

# Resonant dynamics of extreme mass-ratio inspirals in a perturbed Kerr spacetime

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# Outline

- **Phenomenology of near-resonance orbits**
  - Poincare map and rotation number (*3 kinds of orbits*)
- **Unified description in action-angle variables**
  - Simple harmonics
- **Impact for EMRI evolution**
  - Phase shifts

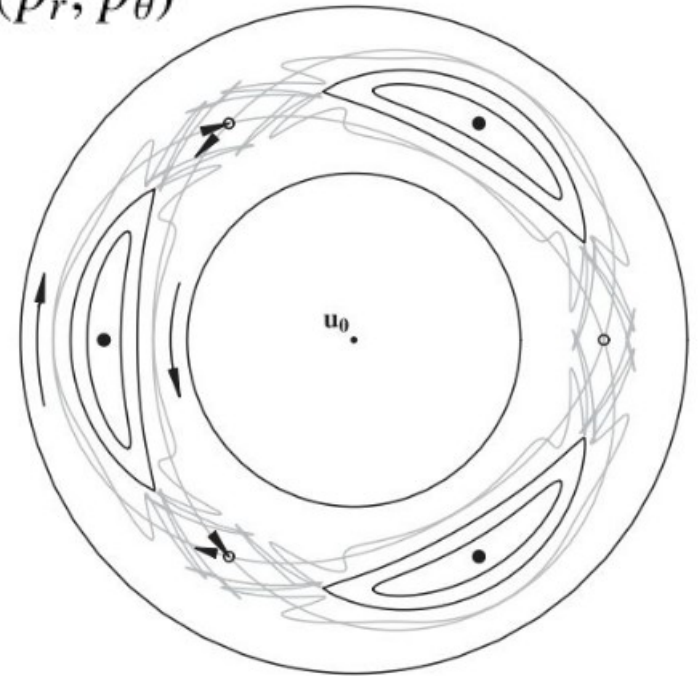
# 1. Phenomenology: Poincare map

A test particle in Kerr  $\rightarrow$  integrable  $(r, \theta) + (p_r, p_\theta)$   
in perturbed Kerr  $\rightarrow$  non-integrable

**Poincare map:**  $(r, p_r)_{\theta=\pi/2}$

## 3 kinds of orbits

- Regular orbits (O1)
- Chaotic transitional orbits (O2)
- On-resonance orbits (O3)



Cartoon Poincare Map  
Lukes-Gerakopoulos+ 2010

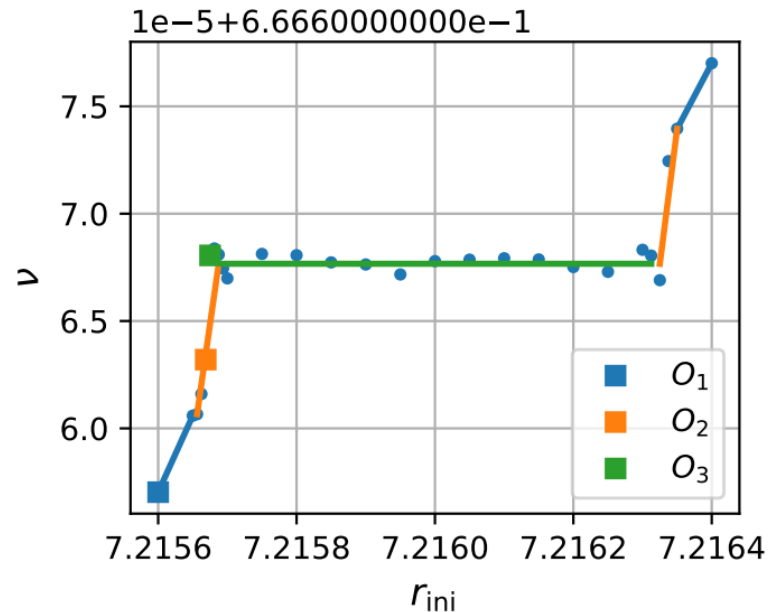
# 1. Phenomenology: rotation number

A test particle in a perturbed Kerr:  $(r, \theta) + (p_r, p_\theta)$

Rotation number:  $\nu := \frac{\langle \Omega^r \rangle}{\langle \Omega^\theta \rangle}$

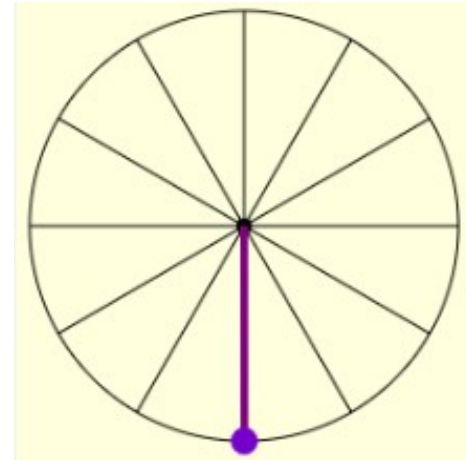
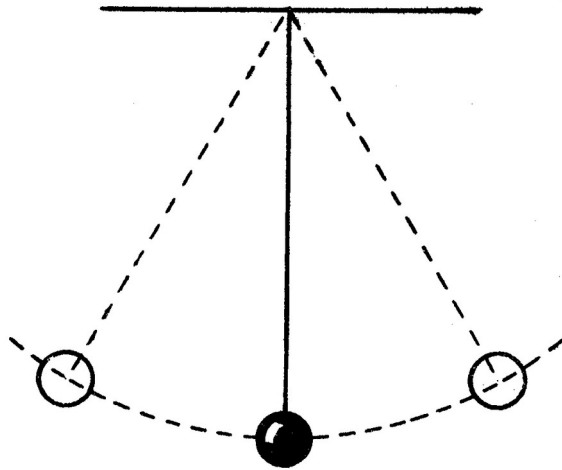
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- Regular orbits (O1)
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- On-resonance orbits (O3)



## 2. Physics and a unified description

Reduced DOF and simple harmonics



**Libration** → regular orbits (O1)      **Rotation** → on-resonance orbits (O3)  
**Transition** → chaotic transitional orbits (O2)

# Unified description: in action-angle variables

- Kerr spacetime  $\mathcal{J}_\alpha = \mathcal{J}_\alpha(\{x^\beta, p_\beta\})$ ,  $q^\alpha = q^\alpha(\{x^\beta, p_\beta\})$ .

$$H_0(x^\mu, p_\nu) = \frac{1}{2} g_{\text{Kerr}}^{\mu\nu} p_\mu p_\nu \longrightarrow H_0^{aa} = H_0^{aa}(\mathcal{J}_\alpha)$$

$$\dot{\mathcal{J}}_\alpha = -\frac{\partial H_0}{\partial q^\alpha} = 0, \quad \dot{q}^\alpha = \frac{\partial H_0}{\partial \mathcal{J}_\alpha} = \Omega^\alpha$$

$$\text{Resonance : } n_r \Omega^r + n_\theta \Omega^\theta = 0$$

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- Perturbed Kerr

$$H(x^\mu, p_\nu) = \frac{1}{2} g^{\mu\nu} p_\mu p_\nu = H_0 + \epsilon H_{\text{int}} \longrightarrow H_0(\mathcal{J}_\alpha) + \epsilon H_{\text{int}}(\mathcal{J}_\alpha, q^\beta)$$

$$\frac{d\mathcal{J}_\alpha}{d\tau} = -\epsilon \frac{\partial \mathcal{H}_{\text{int}}}{\partial q^\alpha}, \quad \frac{dq^\alpha}{d\tau} = \Omega_\alpha + \epsilon \frac{\partial \mathcal{H}_{\text{int}}}{\partial \mathcal{J}_\alpha},$$

# Near Identity Transformation: resonance as a driven force

$$\tilde{q}^\alpha = q^\alpha + \epsilon \sum_{n_\alpha, \Omega^\alpha \neq 0} \frac{1}{n_r \Omega^r + n_\theta \Omega^\theta} \frac{\partial H_{n_j}}{\partial \mathcal{J}_\alpha} e^{i(n_r q^r + n_\theta q^\theta)},$$

$$\tilde{J}_\alpha = J_\alpha + \epsilon \sum_{n_\alpha, \Omega^\alpha \neq 0} \frac{n_\alpha}{n_r \Omega^r + n_\theta \Omega^\theta} H_{n_j} e^{i(n_r q^r + n_\theta q^\theta)}.$$

**We have**

$$\frac{d\tilde{q}^\alpha}{d\tau} = \Omega_\alpha + \epsilon \sum_{k \in \mathbb{Z}} \frac{\partial H_{k\mathbf{N}}}{\partial \mathcal{J}_\alpha} e^{ik(N_r q^r + N_\theta q^\theta)},$$

$$\frac{d\tilde{J}_\alpha}{d\tau} = -i\epsilon N_\alpha \sum_{k \in \mathbb{Z}} k H_{k\mathbf{N}} e^{ik(N_r q^r + N_\theta q^\theta)}.$$

e.g., 2/3 resonance  $\mathbf{N} = (N_r, N_\theta) = (-3, 2)$



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**Not all action angle variables  
are independent.**

$$\frac{dQ}{d\tau} := \frac{d}{d\tau} (N_r \tilde{q}^r + N_\theta \tilde{q}^\theta) = N_\alpha \Omega^\alpha + O(\epsilon) = \Delta\Omega = O(\epsilon)$$

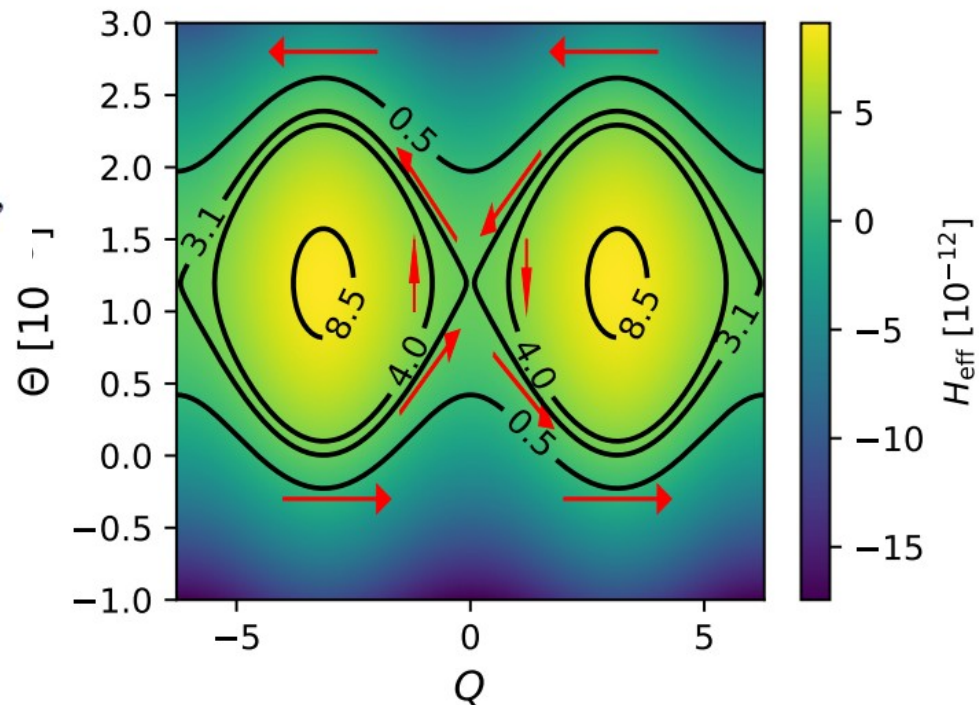
$$\frac{d\Theta}{d\tau} := \frac{1}{N_r} \frac{d\tilde{J}_r}{d\tau} = \frac{1}{N_\theta} \frac{d\tilde{J}_\theta}{d\tau} = \sum H_{k\mathbf{N}} e^{ikQ} = O(\epsilon)$$

# Near-resonance orbits :

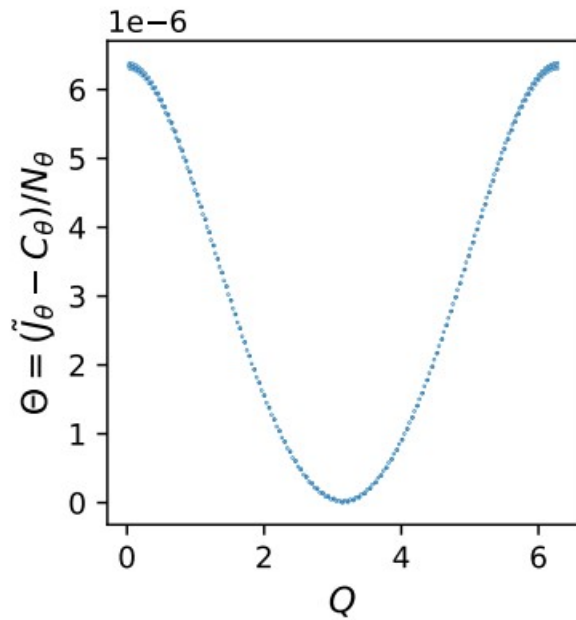
1 d.o.f effective Hamiltonian

$$H_{\text{eff}} := H_{\text{kin}} + H_{\text{pot}} = \int_0^{\Theta} \Delta\Omega(\Theta) d\Theta + \sum_k H_k \mathbf{N} e^{ikQ},$$

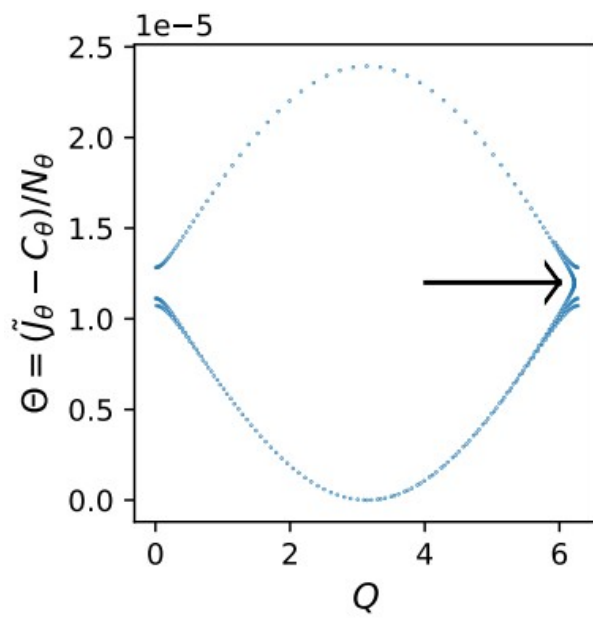
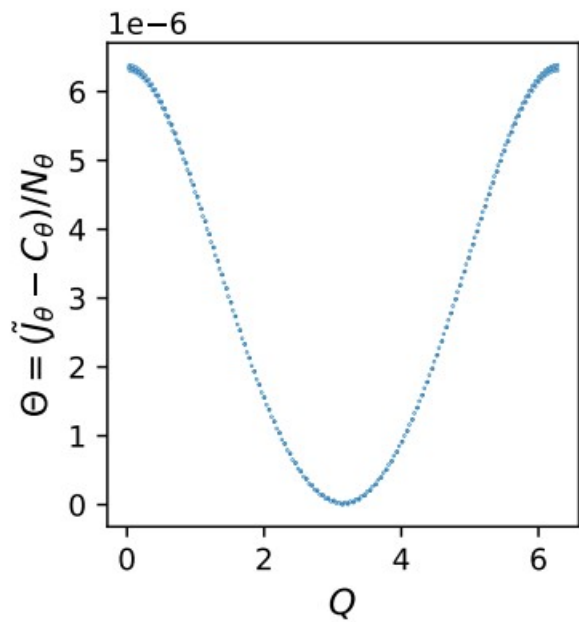
$$\frac{dQ}{d\tau} = \frac{\partial H_{\text{eff}}}{\partial \Theta},$$
$$\frac{d\Theta}{d\tau} = -\frac{\partial H_{\text{eff}}}{\partial Q}.$$



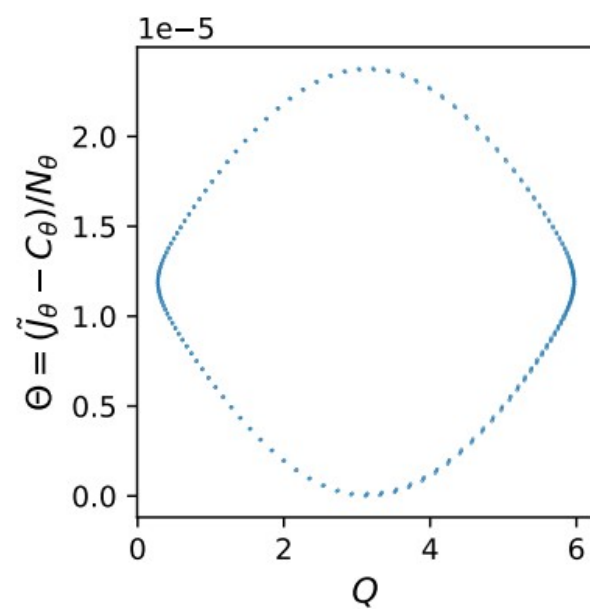
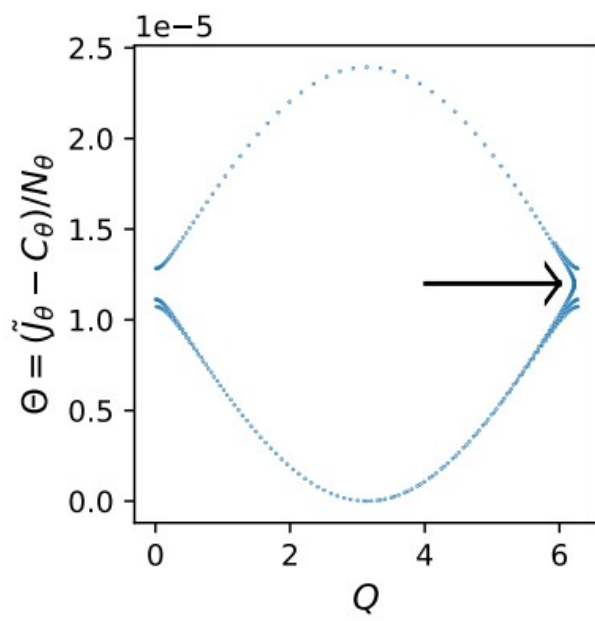
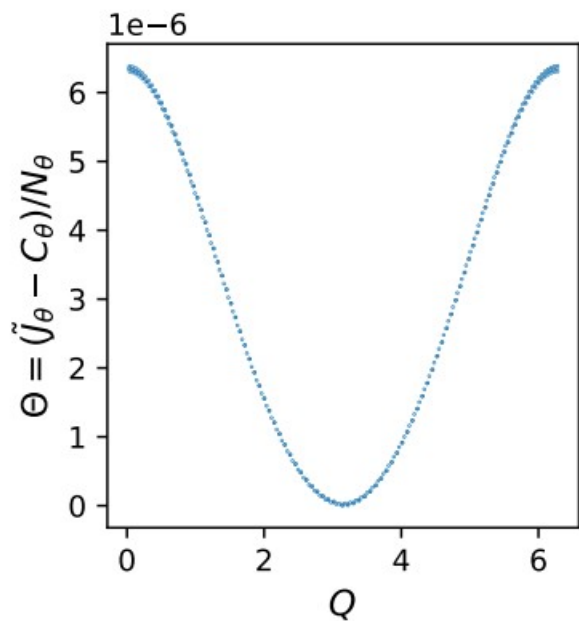
Numerical check:  $(x, p) \rightarrow (q, J) \rightarrow (Q, \Theta)$



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### 3. Impact for EMRI evolution

$$\frac{dQ}{d\tau} = \sum_{\alpha} a^{\alpha} \tilde{\mathcal{J}}_{\alpha}$$

$$\frac{d\tilde{\mathcal{J}}_{\alpha}}{d\tau} = -N_{\alpha} H_{\text{res}} \sin Q + q G_{\alpha}$$

$$\Delta\tilde{\mathcal{J}}_{\alpha} = -N_{\alpha} H_{\text{res}} \int_{-\infty}^{\infty} \sin Q(\tau) d\tau$$

# Non-adiabatic limit:

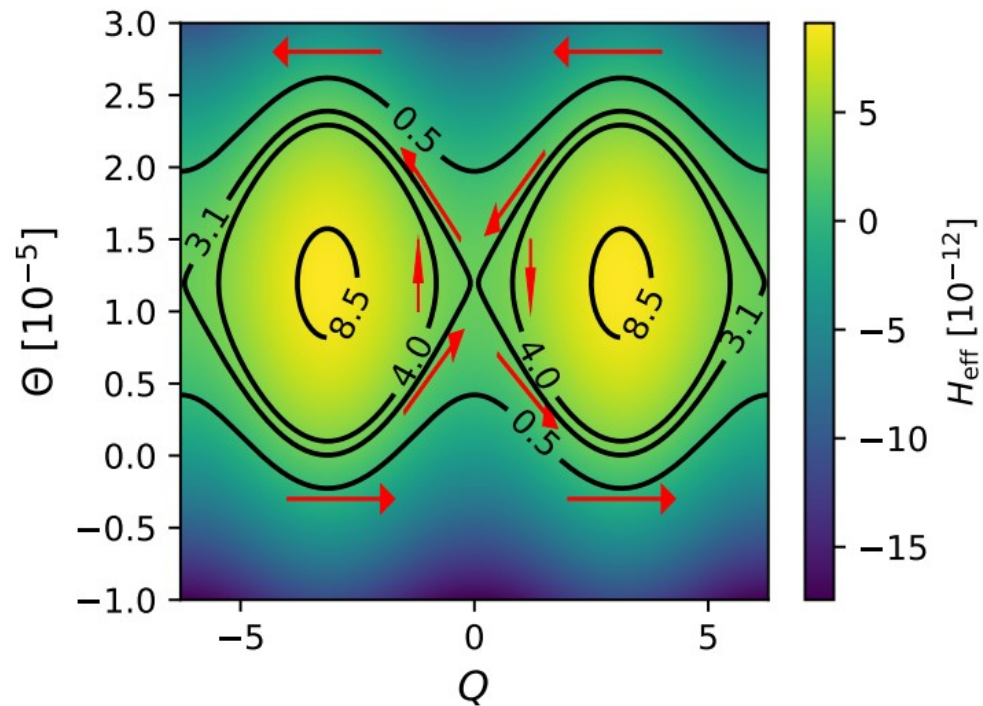
$$t_{\text{cross}} < t_{\text{res}} \quad (h_{\text{res}} < q = m/M)$$

*Crossing the resonance via  
passing through all the orbits*

$$\delta J \sim h_{\text{res}} t_{\text{cross}} \sim h_{\text{res}} / \sqrt{q}$$

$$\delta \Psi \sim \delta J / q \sim h_{\text{res}} / q^{3/2}$$

$$\delta \Psi|_{\text{Kerr}} \sim 1/q^{1/2} \quad (\text{Flanagan+2012})$$



# Adiabatic limit: $t_{\text{cross}} > t_{\text{res}}$ ( $h_{\text{res}} > q$ )

Area conserved as an adiabatic invariant

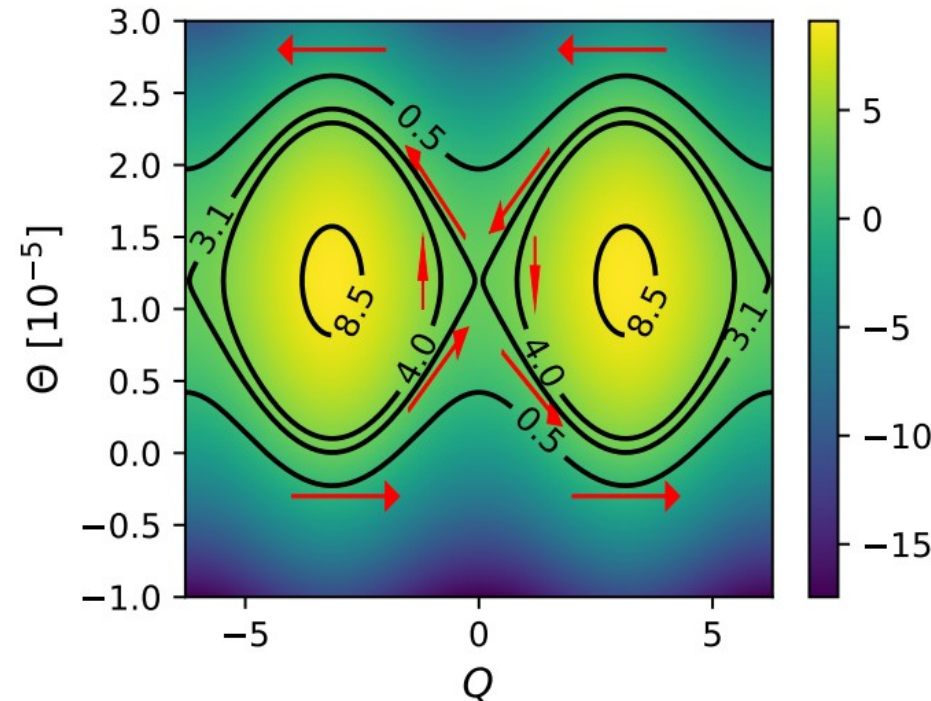
$$\int \Theta dQ$$

Circumventing resonance via chaotic transitional orbits.

$$\Delta \tilde{J}_\alpha = -N_\alpha H_{\text{res}} \int \sin Q(\tau) d\tau + \Delta \tilde{J}_{\alpha, \text{transit}}$$

$$\delta J \approx \delta \Theta \sim \sqrt{h_{\text{res}}}$$

$$\delta \Psi \sim \delta J/q \sim \sqrt{h_{\text{res}}/q}$$





# Summary

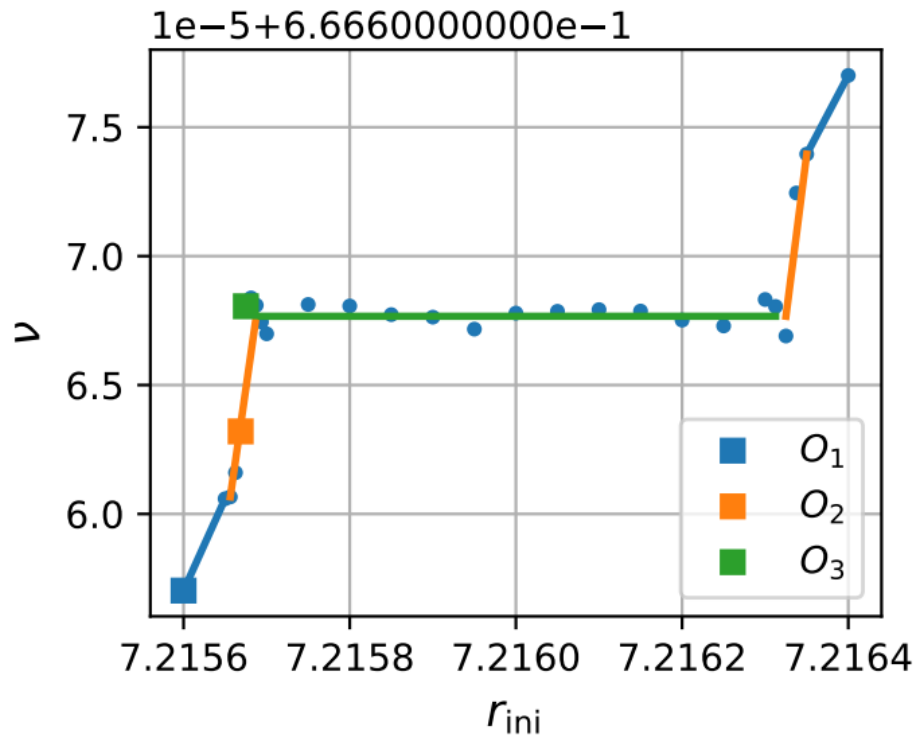
- **Rich phenomena of resonant dynamics can be understood with effective Hamiltonian of simple harmonics.**
- **Only orbits in the transitional regime are chaotic.**
- **Impact for resonance crossing**
  - *In adiabatic regime, resonance crossing does not happen because resonances are avoided by the transitional orbits.*

# 32<sup>nd</sup> Texas Symposium on Relativistic Astrophysics

**Texas in Shanghai**



# Near-resonance orbits : rotation numbers



$$\nu := \frac{\langle \Omega^r \rangle}{\langle \Omega^\theta \rangle}$$

*3 kinds of orbits:*

$O_1, O_2, O_3$

# Integrable systems and KAM theorem

- **Integrable system:** # degrees of freedom = # conserved quantities

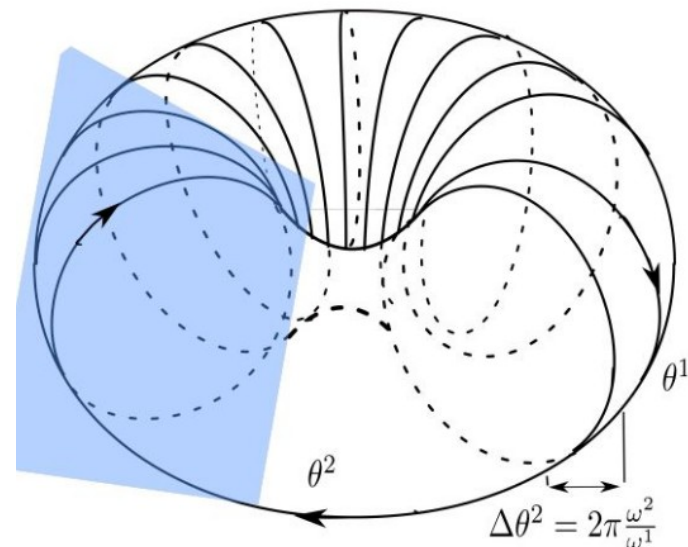
e.g. a test particle in the Kerr spacetime:

$$(t, r, \theta, \phi) \longleftrightarrow (H, E, L, C)$$

e.g. 2-d.o.f. integrable system:

orbit wraps on a 2-torus

(Cardenas-Avendano+2018)

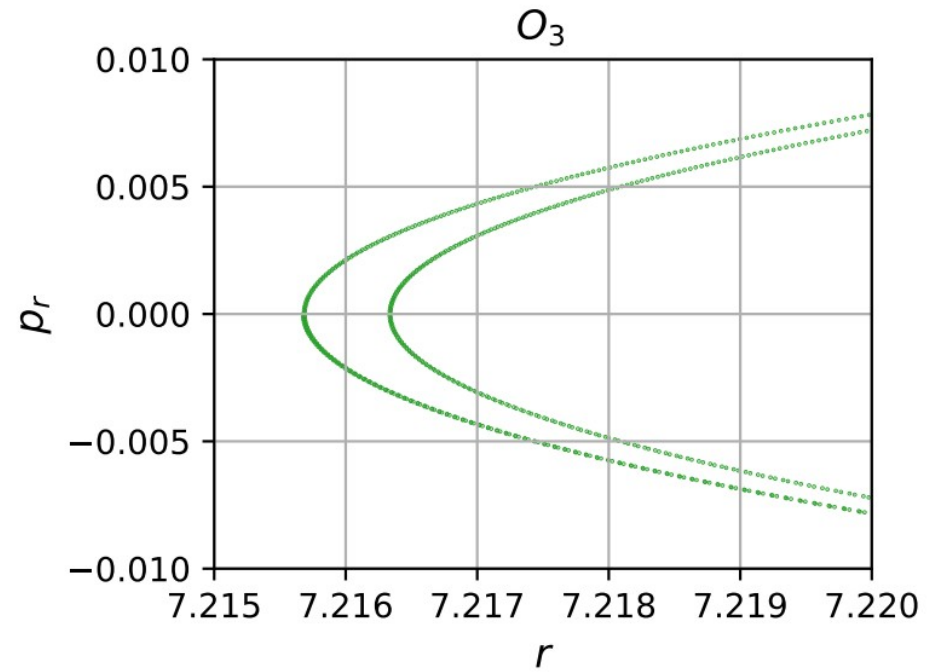
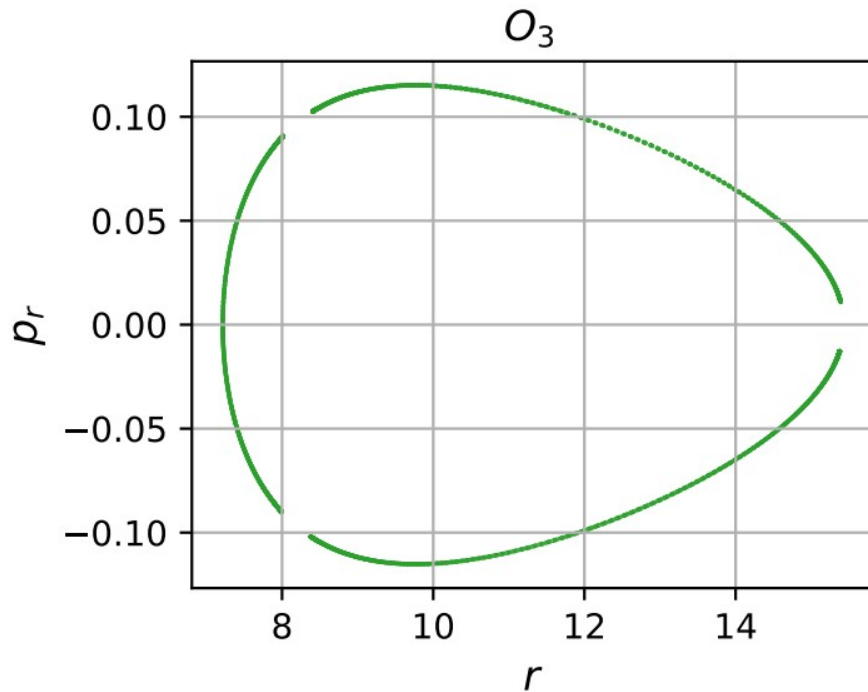


- **KAM theorem:**

an integrable system + a perturbation  $\rightarrow$  non-integrable

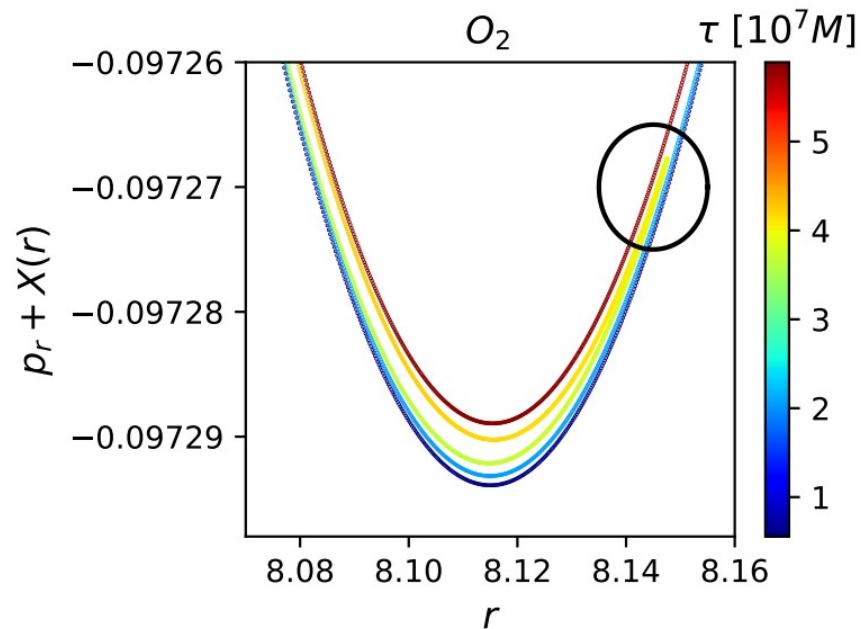
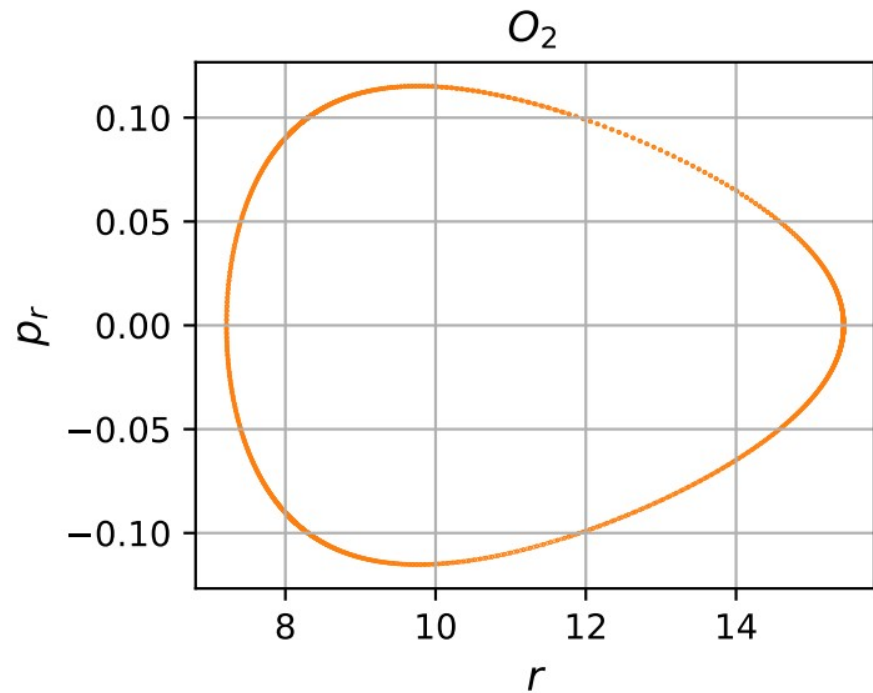
**Chaos can occur at resonances.** e.g. Quadractic Gravity (Donoghue+2021)  $g = g_{\text{Kerr}} + h$ .

# Near-resonance orbits : phenomenology



Poincare maps of  $O_3 \rightarrow$  **On-resonance Orbit**

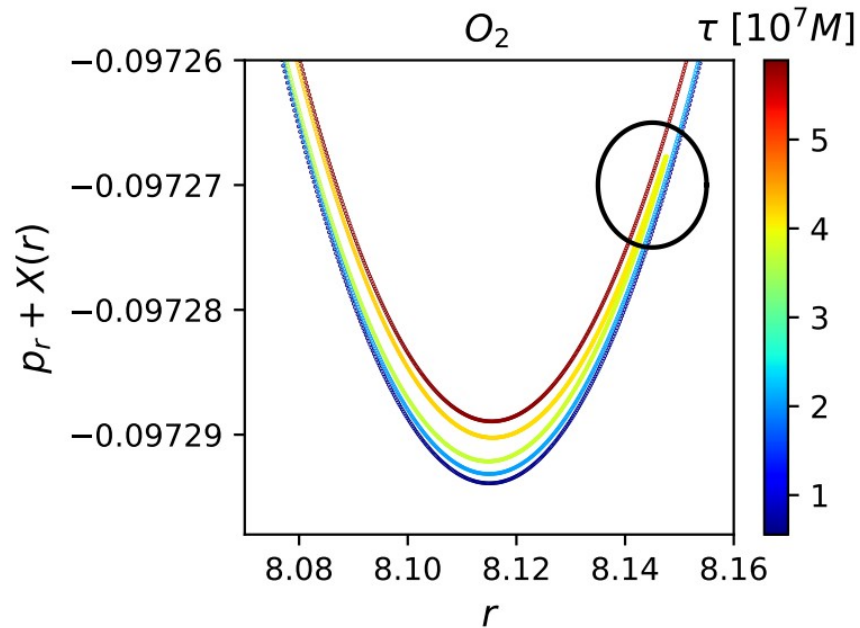
# Near-resonance orbits : phenomenology



Poincare maps of  $O_2 \rightarrow$  **Chaotic Transitional Orbit**

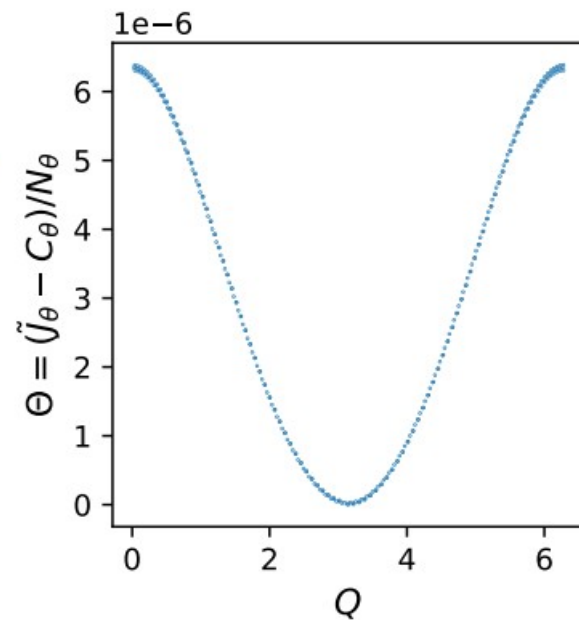
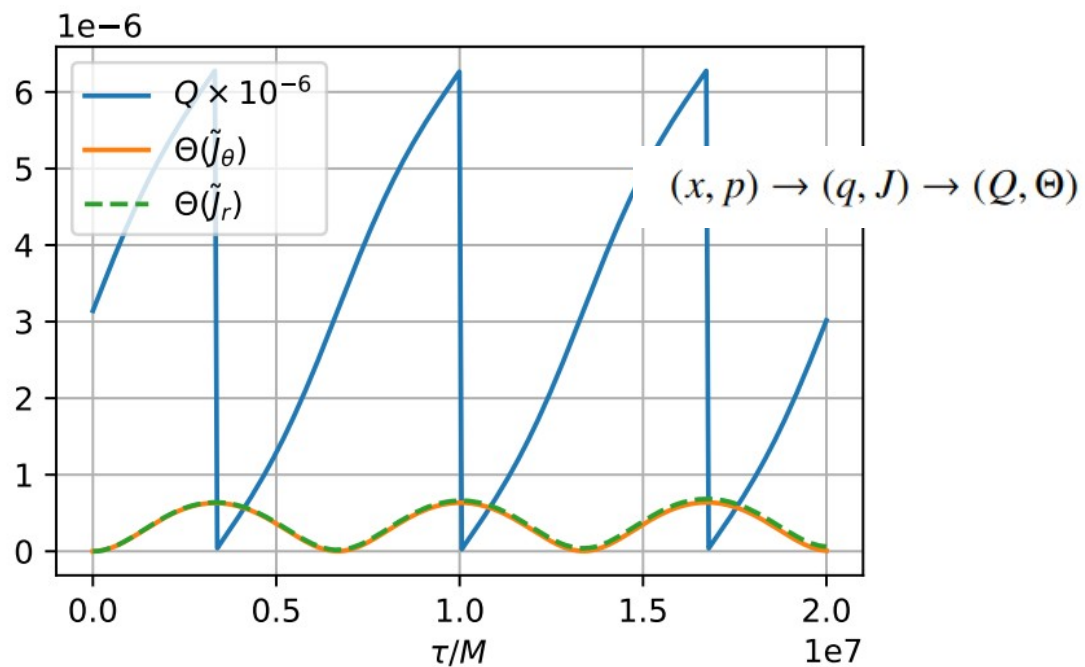
# O2 Poincare map:

e.g. spinning BH in Quadratic Gravity (Donoghue+2021)



***Anti-clockwise* → *clockwise***

# Regular orbit O1





# Chaotic transitional orbit O2

