

Gravity, Twistors & the MHV Formalism

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Based on arXiv:0808.3907 [hep-th] with L. Mason

Hidden Structures in Field Theory Amplitudes

12th September 2008

Introduction

MHV amplitudes $\langle + + - - \dots - - \rangle$ are much simpler than other helicity amplitudes

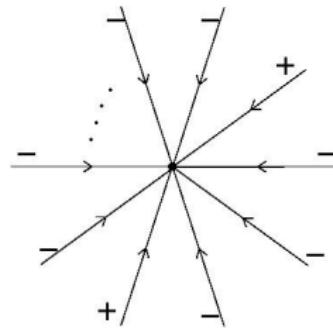
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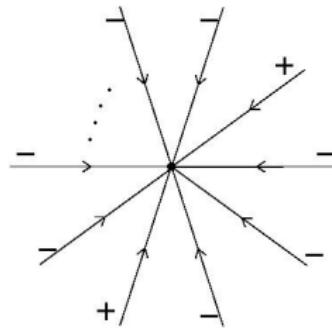
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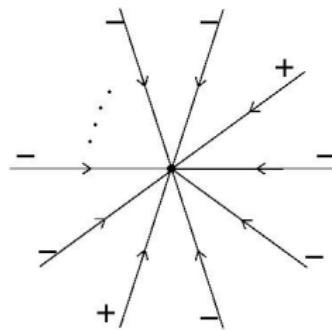


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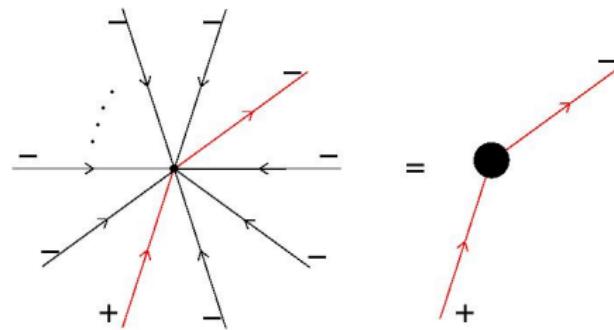


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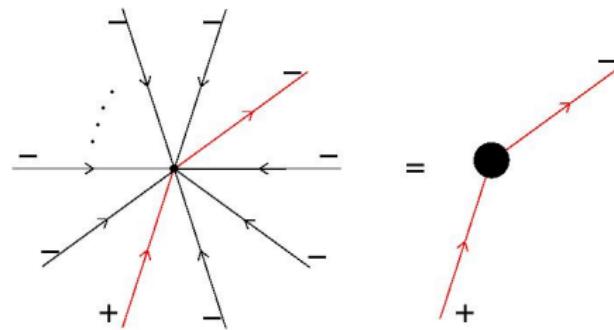


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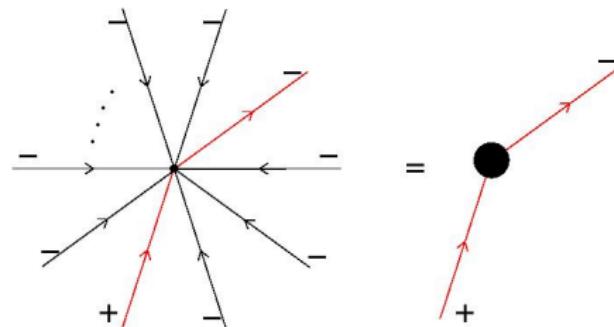
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- Today: focus on gravity

MHV Amplitudes & ASD Spacetimes



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$$S = \frac{1}{\kappa^2} \int_M \Sigma^{\dot{\alpha}\dot{\beta}} \wedge (d\Gamma + \Gamma \wedge \Gamma)_{\dot{\alpha}\dot{\beta}}$$

- $\Sigma^{\dot{\alpha}\dot{\beta}} = e^{\alpha(\dot{\alpha}} \wedge e_\alpha^{\dot{\beta})}$ ($e^{\alpha\dot{\alpha}} = \sigma_a^{\alpha\dot{\alpha}} e_\mu^a dx^\mu$ are vierbein 1-forms)
- $\Gamma^{\dot{\alpha}}_{\dot{\beta}}$ is an **independent** connection on \mathbb{S}^+

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- Can be proved more rigorously using geometric/canonical quantization

Relation to Usual Form of Amplitudes

Claim:

$$\frac{i}{\hbar \kappa^2} \int_M \Sigma_0 \wedge \gamma \wedge \gamma = \sum_n (n\text{-particle BGK amplitude})$$

where BGK amplitude = $i\kappa^{n-2}/\hbar \times \delta(\sum p_i) \mathcal{M}_n^{\text{BGK}}$ and

$$\mathcal{M}_n^{\text{BGK}} = [1n]^8 \left\{ \frac{\langle 12 \rangle \langle n-2 \ n-1 \rangle}{[1 \ n-1]} \frac{F}{N(n)} \prod_{i=1}^{n-3} \prod_{j=i+2}^{n-1} [ij] + P_{(2, \dots, n-2)} \right\}$$

- $N(n) := \prod_{i < j} [ij]$
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(Integral is gauge invariant because of field eqns)

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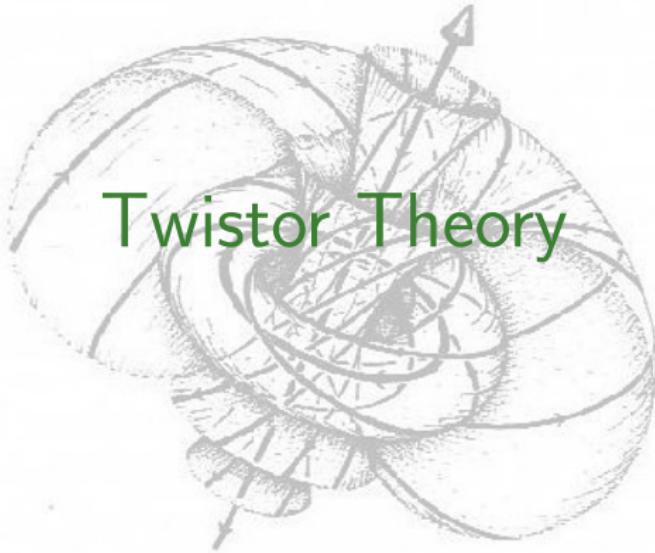
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Twistor space is particularly well-suited to exploiting this integrability (integrability in spacetime = holomorphy in twistor space)

Twistor Theory



Linearized gravity in twistor space

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$$(\delta R)_{\alpha\beta\gamma\delta}(x) = \int_{L_x} [\pi \, d\pi] \wedge \frac{\partial h_2}{\partial \omega^\alpha \partial \omega^\beta \partial \omega^\gamma \partial \omega^\delta} \Big|_{L_x}$$

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- Penrose transform of self-dual spin connection is

$$\gamma^{\dot{\alpha}}_{\dot{\beta}}(x) = \int_{L_x} [\pi \, d\pi] \wedge \pi^{\dot{\alpha}} \pi_{\dot{\beta}} B|_{L_x}$$

where $B \in \Omega^{1,1}(\mathbb{PT}', \mathcal{O}(-4))$, rather than $(0,1)$ -form

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- $\epsilon^{\alpha\beta} \partial_\alpha B_\beta = \tilde{h}_{-6}$ ensures $d\gamma^{\dot{\alpha}}_{\dot{\beta}} = (\delta R)^{\dot{\alpha}}_{\dot{\beta}\dot{\gamma}\dot{\delta}} dx^{\gamma\dot{\gamma}} \wedge dx_{\gamma}^{\dot{\delta}}$

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- Unknown what \tilde{h}_{-6} deforms \Rightarrow only get ASD spacetime

BGK amplitudes in twistor space

$$\int_M \Sigma_0 \wedge \gamma \wedge \gamma = \int_{M \times \mathbb{CP}^1 \times \mathbb{CP}^1} d^4x [\pi_n d\pi_n] [\pi_1 d\pi_1] I(B_n, \cdot) \lrcorner \frac{1}{\bar{\partial} + \mathcal{L}_V} B_1 [\pi_n \pi_1]^4$$

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- ASD background encoded in $\bar{\partial} + \mathcal{L}_V$ ‘propagator’
 - No canonical way to pullback vector field; make gauge choice

$$V = V^\alpha \frac{\partial}{\partial \omega^\alpha} \rightarrow \frac{V^\alpha \xi^{\dot{\alpha}}}{[\pi \xi]} \frac{\partial}{\partial y^{\alpha \dot{\alpha}}} \text{ with } |\xi| = |\eta|$$

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- $I(\cdot, \cdot) := \epsilon^{\alpha\beta} \frac{\partial}{\partial \omega^\alpha} \wedge \frac{\partial}{\partial \omega^\beta}$

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Expanding in powers of V gives

$$\sum_{n=3}^{\infty} \int d^4x [\pi_n d\pi_n] I(B_n, \cdot) \lrcorner \left(\frac{1}{\bar{\partial}} \mathcal{L}_{V_{n-1}} \frac{1}{\bar{\partial}} \mathcal{L}_{V_{n-2}} \cdots \frac{1}{\bar{\partial}} \mathcal{L}_{V_3} \frac{1}{\bar{\partial}} \mathcal{L}_{V_2} \frac{1}{\bar{\partial}} B_1 [\pi_n \pi_1]^4 \right)$$

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- Derivative of δ -fn support \rightarrow perturbative description of support on deformed twistor lines

Odds and Ends



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The spacetime Plebanski action $\int_M \Sigma \wedge d\Gamma + \int_M \Sigma \wedge \Gamma \wedge \Gamma$ suggests we consider the **twistor action**

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- Can we find a 'connected prescription' for Einstein gravity?