

Neutrino Masses, Leptogenesis and Dark Matter from a Scotogenic Model

Edward Wang

Technical University of Munich (TUM)

2022



SFB 1258

Neutrinos
Dark Matter
Messengers



Introduction

Three important shortcomings of the Standard Model:

- 1 Neutrino masses

Introduction

Three important shortcomings of the Standard Model:

- ① Neutrino masses
- ② Baryon Asymmetry of the Universe (BAU)

Introduction

Three important shortcomings of the Standard Model:

- ① Neutrino masses
- ② Baryon Asymmetry of the Universe (BAU)
- ③ Dark Matter

Introduction

Three important shortcomings of the Standard Model:

- ① Neutrino masses
- ② Baryon Asymmetry of the Universe (BAU)
- ③ Dark Matter

Address three issues with single extension to SM

The Model

- To produce two neutrino mass states, usually add two Majorana fermions and one Higgs doublet

The Model

- To produce two neutrino mass states, usually add two Majorana fermions and one Higgs doublet
- Ibarra and Hehn (2012) proposed a Scotogenic model with one Majorana fermion N and two additional complex scalar doublets $\eta_{1,2}$

The Model

- To produce two neutrino mass states, usually add two Majorana fermions and one Higgs doublet
- Ibarra and Hehn (2012) proposed a Scotogenic model with one Majorana fermion N and two additional complex scalar doublets $\eta_{1,2}$

$$\mathcal{L}_{\text{int}}^{\text{fermion}} = -Y_i^{(a)} \bar{N} (\nu_i \eta_a^0 - l_i^- \eta_a^+) + \text{h.c.}$$

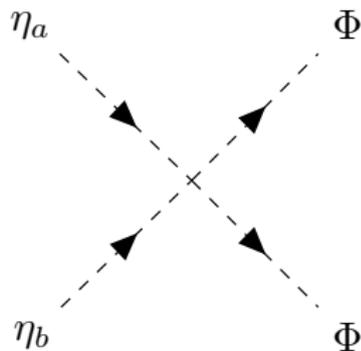
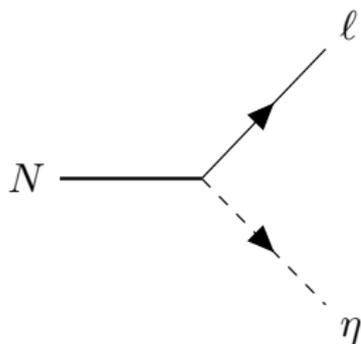
$$\mathcal{L}_{\text{int}}^{\text{scalar}} = -\frac{1}{2} \lambda_3^{(ab)} (\Phi^\dagger \Phi) (\eta_a^\dagger \eta_b) - \frac{1}{2} \lambda_4^{(ab)} (\Phi^\dagger \eta_a) (\eta_b^\dagger \Phi) - \frac{1}{2} \lambda_5^{(ab)} (\Phi^\dagger \eta_a) (\Phi^\dagger \eta_b)$$

The Model

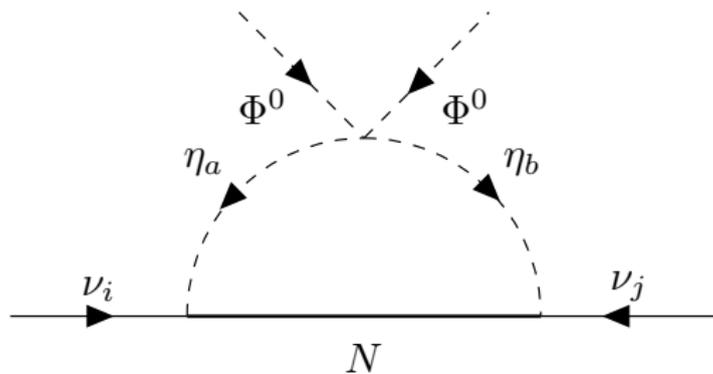
- To produce two neutrino mass states, usually add two Majorana fermions and one Higgs doublet
- Ibarra and Hehn (2012) proposed a Scotogenic model with one Majorana fermion N and two additional complex scalar doublets $\eta_{1,2}$

$$\mathcal{L}_{\text{int}}^{\text{fermion}} = -Y_i^{(a)} \bar{N} (\nu_i \eta_a^0 - l_i^- \eta_a^+) + \text{h.c.}$$

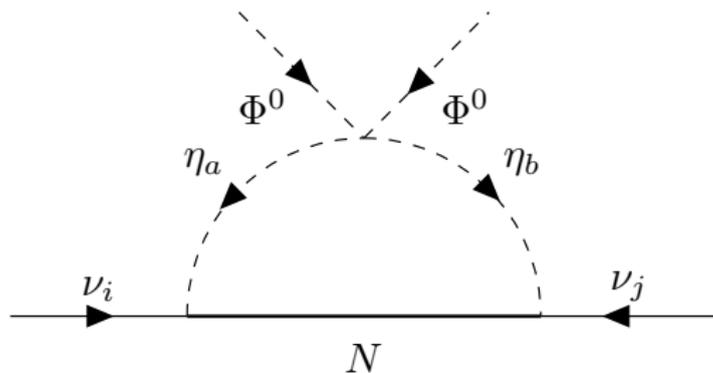
$$\mathcal{L}_{\text{int}}^{\text{scalar}} = -\frac{1}{2} \lambda_3^{(ab)} (\Phi^\dagger \Phi) (\eta_a^\dagger \eta_b) - \frac{1}{2} \lambda_4^{(ab)} (\Phi^\dagger \eta_a) (\eta_b^\dagger \Phi) - \frac{1}{2} \lambda_5^{(ab)} (\Phi^\dagger \eta_a) (\Phi^\dagger \eta_b)$$



Neutrino Masses

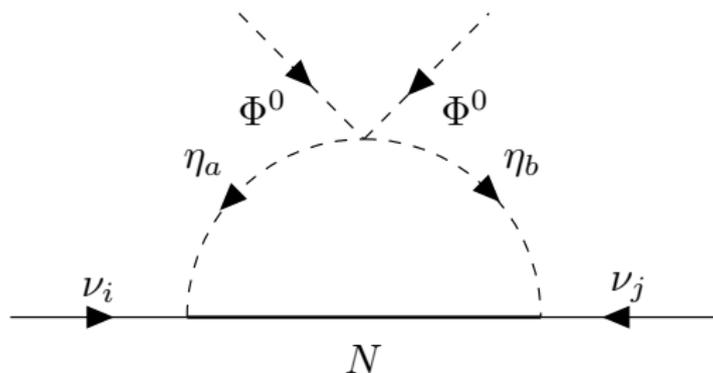


Neutrino Masses



$$(m_\nu)_{ij} = \frac{Y_i^{(a)} Y_j^{(b)} \lambda_5^{(ab)} v^2}{8\pi^2} \frac{M_N}{m_{\eta_b}^2 - M_N^2} \left(\frac{m_{\eta_b}^2}{m_{\eta_a}^2 - m_{\eta_b}^2} \log \left(\frac{m_{\eta_a}^2}{m_{\eta_b}^2} \right) - \frac{M_N^2}{m_{\eta_a}^2 - M_N^2} \log \left(\frac{m_{\eta_a}^2}{M_N^2} \right) \right)$$

Neutrino Masses



$$(m_\nu)_{ij} = \frac{Y_i^{(a)} Y_j^{(b)} \lambda_5^{(ab)} v^2}{8\pi^2} \frac{M_N}{m_{\eta_b}^2 - M_N^2} \left(\frac{m_{\eta_b}^2}{m_{\eta_a}^2 - m_{\eta_b}^2} \log \left(\frac{m_{\eta_a}^2}{m_{\eta_b}^2} \right) - \frac{M_N^2}{m_{\eta_a}^2 - M_N^2} \log \left(\frac{m_{\eta_a}^2}{M_N^2} \right) \right)$$

Note

m_ν has rank 2 \Rightarrow produce two massive neutrino states

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

- Departure from thermal equilibrium

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

- Departure from thermal equilibrium
- Baryon-number violation

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

- Departure from thermal equilibrium
- Baryon-number violation
- CP-violation

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

- Departure from thermal equilibrium \Rightarrow Expansion of the Universe
- Baryon-number violation
- CP-violation

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

- Departure from thermal equilibrium \Rightarrow Expansion of the Universe
- Baryon-number violation \Rightarrow Sphaleron processes
- CP-violation

Baryogenesis

Any Baryogenesis mechanism must contain three ingredients (Sakharov conditions)

- Departure from thermal equilibrium \Rightarrow Expansion of the Universe
- Baryon-number violation \Rightarrow Sphaleron processes
- CP-violation \Rightarrow Majorana fermion decay

CP-violation

Majorana fermion decays into leptons and anti-leptons \Rightarrow CP-violation possible.

CP-violation

Majorana fermion decays into leptons and anti-leptons \Rightarrow CP-violation possible. First attempt: tree-level.

CP-violation

Majorana fermion decays into leptons and anti-leptons \Rightarrow CP-violation possible. First attempt: tree-level. But

$$\left(\begin{array}{c} \text{Diagram 1} \end{array} \right)^* = \text{Diagram 2}$$

The diagrammatic equation shows the complex conjugate of a Majorana fermion decay into a lepton and a neutrino, equal to a Majorana fermion decay into an anti-lepton and an anti-neutrino. On the left, a Majorana fermion N (solid line) decays into a lepton ℓ (solid line) and a neutrino η (dashed line). This diagram is enclosed in large parentheses with an asterisk $*$ to its upper right. An equals sign $=$ follows. On the right, a Majorana fermion N (solid line) decays into an anti-lepton $\bar{\ell}$ (solid line) and an anti-neutrino η^* (dashed line).

CP-violation

Majorana fermion decays into leptons and anti-leptons \Rightarrow CP-violation possible. First attempt: tree-level. But

$$\left(\begin{array}{c} \text{Diagram 1: } N \text{ decaying to } \ell \text{ and } \eta \\ \text{Diagram 2: } N \text{ decaying to } \bar{\ell} \text{ and } \eta^* \end{array} \right)^* = \text{Diagram 2}$$

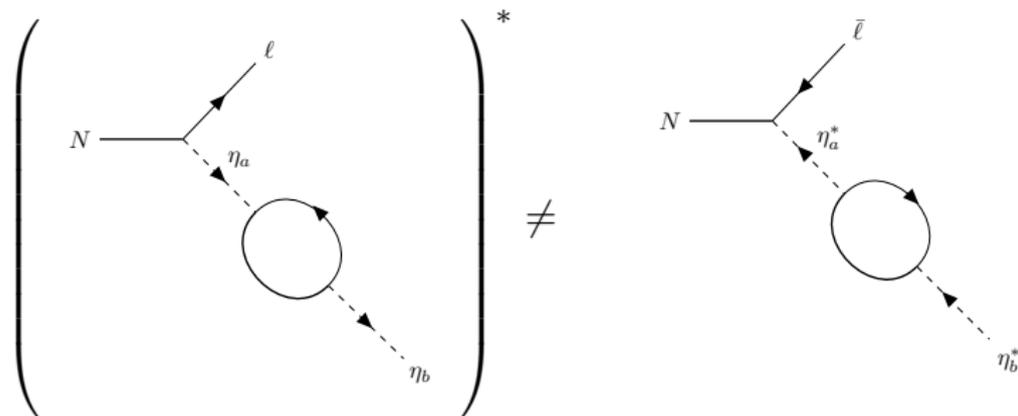
$\Rightarrow |i\mathcal{M}|^2 = |i\mathcal{M}^{CP}|^2$

CP-violation

Next attempt: loop-level.

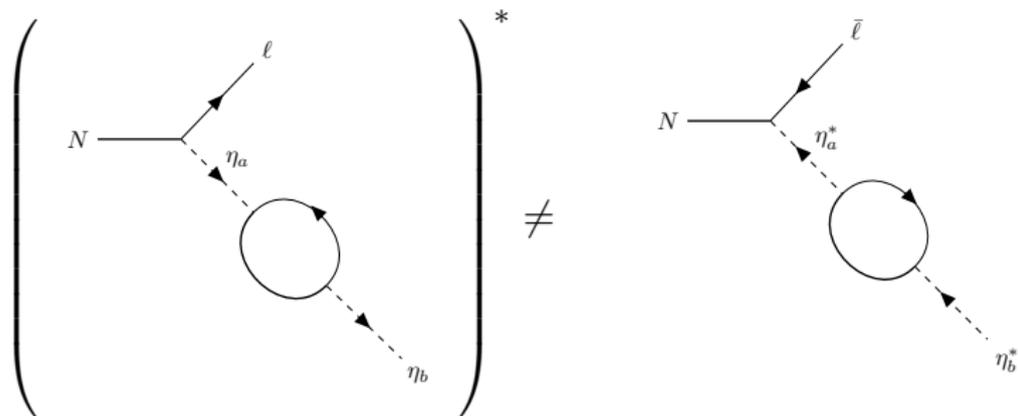
CP-violation

Next attempt: loop-level. Now have



CP-violation

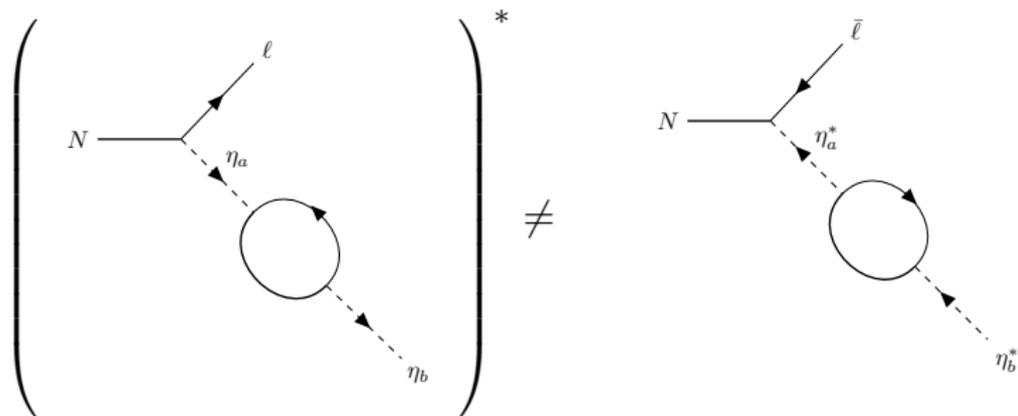
Next attempt: loop-level. Now have



$$\Rightarrow |i\mathcal{M}_{tree} + i\mathcal{M}_{1-loop}|^2 \neq |i\mathcal{M}_{tree}^{CP} + i\mathcal{M}_{1-loop}^{CP}|^2$$

CP-violation

Next attempt: loop-level. Now have



$$\Rightarrow |i\mathcal{M}_{tree} + i\mathcal{M}_{1-loop}|^2 \neq |i\mathcal{M}_{tree}^{CP} + i\mathcal{M}_{1-loop}^{CP}|^2$$

Note

$$\text{Decay asymmetry} \sim \frac{\text{Im}[Y_i^{(1)} Y_j^{(1)*} Y_j^{(2)} Y_i^{(2)*}]}{m_{\eta 1}^2 - m_{\eta 2}^2}$$

Kinetic equations

Assume $M_N \ll m_{\eta 1}, m_{\eta 2}$.

Kinetic equations

Assume $M_N \ll m_{\eta 1}, m_{\eta 2}$. Obtain kinetic equations of the form

$$\begin{aligned}\frac{dY_N}{dz} &= C_N(Y_N - Y_N^{\text{eq}}) \\ \frac{dY_{\ell i}}{dz} &= S_{\ell i}(Y_N - Y_N^{\text{eq}}) - W_{\ell i}Y_{\ell i}\end{aligned}$$

with $z = m_{\eta 1}/T$.

Kinetic equations

Assume $M_N \ll m_{\eta_1}, m_{\eta_2}$. Obtain kinetic equations of the form

$$\begin{aligned}\frac{dY_N}{dz} &= C_N(Y_N - Y_N^{\text{eq}}) \\ \frac{dY_{\ell i}}{dz} &= S_{\ell i}(Y_N - Y_N^{\text{eq}}) - W_{\ell i}Y_{\ell i}\end{aligned}$$

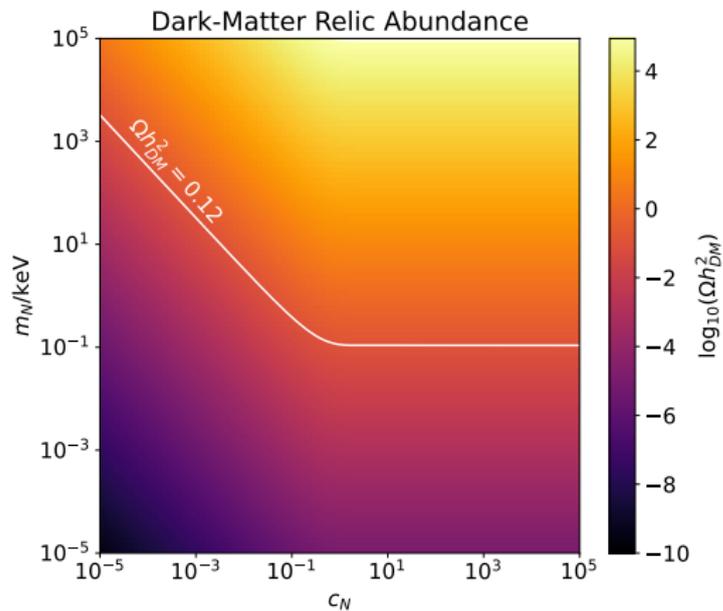
with $z = m_{\eta_1}/T$. We find

$$C_N \approx -c_N z^{5/2} e^{-z}$$

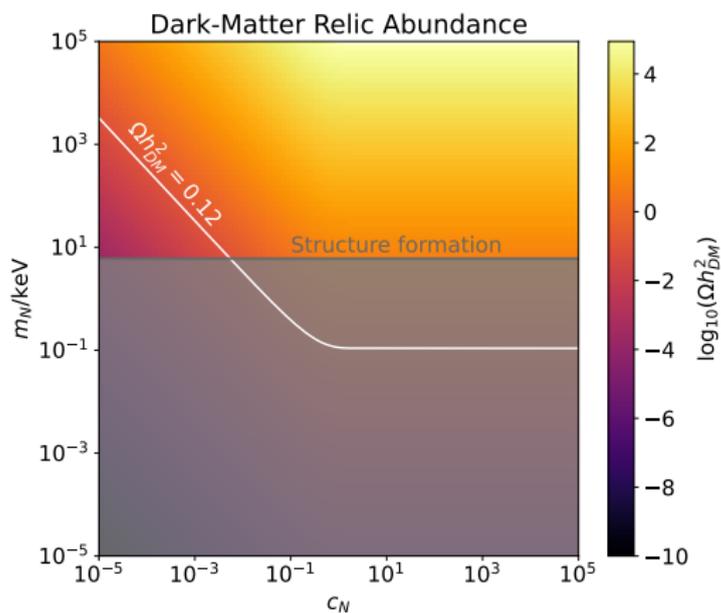
$$W_{\ell i} \approx c_{W i} z^{5/2} e^{-z}$$

$$S_{\ell i} \approx c_{S i} z^{5/2} e^{-z}$$

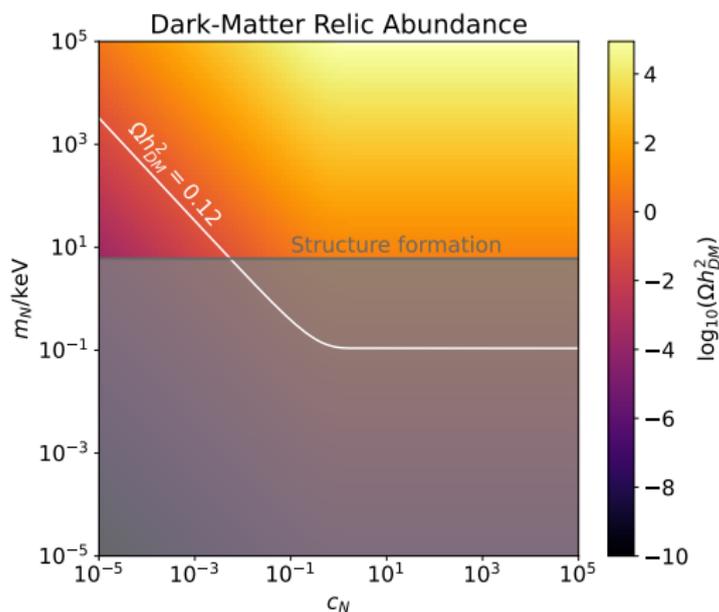
Dark Matter



Dark Matter



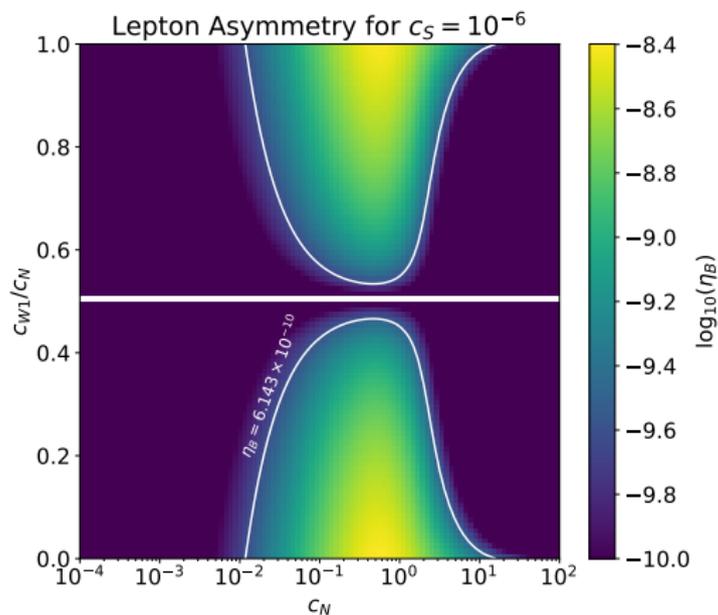
Dark Matter



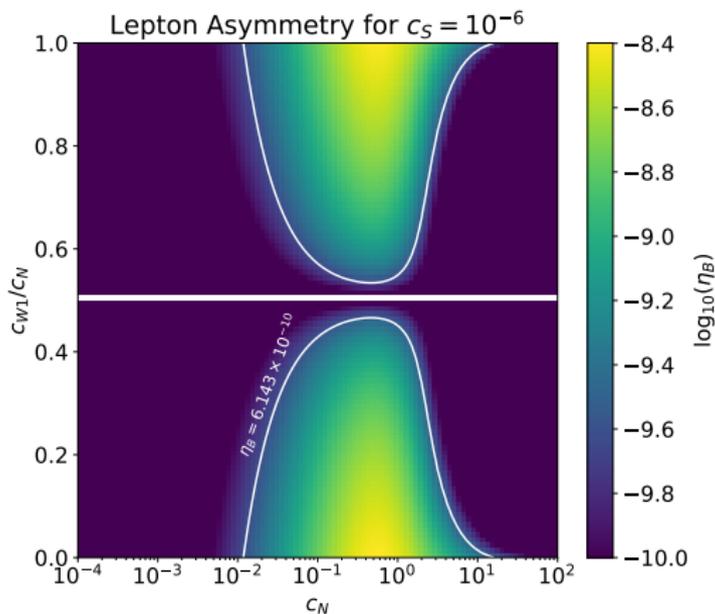
1. Constraint

$$m_{DM} \gtrsim 6 \text{ keV} \Rightarrow c_N \lesssim 5.5 \times 10^{-3}$$

Leptogenesis



Leptogenesis



2. Constraint

$$c_{S1} = -c_{S2}, c_{W1} + c_{W2} = c_N \Rightarrow c_S > 2.106 \times 10^{-6}$$

Neutrino Masses

In neutrino mass eigenbasis,

$$m_\nu = Y\Lambda Y^T = \begin{pmatrix} m_2 & 0 \\ 0 & m_3 \end{pmatrix}$$

where

$$\Lambda^{(ab)} = \frac{\lambda_5^{(ab)} v^2}{8\pi^2} \frac{M_N}{m_{\eta_a}^2 - m_{\eta_b}^2} \log\left(\frac{m_{\eta_a}^2}{m_{\eta_b}^2}\right) \approx \frac{\lambda_5^{(ab)} v^2}{8\pi^2} \frac{M_N}{m_{\eta_a}^2}$$

in the limit $m_{\eta_b} \rightarrow m_{\eta_a}$.

Neutrino Masses

In neutrino mass eigenbasis,

$$m_\nu = Y\Lambda Y^T = \begin{pmatrix} m_2 & 0 \\ 0 & m_3 \end{pmatrix}$$

where

$$\Lambda^{(ab)} = \frac{\lambda_5^{(ab)} v^2}{8\pi^2} \frac{M_N}{m_{\eta_a}^2 - m_{\eta_b}^2} \log\left(\frac{m_{\eta_a}^2}{m_{\eta_b}^2}\right) \approx \frac{\lambda_5^{(ab)} v^2}{8\pi^2} \frac{M_N}{m_{\eta_a}^2}$$

in the limit $m_{\eta_b} \rightarrow m_{\eta_a}$. We can invert this to find

$$\lambda_5^{(ab)} = \left(Y^{-1} m_\nu (Y^T)^{-1} \right)^{(ab)} / \tilde{\Lambda}^{(ab)}$$

with $\tilde{\Lambda}^{(ab)} = \Lambda^{(ab)} / \lambda_5^{(ab)}$

Joint Parameter Region

For the decay parameter,

$$c_N = \sum_{i,a} \left| Y_i^{(a)} \right|^2 2^{-7/2} \pi^{-1/2} \frac{aR}{m_{\eta a}}$$

Joint Parameter Region

For the decay parameter,

$$c_N = \sum_{i,a} \left| Y_i^{(a)} \right|^2 2^{-7/2} \pi^{-1/2} \frac{aR}{m_{\eta a}}$$

\Rightarrow make $Y_i^{(a)}$ tiny and $m_{\eta a}$ large to obtain $c_N < 5.5 \times 10^{-3}$

Joint Parameter Region

For the decay parameter,

$$c_N = \sum_{i,a} \left| Y_i^{(a)} \right|^2 2^{-7/2} \pi^{-1/2} \frac{a_R}{m_{\eta a}}$$

\Rightarrow make $Y_i^{(a)}$ tiny and $m_{\eta a}$ large to obtain $c_N < 5.5 \times 10^{-3}$

For decay asymmetry,

$$c_S = \frac{\text{Im}[Y_i^{(1)} Y_j^{(1)*} Y_j^{(2)} Y_i^{(2)*}]}{m_{\eta 1}^2 - m_{\eta 2}^2} a_R m_{\eta 1} 2^{-5/2} \pi^{1/2}$$

Joint Parameter Region

For the decay parameter,

$$c_N = \sum_{i,a} \left| Y_i^{(a)} \right|^2 2^{-7/2} \pi^{-1/2} \frac{a_R}{m_{\eta a}}$$

\Rightarrow make $Y_i^{(a)}$ tiny and $m_{\eta a}$ large to obtain $c_N < 5.5 \times 10^{-3}$

For decay asymmetry,

$$c_S = \frac{\text{Im}[Y_i^{(1)} Y_j^{(1)*} Y_j^{(2)} Y_i^{(2)*}]}{m_{\eta 1}^2 - m_{\eta 2}^2} a_R m_{\eta 1} 2^{-5/2} \pi^{1/2}$$

\Rightarrow resonant enhancement! Expect $m_{\eta 1} \sim m_{\eta 2}$ to obtain $c_S > 2 \times 10^{-6}$.

Joint Parameter Region

For the decay parameter,

$$c_N = \sum_{i,a} \left| Y_i^{(a)} \right|^2 2^{-7/2} \pi^{-1/2} \frac{a_R}{m_{\eta a}}$$

\Rightarrow make $Y_i^{(a)}$ tiny and $m_{\eta a}$ large to obtain $c_N < 5.5 \times 10^{-3}$

For decay asymmetry,

$$c_S = \frac{\text{Im}[Y_i^{(1)} Y_j^{(1)*} Y_j^{(2)} Y_i^{(2)*}]}{m_{\eta 1}^2 - m_{\eta 2}^2} a_R m_{\eta 1} 2^{-5/2} \pi^{1/2}$$

\Rightarrow resonant enhancement! Expect $m_{\eta 1} \sim m_{\eta 2}$ to obtain $c_S > 2 \times 10^{-6}$.

Set

$$\lambda_5^{(ab)} = \left(Y^{-1} m_\nu \left(Y^T \right)^{-1} \right)^{(ab)} / \tilde{\Lambda}^{(ab)}$$

to generate neutrino masses

Summary & Outlook

- Explored new mechanism for Leptogenesis ✓

Summary & Outlook

- Explored new mechanism for Leptogenesis ✓
- Explain Neutrino Masses, BAU and DM with two Higgs doublets and single Majorana fermion ✓

Summary & Outlook

- Explored new mechanism for Leptogenesis ✓
- Explain Neutrino Masses, BAU and DM with two Higgs doublets and single Majorana fermion ✓
- Need to understand the joint parameter region better

Summary & Outlook

- Explored new mechanism for Leptogenesis ✓
- Explain Neutrino Masses, BAU and DM with two Higgs doublets and single Majorana fermion ✓
- Need to understand the joint parameter region better
- Study other interesting region $M_N \gg m_{\eta 1}, m_{\eta 2}$

Summary & Outlook

- Explored new mechanism for Leptogenesis ✓
- Explain Neutrino Masses, BAU and DM with two Higgs doublets and single Majorana fermion ✓
- Need to understand the joint parameter region better
- Study other interesting region $M_N \gg m_{\eta 1}, m_{\eta 2}$
- Explore phenomenology of the model

Backup Slides

