

Binary evolution & gravitational-wave mergers in second-generation gas-enriched globular-clusters

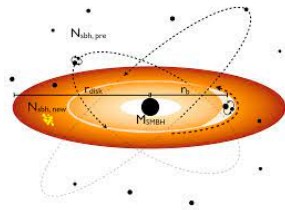
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There are many gaseous media



(a) AGN disks (Ford et al 2019)



(b) protoplanetary disks



(c) Star forming epochs in GCs

Dynamics in gaseous media is rich

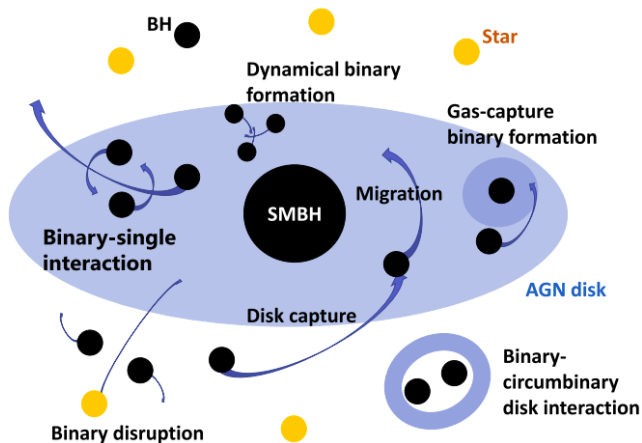


Figure: Borrowed from Tagawa et al 2020

Globular clusters host multiple generations

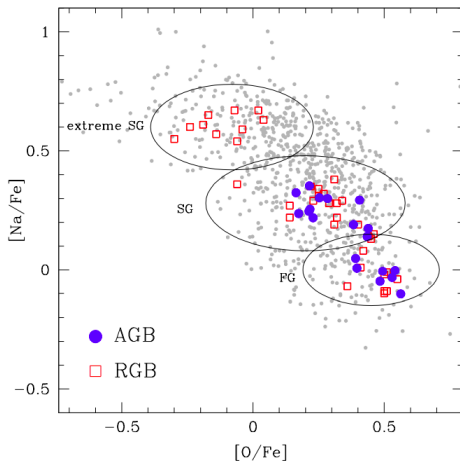
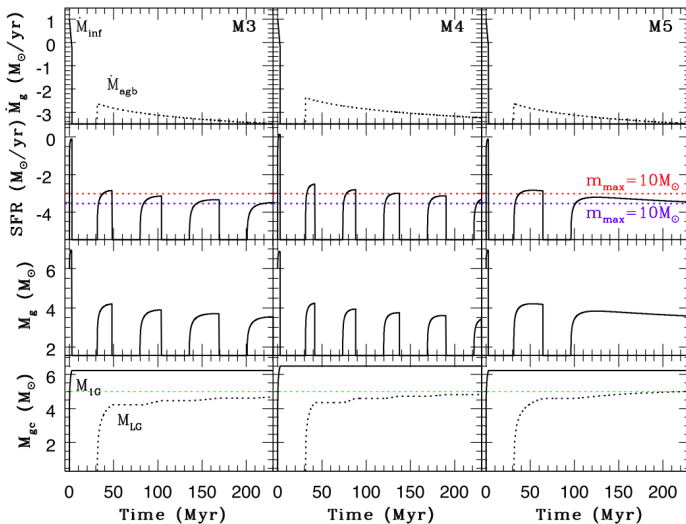


Figure: Behavior of $[Na/Fe]$ as a function of $[O/Fe]$ for the AGB (filled blue circles, this work) and RGB stars of NGC 6752. From Lapenna et al. 2016.

An example of gas from AGB stars – Bekki 2017



Second generation stars could form in disks

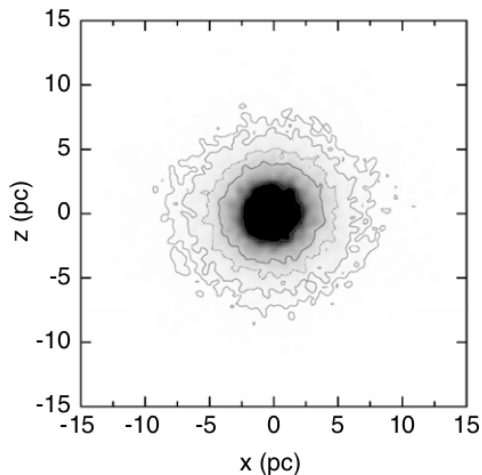


Figure: Projected isodensity contour, borrowed from Mastrobuono-Battisti & Perets 2013

Gaseous media affect on the dynamics

- AGN disks (e.g. McKernan et al 2012, Stone et al 2017, Tagawa et al 2020)
- Protoplanetary disks (e.g. Murray Clay & Perets 2011, Grishin & Perets 2016)
- Mergers between galaxies (e.g. Lin 2008)

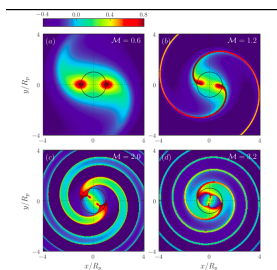
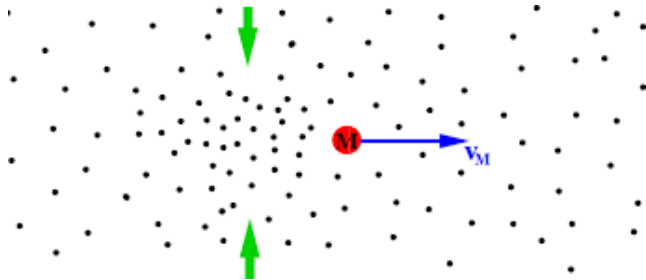


Figure: From Kim et al 2008

- Gas dynamical friction
- Three-body hardening
- Gravitational waves emission

$$\frac{da_{\text{bin}}}{dt} = \frac{da_{\text{bin}}}{dt} \Big|_{\text{hard}, \star} + \frac{da_{\text{bin}}}{dt} \Big|_{\text{GDF}} + \frac{da_{\text{bin}}}{dt} \Big|_{\text{GW}}$$

For large separations – gas dynamical friction



You can think about a duck



and then about a binary-duck!



For large separations – gas dynamical friction

$$\left. \frac{da_{\text{bin}}}{dt} \right|_{\text{GDF}} = - \frac{8\pi G^{3/2} a_{\text{bin}}^{3/2}}{\sqrt{m_1 + m_2}} \rho_g(t) \frac{m}{v_{\text{rel}}^2} f\left(\frac{v_{\text{rel}}}{c_s}\right)$$

a_{bin} – binary separation

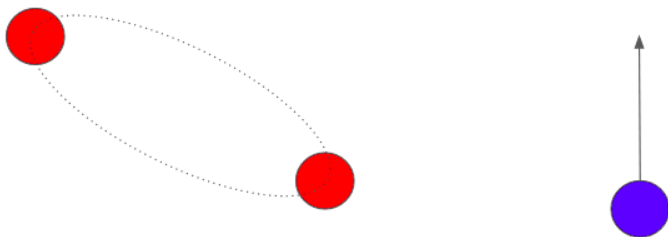
$m = m_1 = m_2$ – masses of the binary companions

$\rho_g = \rho_{g,0} e^{-t/\tau_g}$, ρ_g – gas density, τ_g – gas density

v_{rel} – relative velocity

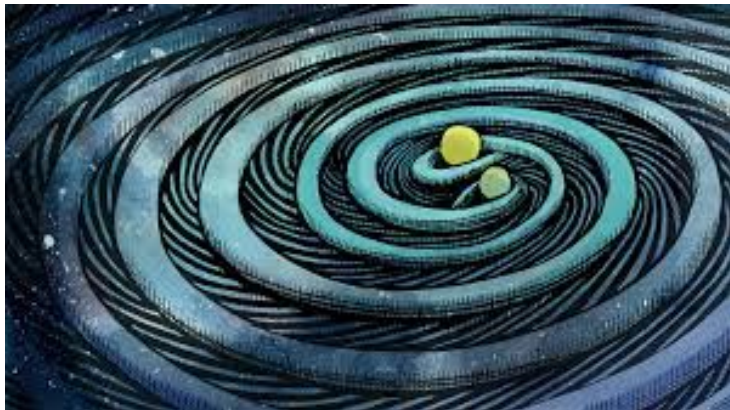
Intermediate separations – three-body hardening

$$\left. \frac{da_{\text{bin}}}{dt} \right|_{3\text{-body}} = - \frac{2\pi G n_{\star} m_{\text{pert}} (2m + m_{\text{pert}}) a_{\text{bin}}^2}{m v_{\infty}}$$



Finally – gravitational waves!

$$\left. \frac{da}{dt} \right|_{\text{GW}} = - \frac{64G^3 m_1 m_2 (m_1 + m_2)}{5c^5 a^3}$$



All together now

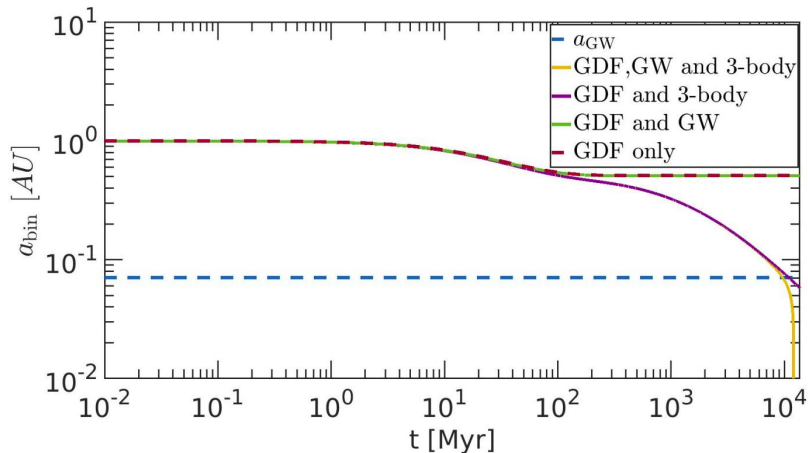


Figure: $m_1 = m_2 = 10 M_{\odot}$ (Rozner & Perets 2022)

Gas hardening is very efficient

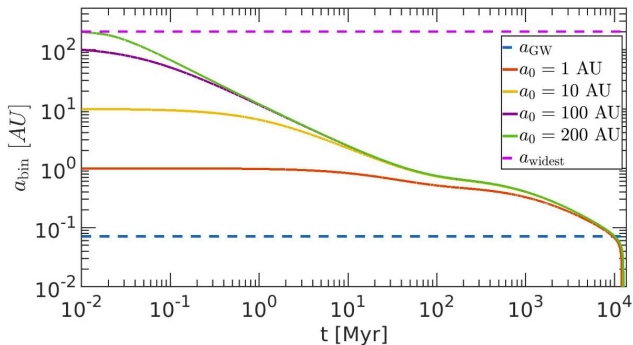
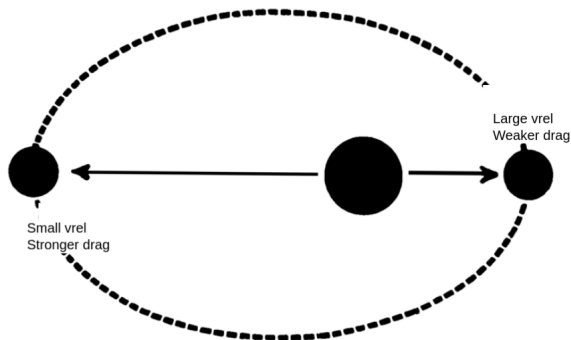


Figure: $m_1 = m_2 = 10 M_{\odot}$ (Rozner & Perets 2022)

What happens in eccentric evolution?



$$\left. \frac{da_{\text{bin}}}{dt} \right|_{\text{GDF}} = - \frac{8\pi G^{3/2} a_{\text{bin}}^{3/2}}{\sqrt{m_1 + m_2}} \rho_g(t) \frac{m}{v_{\text{rel}}^2} f\left(\frac{v_{\text{rel}}}{c_s}\right)$$

We consider orbit-averaged equations

For $v_g = 0$ relative to the center of mass of the binary

$$\mathbf{F}_{\text{drag}} = F_0 \mathbf{v}_{\text{rel}} / v_{\text{rel}}^3,$$

$$\left. \frac{da}{dt} \right|_{\text{GDF}} = \frac{4F_0(1-e^2)^2}{\pi m_{\text{bin}} \Omega^3 a^2} \int_0^{2\pi} \frac{df}{(1+e \cos f)^2 \sqrt{1+2e \cos f + e^2}},$$

$$\left. \frac{de}{dt} \right|_{\text{GDF}} = \frac{F_0(1-e^2)^3}{\pi m_{\text{bin}} \Omega^3 a^3} \int_0^{2\pi} \frac{(e + \cos f) df}{(1+e \cos f)^2 (1+2e \cos f + e^2)^{3/2}}$$

The merger timescales shorten

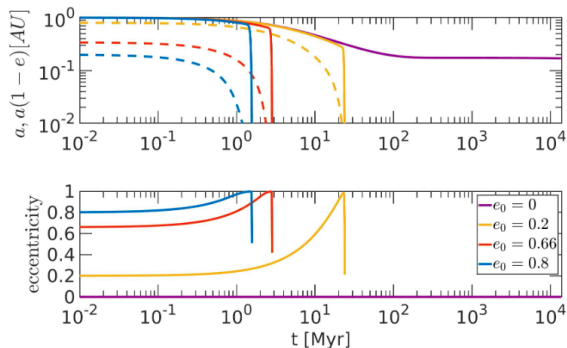


Figure: $m_1 = m_2 = 10 M_\odot$, $a_0 = 1$ AU orbit-averaged calculation. (Rozner & Perets 2022)

Eccentric vs. non-eccentric binaries

	eccentric $e_0 = 0.66$	circular $e_0 = 0$
timescales	rapid merger < 1 Gyr	slower merger > 1 Gyr
regimes of dominance	gas \rightarrow GWs	gas \rightarrow 3-body \rightarrow GWs
observed eccentricity	might be eccentric in LISA	circular

We can estimate the number of mergers

$$N_{\text{merge}} \sim f_{\text{disk}} f_{\text{bin,surv}} f_{\geq 20M_{\odot}} f_{\text{ret}} f_{\text{merge}} N_{\star}$$

f_{disk} – fraction of stars in the disk

$f_{\text{bin,surv}}$ – survival fraction of binaries

$f_{20M_{\odot}}$ – fraction of stars with mass $\geq 20M_{\odot}$

f_{ret} – retention fraction

f_{merge} – merger fraction (from the binaries specified above)

Now we can calculate the rate

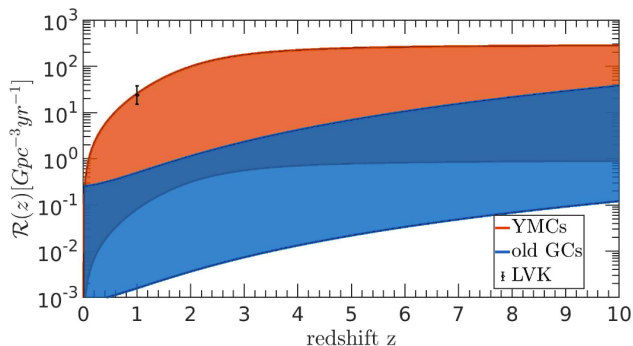


Figure: The cumulative contribution to GWs rate (Rozner & Perets 2022)

- Setting initial conditions for other GWs channels
- Pulsar retention fractions will change (Perets 2022)
- The distribution of binary separations will change
- Setting constraints on the gas amounts in GCs
- We can form short-lived binaries (Rozner, Generozov & Perets in prep.)

Takeaways

- GCs tend to host multiple generations
- Star forming environments are gas-rich
- Gas dissipation assists in hardening – leads to a production of more GWs
- It also affects indirectly on other GWs channels – by setting different initial conditions

Rozner & Perets 2022 <https://arxiv.org/pdf/2203.01330.pdf>