



Leibniz-Institut für  
Astrophysik Potsdam

# Cosmological zoom simulations of merging galaxy clusters

Thomas Berlok

6th ICM Theory and Computation Workshop, 19/8/22

# **Part I**

**Magneto-thermal instability with suppressed heat conductivity in  
mirror-unstable regions**

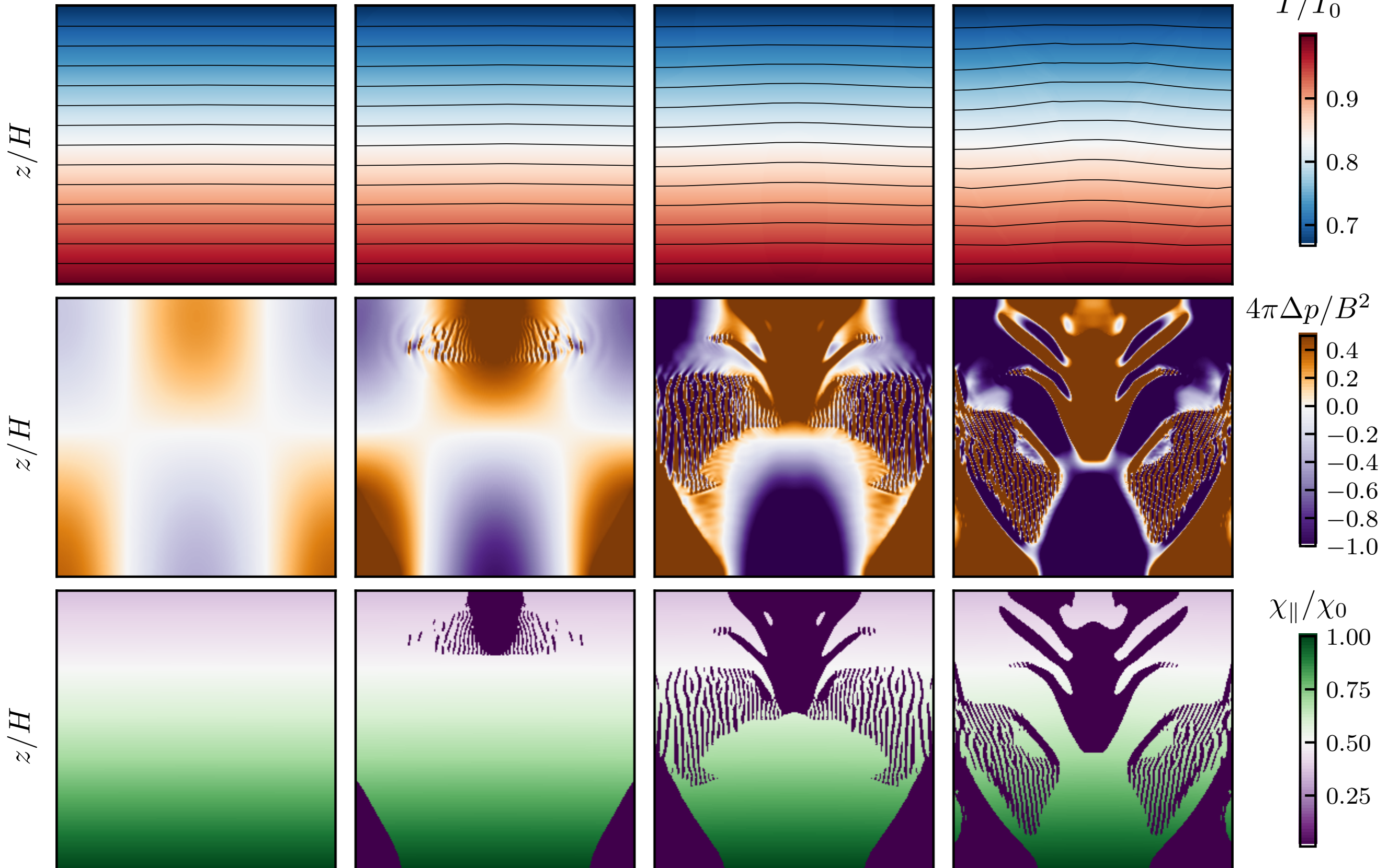
# Magneto-thermal instability

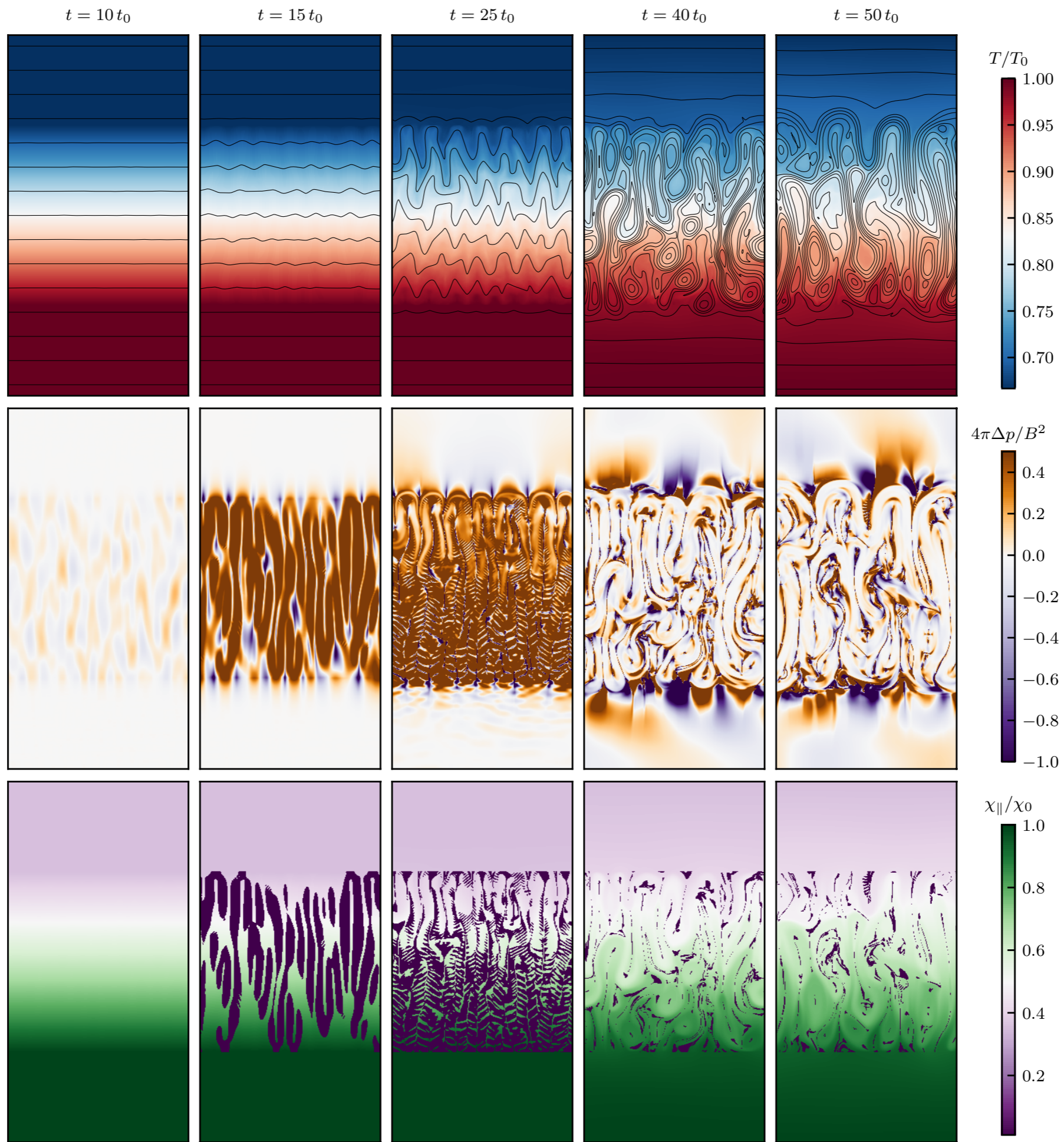
$t = 3t_0$

$t = 6t_0$

$t = 9t_0$

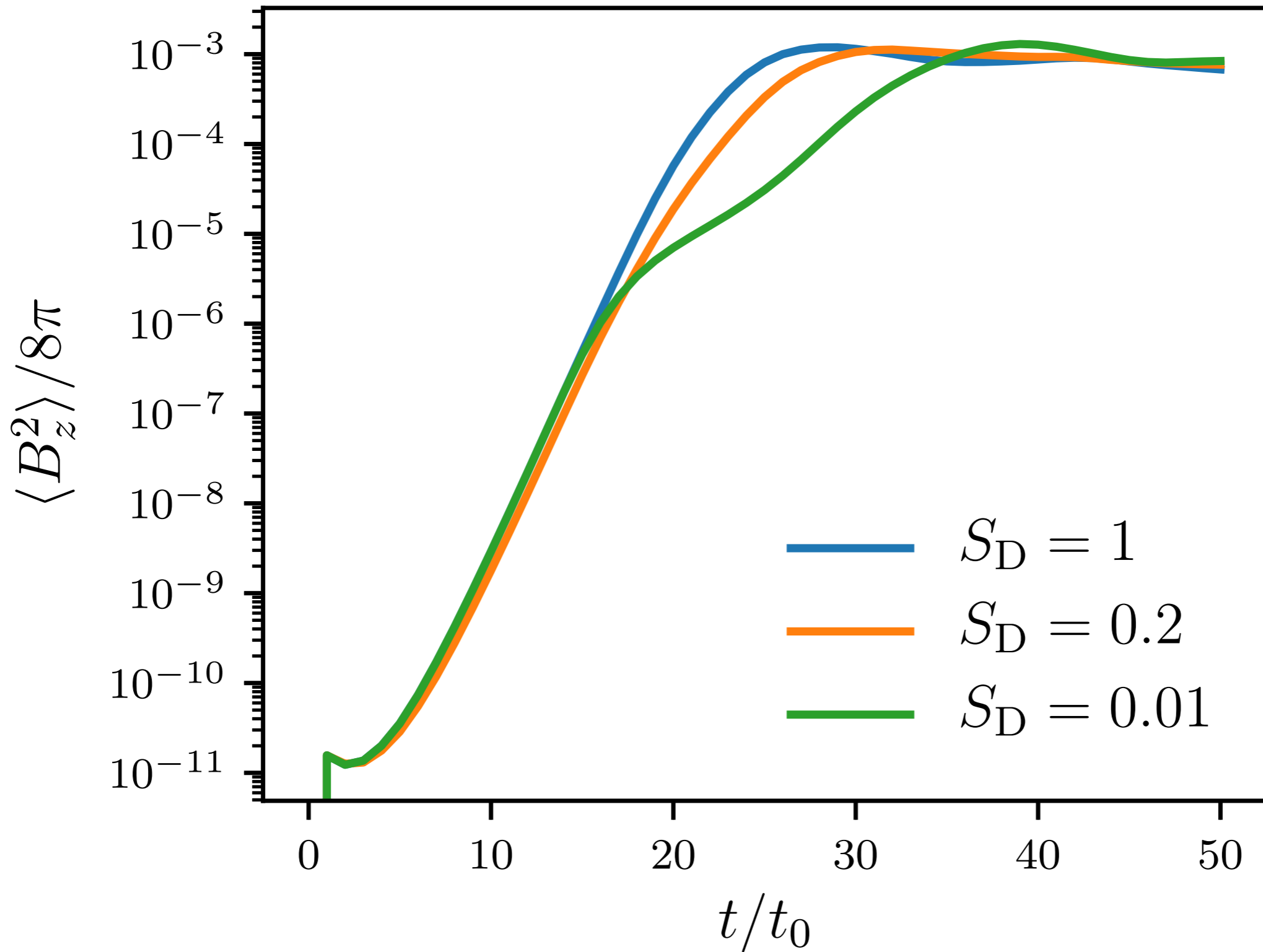
$t = 12t_0$





# Magneto-thermal instability

MTI with  $\chi_{\text{eff}} = S_D \chi_{\parallel}$  where mirror-unstable



# **Part II**

**Hydromagnetic waves in an expanding universe – cosmological  
MHD code tests using analytic solutions**

# New tests of comoving hydrodynamics/MHD

- Berlok 2022, MNRAS
- Python implementation of analytic solutions available:  
[https://github.com/tberlok/comoving\\_mhd\\_waves](https://github.com/tberlok/comoving_mhd_waves)

## Ideal MHD equations

$$\frac{d \ln \rho}{dt} = -\nabla_{\mathbf{r}} \cdot \mathbf{v} ,$$

$$\rho \frac{d\mathbf{v}}{dt} = -\nabla_{\mathbf{r}} p - \nabla_{\mathbf{r}} \cdot \left( \frac{B^2}{2} \mathbf{1} - \mathbf{B}\mathbf{B} \right) - \rho \nabla_{\mathbf{r}} \Phi ,$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla_{\mathbf{r}} \times (\mathbf{v} \times \mathbf{B}) ,$$

$$\frac{p}{\gamma - 1} \frac{d \ln(p \rho^{-\gamma})}{dt} = 0 ,$$

## Comoving MHD equations

$$\frac{d \ln \rho_c}{dt} = -\frac{1}{a} \nabla_{\mathbf{x}} \cdot \mathbf{u} ,$$

$$a \rho_c \frac{d\mathbf{u}}{dt} = -\nabla_{\mathbf{x}} p_c - \frac{1}{a} \nabla_{\mathbf{x}} \cdot \left( \frac{B_c^2}{2} \mathbf{1} - \mathbf{B}_c \mathbf{B}_c \right) - \rho_c \nabla_{\mathbf{x}} \delta \Phi - \rho_c \dot{a} \mathbf{u} ,$$

$$\frac{\partial \mathbf{B}_c}{\partial t} = \frac{1}{a} \nabla_{\mathbf{x}} \times (\mathbf{u} \times \mathbf{B}_c) ,$$

$$\frac{d\varepsilon_c}{dt} = -3 \frac{\dot{a}}{a} (\gamma - 1) \varepsilon_c - \gamma \varepsilon_c \frac{1}{a} \nabla_{\mathbf{x}} \cdot \mathbf{u} ,$$

$$\nabla_{\mathbf{x}}^2 \delta \Phi = \frac{4\pi G}{a} (\rho_{\text{tot},c} - \bar{\rho}_{\text{tot},c})$$

## Substitutions

$$\mathbf{r} = a\mathbf{x}, \quad \rho_c = \rho a^3, \quad \mathbf{u} = a\dot{\mathbf{x}},$$

$$B_c = B a^2, \quad \varepsilon_c = \varepsilon a^3$$

$$\varepsilon = p / (\gamma - 1)$$

## Friedmann equation

$$\left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \bar{\rho}_{\text{tot}}$$

$$\bar{\rho}_{\text{tot}} = \left( \frac{\Omega_{\text{r},0}}{a^4} + \frac{\Omega_{\text{m},0}}{a^3} + \Omega_{\Lambda,0} \right) \rho_{\text{crit},0}$$

$$\rho_{\text{crit},0} = 3H_0^2 / (8\pi G)$$



## Linear theory (Berlok, MNRAS, 2022)

$$c_s \equiv \sqrt{\frac{\gamma p}{\rho}} = \sqrt{\frac{\gamma p_{c,0}}{\rho_c}} a^{-3(\gamma-1)/2},$$

$$v_A \equiv \frac{B}{\sqrt{\rho}} = \frac{B_c}{\sqrt{\rho_c}} a^{-1/2},$$

$$v_g \equiv \frac{\sqrt{4\pi G \rho}}{k'} = \frac{\sqrt{4\pi G \rho_c}}{k} a^{-1/2}$$

$$k' = k/a$$

### Standard MHD result

Alfvén wave

$$\omega = k' v_A$$

Magnetosonic wave

$$\omega = k' (c_s^2 + v_A^2 - v_g^2)^{1/2}$$

## Useful definitions

$$\mathcal{V}_s \equiv \sqrt{\frac{\gamma p_{c,0}}{\rho_c}}, \quad \mathcal{V}_A \equiv \frac{B_c}{\sqrt{\rho_c}}, \quad \mathcal{V}_g \equiv \frac{\sqrt{4\pi G \rho_c}}{k},$$

$$\Omega_s \equiv \frac{k \mathcal{V}_s}{H_0}, \quad \Omega_A \equiv \frac{k \mathcal{V}_A}{H_0}, \quad \Omega_g \equiv \frac{k \mathcal{V}_g}{H_0}.$$

## Comoving Alfvén wave

$$\begin{array}{ccc}
 \frac{\partial}{\partial t} \frac{\delta B_c}{B_c} = \frac{ik}{a} \delta u & & \frac{\partial}{\partial a} \frac{\delta B_c}{B_c} = \frac{ik}{H_0 a^{1/2}} \delta u \\
 \frac{\partial(a\delta u)}{\partial t} = ik \frac{B_c^2}{a\rho_c} \frac{\delta B_c}{B_c} & \xrightarrow[\dot{a} = H_0/\sqrt{a}]{\text{Einstein-de-Sitter}} & \frac{\partial(a\delta u)}{\partial a} = \frac{ik\mathcal{V}_A^2}{H_0 a^{1/2}} \frac{\delta B_c}{B_c}
 \end{array}$$

## Euler equation

$$\frac{\partial^2}{\partial a^2} \frac{\delta B_c}{B_c} + \frac{3}{2a} \frac{\partial}{\partial a} \frac{\delta B_c}{B_c} + \frac{\Omega_A^2}{a^2} \frac{\delta B_c}{B_c} = 0$$

$$\frac{\delta B_c}{B_c} = a^{-1/4} \left( c_1 e^{i\kappa \ln a} + c_2 e^{-i\kappa \ln a} \right)$$

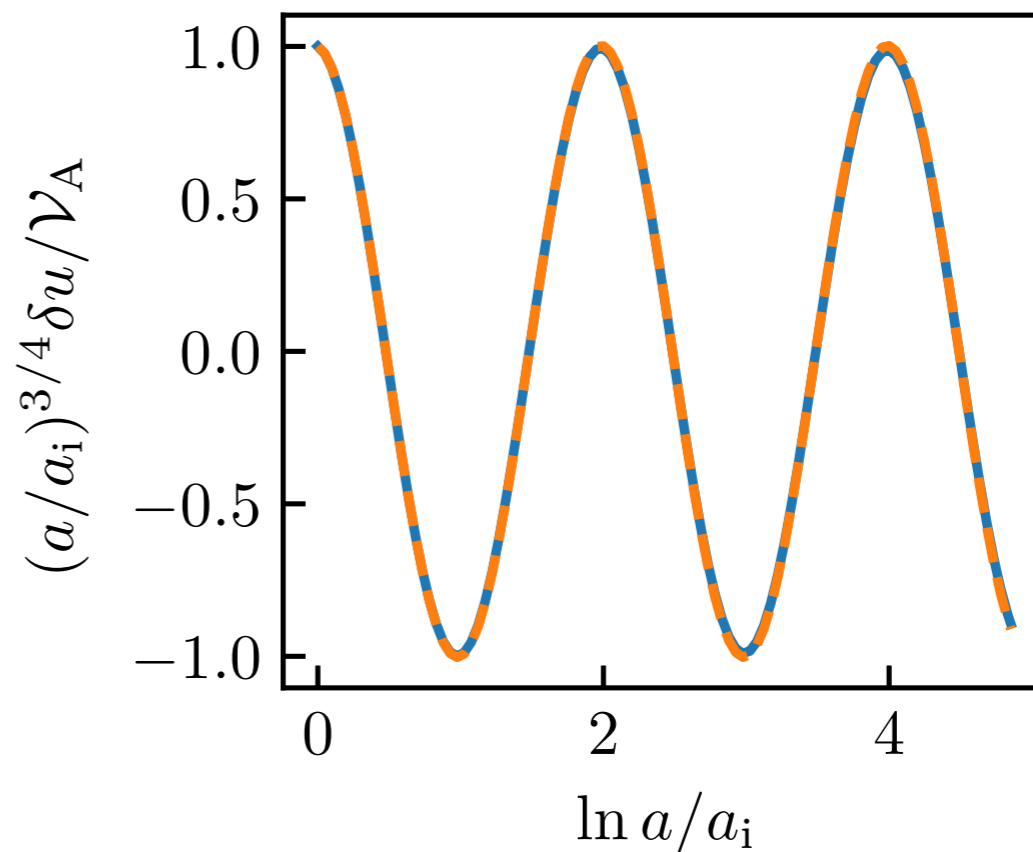
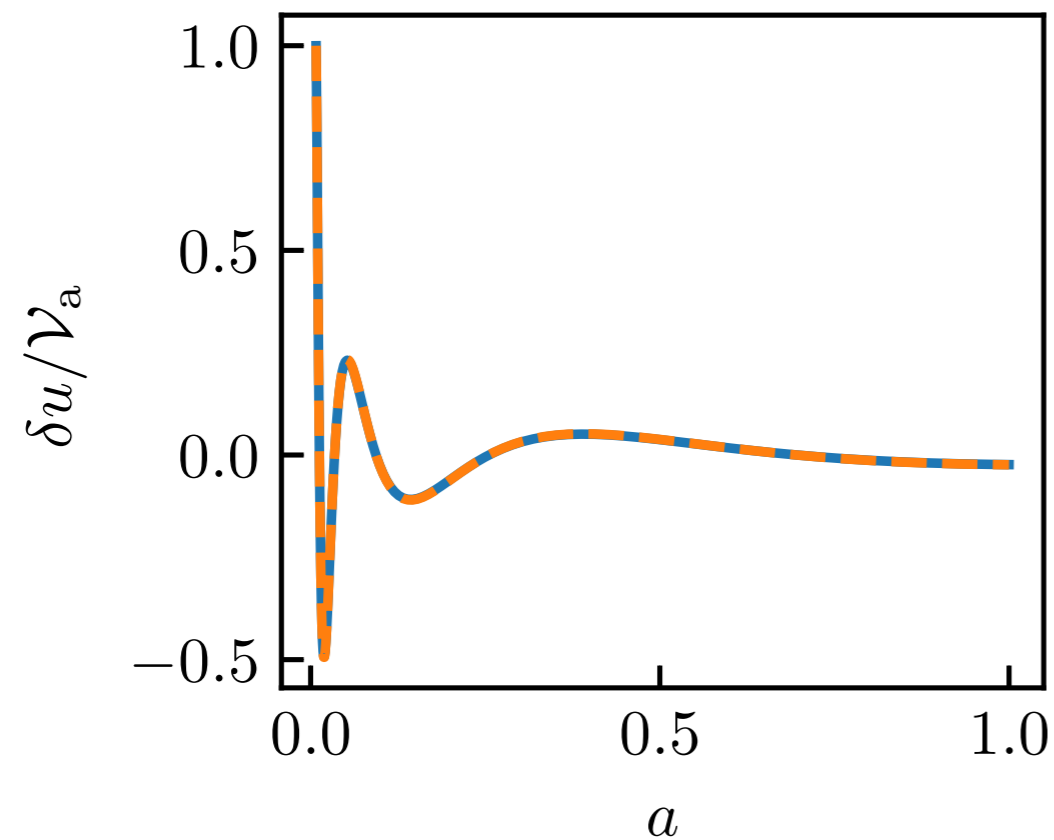
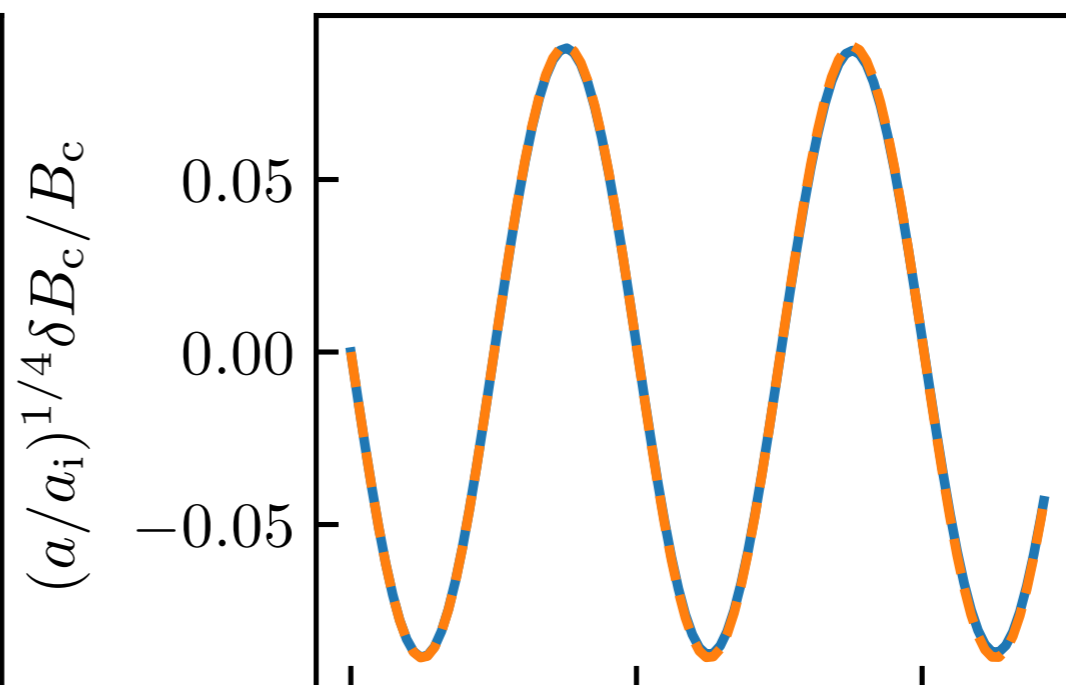
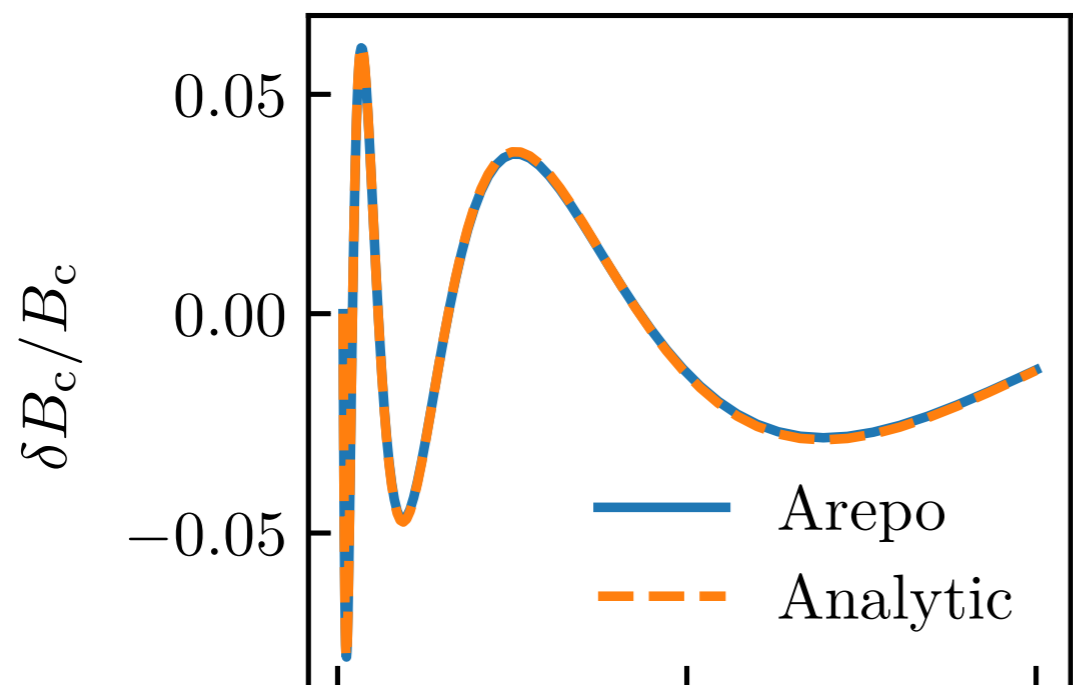
$$\kappa \equiv \sqrt{\Omega_A^2 - \frac{1}{16}}$$

# Comoving Alfvén wave

$$\frac{\delta B_c(x, a)}{B_c} = -A_u \left( \frac{a}{a_i} \right)^{-1/4} \frac{\sqrt{a_i} \Omega_A}{\kappa} \sin(\psi) \sin(kx),$$

$$\psi = \kappa \ln(a/a_i)$$

$$\frac{\delta u(x, a)}{\mathcal{V}_A} = A_u \left( \frac{a}{a_i} \right)^{-3/4} \left( \cos(\psi) - \frac{\sin(\psi)}{4\kappa} \right) \cos(kx)$$



## Magnetosonic wave

$$\frac{\partial}{\partial a} \frac{\delta \rho_c}{\rho_c} = - \frac{ik}{H_0 \sqrt{a}} \delta u ,$$

$$\frac{\partial(a\delta u)}{\partial a} = - \frac{ik\sqrt{a}}{H_0} \left( \frac{\mathcal{V}_s^2}{a^{3(\gamma-1)}} + \frac{\mathcal{V}_A^2 - \mathcal{V}_g^2}{a} \right) \frac{\delta \rho_c}{\rho_c}$$

## Gravitational instability

$\gamma = 1$  Thermal pressure term does not decay

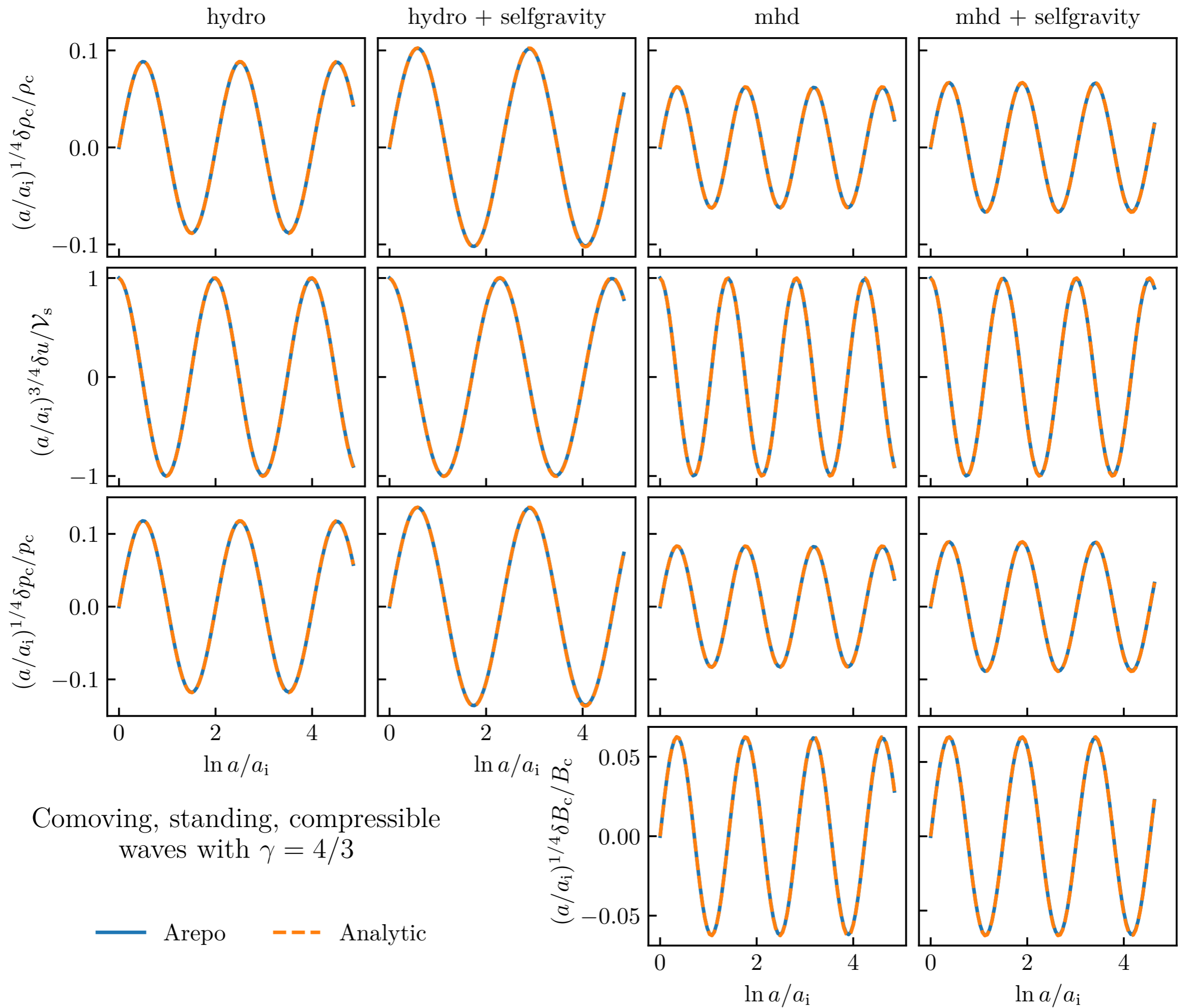
$\gamma = 4/3$  Thermal pressure term decays at same rate as other terms

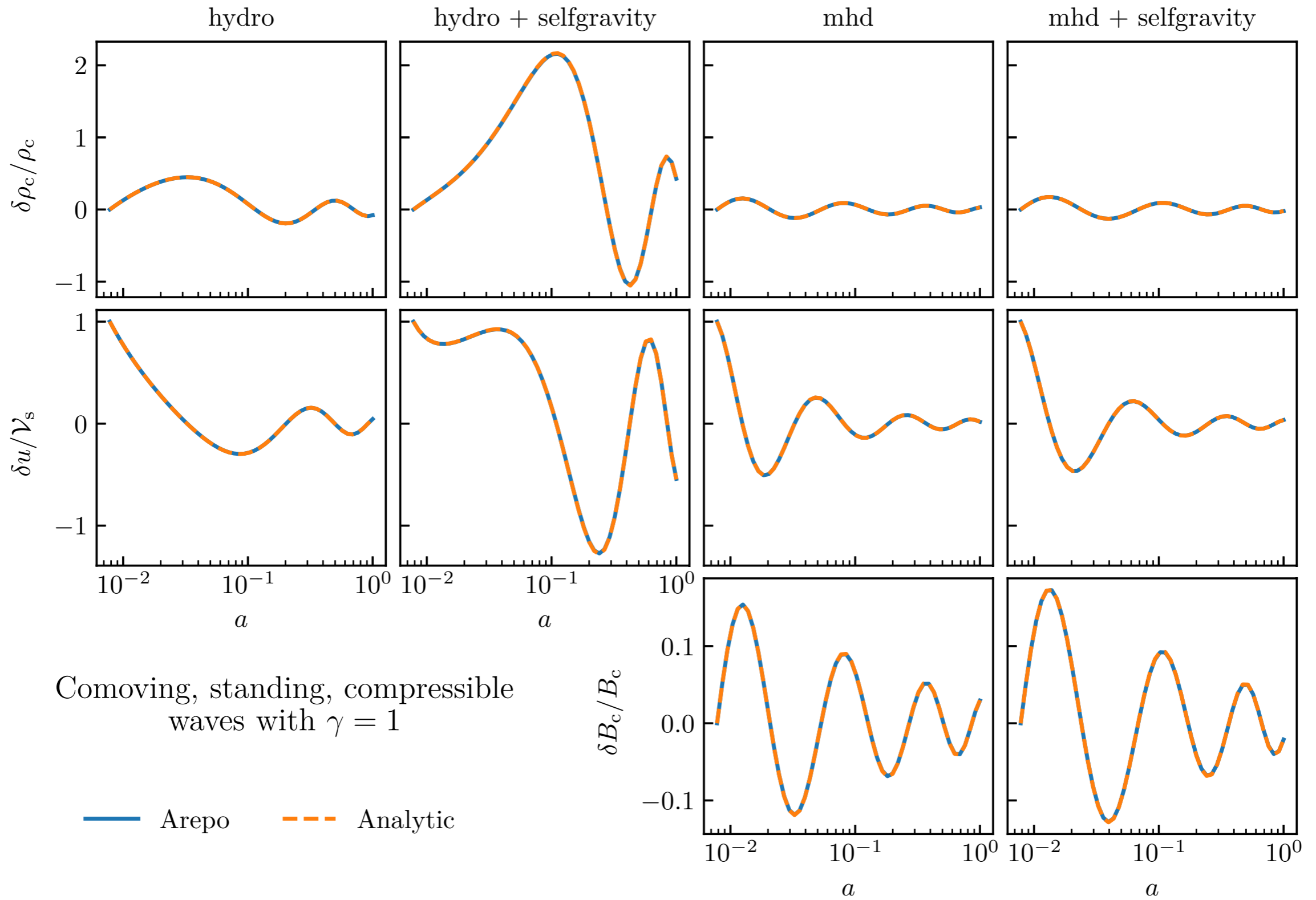
$\gamma = 5/3$  Thermal pressure term decays faster than other terms

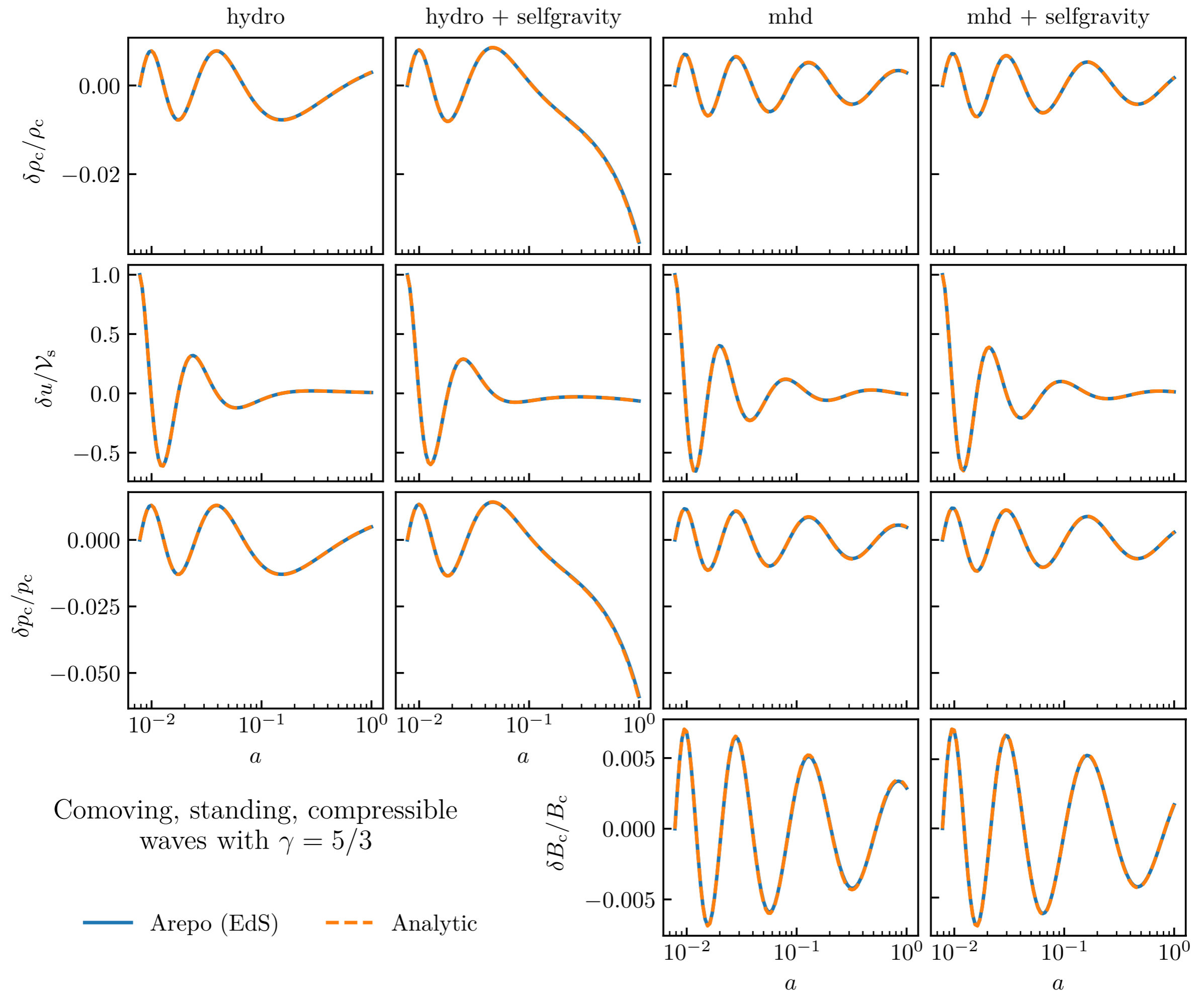
## Differential equation

$$\frac{\partial^2}{\partial a^2} \frac{\delta \rho_c}{\rho_c} + \frac{3}{2a} \frac{\partial}{\partial a} \frac{\delta \rho_c}{\rho_c} + \left( \frac{\Omega_s^2}{a^{3\gamma-2}} + \frac{\Omega_A^2 - \Omega_g^2}{a^2} \right) \frac{\delta \rho_c}{\rho_c} = 0$$

$\gamma = 4/3$  is an Euler ODE, in general a transformed Bessel equation!







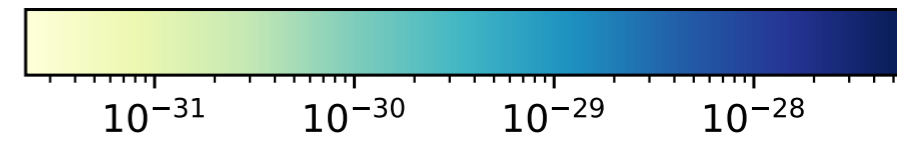
# **Part III**

**Zoom simulations of merging galaxy clusters**



# Zooming in on a galaxy cluster

Gas density [ $\text{g cm}^{-3}$ ]



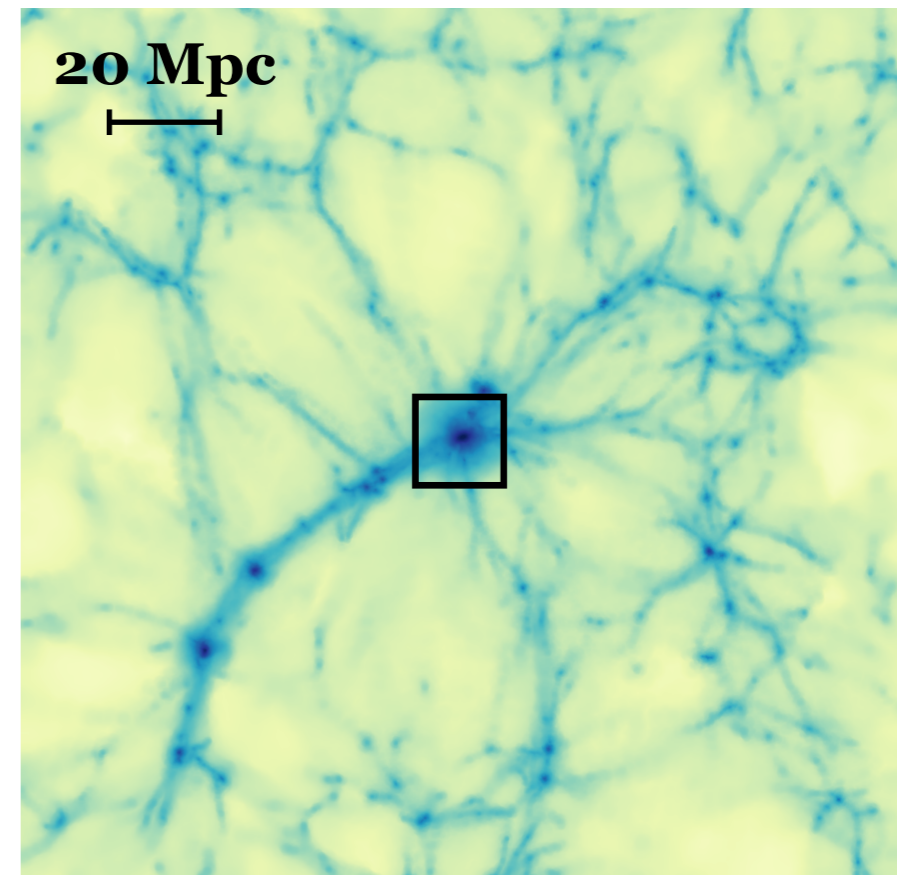
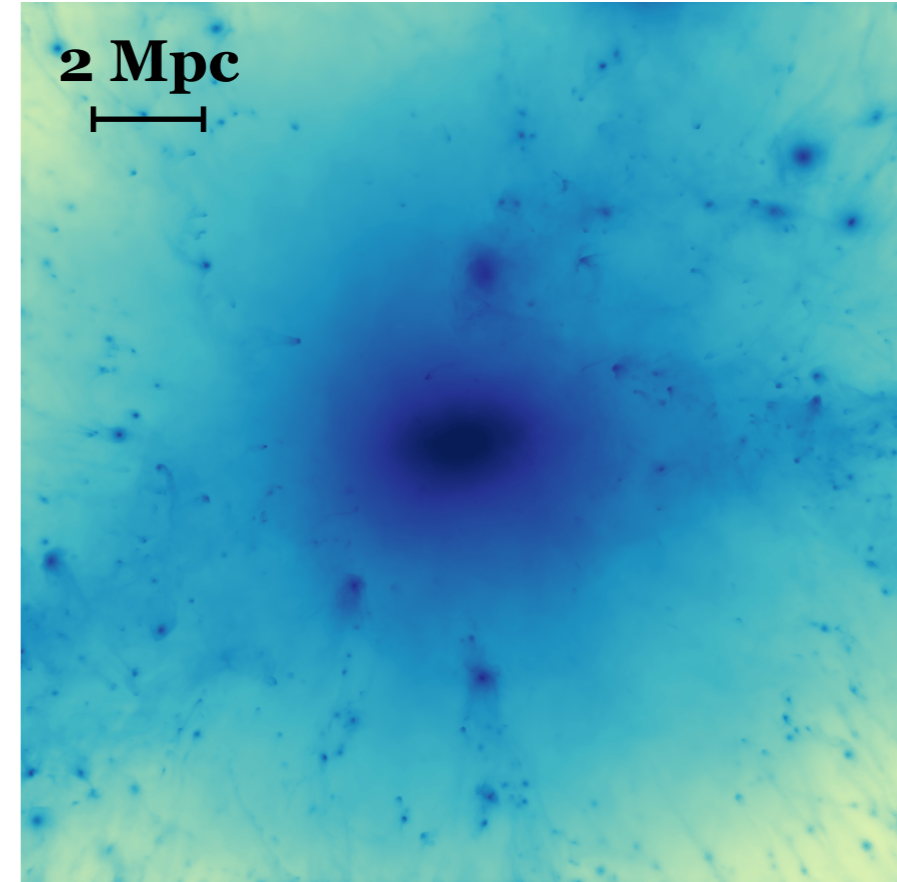
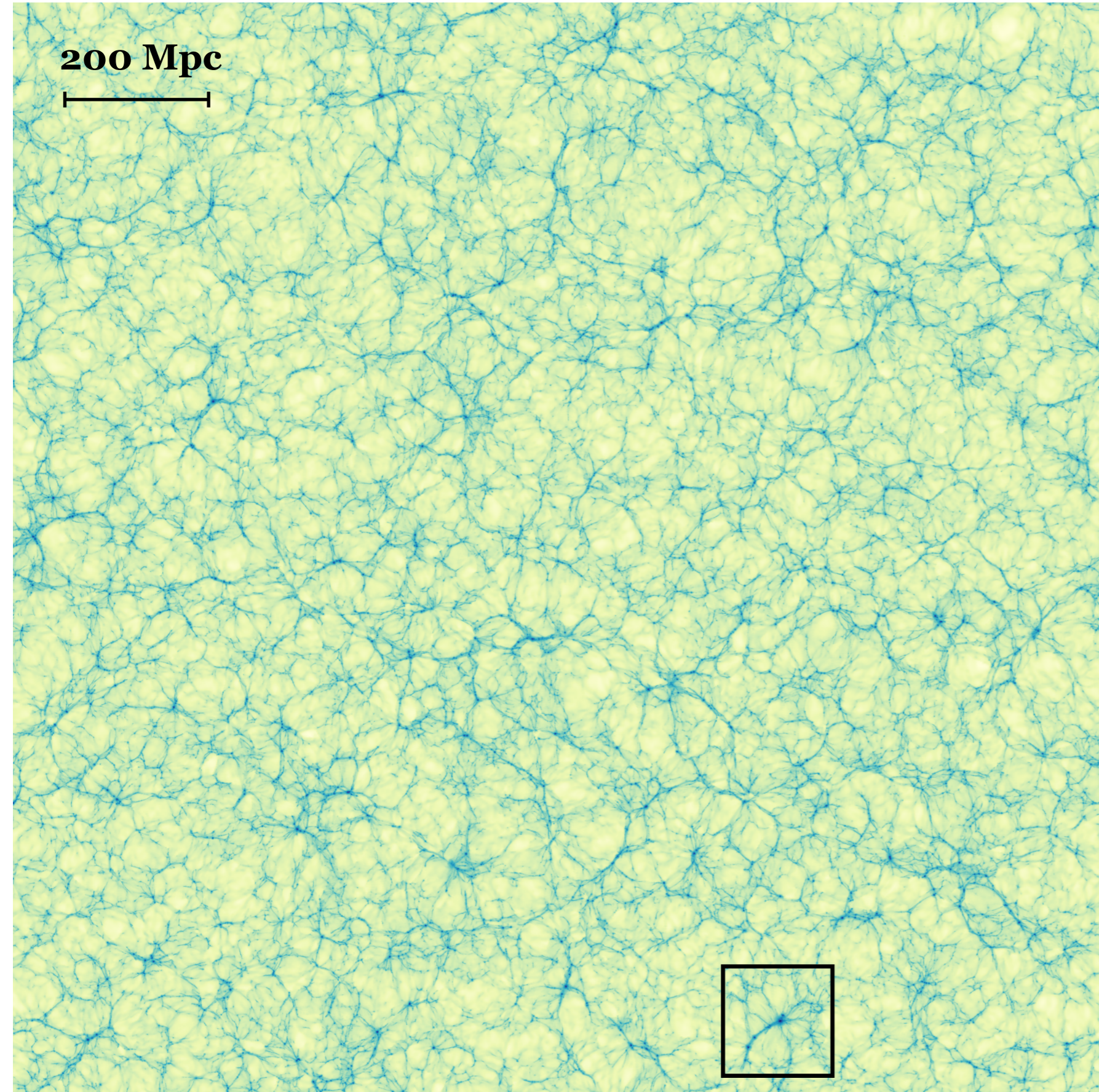
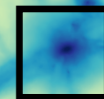
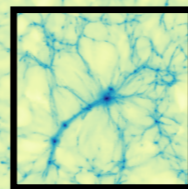
200 Mpc



2 Mpc



20 Mpc

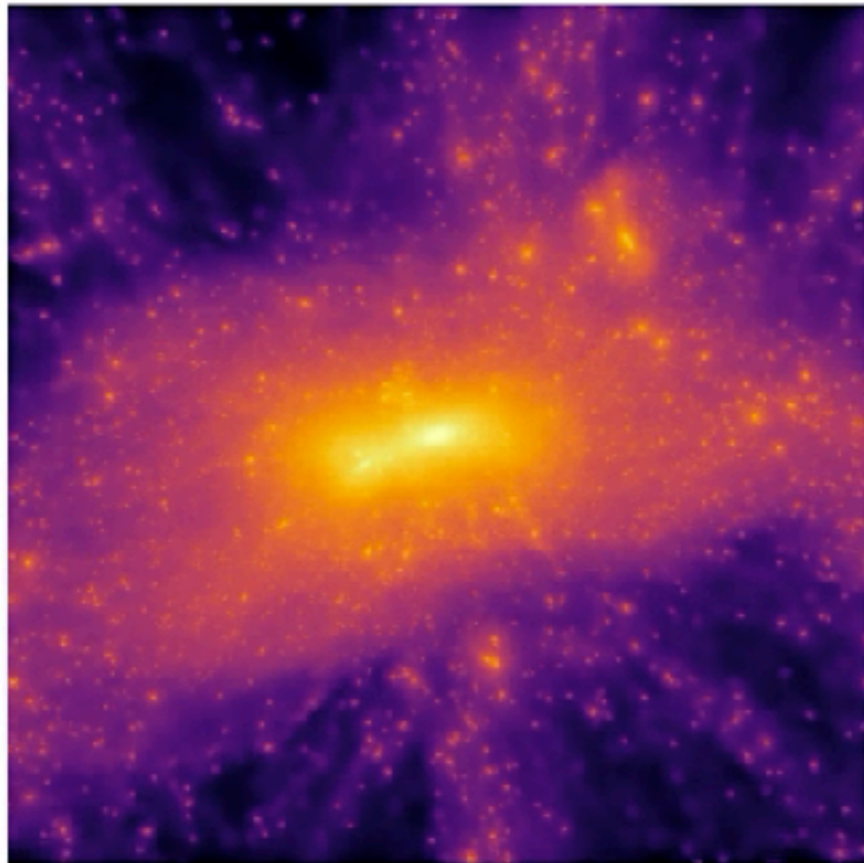


# Major merger including galaxy formation

Size: 39.73 mio. ly

Age: 11.06 Gyr

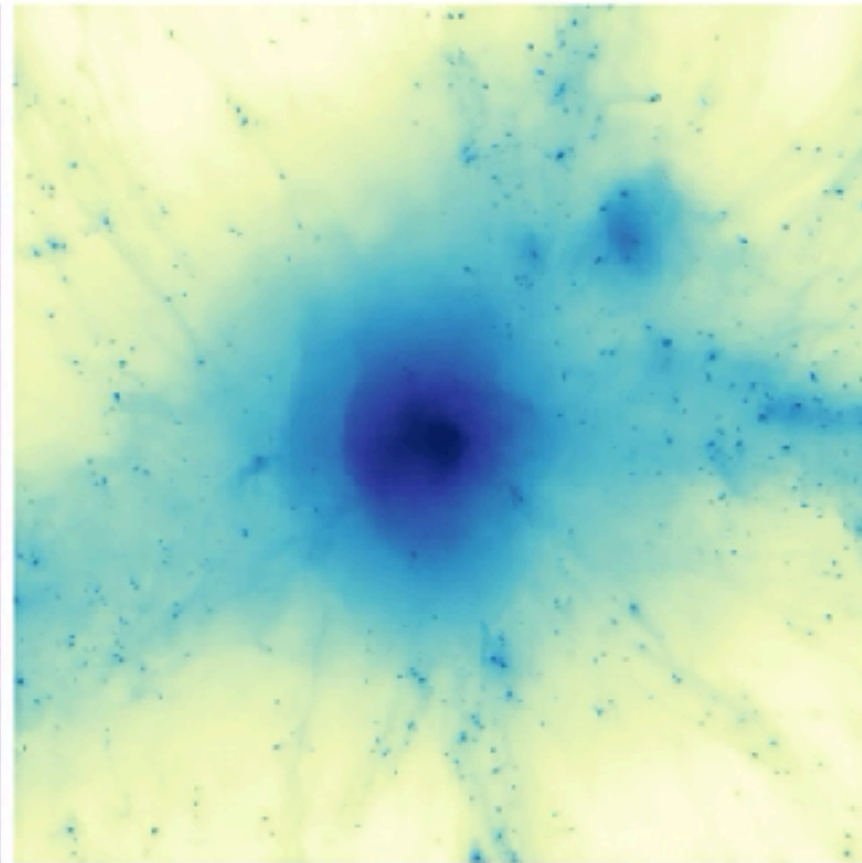
Dark matter density



$10^{-29}$   $10^{-27}$

$\text{g/cm}^3$

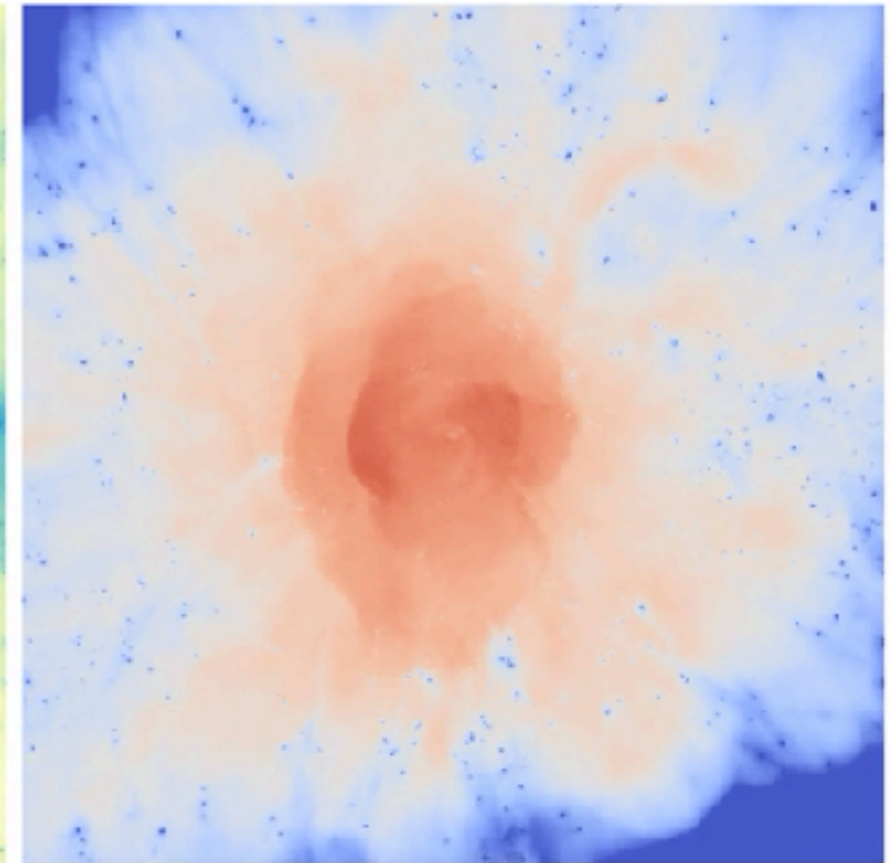
Gas density



$10^{-30}$   $10^{-28}$

$\text{g/cm}^3$

Temperature



$10^6$   $10^8$

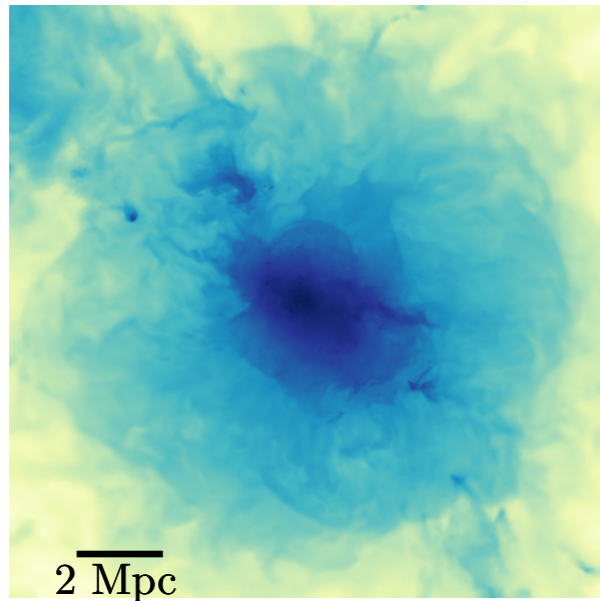
Kelvin/Celsius

11 mio cpu-h on SuperMUC-NG in Germany (PI: T. Berlok)

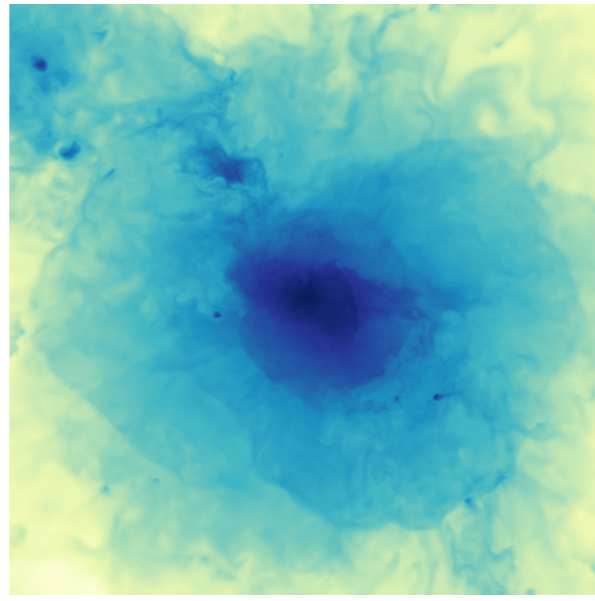
In collaboration with:

Joseph Whittingham, Léna Jlassi, Larissa Tevlin, Martin Sparre, Rainer Weinberger, Ewald Puchwein, Rüdiger Pakmor and Christoph Pfrommer

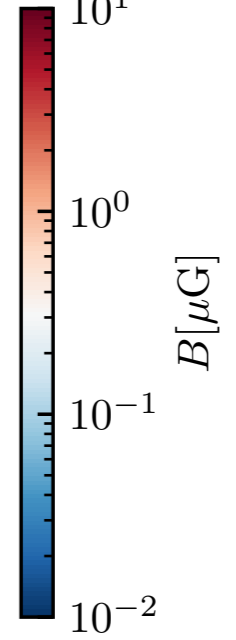
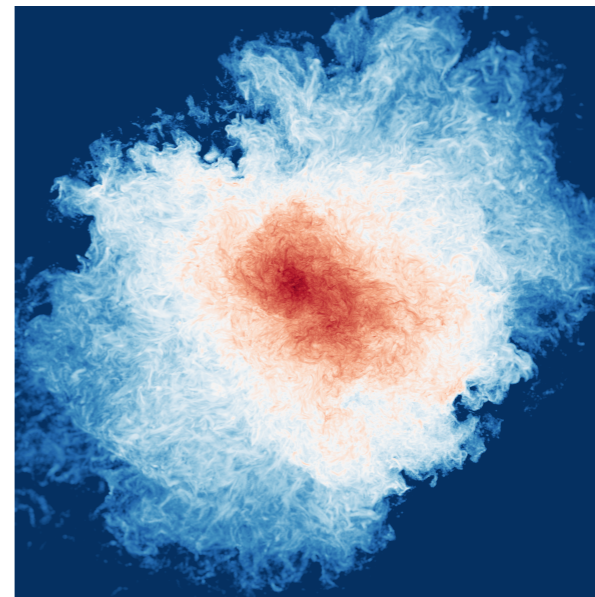
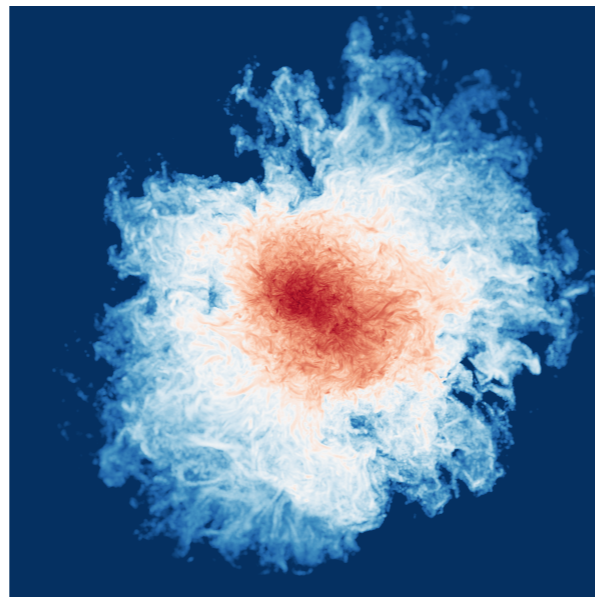
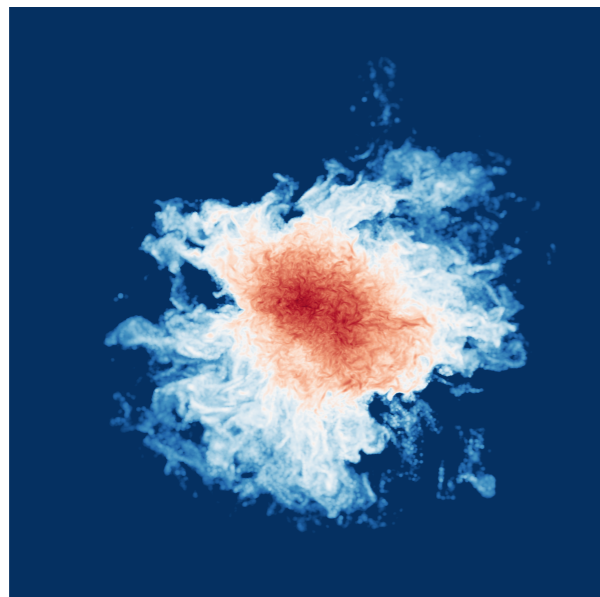
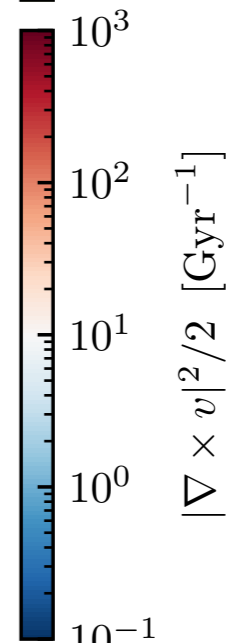
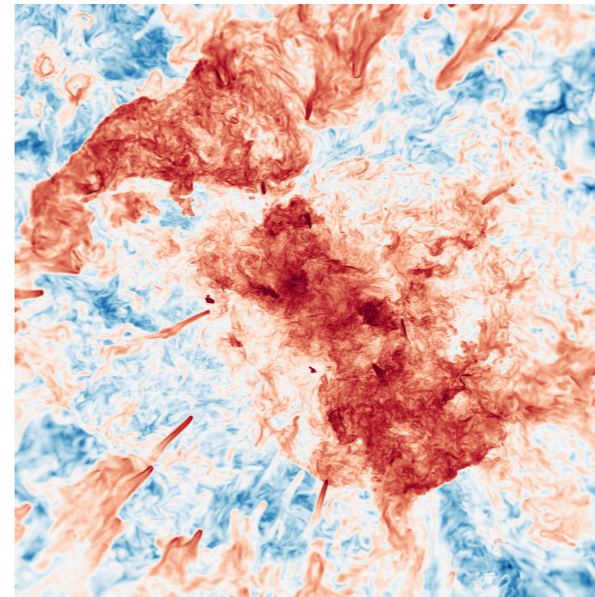
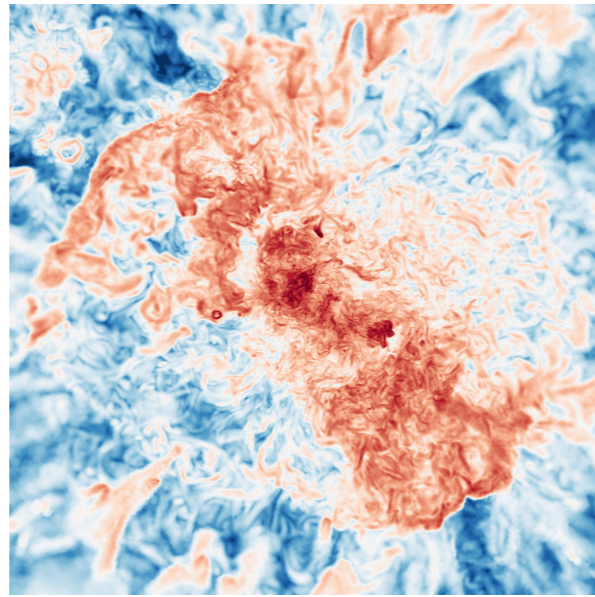
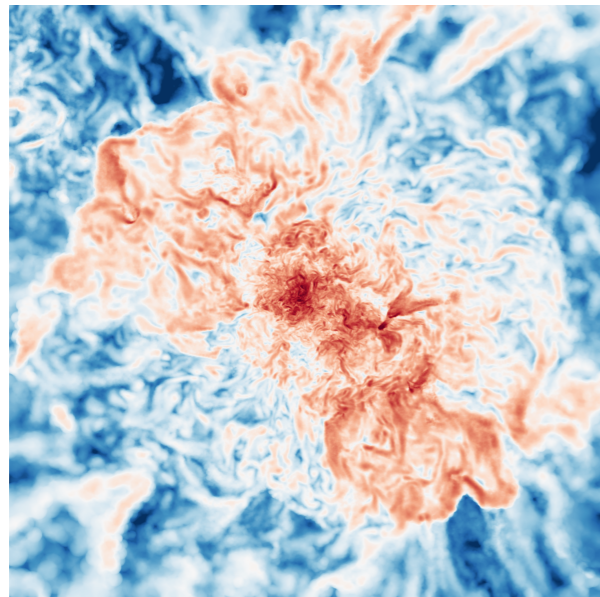
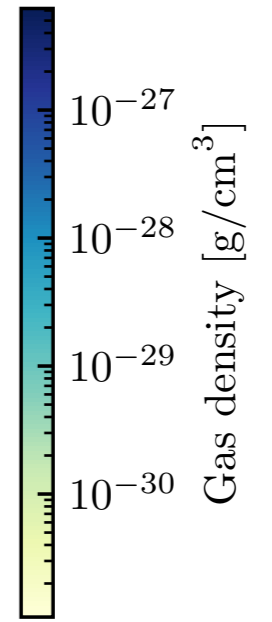
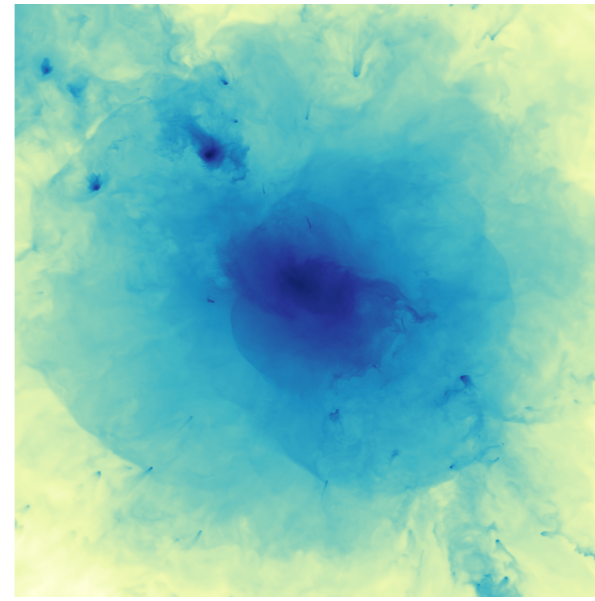
Zoom factor 8



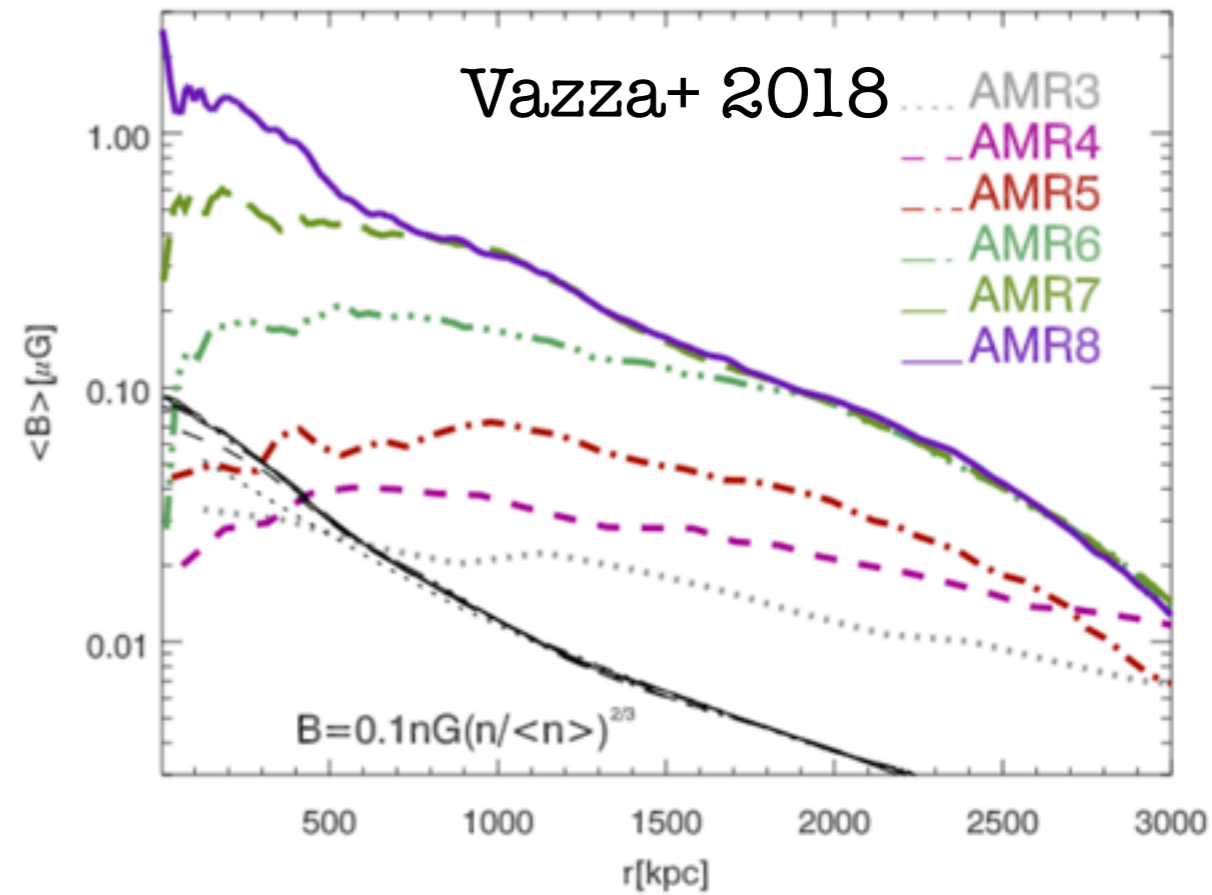
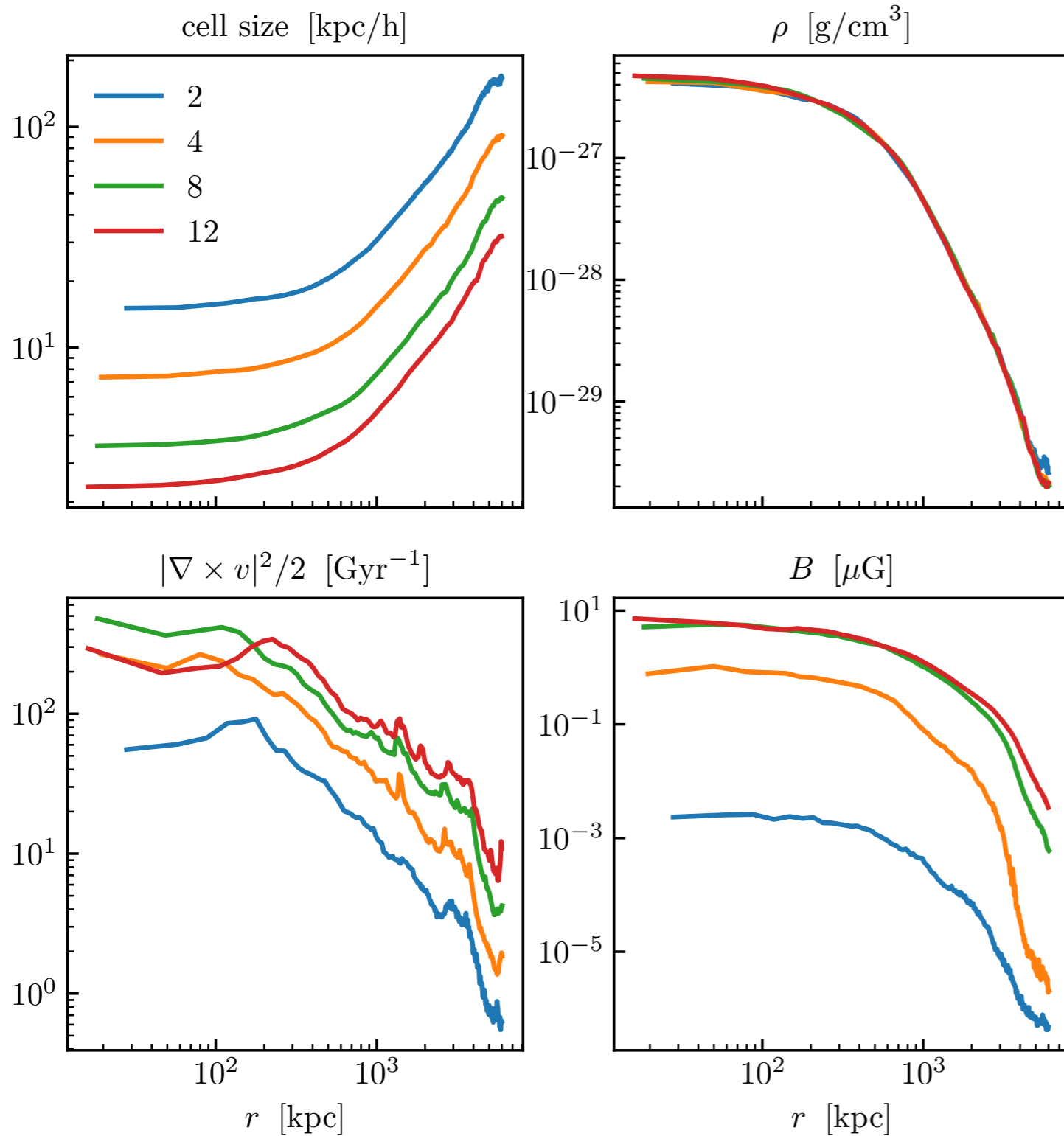
Zoom factor 12



Zoom factor 24



# Resolving the magnetic dynamo is expensive



# **Part IV**

**Braginskii viscosity in Arepo**

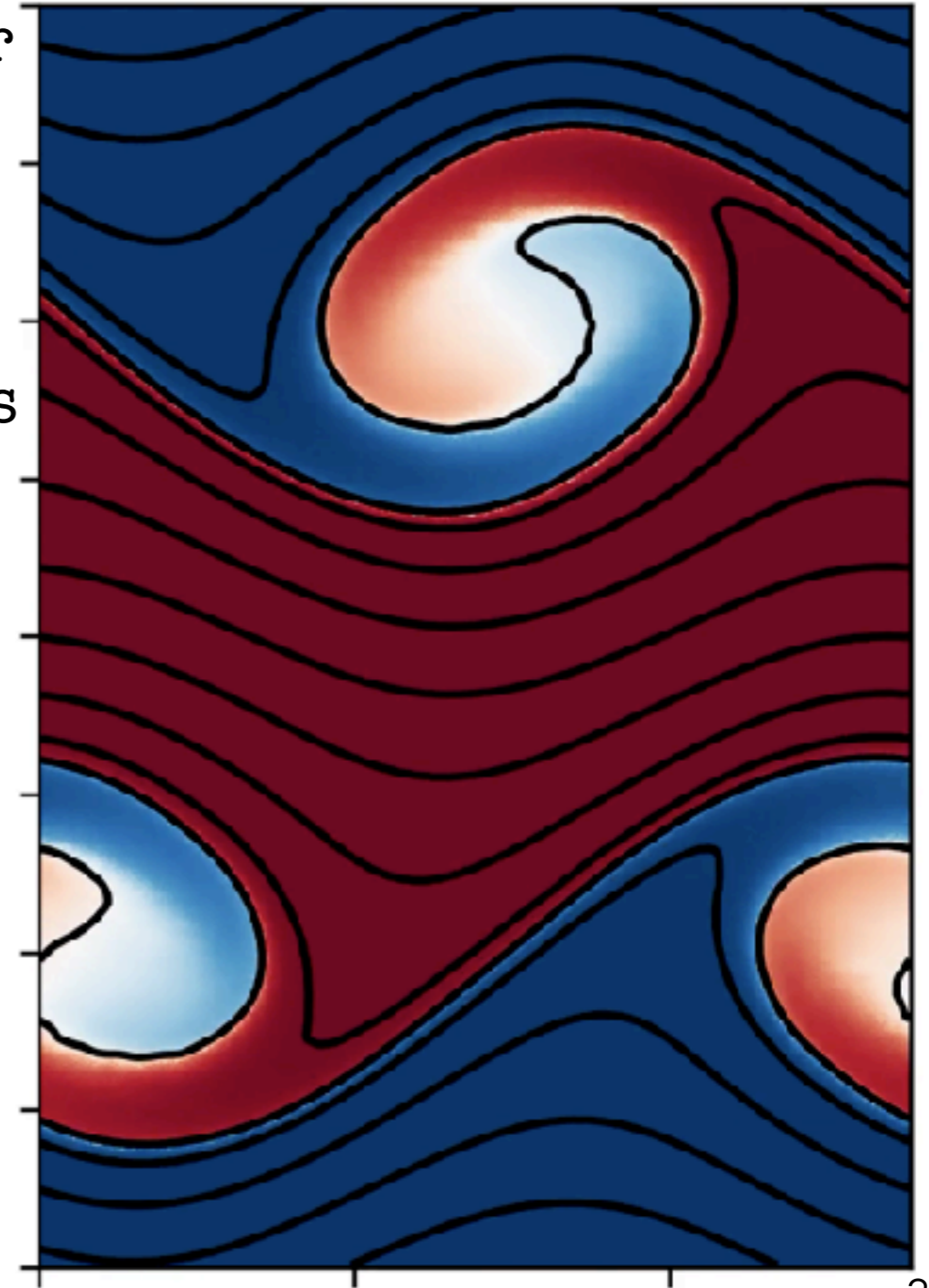
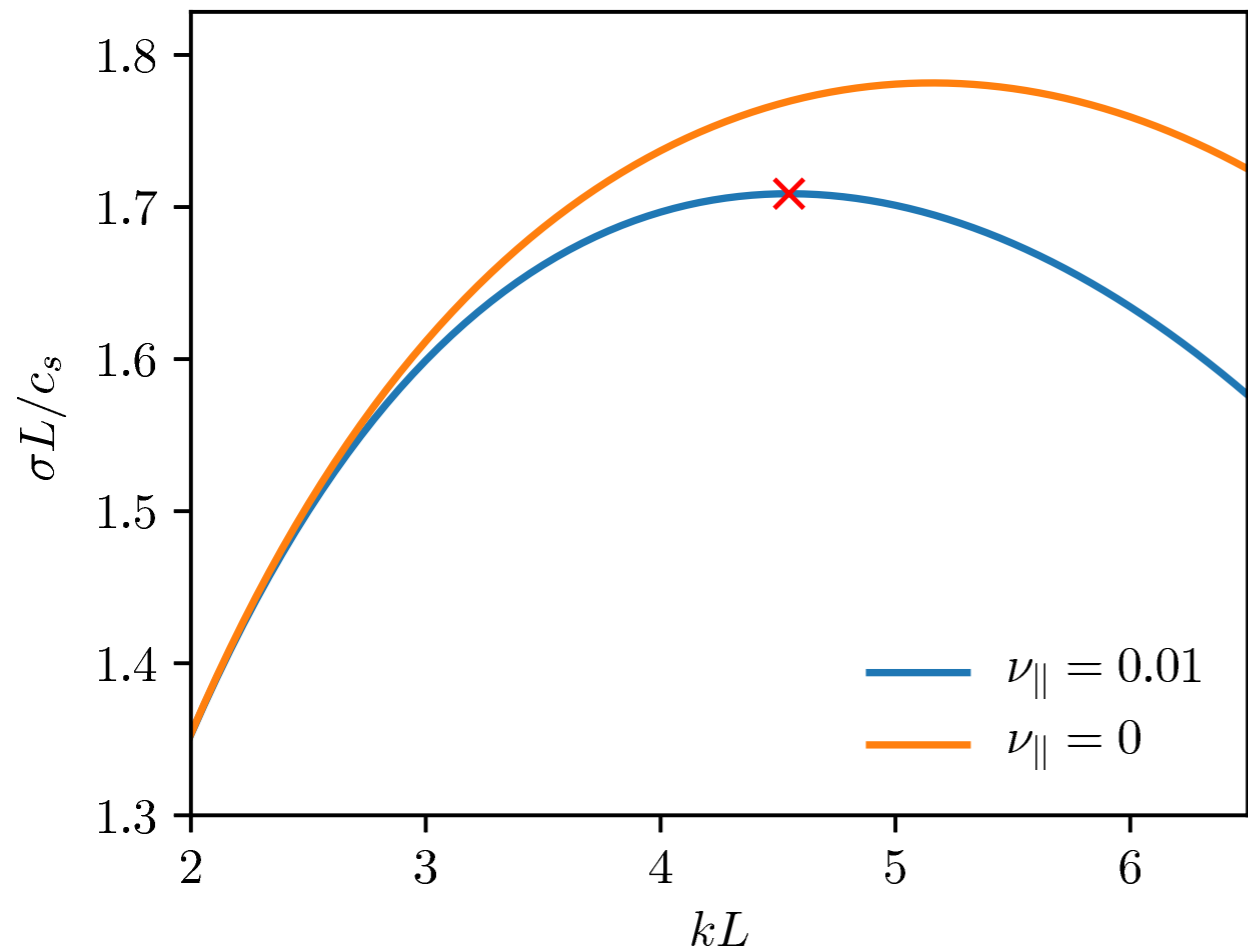
# Braginskii viscosity in Arepo

$$\Pi = -\Delta p \left( bb - \frac{1}{3} \right)$$

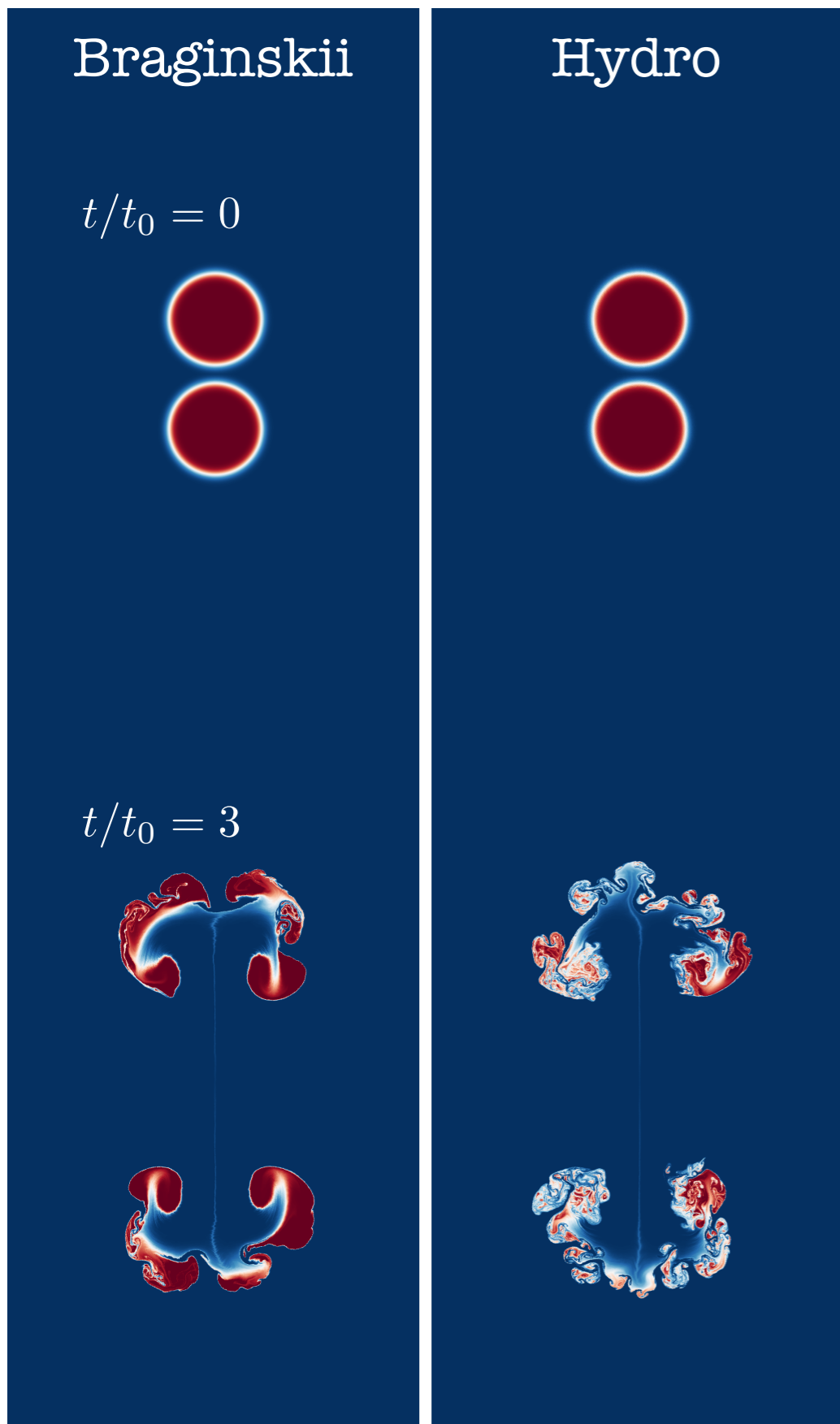
$$\frac{\partial \rho v}{\partial t} = -\nabla \cdot \Pi$$

$$\frac{\partial E}{\partial t} = -\nabla \cdot (\Pi \cdot v)$$

- Extensive test suite (detailed in Berlok, Pakmor & Pfrommer 2020)
- Second order accurate Super timestepping (RKL2)
- Recently extended for cosmological applications

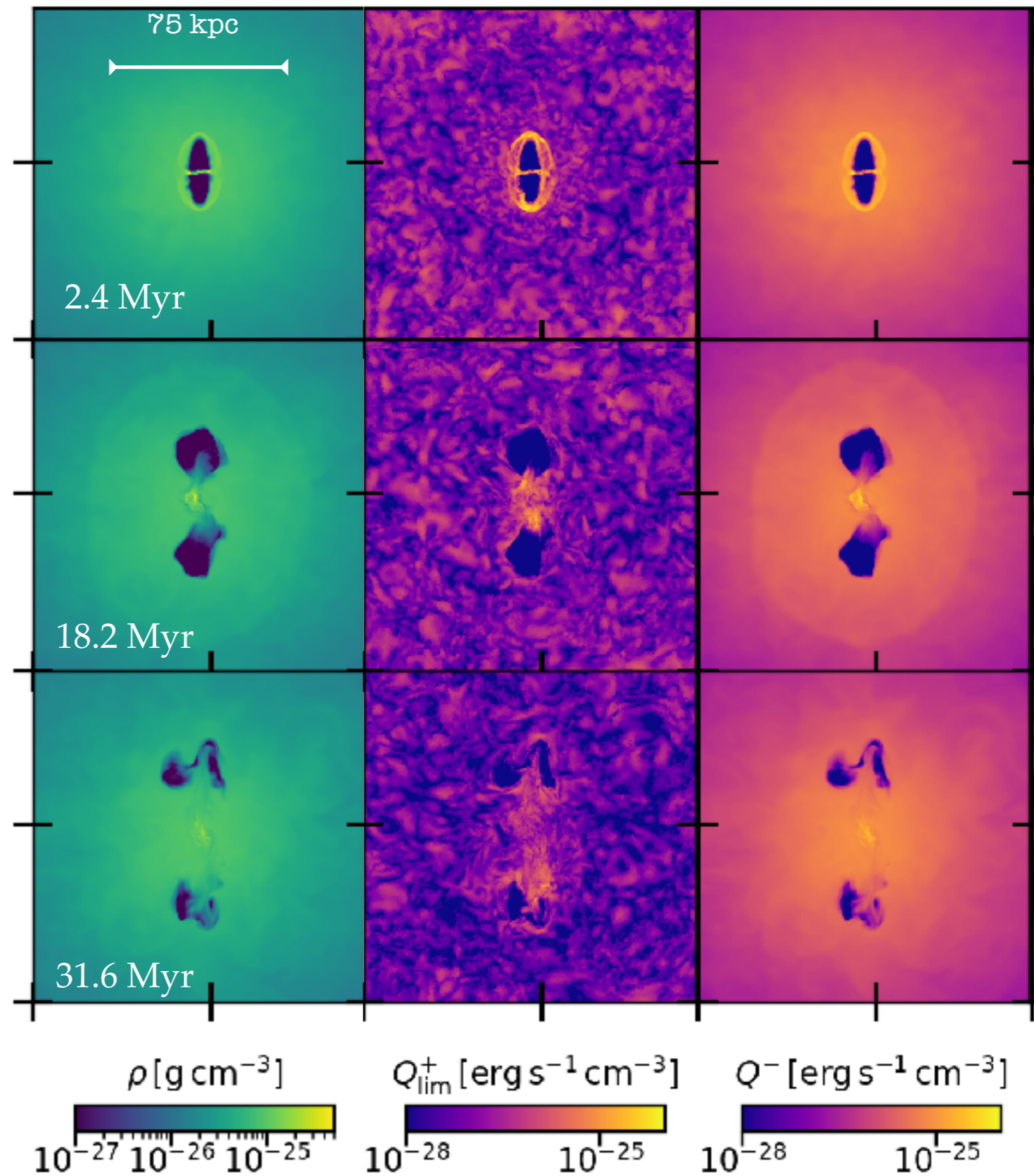


# Idealized 2D bubbles



Berlok+ (unpublished)

# AGN jet with Braginskii viscosity



Berlok+, in prep.

**Analysis to be done!**

