

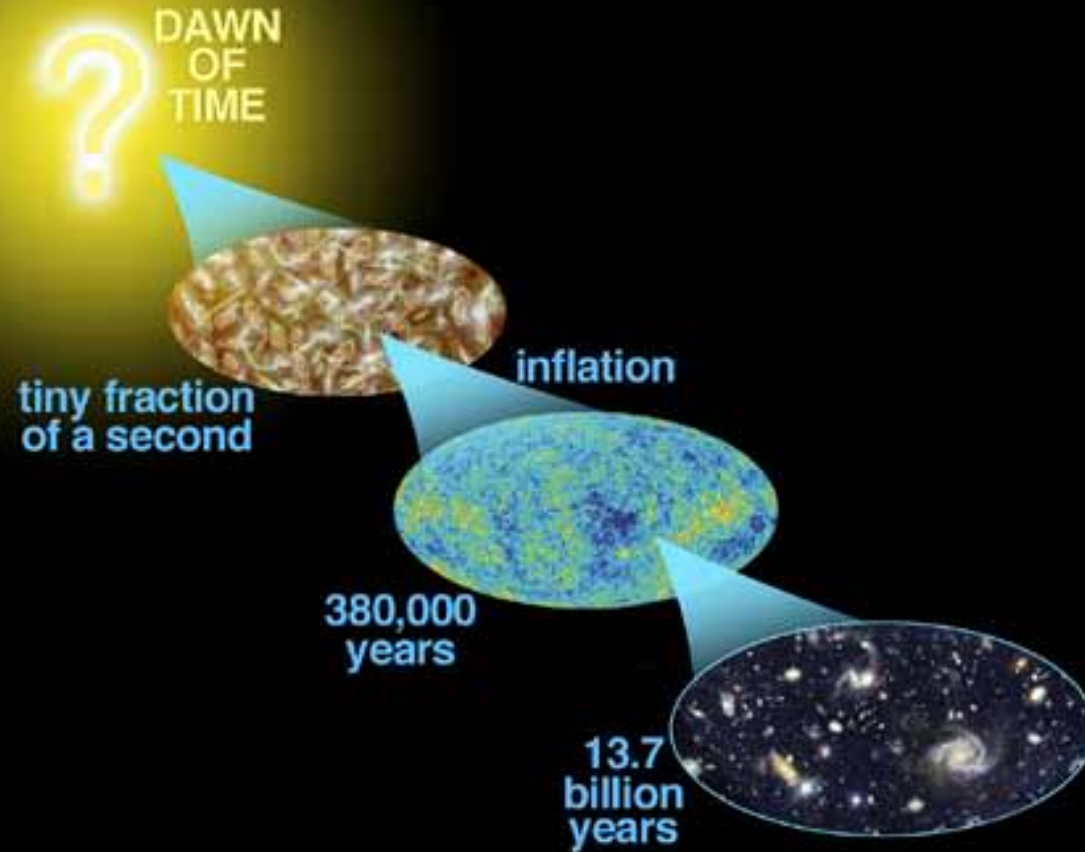
A minimal SM/LCDM cosmology

Neil Turok

Higgs Centre, University of Edinburgh
and
Perimeter Institute for Theoretical Physics

with Latham Boyle

current consensus



Inflation was a groundbreaking idea but

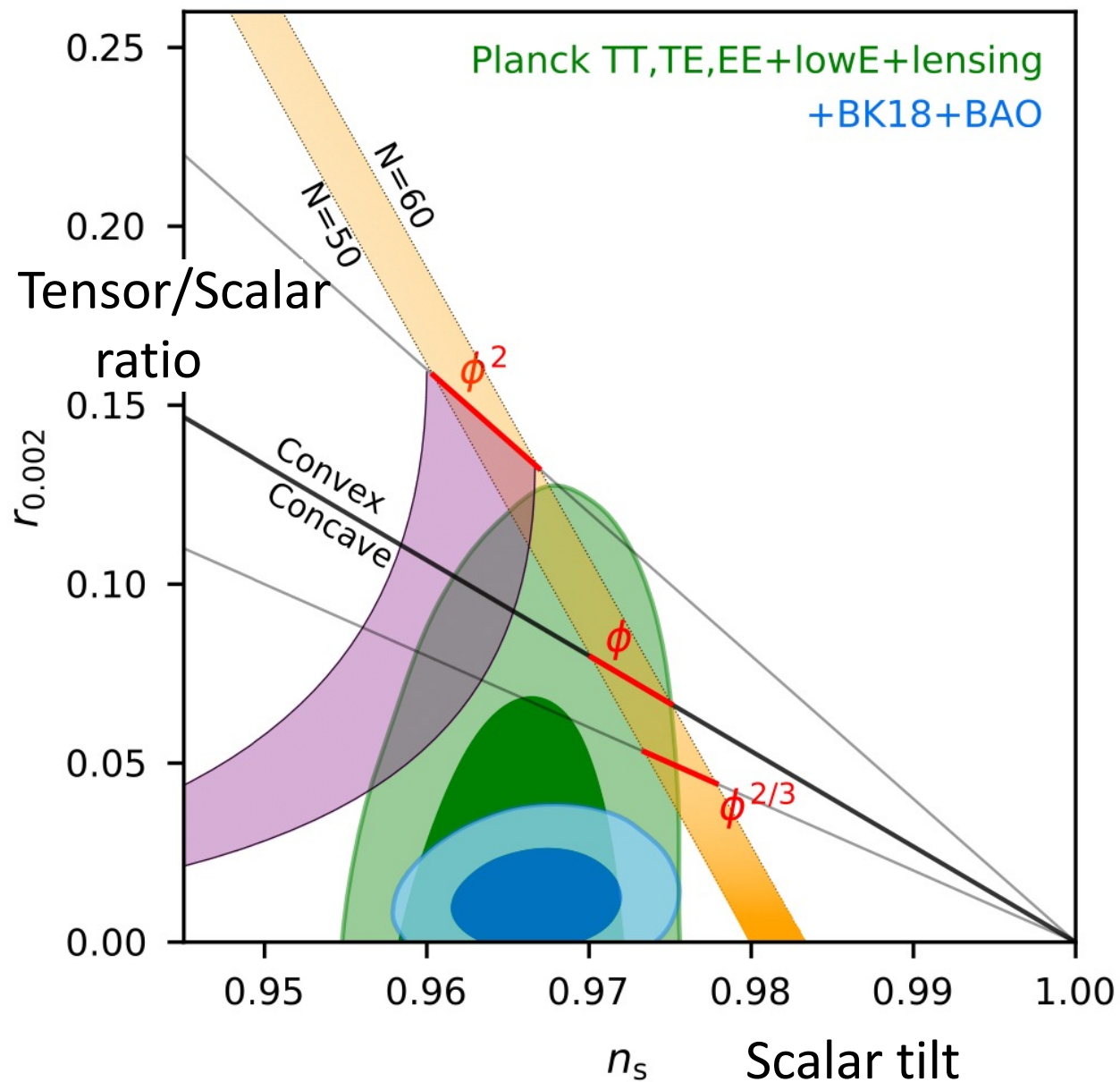
1. Observations are very well fit by vanilla LCDM
cf. overabundance of inflationary models
2. No sign of inflation's “smoking gun” signal: long wavelength gravitational waves
(Current bound is $r < 0.03$; CMB experimenters project $r < 0.003$ by 2027)
3. No satisfactory measure on inflationary universes

In this talk (and in the spirit of Bohr's atomic model)
I will attempt to outline a comprehensive alternative

no sign of
inflationary
tensors

BICEP/Keck
Collaboration
2203.16556
[astro-ph]

anticipated limit
 $r < .003$
using SPT for
“delensing”
(2027)



vanilla LCDM:

just 5 fundamental physics parameters

matter/energy content

1. ρ_Λ cosmological constant
2. ρ_{DM}/ρ_B DM/baryon density
3. n_B/n_γ baryons per photon

fluctuations

Newtonian potential

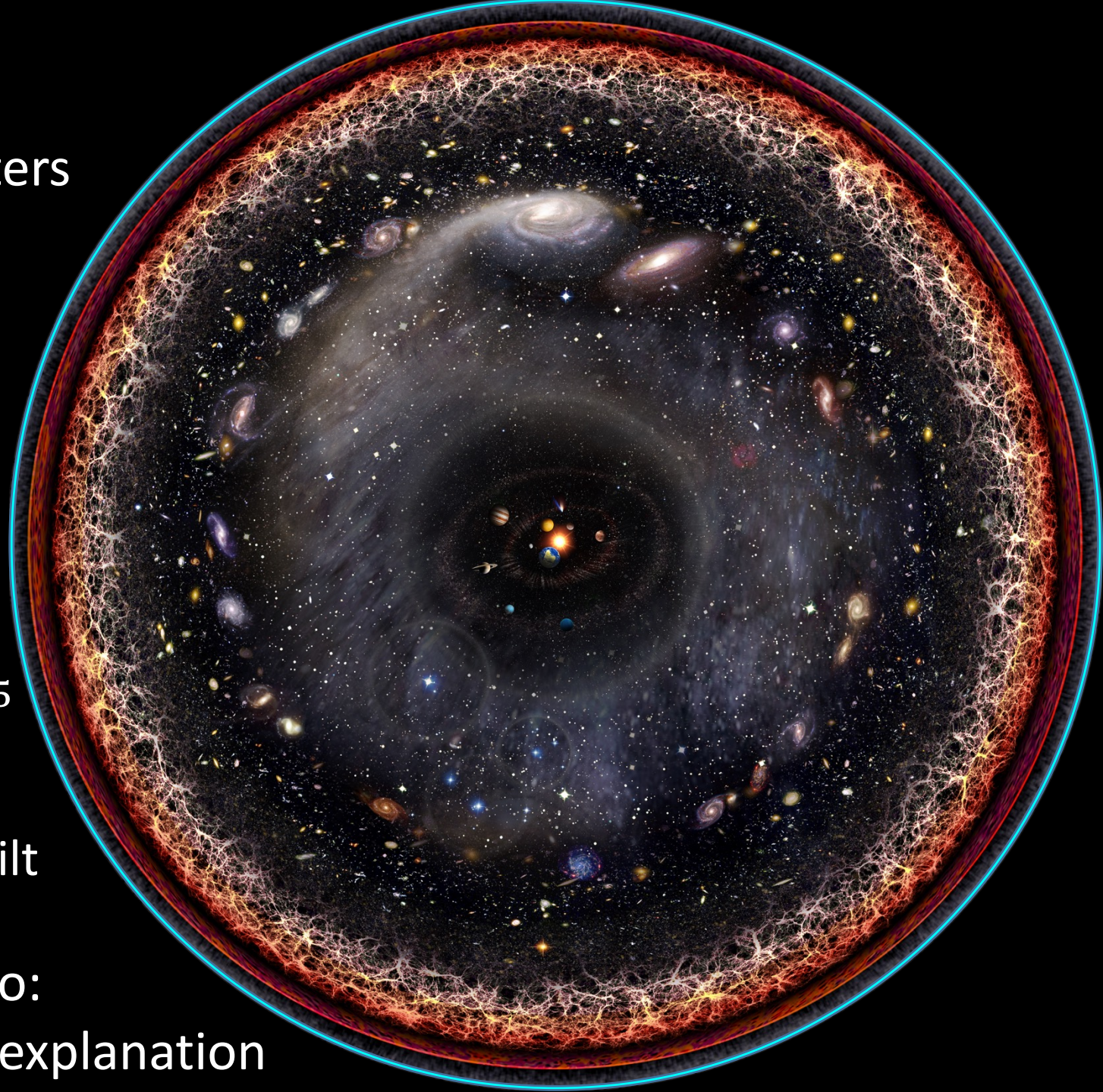
$$\langle \Phi^2 \rangle = \int \frac{dk}{k} A \left(\frac{k}{k_*} \right)^{n_s-1} \quad (k_* \equiv 0.05 \text{Mpc}^{-1})$$

4. large scale $\Phi_{rms} \approx \sqrt{A} \approx 3 \times 10^{-5}$

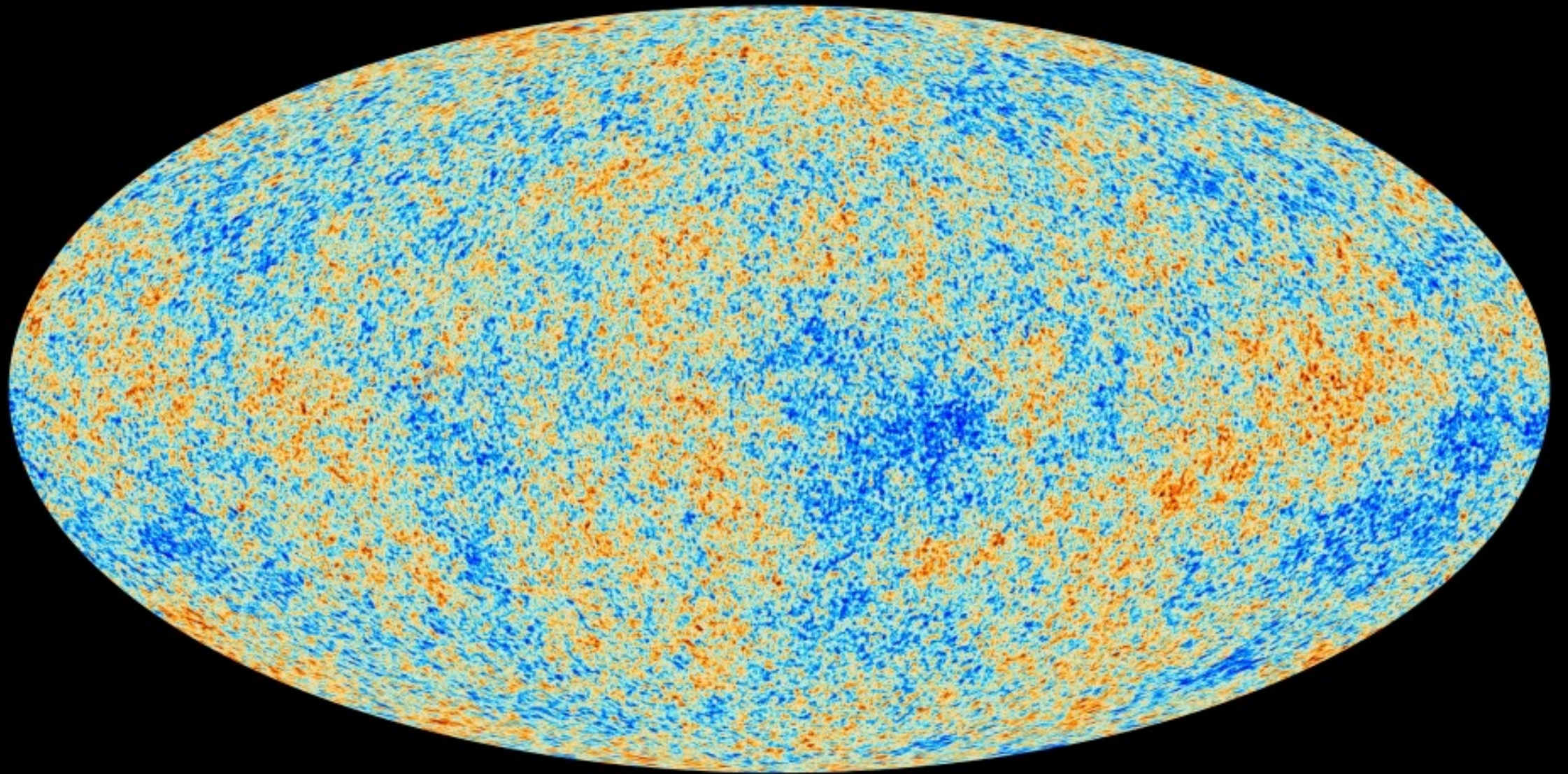
Sachs-Wolfe $\delta T/T \approx \frac{1}{3} \Phi \approx 10^{-5}$

5. $n_s - 1 \approx -0.04 \pm .006$ small red tilt

many quantities consistent with zero:
suggests looking for an economical explanation

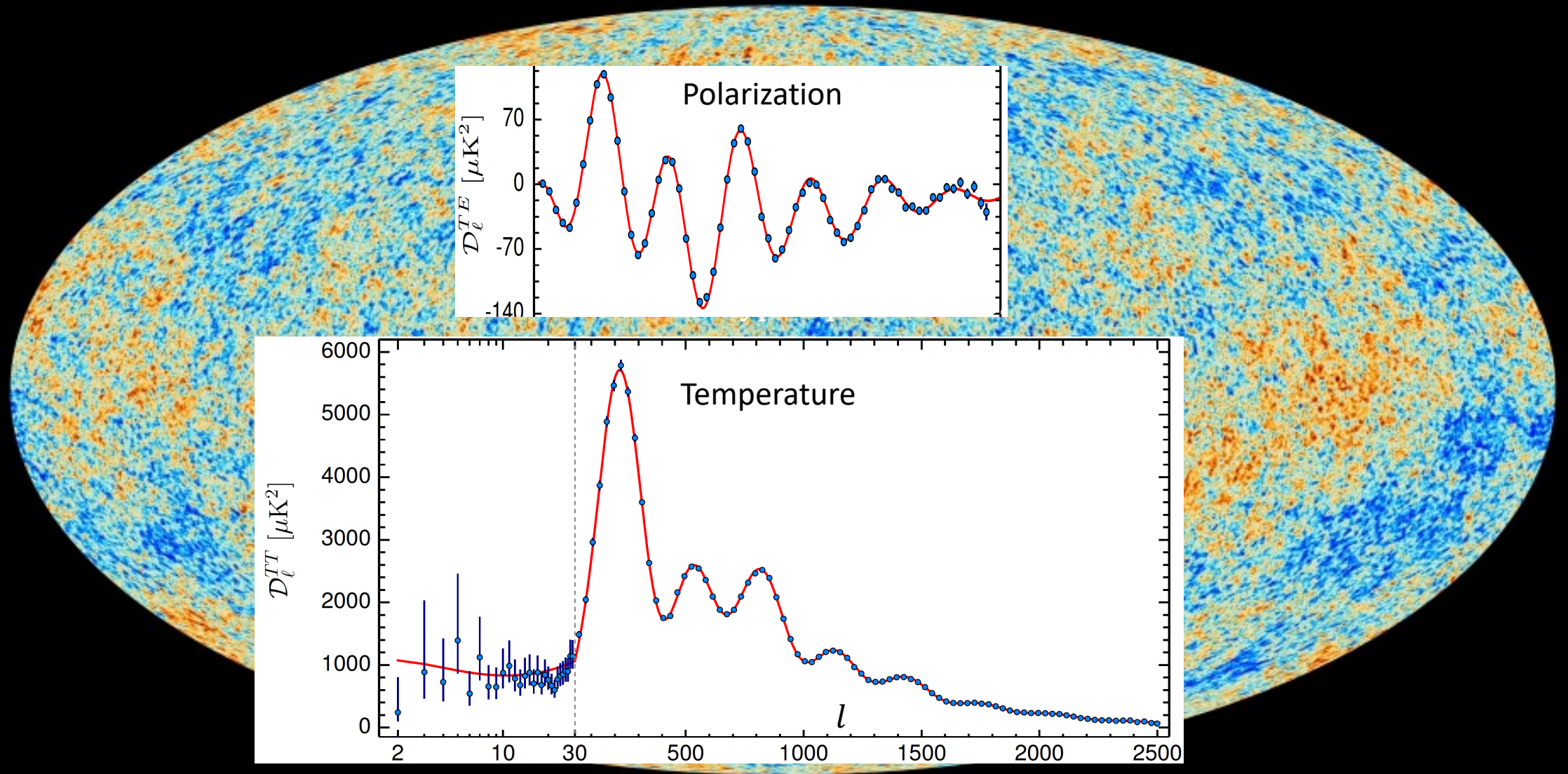


Looking back to the bang



ESA Planck satellite

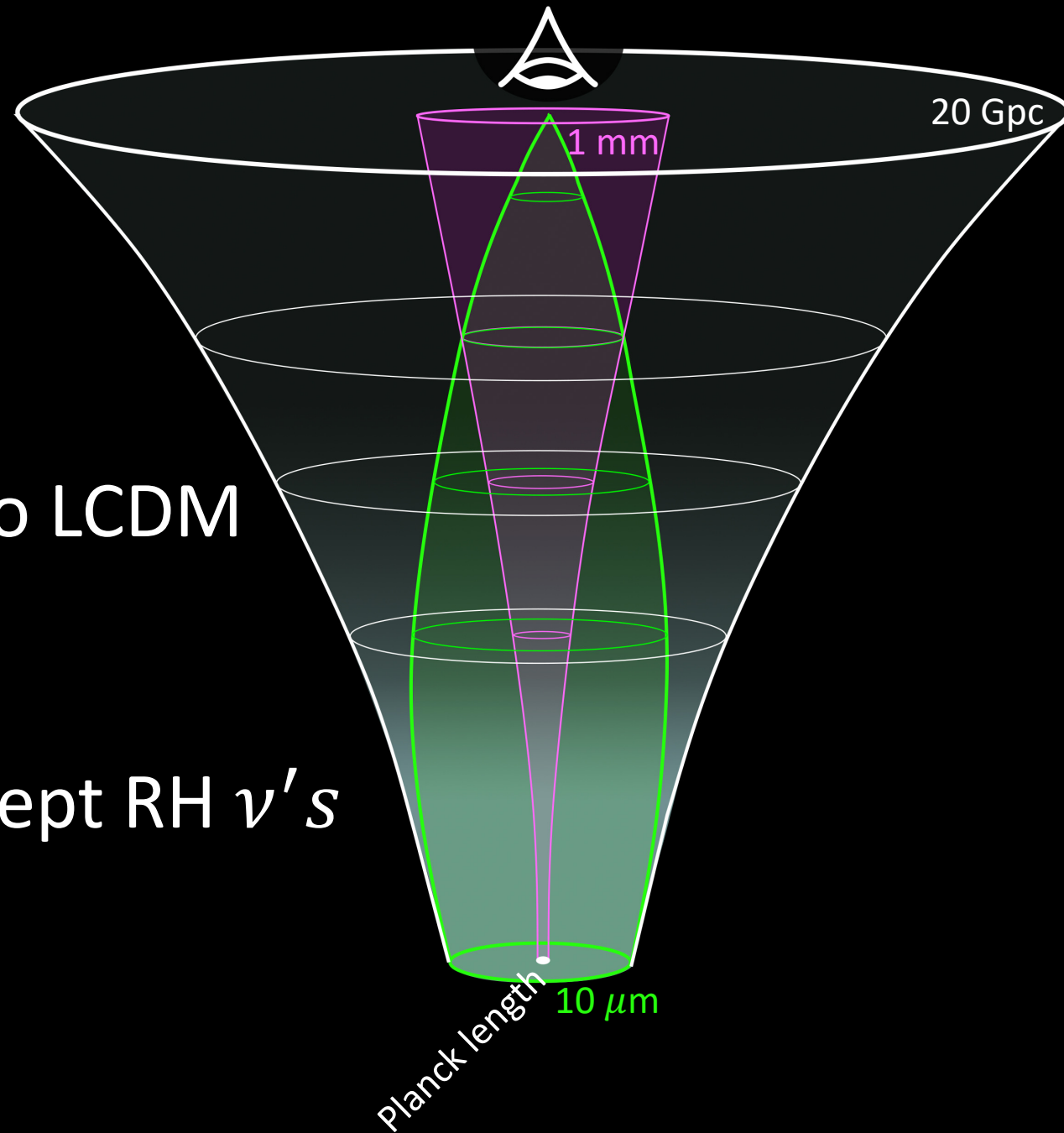
Large scale perturbations



this talk:

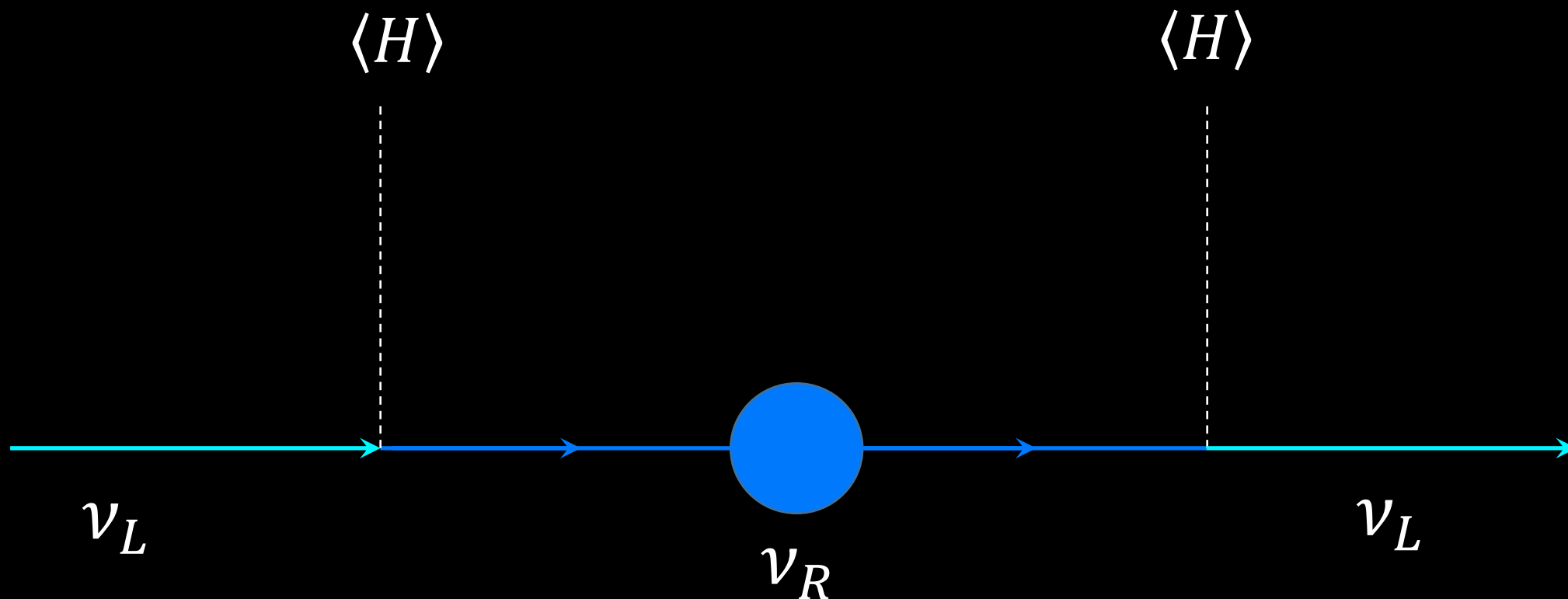
a new framework
connecting the SM to LCDM
(all 5 parameters)

no new particles except RH ν 's
no need for inflation



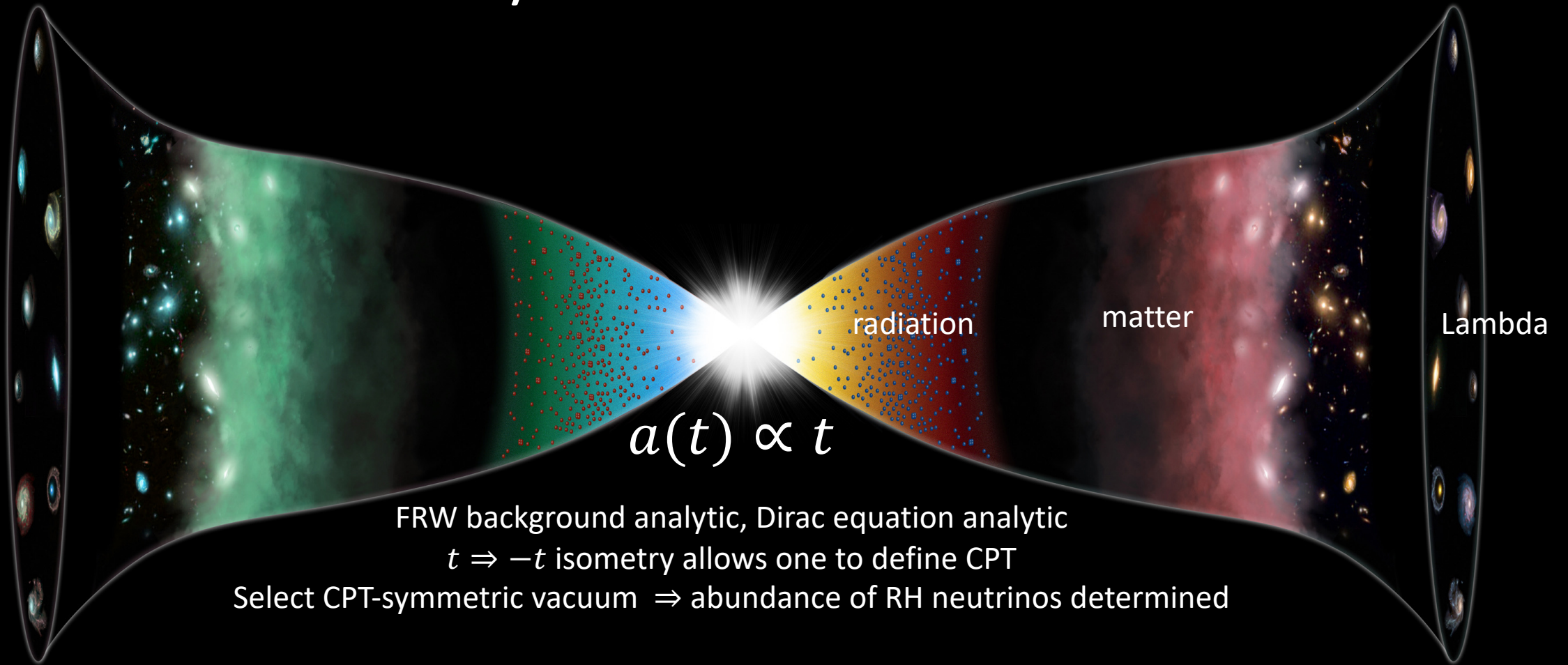
dark matter

Right-handed neutrinos:



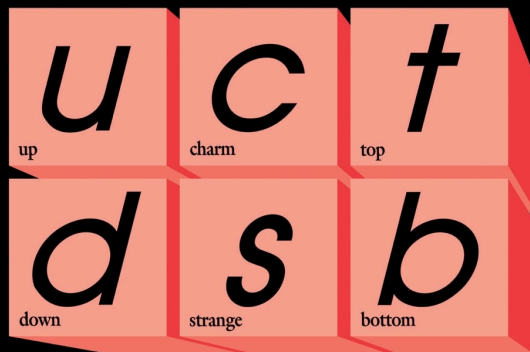
naturally explain observed neutrino masses

CPT symmetric universe

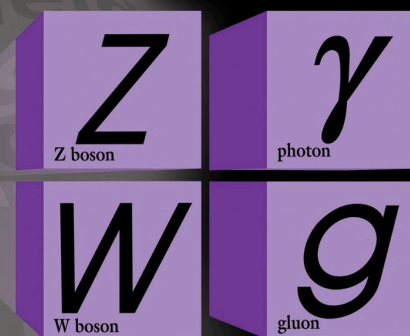


RH neutrinos produced as Hawking radiation

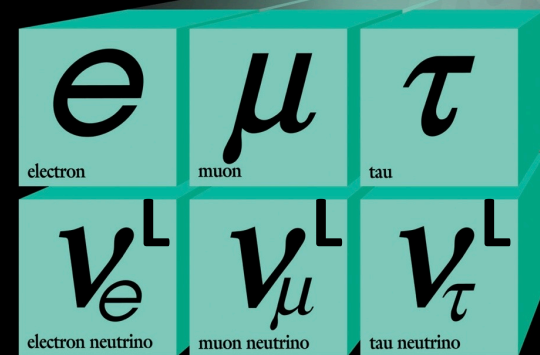
Quarks



Forces



Gravity



Leptons

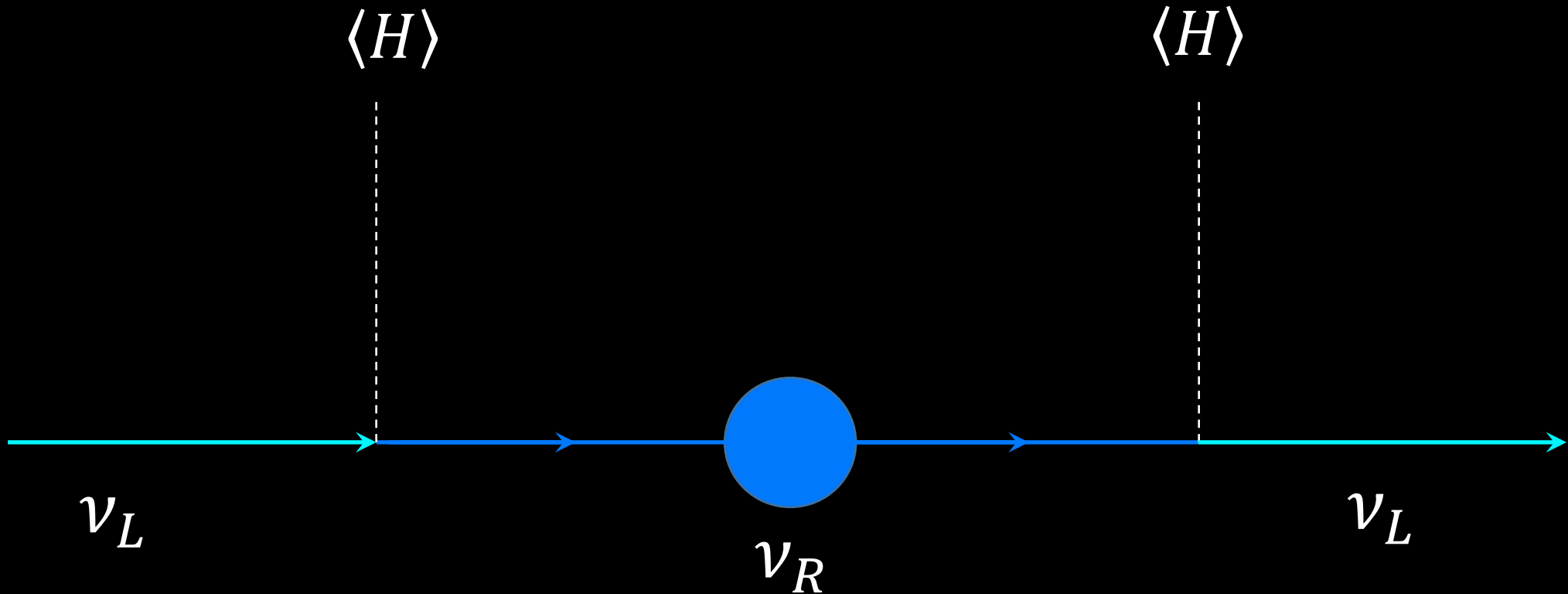


RH neutrinos

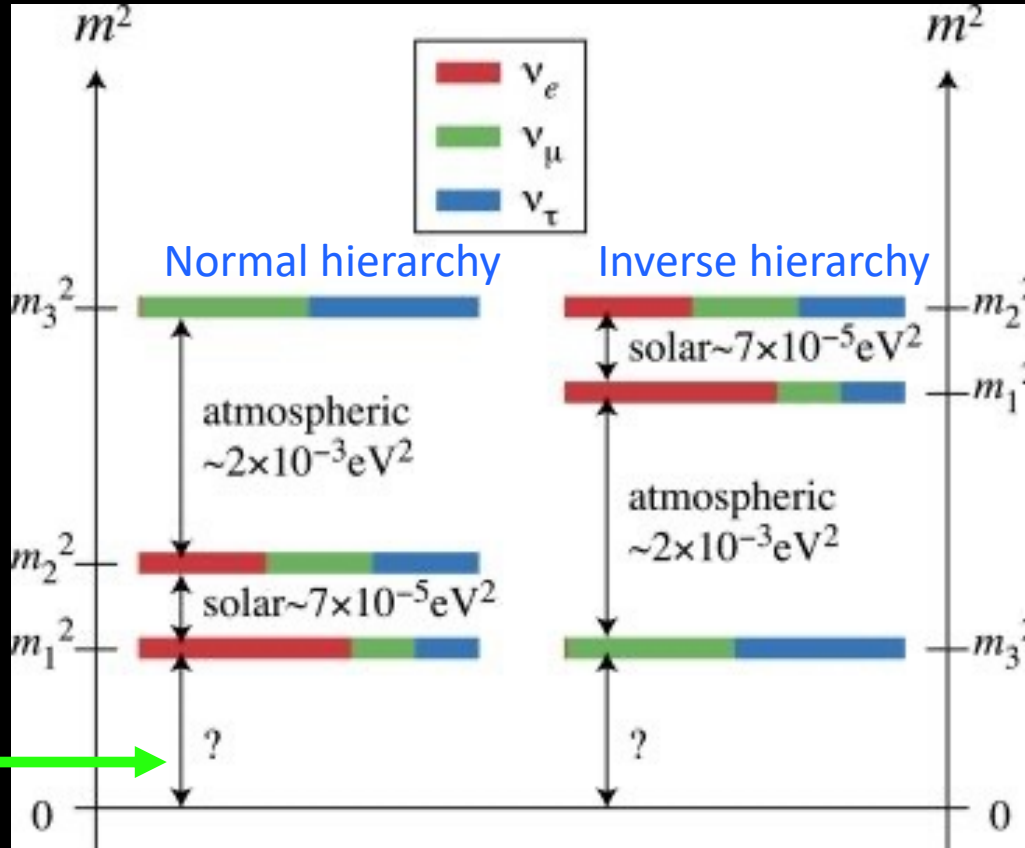
dark matter



Stability of one right-handed neutrino \Rightarrow lightest ν is massless



Light neutrinos: observations



Normal hierarchy: $M_\nu \equiv \sum m_\nu \approx 0.06 \text{ eV}$

Inverted hierarchy: $M_\nu \approx 0.1 \text{ eV}$

current data

eBOSS 2007.08991

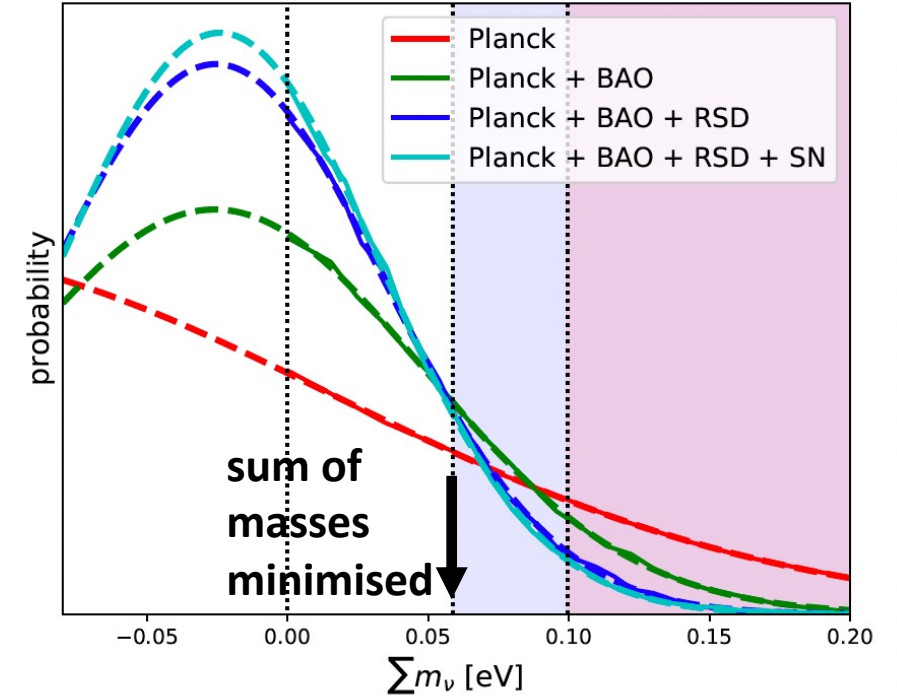
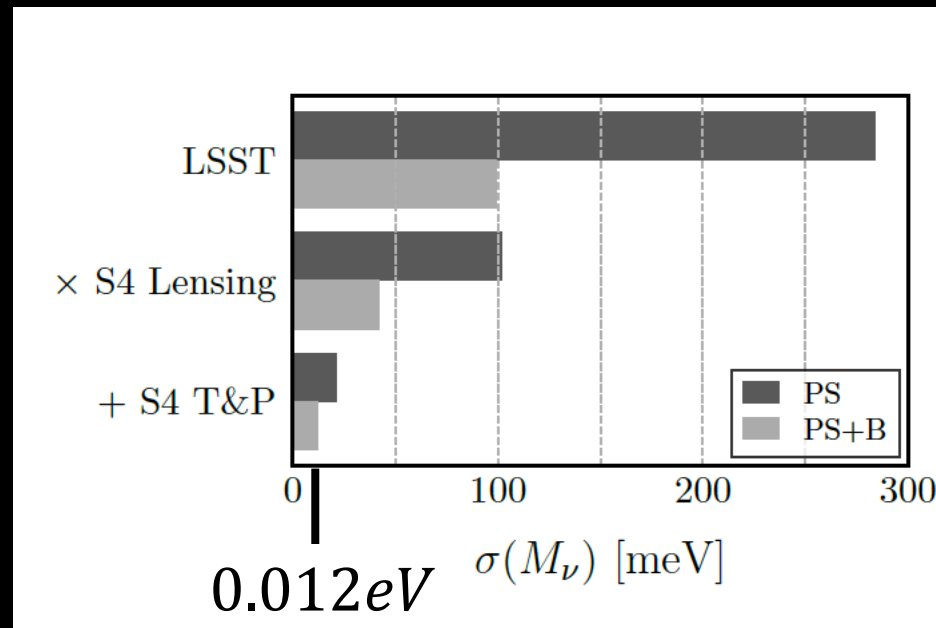
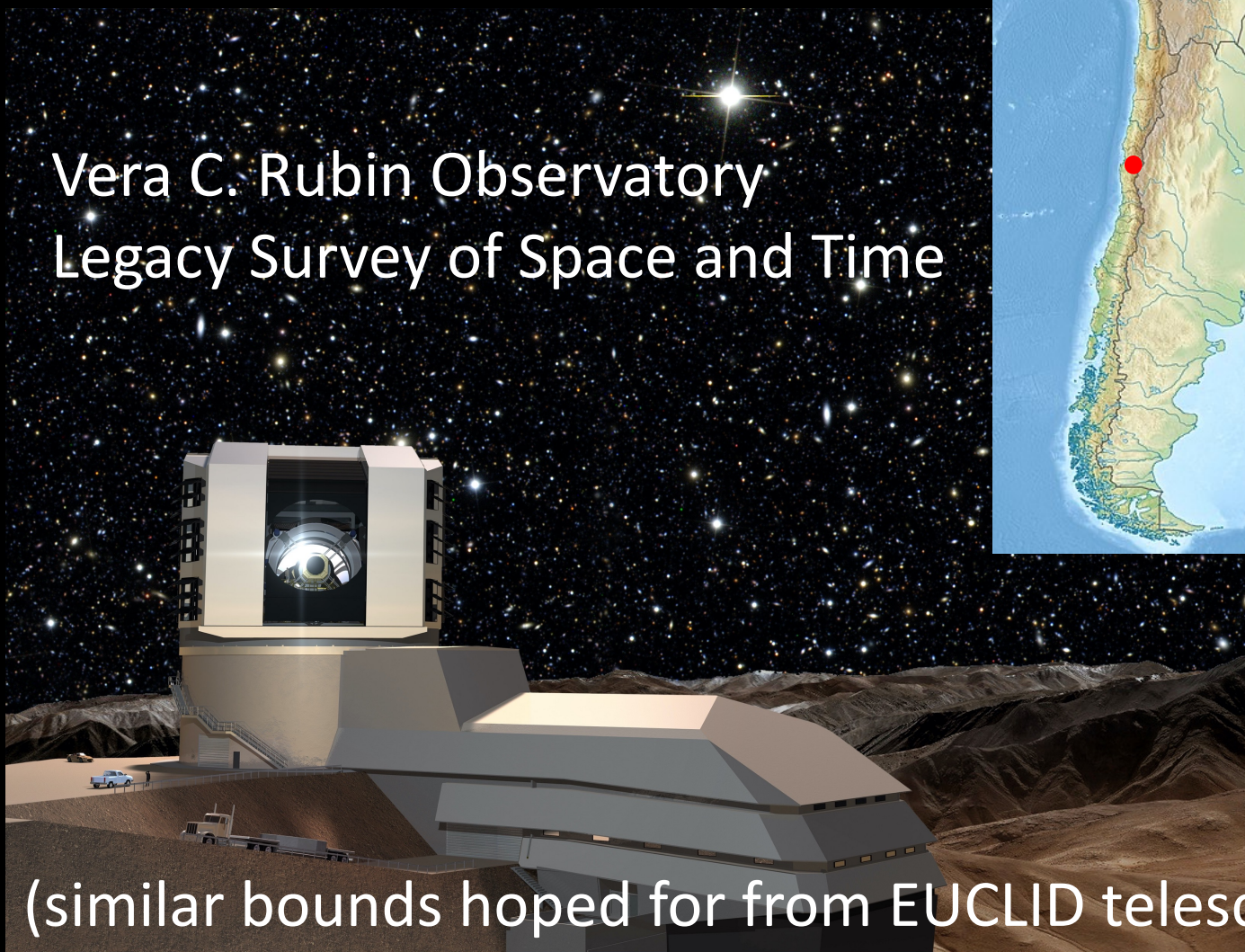
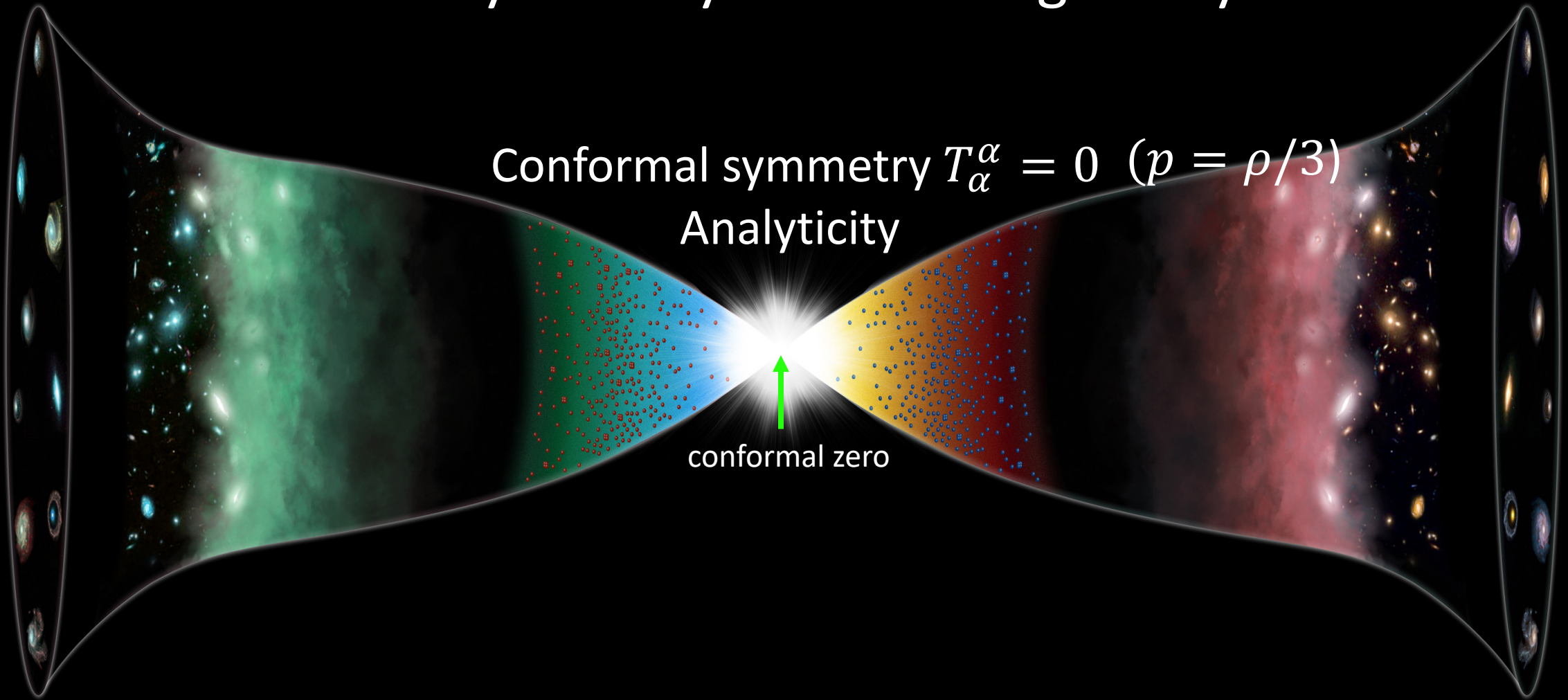


FIG. 13.— Posterior for sum of neutrino masses for selected combinations of data with a $\nu\Lambda\text{CDM}$ cosmology. Dashed curves show the implied Gaussian fits. Shaded regions correspond to lower limits on normal and inverted hierarchies. Likelihood curves are not normalized to have the same area under the curve for $\Sigma m_\nu > 0$.

Vera C. Rubin Observatory Legacy Survey of Space and Time

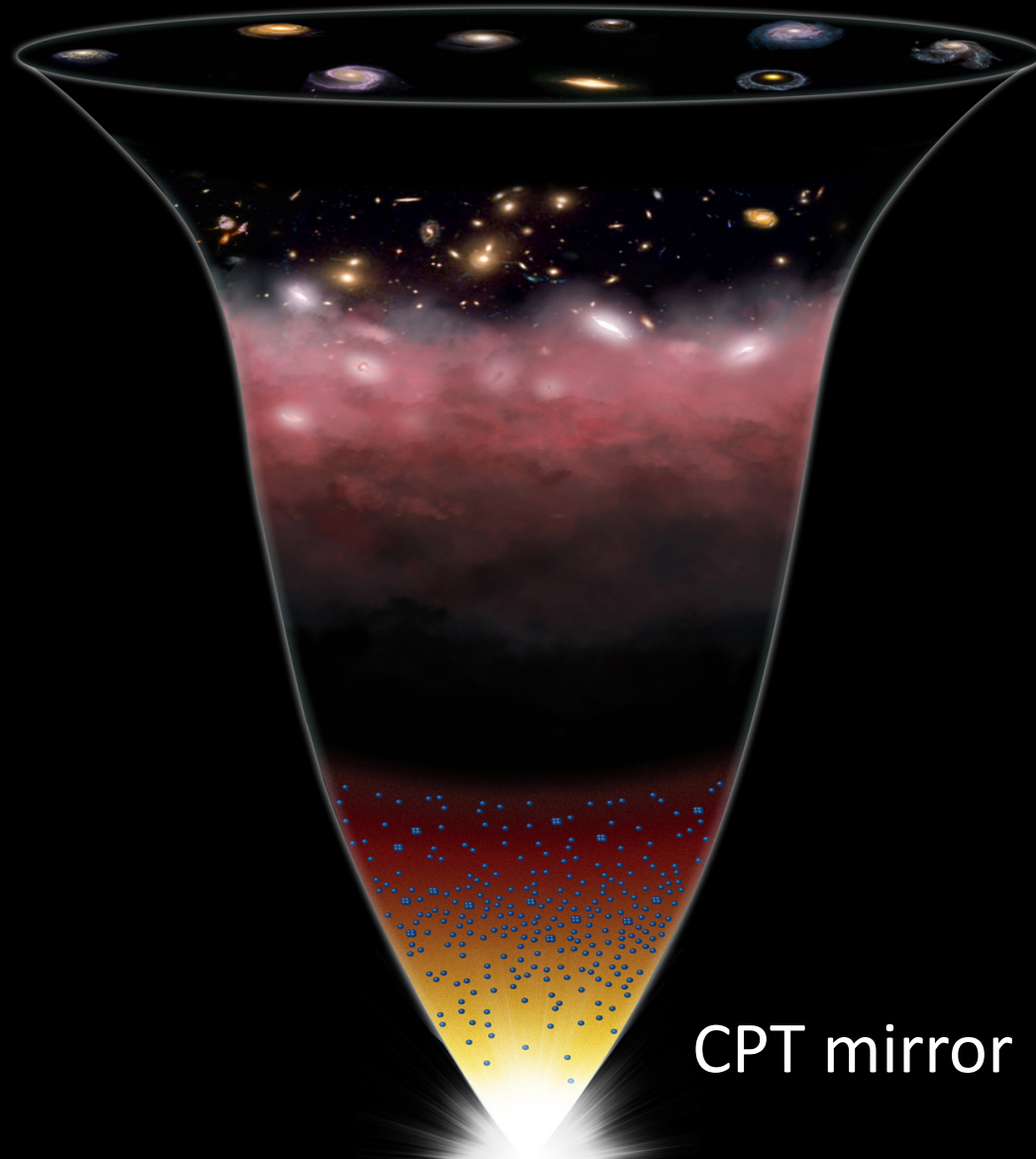


CPT symmetry and the singularity



CPT-symmetry imposed via the “method of images”

classically, the big bang is an “analytic mirror”



Striking fact: for perfect fluid with $T^\mu{}_\mu = 0$, *i.e.*, local conformal symmetry,
 $\exists \infty^3$ solutions to the Einstein equations which are analytic at $t = 0$:

$$ds^2 = t^2(-dt^2 + h_{ij}(t, \mathbf{x})dx^i dx^j); \quad h_{ij}(t, \mathbf{x}) = h_{ij}^0(\mathbf{x}) + t^2 h_{ij}^2(\mathbf{x}) + \dots$$

regular 4-metric

regular 3-metric

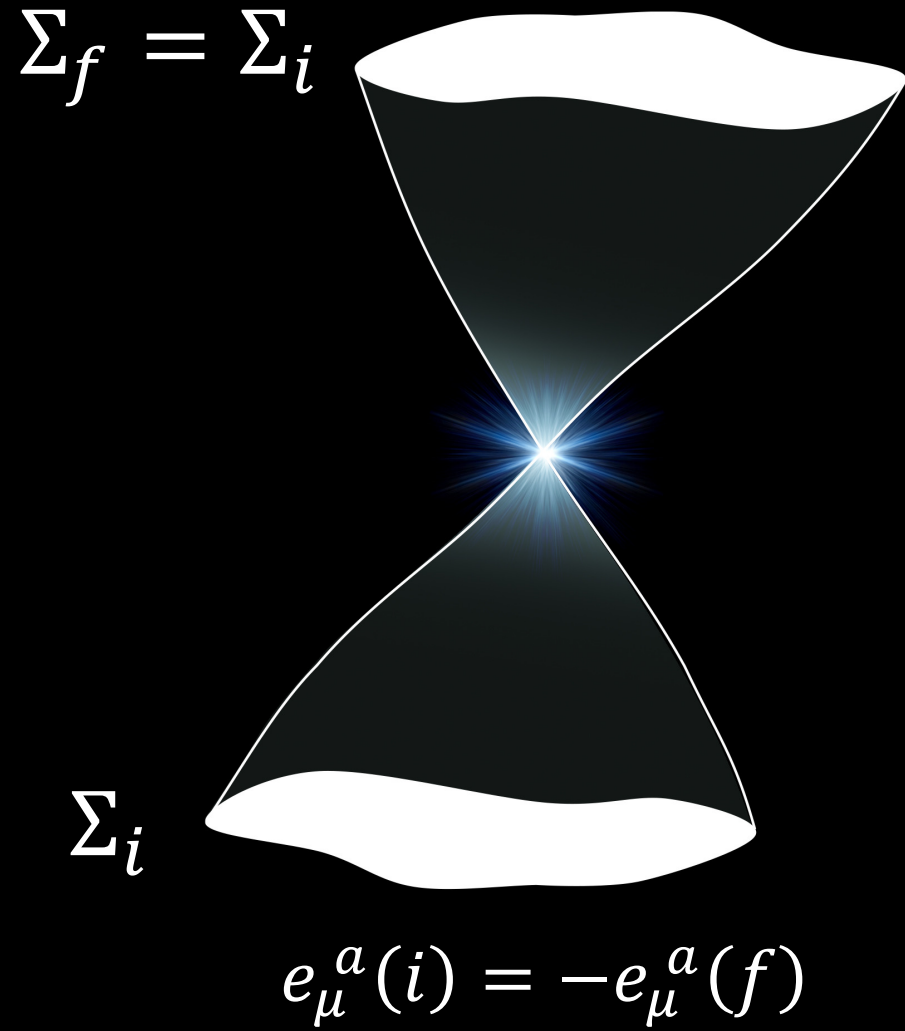
determined by
Einstein eqns

The extended spacetime is symmetric under $t \Rightarrow -t$; provides a saddle to the path integral for gravity with CPT-symmetric boundary conditions

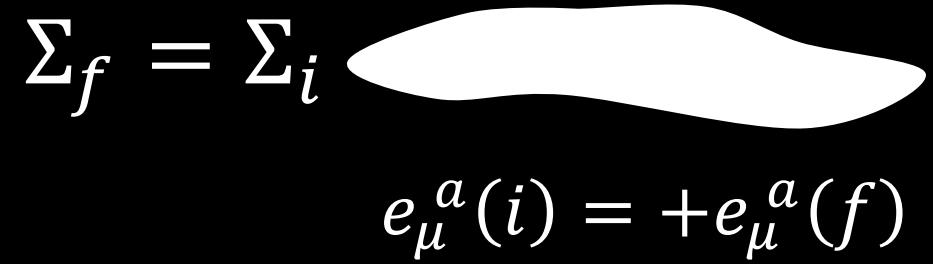
The big bang singularity is purely conformal: the Weyl tensor vanishes there
 Penrose conjecture is a consequence of the CPT-symmetric path integral

BKL or Mixmaster metrics are excluded because they are singular and hence not genuine saddles

Lorentzian path integral for gravity with CPT-symmetric boundary conditions



Past



Present

the puzzling large-scale geometry of the cosmos



Penrose

Path integrals and gravity

$$\int e^{\frac{i}{\hbar} \int \left(\frac{R}{16\pi G} - \frac{1}{4} F^2 + \bar{\psi} i \not{D} \psi - \lambda H \bar{\psi} \psi + |DH|^2 - V(H) \right)}$$

gauge fields

Higgs

gravity

particles

With pbcs in imaginary time, $Z = e^{S_g}$

gravitational entropy

partition function

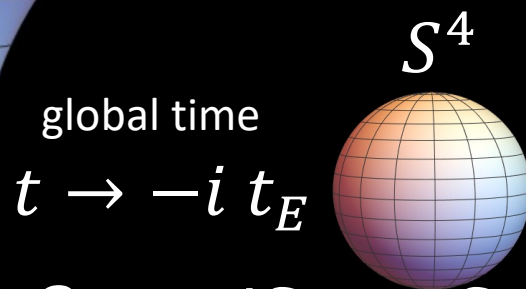
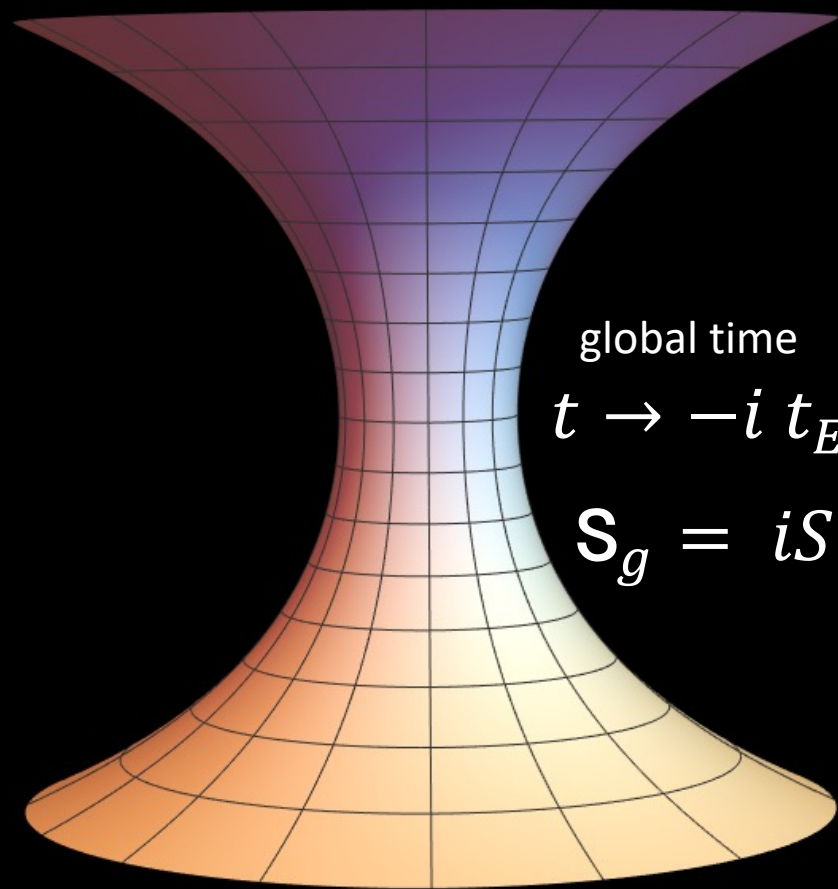
Black hole thermodynamics

Hawking
Bekenstein
Bardeen
Geroch
Gibbons
Hartle
Unruh
Wald

Hawking temperature T_H , gravitational entropy S_g

de Sitter

gravitational entropy from the Euclidean path integral



global time

$$t \rightarrow -i t_E$$

trace of Einstein

$$R = \frac{4\rho_\Lambda}{M_P^2}$$

$$\mathbf{S}_g = iS = -S_E = \int \left(\frac{1}{2} M_P^2 R - \rho_\Lambda \right) = \rho_\Lambda \text{Vol} = \frac{24\pi^2 M_P^4}{\rho_\Lambda}$$
$$\equiv \mathbf{S}_\lambda \approx 3.26 \times 10^{122} \text{ for measured } \rho_\Lambda$$

de Sitter Entropy

cf. entropy of radiation in our Hubble volume $\sim 10^{90}$

realistic cosmology:

$$ds^2 = \overset{\substack{\text{scale} \\ \text{factor}}}{a(t)^2} \left(\underset{\substack{\text{conformal} \\ \text{time}}}{-dt^2} + \underset{\substack{\text{comoving space} \\ \text{(assume compact)}}}{\gamma_{ij} dx^i dx^j} \right) \quad \overset{\text{symmetric space}}{R^{(3)} = 6\kappa}$$

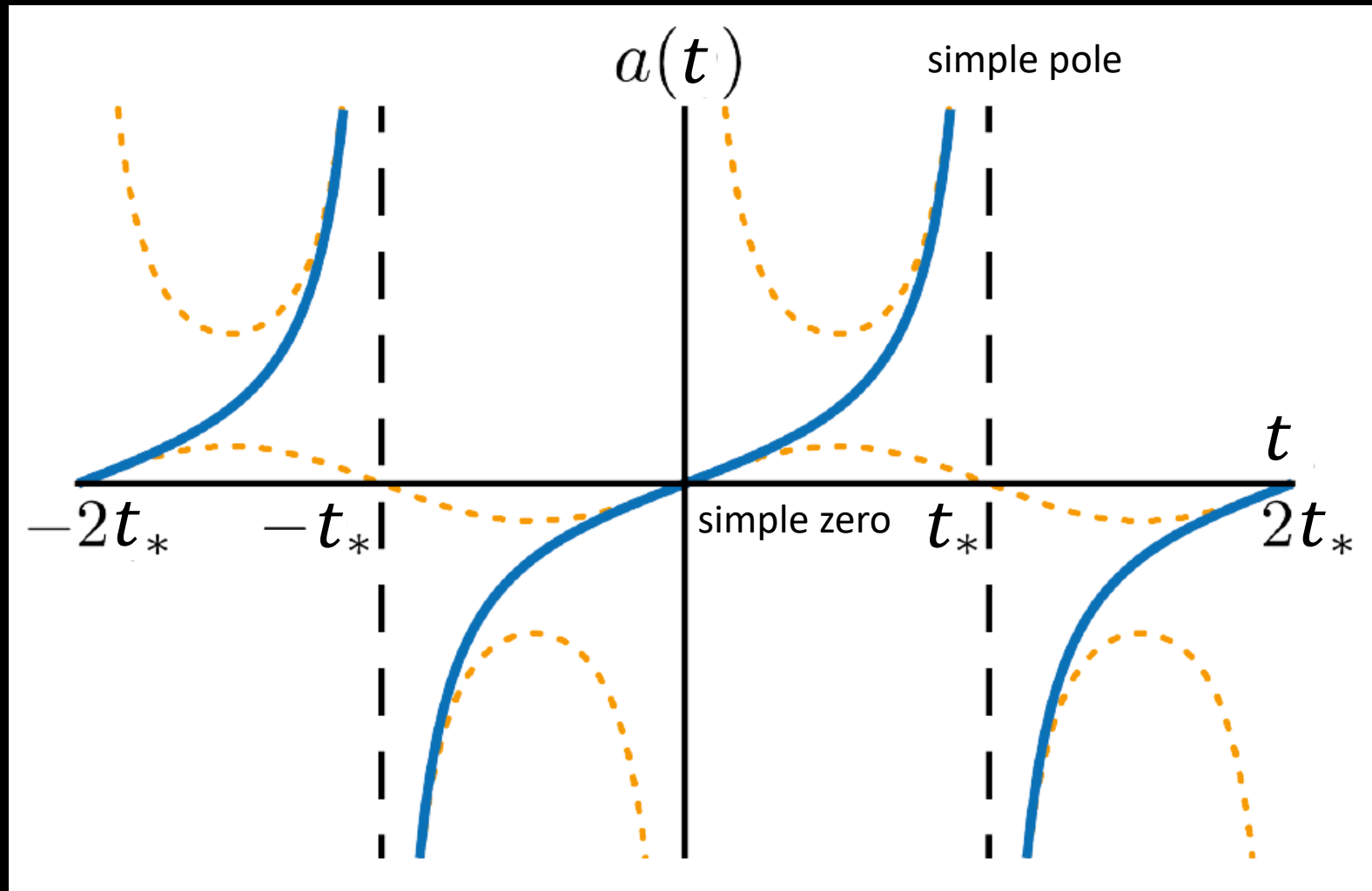
In suitable units

$$\text{Friedmann} \quad 3\dot{a}^2 = \overset{\text{radiation}}{r} + \overset{\text{matter}}{\mu} a - \overset{\text{space curvature}}{3\kappa} a^2 + \overset{\text{Lambda}}{\lambda} a^4$$

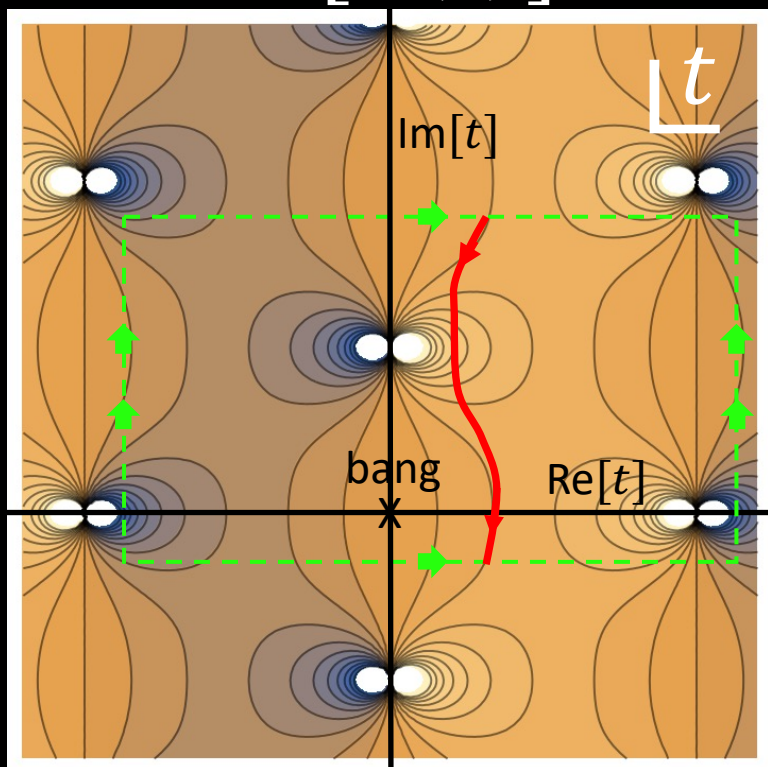
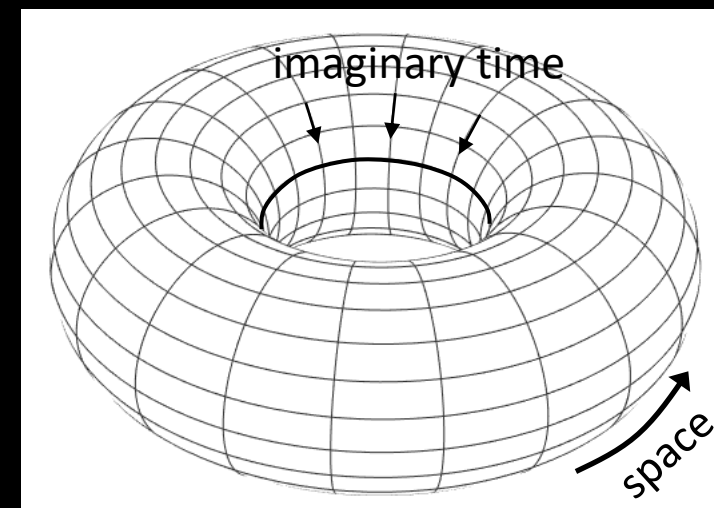
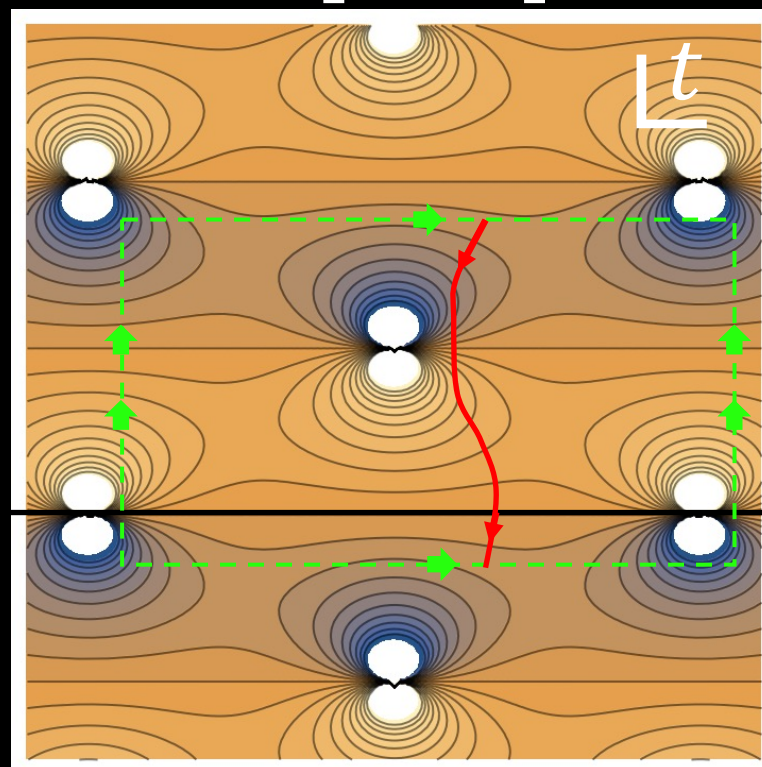
$$\overset{\text{Einstein}}{T^\mu{}_\mu = 0} \Rightarrow R = 0 \Rightarrow a(t) \text{ analytic at } t = 0$$

general solution has remarkable analytical properties

real time



$a(t)$ is single-valued and doubly periodic in the complex t -plane: its only singularities are simple poles. The imaginary time period and the action computed over a period determine T_H and the gravitational entropy S_g

 $\text{Re}[a(t)]$

 $\text{Im}[a(t)]$


Euclidean instanton for a universe w/radiation, matter, curvature, Lambda

We recently computed S_g analytically for a general cosmology with radiation, matter, space curvature and a cosmological constant (*i.e.*, all conserved quantities).

Inhomogeneities and anisotropies treated in cosmological perturbation theory.

We found that S_g is greatest for:

1. a spatially flat, homogeneous, isotropic universe
2. a small, positive cosmological constant

(echoing earlier arguments of Baum, Hawking, Coleman...)

Note:

S_g is the *global* entropy for the entire spacetime. It is a fixed number, independent of Lorentzian time (via Cauchy theorem), depending on the cosmological parameters.

compare with de Sitter entropy

$$S_g/S_\Lambda$$

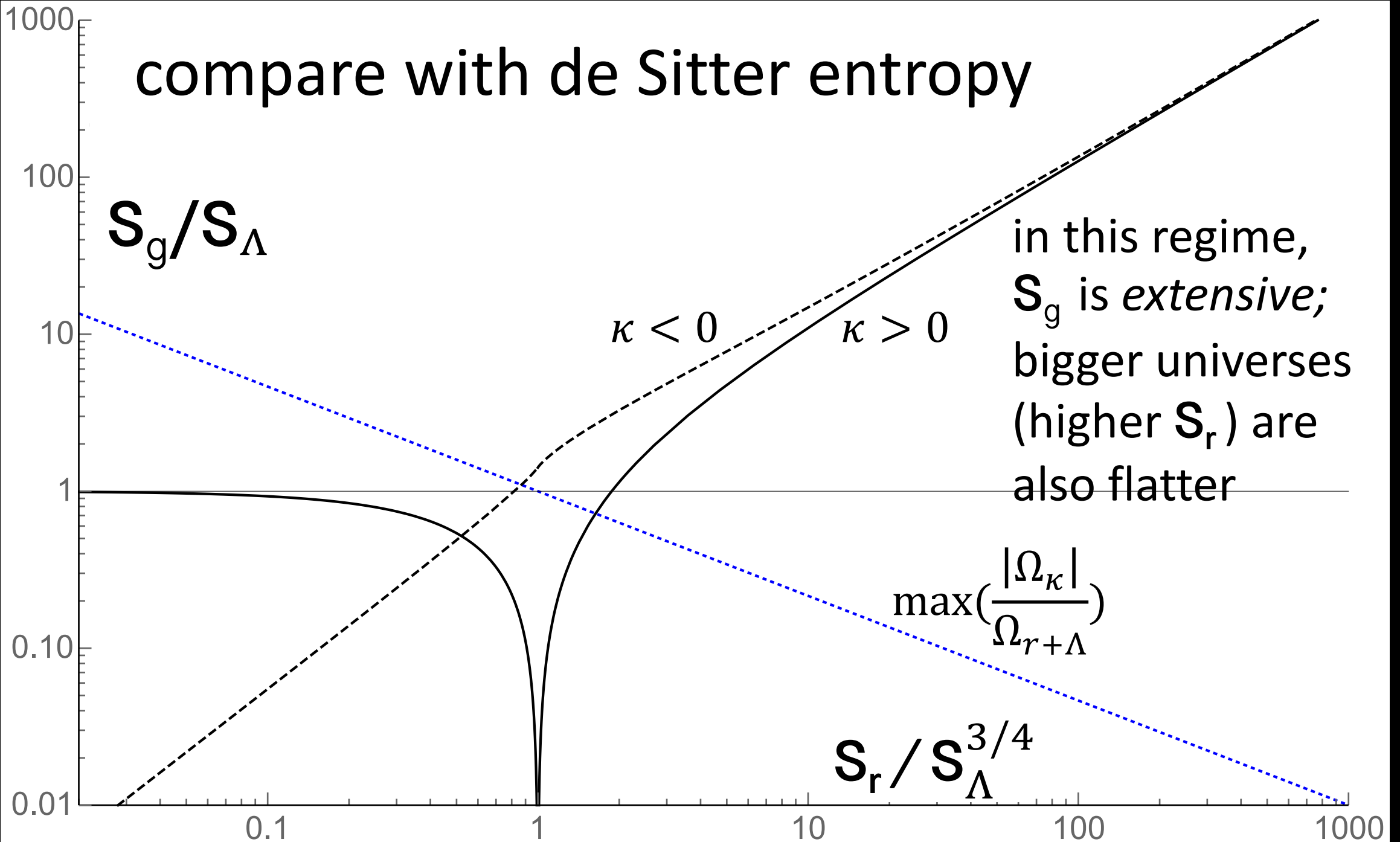
$$\kappa < 0$$

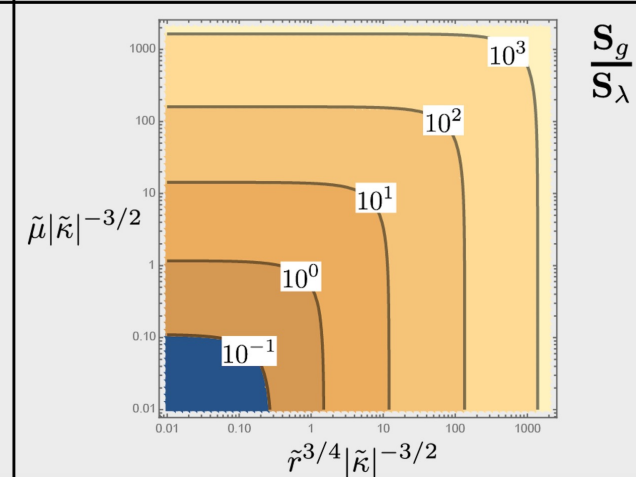
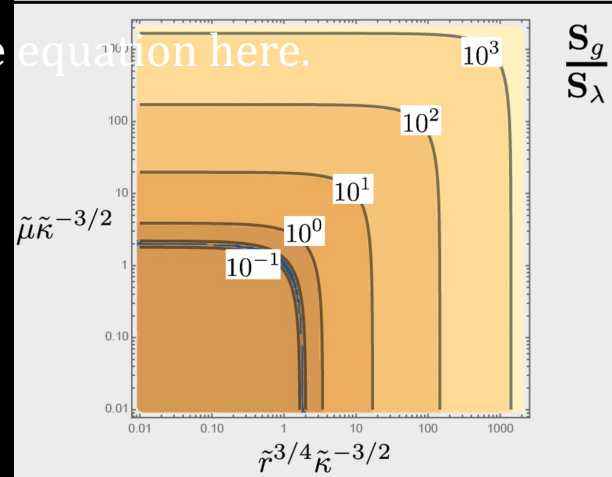
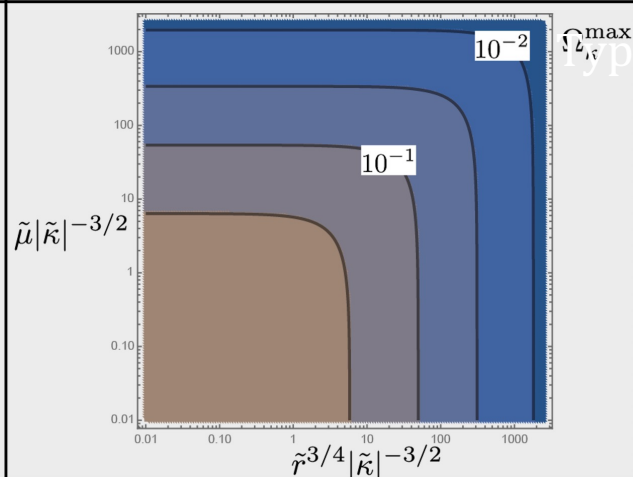
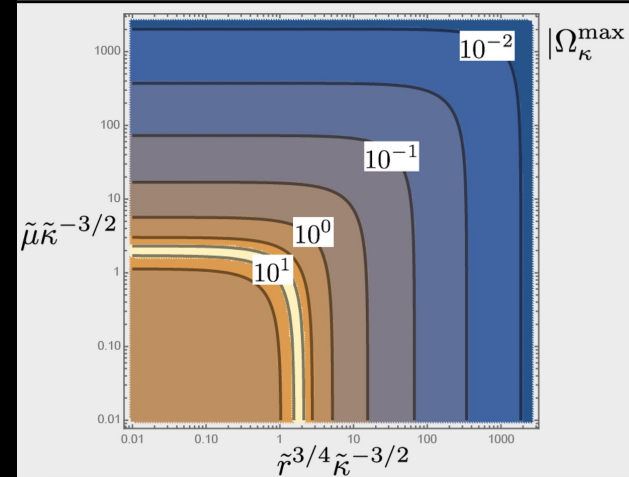
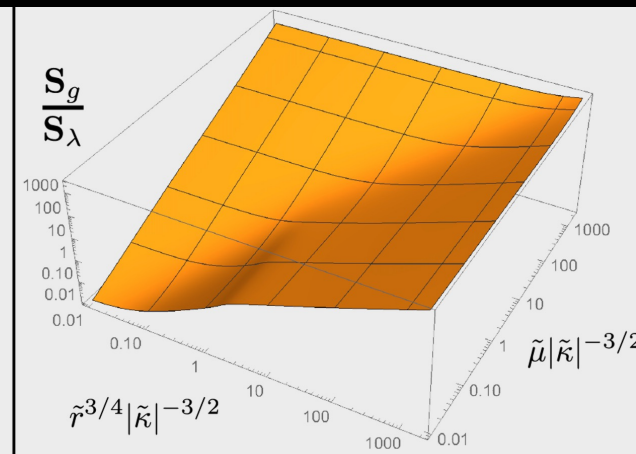
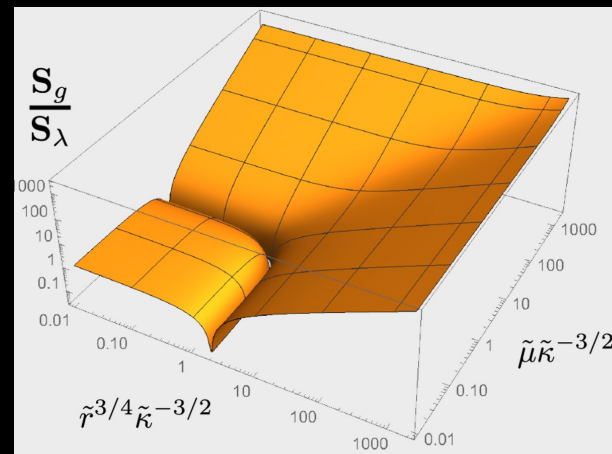
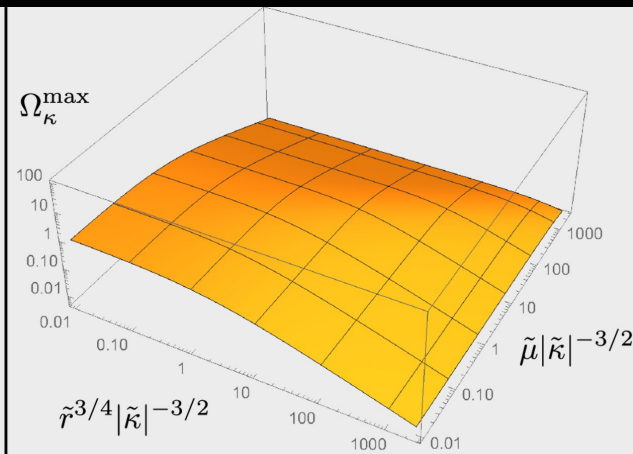
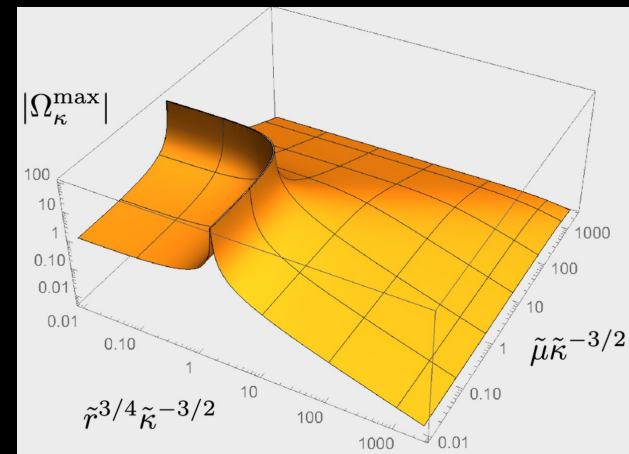
$$\kappa > 0$$

in this regime,
 S_g is *extensive*;
bigger universes
(higher S_r) are
also flatter

$$\max\left(\frac{|\Omega_\kappa|}{\Omega_{r+\Lambda}}\right)$$

$$S_r/S_\Lambda^{3/4}$$



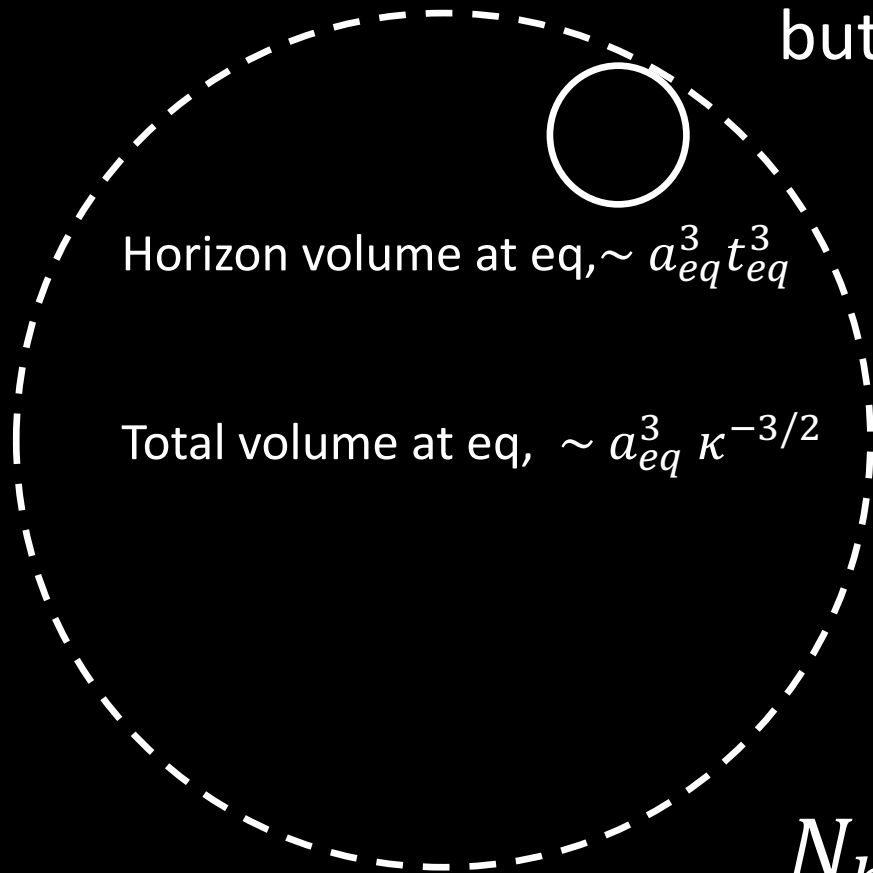
$\kappa > 0$ $\kappa < 0$ $\kappa > 0$ $\kappa < 0$ 

$$\tilde{r} \equiv \frac{r}{\lambda}; \quad \tilde{\mu} \equiv \frac{\mu}{\lambda}; \quad \tilde{\kappa} \equiv \frac{\kappa}{\lambda}; \quad S_{\lambda} = \frac{24\pi^2}{L_{Pl}^2 \lambda}$$

understanding in terms of horizons

$$3\dot{a}^2 = -3\kappa a^2 + r + \lambda a^4; \text{ equal } \Lambda, \text{ radiation density at } a_{eq} = (r/\lambda)^{1/4}$$

$$\text{but for } a < a_{eq}, a \sim r^{1/2} t \text{ so } t_{eq} \sim 1/(r\lambda)^{1/4}$$



Horizon volume at eq, $\sim a_{eq}^3 t_{eq}^3$

Total volume at eq, $\sim a_{eq}^3 \kappa^{-3/2}$

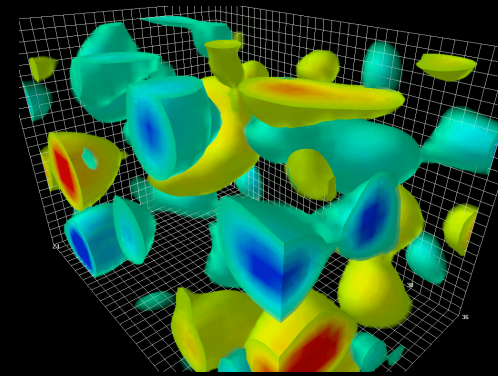
number of horizon volumes at equality

$$N_{hor} \sim (\lambda r / \kappa^2)^{3/4}$$

multiply by de Sitter entropy $\sim \lambda^{-1}$

$$N_{hor} \lambda^{-1} \sim (r / \kappa^2)^{3/4} \lambda^{-1/4} \sim S_r S_\Lambda^{1/4}$$

Quantum fields and gravity



vacuum energy and pressure are divergent,
simple physical regularizations such as point splitting give (for, *e.g.*, Maxwell):

$$\Rightarrow \langle T^{\mu\nu} \rangle_{vac} \sim \frac{3}{\pi^2 \Delta t^4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{3} & 0 & 0 \\ 0 & 0 & \frac{1}{3} & 0 \\ 0 & 0 & 0 & \frac{1}{3} \end{pmatrix}, \quad \text{where } \Delta t^2 = \text{invariant time-like separation}$$

B.S.DeWitt, Phys. Rep. 19 (1975) 295

Breaks Lorentz invariance! Can be renormalized away but leaves us with little physical understanding of the QFT vacuum.

Worse still are Weyl anomalies where quantum divergences spoil the local scale invariance of Maxwell and Dirac fields: violations cannot be renormalized away

Dimension zero scalars

A four-derivative, Weyl-invariant (*i.e.*, locally scale-invariant) action

$$S_4 = -\frac{1}{2} \int d^4x \sqrt{-g} \varphi \Delta_4 \varphi; \quad \Delta_4 = \square^2 + \dots$$

Heisenberg (1957), Pauli, Thirring, Nakanishi, ...
Flato, Fronsdal ('70s, '80s) forerunner of AdS/CFT

φ is Heisenberg's "dipole ghost" or Dirac's "singleton"; a very interesting theory

It has an infinite dimensional symmetry: $\varphi(x) \rightarrow \varphi(x) + \alpha(x)$ with $\square \alpha = 0$

The only physical state is the vacuum: there are no excited states

Bogoliubov et al (1987); Rivelles (2003)

The vacuum fluctuations are scale-invariant

$$\langle \varphi(0, x) \varphi(0, y) \rangle = \int \frac{d^3k}{(2\pi)^3} \frac{e^{ik \cdot (x-y)}}{4k^3}$$

cf. observed Newtonian potential
in cosmology

SM + dim-zero scalar numerology:

the vacuum energy and conformal anomalies (*at lowest order*)

$$E_k = \frac{1}{2} \hbar k (n_{s,1} - 2n_F + 2n_A + 2n_{s,0}) \quad \text{per mode } k$$

dim-one scalars chiral fermions gauge fields dim-zero scalars

$$\langle T^\mu_\mu \rangle = -a E + c C^2 \quad E = R_{\alpha\beta\gamma\delta} R^{\alpha\beta\gamma\delta} - 4R_{\alpha\beta} R^{\alpha\beta} + R^2; \quad C^2 = C^{\alpha\beta\gamma\delta} C^{\alpha\beta\gamma\delta}$$

$$a = \frac{1}{360(4\pi)^2} \left[n_{s,1} + \frac{11}{2} n_F + 62 n_A - 28 n_{s,0} \right]$$

$$c = \frac{1}{120(4\pi)^2} \left[n_{s,1} + 3 n_F + 12 n_A - 8 n_{s,0} \right]$$

Cancellation of all three implies $n_F = 4n_A$; $n_{s,0} = 3n_A$; $n_{s,1} = 0$.

Given $SU3 \times SU2 \times U1$, $n_A = 12$, **predicts** 3 generations of fermions, each with a ν_R

Also **requires** no fundamental dim-1 scalars so the Higgs must be composite
(exponentiating a dim 0 scalar gives an operator with nontrivial scaling dimension)

Promising development in self-dual gravity – twistor formulation of (exact) path integral for gravity also has anomalies but these can be removed *to all orders* using dimension zero scalars

K. Costello 2111.08879, R. Bittleston *in prep.* 2023

Graviton propagator with 1 loop SM corrections

$$\text{wavy line} + \text{wavy line with loop} + \text{wavy line with two loops} + \dots$$

Loop is given by the Fourier transform of the stress-energy correlator: for a CFT,

$$x \underset{\mu\nu}{*} \bigcirc \underset{\rho\lambda}{*} y = \langle T^{\mu\nu}(x) T^{\rho\lambda}(y) \rangle = C^T \frac{1}{4\pi^4 x^8} I^{\mu\nu,\rho\lambda}(x-y)$$

where $I^{\mu\nu,\rho\lambda}(x) = \frac{1}{2}(I^{\mu\rho}(x)I^{\nu\lambda}(x) + I^{\mu\lambda}(x)I^{\rho\nu}(x)) - \frac{1}{4}\eta^{\mu\nu}\eta^{\rho\lambda}$ and $I^{\mu\nu}(x) = \eta^{\mu\nu} - 2\frac{x^\mu x^\nu}{x^2}$

$C^T = \frac{4}{3}[n_{s,1} + 3n_F + 12n_A - 8n_{s,0}] \equiv \frac{4}{3}n_{eff}$ (\propto coefft of Weyl squared in the trace anomaly)

$$\text{Dim reg and min subt} \Rightarrow D^{\alpha\beta,\mu\nu}(k) = \frac{P^{\alpha\beta,\mu\nu}(k)}{k^2 \left(\left(1 - \frac{n_{eff}}{240\pi} G k^2 \ln\left(-\frac{k^2}{\mu^2}\right)\right) \right)}$$

Projector onto spin 2 component
- gauge invariant

SM corrections to the graviton propagator are problematic:

1. Inconsistent with Källén-Lehmann repn. $D(k) = \int_0^\infty dm^2 \rho(m^2) \frac{1}{k^2 - m^2 + i\varepsilon}$

(follows from Poincare invariance and positivity of the physical Hilbert space)

2. Specifically, resummed $D(k)$ (i) falls off as $|k|^{-4}$ at large $|k|$

(ii) has complex (acausal) poles on physical sheet

Similarly, Dim-0 scalar loops alone violate K-L: (i) $|k|^{-4}$ fall off; (ii) a tachyonic pole

BUT:

SM + Dim-0 combination is consistent with Poincaré, positivity and microcausality (at one loop in SM gauge+fermion fields: we are examining higher orders)

A Minimal Explanation of the Primordial Cosmological Perturbations

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²*Perimeter Institute for Theoretical Physics, Waterloo, Ontario, Canada, N2L 2Y5*

We outline a new explanation for the primordial density perturbations in cosmology. Dimension zero fields are a minimal addition to the Standard Model of particle physics: if the Higgs doublet is emergent, they cancel the vacuum energy and both Weyl anomalies without introducing any new particles. Furthermore, the cancellation explains why there are three generations of elementary particles, including RH neutrinos. We show how quantum zero point fluctuations of dimension zero fields seed nearly scale-invariant, Gaussian, adiabatic density perturbations. We determine their amplitude in terms of Standard Model couplings and find it is consistent with observation. Subject to two simple theoretical assumptions, both the amplitude and the tilt we compute *ab initio* agree with the measured values inferred from large scale structure observations, with no free parameters.


primordial perturbations from dim-0 fields and the SM

Running couplings violate scale symmetry: at high temperature,

$$T_{\beta}^{SM} \equiv \langle T_{\mu}^{SM\mu} \rangle_{\beta} = 3P - \rho \approx \sum c_i \alpha_i^2 T^4 \equiv c_{\beta}^{SM} T^4$$

This anomalous trace can be cancelled by introducing a linear coupling in the effective action,

$$\Gamma^{\varphi} = \sum_{j=1}^{n_{s,0}} \frac{1}{2} \int -a \varphi_j \Delta_4 \varphi_j + \left[a \left(E - \frac{2}{3} \square R \right) + c C^2 - n_{s,0}^{-1} T_{\beta}^{SM} \right] \varphi_j$$

 non-Weyl invariant term used to cancel anomalies
(generalizing sigma models in string theory)

The final linear term is chosen to cancel the trace anomaly due to running couplings at high T
It corrects the Einstein-fluid equations, converting quantum correlations in the dim-0 fields into large scale curvature fluctuations: Friedmann equation becomes

$$\dot{a}^2 = \frac{8\pi G}{3} \rho_r a^4 (1 + c_{\varphi} \bar{\varphi}(x)) \text{ with } \bar{\varphi}(x) = n_{s,0}^{-1} \sum \varphi_j(x), \quad c_{\varphi} = c_{\beta}^{SM} / \left(\frac{\pi^2}{30} \mathcal{N}_{eff} \right), \quad \mathcal{N}_{eff} \approx 106\frac{1}{4}$$

Conformal factor translates directly into “comoving curvature perturbation” $\mathcal{R}(x) = \frac{1}{4} c_{\varphi} \bar{\varphi}(x)$
(adiabatic, Gaussian, scalar: no primordial long-wavelength gravitational waves)

Spectral tilt

Dominated by QCD: asymptotic freedom \Rightarrow red tilt!

To understand quantitatively, consider the trace anomaly (for QCD)

$$S = -\int \frac{1}{4} F^2 \Rightarrow -\int \frac{1}{4g^2} F^2. \quad \alpha \equiv \frac{g^2}{4\pi}; \quad \mu \partial_\mu \alpha \equiv \beta_\alpha; \Rightarrow \mu \partial_\mu S = \int \frac{\beta_\alpha}{\alpha} \frac{1}{4g^2} F^2 \Rightarrow \int \frac{\beta_\alpha}{4\alpha} F^2$$

$$\Rightarrow T_\lambda^\lambda = \frac{\beta_\alpha}{4\alpha} F^2; \quad \frac{\beta_\alpha}{\alpha} = -(11 - \frac{2}{3} n_f) \frac{\alpha}{2\pi} T^4; \quad \langle F^2 \rangle_\beta = \frac{2\pi\alpha}{9} (12 + 5n_f) T^4$$

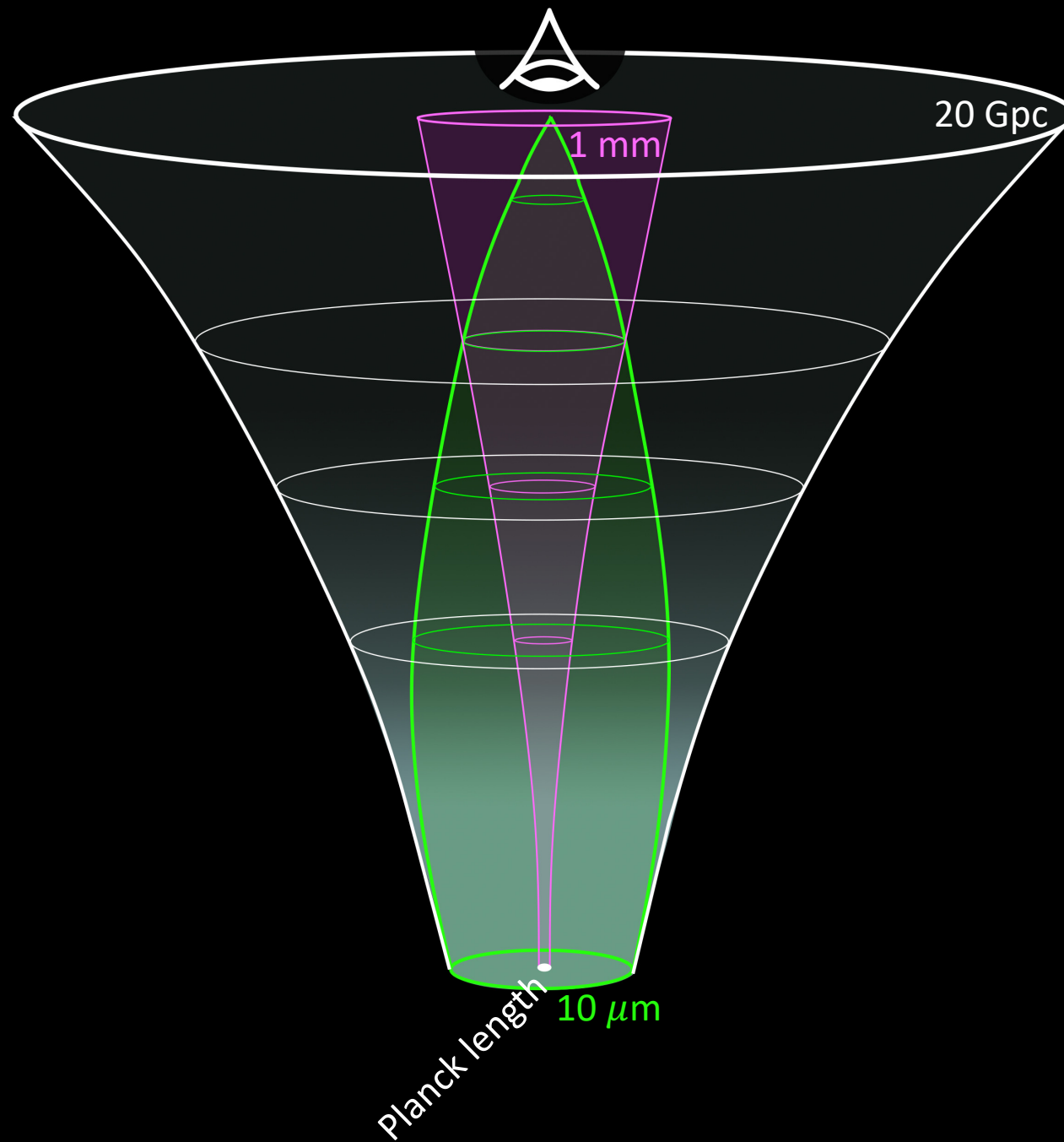
running coupling: energy scale of φ

plasma interactions: energy scale of T

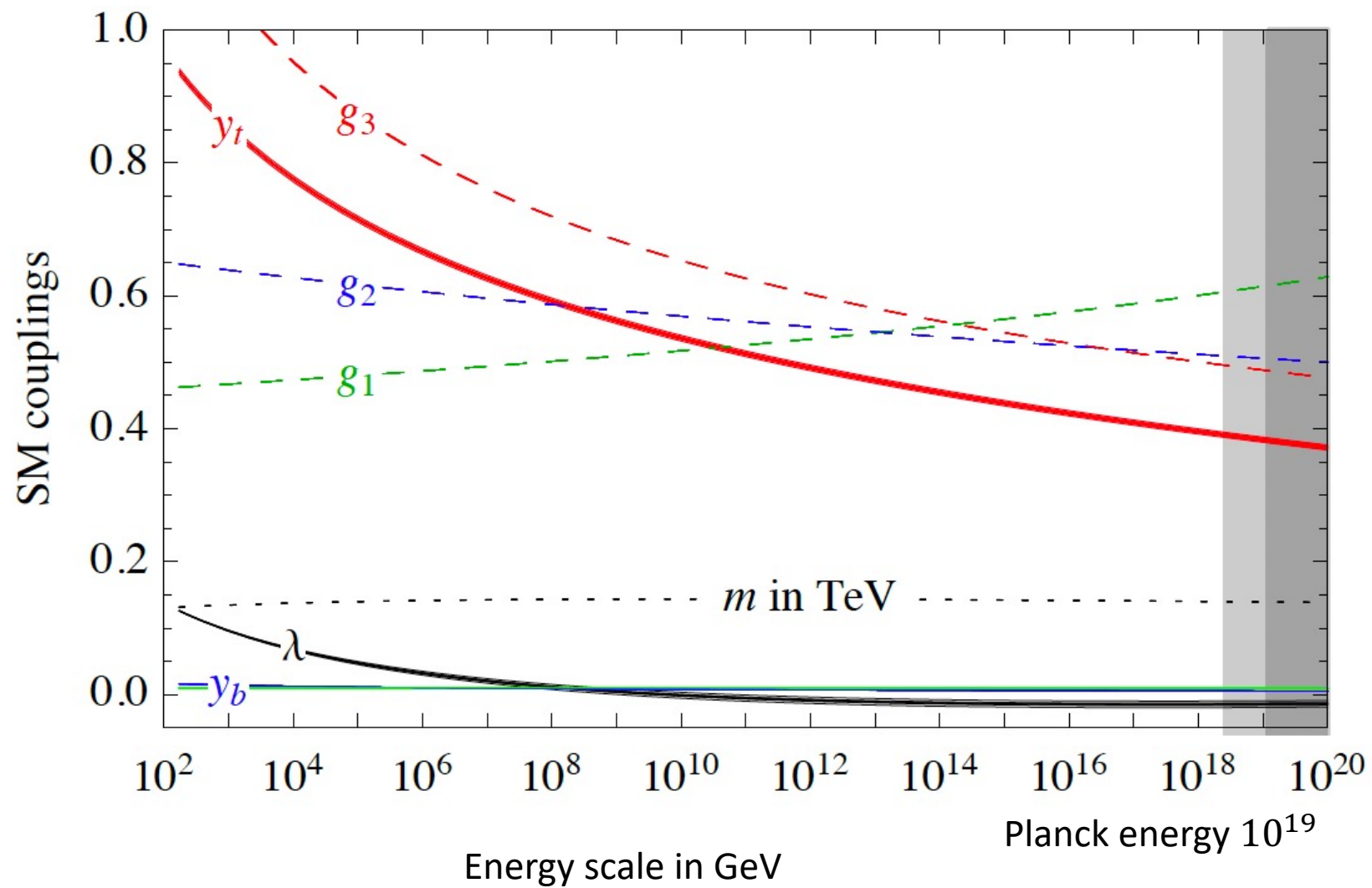
Thus, $\mathcal{P}_\mathcal{R}(k)$ scales with k as $\alpha^2(k)$; $n_s - 1 \equiv \frac{d \ln \mathcal{P}_\mathcal{R}(k)}{d \ln k} = 2 \frac{\beta_\alpha}{\alpha} = -\frac{7}{\pi} \alpha_{QCD}(M_P)$

The red tilt is a **critical exponent** which can be computed perturbatively

If so, we can extrapolate over 30 orders of magnitude in length scale...



Buttazzo et al
1307.3536
[hep-ph]



Comparison with observation

$$\mathcal{P}_{\mathcal{R}}(k) = \frac{3^2 5^2}{7(2\pi)^4} \left(\frac{c_{\beta}^{SM}}{\mathcal{N}_{eff}} \right)^2 \left(\frac{k}{k_P} \right)^{-\frac{7\alpha_3}{\pi}}; \quad k_P = \text{comoving Planck wavenumber}$$

with $c_{\beta}^{SM} \equiv \frac{125}{108} \alpha_Y^2 - \frac{95}{72} \alpha_2^2 - \frac{49}{6} \alpha_3^2$ and $\mathcal{N}_{eff} = 106\frac{1}{4}$ (to lowest order, neglect Higgs)

Now use $(k_P/k_*)^{1-n_s} = 14.8 \pm 5.1$, $k_* \equiv 0.05 \text{ Mpc}^{-1}$

Thus, we find $\mathcal{P}_{\mathcal{R}} = A \left(\frac{k}{k_*} \right)^{n_s-1}$, $A = (13 \pm 5) \times 10^{-10}$; $n_s = 0.958$

cf. Planck satellite: $A = (21 \pm 0.3) \times 10^{-10}$; $n_s = 0.959 \pm 0.006$

Comparison with observation

$$\mathcal{P}_{\mathcal{R}}(k) = \frac{3^2 5^2}{7(2\pi)^4} \left(\frac{c_{\beta}^{SM}}{\mathcal{N}_{eff}} \right)^2 \left(\frac{k}{k_P} \right)^{-\frac{7\alpha_3}{\pi}}; \quad k_P = \text{comoving Planck wavenumber}$$

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Now use $(k_*/k_P)^{n_s-1} = 14.8 \pm 5.1$, $k_* \equiv 0.05 \text{ Mpc}^{-1}$

Thus, we predict $\mathcal{P}_{\mathcal{R}} = A \left(\frac{k}{k_*} \right)^{n_s-1}$, $A = (13 \pm 5) \times 10^{-10}$; $n_s = 0.958$

cf. Planck satellite: $A = (21 \pm 0.3) \times 10^{-10}$; $n_s = 0.959 \pm 0.006$

prediction will be tested further as observations and theory improve

summary

analytic extension of cosmological solutions of the Einstein equations lead to

- a new picture of the big bang singularity as a CPT “mirror”
- a calculation of the gravitational entropy for cosmologies

new explanations and predictions for

- the large-scale homogeneity, isotropy and flatness of the cosmos (and a hint about Lambda)
- the dark matter
- the arrow of time and the strong CP problem

Dimension zero scalars

- cancel the vacuum energy and both Weyl anomalies at leading (free field) order
- explain why there are 3 generations of SM fermions, including RH neutrinos
- explain the amplitude, tilt and character of the primordial perturbations
- require the Higgs to be emergent/composite, a new approach to the gauge-gravity hierarchy?

All without adding any new propagating degrees of freedom to the SM and Einstein gravity

These are encouraging signs but much remains to be understood

Thank You!

Boyle, Finn, NT Phys. Rev. Lett. 121 (2018) 251301; Annals of Physics 438 (2022) 168767
arXiv: 2109.06204, 2110.06258, 2201.07279, 2208.10396, 2210.01142, [2302.00344](#)