

Tau lepton asymmetry by sterile neutrino emission – Moving beyond one-zone supernova model

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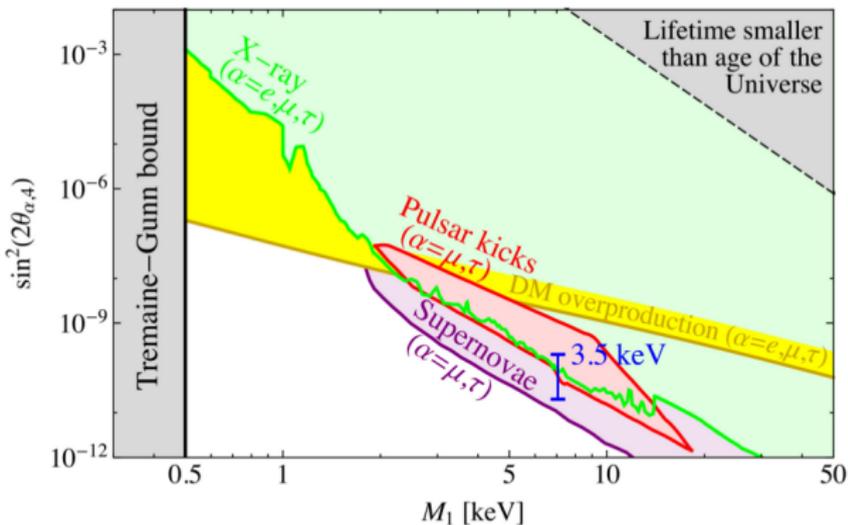


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- ① Motivation
- ② Sterile neutrino conversions in the stellar core
- ③ Development of the neutrino lepton asymmetry
- ④ Conclusions

Motivation

Sterile neutrino as Dark Matter candidate



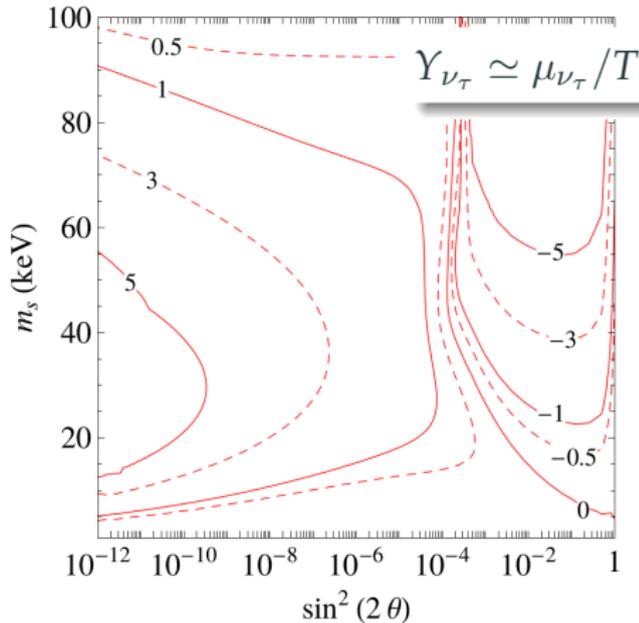
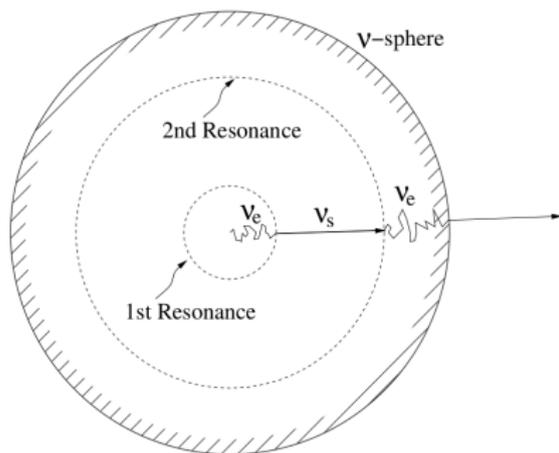
Favorable regions

- Pulsar kicks
(A. Kusenko (2004))
- 3.5 keV line
(A. Boyarsky et al. (2014))

Constraints

- Supernovae energy bounds (X. Shi & G. Sigl (1994))
- DM overproduction (S. Dodelson, L. M. Widrow (1994), X. Shi, G. M. Fuller (1999))
- Radiative decay (NuSTAR, XMM, Chandra)
- Tremaine-Gunn bound

The role of sterile neutrinos in SNe



J. Hidaka and G. M. Fuller (2006)

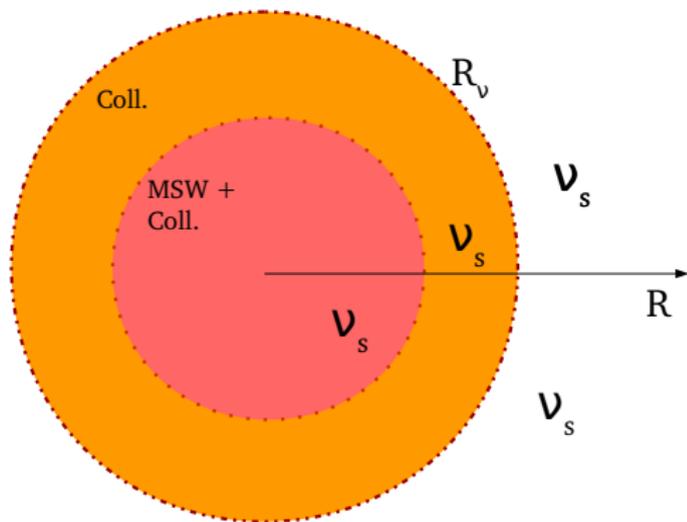
G. G. Raffelt and S. Zhou (2011)

- Suppression / enhancement of the SN explosion
- Change of the electron or neutrino (ν_e, ν_μ, ν_τ) fractions

H. Nunokawa et al. (1997), M. L. Warren et al. (2016), C. A. Argüelles et al. (2016) ...

Sterile neutrino conversions in the stellar core

Sterile neutrino conversions in the stellar core



$\nu_\tau - \nu_s$
mixing

1D SN model
Garching group
archive

Collisions

$$\Gamma_{\nu_s} = \sin^2 2\theta \Gamma_{\nu_\tau}$$

MSW

$$V_{\text{eff}} = \sqrt{2}G_F n_B \left[-\frac{1}{2}Y_n + Y_{\nu_e} + Y_{\nu_\mu} + 2Y_{\nu_\tau} \right]$$

L. Stodolsky (1987), H. Nunokawa et al. (1997), K. Abazajian et al. (2001)

Sterile neutrino conversions in the stellar core

Collisional production

$$\langle P_{\nu_\tau \rightarrow \nu_s}(E) \rangle \approx \frac{1}{2} \frac{\sin^2 2\theta}{(\cos 2\theta - 2V_{\text{eff}}E/\Delta m_s^2)^2 + \sin 2\theta^2 + D^2}$$

MSW production

$$P_{\nu_\tau \rightarrow \nu_s}(E_{\text{res}}) = 1 - \exp\left(-\frac{\pi^2}{2}\gamma\right), \gamma = \Delta_{\text{res}}/l_{\text{osc}}$$

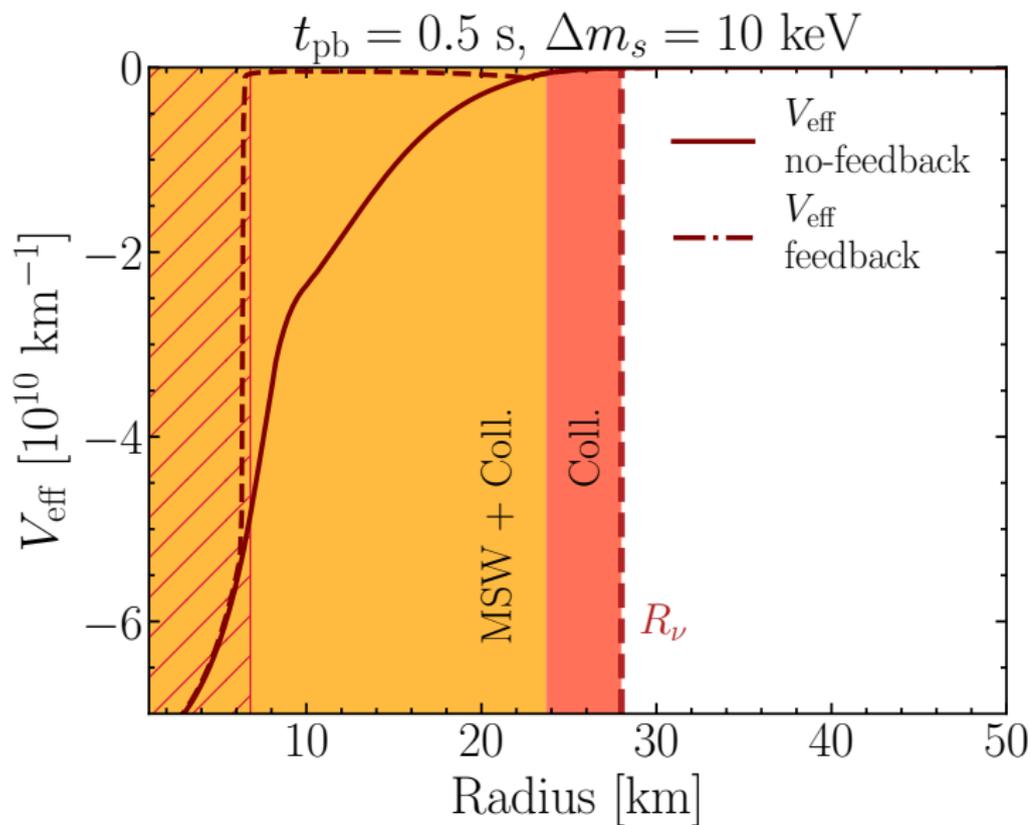
$$\Gamma_\nu(E) \simeq n(r)\sigma(E, r)$$

$$D = \frac{E\Gamma_\nu(E)}{\Delta m_s^2}$$

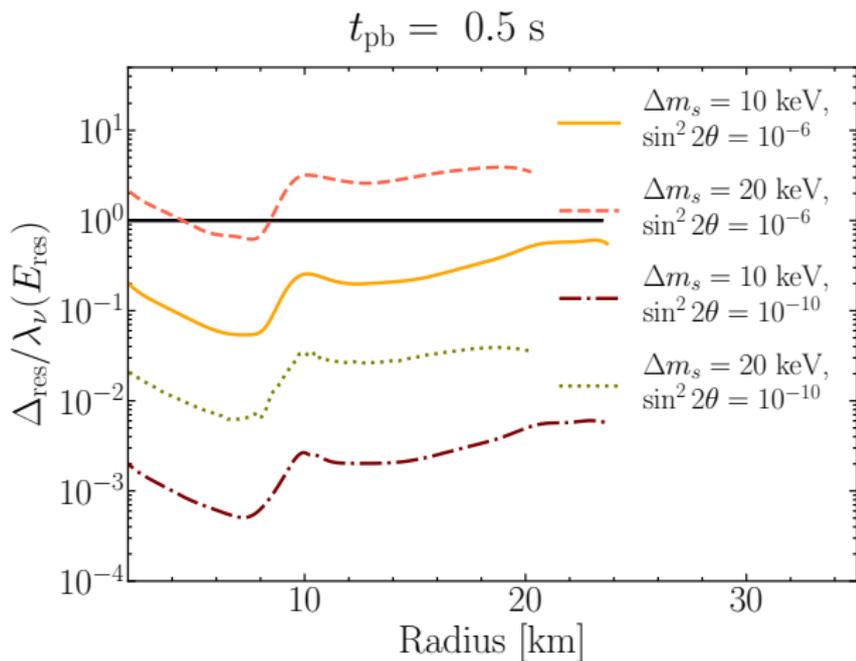
$$\Delta_{\text{res}} = \tan 2\theta \left| \frac{dV/dr}{V} \right|^{-1}$$

$$l_{\text{osc}}(E_{\text{res}}) = (2\pi E_{\text{res}})/(\Delta m_s^2 \sin 2\theta)$$

Conversion regions



Will they collide or undergo MSW resonance?

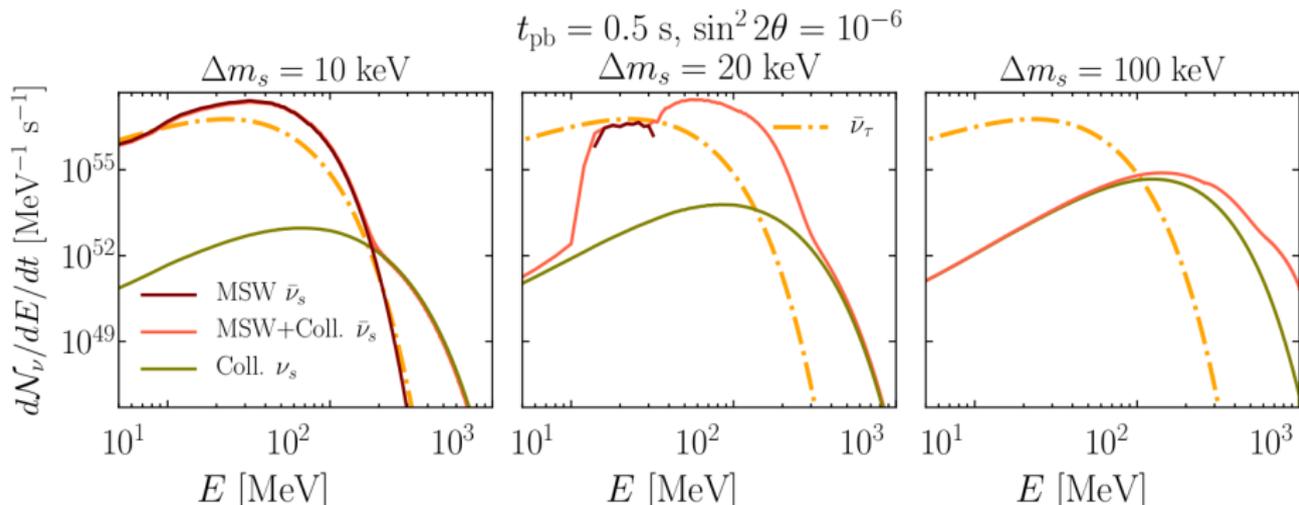


$$\Delta_{\text{res}} = \tan 2\theta \left| \frac{dV/dr}{V} \right|^{-1}$$

$$\lambda_{\nu}(E_{\text{res}}) \simeq \frac{1}{n(r)\sigma(E,r)}$$

$$\Delta_{\text{res}} < \lambda_{\nu}(E_{\text{res}}) ?$$

Sterile neutrino energy distribution



- antineutrinos - collisional and MSW production
- neutrinos - only collisional production
- $\Delta m_s \uparrow$ - collisions are important, more conversions deep in the core, where λ_ν is small

Development of the neutrino lepton asymmetry

Development of the neutrino lepton asymmetry

Only active neutrinos

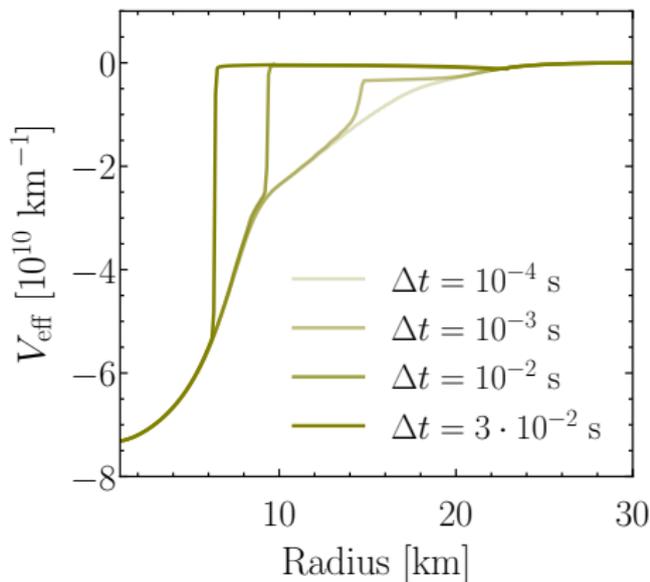
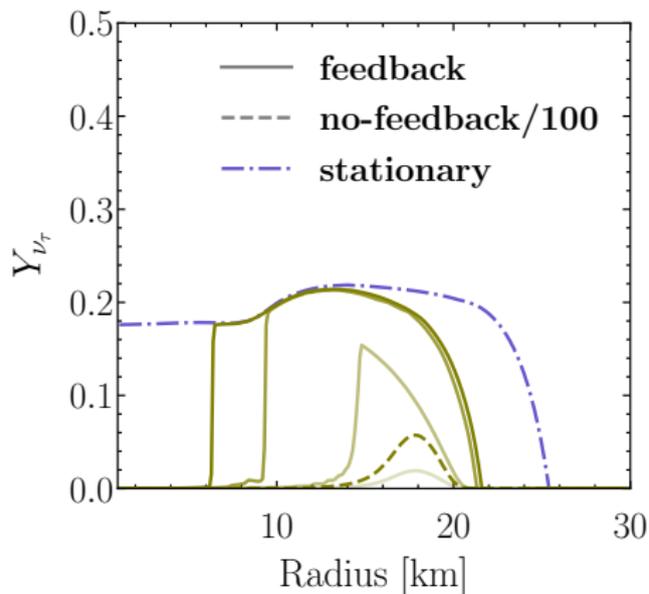
$$Y_{\nu_\tau}(r, t) \equiv 0$$

Active + **sterile** neutrinos

$$Y_{\nu_\tau}(r, t) = \frac{1}{n_b(r)} \int_0^t dt' \frac{d(P_{\nu_\tau \rightarrow \nu_s} n_{\nu_\tau}(r, t') - P_{\bar{\nu}_\tau \rightarrow \bar{\nu}_s} n_{\bar{\nu}_\tau}(r, t'))}{dt'}$$

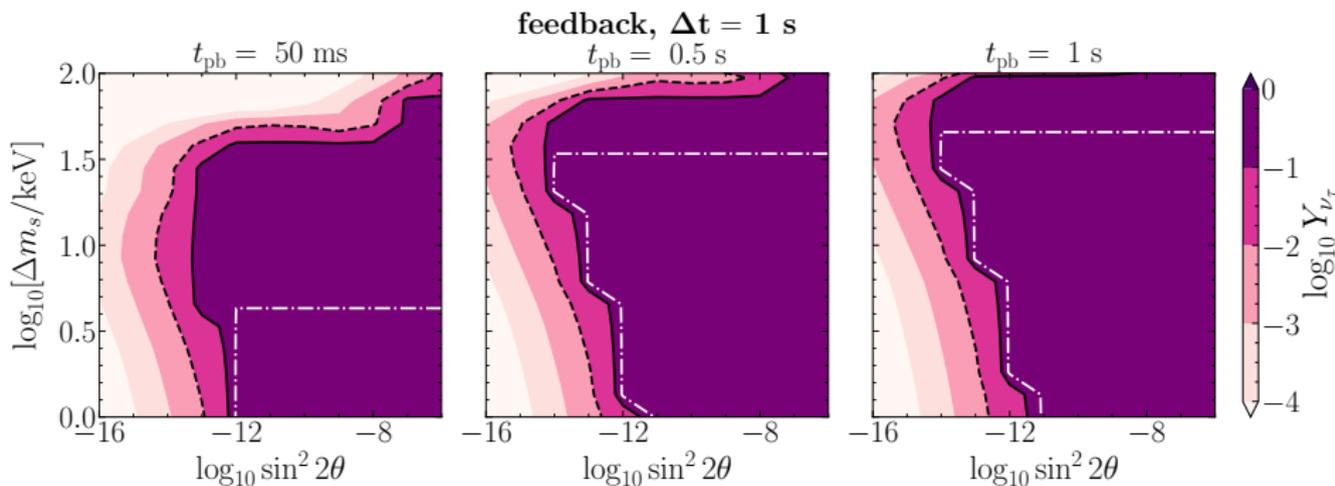
Radial evolution of the asymmetry w and w/o feedback

$$t_{\text{pb}} = 0.5 + \Delta t \text{ s}, \quad \Delta m_s = 10 \text{ keV}, \quad \sin^2 2\theta = 10^{-10}$$



- Feedback inhibits Y_{ν_τ} from unphysical growth.
- Stationary value of Y_{ν_τ} can be reached very quickly.

Contour plot



- Higher mixing angles reach the saturation value faster.
- More massive sterile neutrinos reach smaller saturation values, less energy modes has enhanced conversion probability.

Conclusions

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- **Sterile neutrinos with keV mass**
 - have a major impact on SN physics.
 - Their production leads to the growth of Y_{ν_τ} asymmetry.
- **Feedback is crucial.**
- **Large Y_{ν_τ} asymmetry relevant for SN physics.**
- **SN bounds on the sterile neutrino DM must be updated.**

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Thank you!