



A Selberg zeta function for warped AdS3 black holes

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Selberg zeta: What and why

Let's start with the Riemann zeta function for comparison:

Selberg zeta functions are defined for hyperbolic quotient spacetimes.

Example: The BTZ black hole Perry, Williams 2003

IMPORTANT PART: The Selberg zeta function for BTZ encodes quantum corrections and quasinormal modes. Keeler, Martin, Svesko 2018

$$\log Z_{\Gamma}(\Delta) \propto \log \det \nabla_{reg}^2$$
$$s^* \leftrightarrow \omega_{QN}$$

Krasnov 2000

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s} = \prod_{p \text{ prime}} \frac{1}{1 - p^{-s}}$$

$$Z_{\Gamma}(s) = \prod_{p} \prod_{n=0}^{\infty} \left(1 - N(p)^{-s-n} \right)$$

$$\mathbb{H}^n/\Gamma \qquad \partial_\phi$$

$$(s) = \prod_{k_1, k_2=0}^{\infty} \left[1 - e^{2ibk_1} e^{-2ibk_2} e^{-2a(k_1 + k_2 + s)} \right]$$

 Z_{Γ}

Warped spacetimes: What and why

Spacelike WAdS3 is constructed by stretching or squashing the length of a fibre over AdS2. The symmetry group is SL(2,R)xU(1).

$$ds^{2} = \frac{L^{2}}{\nu^{2} + 3} \left(-\cosh^{2} \sigma d\tau^{2} + d\sigma^{2} + \frac{4\nu^{2}}{\nu^{2} + 3} (du + \sinh \sigma d\tau)^{2} \right)$$

Just like the BTZ black hole is a discrete quotient of AdS3, WAdS3 black holes are discrete quotients of WAdS3.

Anninos, Li, Padi, Song, Strominger 2008

$$\partial_{\phi} = \frac{\nu^2 + 3}{8} \left[\left(r_+ - r_- - \frac{\sqrt{(\nu^2 + 3)r_+ r_-}}{\nu} \right) J_2 - (r_+ - r_-) \tilde{J}_2 \right]$$



Why interesting?: WAdS quotients appear in many important contexts, such Kerr/CFT.



How to make Z_{Γ} from quotient

Perry, Williams 2003

E

xample: Euclidean BTZ
$$ds^2 = N(r)^2 d\tau^2 + \frac{dr^2}{N(r)^2} + r^2 \left(N^{\phi}(r)d\tau + d\phi\right)^2$$

To a point $(\tau, r, \phi) \in \mathbb{R}^3$ one associates the point $(x, y, z) \in \mathbb{H}^3$ given by

$$\begin{aligned} x &= A(r)\cos f(\phi,\tau)\exp(r_{+}\phi - |r_{-}|\tau) \\ y &= A(r)\sin f(\phi,\tau)\exp(r_{+}\phi - |r_{-}|\tau) \\ z &= B(r)\exp(r_{+}\phi - |r_{-}|\tau) \end{aligned} \qquad ds^{2} = \frac{L^{2}}{z^{2}}\left(dx^{2} + dy^{2} + dz^{2}\right) \end{aligned}$$

When $\phi \to \phi + 2\pi n$ we have $(x,y,z) \to (x',y',z')$, where

$$\begin{aligned} x' &= e^{2\pi n r_{+}} \left[x \cos \left(2\pi n |r_{-}| \right) - y \sin \left(2\pi n |r_{-}| \right) \right] \\ y' &= e^{2\pi n r_{+}} \left[y \cos \left(2\pi n |r_{-}| \right) + x \sin \left(2\pi n |r_{-}| \right) \right] \\ z' &= e^{2\pi n r_{+}} z \end{aligned}$$



The zeta function parameters are the coefficients of the dilation and rotation.

$$a = \pi r_{+}$$

$$b = \pi |r_{-}|$$

$$Z_{\Gamma}(s) = \prod_{k_{1},k_{2}=0}^{\infty} \left[1 - e^{2ibk_{1}}e^{-2ibk_{2}}e^{-2a(k_{1}+k_{2}+s)}\right]$$

Warped Poincare Patch

Build conformal coordinates $w^+ = x + iy$ $w^- = x - iy$

$$w^{+} = \sqrt{\frac{r^{2} - r_{+}^{2}}{r^{2} - r_{-}^{2}}} e^{\alpha \phi + \beta t} \qquad w^{-} = \sqrt{\frac{r^{2} - r_{+}^{2}}{r^{2} - r_{-}^{2}}} e^{\gamma \phi + \delta t} \qquad z = \sqrt{\frac{r_{+}^{2} - r_{-}^{2}}{r^{2} - r_{-}^{2}}} e^{1/2((\alpha + \gamma)\phi + (\beta + \delta)t)}$$

From these coordinates we can make generators of SL(2,R)xSL(2,R):

$$H_{1} = i\partial_{+}, \qquad H_{0} = i\left(w^{+}\partial_{+} + \frac{1}{2}z\partial_{z}\right), \qquad H_{-1} = i((w^{+})^{2}\partial_{+} + w^{+}z\partial_{z} - z^{2}\partial_{-})$$
$$\bar{H}_{1} = i\partial_{-}, \qquad \bar{H}_{0} = i\left(w^{-}\partial_{-} + \frac{1}{2}z\partial_{z}\right), \qquad \bar{H}_{-1} = i((w^{-})^{2}\partial_{-} + w^{-}z\partial_{z} - z^{2}\partial_{+})$$

To find $(\alpha, \beta, \gamma, \delta)$, we require that the SL(2,R)xU(1) Casimir is proportional to the scalar Laplacian:

$$(\mathcal{H}^2 + \lambda H_0^2) \Phi \propto \nabla^2 \Phi$$

 $\mathcal{H}^2 = -\bar{H}_0^2 + \frac{1}{2} (\bar{H}_1 \bar{H}_{-1} + \bar{H}_{-1} \bar{H}_2)$

Result 1:
$$ds^2 = \frac{4}{(3+\nu)^2 z^2} \left((3+\nu^2) dw_+ dw_- + 4\nu^2 dz^2 + \frac{3(\nu^2-1)w_+}{z^2} (dw_-^2 + 2z dw_- dz) \right)$$

Selberg zeta function for WAdS3

The rotation and dilation generators are:

The quotient is generated by the group element:

In the dilation/rotation basis:

$$H_0 - \bar{H}_0 = i(w_+\partial_+ - w_-\partial_-)$$
$$H_0 + \bar{H}_0 = i(w_+\partial_+ + w_-\partial_- + z\partial_z)$$

$$e^{-2\pi i(\gamma H_0 + \alpha \bar{H}_0)} = e^{2\pi \partial_\phi}$$

$$e^{-2\pi i \left(\frac{\gamma+\alpha}{2} (H_0 + \bar{H}_0) + \frac{\gamma-\alpha}{2} (H_0 - \bar{H}_0) \right)}$$

Result 2:
$$Z_{\Gamma}(s) = \prod_{k_1,k_2=0}^{\infty} \left[1 - e^{2ibk_1} e^{-2ibk_2} e^{-2a(k_1+k_2+s)} \right]$$
 $2a = \gamma + \alpha$
 $2b = \gamma - \alpha$

Result 3: We establish a map between Selberg zeros and QNMs.

$$s^* \leftrightarrow \Delta$$

$$\omega_{QN} \leftrightarrow \omega_n$$

Keeler, Martin, Svesko 2018

Result 4: We calculate the thermal frequencies for WAdS3.

Castro, Keeler, Szepietowski 2017

Heat Kernel for (Euclidean) WAdS3

We would like to see whether a similar relationship holds in the warped case.

 $\log Z_{\Gamma}(\Delta) \propto \log \det \nabla_{reg}^2$

To accomplish this, $\nabla^2 \psi_n = \lambda_n \psi_n$ $(\partial_t + \nabla_x^2) K(t; x, y) = 0$ $K(0; x, y) = \delta(x, y)$ we set about computing the heat kernel for WAdS3. $S^{(1)} = -\frac{1}{2} \log \det \nabla^2 = -\frac{1}{2} \sum_n \log \lambda_n = \frac{1}{2} \int_{0^+}^{\infty} \frac{dt}{t} \int d^3x \sqrt{g} K(t; x, x)$

For warped black holes, we employ the method of images. Giombi, Maloney, Yin 2008

$$K^{\mathbb{H}^3/\Gamma}(t;x,y) = \sum_{\gamma \in \Gamma} K^{\mathbb{H}^3}(t;x,\gamma y)$$

We can use our warped upper-half space metric to do this calculation:

$$ds^{2} = \frac{4}{(3+\nu)^{2}z^{2}} \left((3+\nu^{2})dw_{+}dw_{-} + 4\nu^{2}dz^{2} + \frac{3(\nu^{2}-1)w_{+}}{z^{2}}(dw_{-}^{2}+2zdw_{-}dz) \right)$$

Future directions

Quotient spacetimes are everywhere! Lots of future directions.

Quantum corrections in the context of Kerr/CFT (extremal and nonextremal versions).

Study quantum corrections of more generic quotients:

1. Wormholes

2. Warped de Sitter black holes

3. Holographic entanglement



Selberg trace formula as a connection between two differerent methods of calculating functional determinants: The heat kernel method and the quasinormal mode method.



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