## 2nd-order Gravitational Self-Force in Schwarzschild: Mode Decomposition of the 1st-Order Puncture Jonathan Thornburg

Department of Astronomy and Center for Spacetime Symmetries Indiana University


Work done as part of 2SF group:

## 2nd-order GSF in Schwarzschild: the Big Picture

Puncture scheme near the particle:

1st order puncture:
${ }^{(1)} h_{a b}^{(\text {puncture })}$

Terminology:

- Penrose abstract-index notation
- $a b$ are tensor indices
- ${ }^{(n)} h_{a b}$ is the $n$ th-order metric perturbation


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compute mode decomposition
of 1st-order puncture
Barack-Lousto-Sago modes of the 1st-order puncture: $\left.{ }^{(1)} h_{a b}^{(\text {puncture })}\right)^{\mathrm{I} \ell m}$

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## Mode decomposition of the 1st-order puncture

Conceptually, computing the mode decomposition is easy: the Barack-Lousto-Sago modes $Y_{a b}^{\mathrm{I} \ell m}$ are orthogonal, so

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- Compute it analytically $\Rightarrow$ this talk


## How many $\beta$ integrals are there?

Each " $\beta$ integral" is actually a set of integrals, one for each Barack-Lousto-Sago tensor mode and ( $\ell, m$ ):

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- we typically compute for $(\ell, m)$ in the range $0 \leq \ell \lesssim 50,0 \leq m \lesssim 10$
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Each $\beta$ integral depends on 2 parameters:

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Each self-force computation requires numerically evaluating the " $\beta$ integral" set on a grid of 100-1000 $\Delta r$ values.
$\Rightarrow$ Need $10^{3}-10^{5}$ numerical evaluations of each of the $\sim 2500$ individual integrals

## Typical form of an individual integrand

For the $\mathrm{I}=1$ Barack-Lousto-Sago mode, the $\ell=0, m=0$ integrand is:

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\begin{aligned}
I_{1,00}= & \frac{P_{3}\left(\sin ^{2} \beta\right) P_{6}^{(1)}\left(1-\frac{M \sin ^{2} \beta}{r_{0}-2 M}\right)}{\left(r_{0}-2 M-M \sin ^{2} \beta\right)^{5 / 2}} \\
& \times\left[P_{6}^{(2)}\left(\frac{1}{1-\frac{M \sin ^{2} \beta}{r_{0}-2 M}}\right)\left(\frac{K_{1}}{\left(1-\frac{M \sin ^{2} \beta}{r_{0}-2 M}\right)^{3 / 2}}+\left(K_{2}+\frac{K_{3}}{1-\frac{M \sin ^{2} \beta}{r_{0}-2 M}}\right)^{3 / 2}\right)\right.
\end{aligned}
$$

where each $K_{i}$ is a "constant" and each $P_{k}$ or $P_{k}^{(i)}$ is a polynomial of degree $k$. The "constants" $K_{i}$ and the polynomial coefficients all depend on the parameters $r_{0}$, and $\Delta r$.

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This is the simplest of the integrands; the integrands rapidly become more complicated with increasing $\ell$ and/or $m$.

## Overall strategy for doing the $\beta$ integrals

None of the symbolic algebra systems I tried (Mathematica, Mathematica with the RUBI rules-based-integration package, Maple, Sage) could do the $I=1$, $\ell=0, m=0$ integral directly.

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- "flatten" the integrand into a single linear combination $K+\sum_{k} c_{k} X_{k}$, where the coefficients $K$ and $\left\{c_{k}\right\}$ depend on $r_{0}$ and $\Delta r$, but not on $\beta$ :


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- assemble the final result from $K, c_{k}$, and the $X_{k}$ integrals


## Example of flattening into a linear combination

For the $\mathrm{I}=1$ Barack-Lousto-Sago mode, the $\ell=0, m=0$ integrand is a linear combination of 251 components. Some examples:

$$
\begin{aligned}
& X_{1}=\frac{1}{\left(M \cos ^{2} \beta+r_{0}-3 M\right)^{4}} \\
& X_{10}=\frac{\left[\left(4 M r_{0}^{2}-8 M^{2} r_{0}\right) \cos ^{2} \beta+\left(r_{0}-3 M\right)(\Delta r)^{2}+4 r_{0}^{3}-20 M r_{0}^{2}+24 M^{2} r_{0}\right]^{3 / 2}}{\left(M \cos ^{2} \beta+r_{0}-3 M\right)^{7}} \\
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& X_{200}=\frac{1}{\left(M \cos ^{2} \beta+r_{0}-3 M\right)^{8}\left[\left(4 M r_{0}^{2}-8 M^{2} r_{0}\right) \cos ^{2} \beta+\left(r_{0}-3 M\right)(\Delta r)^{2}+4 r_{0}^{3}-20 M r_{0}^{2}+24 M^{2} r_{0}\right]^{3 / 2}} \\
& X_{251}=\frac{\sin ^{6} \beta}{\left[\left(4 M r_{0}^{2}-8 M^{2} r_{0}\right) \cos ^{2} \beta+\left(r_{0}-3 M\right)(\Delta r)^{2}+4 r_{0}^{3}-20 M r_{0}^{2}+24 M^{2} r_{0}\right]^{1 / 2}\left(M \cos ^{2} \beta+r_{0}-3 M\right)^{6}}
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Substituting $x=\sin \beta$ converts each of our component integrals $\int_{0}^{2 \pi} X_{k} d \beta$ into an elliptic integral.

## Elliptic integrals

Formally, an elliptic integral is an integral

$$
\int_{a}^{b} R\left(x, \sqrt{P_{3 \mid 4}(x)}\right) d x
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where $R$ is a rational function and $P_{3 \mid 4}$ is a polynomial of degree 3 or 4 .

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- numerical computation of $E, K$, and $\Pi$ is (can be) very efficient
- but the (symbolic) reduction of an arbitrary elliptic integral to Legendre form can be very complicated
- Maple has excellent code built-in to do this reduction (better than Mathematica or Mathematica/RUBI; alas Sage is very poor at this) $\Rightarrow$ do the elliptic integrals in Maple


## Counting the $\boldsymbol{\beta}$ integrals for multiple $(\ell, m)$

How many elliptic integrals do we need to do?
For $I=1, \ell \in\{0,2,4, \ldots, 48\}, m=0$, we have:
$\ell \quad$ number of components $X_{k}$
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|  |  |
| ---: | ---: |
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In total, to do all of $\ell \in\{0,2,4, \ldots, 48\}$ (again just for $I=1, m=0$ ) requires 37,763 elliptic integrals.

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| $\ell$ | number of components $X_{k}$ |  | Number of components $\mathrm{X}_{\mathrm{k}}$ for each ell ( $\mathrm{l}=1, \mathrm{emm}=0$ ) |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | total | unique to this $\ell$ | 3000 | - number of components |  |  |  |
| 0 | 251 | 251 | $\times$ | + number of components | to th |  |  |
| 2 | 367 | 337 | $\stackrel{\text { ¢ }}{\text { ¢ }}$ |  |  |  |  |
| 4 | 471 | 432 |  |  |  |  |  |
| 8 | 679 | 618 | $\stackrel{8}{\square}$ |  |  |  |  |
| 16 | 1095 | 128 | - 1000 |  |  |  |  |
| 24 | 1511 | 158 |  | $\cdots 4^{+}$ |  |  |  |
| 36 | 2135 | 163 | 0 | ++*++++ | + + | + + |  |
| 48 | 2759 | 224 |  | 1020 | 30 | 40 | 50 |

In total, to do all of $\ell \in\{0,2,4, \ldots, 48\}$ (again just for $I=1, m=0$ ) requires 37,763 elliptic integrals.

Fortunately, many integrands are common to multiple $\ell$; for this same set of $\ell$ we "only" need to integrate 4518 unique integrands $X_{k}$.

## Cost of computing $\beta$ integrals for multiple $(\ell, m)$

Test computation:
$I=1, \ell \in\{0,2,4,6,8,10,12,16,20\}, m=0$
(each $\ell$ computed independently; duplicate integrals not removed)
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## An extended divide-and-conquer strategy

To make the computation more efficient, and extend to larger sets of $\mathrm{I}, \ell, m$, we extend our divide-and-conquer strategy to keep a database of components and integrals:

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Working on $I=1, \ell=0,2,4, \ldots, 48, m=0$ (Maple technical limitation prevents doing $\ell=50$ )

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- many integrals are very large: median size (printed out) is 116,000 characters, maximum size 6.7 million characters
- database size is currently 3.1 GB

[^3]
## Largest integral completed so far

\#3985 in the database:

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X_{k}=\frac{\left[\left(4 M r_{0}^{2}-8 M^{2} r_{0}\right) \cos ^{2} \beta+\left(r_{0}-3 M\right)(\Delta r)^{2}+4 r_{0}^{3}-20 M r_{0}^{2}+24 M^{2} r_{0}\right]^{1 / 2} \sin ^{2} \beta}{\left(M \cos ^{2} \beta+r_{0}-3 M\right)^{57}}
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## Integral took about 2 hours CPU time:

$\int_{0}^{2 \pi} x_{k} d \beta=$
$-1 / 832843192699583869114919332328039753766781580004556800 * r 0 *(380417909501120488$ 66012040380360645394177697720438803382086294874633847354336080936367189009818541 $51282803816058060800 * M^{\wedge} 110 * \mathrm{rO}^{\wedge} 55-20922985022561626876306622209198354966797733746$ $2413418601474621810486160448848445150019539554001978320554209883193344000 * \mathrm{M}^{\wedge} 109 *$ r0^56-42898808601999443484226115872348587122639223346825103202276211727893673430 $921803808212235313059899641759446702489600 * M^{\wedge} 109 * r 0^{\wedge} 54 * D e l t a \_r \wedge 2+570151341864804$ 33237935545520065517284523824458507656568901834443357478722311201303380324528465 $53909235102219317018624000 * M^{\wedge} 108 * r^{\wedge}{ }^{\wedge} 57+23165356645079699481482102571068237046225$ $18060728555572922915433306258365269777405643460706905234580655010121934438400 * \mathrm{M}^{-}$ $108 *$ r0^55*Delta_r^2-930677522987342195516434903727549826799107354752412894535838 ... skip 83967 lines
$8930131579947865654427648 * \mathrm{M} * \mathrm{rO}^{\wedge} 71 *$ Delta_r^96-85672898165824655870753189432131584 *M*r0^69*Delta_r^98+27259558507307845049785105728405504*M*r0^67*Delta_r^100-9086 $519502435948349928368576135168 * M *$ O~ $65 *$ Delta_r^102 +32451855365842672678315602057 $62560 * M *$ r0^63*Delta_r^104-1298074214633706907132624082305024*M*r0^61*Delta_r^106 $+649037107316853453566312041152512 * \mathrm{M} * \mathrm{rO}^{\wedge} 59 *$ Delta_r^108-6490371073168534535663120 $41152512 * M *$ r0^57*Delta_r^110-324518553658426726783156020576256*M*r0^55*Delta_r^1 $\left.12+2283850746669557096487036899494838068356102178363870539392483328 * \mathrm{r} 0^{\wedge} 168\right) / \mathrm{M} /(2$ $* \mathrm{M}-\mathrm{r} 0)^{\wedge} 56 / D e l t a \_r \wedge 110 /(-\mathrm{r} 0+3 * \mathrm{M}) /\left(9 * \mathrm{M}^{\wedge} 2-6 * \mathrm{M} * \mathrm{r} 0+\mathrm{rO}{ }^{\wedge} 2\right)^{\wedge} 27 /\left(16 * \mathrm{M}^{\wedge} 2 * \mathrm{rO}-16 * \mathrm{M} * \mathrm{r} \mathrm{O}^{\wedge} 2-3 * \mathrm{M} *\right.$ Delta_r^2+4*r0^3+r0*Delta_r^2) ^(1/2) *EllipticPi $\left(-M /(2 * M-r 0), 2 *\left(1 /\left(16 * M^{\wedge} 2 * r 0-16 * M\right.\right.\right.$ *r0^2-3*M*Delta_r $\left.\left.\left.{ }^{\wedge} 2+4 * r 0^{\wedge} 3+r 0 * D e l t a \_r \wedge 2\right) * M * r 0 *(-2 * M+r 0)\right)^{\wedge}(1 / 2)\right)$

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Still lots to explore here!

## Conclusions

Several parts of our 2nd-order self-force calculation require computing the Barack-Lousto-Sago tensor-spherical-harmonic modes of the 1st-order puncture. The main difficulty in doing this is the $\beta$ integrals.

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[^0]:    * Funding: Royal Society

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