



Bethe Center for  
Theoretical Physics

UNIVERSITÄT **BONN**

# Feynman integrals beyond modular forms

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based on work in collaboration with

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and

Johannes Brödel, Nils Matthes [2109.xxxx]

Amplitudes 2021

20 August 2021

- Feynman integrals are the building blocks for multi-loop scattering amplitudes.
  - ➔ Important both for collider and gravitational wave phenomenology.
- They exhibit a rich mathematical structure. [See talks by S. Weinzierl, F. Brown, R. Britto, ...]
  - ➔ Connections to algebraic geometry, modular forms,...
- The understanding of the case of polylogarithms and elliptic curves is now relatively well advanced. [See talks by S. Weinzierl & C. Zhang]
  - ➔ The spaces of special functions are (relatively) well understood.

- We know that elliptic curves are not the end of the story!

➔ Also Calabi-Yau varieties appear.

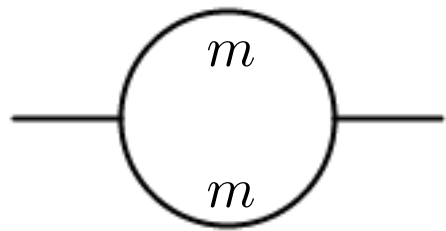
[See talk by A. Klemm]

[Brown, Schnetz; Bloch, Kerr, Vanhove;  
Bourjaily, He, McLeod, von Hippel,  
Wilhelm; Bourjaily, McLeod, von Hippel,  
Wilhelm; Bourjaily, McLeod, Vergu, Volk,  
von Hippel; Klemm, Nega, Safari; Bönisch,  
Fischbach, Klemm, Nega; ... ]

- This case is still poorly explored and understood.

- Goals of this talk:

- ➔ Study the simplest examples of an infinite family of  $l$ -loop integrals associated to Calabi-Yau  $(l - 1)$ -folds.
- ➔ Show how the geometric concepts generalise from the known elliptic cases to Calabi-Yau cases.
- ➔ Present an all-loop generalisation of known results for  $l \leq 3$ .

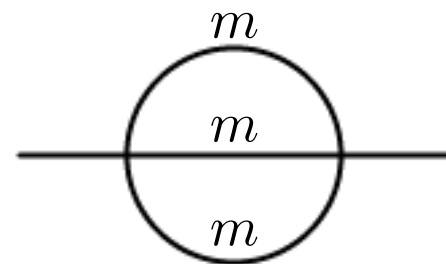


Riemann sphere

$$z = \frac{m^2}{p^2}$$

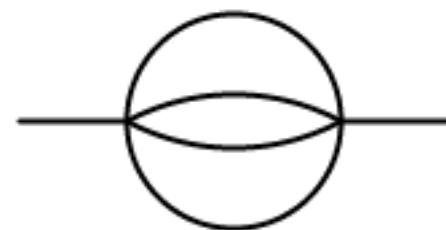
$$\sim \frac{z}{\sqrt{1-4z}} \int_0^z \frac{dz'}{z' \sqrt{1-4z'}}$$

$$\sim \frac{z}{\sqrt{1-4z}} \log \frac{1 - \sqrt{1-4z}}{1 + \sqrt{1-4z}}$$



Elliptic curve

$$\sim h_1(\tau) \int_{i\infty}^{\tau} \frac{d\tau'}{2\pi i} h_3(\tau') \tau'$$



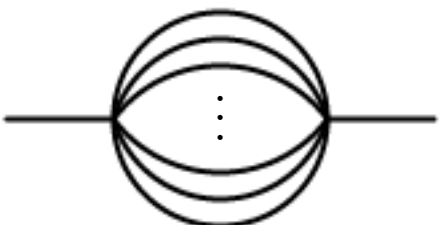
K3 surface  
~ “elliptic x elliptic”  
(only here!)

$$\sim h_2(\tau) \int_{i\infty}^{\tau} \frac{d\tau'}{2\pi i} h_4(\tau') \tau'^2$$

⋮

$h_n$  = ‘modular form of weight  $n$  for  $\Gamma_1(6)$ ’

[See talk by S. Weinzierl]



Calabi-Yau  
 $(l-1)$ -fold

?



# Modular forms

- $\Gamma \subseteq \mathrm{SL}_2(\mathbb{Z}) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} : a, b, c, d \in \mathbb{Z} \text{ and } \det \begin{pmatrix} a & b \\ c & d \end{pmatrix} = 1 \right\}$

$$\mathbb{H} = \{\tau \in \mathbb{C} : \mathrm{Im} \tau > 0\}$$

- **Modular function for  $\Gamma$  :** meromorphic function  $f : \mathbb{H} \rightarrow \mathbb{C}$  s.t.

$$f\left(\frac{a\tau + b}{c\tau + d}\right) = f(\tau)$$

- **Modular of weight  $k$  for  $\Gamma$  :** holomorphic function  $h_k : \mathbb{H} \rightarrow \mathbb{C}$  s.t.

$$h_k\left(\frac{a\tau + b}{c\tau + d}\right) = (c\tau + d)^k h_k(\tau)$$

[+ holomorphicity at the cusps]

- Master integrals for  $l$ -loop equal-mass bananas ( $D = 2 - 2\epsilon$ ) :

$$J_{l,1}(z; \epsilon) \sim \text{---} \circlearrowleft \text{---} \quad J_{l,k}(z; \epsilon) \sim \partial_z^{k-1} J_{l,1}(z; \epsilon) \quad 2 \leq k \leq l$$

- The vector of MIs satisfies a system of differential equations:

$$\partial_z \underline{J}_l(z; \epsilon) = \underline{\mathbf{B}}_l(z; \epsilon) \underline{J}_l(z; \epsilon) + \underline{N}_l(z, \epsilon)$$

## Rational matrix

## Inhomogeneity from integrals with fewer propagators

$$\mathbf{B}_l(z; \epsilon) = \mathbf{B}_{l,0}(z) + \sum_{k=1}^l \mathbf{B}_{l,k}(z) \epsilon^k$$

$$\underline{N}_l(z, \epsilon) = \left( 0, \dots, 0, (-1)^{l+1} (l+1)! \frac{z}{z^l \prod_{k \in \Delta^{(l)}} (1 - kz)} \frac{\Gamma(1 + \epsilon)^l}{\Gamma(1 + l\epsilon)} \right)^T$$

- **Step 1:** Solve the homogeneous system at  $\epsilon = 0$ :

$$\partial_z \underline{J}_l^\Gamma(z) = \mathbf{B}_{l,0}(z) \underline{J}_l^\Gamma(z) \quad \longleftrightarrow \quad \mathcal{L}_l J_{l,1}^\Gamma = 0$$

- ➔ Solution space spanned by maximal cuts:

[Tancredi, Primo;  
Frellesvig, Papadopoulos;  
Bosma, Larsen, Zhang]

$$\mathbf{W}_l(z) := \left( \underline{J}_l^{\Gamma^1}(z), \dots, \underline{J}_l^{\Gamma^l}(z) \right)$$

- **Step 2:** Change basis:  $\underline{J}_l(z, \epsilon) = \mathbf{W}_l(z) \underline{L}_l(z, \epsilon)$

$$\partial_z \underline{L}_l(z, \epsilon) = \tilde{\mathbf{B}}_l(z, \epsilon) \underline{L}_l(z, \epsilon) + \tilde{\underline{N}}_l(z, \epsilon)$$

$$\tilde{\mathbf{B}}_l(z, \epsilon) = \mathbf{W}_l(z)^{-1} [\mathbf{B}_l(z, \epsilon) - \mathbf{B}_{l,0}(z)] \mathbf{W}_l(z) = \mathcal{O}(\epsilon)$$

$$\tilde{\underline{N}}_l(z, \epsilon) = \mathbf{W}_l(z)^{-1} \underline{N}_l(z, \epsilon)$$

Polynomial of degree  $l - 1$  and  $l$   
in the entries of  $\mathbf{W}_l(z)$ .

- Step 3: The new system can easily be solved order by order in  $\epsilon$ .

➔ Example:  $\epsilon = 0$

$$\begin{aligned}\underline{L}_l(z, 0) &= \underline{L}_l(0, 0) + \int_0^z dz' \tilde{\underline{N}}_l(z', 0) && \text{Integrand involves polynomials of} \\ &&& \text{degree } l - 1 \text{ in the entries of } \mathbf{W}_l(z). \\ &= \underline{L}_l(0, 0) + \int_0^z dz' \mathbf{W}_l(z')^{-1} \underline{N}_l(z', 0)\end{aligned}$$

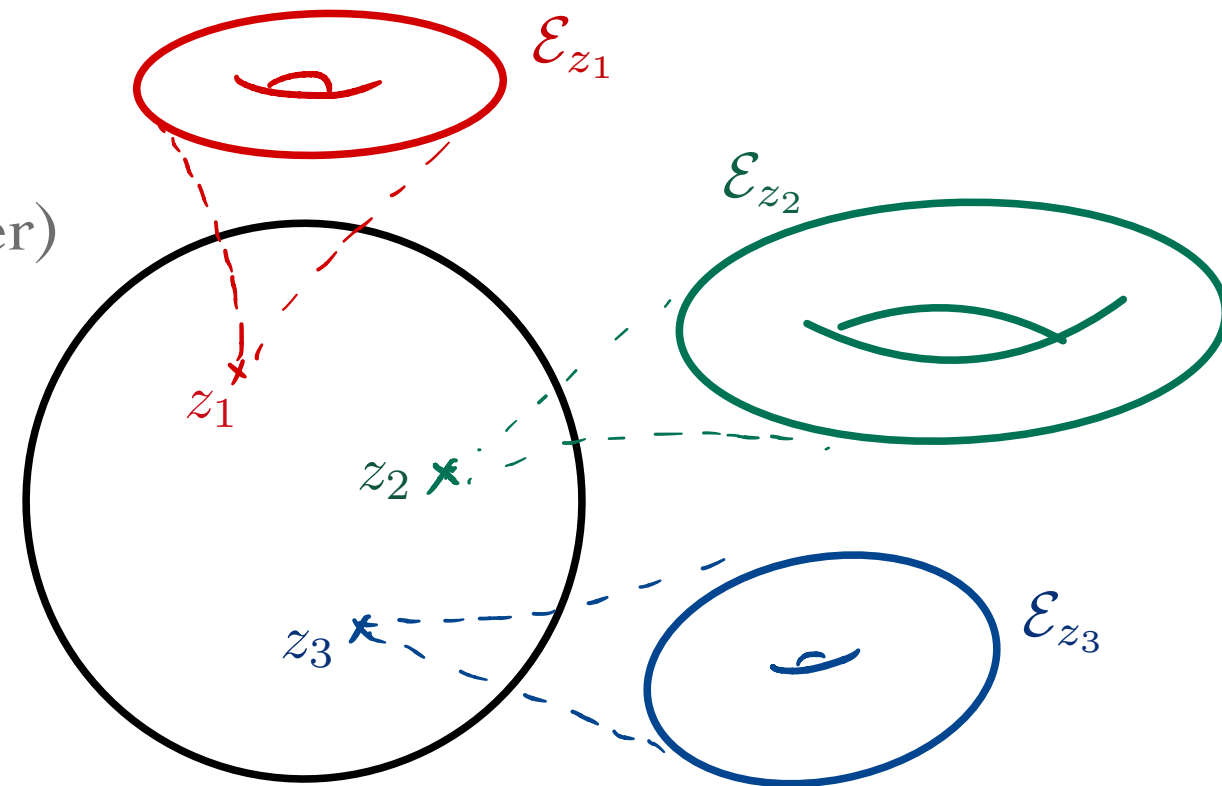
- Questions:

- ➔ What are the entries of  $\mathbf{W}_l(z)$ ?
- ➔ What are their properties, relations, etc.?
- ➔ Why do we get modular forms for 2 (and 3) loops?
- ➔ What about the integrals we get?

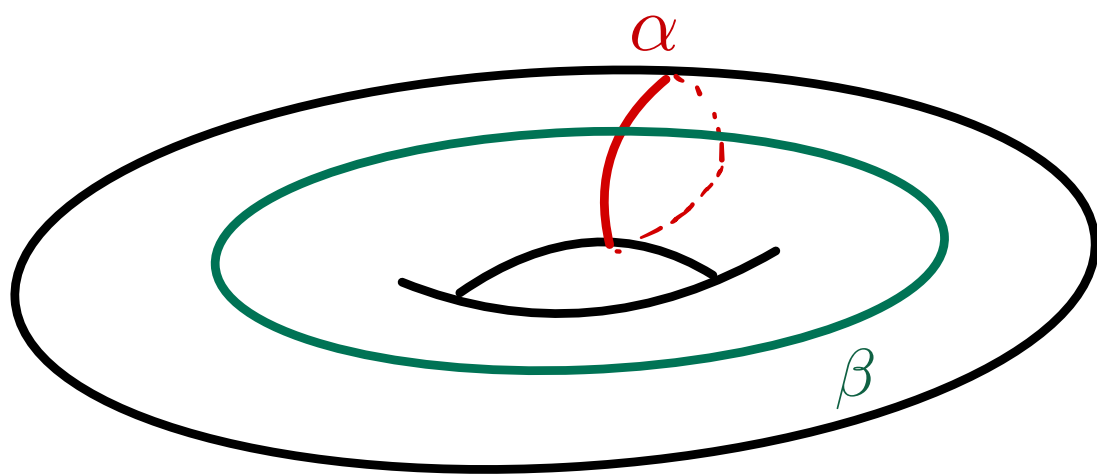
# 2-loop banana integrals and elliptic curves

- The 2-loop (equal-mass) banana family is associated to a (1-parameter) family of elliptic curves

→ For every  $z \in \mathbb{CP}^1$  there is a (possibly degenerate) elliptic curve  $\mathcal{E}_z$ .



- **Elliptic curve:** Riemann surface of genus 1 = Torus



→ There is a unique holomorphic differential 1-form  $\frac{dx}{y}$ .

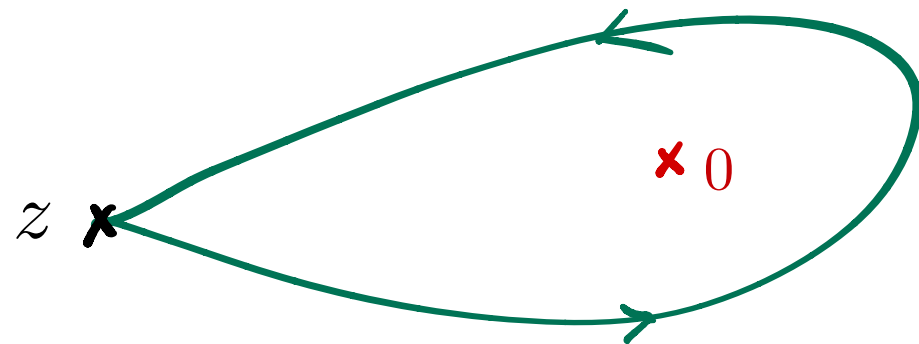
→ Characterised by its periods:

$$\Pi_\alpha = \int_\alpha \frac{dx}{y} \quad \Pi_\beta = \int_\beta \frac{dx}{y}$$

- Here: Periods will be functions of  $z$ :
  - ➔  $z$ -dependence governed by Picard-Fuchs equation ( $\theta = z\partial_z$ ):
$$\mathcal{L}_2 \Pi_\gamma(z) = 0 \quad \mathcal{L}_2 = (1-z)(1-9z)\theta^2 - 2(1-5z)\theta + (1-3z)$$
  - ➔ Same differential operator that annihilates maximal cuts!
  - ➔ Maximal cuts at two-loop can be associated to the periods of the family of elliptic curves.
  - ➔ Singular points of  $\mathcal{L}_2$  :  $0, 1/9, 1, \infty$ .
- We get multi-valued functions on  $X_2 = \mathbb{CP}^1 \setminus \{0, 1/9, 1, \infty\}$ .

$$\Pi_\alpha(z) = 2\pi i z + \mathcal{O}(z^2) \quad \Pi_\beta(z) = \Pi_\alpha(z) \frac{\log z}{2\pi i} + c_2 z^2 + \mathcal{O}(z^3)$$

- Analytic continuation around 0:



$$\begin{pmatrix} \Pi_\beta(z) \\ \Pi_\alpha(z) \end{pmatrix} \longrightarrow \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \Pi_\beta(z) \\ \Pi_\alpha(z) \end{pmatrix}$$

- We can compute such a matrix for every small loop around a singular point (N.B.  $\mathbf{T}_0 \mathbf{T}_{1/9} \mathbf{T}_1 \mathbf{T}_\infty = 1$ ) :

$$\mathbf{T}_0 = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \quad \mathbf{T}_1 = \begin{pmatrix} 1 & 0 \\ -6 & 1 \end{pmatrix} \quad \mathbf{T}_{1/9} = \begin{pmatrix} 7 & 2 \\ -18 & -5 \end{pmatrix}$$

- ➡ These matrices generate the monodromy group of  $\mathcal{L}_2$  :

$$\Gamma_1(6) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathrm{SL}_2(\mathbb{Z}) : a, d = 1 \bmod 6, c = 0 \bmod 6 \right\}$$



- Define:  $\tau = \frac{\Pi_\beta(z)}{\Pi_\alpha(z)} \qquad q = e^{2\pi i\tau} = z + 4z^2 + \mathcal{O}(z^3)$

➔ Can be inverted:  $z(\tau) = q - 4q^2 + \mathcal{O}(q^3)$

- Action of the monodromy group:

$$\tau \longrightarrow \frac{a\Pi_\beta(z) + b\Pi_\alpha(z)}{c\Pi_\beta(z) + d\Pi_\alpha(z)} = \frac{a\tau + b}{c\tau + d}$$

$$z(\tau) \longrightarrow z\left(\frac{a\tau + b}{c\tau + d}\right) = z(\tau) \quad \text{Modular function for } \Gamma_1(6)$$

$$h_1(\tau) = \Pi_\alpha(z(\tau)) \longrightarrow (c\tau + d) h_1(\tau) \quad \text{Modular form of weight 1 for } \Gamma_1(6)$$

$$\begin{aligned}
 & \text{Diagram: a circle with three horizontal lines passing through it, labeled } m \text{ at the top, middle, and bottom.} \\
 & \sim h_1(\tau) \int_{i\infty}^{\tau} \frac{d\tau'}{2\pi i} h_3(\tau') \tau' \sim \underline{\Pi}(z)^T \Sigma_2 \int_{\infty}^z \frac{dz'}{z'^2} \underline{\Pi}(z') \\
 & \quad = \Pi_{\alpha}(z) \quad \sim \Pi_{\alpha}(z)^3 \quad \underline{\Pi}(z) = (\Pi_{\alpha}(z), \Pi_{\beta}(z))^T \\
 & \text{=maximal cut of the sunrise} \\
 & \quad d\tau = \frac{z dz}{64(1-z)(1-9z) \Pi_{\alpha}(z)^2} \quad \Sigma_2 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}
 \end{aligned}$$

➔ Matches structure from differential equations:

$$\underline{J}_l(z, 0) \sim \underline{\mathbf{W}}_l(z) \int_0^z dz' \underline{\mathbf{W}}_l(z)^{-1} \underline{N}_l(z, 0)$$

➔ Also valid at 1-loop (and also 3-loop, but this is special...)

$$\text{Diagram: a circle with two horizontal lines passing through it, labeled } m \text{ at the top and bottom.} \sim \frac{z}{\sqrt{1-4z}} \int_0^z \frac{dz'}{z' \sqrt{1-4z'}}$$

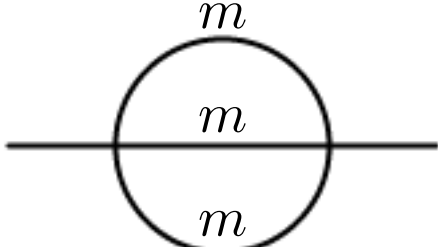
Maximal cut; Monodromy group =  $\mathbb{Z}_2$

- Maximal cuts of  $J_{l,1} \longleftrightarrow$  periods of the elliptic curve.
- Maximal cuts are annihilated by Picard-Fuchs operator  $\mathcal{L}_2$ .  
 ➔ Monodromy group:  $\Gamma_1(6)$

- Close to  $z = 0 : 1$  holomorphic and 1 single-log solution:

$$\Pi_\alpha(z) = 2\pi i z + \mathcal{O}(z^2) \quad \Pi_\beta(z) = \Pi_\alpha(z) \frac{\log z}{2\pi i} + c_2 z^2 + \mathcal{O}(z^3)$$

- General structure of the full Feynman integral:

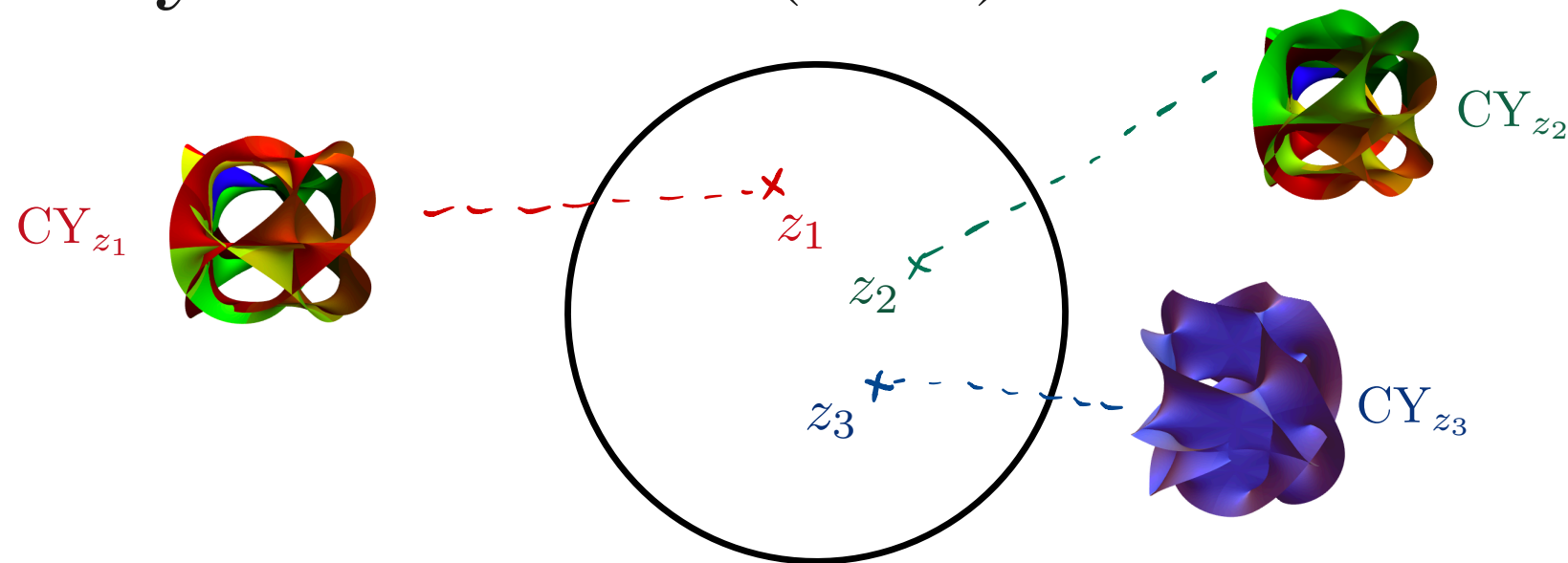


$$\sim \underline{\Pi}(z)^T \Sigma_2 \int_\infty^z \frac{dz'}{z'^2} \underline{\Pi}(z')$$

- For 2 loops: Period is a modular form.

Higher-loop banana  
integrals  
and  
Calabi-Yau varieties

- The  $l$ -loop (equal-mass) banana family is characterised by a (1-parameter) family of Calabi-Yau  $(l - 1)$ -folds:



- **Calabi-Yau  $n$ -fold:**  $n$ -dim. Kähler manifold with a unique holomorphic differential  $n$ -form  $\Omega(z)$ .

➔ Characterised by its periods:  $\underline{\Pi}(z) = (\Pi_1(z), \dots, \Pi_r(z))^T$

$$\Pi_i(z) = \int_{\Gamma_i} \Omega(z) \quad \Gamma_i = \text{Basis of independent cycles.}$$

➔ For the equal-mass banana:  $r = l = \# \text{MIs}$

- The periods provide a basis for the maximal cuts in 2 dimensions.
- They are annihilated by the PF-operator  $\mathcal{L}_l$  of degree  $l$ .
- ➔ They can be constructed explicitly for all  $l$ . [Bönisch, Fischbach  
Klemm, Nega, Safari]

- Set of singular points:  $\Delta^{(l)} := \bigcup_{j=0}^{\lceil \frac{l-1}{2} \rceil} \{(l+1-2j)^2\}$

- Structure of the solutions: close to  $z = 0$ :

$$\Pi_{l,1}(z) = \mathcal{O}(z) \qquad \Pi_{l,k}(z) \sim \Pi_{l,1}(z) \frac{1}{(k-1)!} \log^{k-1} z + \mathcal{O}(z^2)$$

- ➔ Point of maximal unipotent monodromy (MUM-point).
- ➔ Not true for other singular points!
- ➔ (Expected to be) genuinely new transcendental functions.
- ➔ Can be evaluated in a fast and efficient way for all  $z$ .

- General expected form:

$$\text{Maximal cuts} \quad \sim \underline{\underline{\Pi}}_l(z)^T \Sigma_l \int_0^z \frac{dz'}{z'^2} \underline{\underline{\Pi}}_l(z') \quad \text{Linear in the periods}$$

- $\Sigma_l$  is uniquely determined from the geometry (must be invariant under action of the monodromy group):

$$\Sigma_l = \begin{pmatrix} & & & & 1 \\ & & & & \\ & & & -1 & \\ & & 1 & & \\ & \ddots & & & \\ \ddots & & & & \end{pmatrix} \quad [\text{e.g., } \mathbf{M}^T \Sigma_2 \mathbf{M} = \Sigma_2 \text{ for all } \mathbf{M} \in \text{SL}_2(\mathbb{Z})]$$

- But it does not yet match what we (seem to) get from the differential equations....

$$\underline{J}_l(z, 0) \sim \mathbf{W}_l(z) \int_0^z dz' \mathbf{W}_l(z)^{-1} \underline{N}_l(z, 0)$$

Maximal cuts

Polynomial of degree  $l - 1$  in the periods

- Are there relations among the periods?

[See talk by A. Klemm]

$$\underline{\Pi}_l(z)^T \underline{\Sigma}_l \underline{\Pi}_l(z) = \int_M \Omega(z) \wedge \Omega(z) = 0!$$

➔ This is a quadratic relation among maximal cuts!

- Are there more of these relations?

$$\underline{\Pi}_l(z)^T \underline{\Sigma}_l \partial_z \underline{\Pi}_l(z) = \int_M \Omega(z) \wedge \partial_z \Omega(z) = 0$$

➔ This is a quadratic relation among maximal cuts of different MIs.

- What is their origin? How can we find these relations?



- The cohomology group  $H_{dR}^n(M)$  comes with a Hodge filtration ( $n = l - 1$ ).

$$\begin{array}{ccccccc}
 F^{l-1} & \subseteq & F^{l-2} & \subseteq & F^{l-3} & \subseteq & \dots \subseteq F^1 \subseteq F^0 = H_{dR}^n(M) \\
 \hline
 \Omega & & \partial_z \Omega & & \partial_z^2 \Omega & & \partial_z^{l-2} \Omega & & \partial_z^{l-1} \Omega
 \end{array}$$

$$F^p = \bigoplus_{p \leq q \leq l-1} H^{q,p-q} \quad H^{p,q} = \text{cohomology class of } (p,q)\text{-forms}$$

e.g.  $dz_1 \wedge \dots \wedge dz_p \wedge d\bar{z}_1 \wedge \dots \wedge d\bar{z}_q$

- Griffiths transversality:  $\partial_z F^p \subseteq F^{p-1}$

$\rightarrow$  cf. for the elliptic case:  $\partial_z \int_{\alpha} \frac{dx}{y} = A \int_{\alpha} \frac{dx}{y} + B \int_{\alpha} \frac{x dx}{y}$

$\in F^1$                        $\in F^1$                        $\in F^0$

$$\underline{\Pi}_l(z)^T \Sigma_l \underline{\Pi}_l(z) = \int_M \Omega(z) \wedge \Omega(z) = 0$$

‘ $(n, 0) \wedge (n, 0)$ ’,

$$\underline{\Pi}_l(z)^T \Sigma_l \partial_z \underline{\Pi}_l(z) = \int_M \Omega(z) \wedge \partial_z \Omega(z) = 0$$

‘ $(n, 0) \wedge (n-1, 1)$ ’,

⋮

$$\underline{\Pi}_l(z)^T \Sigma_l \partial_z^{l-2} \underline{\Pi}_l(z) = \int_M \Omega(z) \wedge \partial_z^{l-2} \Omega(z) = 0$$

‘ $(n, 0) \wedge (1, n-1)$ ’,

$$\underline{\Pi}_l(z)^T \Sigma_l \partial_z^{l-1} \underline{\Pi}_l(z) = \int_M \Omega(z) \wedge \partial_z^{l-1} \Omega(z) = C_l(z)$$

‘ $(n, 0) \wedge (0, n)$ ’,

$$C_l(z) = \frac{1}{z^{l-3} \prod_{k \in \Delta^{(l)}} (1 - kz)}$$

- We obtain a matrix of quadratic relations among maximal cuts!

$$\mathbf{Z}_l(z) = \begin{pmatrix} \underline{\Pi}_l(z)^T \boldsymbol{\Sigma}_l \underline{\Pi}_l(z) & \cdots & \underline{\Pi}_l(z)^T \boldsymbol{\Sigma}_l \partial_z^{l-1} \underline{\Pi}_l(z) \\ \vdots & \ddots & \vdots \\ \partial_z^{l-1} \underline{\Pi}_l(z)^T \boldsymbol{\Sigma}_l \underline{\Pi}_l(z) & \cdots & \partial_z^{l-1} \underline{\Pi}_l(z)^T \boldsymbol{\Sigma}_l \partial_z^{l-1} \underline{\Pi}_l(z) \end{pmatrix}$$

$$= \mathbf{W}_l(z) \boldsymbol{\Sigma}_l \mathbf{W}_l(z)^T$$

- ➔ All entries are explicitly calculable rational functions, e.g.,

$$\mathbf{Z}_3(z)^{-1} = \begin{pmatrix} \frac{1}{z^2} - \frac{8}{z} & -10 + 64z & 1 - 20z + 64z^2 \\ -10 + 64z & -1 + 20z - 64z^2 & 0 \\ 1 - 20z + 64z^2 & 0 & 0 \end{pmatrix}$$

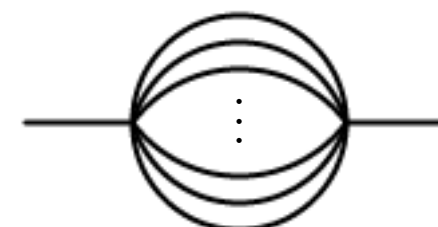
- ➔ We have searched for more quadratic relations, but we could not find any (for generic  $z$ ).
- ➔ We conjecture that these are the only quadratic relations.

- It is now easy to show that

$$\mathbf{W}_l(z)^{-1} = \mathbf{\Sigma}_l \mathbf{W}_l(z)^T \mathbf{Z}_l(z)^{-1}$$

➔ The inverse is linear in the periods!

- Putting it all together, solution of the differential equation takes the form:



$$= \underline{\Pi}_l(z)^T \underline{L}_l^{(0)} + (l+1)! \underline{\Pi}_l(z)^T \mathbf{\Sigma}_l \int_{\vec{1}_0}^z \frac{dw}{w^2} \underline{\Pi}_l(w)$$

[Bönisch, CD, Fischbach  
Klemm, Nega]

➔ Compact formula valid for arbitrary loops!

➔ We still need the initial condition (related to  $z \rightarrow 0$ ):

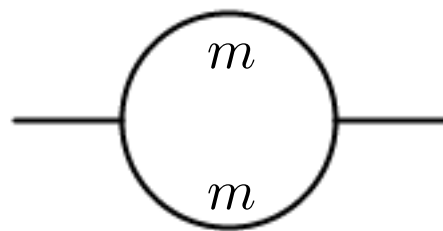
$$\underline{L}_l^{(0)} = (L_{l,1}^{(0)}, \dots, L_{l,l}^{(0)})^T$$

$$\text{---} \circ \text{---} = \underline{\Pi}_l(z)^T \underline{L}_l^{(0)} + (l+1)! \underline{\Pi}_l(z)^T \underline{\Sigma}_l \int_{\vec{1}_0}^z \frac{dw}{w^2} \underline{\Pi}_l(w)$$

- The initial condition can be obtained as a generating series using techniques from mirror symmetry: [Bönisch, Fischbach Klemm, Nega, Safari]

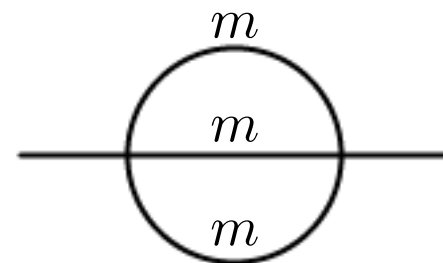
$$L_{l,k}^{(0)} = (l+1)! r_k \quad \sum_{k=0}^{\infty} r_k x^k = -\frac{\Gamma(1-x)}{\Gamma(1+x)}$$

- Compact analytic formula for arbitrary number of loops!
  - ➔ Integrals that appear in it are the natural generalisation of the integrals for  $l \leq 3$ .
  - ➔ Ingredients are of geometric origin (periods, intersection pairing, initial condition from mirror symmetry).



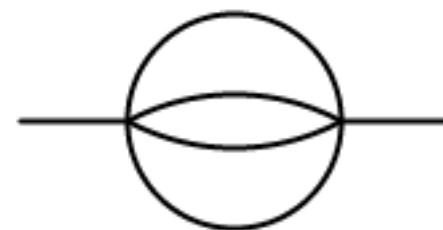
Riemann sphere

$$\sim \frac{z}{\sqrt{1-4z}} \int_0^z \frac{dz'}{z' \sqrt{1-4z'}}$$



Elliptic curve

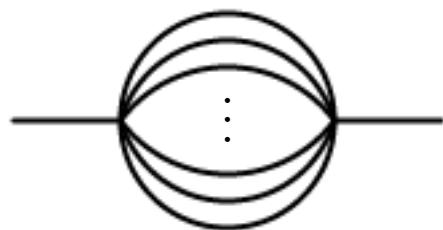
$$\sim h_1(\tau) \int_{i\infty}^{\tau} \frac{d\tau'}{2\pi i} h_3(\tau') \tau'$$



K3 surface  
~ “elliptic x elliptic”  
(only here!)

$$\sim h_2(\tau) \int_{i\infty}^{\tau} \frac{d\tau'}{2\pi i} h_4(\tau') \tau'^2$$

⋮



Calabi-Yau  
( $l-1$ )-fold

$$\sim \underline{\Pi}_l(z)^T \underline{\Sigma}_l \int_{\vec{1}_0}^z \frac{dw}{w^2} \underline{\Pi}_l(w)$$

- The three-loop case is special!
- The geometry is a 1-parameter family of K3 surfaces (=CY 2-folds).
- The Picard-Fuchs operator  $\mathcal{L}_3$  has degree 3, but it is the symmetric square of  $\mathcal{L}_2$  (after a suitable change of variables).

[Verrill; Bloch, Kerr, Vanhove; Amadeo, Primo]

➔ Solution of  $\mathcal{L}_3$  are products of 2 solutions of  $\mathcal{L}_2$ :

$$(\Pi_{3,1}, \Pi_{3,2}, \Pi_{3,3}) \sim (\Pi_{2,1}^2, \Pi_{2,1}\Pi_{2,2}, \Pi_{2,2}^2)$$

Modular forms

- ➔ We get the same class of functions as at 2 loops!
- This is only true at 3 loops in 2D for equal masses!
- ➔ One can check this explicitly for  $l = 4, 5, \dots$

● Question 1: What about  $D = 2 - 2\epsilon$  ?

- ➔ We can get the differential equations and the initial conditions for arbitrary loops. [Bönisch, CD, Fischbach Klemm, Nega]
- ➔ We can write down similar all-loop formulas for higher terms in the expansion.
- ➔ We get iterated integrals involving Calabi-Yau periods, and their derivatives.

$$\partial_z \underline{L}_l(z, \epsilon) = \tilde{\mathbf{B}}_l(z, \epsilon) \underline{L}_l(z, \epsilon) + \tilde{\underline{N}}_l(z, \epsilon)$$

$$\tilde{\underline{N}}_l(z, \epsilon) = \mathbf{W}_l(z)^{-1} \underline{N}_l(z, \epsilon) \quad \text{linear in periods}$$

$$\tilde{\mathbf{B}}_l(z, \epsilon) = \mathbf{W}_l(z)^{-1} [\mathbf{B}_l(z, \epsilon) - \mathbf{B}_{l,0}(z)] \mathbf{W}_l(z) = \mathcal{O}(\epsilon)$$

bilinear in periods



- The resulting integrals may even be structurally different.  
**Example:**
  - ➔ At 2 loops we get iterated integrals of Eisenstein series for  $\Gamma_1(6)$ , to all orders in  $\epsilon$ . [Adams, Weinzierl]
  - ➔ At 3 loops in 2D we get iterated integrals of Eisenstein series for  $\Gamma_1(6)$ . [Bloch, Kerr, Vanhove; Brödel, CD, Dulat, Penante, Tancredi]
  - ➔ At 3 loops we get iterated integrals of meromorphic modular forms for  $\Gamma_1(6)$ , to all order in  $\epsilon$ . [Brödel, CD, Matthes, to appear]
- Iterated integrals of meromorphic modular forms for  $SL_2(\mathbb{Z})$  were very recently (2021!) introduced by Nils Matthes.
  - ➔ The theory can be generalised to subgroups  $\Gamma \subseteq SL_2(\mathbb{Z})$  (where  $\Gamma$  has genus 0) [Brödel, CD, Matthes, to appear]

- **Question 2:** Is there a unique Calabi-Yau variety attached to a Feynman integral, and what is it?

➔ There are different families of Calabi-Yau varieties that describe the banana integrals: [Bönisch, Fischbach Klemm, Nega, Safari]

$$M_{l-1}^{\text{HS}} = \{ \underline{x} \in \mathbb{P}^l \mid \mathcal{F}(1, \underline{z}; \underline{x}) = 0 \} \quad \text{F-polynomial from Feynman parameter integral}$$

$$M_{l-1}^{\text{CI}} = \left\{ \left( w_1^{(i)} : w_2^{(i)} \right) \in \mathbb{P}_{(i)}^1, \forall i \mid P_1 := \sum_{i=1}^{l+1} a^{(i)} w_1^{(i)} + b^{(i)} w_2^{(i)} = \sum_{i=1}^{l+1} c^{(i)} w_1^{(i)} + d^{(i)} w_2^{(i)} =: P_2 = 0 \right\}$$

- ➔ For  $l = 2$  they are the same, but for higher  $l$  they are distinct (e.g., they have distinct Euler characteristic).
- ➔ But they do define the same Calabi-Yau motive!
- ➔ Calabi-Yau motive  $\sim$  linear subspace of a cohomology group (+other conditions).

- **Question 3:** Which Feynman integrals can be attached to families of Calabi-Yau varieties/motives? All of them?
  - ➔ There are examples of Feynman integrals whose maximal cuts give rise to Riemann surfaces of genus  $g > 1$ .

[Huang, Zhang; Hauenstein, Huang, Mehta, Zhang]
  - ➔ They cannot be Calabi-Yau varieties! (because they have  $g > 1$  holomorphic differentials)
  - ➔ There are examples of Calabi-Yau motives that describe Riemann surfaces of higher genus.
- **Interesting question for the future:** Are the examples of higher-genus curves from Feynman integrals associated to Calabi-Yau motives?

- After elliptic curves Calabi-Yau varieties are the next frontier.
- We have shown how results from the elliptic case generalise nicely to all loops for equal-mass banana integrals.
  - ➔ Compact analytic formula valid for arbitrary number of loops!
- Interesting questions for the future:
  - ➔ What about other integrals (e.g., with more scales)?
  - ➔ Which integrals are associated to Calabi-Yau motives?
  - ➔ What are the properties of the (iterated) integrals of Calabi-Yau periods?