

# Comparison of post-Minkowskian and self-force expansions: Scattering in a scalar charge toy model

arXiv:2304.09200



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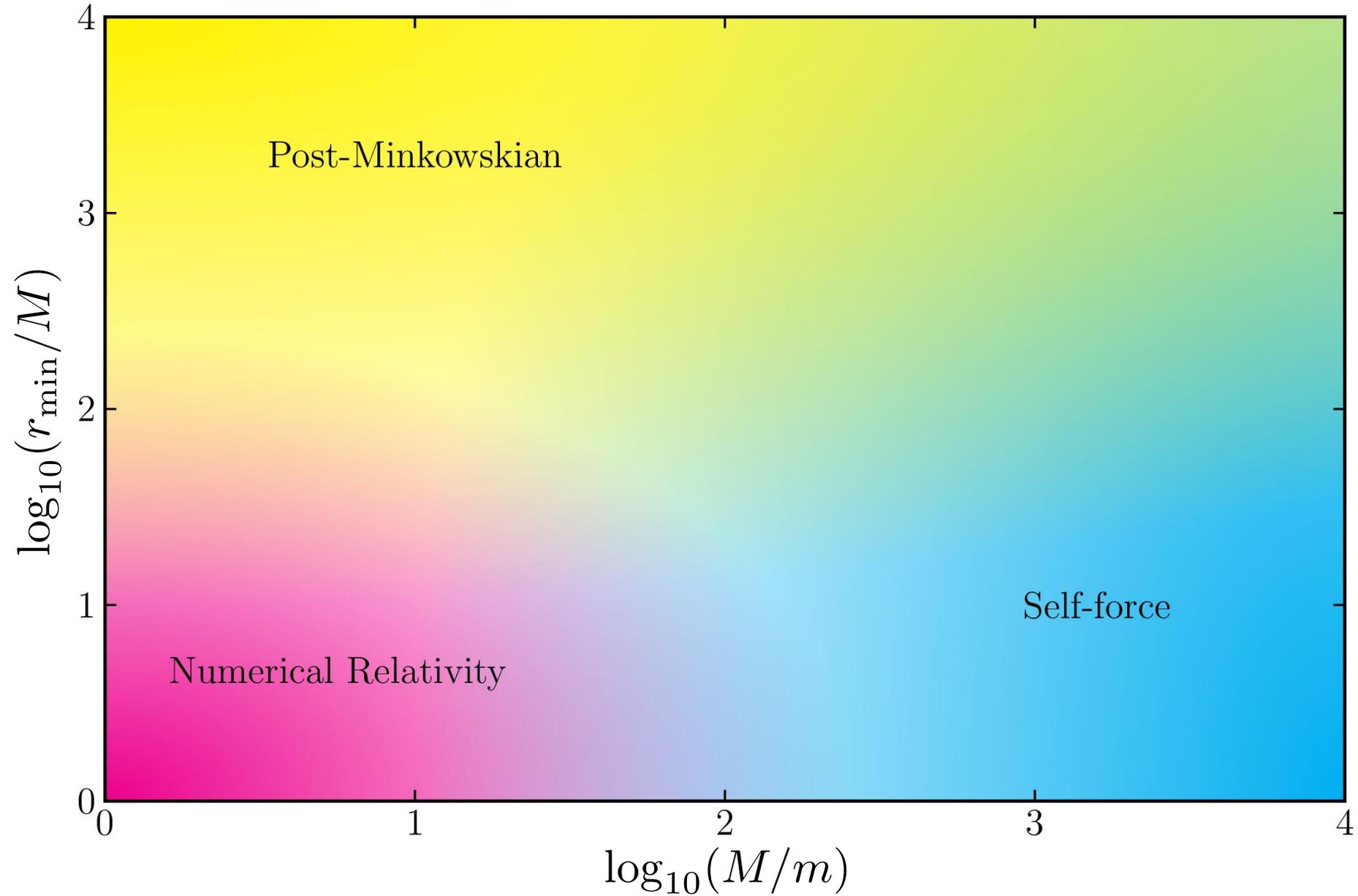
with L. Barack, Z. Bern, E. Herrmann, J. Parra-Martinez, R. Roiban,  
M. S. Ruf, C.-H. Shen, M. P. Solon, F. Teng, and M. Zeng.

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# Two-body parameter space for scattering



# Self-force correction to the scattering angle

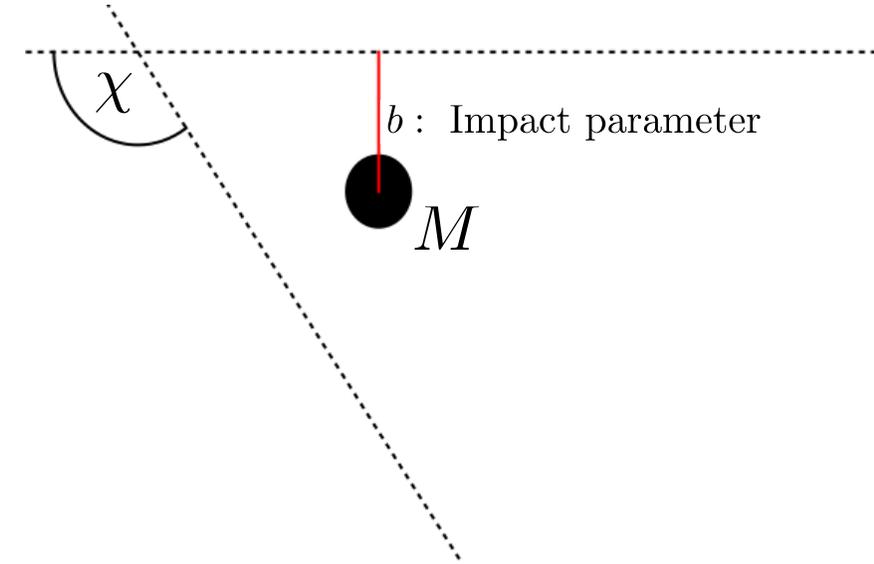


Scattering angle [Barack & OL '22]:

$$\chi = \chi^{(0)} + \eta \delta \chi$$

$$\delta \chi = \sum_{\pm} \int_{R_{\min}}^{\infty} [\mathcal{G}_E^{\pm}(r) F_t^{\pm} - \mathcal{G}_L^{\pm}(r) F_{\varphi}^{\pm}] dr$$

Functions of geodesics



Can split into **conservative** and **dissipative** pieces on outgoing leg:

$$F_{\alpha}^{\text{cons}}(r, \dot{r}) = -F_{\alpha}^{\text{cons}}(r, -\dot{r})$$

$$F_{\alpha}^{\text{diss}}(r, \dot{r}) = F_{\alpha}^{\text{diss}}(r, -\dot{r})$$

$$\alpha = t, \varphi$$

$$\delta \varphi_{\text{cons}}^{(1)} = \int_{R_{\min}}^{\infty} [\mathcal{G}_E^{\text{cons}} F_t^{\text{cons}} - \mathcal{G}_L^{\text{cons}} F_{\varphi}^{\text{cons}}] dr$$

$$\delta \varphi_{\text{diss}}^{(1)} = \int_{R_{\min}}^{\infty} [\mathcal{G}_E^{\text{diss}} F_t^{\text{diss}} - \mathcal{G}_L^{\text{diss}} F_{\varphi}^{\text{diss}}] dr$$

# 1+1D scalar field evolution scheme



Scalar field obeys the **Klein-Gordon** equation.

Decompose in the **time domain**:

$$\Phi = \frac{2\pi q}{r} \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \psi_{\ell m}(t, r) Y_{\ell m}(\theta, \varphi)$$

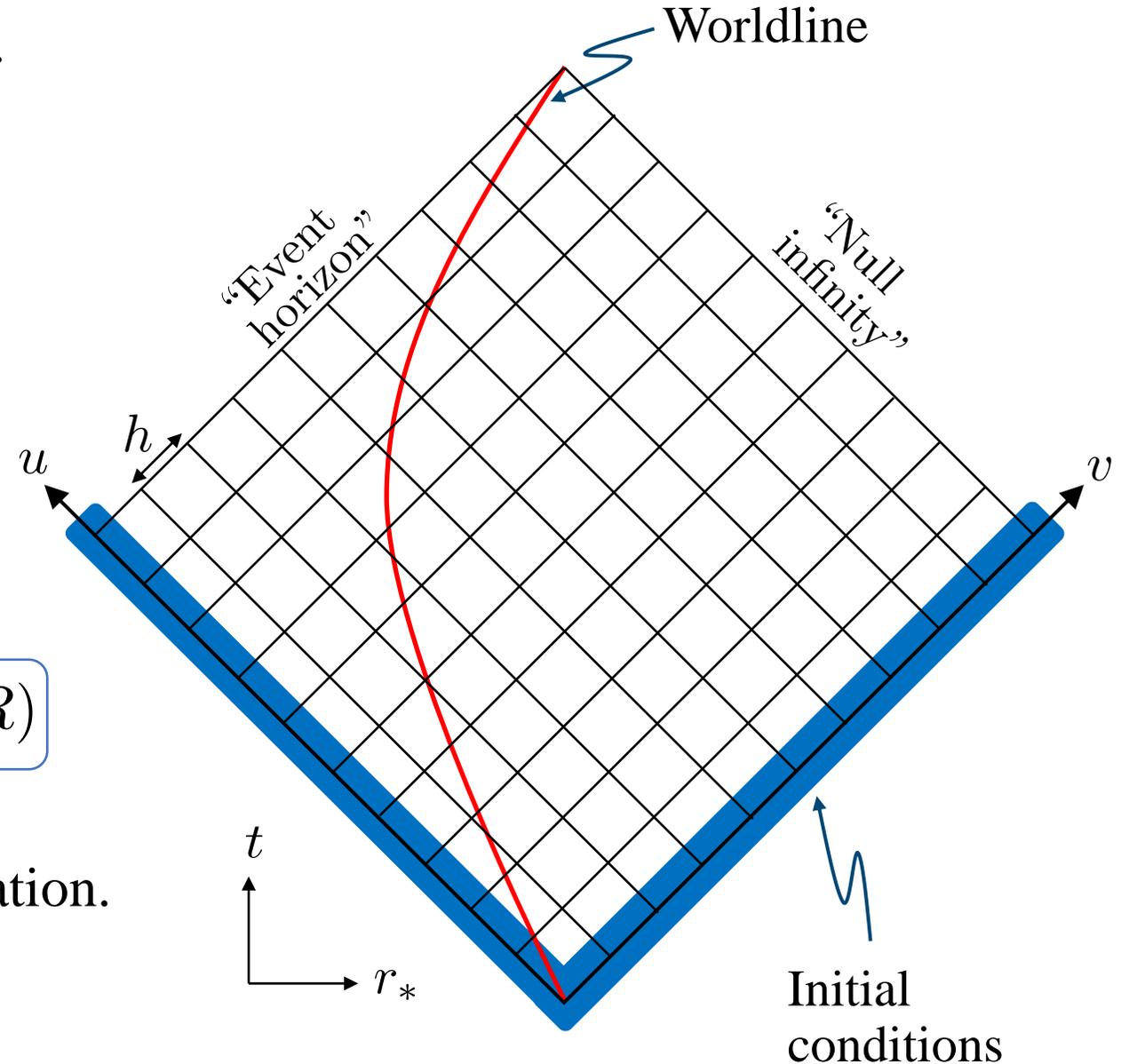
$q$  : Scalar charge

1+1D scalar wave equation:

$$\psi_{,uv} + V(\ell; r)\psi = S_{\psi}(\ell; x_p^{\mu}) \delta(r - R)$$

Evolve **finite-difference** version of 1+1D equation.

Extract SF via mode-sum regularisation.



# Post-processing: Truncation at finite radius



Can only numerically determine the self-force up to a **finite radius**  $R = R_{\text{final}}$  :

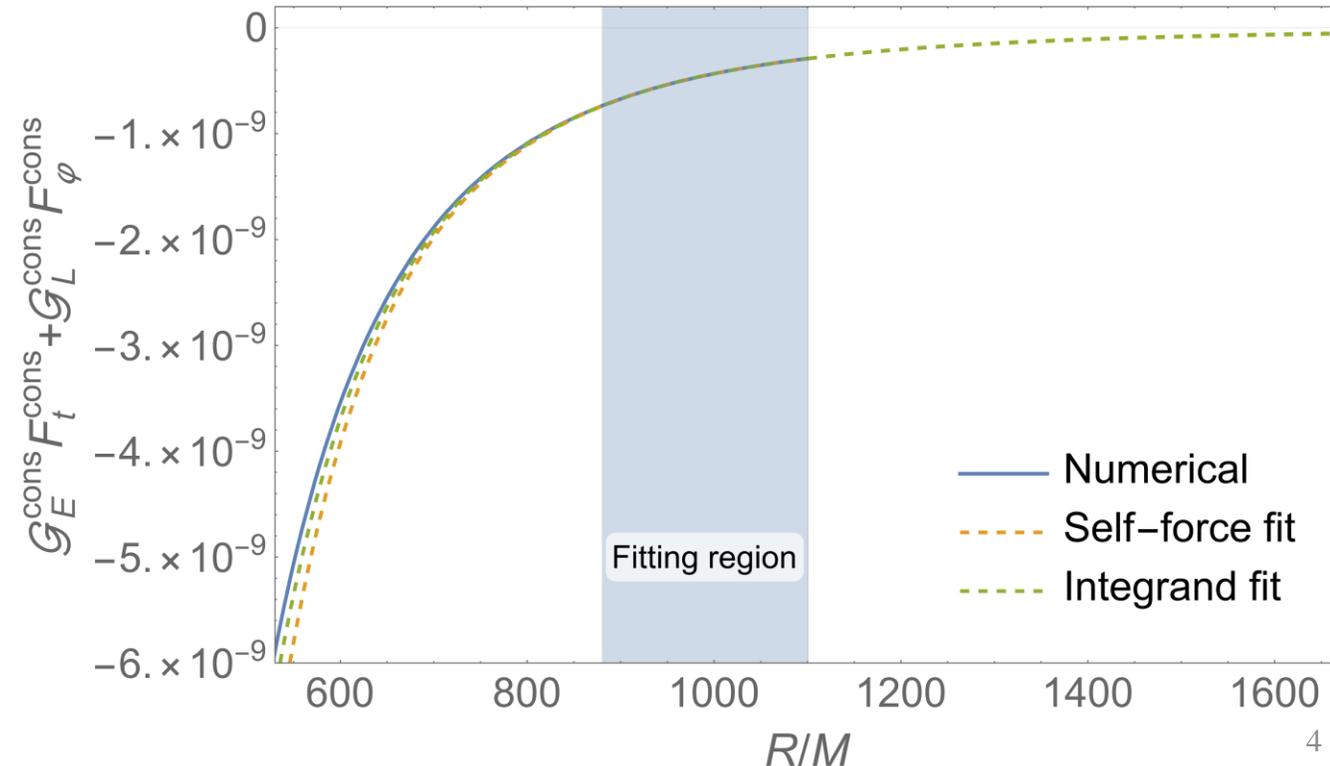
$$\delta\varphi^{(1)} = \int_{R_{\text{min}}}^{R_{\text{final}}} [\mathcal{G}_E F_t - \mathcal{G}_L F_\varphi] dr + \int_{R_{\text{final}}}^{\infty} [\mathcal{G}_E F_t - \mathcal{G}_L F_\varphi] dr$$

Numerical  $\rightsquigarrow$  Error ( $\sim 1\%$ )  $\rightsquigarrow$

Form an analytic tail by fitting to the data:

- Fit the **self-force** data.
- Fit the **integrand** directly.

Tail contributes an error  $\sim 0.01\%$ .



# Post-processing: Richardson extrapolation



Next dominant error due to **finite resolution**  $\sim 0.1\%$ .

Can increase the convergence from quadratic to **cubic** using Richardson extrapolation.

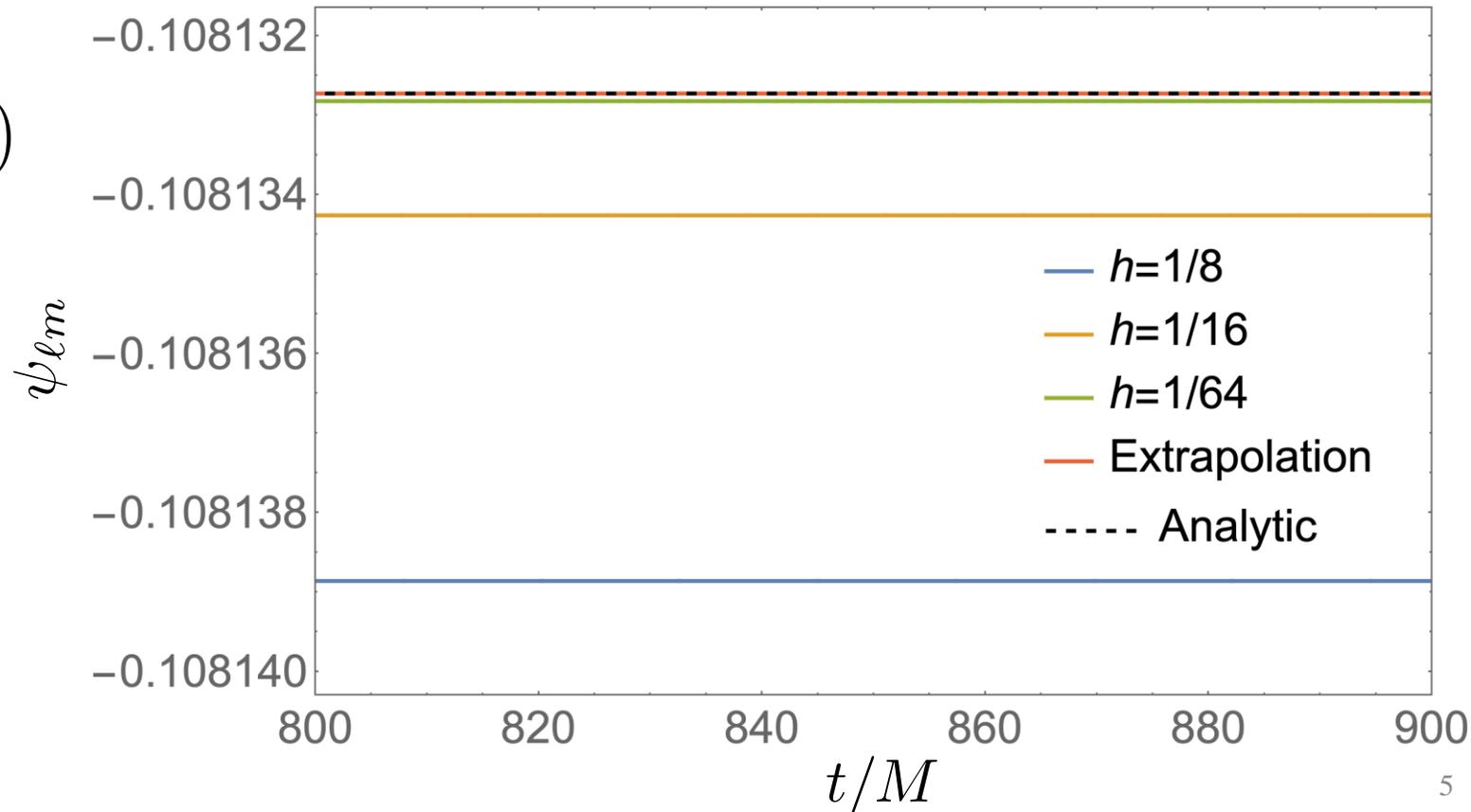
Model:

$$A(h) = A_{\text{exact}} + Ch^n + O(h^{n+1})$$

Extrapolation:

$$\begin{aligned} A_{\text{Extr}} &= \frac{t^n A\left(\frac{h}{t}\right) - A(h)}{t^n - 1} \\ &= A_{\text{exact}} + O(h^{n+1}) \end{aligned}$$

Error in extrapolation  $< 0.001\%$ .



# Scattering angle correction: PM expansion



Expansion around flat space:

$$\delta\chi^{\text{PM}} = \sum_{i=0}^{\infty} \delta\chi_i \left(\frac{M}{b}\right)^i$$

2PM [Gralla & Lobo '22]:

$$\delta\chi_2^{\text{cons}} = -\frac{\pi}{4} \left(\frac{M}{b}\right)^2$$

$$\delta\chi_2^{\text{diss}} = 0$$

$v$  : Velocity at infinity

$b$  : Impact parameter

3PM:

$$\delta\chi_3^{\text{cons}} = -\frac{4(3-v^2)}{3v^2\sqrt{1-v^2}} \left(\frac{M}{b}\right)^3$$

$$\delta\chi_3^{\text{diss}} = \frac{2(v^2+1)^2}{3v^3\sqrt{1-v^2}} \left(\frac{M}{b}\right)^3$$

LO

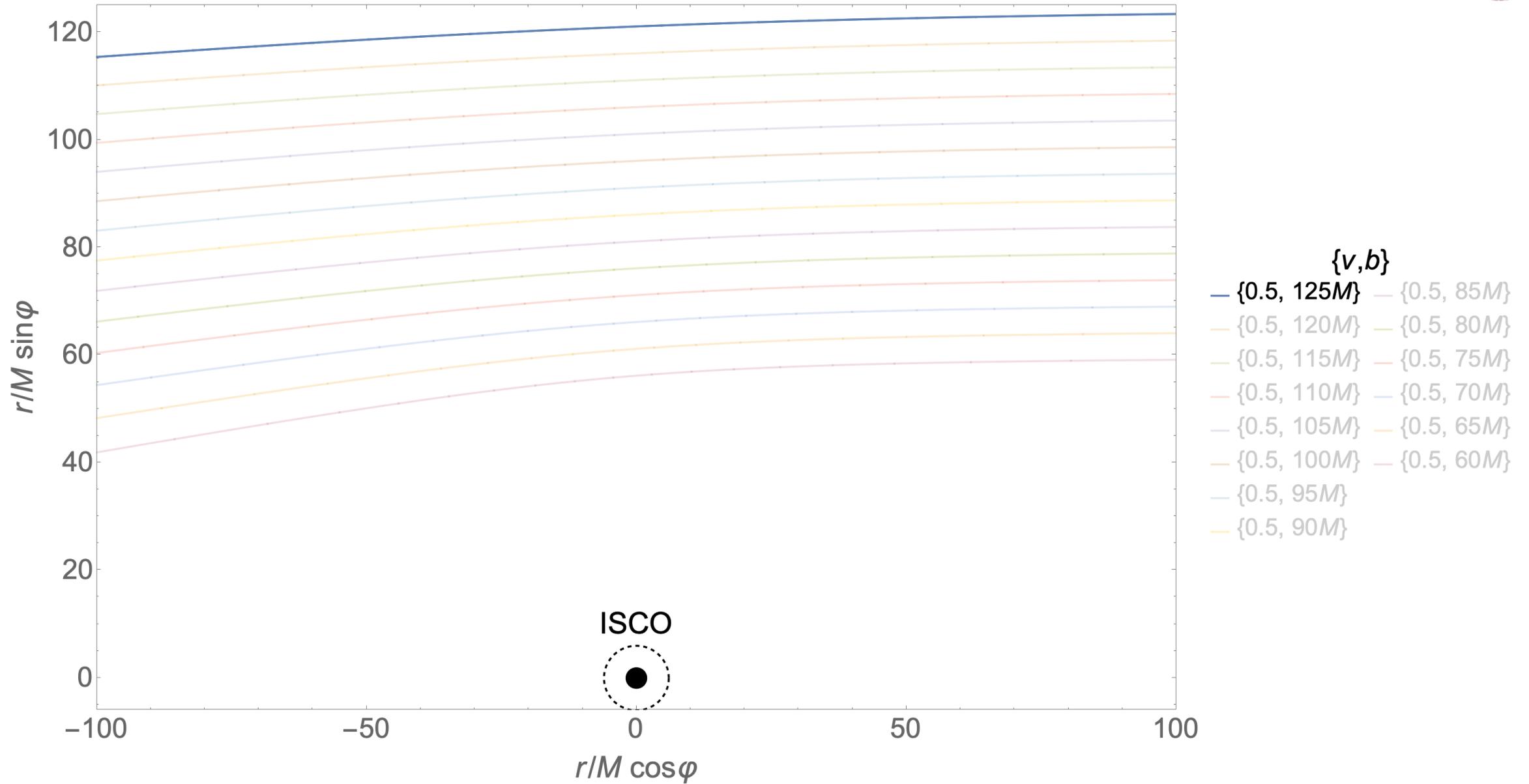
NLO

4PM dissipative:

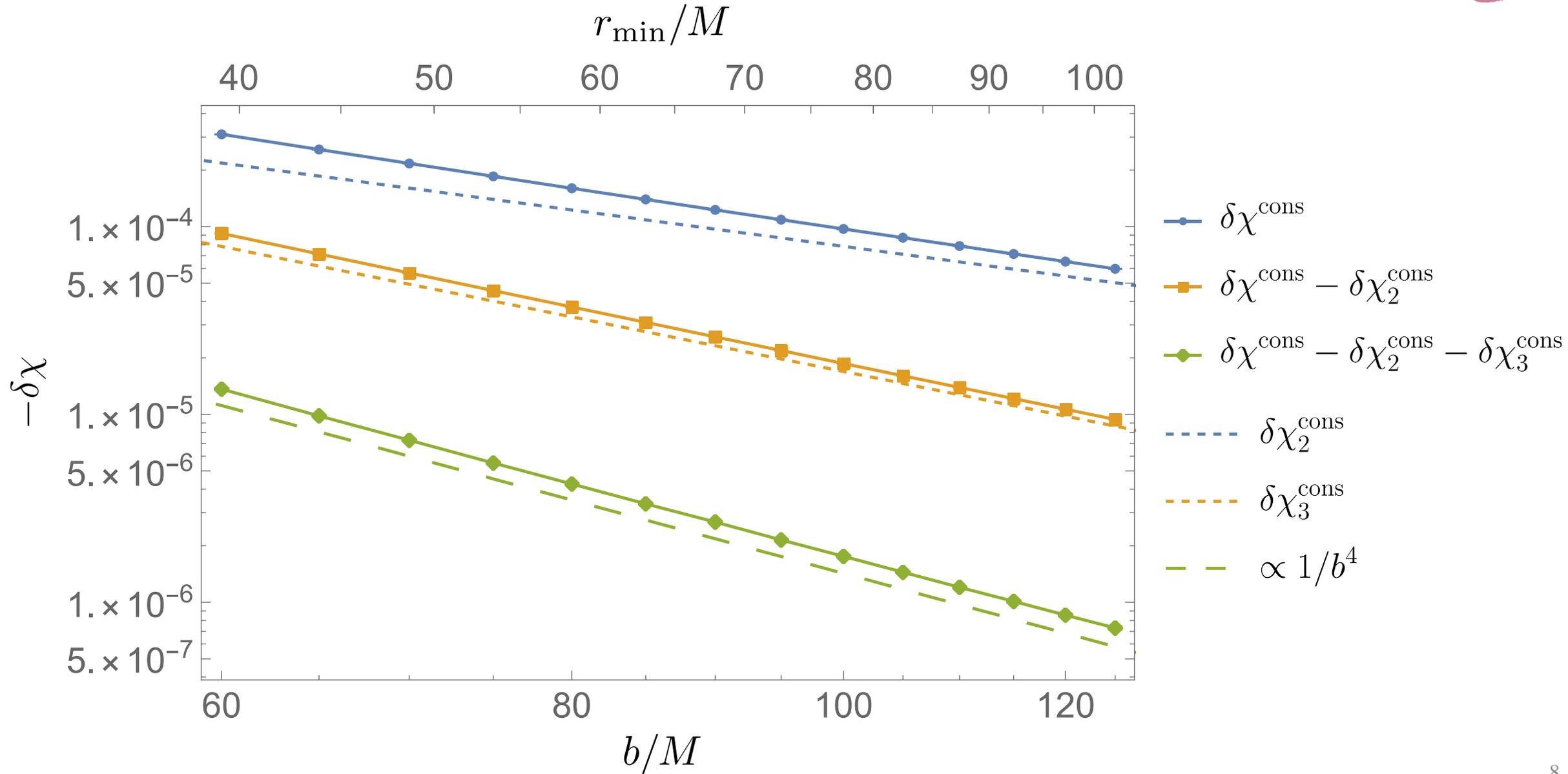
$$\delta\chi_4^{\text{diss}} = \left( r_1 + r_2 \operatorname{arcsech} \left( \sqrt{1-v^2} \right) + r_3 \log \left[ \frac{1}{2} \left( \frac{1}{\sqrt{1-v^2}} + 1 \right) \right] \right) \left(\frac{M}{b}\right)^4$$

$r_i = \text{rational coefficients}$

# Sample orbits



# Conservative: $v = 0.5$



# Extraction of high-order conservative PM results



PM expansion with free parameters:

$$\delta\varphi_{\text{cons}}^{\text{PM}} = \frac{a_2}{b^2} + \frac{a_3}{b^3} + \frac{a_4}{b^4} + \frac{a_5}{b^5} + \dots$$

Up to **3PM** can fit value or use **analytic value**.

$a_2$	$a_3$	$a_4$	$a_5$
-1.0886	—	—	—
-0.7535	-21.77	—	—
-0.7899	-16.17	-206.5	—
-0.7803	-18.49	-25.0	-4620
<b>-0.785398</b>	-19.18	—	—
<b>-0.785398</b>	-16.93	-176.2	—
<b>-0.785398</b>	-17.20	-131.1	-1793
<b>-0.785398</b>	<b>-16.9356</b>	-175.9	—
<b>-0.785398</b>	<b>-16.9356</b>	-174.4	-107
< 1%	~ 1%	~ -175	< 0(?)

# Scattering angle correction: 4PM conservative



$$\begin{aligned}
 \delta\varphi_4^{\text{cons}} = & \left( r_1 + r_2 \operatorname{arccosh} \left( \frac{1}{\sqrt{1-v^2}} \right) + r_3 \operatorname{arccosh} \left( \frac{1}{\sqrt{1-v^2}} \right)^2 + r_4 \operatorname{E} \left( -\frac{v^2 + 2\sqrt{1-v^2} - 2}{v^2} \right)^2 \right. \\
 & \left. + r_5 \operatorname{K} \left( -\frac{v^2 + 2\sqrt{1-v^2} - 2}{v^2} \right) \operatorname{E} \left( -\frac{v^2 + 2\sqrt{1-v^2} - 2}{v^2} \right) + r_6 \operatorname{K} \left( -\frac{v^2 + 2\sqrt{1-v^2} - 2}{v^2} \right)^2 \right. \\
 & + r_7 \log \left( \frac{v}{2\sqrt{1-v^2}} \right) + r_8 \log \left( \frac{v}{2\sqrt{1-v^2}} \right) \operatorname{arccosh} \left( \frac{1}{\sqrt{1-v^2}} \right) \\
 & + r_9 \log \left( \frac{v}{2\sqrt{1-v^2}} \right) \log \left( \frac{1}{2} \left( \frac{1}{\sqrt{1-v^2}} + 1 \right) \right) + r_{10} \log \left( \frac{1}{2} \left( \frac{1}{\sqrt{1-v^2}} + 1 \right) \right) \\
 & \left. + r_{11} \log^2 \left( \frac{1}{2} \left( \frac{1}{\sqrt{1-v^2}} + 1 \right) \right) + r_{12} \alpha + r_{13} \frac{\beta}{v^2} + r_4 \log(b) \right) \left( \frac{M}{b} \right)^4
 \end{aligned}$$

Elliptic integrals

Free coefficients

Log term

$r_i =$  rational coefficients

# Extraction of high-order conservative PM results



Subtract known **analytic** parts of conservative 4PM:

$$\Delta_4(v) := (\delta\chi_4^{\text{cons}} - \delta\chi_4^{\text{known}})b^4 = \frac{3}{8}\pi M^4 [c_2 + c_1(5 - 4/v^2)] + \mathcal{O}(1/b)$$

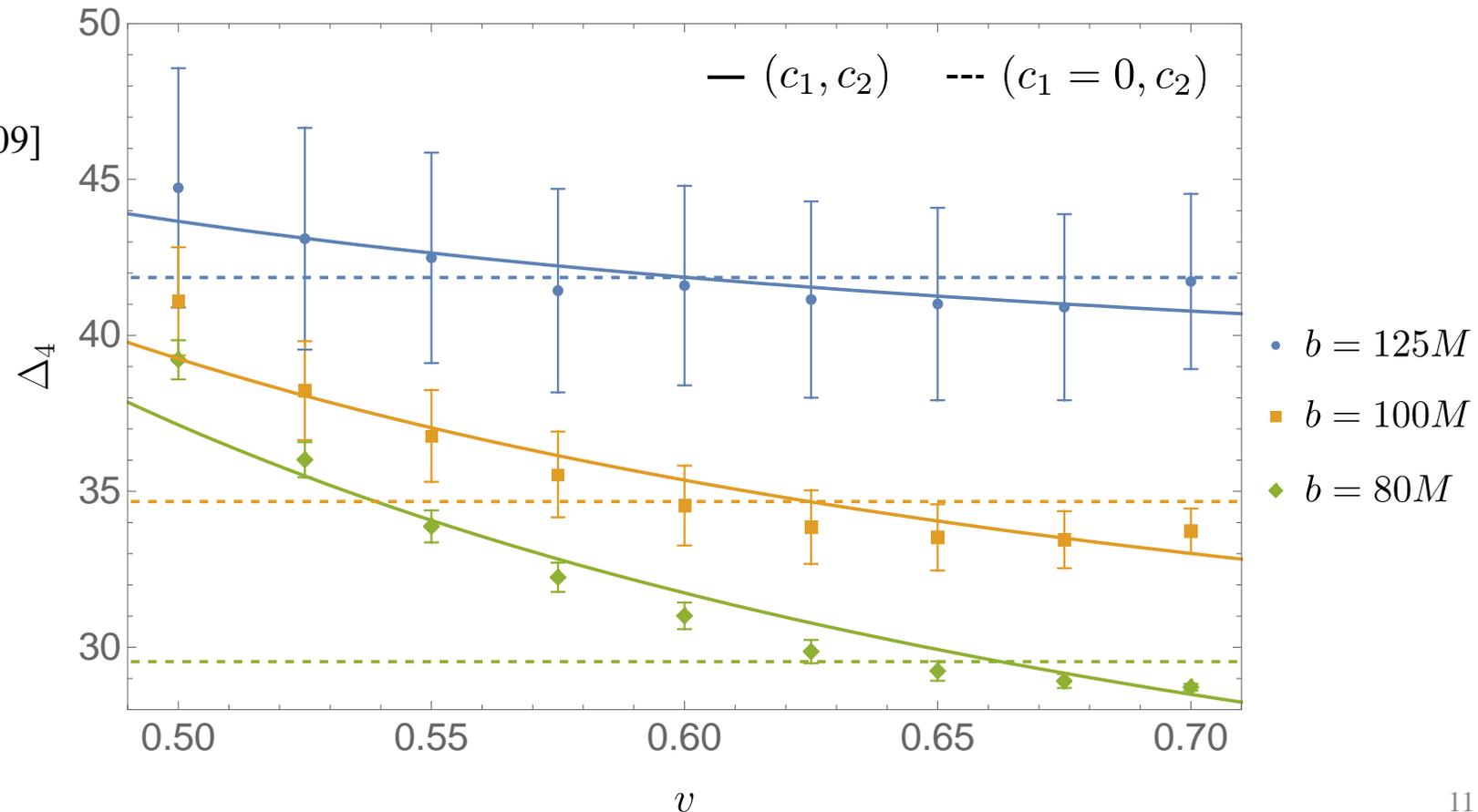
$c_1$  &  $c_2$  are **Love numbers**.

Expect  $c_1 = 0$  [Binnington & Poisson '09]

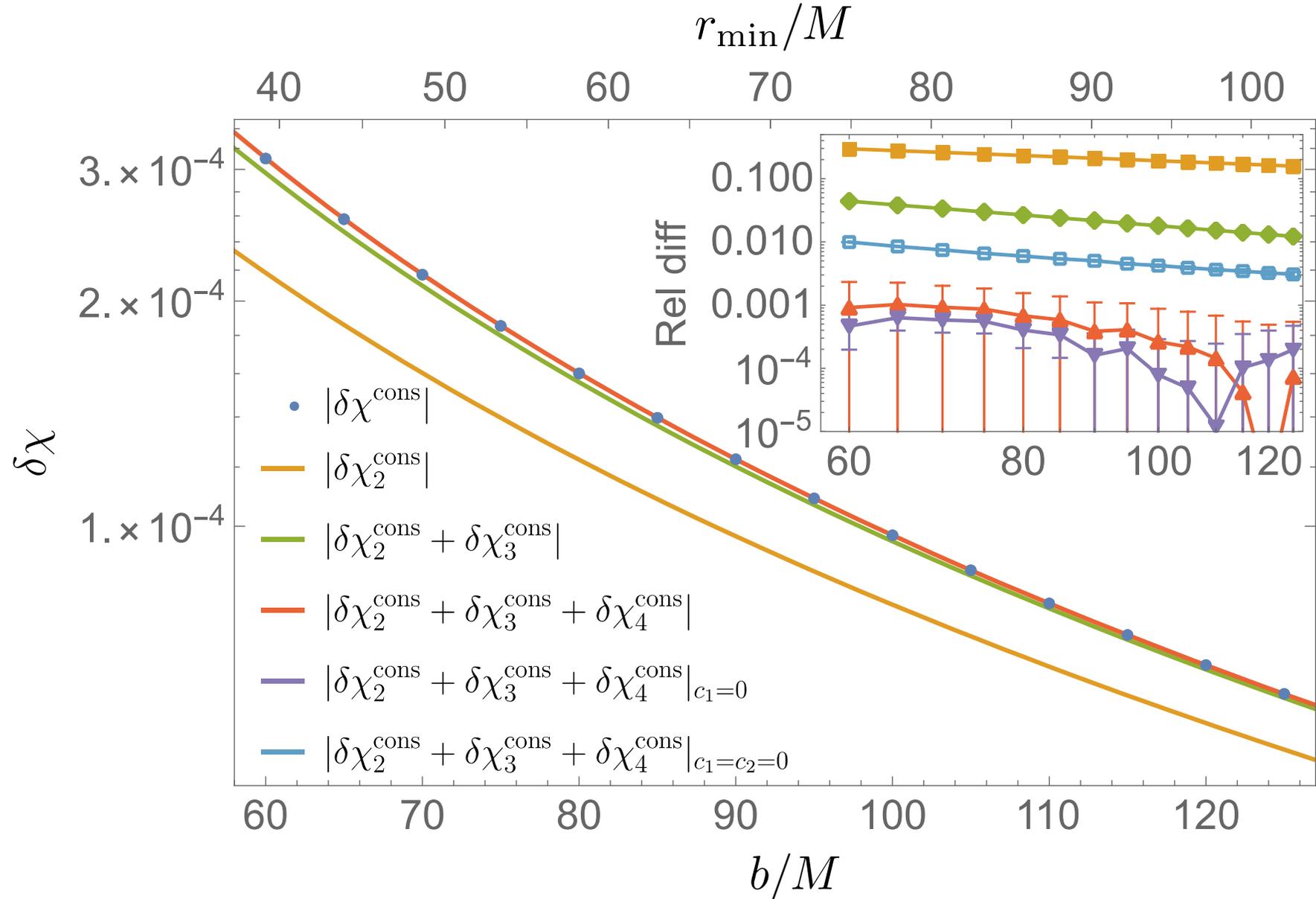
$b/M$	$c_1$	$c_2$
80	0.94	-21.2
100	0.68	-25.9
125	0.31(±0.38)	-33.6
80	0	-25.1
100	0	-29.4
125	0	-35.5

**Fixed**

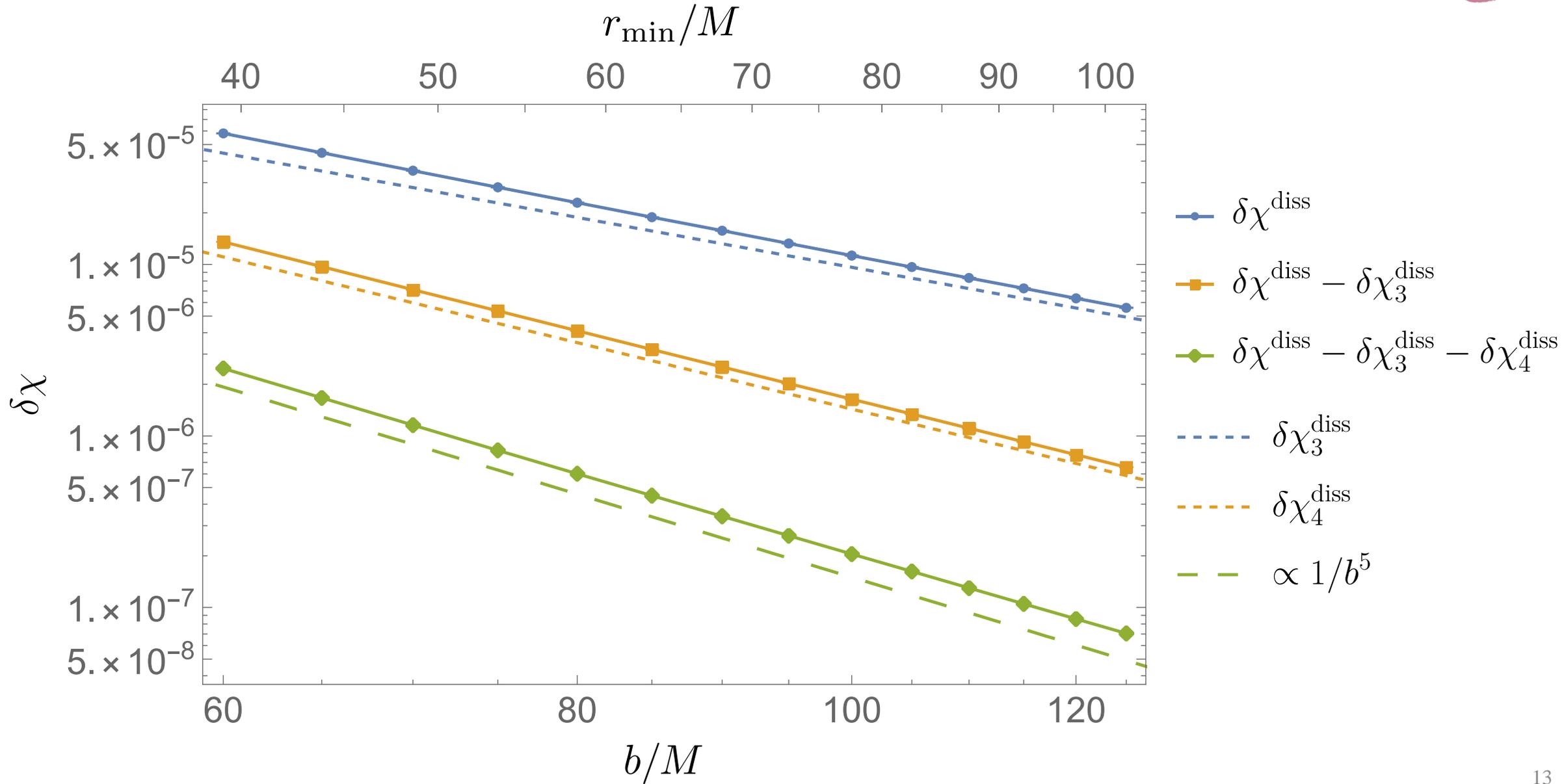
(Fitting error)



# PM comparison: Conservative $v = 0.5$



# Dissipative: $v = 0.5$



# Extraction of high-order dissipative PM results



PM expansion with free parameters:

$$\delta\varphi_{\text{diss}}^{\text{PM}} = \frac{\alpha_3}{b^3} + \frac{\alpha_4}{b^4} + \frac{\alpha_5}{b^5} + \frac{\alpha_6}{b^6} + \dots$$

Up to **4PM** can fit value or use **analytic value**.

$\alpha_3$	$\alpha_4$	$\alpha_5$	$\alpha_6$
11.19	—	—	—
9.44	188	—	—
9.64	142	1900	—
9.61	154	920	26615
<b>9.6225</b>	169	—	—
<b>9.6225</b>	147	1720	—
<b>9.6225</b>	149	1321	15859
<b>9.6225</b>	<b>143.344</b>	1965	—
<b>9.6225</b>	<b>143.344</b>	2248	−20216
< 1%	~ 1%	~ 2000(?)	???

# Summary and future work

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Compared numerical **scalar self-force** results with analytic PM results.

Good **agreement** for both conservative and dissipative:

- $\sim 10\%$  for LO.
- $\sim 1\%$  for NLO.

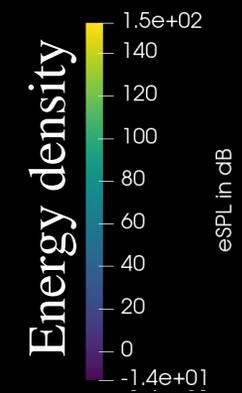
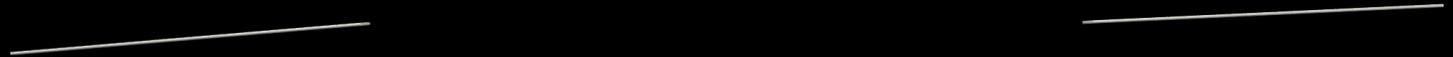
Have investigated extractions of higher-order components from numerics:

- Fixed free parameters in **4PM conservative**: agreement with numerics to  $< 0.5\%$ .
- Extraction of **5PM dissipative** coefficient.

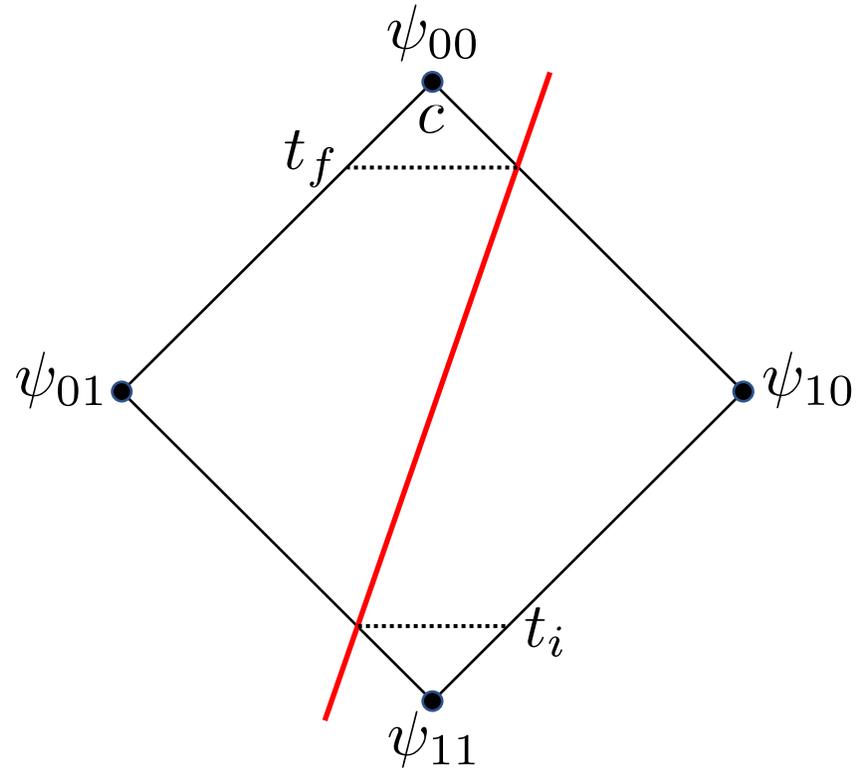
Future work:

- Hybrid **FD/TD** model for more accurate calculations (with **C. Whittall**).
- 1+1D with **hyperboloidal** slicing and **compactification** (with **R. Panosso Macedo**).
- Calculate the **gravitational self-force** correction to the scattering angle.
- Compare to **PM/NR/EOB** in the gravitational case.

Time: -280



# Scalar field evolution scheme



Finite-difference version of scalar wave equation:

$$\psi_{00} = -\psi_{11} + (\psi_{01} + \psi_{10}) \left( 1 - \frac{1}{2} h^2 V \right) + Z + \mathcal{O}(h^4)$$

$Z = 0$  if particle does not enter cell.

