

# Branched $SL(2, \mathbb{Z})$ Duality

Ziqi Yan

Nordita

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2208.13815 w/ Eric Bergshoeff, Kevin Grosvenor, Johannes Lahnsteiner, Utku Zorba

2112.09316 w/ Stephen Ebert, Hao-Yu Sun

## 4D Yang-Mills

$$S_{\text{YM}} = -\frac{1}{4g_{\text{YM}}^2} \int d^4x F_{\mu\nu} F^{\mu\nu} \quad F = dA$$

S-duality:  $S_{\text{parent}} = S_{\text{YM}} + \int \tilde{F} \wedge (F - dA)$

e.o.m. from varying  $A_\mu$ :  $\tilde{F} = d\tilde{A}$

$$\tilde{S}_{\text{YM}} = -\frac{g_{\text{YM}}^2}{4} \int d^4x \tilde{F}_\mu \tilde{F}^{\mu}$$

Strong-weak duality:  $g_{\text{YM}} \rightarrow \frac{1}{g_{\text{YM}}}$

# D3-Brane

[Tseytlin '96]

$$S_{D3} = - \int d^4x \sqrt{-\det(\eta_{\mu\nu} + e^{-\frac{\Phi}{2}} F_{\mu\nu}^B)} \quad g_{\mu\nu} \sim e^{\frac{\Phi}{2}}$$

$$F^B = dA^B$$

S-dual  $\left\{ \begin{array}{l} \downarrow \end{array} \right.$

$$\tilde{S}_{D3} = - \int d^4x \sqrt{-\det(\eta_{\mu\nu} + e^{\frac{\Phi}{2}} F_{\mu\nu}^C)} \quad \tilde{g}_{\mu\nu} \sim e^{-\frac{\Phi}{2}}$$

$$F^C = dA^C$$

•  $SL(2, \mathbb{Z})$  duality  $\Lambda = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$   $a, b, c, d \in \mathbb{Z}$   $ad - bc = 1$

$$S_{D3} = -T_3 \int d^4x \sqrt{-\det \left( \eta_{\mu\nu} + \frac{\Theta^\top \mathcal{F}_{\mu\nu}}{\sqrt{\Theta^\top \mathcal{M}^{-1} \Theta}} \right)}$$

$$\Theta = \begin{pmatrix} p \\ q \end{pmatrix} \leftarrow (p, q)\text{-string}$$

$$\Theta \rightarrow \Lambda \Theta$$

$$\mathcal{F} = \begin{pmatrix} B + dA^B \\ C^{(2)} + dA^C \end{pmatrix}$$

$$\mathcal{F} \rightarrow (\Lambda^{-1})^\top \mathcal{F}$$

$$\mathcal{M} = e^{\Phi} \begin{pmatrix} |\tau|^2 & c^{(0)} \\ c^{(0)} & 1 \end{pmatrix}$$

$$\mathcal{M} \rightarrow \Lambda \mathcal{M} \Lambda^\top$$

$$\tau = c^{(0)} + i e^{-\Phi}$$

$$\tau \rightarrow \frac{a\tau + b}{c\tau + d}$$



Open Strings in B-Field

• magnetic

$$B_{23} \sim \frac{\alpha'}{\theta} \longrightarrow \text{noncommutativity in } x^2 \text{ and } x^3$$

$$[x^i, x^j] \sim \alpha' \frac{1}{G+B} B \frac{1}{G-B} \quad \text{Seiberg-Witten map}$$

particle limit:  $\alpha' \rightarrow 0$ 

Non-Commutative Yang-Mills

$$[x^2, x^3] \sim \theta \quad g_{\text{NeyM}}^2 \sim \frac{g_s}{\theta}$$

## S-dual: Non-Commutative Open String

• electric

$$\tilde{B}_{01} \sim 1 - \frac{\alpha'}{\alpha'_{\text{eff}}} e \longrightarrow \text{space/time noncommutativity}$$

$\alpha' \rightarrow 0$ : still a string theory!

$$[x^0, x^1] \sim \frac{\alpha'_{\text{eff}}}{e} \quad g^2_{\text{NCOS}} \sim g_s \sqrt{\frac{e}{\alpha'_{\text{eff}}}}$$

$SL(2, \mathbb{Z})$  generalization?

# Generalized NCOS Limit $\omega \rightarrow \infty$

$$\hat{G}_{MN} = \omega^2 \overbrace{\tau_M^A \tau_N^B \eta_{AB}}^{\tau_{MN}} + \overbrace{E_M^{A'} E_N^{A'}}^{E_{MN}}$$

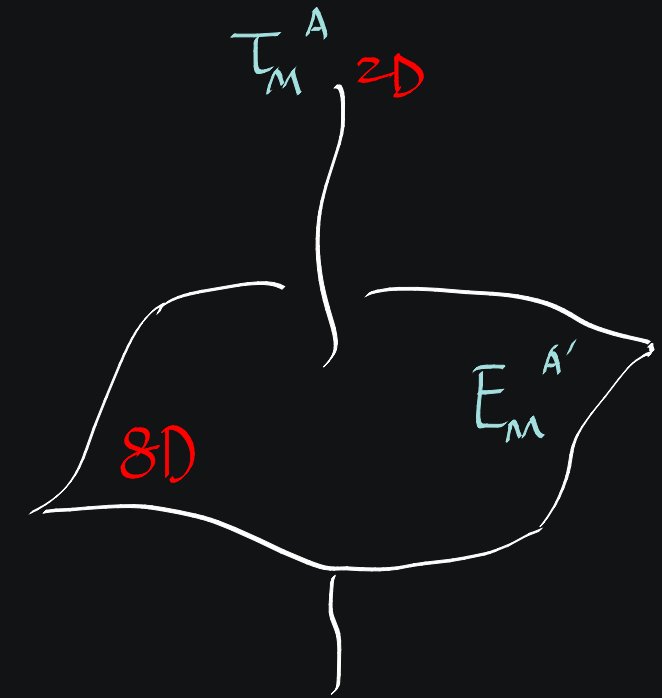
$$\hat{B} = -\omega^2 \ell + B$$

$$\hat{\Phi} = \bar{\Phi} + \ln \omega$$

$$\hat{C}^{(q)} = \omega^2 \ell \wedge C^{(q-2)} + C^{(q)} \quad [\text{Ebert, Sun, ZY '21}]$$

longitudinal  $A = 0, 1$

transverse  $A' = 2, \dots, 9$



$$M, N = 0, 1, \dots, 9$$

→ nonrelativistic string theory [Review by Oling, ZY '22]

[Klebanov, Maldacena; Gomis, Ooguri; Danielsson, Guijosa, Kruczenski '00]

# Branched $SL(2, \mathbb{Z})$

- $u \rightarrow \infty$  limit of D3-brane
- D3-brane in nonrelativistic string theory

$$S_{D3} = -T_3 \int d^4x \sqrt{-\det \begin{pmatrix} 0 & \tau_\beta^E \\ \bar{\tau}_\alpha^E & E_{\alpha\beta}^E + \frac{\Theta^\top F_{\alpha\beta}}{\Theta^\top \zeta} \end{pmatrix}}$$

$$\tau_\alpha = \tau_\alpha^0 + \tau_\alpha^1$$

$$\bar{\tau}_\alpha = \tau_\alpha^0 - \tau_\alpha^1$$

$$\Theta = \begin{pmatrix} p \\ q \end{pmatrix}$$

$$F = \Sigma + \begin{pmatrix} dA^B \\ dA^C \end{pmatrix}$$

$$\zeta = e^{\frac{\Phi}{2}} \begin{pmatrix} 1 \\ -c^{(0)} \end{pmatrix}$$

$$\Sigma = \begin{pmatrix} B \\ c^{(2)} \end{pmatrix} + \frac{q}{p - qc^{(0)}} \begin{pmatrix} 0 \\ l e^{-2\Phi/2} \end{pmatrix}$$

$$l = \frac{1}{2} \epsilon_{AB} \tau^A \wedge \tau^B$$

•  $\omega \rightarrow \infty$  limit of  $SL(2, \mathbb{Z})$  transformations

finite  $SL(2, \mathbb{Z})$  transformations only if  $c\tau_0 + d > 0$  !

$$\mathbb{Q} \rightarrow \Lambda \mathbb{Q}$$

$$\mathcal{L} \rightarrow (\Lambda^{-1})^T \mathcal{L}$$

$$\Lambda = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

quasi-doublet

$$\Sigma \rightarrow (\Lambda^{-1})^T \left\{ \Sigma + \alpha \begin{pmatrix} q \\ -p \end{pmatrix} \right\} \quad \mathbb{Q}^T \Sigma \text{ is invariant}$$

$\downarrow$   
orthogonal to  $\mathbb{Q}$

- branched  $SL(2, \mathbb{Z})$  transformations

$$\mathcal{G} \rightarrow \wedge \mathcal{G}$$

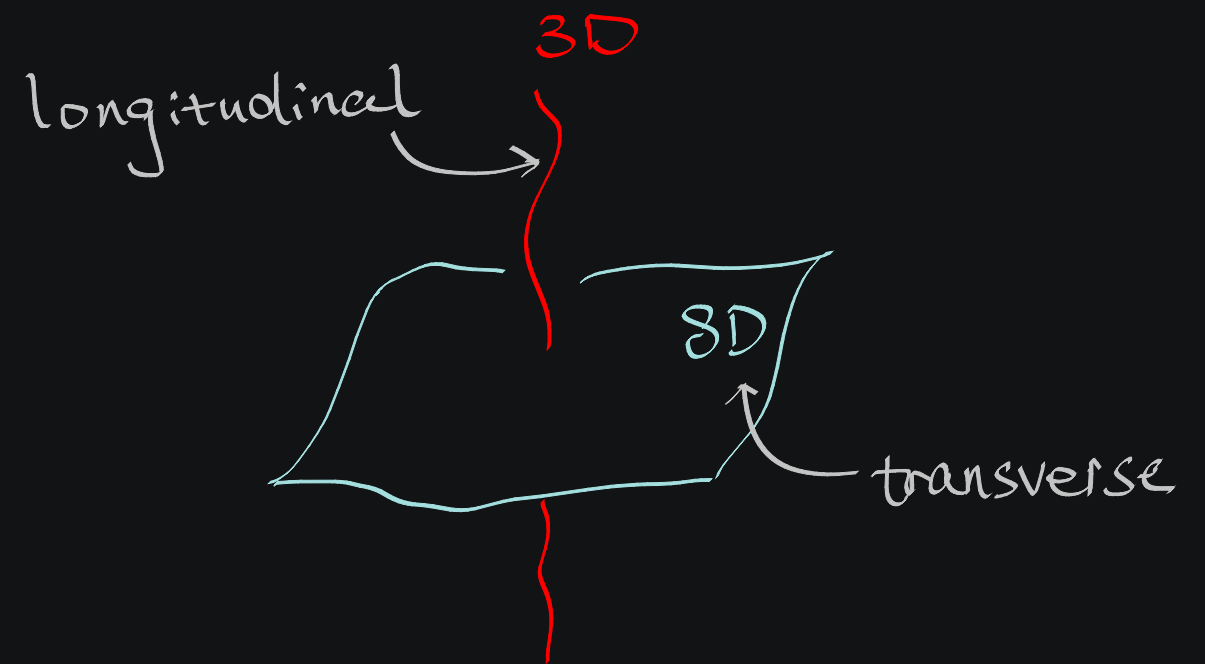
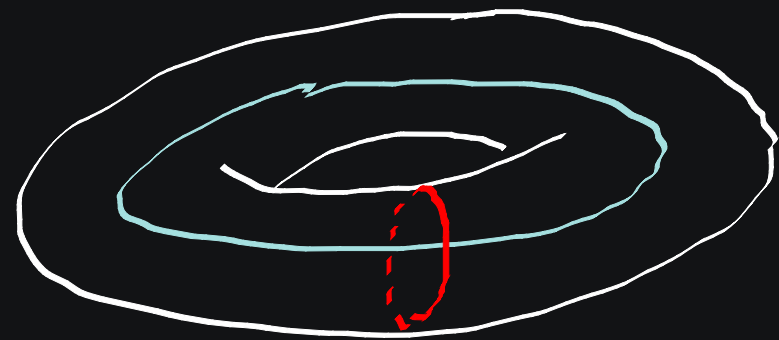
$$\mathcal{G} \rightarrow \text{sgn}(c\mathcal{G}_0 + d) (\Lambda^{-1})^T \mathcal{G}$$

$$\mathcal{F} \rightarrow \text{sgn}(c\mathcal{G}_0 + d) (\Lambda^{-1})^T \left\{ \mathcal{F} + \alpha \begin{pmatrix} q \\ -p \end{pmatrix} \right\}$$

[Berman '98]  
[Ebert, ZY, in progress]

# Geometrical Interpretation

nonrelativistic M5-brane on a split torus



$\Rightarrow$  branched  $SL(2, \mathbb{Z})$  of D3-brane

Branched (p, q) - Strings

- $\omega \rightarrow \infty$  limit of (p, q) - strings

finite action only if  $p - q c^{(0)} > 0$  !

- (p, q) - strings in nonrelativistic string theory

$$S_+ = - \int d^2x (\Theta^T \mathcal{L}) \sqrt{-\tau^E} \tau_E^{\mu\nu} E_{\mu\nu}^E - \int \Theta^T \Sigma$$

$$\uparrow$$

$$p - q c^{(0)} > 0$$

$$p - q c^{(0)} \rightarrow \frac{p - q c^{(0)}}{c c^{(0)} + d} \rightsquigarrow \text{sign change if } c c_0 + d < 0$$



$$p-q C^{(0)} > 0$$

$$S_+ = - \left\{ \int d^2x (\Theta^T \mathcal{L}) \sqrt{-\tau^E} \tau_E^{\mu\nu} E_{\mu\nu}^E + \int \Theta^T \Sigma \right\} \Leftrightarrow c C_0 + d > 0$$


$$c C_0 + d < 0$$

$$S_- = + \left\{ \int d^2x (\Theta^T \mathcal{L}) \sqrt{-\tau^E} \tau_E^{\mu\nu} E_{\mu\nu}^E + \int \Theta^T \Sigma \right\} \Leftrightarrow c C_0 + d > 0$$

$$p-q C^{(0)} < 0$$

To Be Continued:

Unbranched in nonrelativistic IIB supergravity?

[Bergshoeff, Grosvenor, Lahnsteiner, ZY, Zorba, to appear]

D-brane solutions?

Nonrelativistic holography? [Harmark, Hartong, Obers '17]

DLCQ supergravity and Matrix theory?

[Bergshoeff, Lahnsteiner, Romano, Rosseel '22]

Thank You!