

Nonrelativistic Open String Theory

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Collaborations

1. Nonrelativistic Open Strings and Yang-Mills Theory ([2020.xxxxx](#))
2. A Triality for Relativistic, Nonrelativistic
and Noncommutative Open Strings ([2020.xxxxx](#))

with **Jaume Gomis, Matthew Yu** ([Perimeter Institute](#))

Open Strings in Nonrelativistic String Theory

- ▶ longitudinal D-brane: noncommutative open string (NCOS)
[Seiberg, Susskind, Toumbas '00] [Gopakumar, Maldacena, Minwalla, Strominger '01]
- ▶ transverse D-brane: nonrelativistic open string (NROS)
[Danielsson, Guijosa, Kruczenski '01]

We will focus on NROS.

gauge theories from NROS?

- ▶ Galilean electrodynamics
- ▶ Galilean Yang-Mills
- ▶ Galilean DBI
- ▶ ...

Nonrelativistic Closed Strings

- ▶ **2D foliation**

$$X^\mu = (X^A, X^{A'}) \quad A = 0, 1 \quad A' = 2, \dots, d$$

- ▶ **light-cone coordinates**

- ▶ **spacetime** $X = X^0 + X^1 \quad \bar{X} = X^0 - X^1$

- ▶ **worldsheet** $\partial = \partial_\sigma - i\partial_\tau \quad \bar{\partial} = \partial_\sigma + i\partial_\tau$

- ▶ **nonrel. string theory** [Gomis, Ooguri '00] [Danielsson, Guijosa, Kruczenski '00]

$$S_0 = \frac{1}{4\pi\alpha'} \int_{\Sigma} d^2\sigma \left(\partial_\alpha X^{A'} \partial^\alpha X^{A'} + \lambda \bar{\partial} X + \bar{\lambda} \partial \bar{X} \right)$$

- ▶ **closed string spectrum** $p_0 = \frac{1}{wR} (\alpha' p^{A'} p^{A'} + N + \bar{N}) , \quad N - \bar{N} = wK$

Open String Sector

- ▶ transverse sector is in form the same as in rel. string theory
- ▶ boundary conditions in longitudinal sector?
- ▶ to be concrete
 - ▶ boundary at $\sigma = 0$
 - ▶ $\partial_\sigma X^{A'}|_{\partial\Sigma} = 0$ (**N**)

Noncommutative Open String Theory (NCOS)

- ▶ Neumann in X^A
- ▶ nonzero electric field

$$\frac{E}{2\pi\alpha'} \int_{\partial\Sigma} d\tau i X^0 \partial_\tau X^1$$

- ▶ noncommutativity $[X^A(\tau), X^B(\tau')] \sim iE^{-1} \epsilon^{AB} sgn(\tau - \tau')$
- ▶ relativistic open string spectrum

[Gomis, Ooguri '00]

Nonrelativistic Open String Theory (NROS)

- ▶ Dirichlet in X^1 and Neumann in X^0
- ▶ boundary variation of the action

$$\delta S = \frac{1}{2\pi\alpha'} \int_{\partial\Sigma} d\tau \left[\delta X^{A'} \partial_\sigma X^{A'} + \frac{1}{2}(\lambda + \bar{\lambda}) \delta X^0 + \frac{1}{2}(\lambda - \bar{\lambda}) \delta X^1 \right]$$

- ▶ Dirichlet boundary condition: $\delta X^1 = 0$ (**D**)
- ▶ varying X^0 : $\lambda + \bar{\lambda} = 0$
- ▶ open string spectrum $p_0 = \frac{1}{wR} \left(\alpha' p^{A'} p^{A'} + \frac{1}{2} N \right)$
- ▶ global spacetime symmetries on D-brane:

$$\delta X^{A'} = \Xi^{A'} + \Lambda^{A'} X^0 - \Lambda^{A'} {}_{B'} X^{B'} \quad \delta \lambda = -\Lambda_{A'} \partial X^{A'}$$

$$\delta X^0 = \Xi^0 \quad \delta X^1 = 0 \quad \delta \bar{\lambda} = -\Lambda_{A'} \bar{\partial} X^{A'}$$

⇒ Bargmann algebra

Open String Vertex Operators in NROS

- ▶ independent vertex operators on D-brane:

$$V = \frac{1}{2\pi\alpha'} \int_{\partial\Sigma} d\tau \left(\lambda N + A_0 \partial_\tau X^0 + A_{A'} \partial_\tau X^{A'} \right)$$

- ▶ $U(1)$ gauge transformation: $\delta_\epsilon A_0 = \partial_0 \epsilon$, $\delta_\epsilon A_{A'} = \partial_{A'} \epsilon$

- ▶ invariant under Bargmann:

$$N \rightarrow N \quad A_0 \rightarrow A_0 - \Lambda^{A'} A_{A'} \quad A^{A'} \rightarrow A^{A'} - \Lambda^{A'} N + \Lambda^{A' B'} A_{B'}$$

BRST Invariance (Zero Winding)

- BRST invariance $0 = [Q, V] = -\frac{\alpha'}{4} \int_{\partial\Sigma} dz \partial_z c \mathcal{E}(z)$

$$\begin{aligned}\mathcal{E}(z) = & \frac{1}{2} \left(\partial^{A'} \partial_{A'} N \right) \lambda - \left(\partial_0^2 N - \partial^{A'} E_{A'} \right) \partial X^0 \\ & - \left(\partial_0 \partial_{A'} N + \partial^{B'} F_{B'A'} \right) \partial X^{A'}\end{aligned}$$

- e.o.m.

$$\partial^{A'} \partial_{A'} N = 0 \quad \partial_0^2 N - \partial^{A'} E_{A'} = 0 \quad \partial_{A'} \partial_0 N + \partial^{B'} F_{B'A'} = 0$$

Galilean Electrodynamics

- ▶ Galilean electrodynamics (GED) on D-brane

$$S = -\frac{1}{g^2} \int dX^0 dX^{A'} \left(\frac{1}{2} \partial_0 N \partial_0 N - E_{A'} \partial_{A'} N - \frac{1}{4} F_{A'B'} F_{A'B'} \right)$$

- ▶ in the literature, GED from null reduction and nonrel. limit

[Santos, de Montigny, Khanna, Santana '04] [Bergshoeff, Rosseel, Zojer '15]

[Festuccia, Hansen, Hartong, Obers '16] ...

- ▶ quantum properties [Chapman, Di Pietro, Grosvenor, ZY, to appear]

Winding Modes

- ▶ open string vertex operator with a fixed winding number w

$$V^w(z) = \int dz \left[\frac{1}{2} N^w \lambda(z) - A_0^w \partial_z X^0(z) - A_{A'}^w \partial_z X^{A'}(z) \right] e^{iq \int^z dz' \lambda(z')}$$

where $q = wR/(4\alpha')$

- ▶ define $X^r(z) \equiv \frac{1}{2} \int^z dz' \lambda(z)$, $A_r = -N$, and $\mathcal{A} = (r, 0, A')$

$$A_{\mathcal{A}}(X^{\mathcal{B}}) = \sum_w A_{\mathcal{A}}^w(X^0, X^{A'}) e^{2iqX^r}$$

- ▶ BRST invariance $\sum_w [Q, V^w(z)] = 0$ implies

$$G^{\mathcal{B}\mathcal{C}} \partial_{\mathcal{B}} F_{\mathcal{C}\mathcal{A}} = 0 \quad G^{AB} = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & \delta^{A'B'} \end{pmatrix}$$

- ▶ $d-2$ propagating d.o.f. with $p_0 = \frac{\alpha'}{wR} p_{A'} p_{A'}$

Coinciding D-branes

- ▶ **M coinciding D-branes; zero winding**
- ▶ **adjoint gauge field $(A_0, A_{A'})$ in $U(M)$**
- ▶ **Wilson line coupled to adjoint scalar N**

$$W = \mathbf{tr} \mathcal{P} \exp \left[\frac{1}{2\pi\alpha'} \int_{\partial\Sigma} d\tau \left(\lambda N + iA_0 \partial_\tau X^0 + iA_{A'} \partial_\tau X^{A'} \right) \right]$$

- ▶ **At the lowest order in field strength: Galilean Yang-Mills**

$$S_{\text{YM}} = -\frac{1}{g_{\text{YM}}^2} \int dX^0 dX^{A'} \mathbf{tr} \left(\frac{1}{2} D_0 N D_0 N - E_{A'} D_{A'} N - \frac{1}{4} F_{A'B'} F_{A'B'} \right)$$

Curved Backgrounds (Abelian)

- ▶ target space coordinates $X^\mu \quad \mu = 0, 1, \dots d - 1$
- ▶ D-brane worldvolume coordinates $Y^i \quad i = 0, 1, \dots d - 2$
- ▶ embedding of D-brane in target space $X^\mu|_{\partial\Sigma} = f^\mu(Y^i)$
- ▶ Dirichlet nonlinear sigma model on flat worldsheet

$$S = \frac{1}{4\pi\alpha'} \int_{\Sigma} d^2\sigma \left(\partial X^\mu \bar{\partial} X^\nu (H_{\mu\nu} + B_{\mu\nu}) + \lambda \bar{\partial} X^\mu \tau_\mu + \bar{\lambda} \partial X^\mu \bar{\tau}_\mu \right)$$

$$+ \frac{1}{2\pi\alpha'} \int_{\partial\Sigma} d\tau \left[\frac{1}{2} N(Y) (\lambda - \bar{\lambda}) + i A_i(Y) \partial_\tau Y^i \right]$$

$$H_{\mu\nu} = E_\mu{}^{A'} E_\nu{}^{A'} + (\tau_\mu{}^A m_\nu{}^B + \tau_\nu{}^A m_\mu{}^B) \eta_{AB}$$

- ▶ dilaton only contributes classically at the lowest order in α'

Spontaneous Symmetry Breaking: Unbroken Phase

- **boundary conditions** ($t_i{}^A = \tau_\mu{}^A \partial_i f^\mu$, $\mathcal{F}_{ij} = \partial_i f^\mu B_{\mu\nu} \partial_j f^\nu + F_{ij}$)

$$N(Y) = 0 \quad \partial_\sigma X^\mu \tau_\mu{}^A = i \epsilon^A{}_B \partial_\tau Y^i t_i{}^B$$

$$\partial_i f^\mu H_{\mu\nu} \partial_\sigma X^\nu + i \mathcal{F}_{ij} \partial_\tau X^j + \frac{1}{2} (\lambda t_i + \bar{\lambda} \bar{t}_i) = 0$$

- **string Newton-Cartan symmetry**

$$\delta \tau_\mu{}^A = \Lambda \epsilon^A{}_B \tau_\mu{}^B \quad \delta E_\mu{}^{A'} = -\Lambda_A{}^{A'} \tau_\mu{}^A + \Lambda^{A'}{}_{B'} E_\mu{}^{B'}$$

$$\delta m_\mu{}^A = D_\mu \sigma^A + \Lambda \epsilon^A{}_B m_\mu{}^B + \Lambda^{AA'} E_\mu{}^{A'} - \tau_\mu{}^B \sigma^A{}_B, \quad \sigma^A{}_A = 0$$

$$\delta A_i = -\epsilon_{AB} \sigma^A t_i{}^B$$

Spontaneous Symmetry Breaking: Broken Phase

- ▶ coordinates adapted to the submanifold $X^\mu = (y, Y^i)$
- ▶ for NROS

$$\tau_y{}^0 = E_y{}^{A'} = 0 \quad \tau_y{}^1 \neq 0$$

- ▶ uniformly broken phase

$$\langle f^\mu(Y^i) \rangle = f_0^\mu(Y^i) \quad f_0^y = y_0 \quad f_0^i = Y^i$$

- ▶ Nambu-Goldstone boson $N(Y^i)$

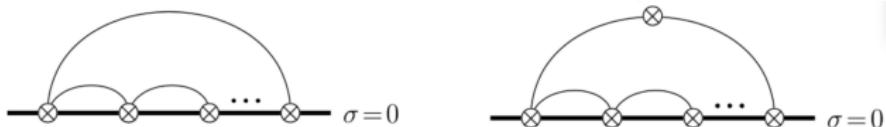
$$f^y(Y^i) = y_0 + N(Y^i) \quad f^i(Y^i) = Y^i$$

- ▶ Bargmann symmetry on D-brane

$$\begin{aligned} \delta E_i{}^{A'} &= \Lambda^{A'} \tau_i{}^0 + \Lambda^{A'} {}_{B'} E_i{}^{B'} & \delta m_i{}^0 &= \partial_i \sigma^0 + \Lambda^{A'} E_i{}^{A'} \\ \delta A_i &= -\sigma^0 (\tau_i{}^1 + \partial_i N) & \delta m_y{}^0 &= \partial_y \sigma^0 \end{aligned}$$

Boundary Beta-Functions and D-brane Action

- beta-functions for N and A_i



$$\beta(N) = \alpha' [J^{ij} K^\mu{}_{ij} \tau_\mu^{-1} + t_{A'}^{-1} (\frac{1}{2} J^{ia} \mathcal{F}_a{}^j \mathcal{H}_{ijA'} - \partial_{A'} \Phi)]$$

$$\begin{aligned} \beta_k(A) = \alpha' & \left\{ J^{ri} \nabla_k t_i^{-1} + J^{ij} \left[\nabla_{(i} \mathcal{F}_{j)k} + (K^\mu{}_{ij} \mathcal{B}_{\mu k} + K^\mu{}_{k(i} \mathcal{B}_{j)\mu}) \right] \right. \\ & \left. - \left(\frac{1}{2} J^{ia} \mathcal{F}_a{}^j \mathcal{H}_{ij\rho} - \partial_\rho \Phi \right) (\epsilon^A{}_B \tau^\rho{}_A t_k{}^B - E^\rho{}_{A'} \mathcal{F}_{A'k}) \right\} \end{aligned}$$

where $\mathcal{F}_{ij} \equiv \mathcal{B}_{ij} + F_{ij}$, $\mathcal{B}_{\mu\nu} = B_{\mu\nu} + 2 m_{[\mu}{}^A \tau_{\nu]}{}^B \epsilon_{AB}$ and

$$J^{ab} = (G_{ab} - \mathcal{F}_{ac} \mathcal{F}^c{}_b)^{-1}, \quad G_{ab} = \begin{pmatrix} 0 & t_j{}^0 \\ t_i{}^0 & e_{ij} \end{pmatrix}, \quad \mathcal{F}_{ab} = \begin{pmatrix} 0 & t_j{}^1 \\ -t_i{}^1 & \mathcal{F}_{ij} \end{pmatrix}$$

- worldvolume action for D-brane

$$S_{\text{brane}} \propto \int d^{d-1} Y e^{-\Phi} \sqrt{-\det \begin{pmatrix} 0 & \partial_i f^\mu \tau_\mu \\ \partial_i f^\mu \bar{\tau}_\mu & \partial_i f^\mu \partial_j f^\nu (H_{\mu\nu} + B_{\mu\nu}) + F_{ij} \end{pmatrix}}$$

Outlooks

- ▶ nonabelian generalization of Galilean DBI
- ▶ SUSY and analog of Dirac's equation?
- ▶ quantum properties: coupling GED/GYM to fermions?
- ▶ black hole/brane, AdS/CFT...

Thank you!