Analytic Solutions to Plunging Geodesics in Kerr



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- Geodesic motion is the 0th order motion in the self-force expansion
- Transition to Plunge smoothly connects inspiral to a plunging geodesic
- For IMRI/EMRIs plunging geodesic sources the ringdown

The Plunge in Self-Force





Previous Work on Analytic Solutions

- Extensive history of solutions to special cases, Chandrasekhar, Wilkins etc...
- Introduction of Mino time, $d\tau = \Sigma d\lambda$ [Mino: 03020]
- Analytic solutions to generic bound geodesics [Fujita, Hikida: 0906]
- And much more...



Capra meeting on Radiation Reaction in General Relativity

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• Plunge following the approach of Fujita and Hikida

• Easily related back to the action angle parameterisation

• Efficiently implement in the Black Hole Perturbation Toolkit

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Bound Plunges

• Write geodesic equations in terms of conserved quantities $(\mathscr{E},\mathscr{L},Q)$

$$\left(\frac{dr}{d\lambda}\right)^2 = R(r), \quad \text{- fourth of}$$
$$\left(\frac{dz}{d\lambda}\right)^2 = Z(z^2), \quad \text{- second}$$
$$\frac{dt}{d\lambda} = T_r(r) + T_z(z),$$
$$\frac{d\phi}{d\lambda} = \Phi_r(r) + \Phi_z(z).$$

• Restrict to $\mathscr{E} < 1$, Bound Geodesics

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- order poly in r
- order order z^2

- Need to solve new sets of radial elliptic integrals
- With $P_{0,1}(x)$ rational polynomials in x
- y^2 is a polynomial in x with order three or four

$$\mathcal{I} = \int_{a}^{b} \frac{f_{0}(x) + f_{1}(x)y}{g_{0}(x) + g_{1}(x)}$$

[Labahn and Mutrie: 0197]

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Elliptic Integrals

 $\frac{\partial y}{\partial y} dx = \int_{a}^{b} P_{0}(x) + \int_{a}^{b} \frac{P_{1}(x)}{y} dx$

Reducing the Problem

• Reduce $P_1(x)$ to sum over rational polynomials with simple poles

• One can then decomposed into four fundamental integrals

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$$egin{aligned} \lambda &= -\int rac{dr}{\sqrt{R(r)}}, & \mathcal{I}_r = \int rac{rdr}{\sqrt{R(r)}}, \ \mathcal{I}_{r^2} &= \int rac{r^2dr}{\sqrt{R(r)}} ext{ and } \mathcal{I}_{r_\pm} = \int rac{dr}{(r-r_\pm)\sqrt{R(r)}} \end{aligned}$$

- Geodesic type given by root structure of radial potential, $\frac{dr}{d\lambda} = \pm \sqrt{R(r)}$.
- Roots correspond of turning points of the trajectories

Four real roots

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Bound Plunges: Two Cases

Complex Roots: Solving the Integrals

- Four real roots, $R(r) = A(r r_1)(r r_2)(r r_3)(r r_4)$
- Two complex two real roots, $R(r) = B(r r_1)$

• Make manifestly real substitution r(y) such the

- Allows one to solve explicitly for the four fundamental integrals
- Analytical continue everything!

$$(r-r_2)(r^2-2\Re(r_3)r+\Re(r_3)-\Im(r_4))$$

hat
$$\frac{dr}{\sqrt{R(r)}} \to C \frac{dy}{\sqrt{(1-k_r^2 y^2)(1-y^2)}}$$

The Plunging Kerr Solutions

- Solutions in terms of Elementary and Elliptic functions
- All function manifestly real parameterised by - $(r_1, r_2, \mathfrak{R}(r_3), \mathfrak{S}(r_4))$
- Immediately obtain Mino frequencies - $\Upsilon_r = \Upsilon_r(\mathscr{E}, r_1, r_2, \Re(r_3), \Im(r_4)).$
- Can be easily recast in action angles parameterisation
- Implemented in the Black Hole Perturbation Toolkit

Bonus: Approach Works Generally

$$u^{a} \nabla_{a} u_{b} = \frac{1}{m} \left(eF_{ba} - g \star F_{ba} \right) u^{a}$$

- Same procedure can be used generically to find geodesics solutions analytically
- Solve for generic Bounded and Plunging geodesics in Dyonic Kerr-Newman

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• Recent work exploring phenomenology of dyonically charged particles in Kerr-Newman. [Dyson and Pereñiguez: arXiv:2306.15751]

• Solutions in terms of Elementary and Elliptic functions

• Can be easily recast in action angles parameterisation

• Implemented in the Black Hole Perturbation Toolkit

Outcome

