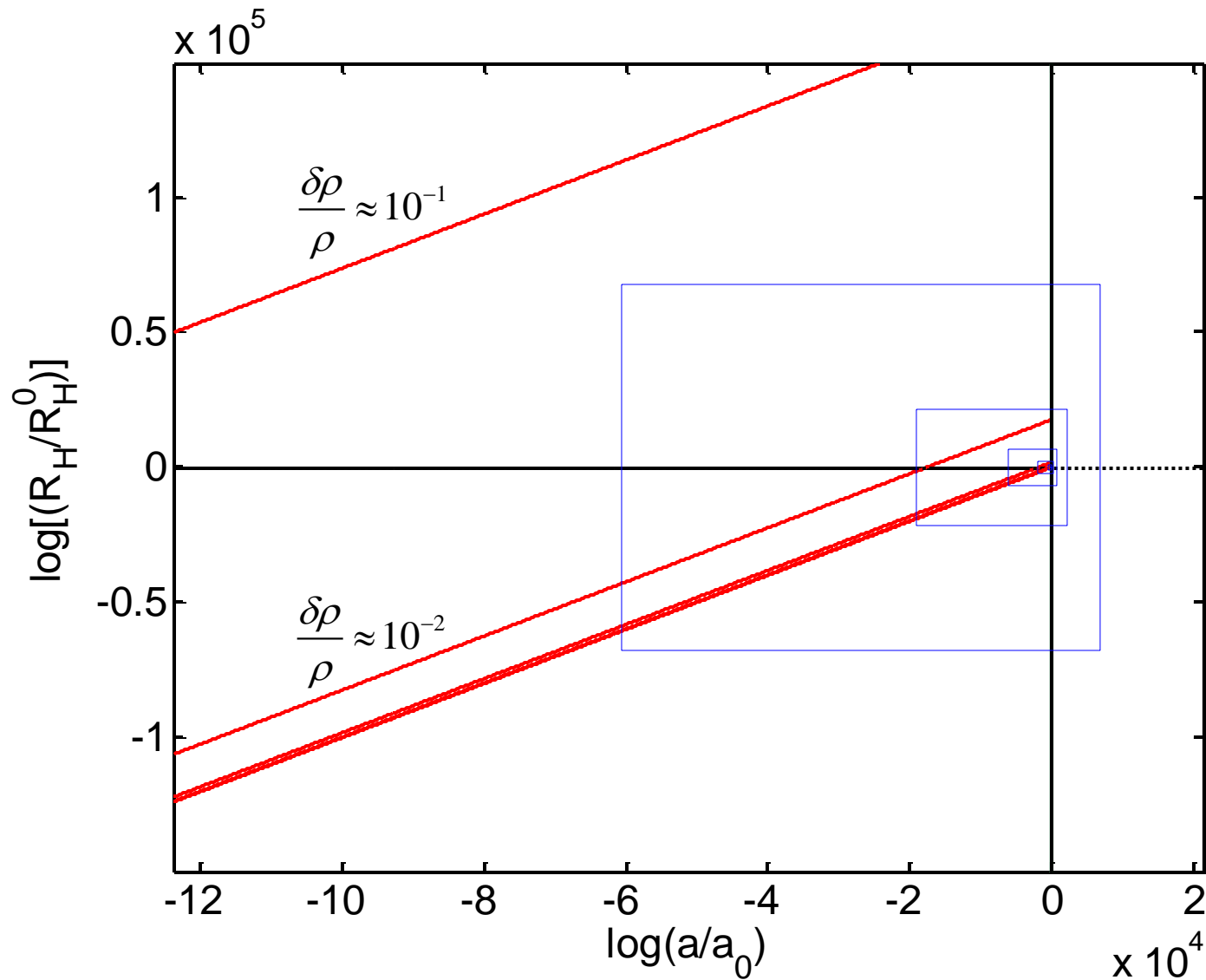




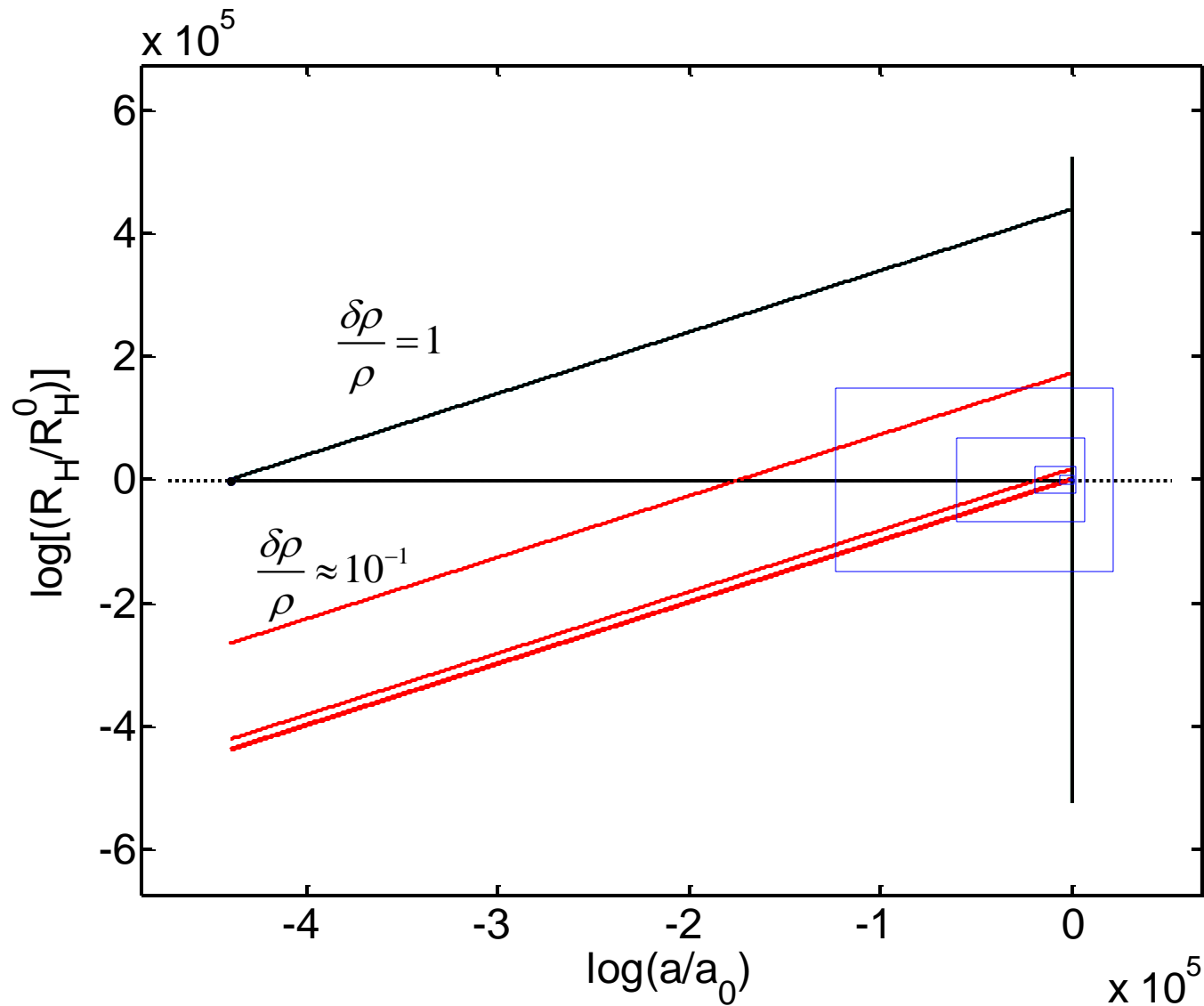
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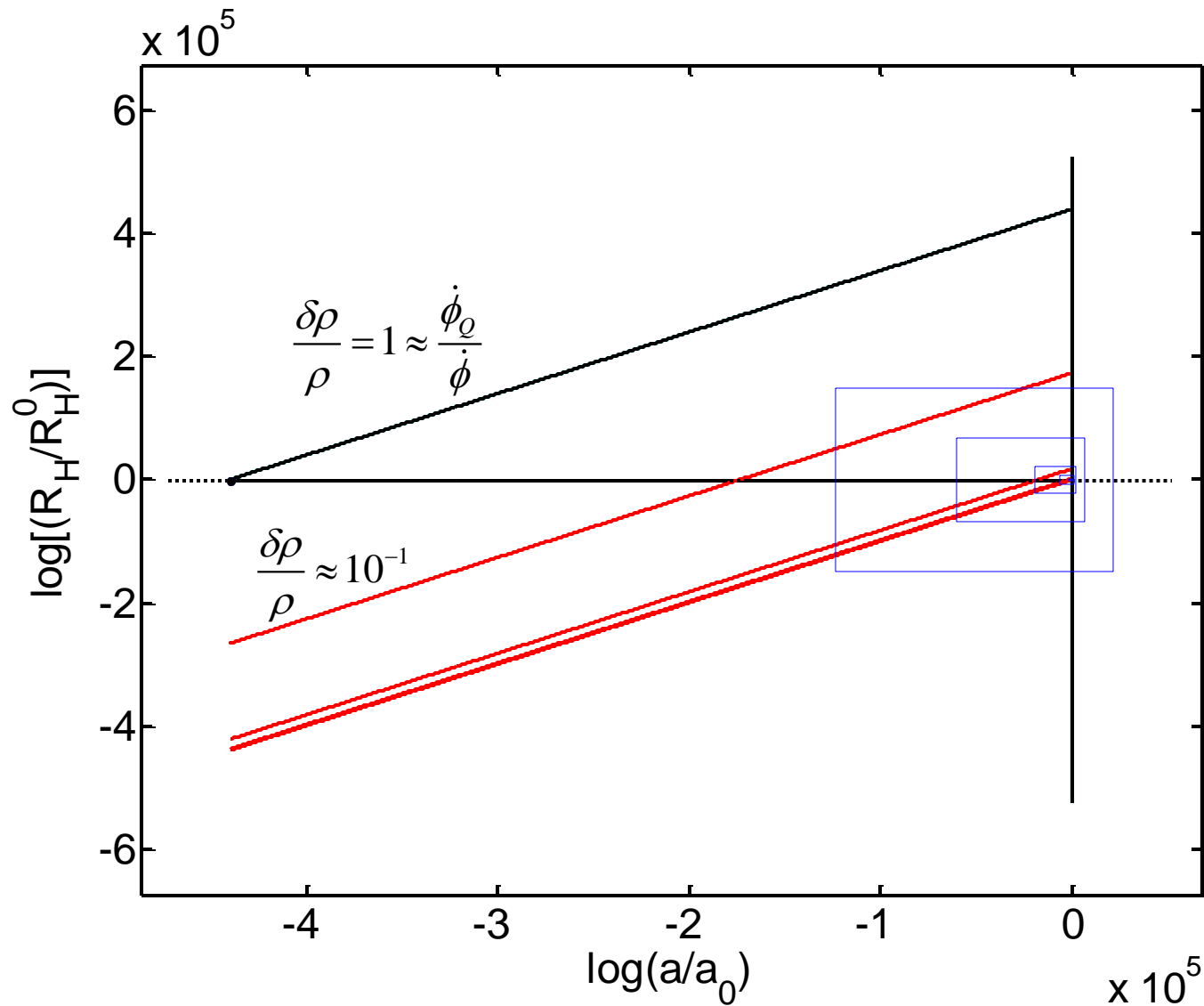
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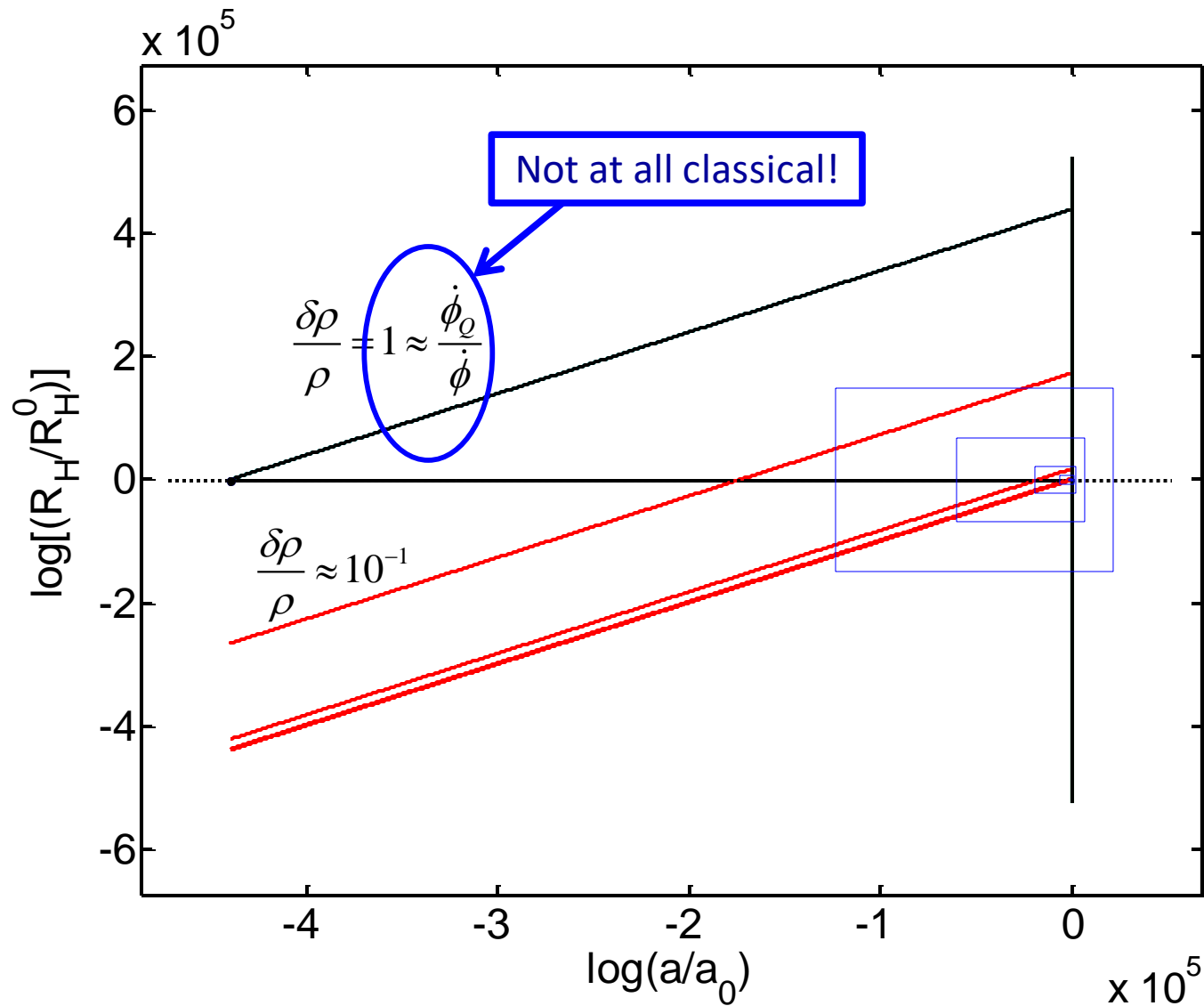
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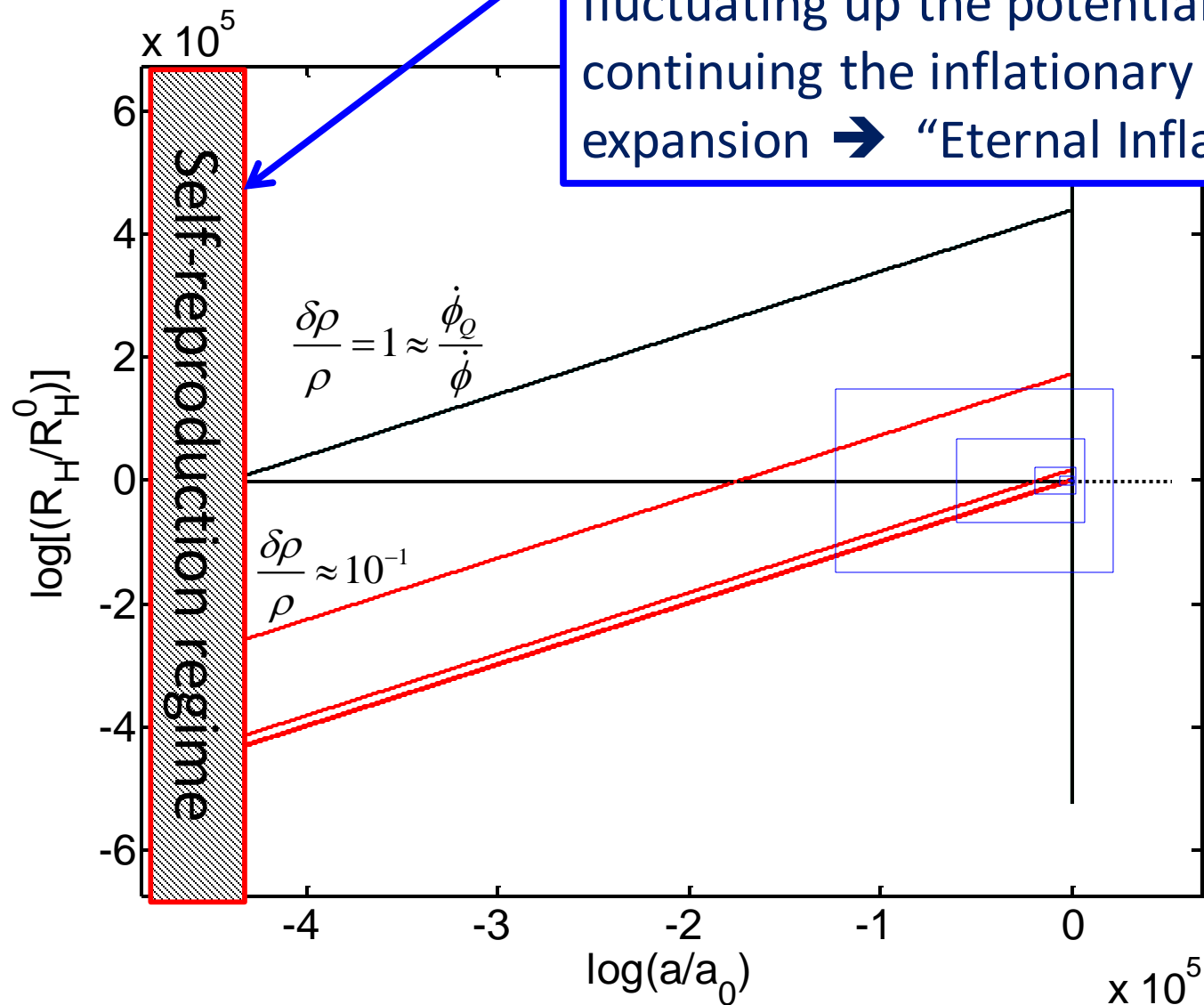


# Evolution of Cosmic Length (zooming out)

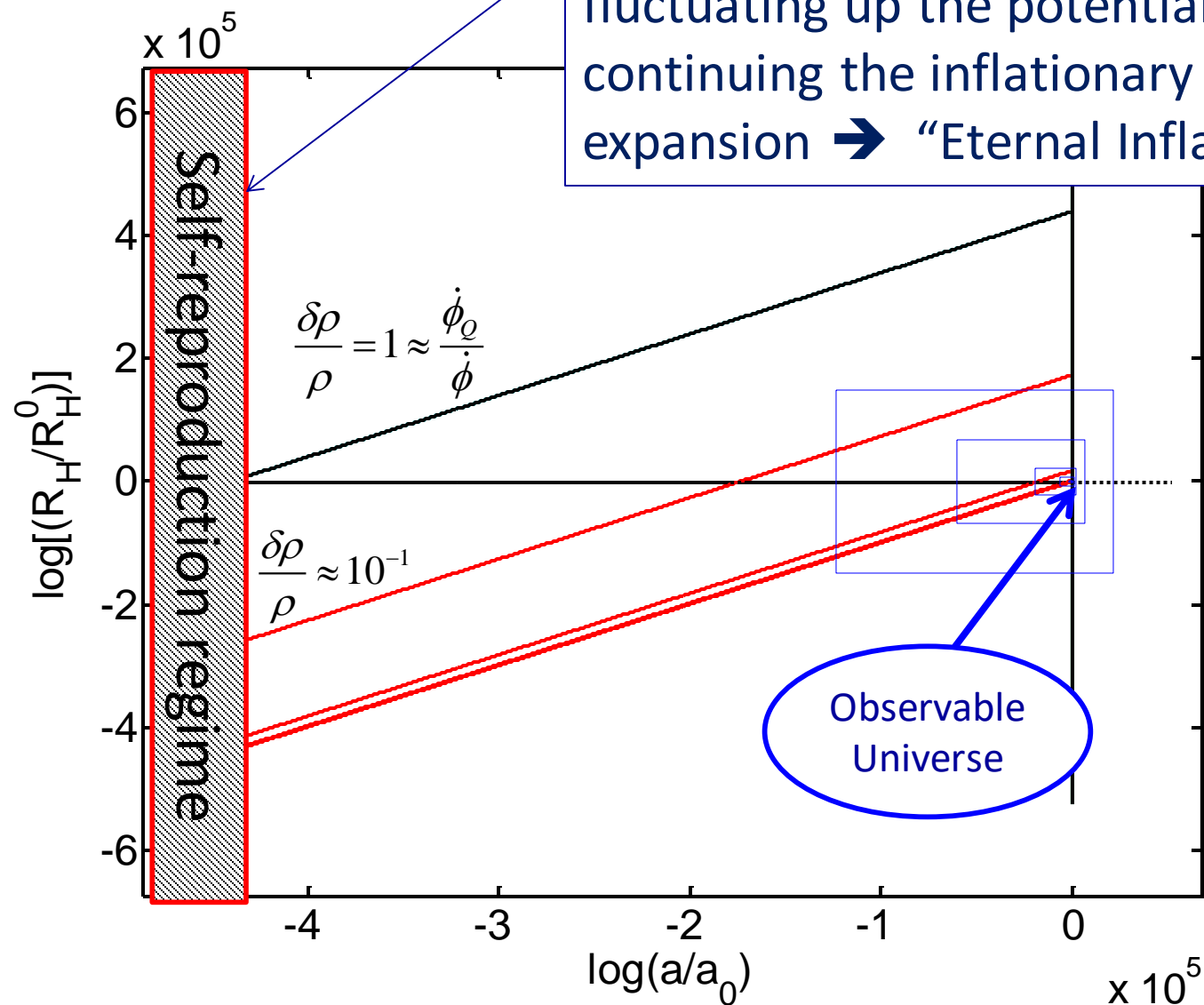


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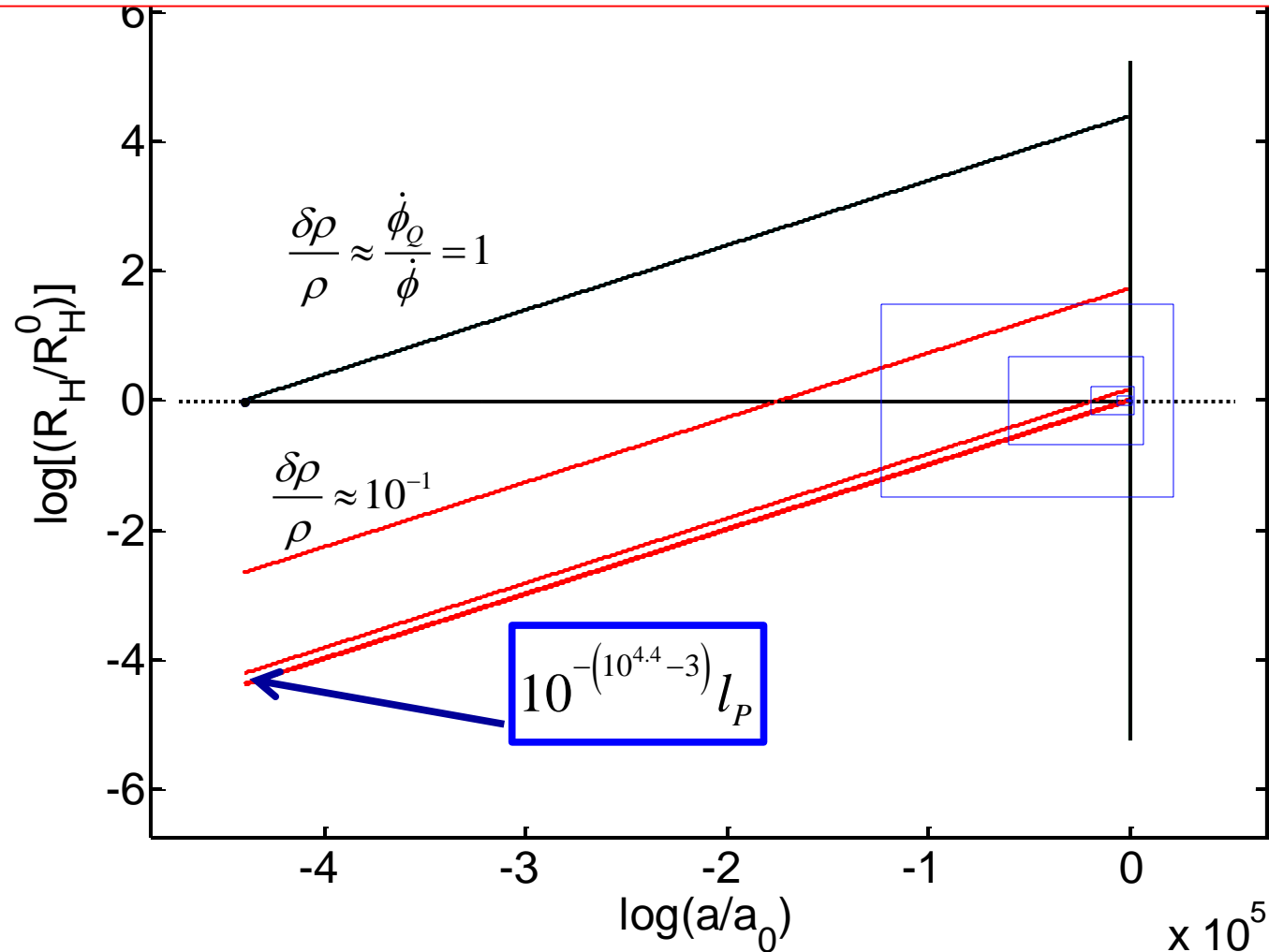


Steinhardt 1982, Linde 1982, Vilenkin 1983, and (then) many others



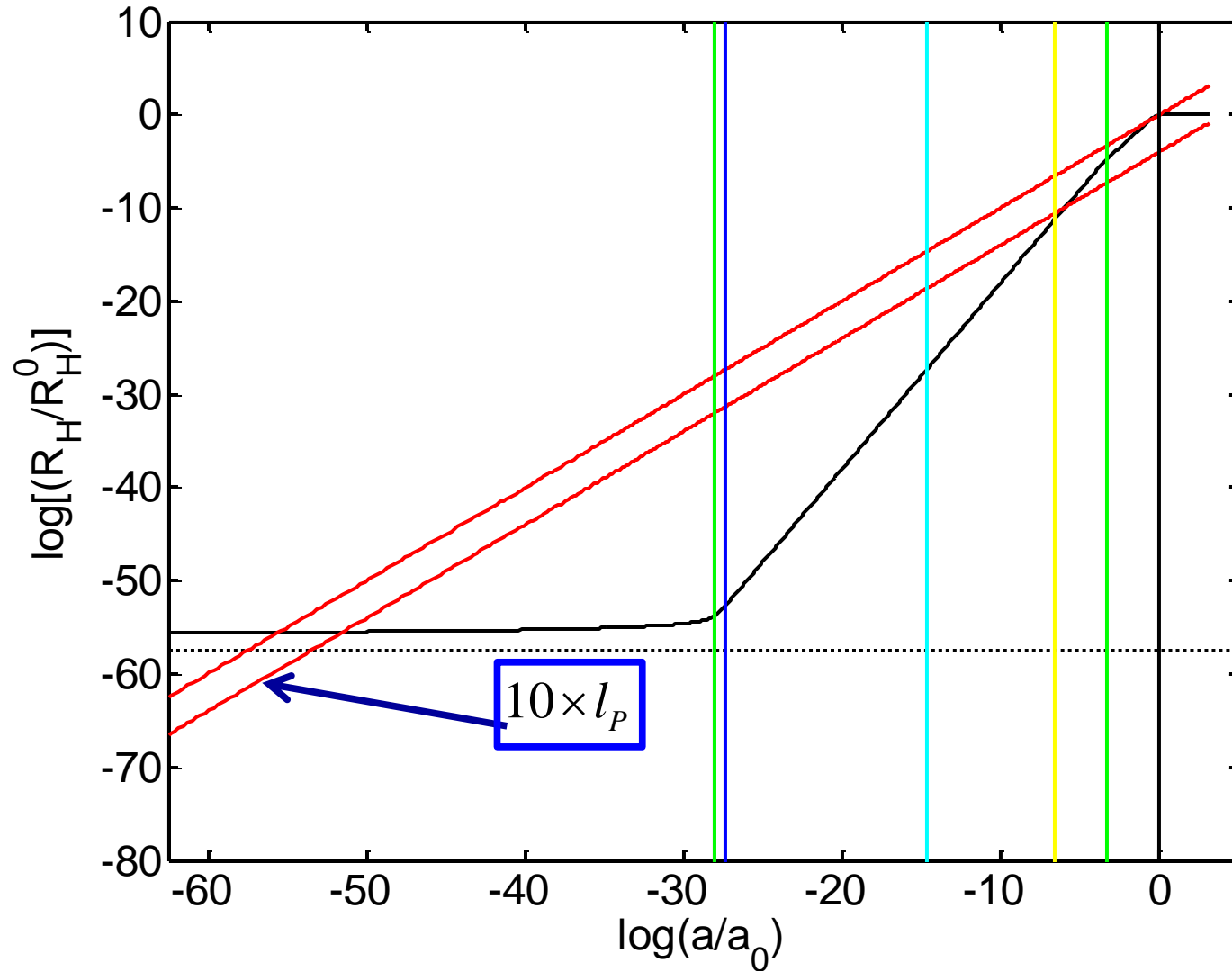
Substantial probability of fluctuating up the potential and continuing the inflationary expansion → “Eternal Inflation”

At end of self-reproduction our observable length scales were exponentially below the Plank length (and much smaller than that \*during\* self-reproduction)!

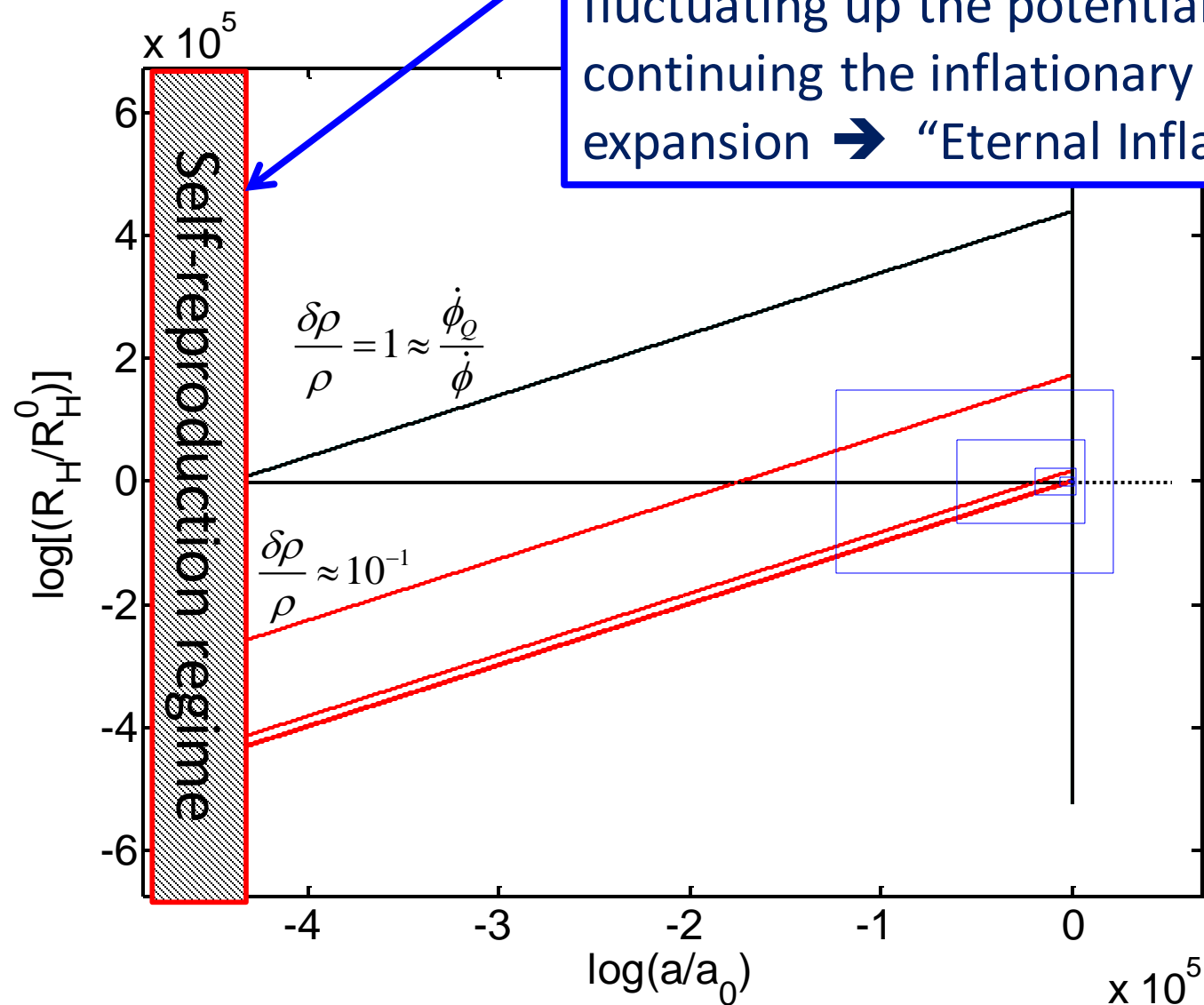


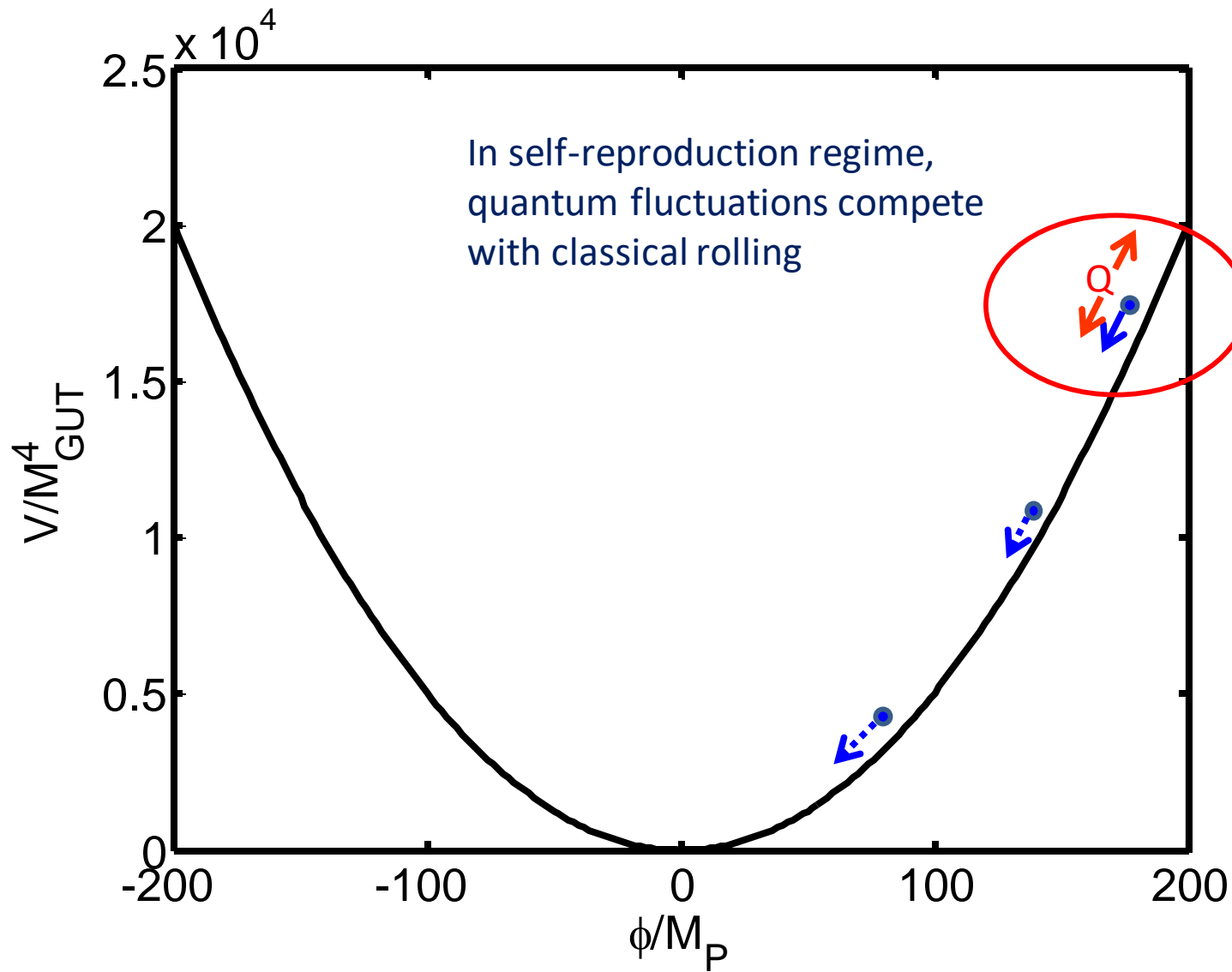


At “formation” (Hubble length crossing) observable scales were just above the Planck length



(Bunch Davies Vacuum)







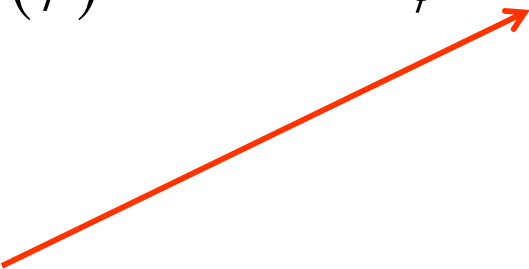
Self-reproduction is a generic feature of almost any inflaton potential:

During inflation

$$\ddot{\phi} + 3H\dot{\phi} = -\Gamma_{\phi}\dot{\phi} - V'(\phi)$$


$$3H\dot{\phi} \approx -V'(\phi)$$



$$\dot{\phi} \approx \frac{-V'(\phi)}{3H}$$



$$\frac{\dot{\phi}_Q}{\dot{\phi}} = \frac{H^2}{\dot{\phi}} \approx \frac{H^3}{-V'(\phi)} \propto \frac{V^{3/2}}{V'}$$

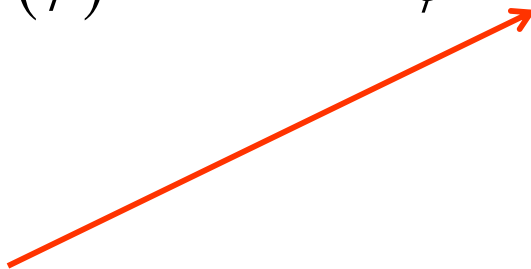
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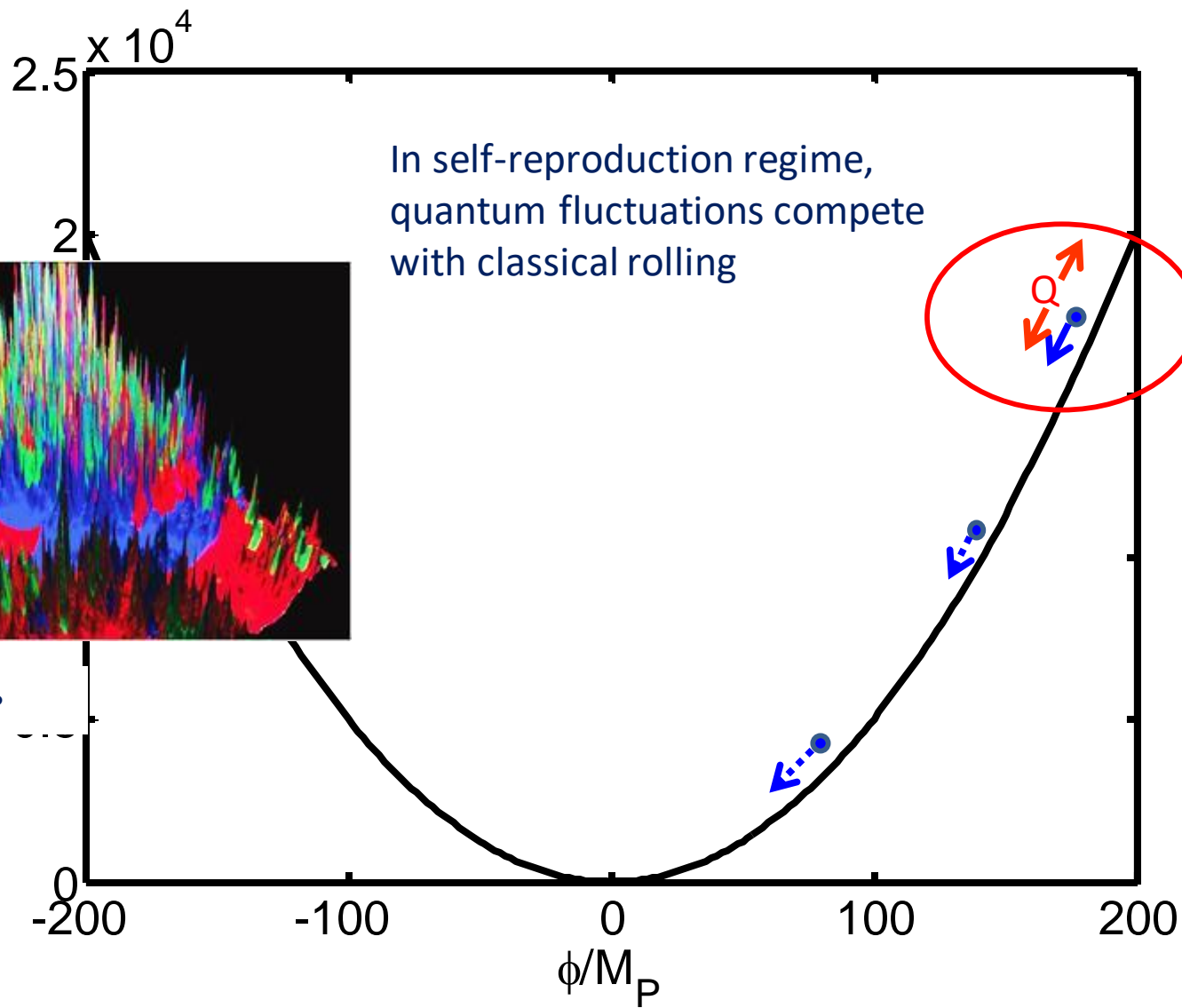
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$\geq 1$  for self-reproduction




$$d \approx 5R_H^S$$

Self-reproduction regime




Classically  
Rolling



$t = 0$

$t$


$$d \approx 5R_H^S$$

*Self-reproduction regime*



*Classically  
Rolling*

NB: shifting focus to  $l(t)$



$t = 0$

$t$



$$d \approx e^2 \times 5R_H^S$$

Self-reproduction regime

Classically  
Rolling

New pocket (elsewhere)

$$t = 2R_H^S / c$$

$t$

$$d \approx e^3 \times 5R_H^S$$

Self-reproduction regime

Classically  
Rolling

New pocket (elsewhere)

$$t = 3R_H^S / c$$

$t$

$$d \approx e^{500} \times 5R_H^S$$

Self-reproduction regime

Classically  
Rolling

New pocket (elsewhere)

$$r \approx e^{-502} d$$

$$t = 500R_H^S / c$$



$$d \approx e^{1000} \times 5R_H^S$$

Self-reproduction regime



New pocket (elsewhere)

$$r \approx e^{-1002} d$$

$$t = 1000R_H^S / c$$



$$d \approx e^{1400} \times 5R_H^S$$

Self-reproduction regime

New pocket (elsewhere)

$$r \approx e^{-1402} d$$

$$t = 1400R_H^S / c$$



$$d \approx e^{1395} \times 5R_H^S$$

Self-reproduction regime

Classically  
Rolling

New pocket (elsewhere)

$$r \approx e^{-1393} d$$

$$t = 1400R_H^S / c$$



$$d \approx e^{1991} \times 5R_H^S$$

Self-reproduction regime

Classically  
Rolling

New pocket (elsewhere)

$$r \approx e^{-1989} d$$

$$t = 2000R_H^S / c$$



$$d \approx e^{534395} \times 5R_H^S \equiv R_H^{lend}$$

Self-reproduction regime

Classically  
Rolling

New pocket (elsewhere)

$$r \approx e^{-534393} d$$

$$t = (602,785) R_H^S / c$$





$$d \approx e^{534395} \times 5R_H^S \equiv R_H^{lend}$$

Self-reproduction regime

● ← Reheating

New pocket (elsewhere)

$$r \approx e^{-534393} d$$

$$t = 2R_H^{lend} / c$$



$$d \approx e^{534395} \times 5R_H^S \equiv R_H^{lend}$$

Self-reproduction regime

•  $\nwarrow$  Radiation  
Era

New pocket (elsewhere)

$$r \approx e^{-534393} d$$

$$t = 3.2 R_H^{lend} / c$$



## Eternal inflation features

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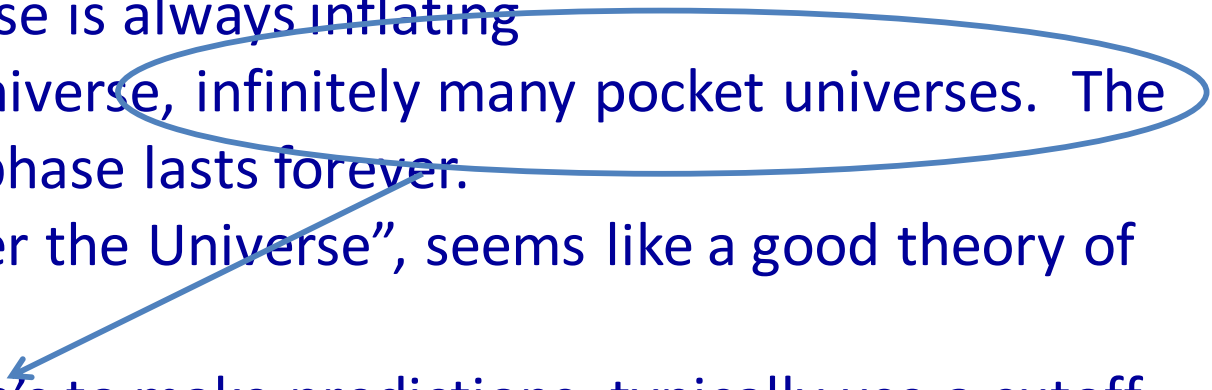
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- End of time problem
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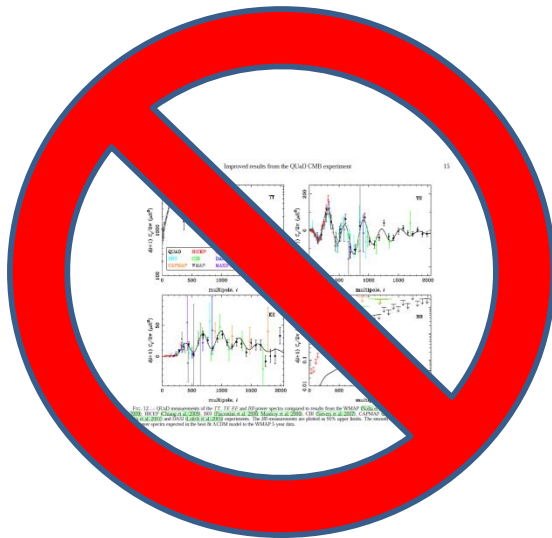
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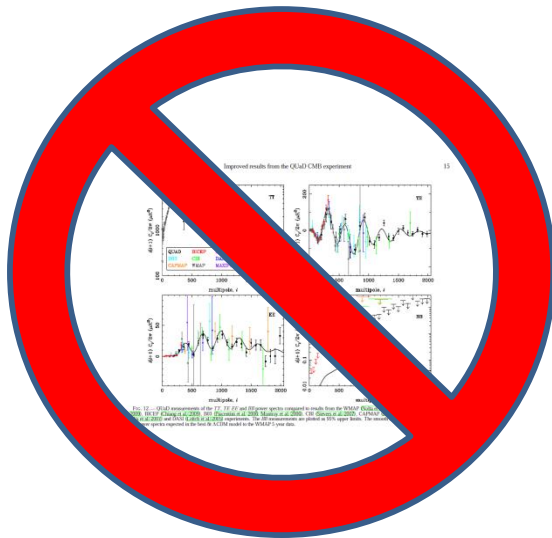
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- Most of the Universe is always inflating
- Leads to infinite Universe, infinitely many pocket universes. The self-renewal of the Universe
- In a good theory of quantum gravity, anything that can happen will happen infinitely many times! (A. Guth)
- No one can actually use a cutoff.
- For a specific region, exponentially produced pocket universes are exponentially favored (produced in an exponentially larger region).

“Anything that can happen will happen infinitely many times! (A. Guth)”

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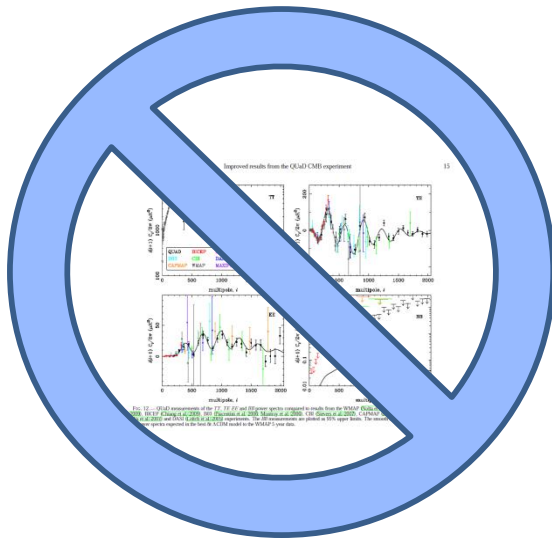
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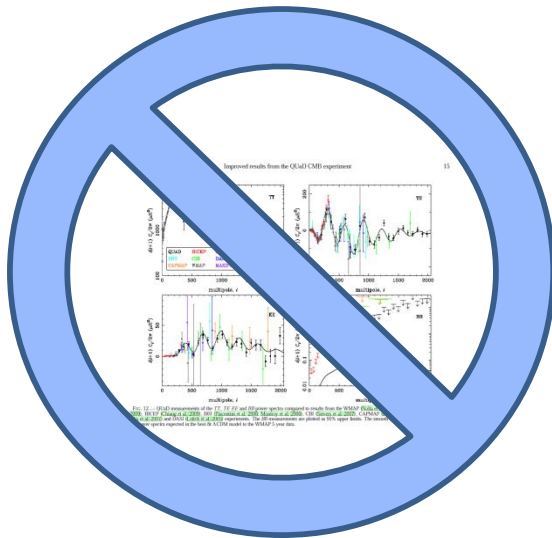
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Hernley, AA & Dray

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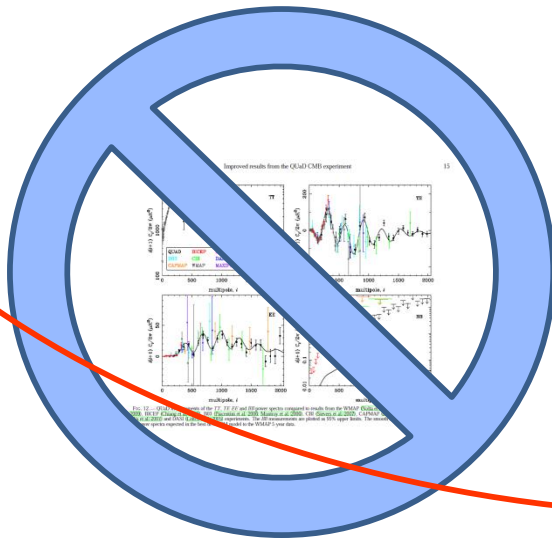
Multiply by  $10^{500}$  to get landscape story!

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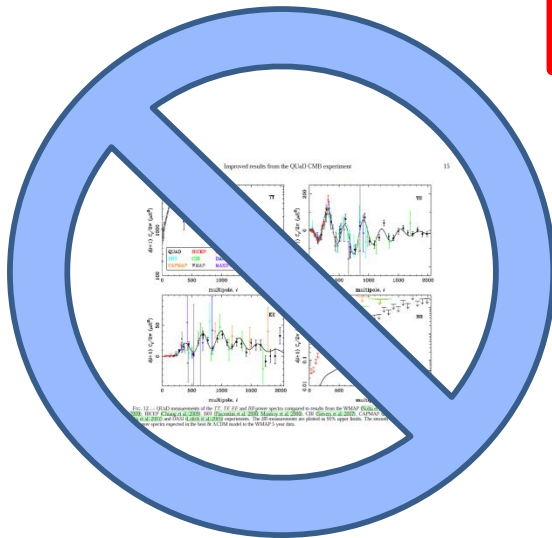
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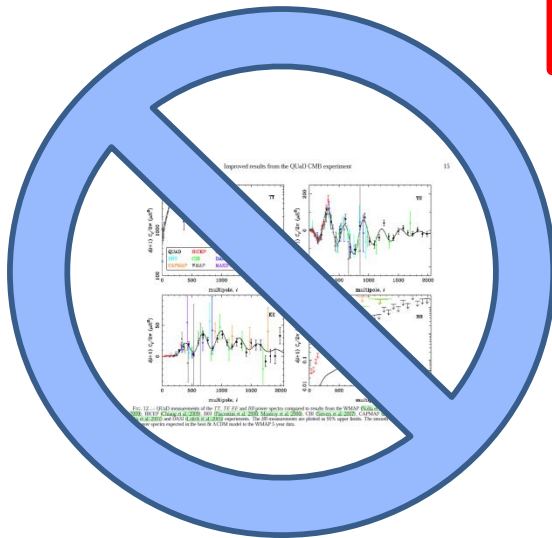
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Be careful about counting arguments. Many are just wrong!

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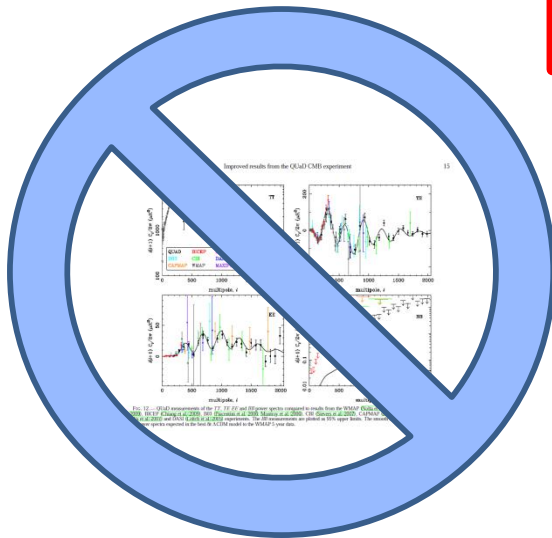
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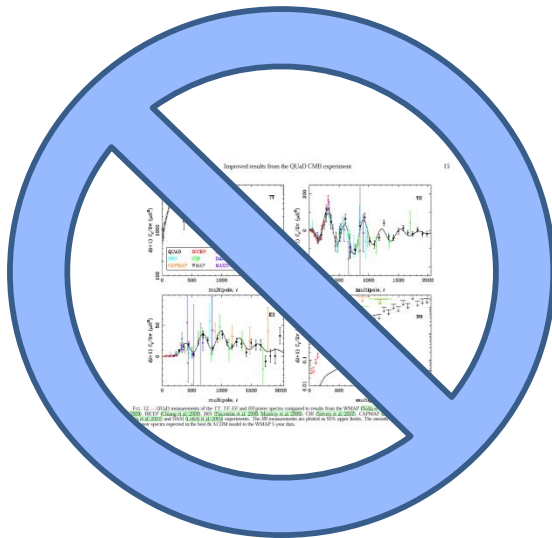


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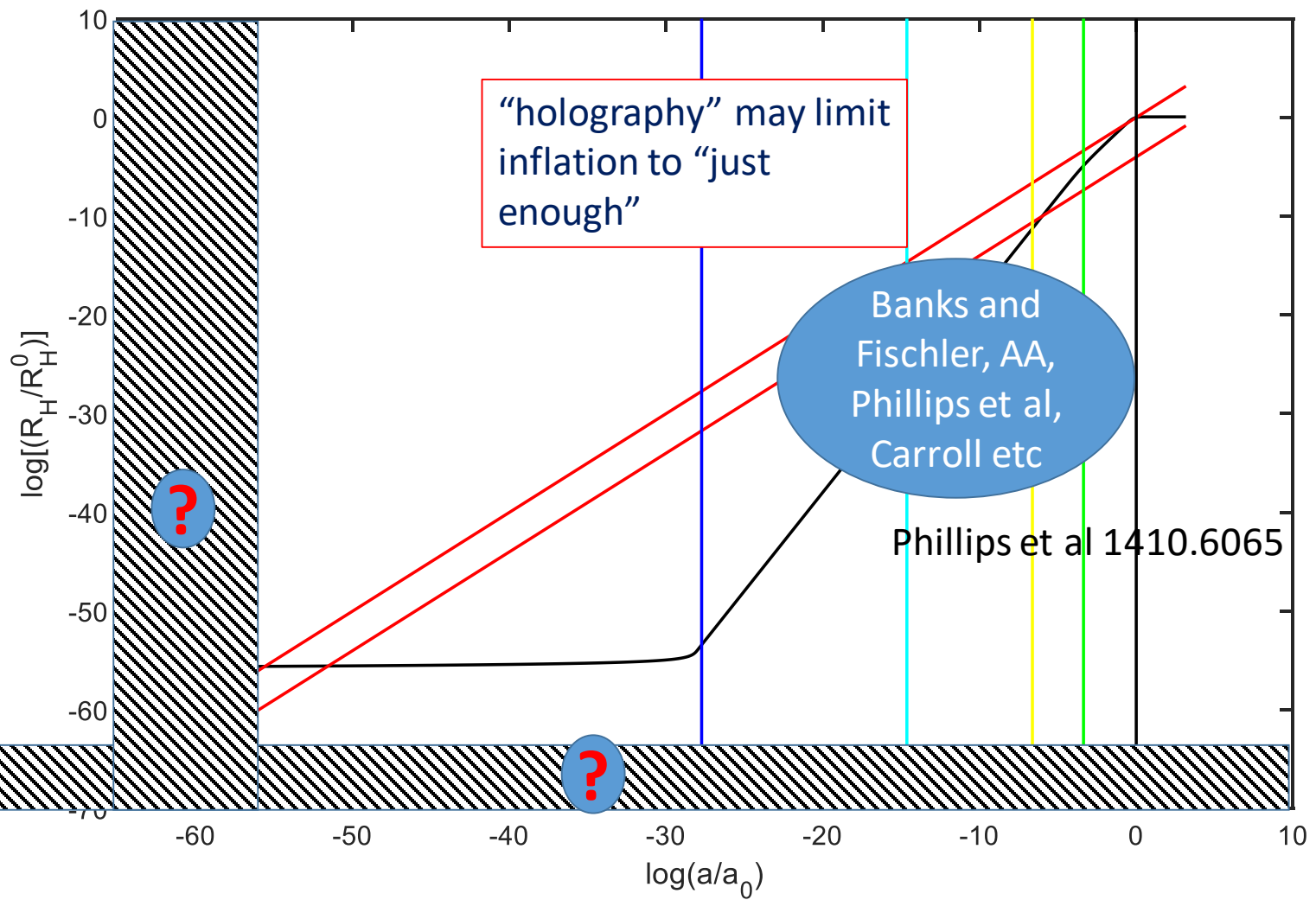
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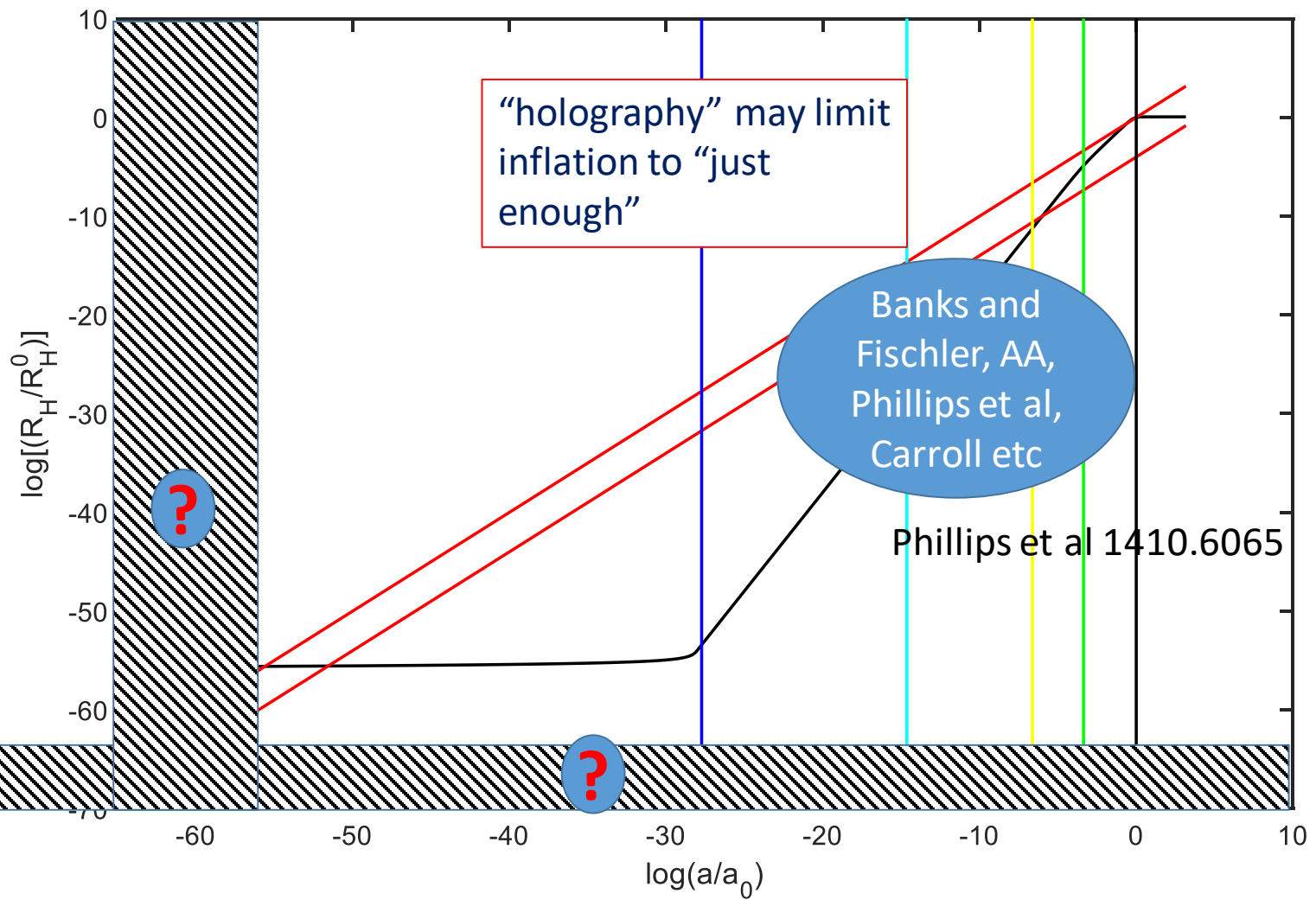
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2. Inflation basics
3. Eternal inflation
4. An alternative to Eternal Inflation
5. Further thoughts

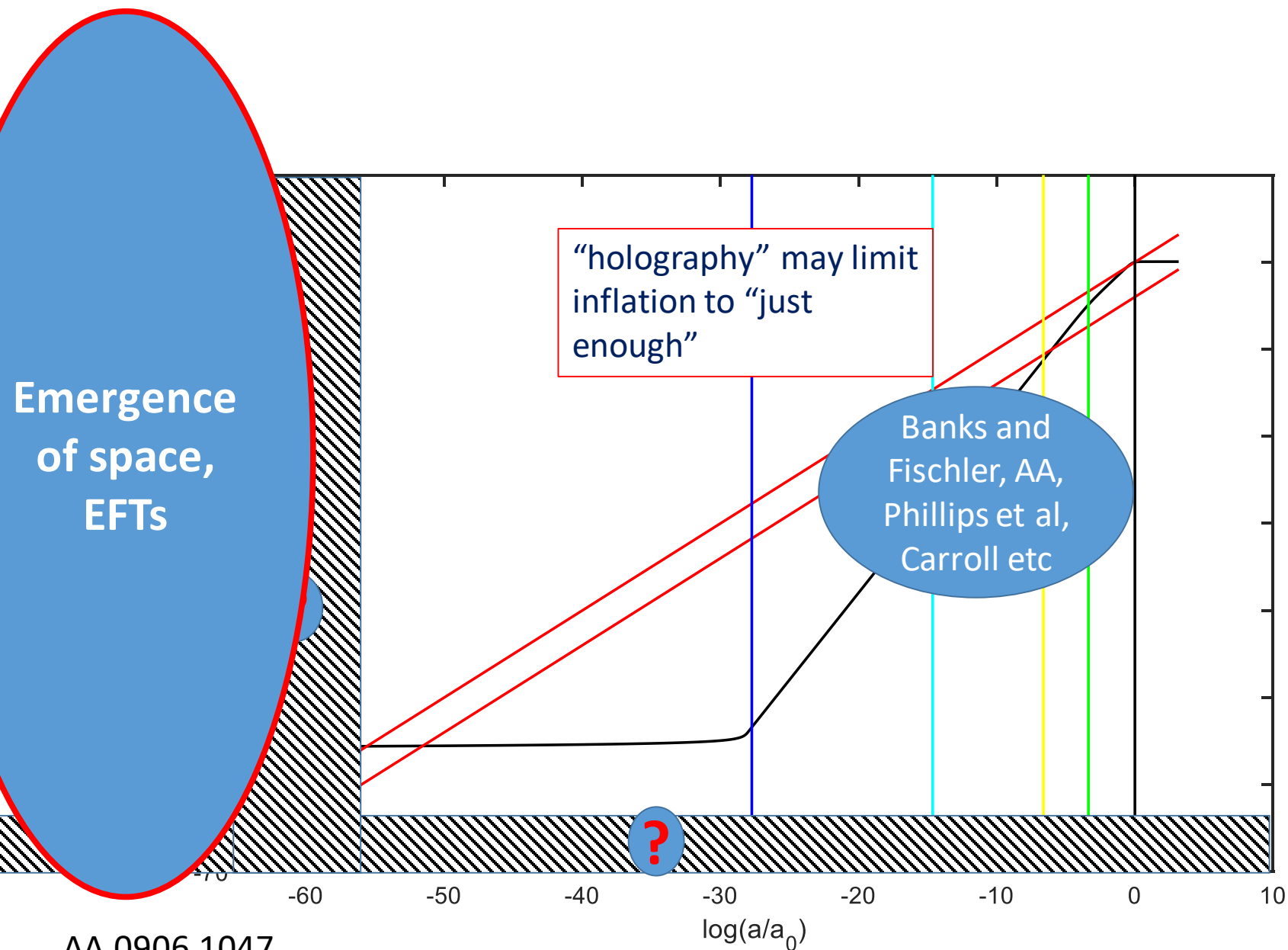
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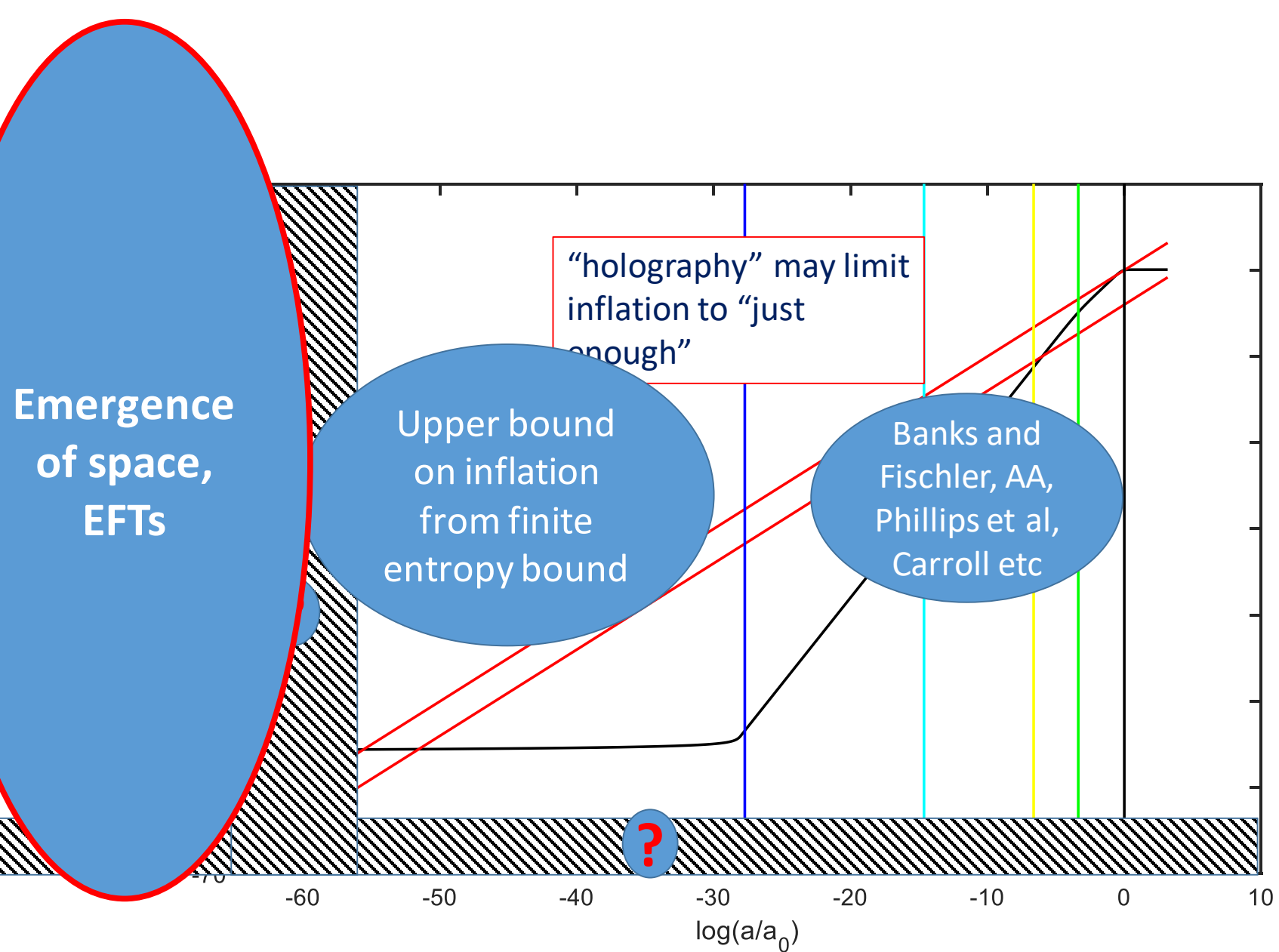


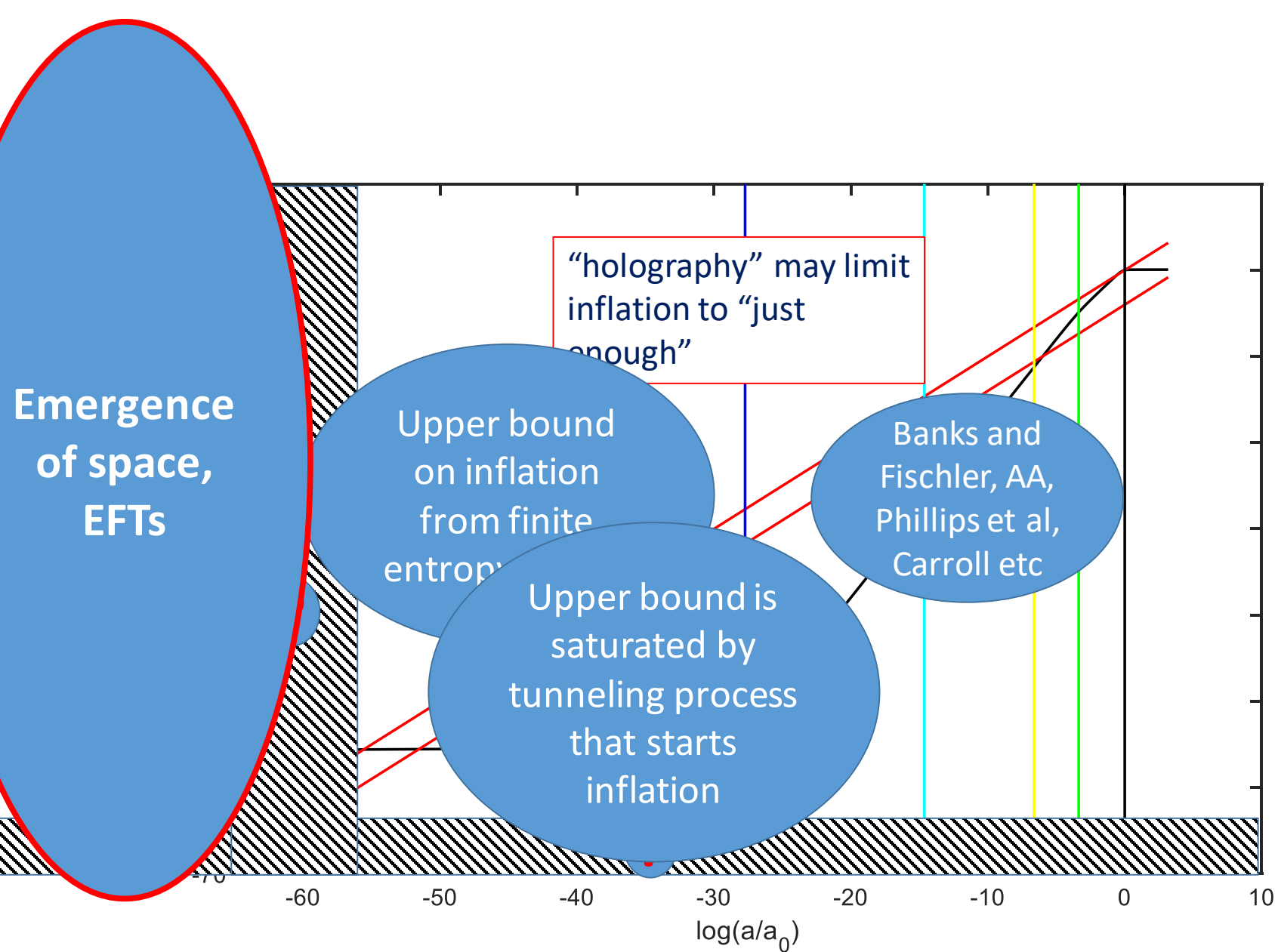


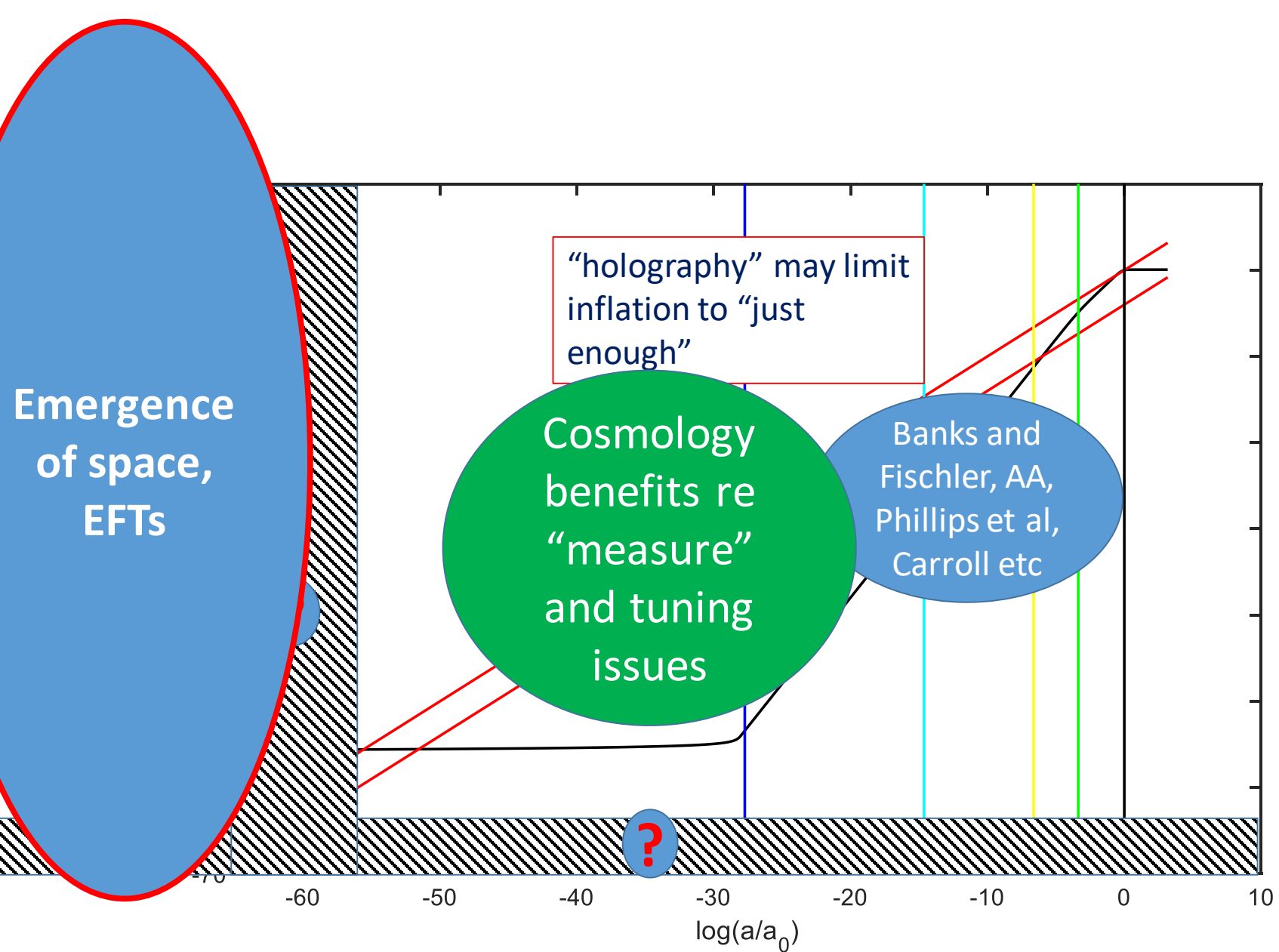
AA 0906.1047

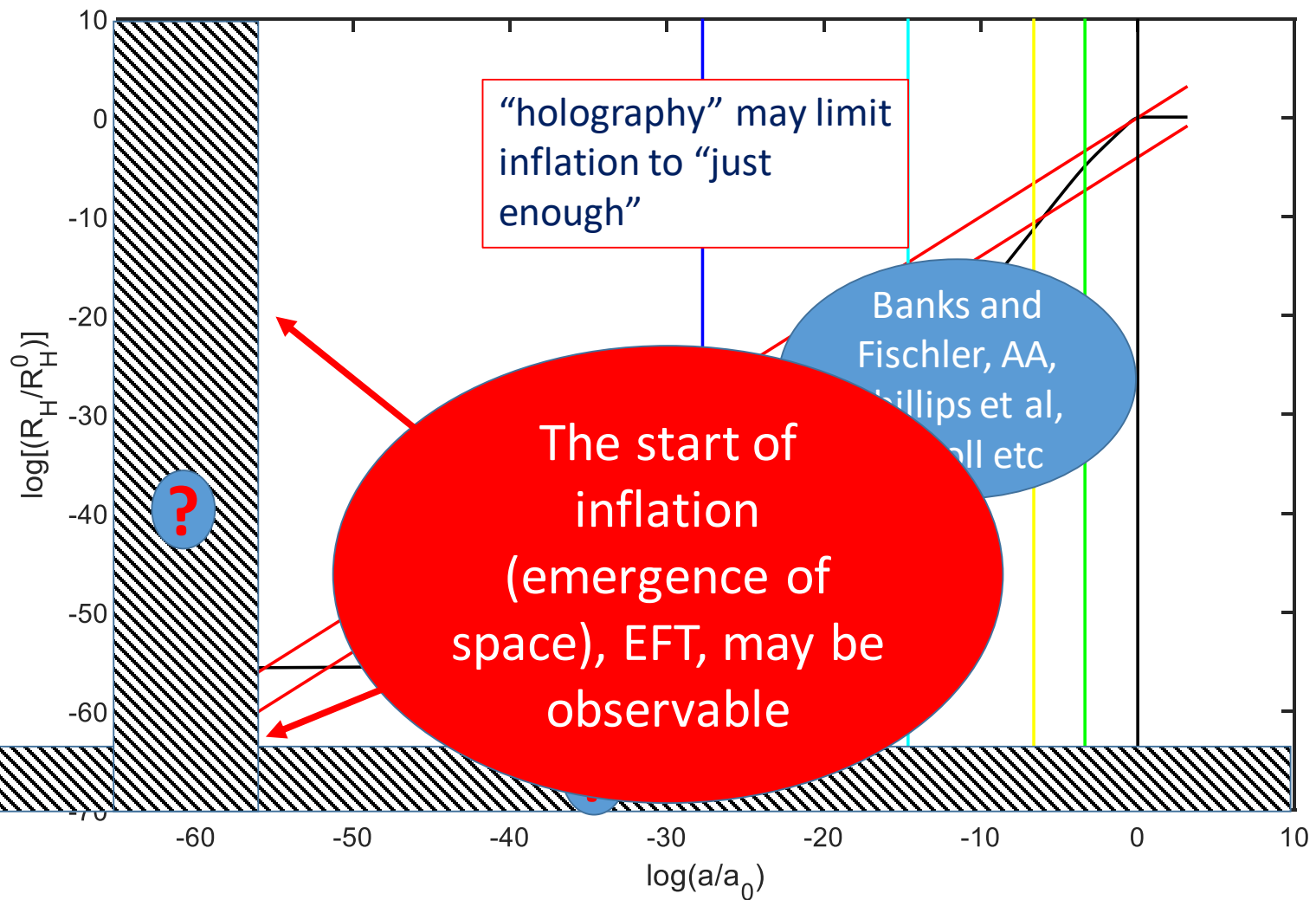
AA and Sorbo, hep-th/0405270

AA 1401.7309

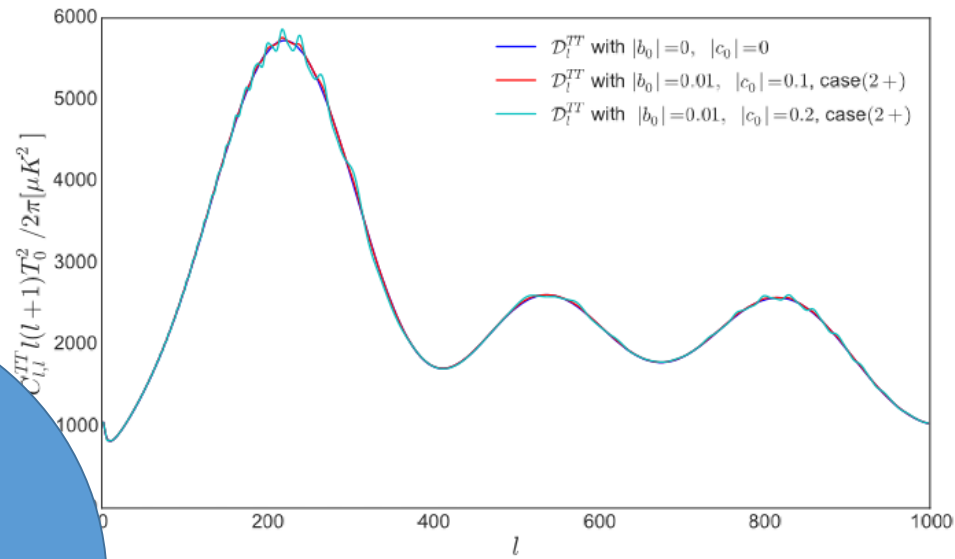




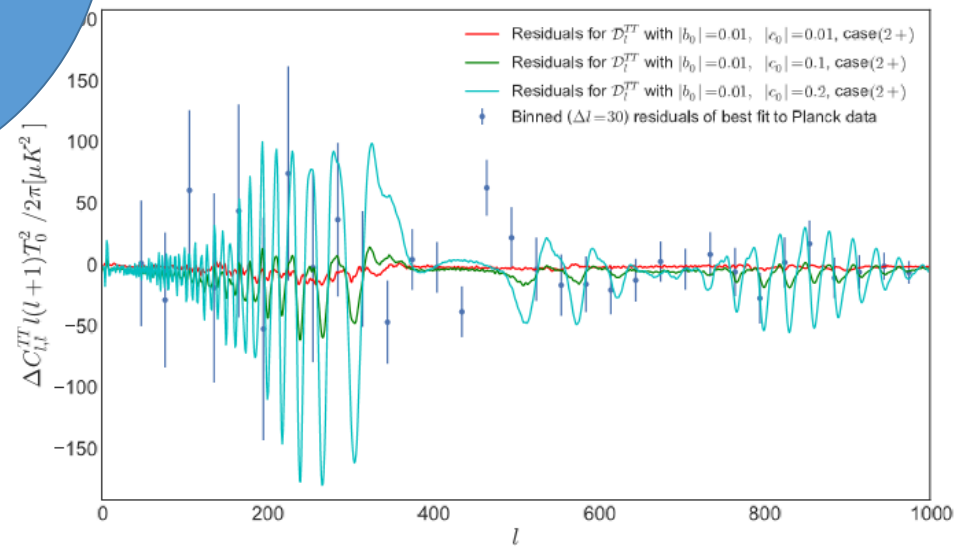




Possible  
signatures of  
“short  
inflation” and  
quantum  
entanglement



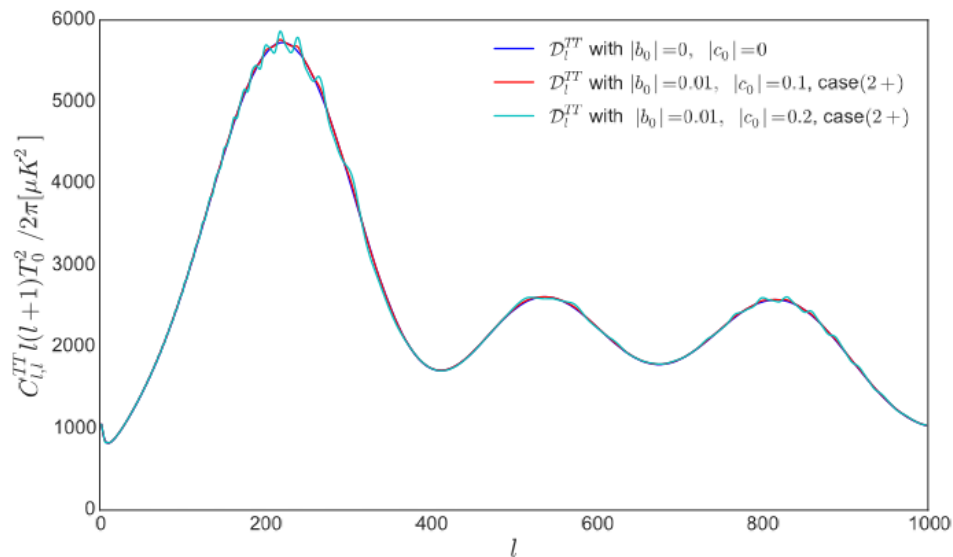
(a)



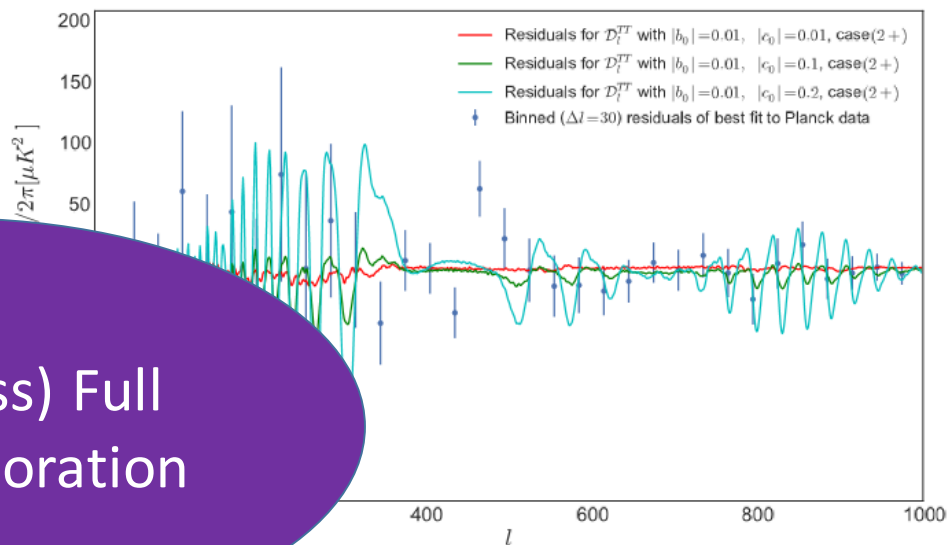
(b)

[Bolis et al](#)  
[arXiv:1605.01008](#)

**Figure 1.** (a) Temperature fluctuation angular power spectrum  $C_l^{TT}$  for different values of entanglement parameter  $|\tilde{C}_{k0}^+|$  ( $|c_0|$  on plot to simplify labeling), keeping  $|\tilde{b}_{k3}(\tau_0)|$  ( $|b_0|$  on plot) constant for



(a)



(In progress) Full  
MCMC exploration

With Andrew Arrasmith

Figure 1. (a) Temperature fluctuation angular power spectrum  $C_l^{TT}$  for different values of entangle-

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Related to some more contentious topics:

- How well does inflation “solve” tuning problems?
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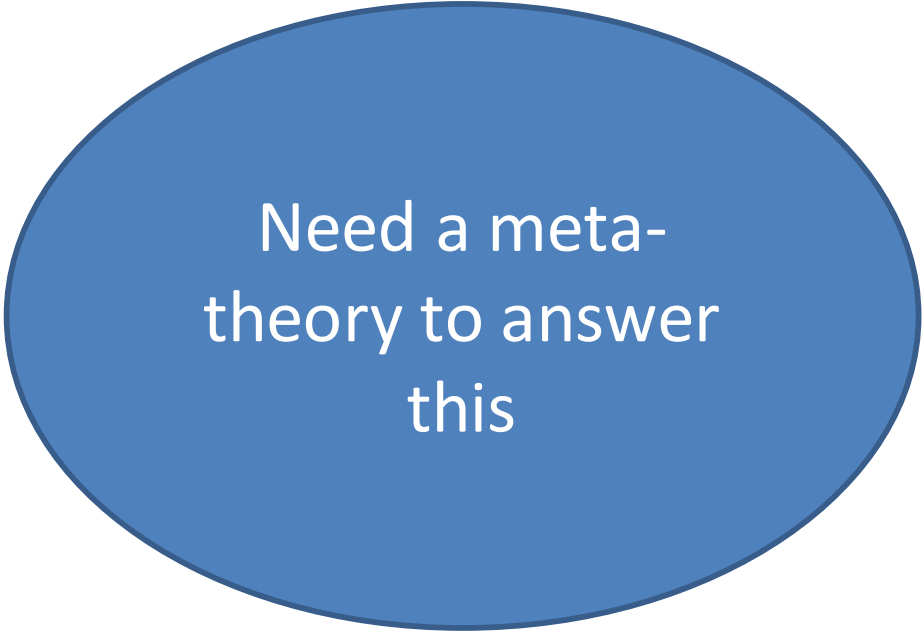
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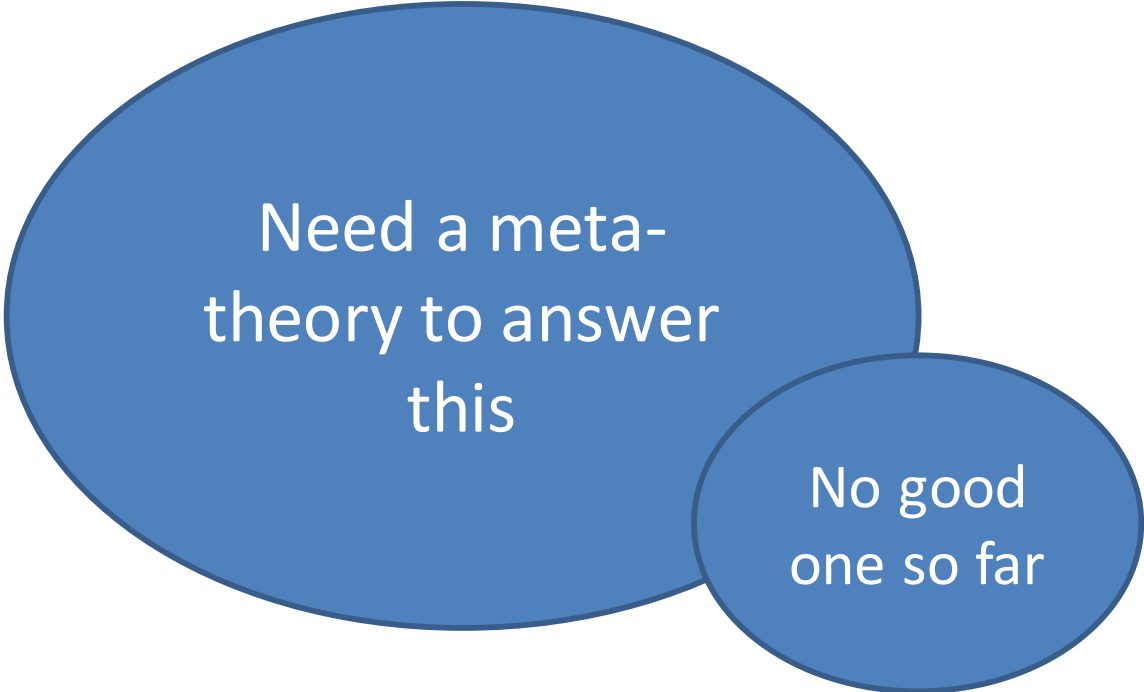


Need a meta-  
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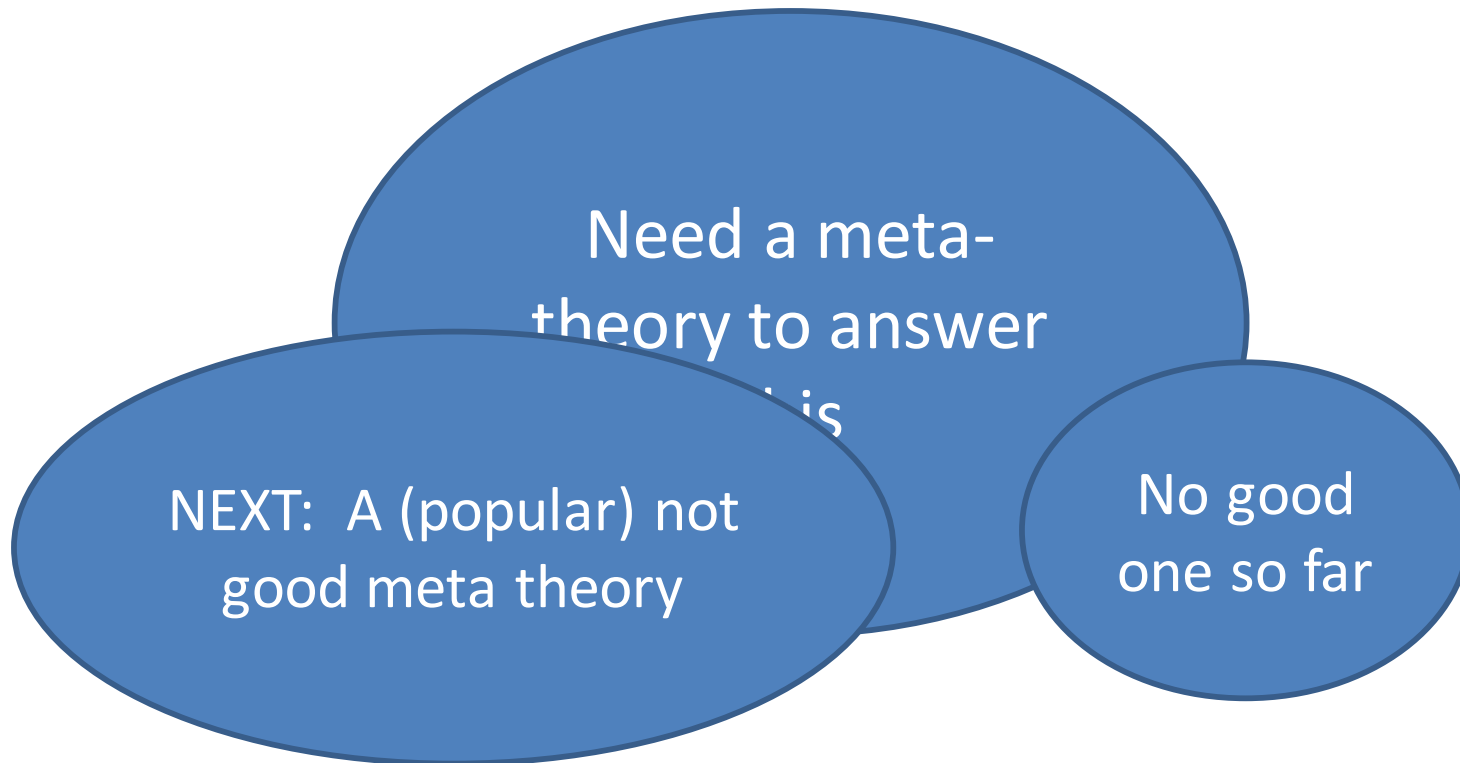
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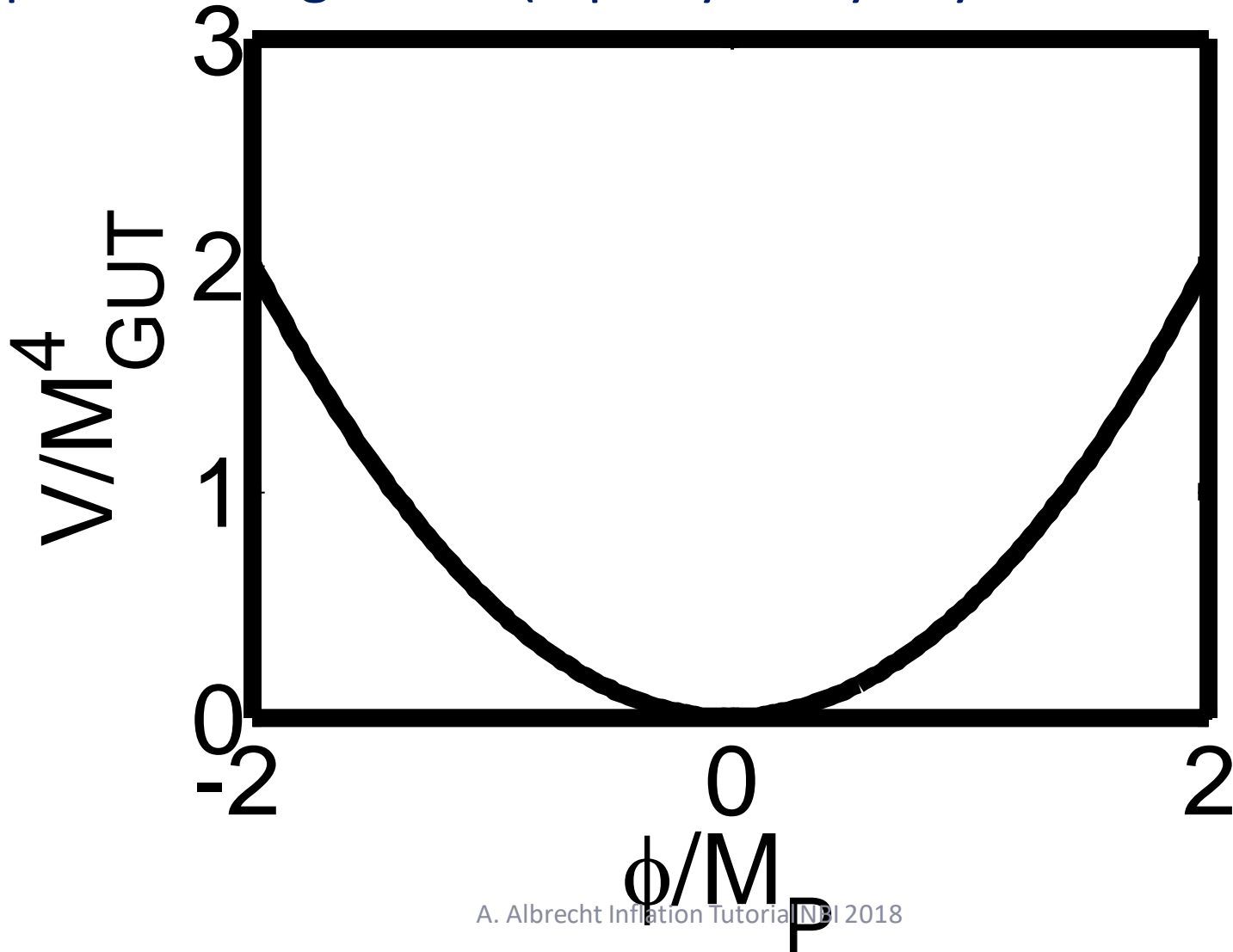
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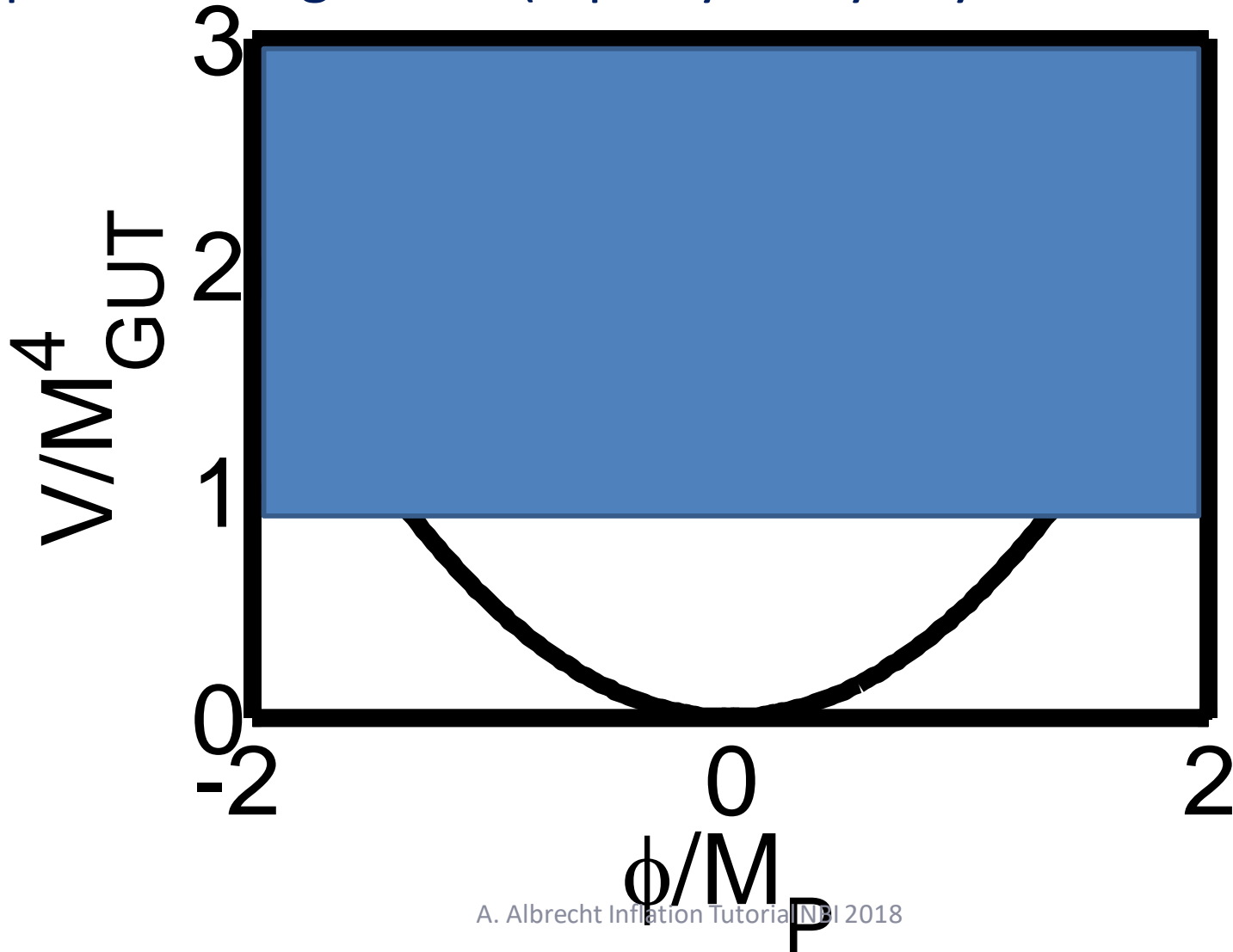
# Beware “temporal provincialism”:

Equipartition argument (equally likely anywhere on potential)



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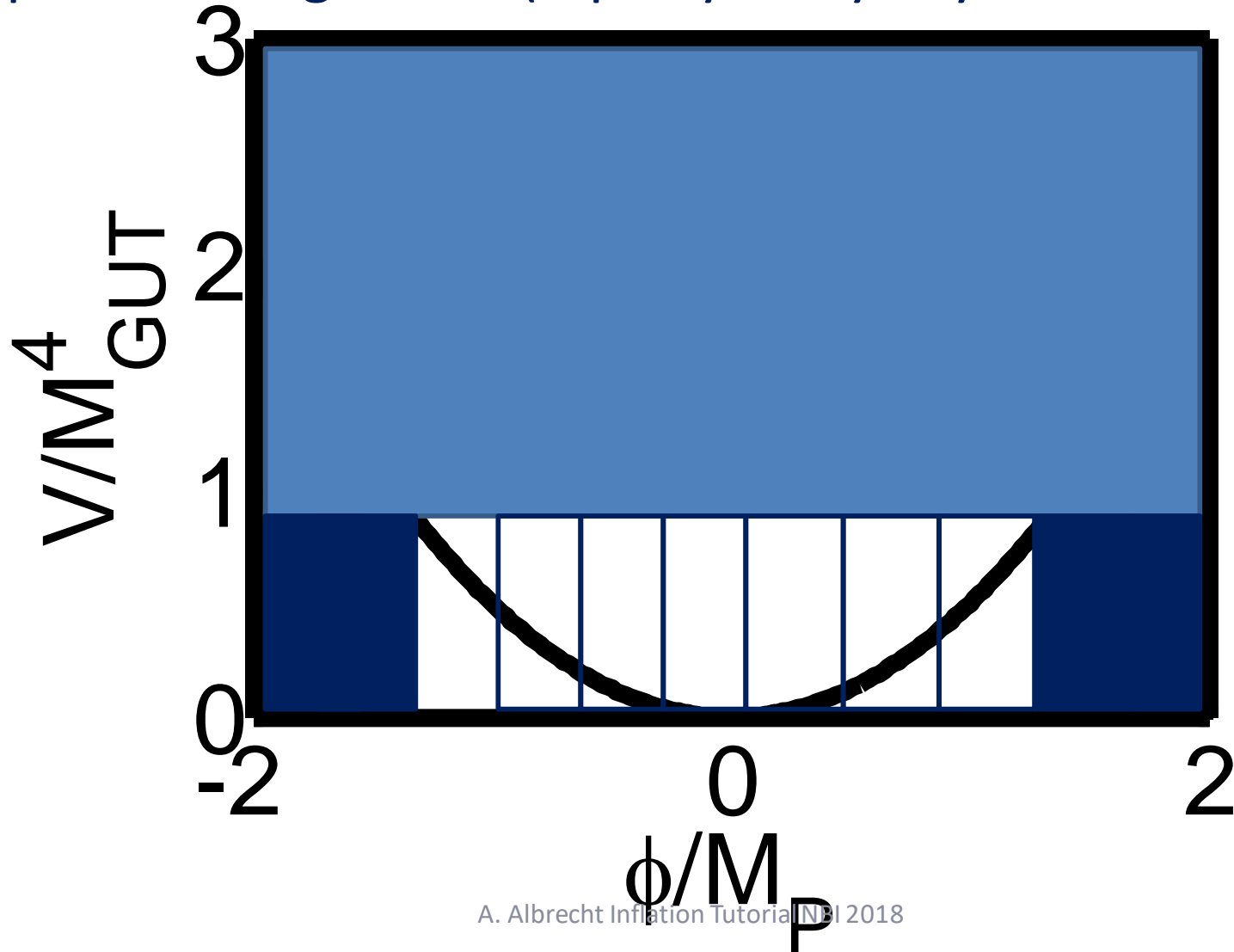
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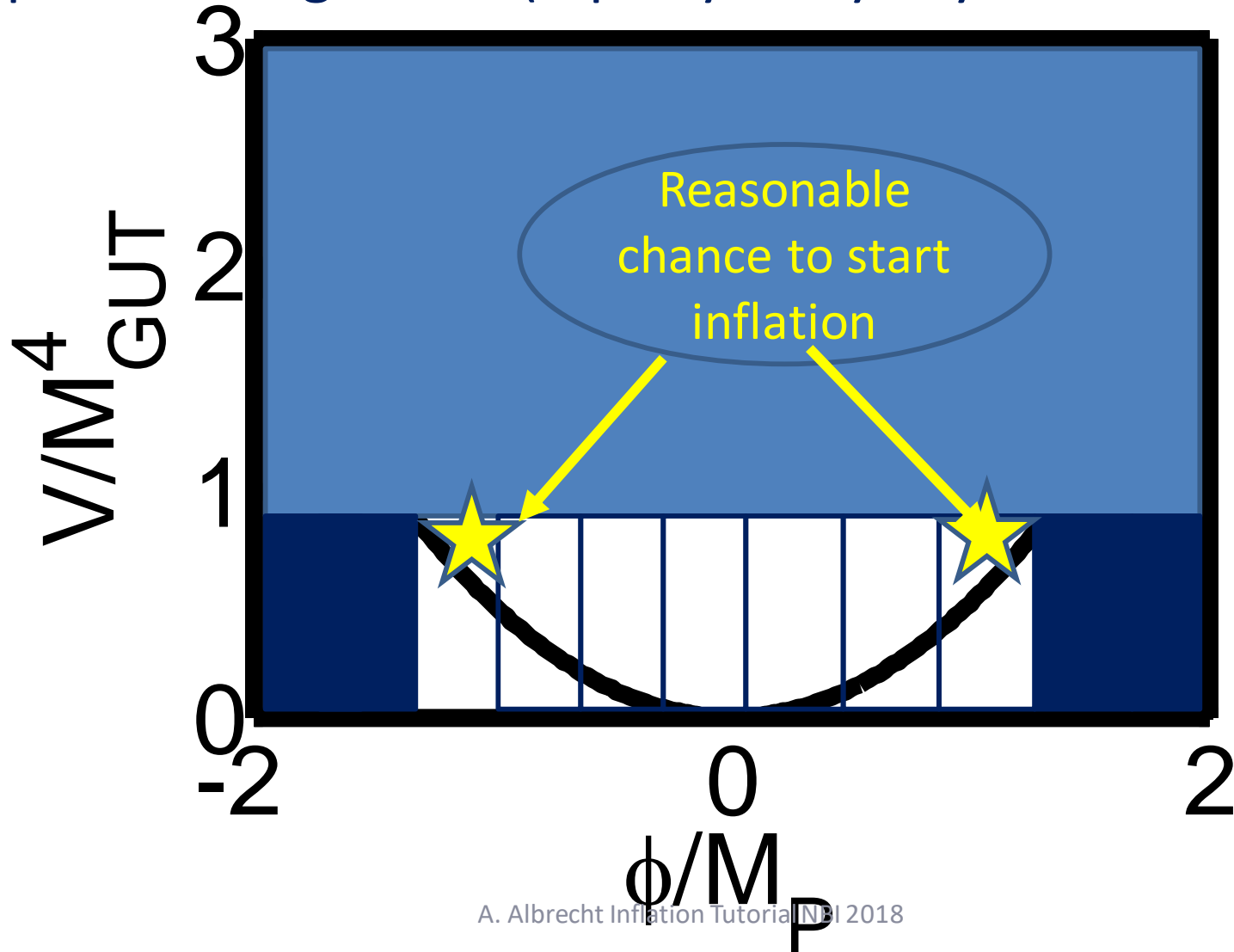
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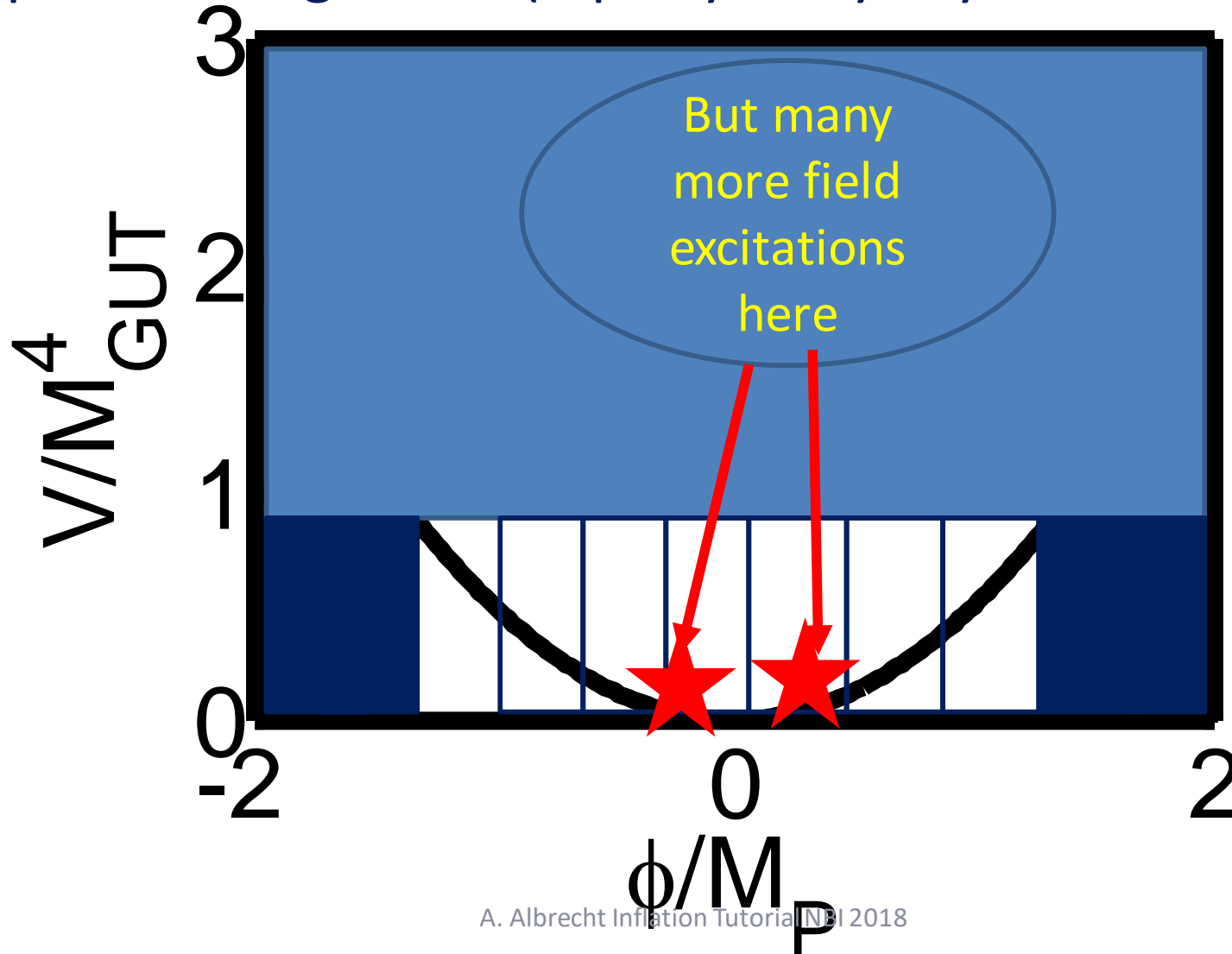
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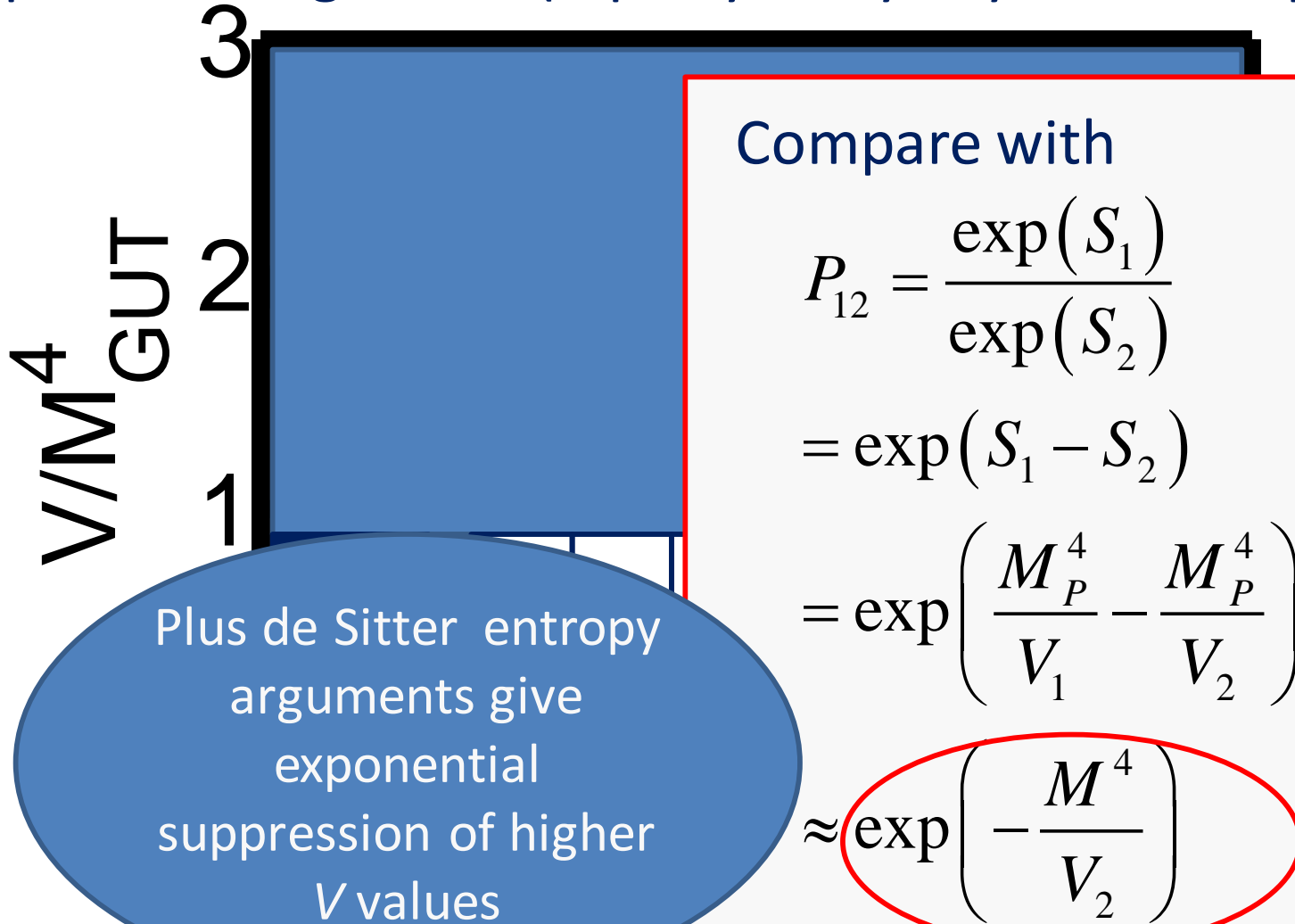
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3

Compare with  
numerical  
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initial conditions  
problem

Entropy arguments  
give exponential  
suppression of higher  
 $V$  values

Compare with

$$P_{12} = \frac{\exp(S_1)}{\exp(S_2)}$$

$$= \frac{\text{East et al } 1511.05143}{\text{Braden et al } 1604.04001}$$
$$= \frac{\text{Clough et al } 1608.04408}{\text{Clough et al } 1608.04408}$$

$$= \exp\left(\frac{M_P^4}{V_1} - \frac{M_P^4}{V_2}\right)$$

$$\approx \exp\left(-\frac{M^4}{V_2}\right)$$

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All are temporally provincial in some way

Compare with numerical “solution” of initial conditions problem

Entropy arguments give exponential suppression of higher  $V$  values

$1.05143$   
 $=$  Braden et al 1604.04001  
 Clough et al 1608.04408

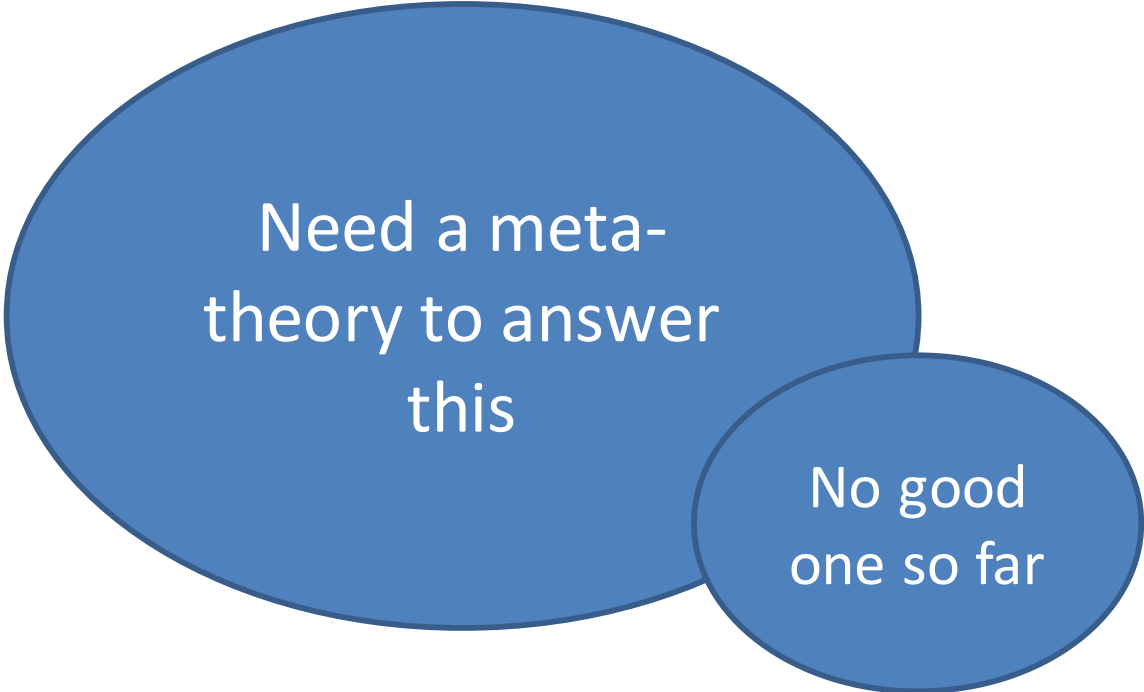
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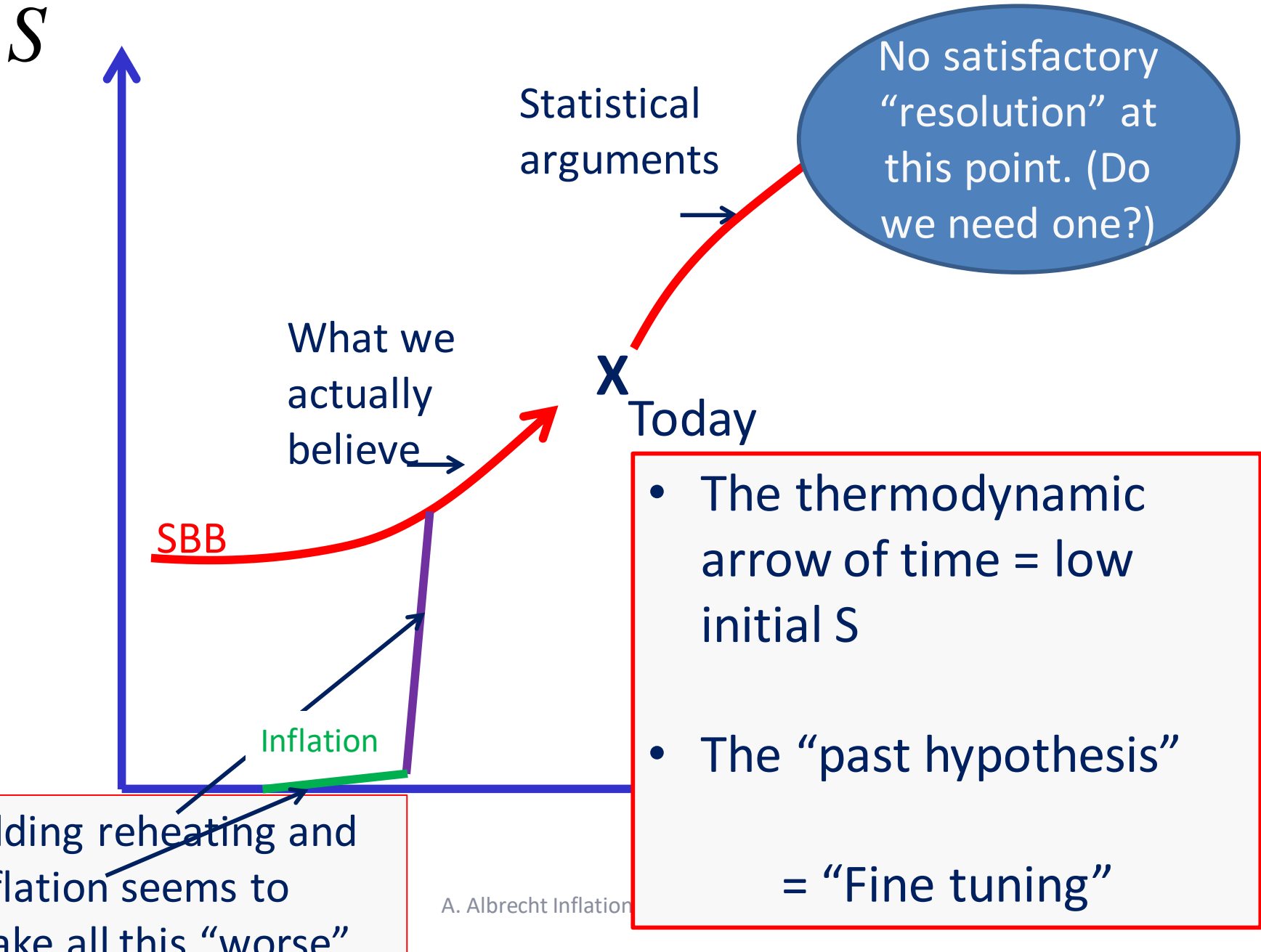
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Not clear whether a  
good meta theory  
would favor inflation  
over just the SBB with  
correct initial conditions

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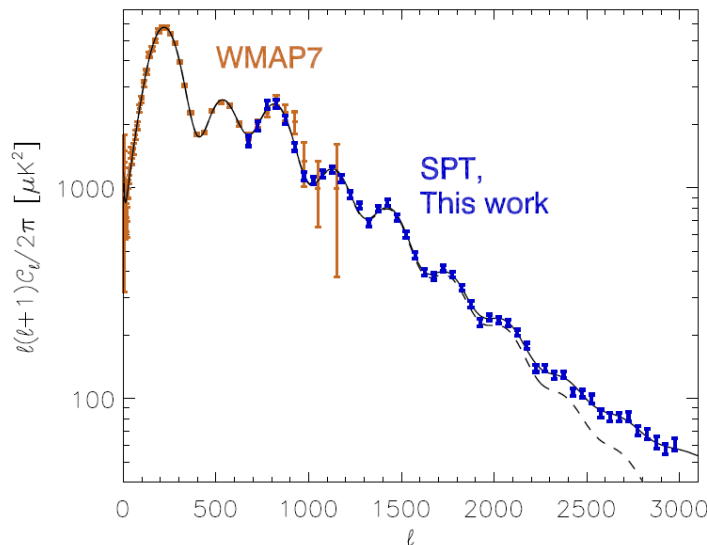




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“True infinity” needed here



Or, just be happy we have equations to solve?

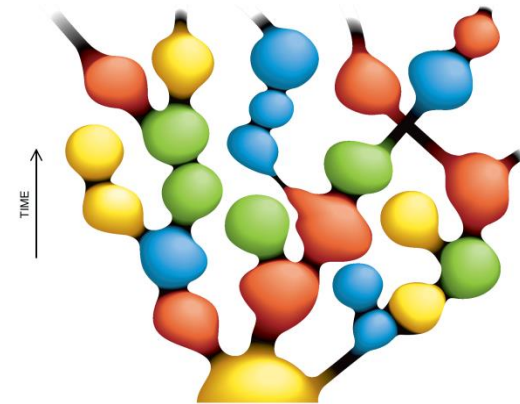
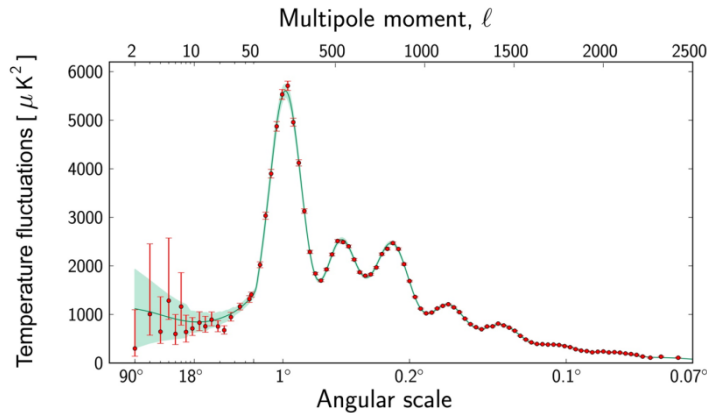
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# Cosmic Inflation:

Consumers

&

Producers



The multiverse of eternal inflation with multiple classical rolling directions

Self-reproduction regime

Classically Rolling A

Classically Rolling B

Classically Rolling C

Where are we? (Young universe, old universe, curvature, physical properties A, B, C, D, etc)

Classically Rolling D

adapted from: [reference]

10



# Conclusions

- A) Inflation is a technical tool for connecting cosmological observables with high energy physics. Impressive successes.
- B) However, without a meta theory about how inflation started (and how it “competes” with other scenarios, such as the SBB) big questions are unresolved.
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My priors