

# Dirac vs. Majorana HNLs (and their oscillations) at SHiP

arXiv: 1912.05520

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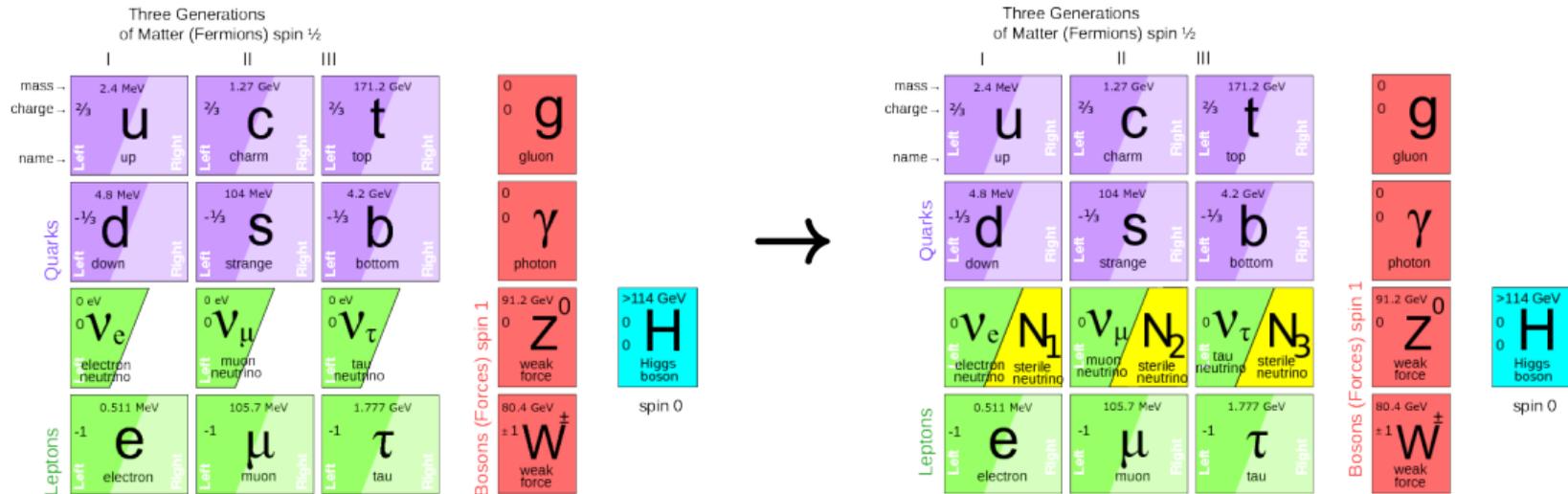
# Feebly interacting particles

- While we wait for the next hadron collider (FCC-hh: 2040–2060) to probe the *energy frontier*, let's explore the *intensity frontier* using low-energy, high-intensity experiments.  
→ C.f. Oleg's talk this morning.
- **Feebly interacting particles (FIPs)**: particles interacting with the SM with a suppressed coupling. The new degrees of freedom are typically *SM singlets*.

## FIP candidates

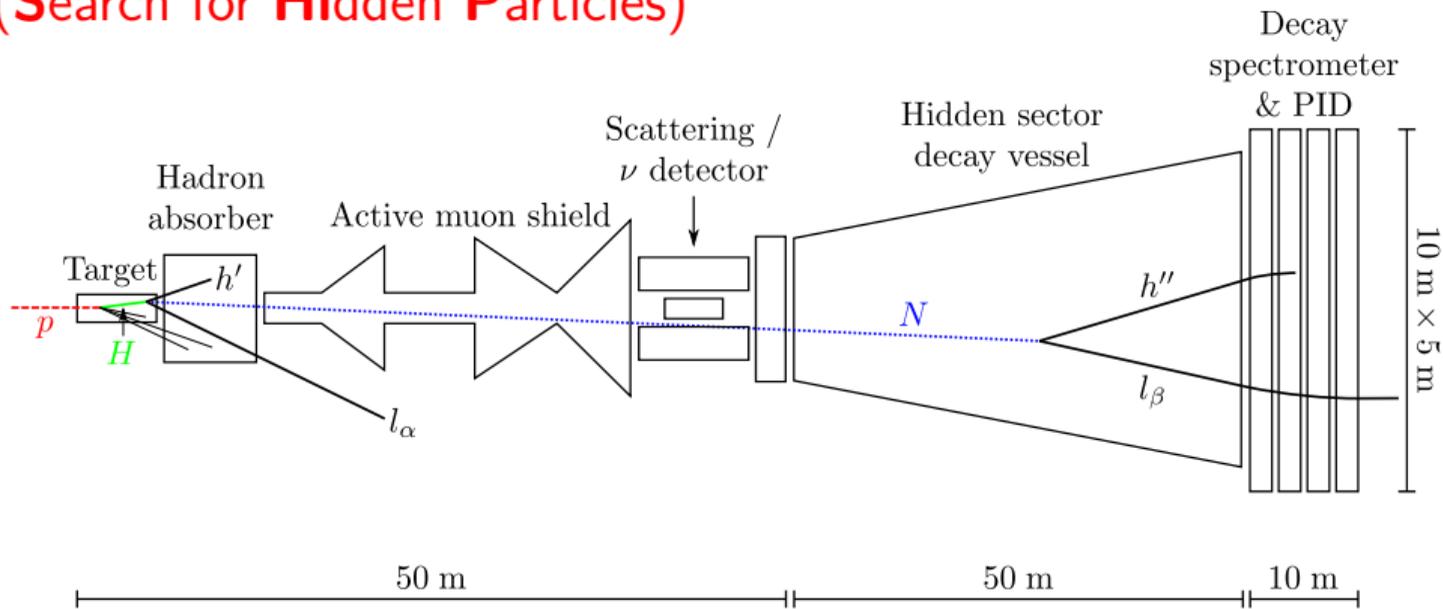
- Renormalizable portals (mix with interacting SM states, or interact with small coupling):
  - 1 Spin 0: scalar portal (dark Higgs).
  - 2 Spin  $\frac{1}{2}$ : neutrino portal (heavy neutral lepton).
  - 3 Spin 1: vector portal (dark photon).
- Non-renormalizable portals (interact through higher dimensional operators):
  - Axion-like particles.
  - ...

# Heavy Neutral Leptons (HNLs)



- HNLs can explain **neutrino masses** and **oscillations** (maybe: baryogenesis, dark matter).
- They interact via **mixing** with flavor eigenstates:  $\nu_\alpha = U_{\alpha i}^{\text{PMNS}} \nu_i + \Theta_{\alpha I} N_I$ ,  $\Theta \ll 1$ .
- Largely constrained below the kaon mass, the neutrino portal will be probed at the GeV scale by the proposed **SHiP experiment**.

# SHiP (Search for Hidden Particles)

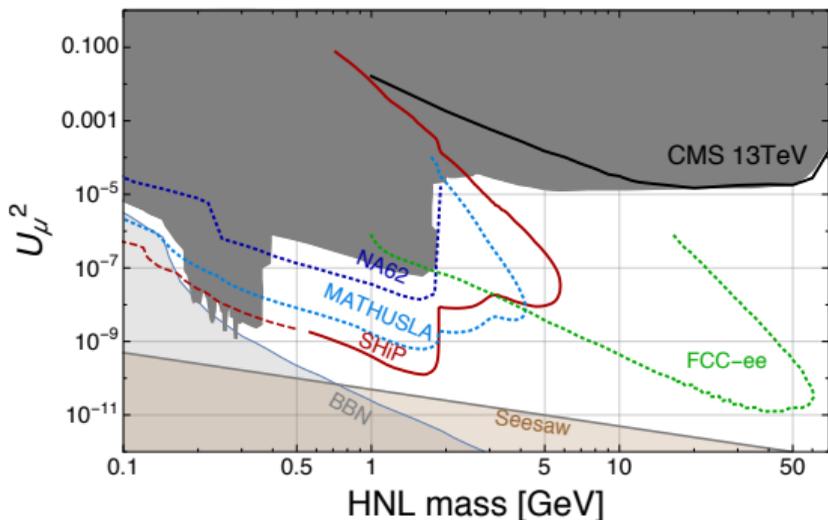


- Low-background (0.1 evts.) beam-dump experiment @ 400 GeV SPS;  $2 \cdot 10^{20}$  POT in 5 yr.
- Comprehensive Design Study for SHiP and Beam Dump Facility submitted last December.
- SHiP aims to *observe* HNLs, and *measure* their mass and mixing angles.

What else can we learn about the *properties* of HNLs at SHiP?

## Detour: realistic HNL benchmarks

Sensitivity study [1811.00930] / PBC  
[1901.09966] assume **one** Majorana HNL,  
mixing with **one** generation only.



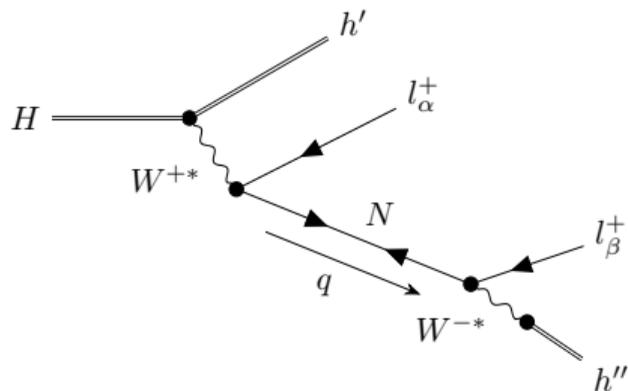
**But:**

- $\nu$  masses generated by **see-saw** mechanism:

$$m_{\alpha\beta} \cong - \sum_I M_I \Theta_{\alpha I} \Theta_{\beta I}$$

- For **one** HNL, the seesaw limit is a **prediction**:  
E.g. for a 1 GeV HNL, we expect  $|\Theta|^2 \sim 10^{-10}$ !
- To generate two distinct  $\Delta m^2$ , at least **two** HNLs are needed, mixing with at least **two** generations.
- If multiple HNLs are **degenerate** as in the  $\nu$ MSM, their mixing angles can be large.

# Majorana HNLs



- New states: SM singlets w/ Majorana mass term.
- Massive states: Majorana *particles*.  
 $\Rightarrow$  Can violate lepton number.
- If we want large mixing angles and correct neutrino masses, lepton number violating (LNV) effects may be **suppressed** (Shaposhnikov [hep-ph/0605047], Kersten and Smirnov [0705.3221]).

- *Is there any hope of observing LNV at all? At SHiP?*
- **Yes & yes!**
- We might even measure the **mass splitting!**

## Main idea

- If there are two quasi-degenerate HNLs, they can **oscillate** among themselves.
- Oscillations in the sterile sector can be **lepton number violating**. For  $|\Theta|^2 \gg m_\nu/M_M$ ,

$$\text{LNC rate} \propto 1 + \cos(\delta M \tau)$$

$$\text{LNV rate} \propto 1 - \cos(\delta M \tau)$$

- To observe them, we need to remember that HNLs are *long-lived*.
- Whether LNV is observable depends on the **mass splitting**  $\delta M$  and **proper lifetime**  $\tau$ :

$$\delta M \tau \ll 2\pi \Rightarrow \text{LNC only}$$

$$\delta M \tau \gg 2\pi \Rightarrow \text{LNC} + \text{LNV with equal integrated rates}$$

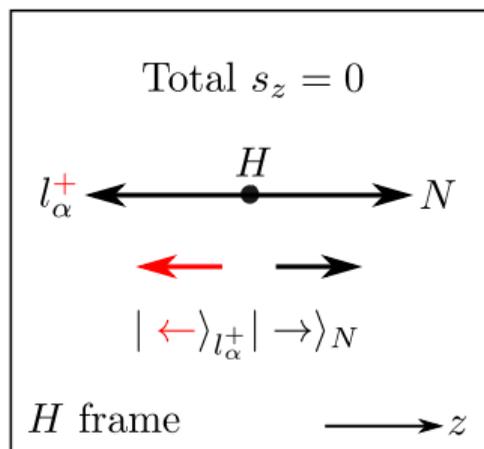
$$\delta M \tau \sim 2\pi \Rightarrow \text{Potentially resolvable oscillations}$$

### Consequences of HNL oscillations

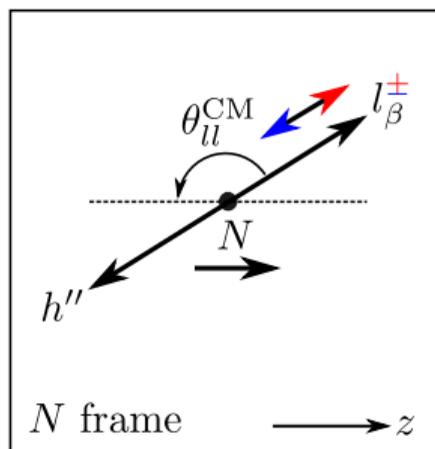
- LNV **may** be suppressed (especially at large mass, cf. Drewes, Klarić, Klose [1907.13034]).  
 $\Rightarrow$  existing bounds relying on LNV might not be valid.
- Observation of LNV (or LNC only) constrains the number and **mass splitting** of HNLs.

# Distinguishing LNC / LNV events at SHiP

- Most production processes are  $H \rightarrow [h']l_\alpha N$ .
- We select the fully reconstructible decay channels  $N \rightarrow l_\beta \pi$ .
- Can we compare the lepton charges?  
→ No! Because the primary decay takes place **inside the target**.
- HNLs carry not only lepton number, but also spin  $\frac{1}{2}$  → look at angular distributions.
- It turns out LNC / LNV processes have very **different kinematics**! E.g. for 2-body decays:



boost  
----->



$$l_\beta^+ \langle \nearrow | \rightarrow \rangle_N = \sin \left( \frac{\theta_{ll}^{\text{CM}}}{2} \right)$$

$$l_\beta^- \langle \searrow | \rightarrow \rangle_N = \cos \left( \frac{\theta_{ll}^{\text{CM}}}{2} \right)$$

$$\mathcal{P} \propto 1 \mp \cos(\theta_{ll}^{\text{CM}})$$

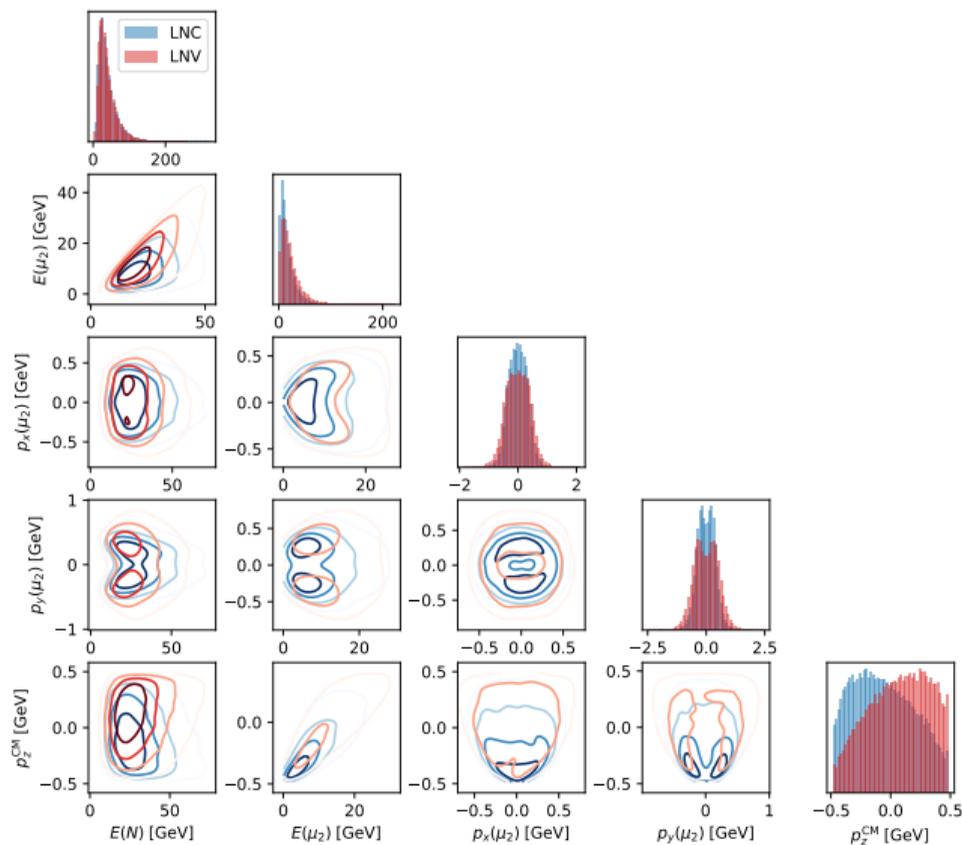
# Complications

- Not all production processes are 2-body decays.
- Decay products ( $l_\alpha, l_\beta, \pi$ ) are not massless  $\Rightarrow$  helicity flips are possible.
- Heavy mesons are not monochromatic  $\Rightarrow$  smears out the distribution of decay products.
- We need to take geometrical acceptance into account.

- **To handle these complications, we need a Monte-Carlo simulation!**
- We use our own Monte-Carlo because we need finer control (tracking spin correlations) over matrix elements compared to what Pythia provides.

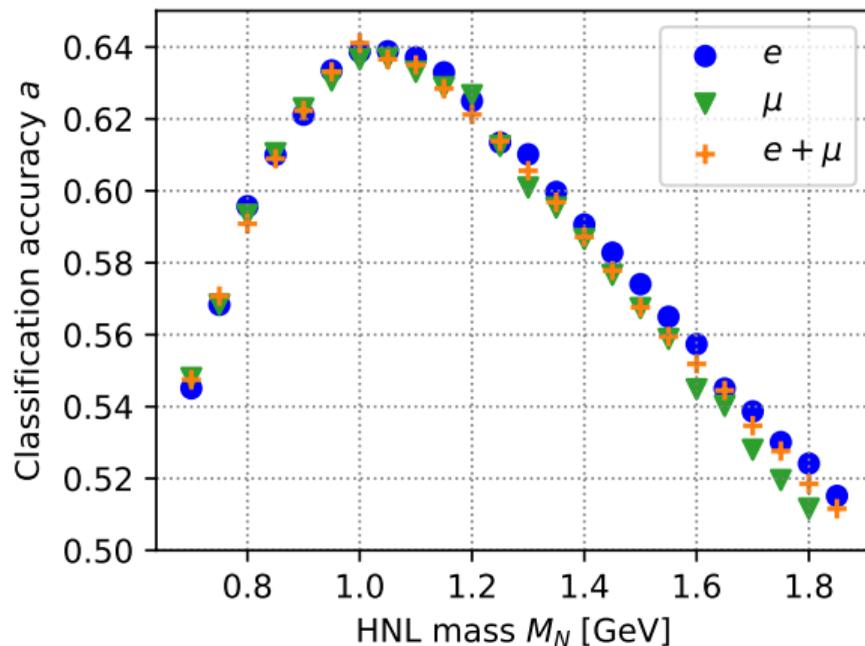
# LNC / LNV distributions

- Most 2/3-body decays implemented.
- $D$ -meson spectra from the LEBC-EHS experiment at the SPS @ 400 GeV.
- Basic propagation and geometrical acceptance cuts.
- Different distributions  $\implies$  can be distinguished given **enough events**.



## We can discriminate these processes using boosted decision trees

- Generate  $3 \cdot 10^6$  events for each mass, split 0.5 : 0.2 : 0.3 into training / validation / test.
- We use the LightGBM gradient boosting algorithm.
- Accuracy is highest when the HNL kinetic energy in CM  $\gtrsim$  heavy meson  $p_T$  spread.



# How to quantitatively distinguish Majorana / Dirac?

## Hypotheses we want to distinguish

- 1  $\mathcal{H}_1$  (**Dirac-like**): HNLs are Dirac or quasi-Dirac with  $\delta M\tau \ll 2\pi$  (LNC only).
- 2  $\mathcal{H}_2$  (**Majorana-like**): HNLs are Majorana or quasi-Dirac with  $\delta M\tau \gg 2\pi$  (LNC + LNV).

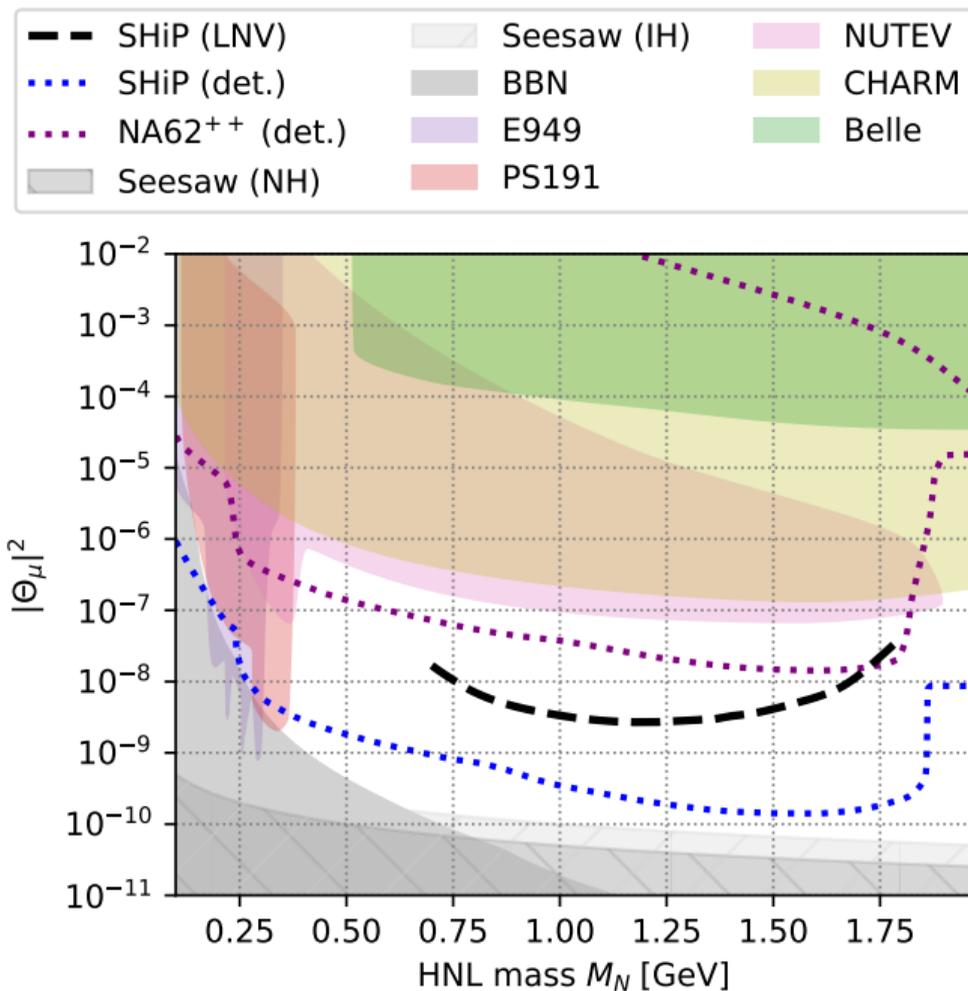
## Model-selection sensitivity

- **Assumptions:** The mass  $M_N$  and  $U_e^2 : U_\mu^2$  ratio have roughly been measured.
- Compute the likelihood of each hypothesis based on the classifier decisions and accuracy.
- Considering in turn each hypothesis as the null hypothesis, draw the “**model-selection sensitivity**” curve where SHiP has a 1/2 probability of excluding this hypothesis at 90% CL if the other is true, after 5 years of nominal operation i.e.  $2 \cdot 10^{20}$  POT.

# Model-selection sensitivity

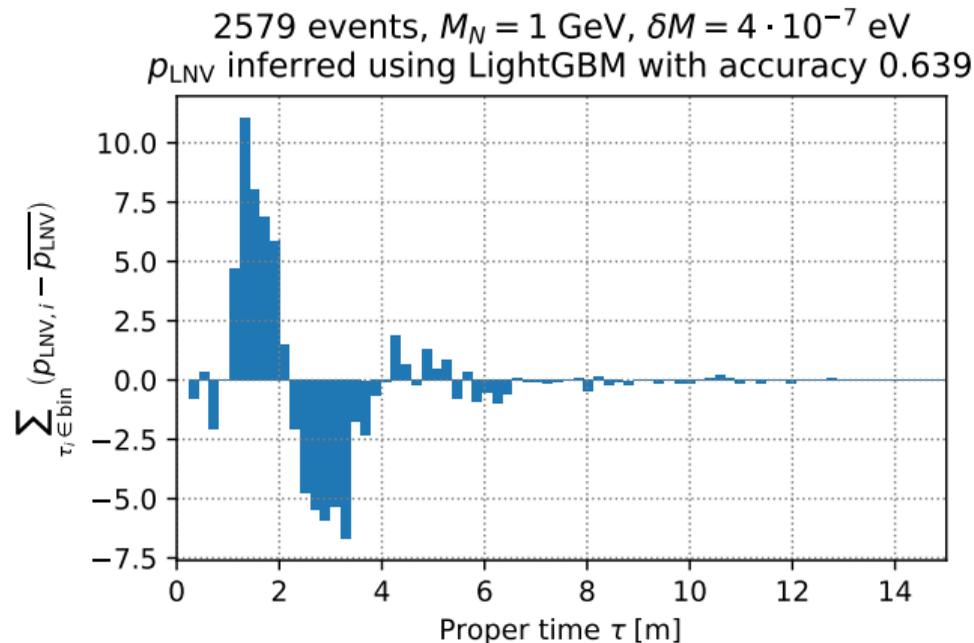
- **Dashed line:** model-selection sensitivity.
- **Colored areas:** existing exclusion bounds
- **Dotted lines:** future experiments that can reconstruct the HNL mass.
- **Hatched areas:** seesaw lower bound.

Source: Physics Beyond Colliders report (arXiv: 1901.09966)



## Resolving HNL oscillations

- Simultaneous requirement of BAU and DM production in the  $\nu$ MSM suggests  $\delta M$  that could be resolved at SHiP (Canetti and Shaposhnikov [1208.4607]).
- Bin events in proper time, weight them by  $P(\text{LNV})$  and subtract the sample average:
- Period of oscillations is  $2\pi/\delta M$ . Allows **measuring** the mass splitting.



# Conclusion

For mixing angles  $|\Theta|^2 \gtrsim 10^{-9}-10^{-8}$ , we can expect many fully reconstructed HNL events.

In this region, SHiP can:

- Test the Majorana nature of HNLs,
- If we are lucky, resolve the mass splitting  $\delta M$ ,

... even if current / next-generation experiments like NA62<sup>++</sup> do not observe any HNLs.

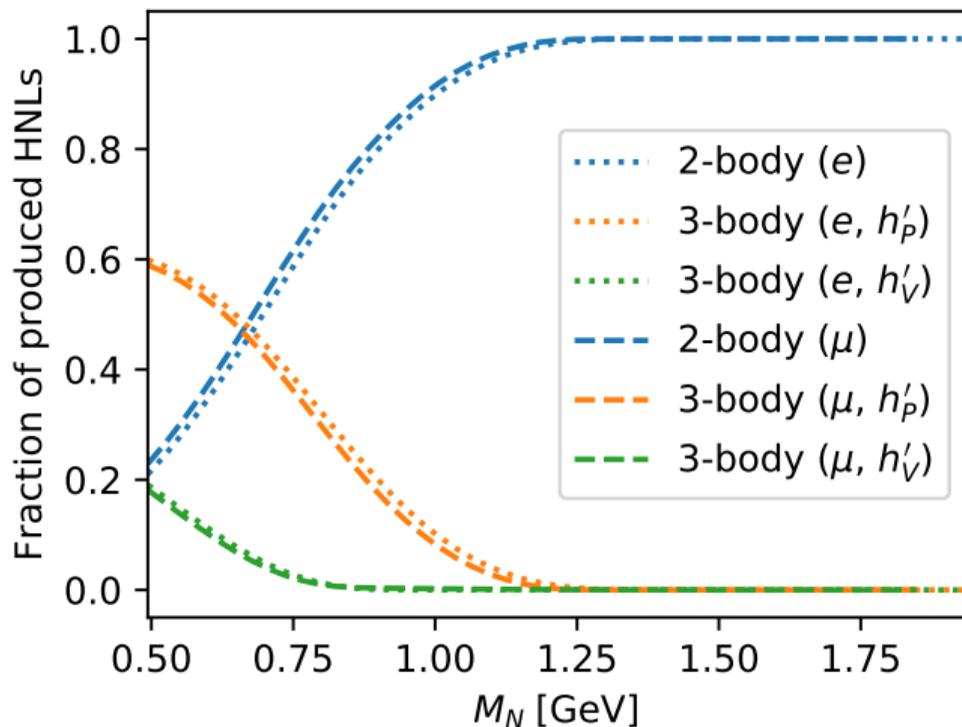
This could help determine the **number** of nearly-degenerate HNLs (needed to measure  $|\Theta_\alpha|^2$ ).

Along with the HNL mass / mixing angles, this would make the  $\nu$ MSM cosmology **predictive**.

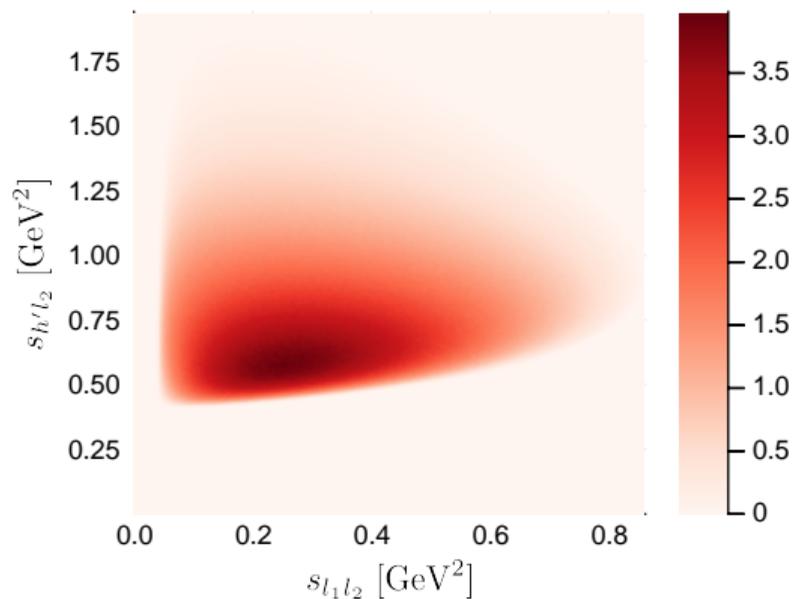
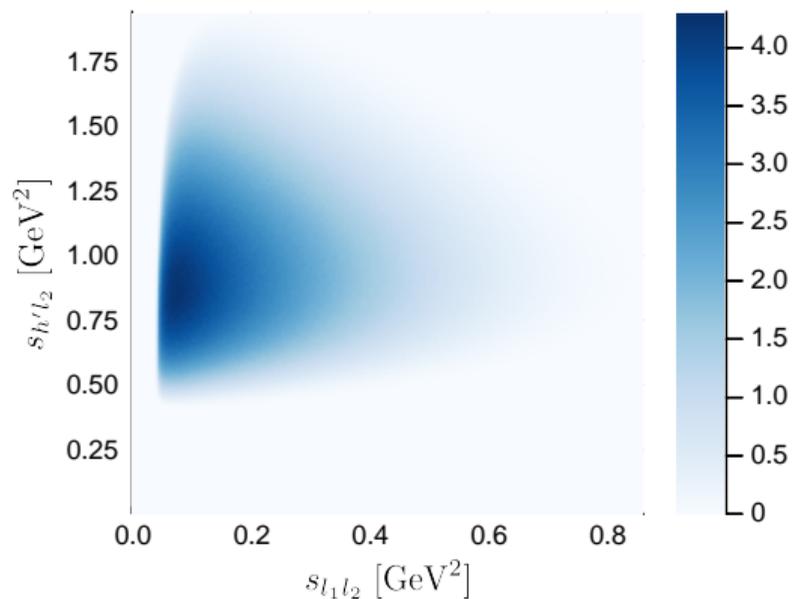
## Related works

- Gorazd Cvetič, Claudio Dib, and C. S. Kim. “Probing Majorana neutrinos in rare  $\pi^+$  to  $e^+ e^- \mu^- \nu$  decays”. In: *Journal of High Energy Physics* 2012.6 (June 2012), p. 149. arXiv: 1203.0573
- Gorazd Cvetič and C. S. Kim. “Rare decays of B mesons via on-shell sterile neutrinos”. In: *Physical Review D* 94.5 (Sept. 2, 2016), p. 053001. arXiv: 1606.04140
- Carolina Arbelaéz et al. “Probing the Dirac or Majorana nature of the Heavy Neutrinos in pure leptonic decays at the LHC”. In: *Physical Review D* 97.5 (Mar. 7, 2018), p. 055011. arXiv: 1712.08704
- Claudio O. Dib, C. S. Kim, and Kechen Wang. “Signatures of Dirac and Majorana Sterile Neutrinos in Trilepton Events at the LHC”. In: *Physical Review D* 95.11 (June 16, 2017), p. 115020. arXiv: 1703.01934
- A. Baha Balantekin, André de Gouvêa, and Boris Kayser. “Addressing the Majorana vs. Dirac Question with Neutrino Decays”. In: *Physics Letters B* 789 (Feb. 2019), pp. 488–495. arXiv: 1808.10518
- P. Hernández, J. Jones-Pérez, and O. Suarez-Navarro. “Majorana vs Pseudo-Dirac Neutrinos at the ILC”. In: *The European Physical Journal C* 79.3 (Mar. 2019), p. 220. arXiv: 1810.07210
- Peter Ballett, Tommaso Boschi, and Silvia Pascoli. “Heavy Neutral Leptons from low-scale seesaws at the DUNE Near Detector”. In: *arXiv:1905.00284 [hep-ph]* (May 1, 2019). arXiv: 1905.00284

## Fraction of produced HNLs by multiplicity and spin

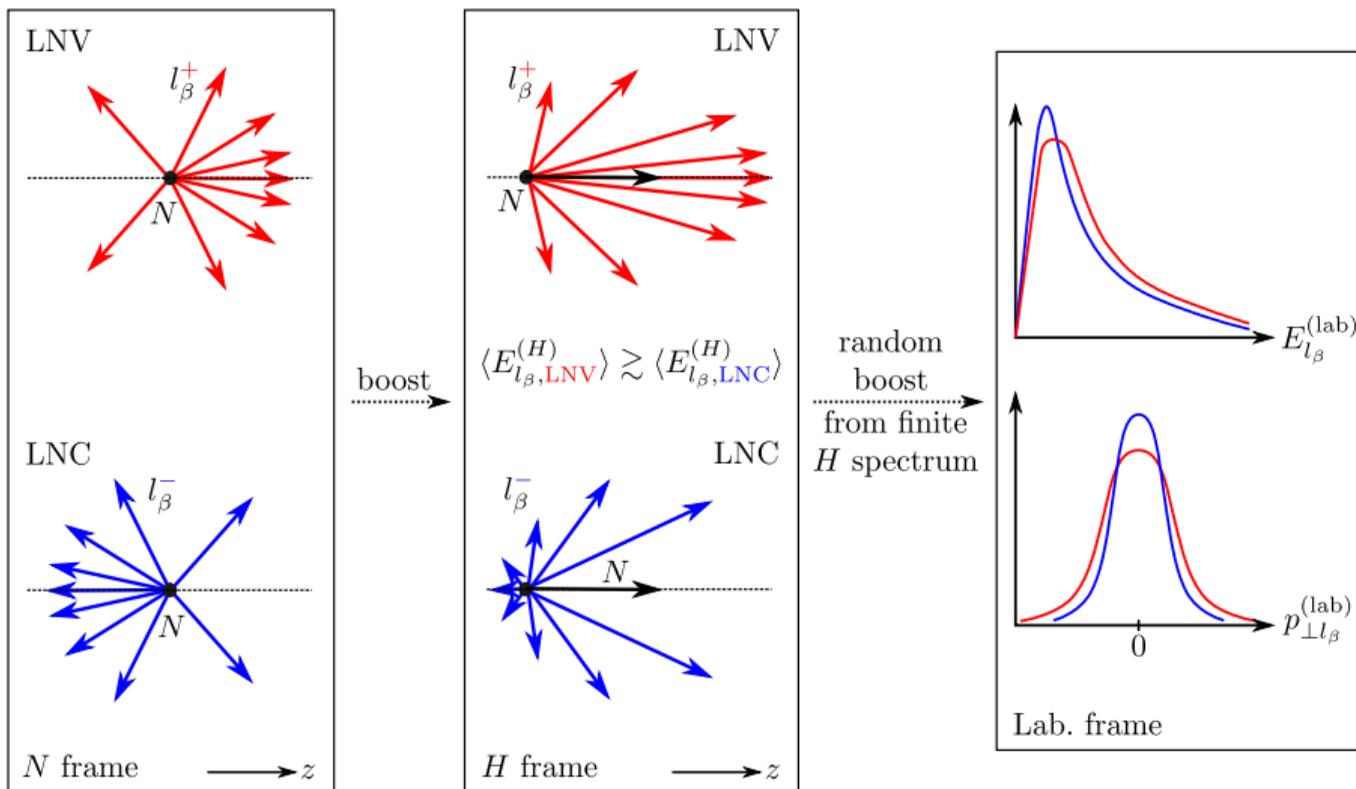


# LNC vs. LNV



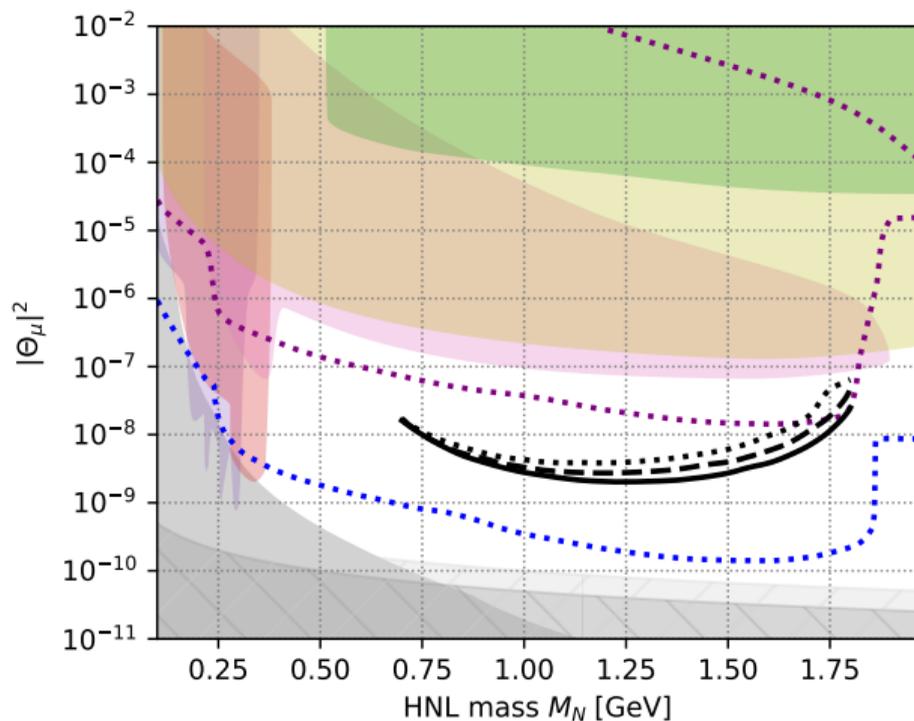
## Angular distribution in the lab frame

- In the **lab frame**, the meson spectrum **smears out** the effect along  $z$ , but not necessarily  $p_T$ .
- If the HNL  $p_T$  (CM) is larger than the heavy hadron  $p_T$  spread (lab), a difference is visible.



# Impact of meson $p_T$ spread

- Higher  $\langle p_T^2 \rangle$   
 $\implies$  lower accuracy  
 $\implies$  curve moves upward
- Solid line:  
 best fit from LEBC-EHS



# Systematic uncertainty on $\langle p_T^2 \rangle$



- What if the real spectrum is different from the simulated one used for training?
- The accuracy mostly depends on the real spectrum, not the one used for training.
- Solid line: best fit from LEBC-EHS

