

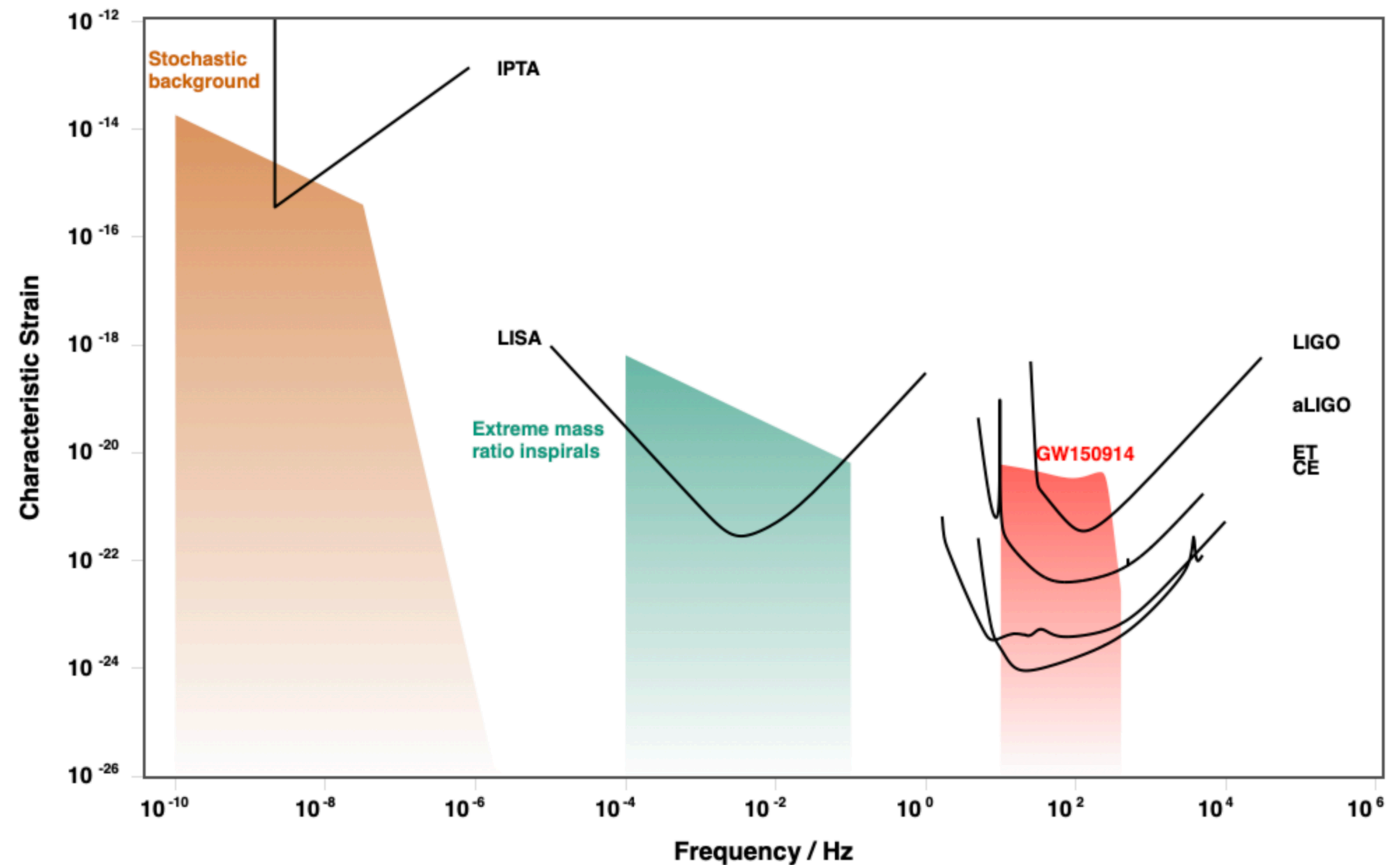
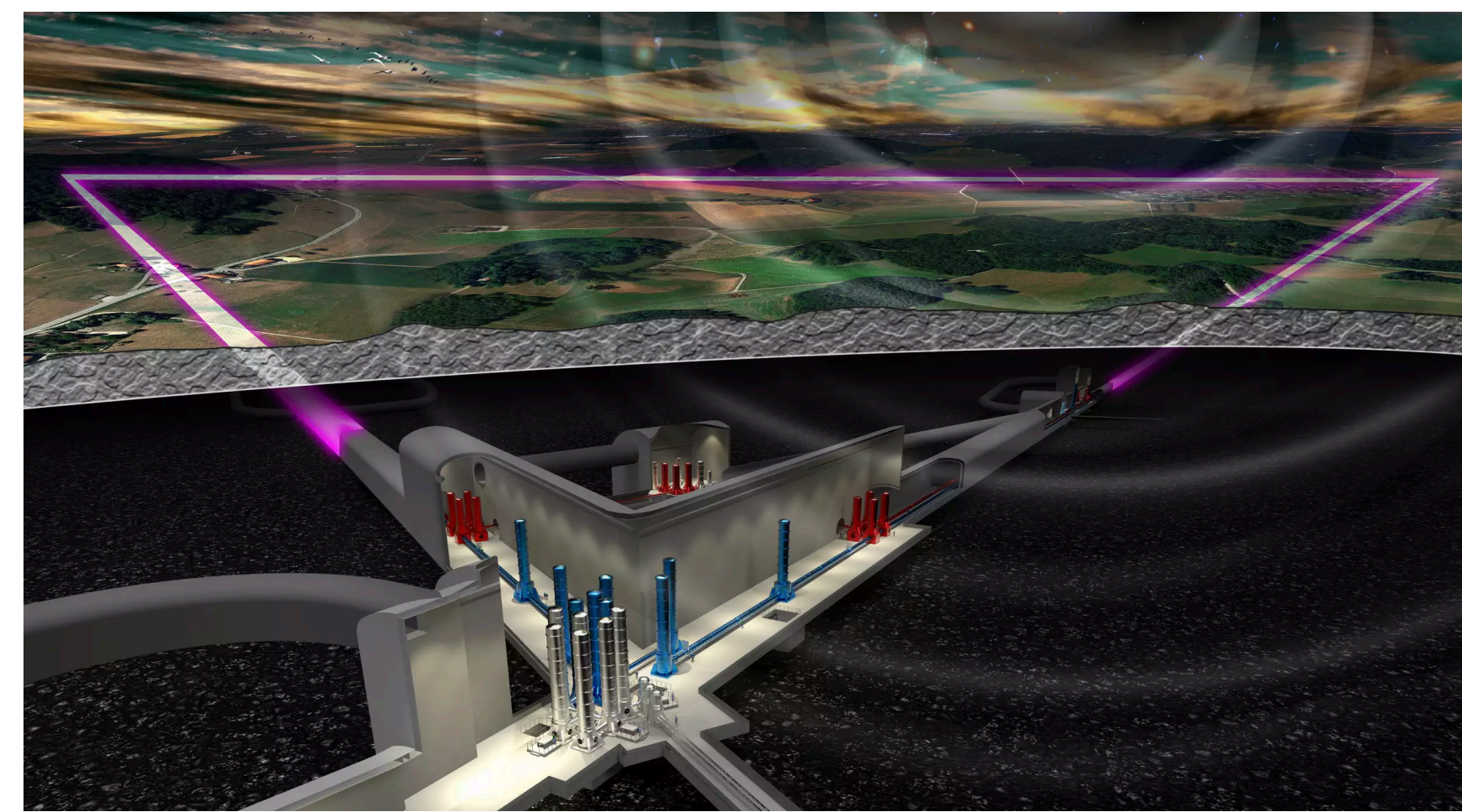
Tidal Heating in eccentric orbits

Sayak Datta

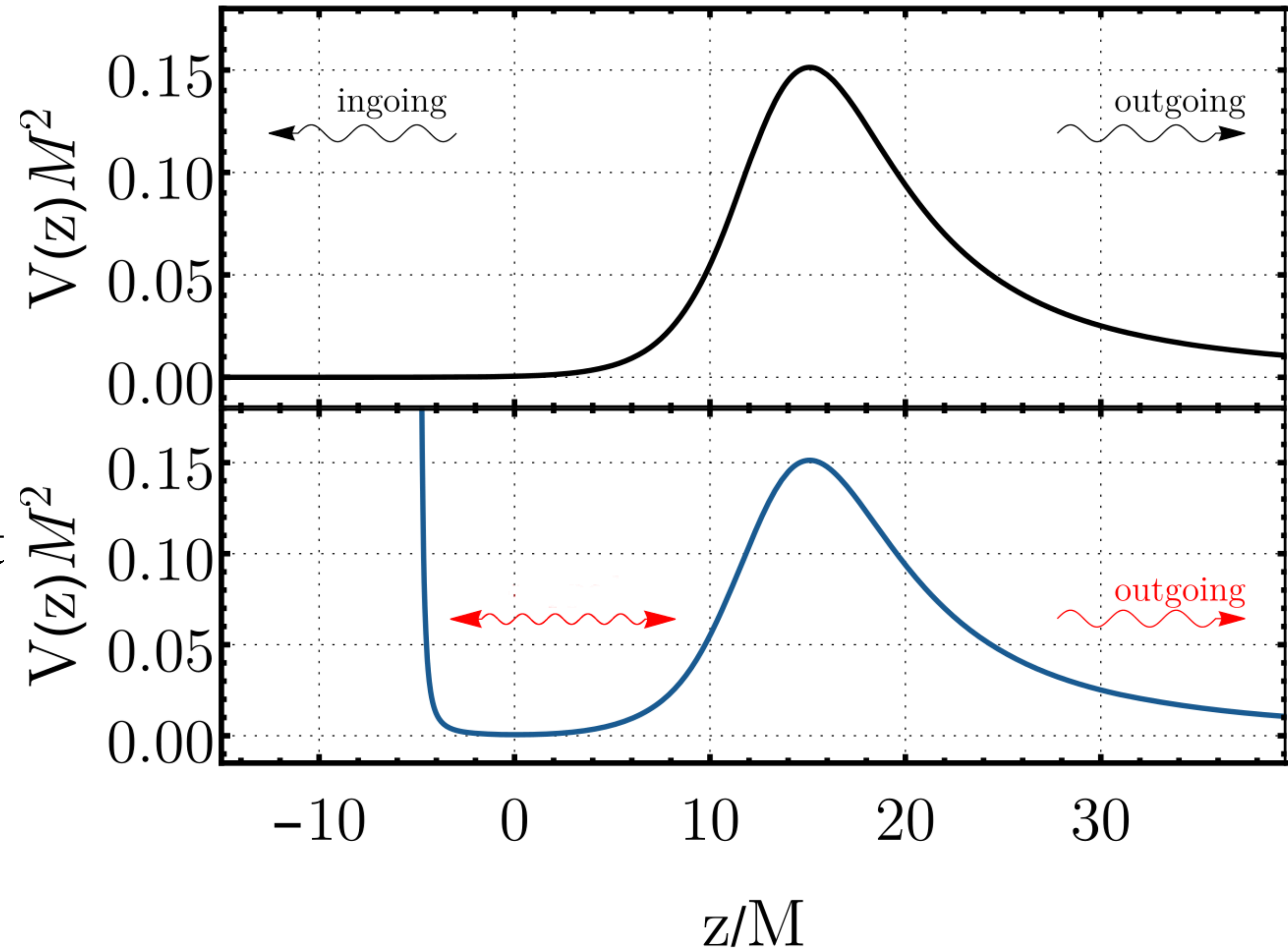
AEI, Hannover



- Current GBDs are being upgraded.
- New detectors **Cosmic Explorer**, **Einstein telescope**, and space based **LISA** is also coming.
- These will be **more sensitive** detectors.
- This opportunity can be used to **test GR**.
- Also the **nature** of the compact objects.
- Exotic compact object (ECO), **quantum effects near BH**.

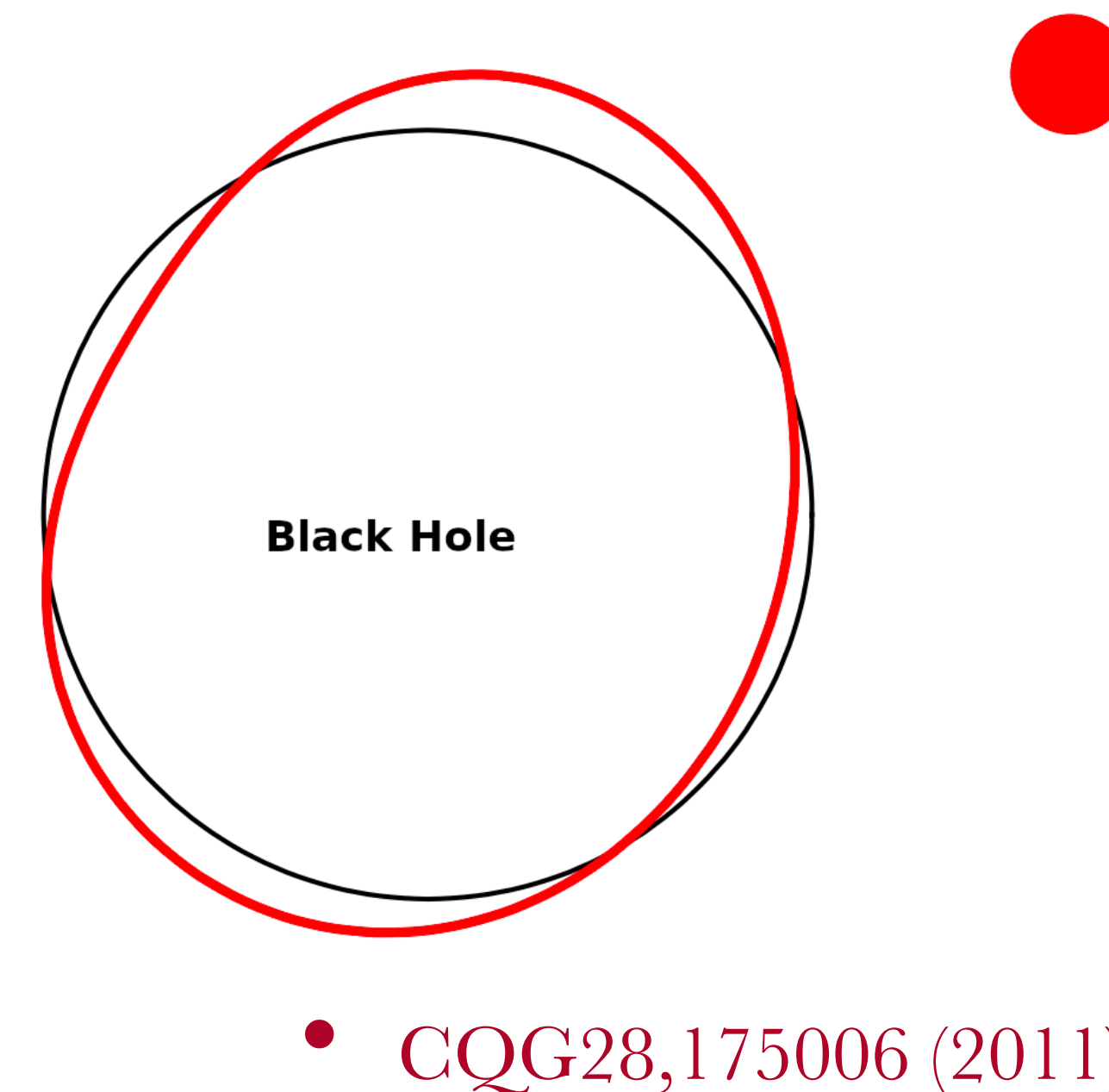


- Classical BH's horizon is perfect absorber due to causality.
- **Absence** (modification) of this implies **imperfect absorption**.
- **Measuring nonzero reflectivity** of compact object surface will be **signature of deviation**.
- Tidal heating is one such effects.



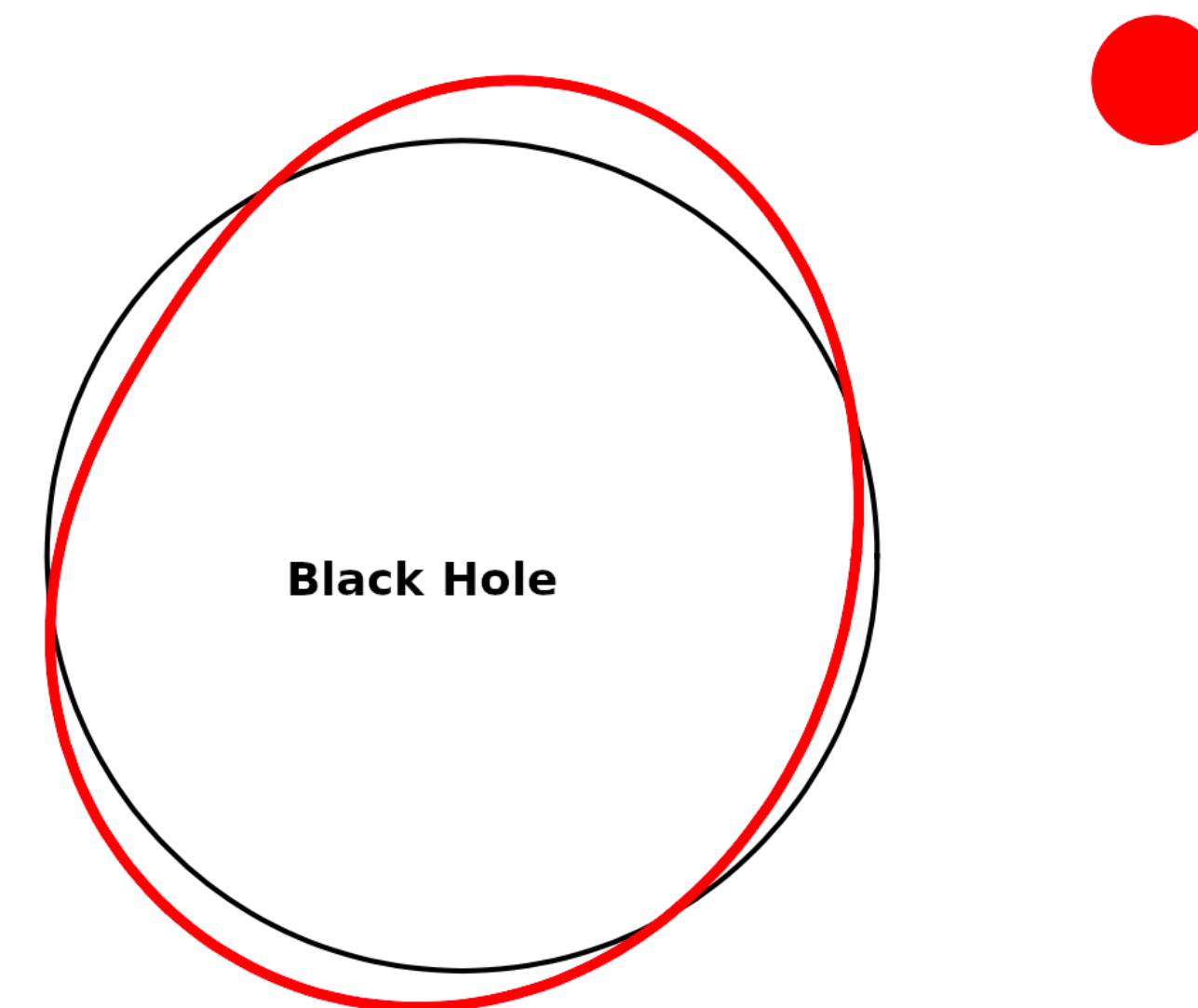
• Living Rev.Rel. 22 (2019) 1, 4

- Components in a binary feel each others' tidal fields (strongly in the late inspiral).
- If the bodies are (at least partially) **absorbing**, these **backreact on the orbit**, exchanging energy and angular momentum with the orbit.
- This effect is called tidal heating [J. B. Hartle, PRD8, 1010 \(1973\)](#),
[S. A. Hughes, PRD64,064004 \(2001\)](#).



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- **In stars** this absorption comes due to **viscous heating** in the material.
- **In BHs** it caused by the **change in the BH mass**.



- *CQG28,175006 (2011)*



- *Tanja Hinderer*

- **Expression for TH** of a star and BH can be brought **into same footing** with viscosity coefficient ($\nu_{BH} \sim M$). **K. Glampedakis+ PRD89,024007(2014)**

- For NS,
$$\nu_{NS} = 10^4 \left(\frac{\rho}{10^{14} \text{ gm cm}^{-3}} \right)^{\frac{5}{4}} \left(\frac{10^8 \text{ K}}{T} \right)^2 \text{ cm}^2 \text{ s}^{-1}$$

- $$\nu_{BH} = 8.6 \times 10^{14} \left(\frac{M}{M_{\odot}} \right) \text{ cm}^2 \text{ s}^{-1}$$

- Even for $M_{BH} \sim M_{NS}$, $\nu_{NS} \ll \nu_{BH}$, resulting in ignorable TH compared to BH.

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- Distinguish BH and NS in this range can change NS mass upperbound and BH mass lower bound.

TH in eccentric orbits

- TH depends on the tidal environment, $\mathcal{E}_{ab}(r, \dot{r}, \phi, \dot{\phi})$ and $\mathcal{B}_{ab}(r, \dot{r}, \phi, \dot{\phi})$. Taylor+ PRD 78, 084016 (2008)
- Orbital quantities depend on, $\xi(f, e_t), u, e_t$. Moore+ PRD93, 124061 (2016)
- $\dot{m} = \dot{m}(\bar{\mathcal{E}}_{ab}, \bar{\mathcal{B}}_{ab})$ and similarly \dot{J} . Taylor+ PRD 78, 084016 (2008), Poisson, PRD70, 084044 (2004)
- Upon averaging over the orbit,
 $\langle \dot{m} \rangle = \dot{m}_{circ} \mathcal{M}(e_t), \langle \dot{m} \rangle_\chi = \dot{m}_{circ, \chi} \mathcal{M}_\chi(e_t), \langle \dot{J} \rangle_\chi = \dot{J}_{circ, \chi} \mathcal{J}_\chi(e_t)$
SD, 2305.03771

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 SD, 2305.03771, Munna+ 2306.12481, Jezereel talk

- During inspiral ξ , e_t evolves with frequency.
- Inspiral gives, $\frac{de_t}{dt}$ and $\frac{dx}{dt}$, $x \sim f^{2/3}$. Arun+ PRD80, 124018, Klein+ PRD81, 124001
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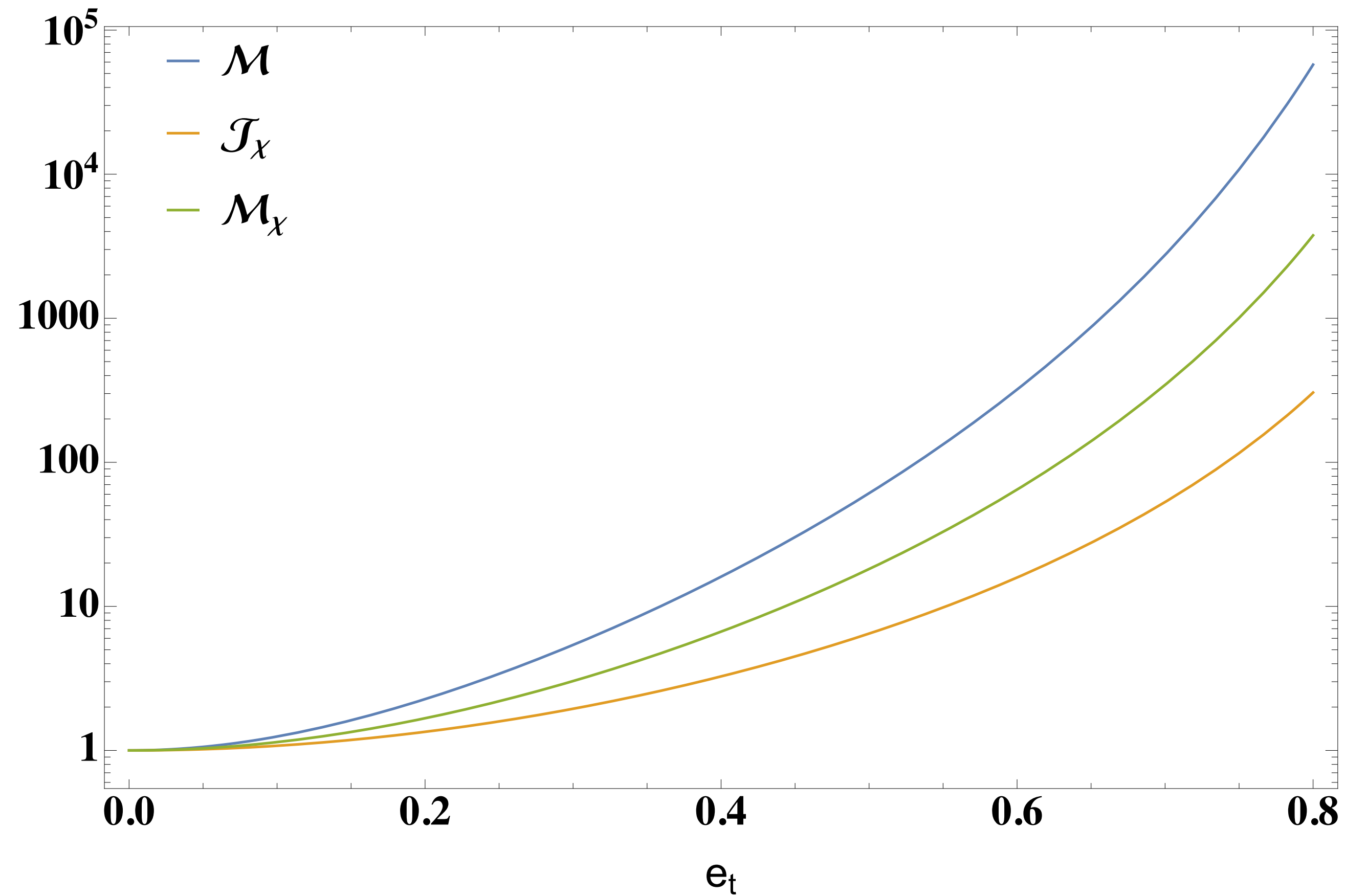
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- $e_t = e_0 \epsilon_1(x, x_0) + e_0^3 \epsilon_3(x, x_0) + \mathcal{O}(e_0^5)$. [SD, 2305.03771](#), [SD 2306.12522](#)
- $\langle \dot{m} \rangle = \mathcal{M}_0 + e_0^2 \mathcal{M}_2 + e_0^4 \mathcal{M}_4 + \mathcal{O}(e_0^6)$. [SD, 2305.03771](#)

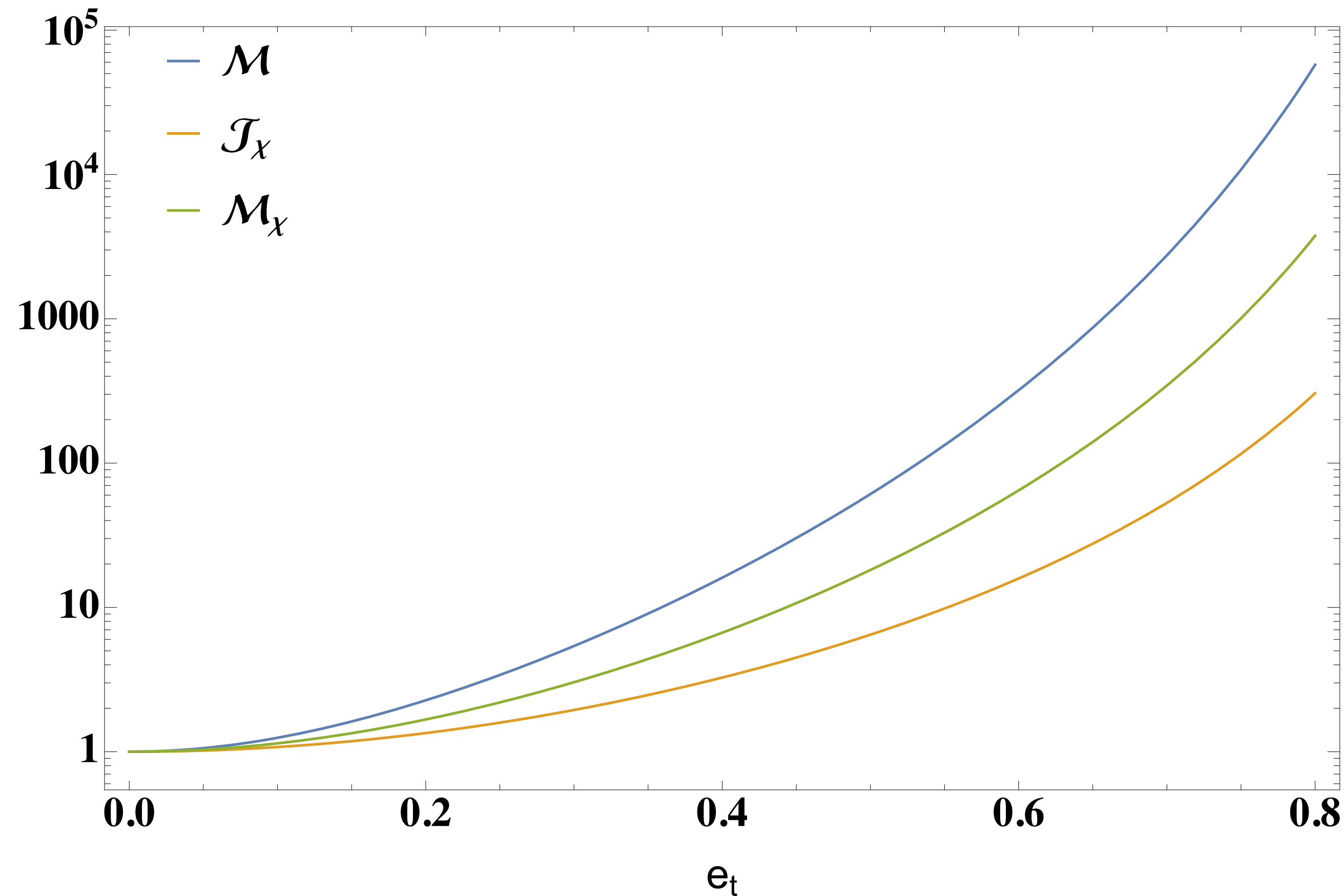
- $\langle \dot{J}_1 \rangle = - \frac{8m_1^5 m_2^2 \xi^4 \chi (3\chi^2 + 1) (3e_t^4 + 24e_t^2 + 8)}{5M^6} \frac{8(1 - e_t^2)^{9/2}}{16(e_t^2 - 1)^6}$

- $\langle \dot{m}_1 \rangle = - \frac{8\epsilon m_1^5 m_2^2 \xi^5 \chi (3\chi^2 + 1) (5e_t^6 + 90e_t^4 + 120e_t^2 + 16)}{5M^7} \frac{8(1 - e_t^2)^{9/2}}{16(e_t^2 - 1)^6}$



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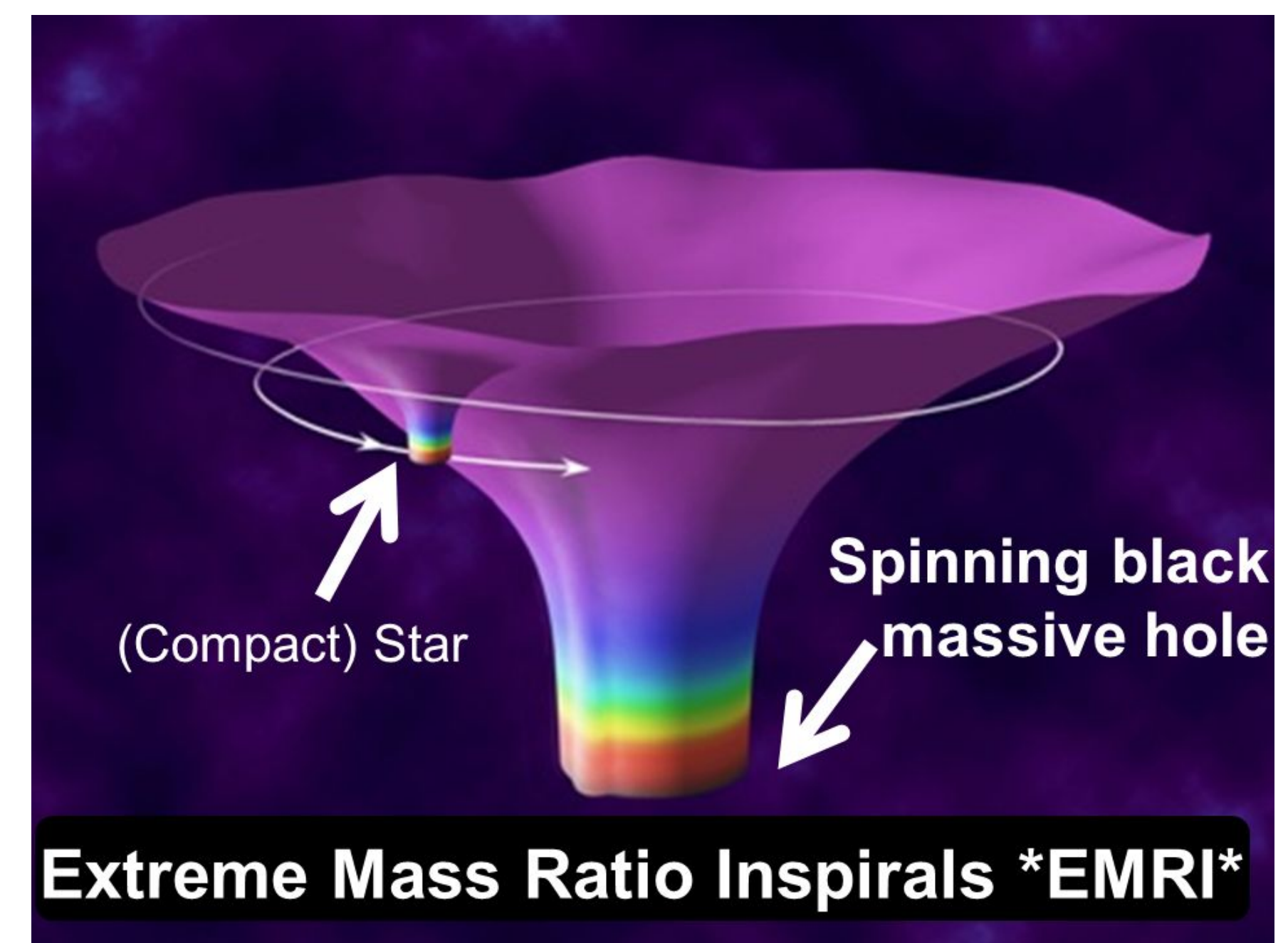
$$F^H \propto \Omega(\Omega - \Omega_H)$$

$$\Omega \sim \nu^3$$

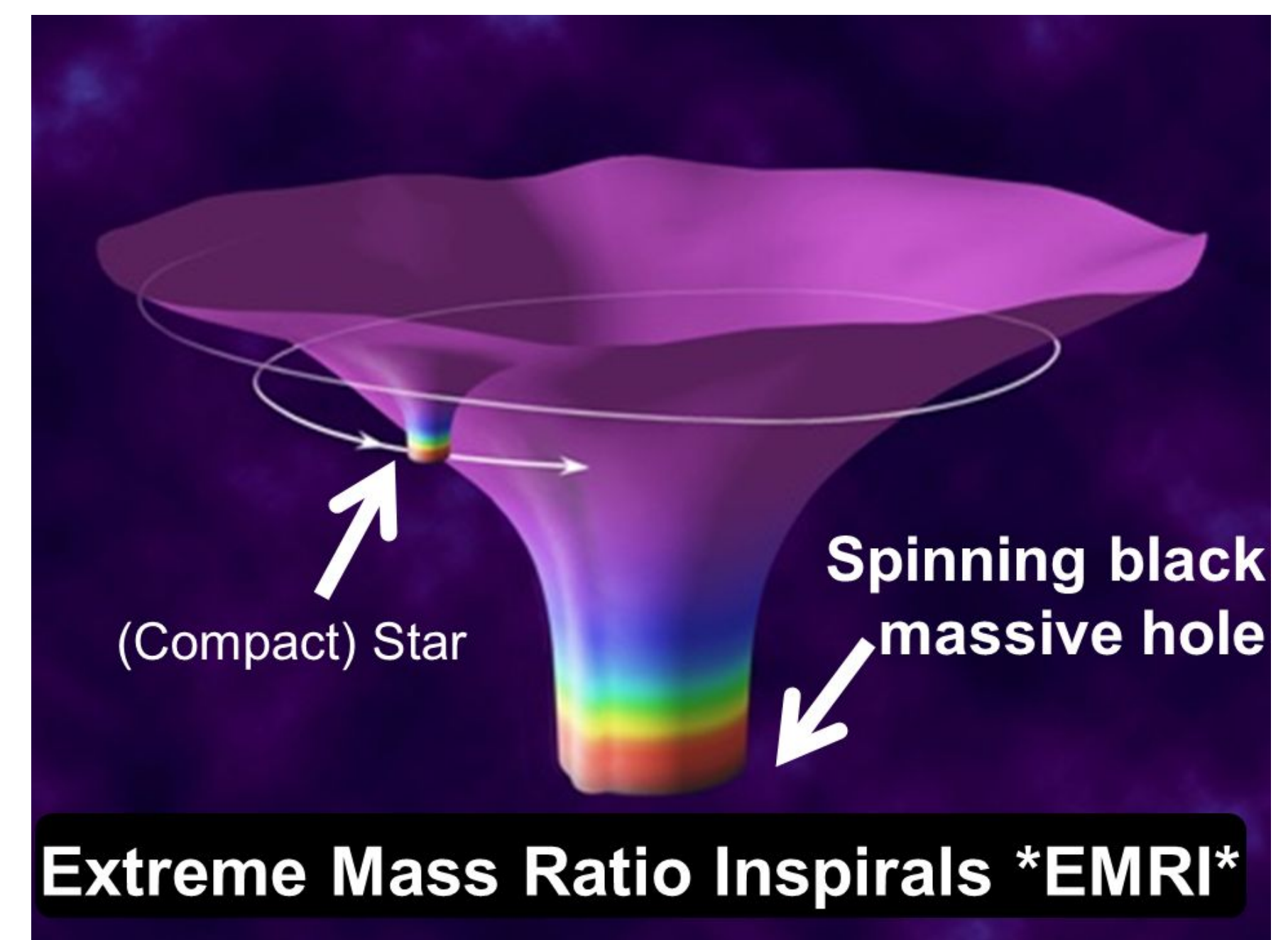
Strength in EMRI

SD, Scott Hughes, Richard Brito, Paolo Pani

- We will focus on EMRI, where a stellar mass $\sim 10 - 100M_{\odot}$ Inspirals around SMBH of $\sim 10^5 - 10^7M_{\odot}$, observable in LISA.
- Hence we calculate perturbation around Kerr BH by a small particle.



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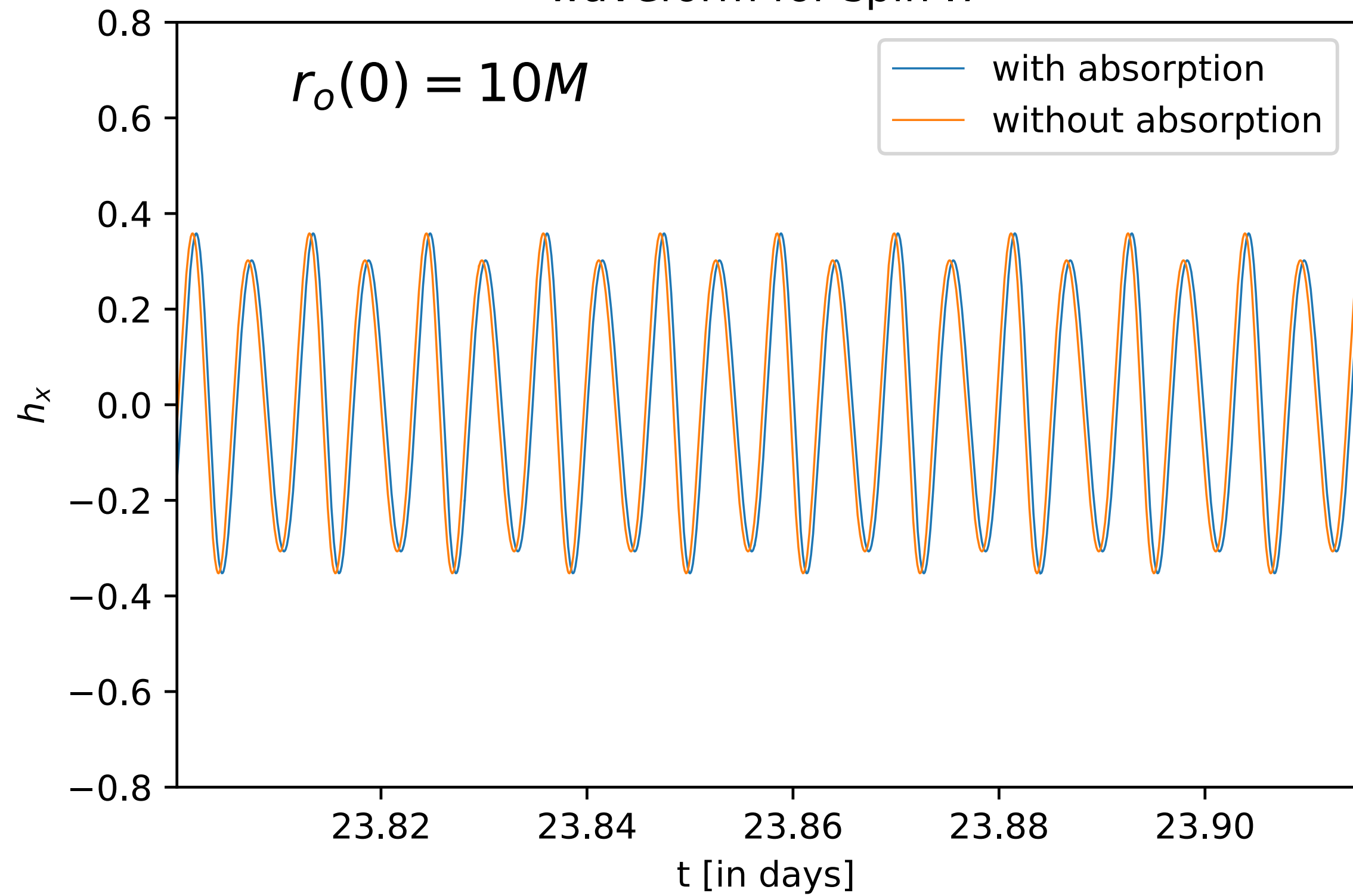
- We solve BHP equation.
- Energy fluxes at infinity and the horizon can be calculated from the perturbation GW waveform.

- $\dot{E} = \dot{E}_{\infty} + \dot{E}_H$

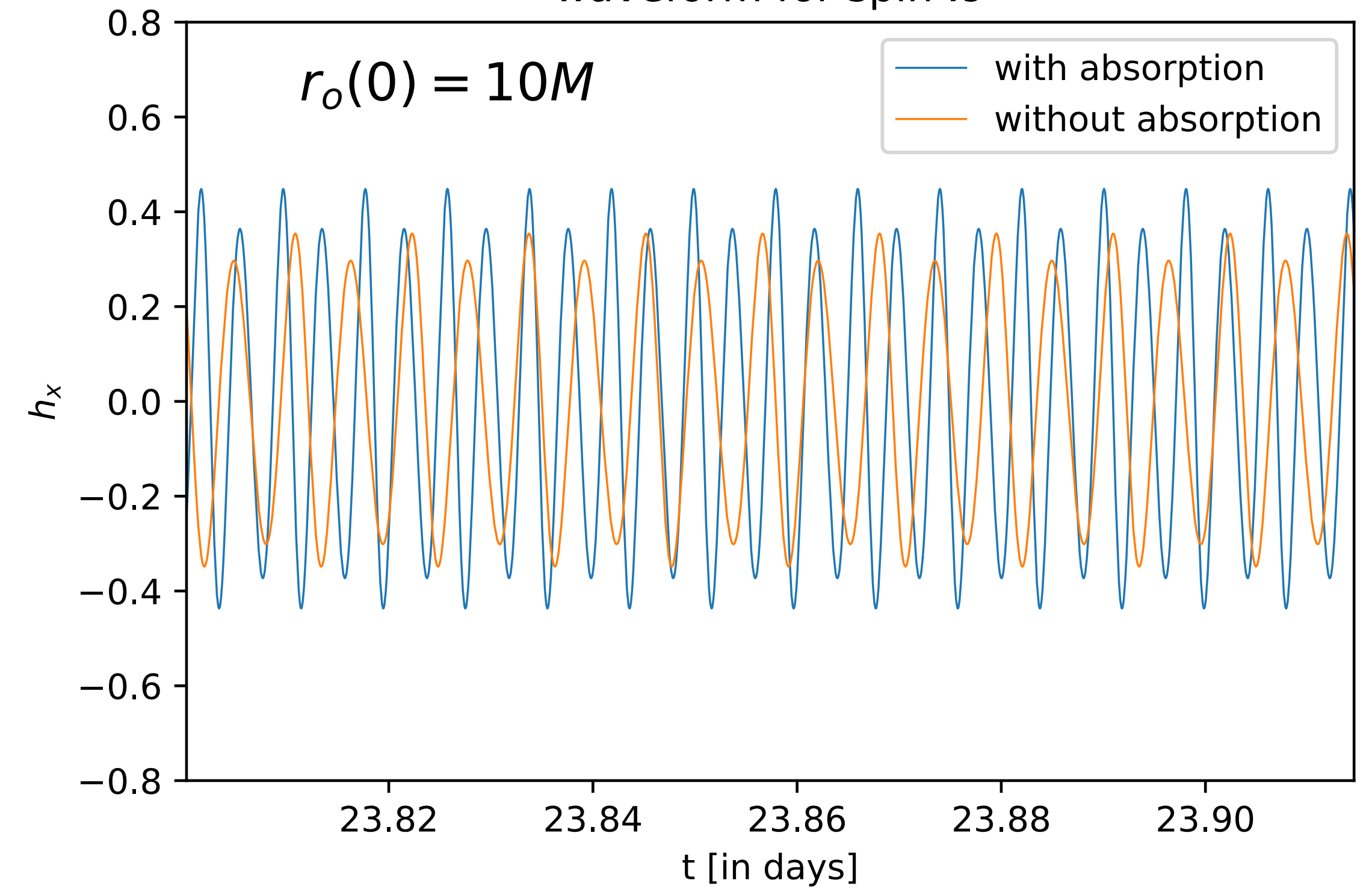
- $\dot{E}_{TH} = (1 - |\mathcal{R}|^2)\dot{E}_H$

- $M = 10^6M_{\odot}, M = 30000\mu$

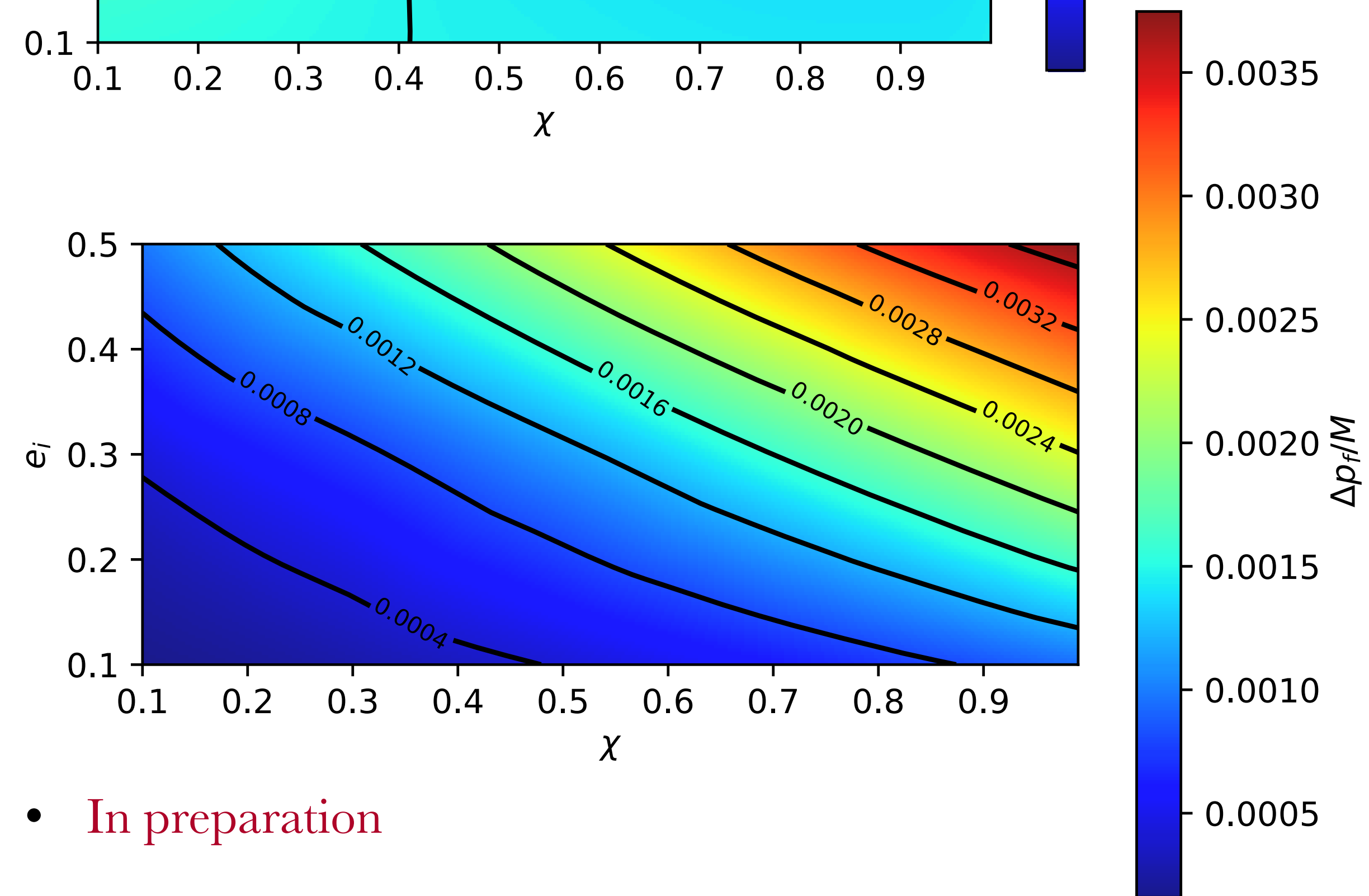
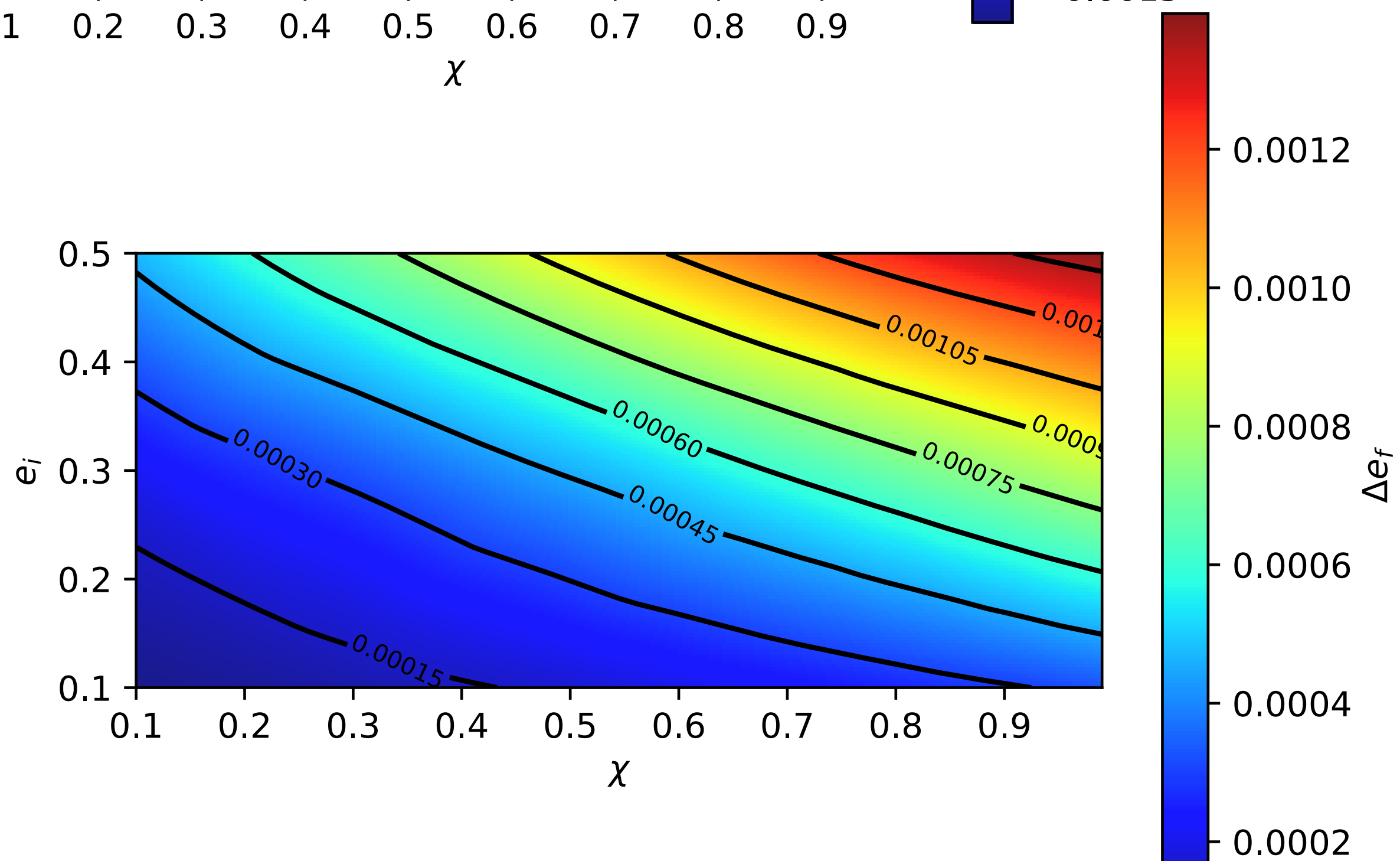
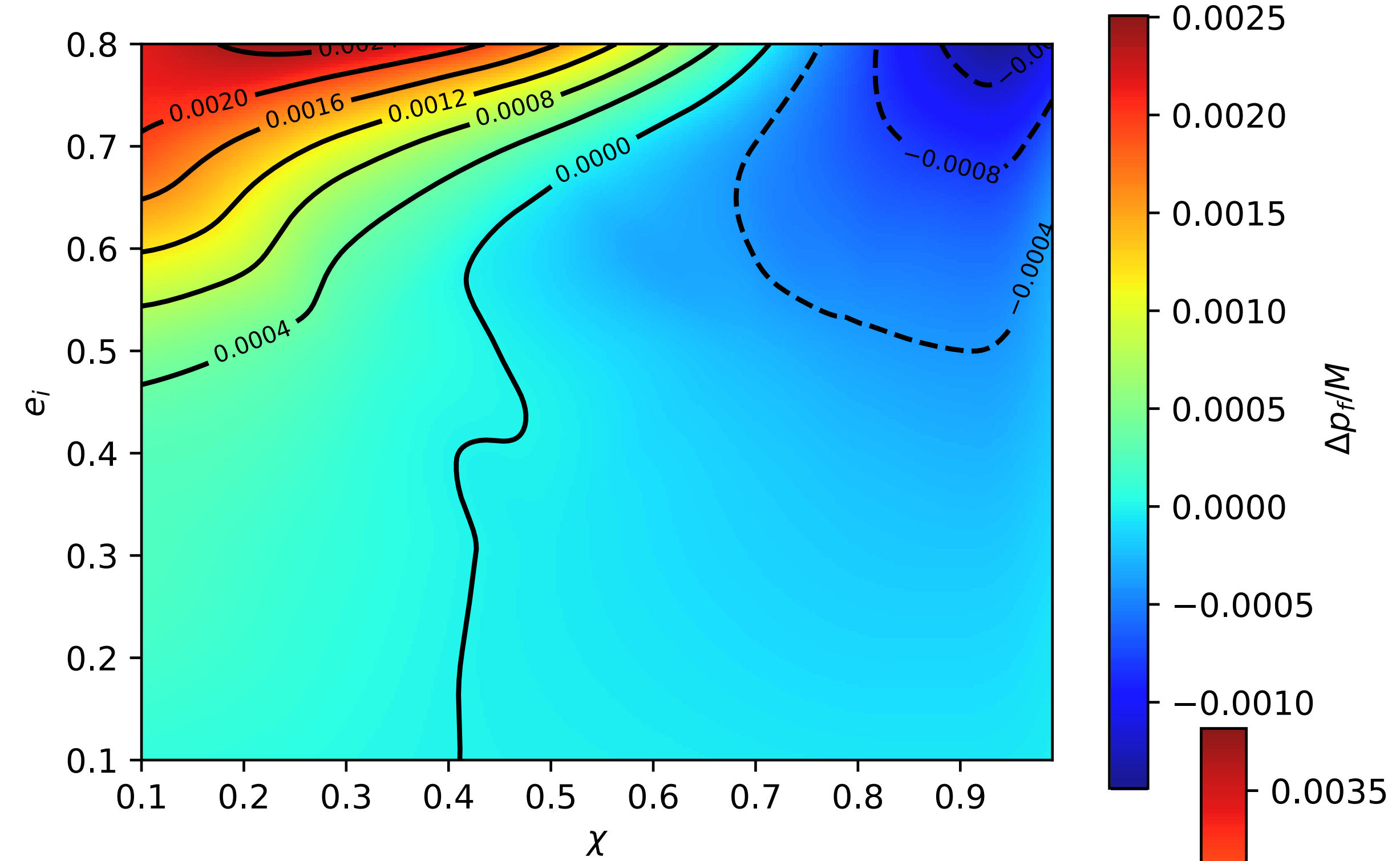
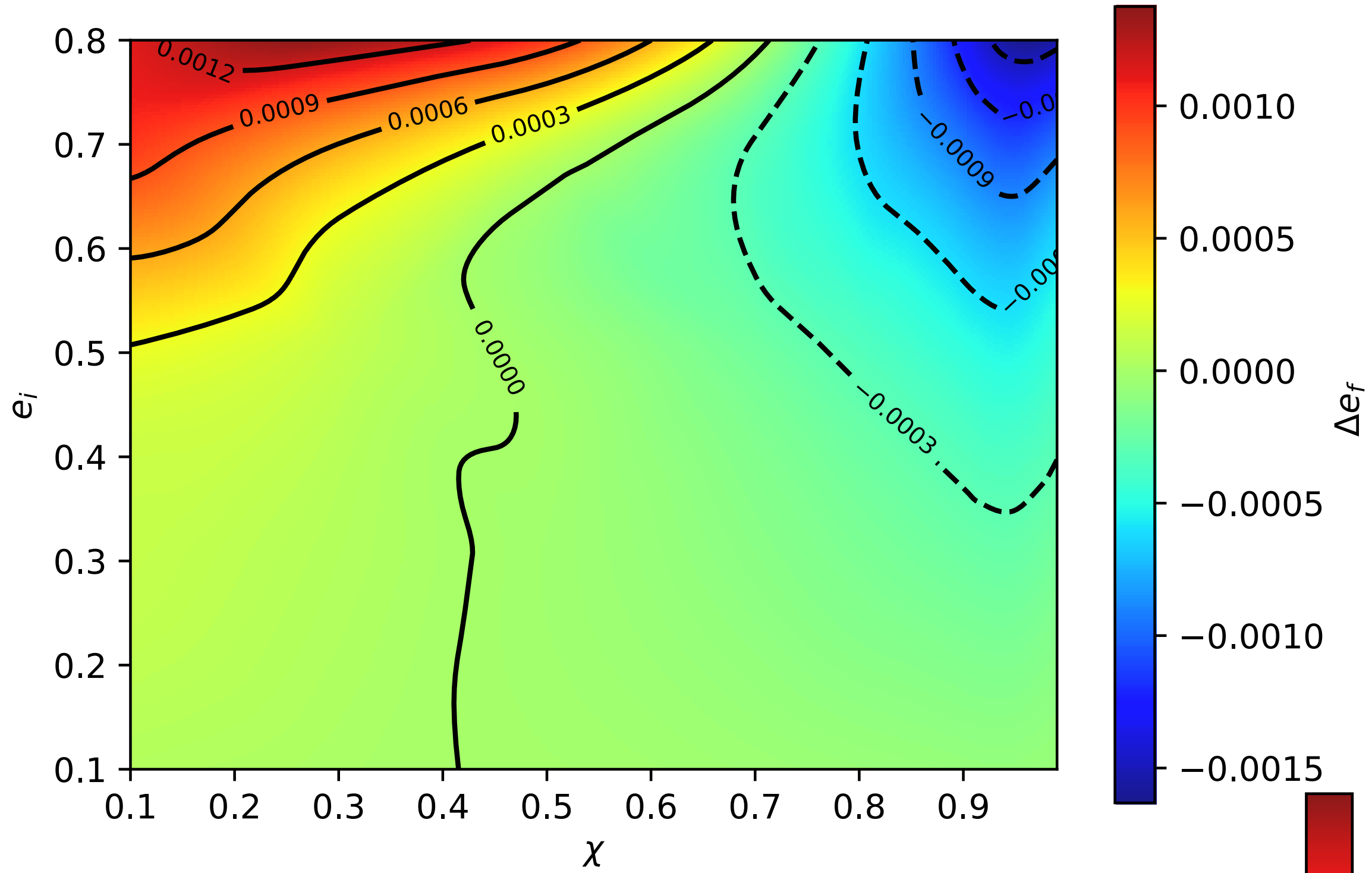
waveform for spin .7



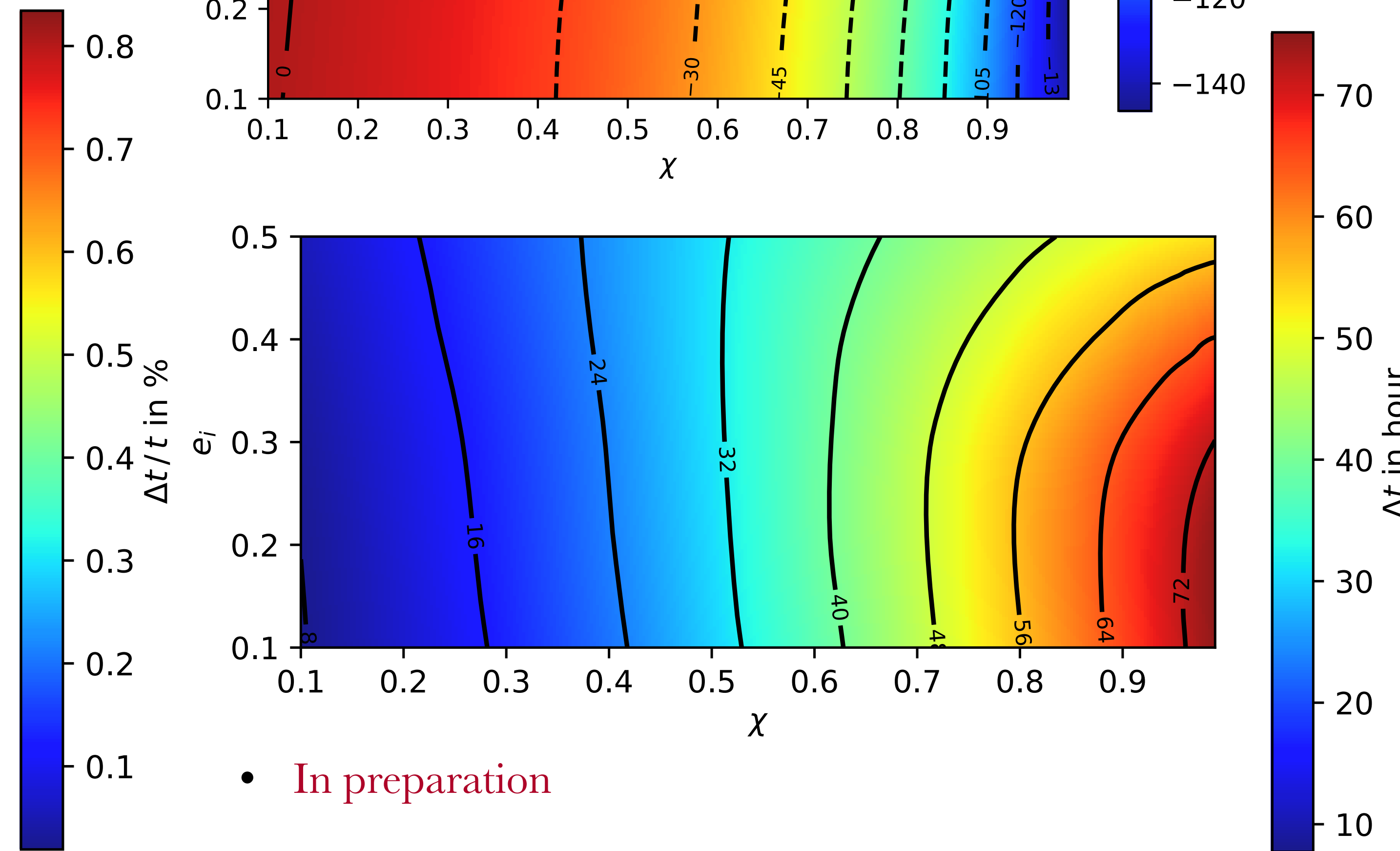
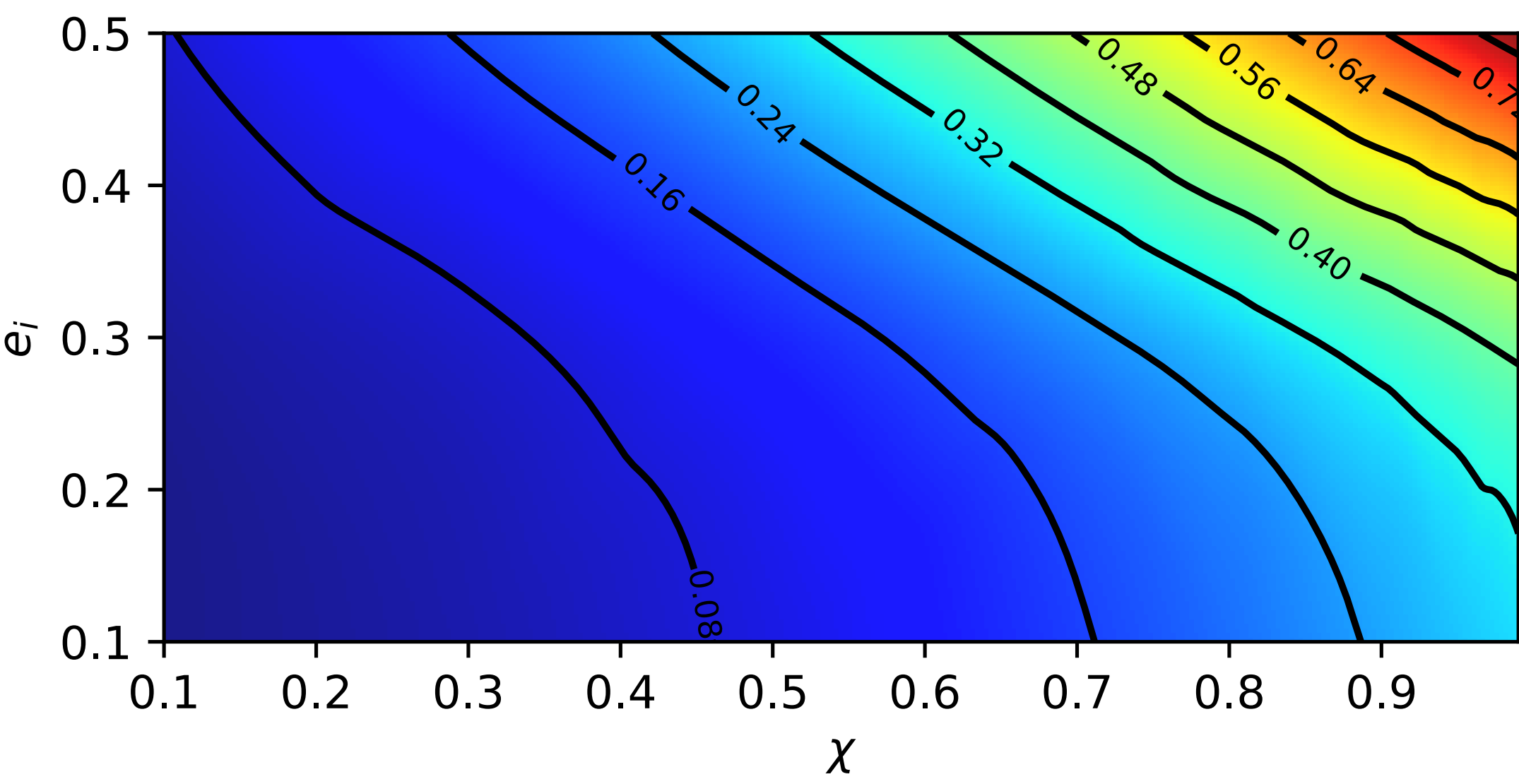
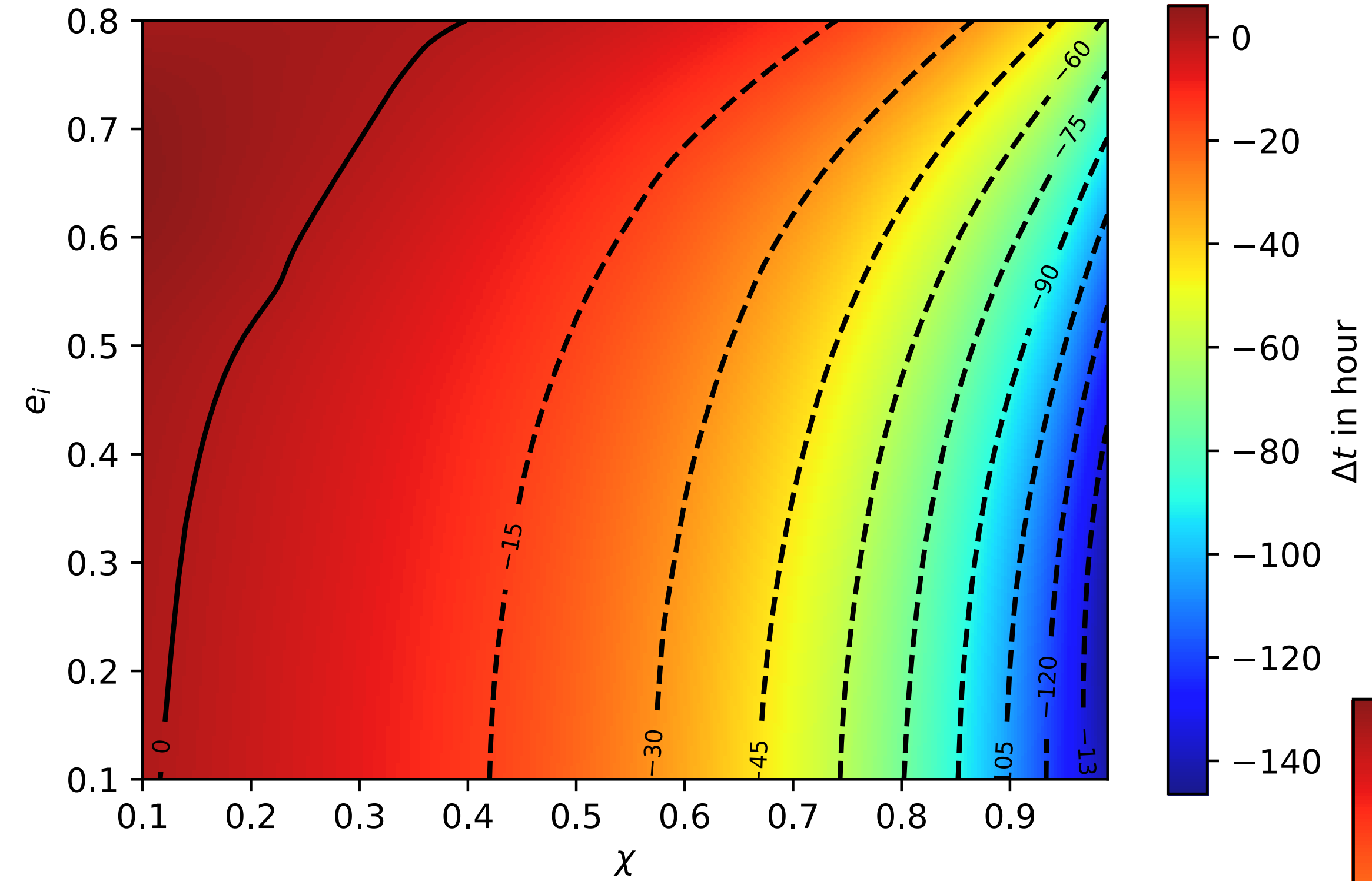
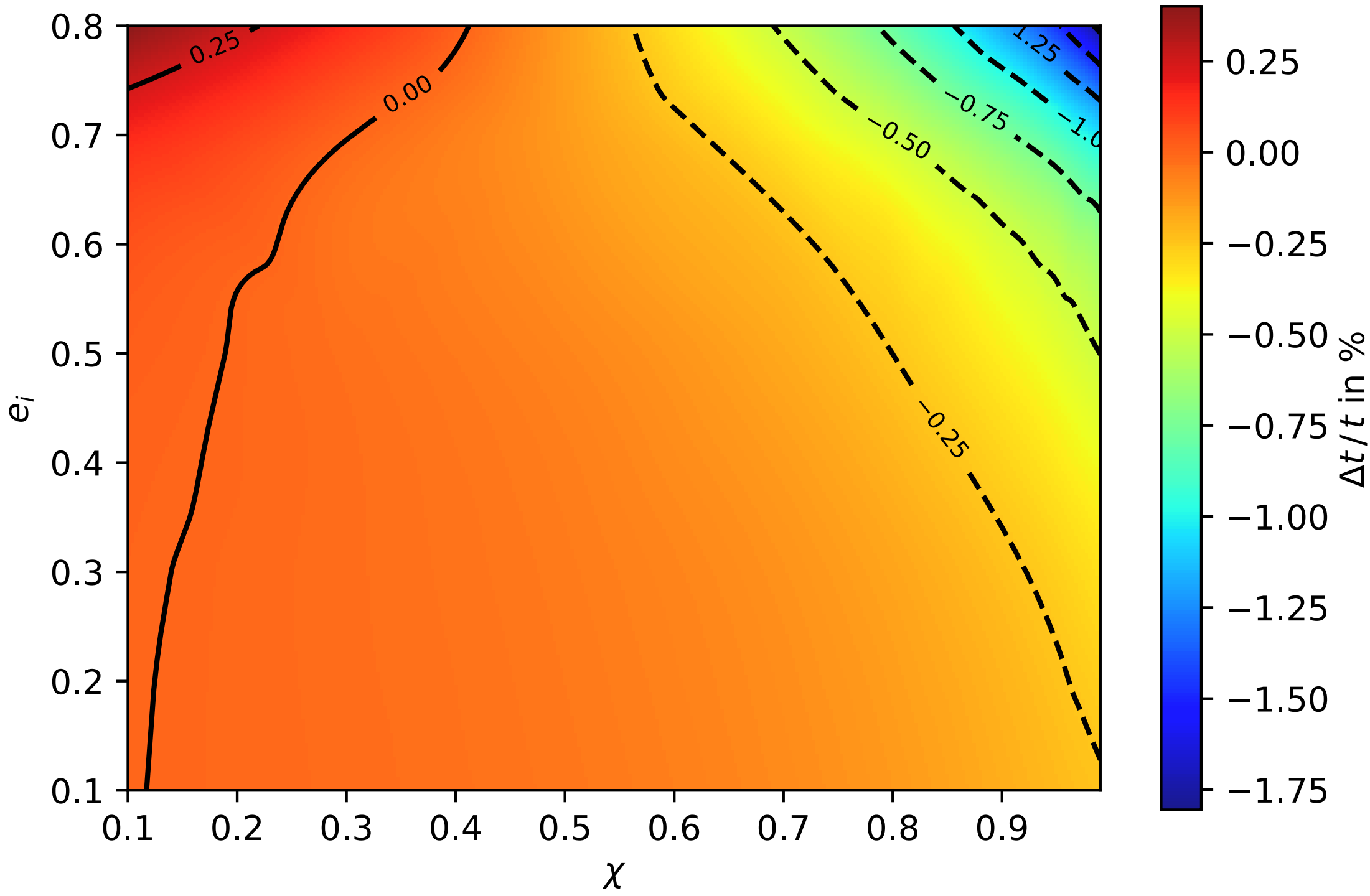
waveform for spin .9



- Inspiral of binary is driven by **energy loss at infinity** and **at horizon**.
- **Switching off** energy exchange due to **heating** modifies inspiral rate, resulting in **change in GW**.
- Calculate GW with and without TH and calculate dephasing [PRD.101.044004 SD, Brito, Bose, Pani, Hughes](#).



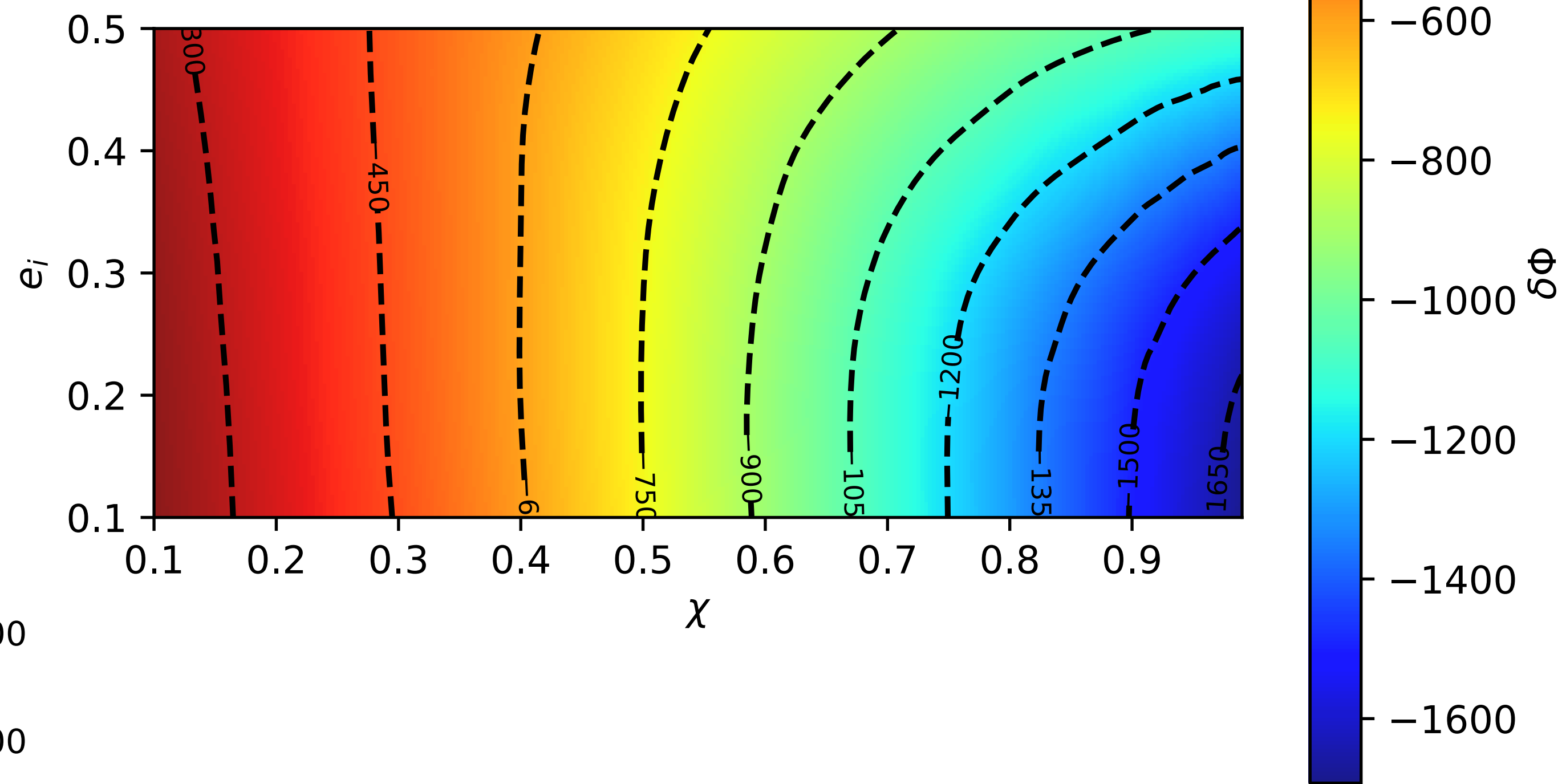
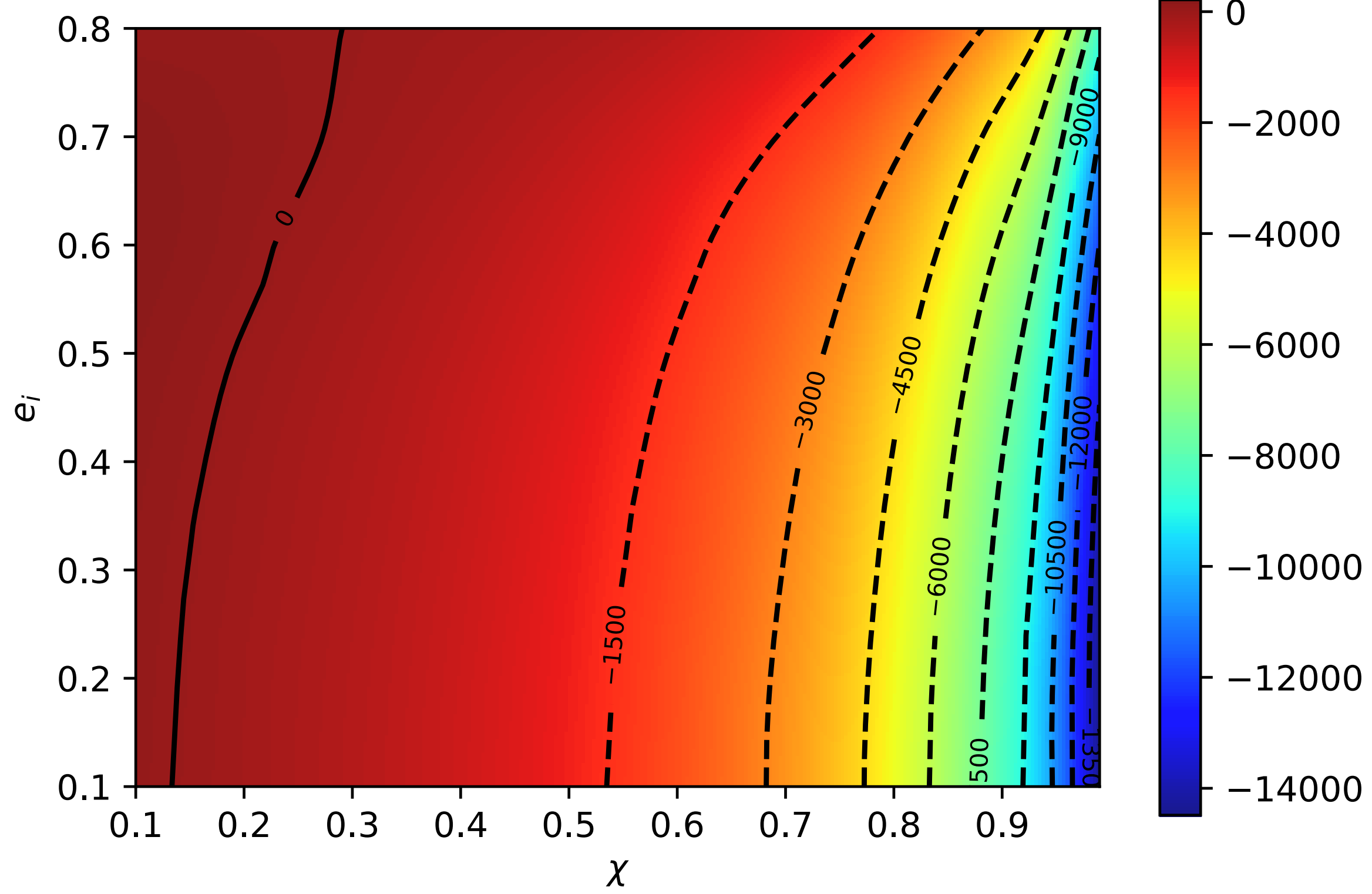
• In preparation



• In preparation

- $\delta\Phi_{m,n} = m\delta\Phi_\phi + n\delta\Phi_r$

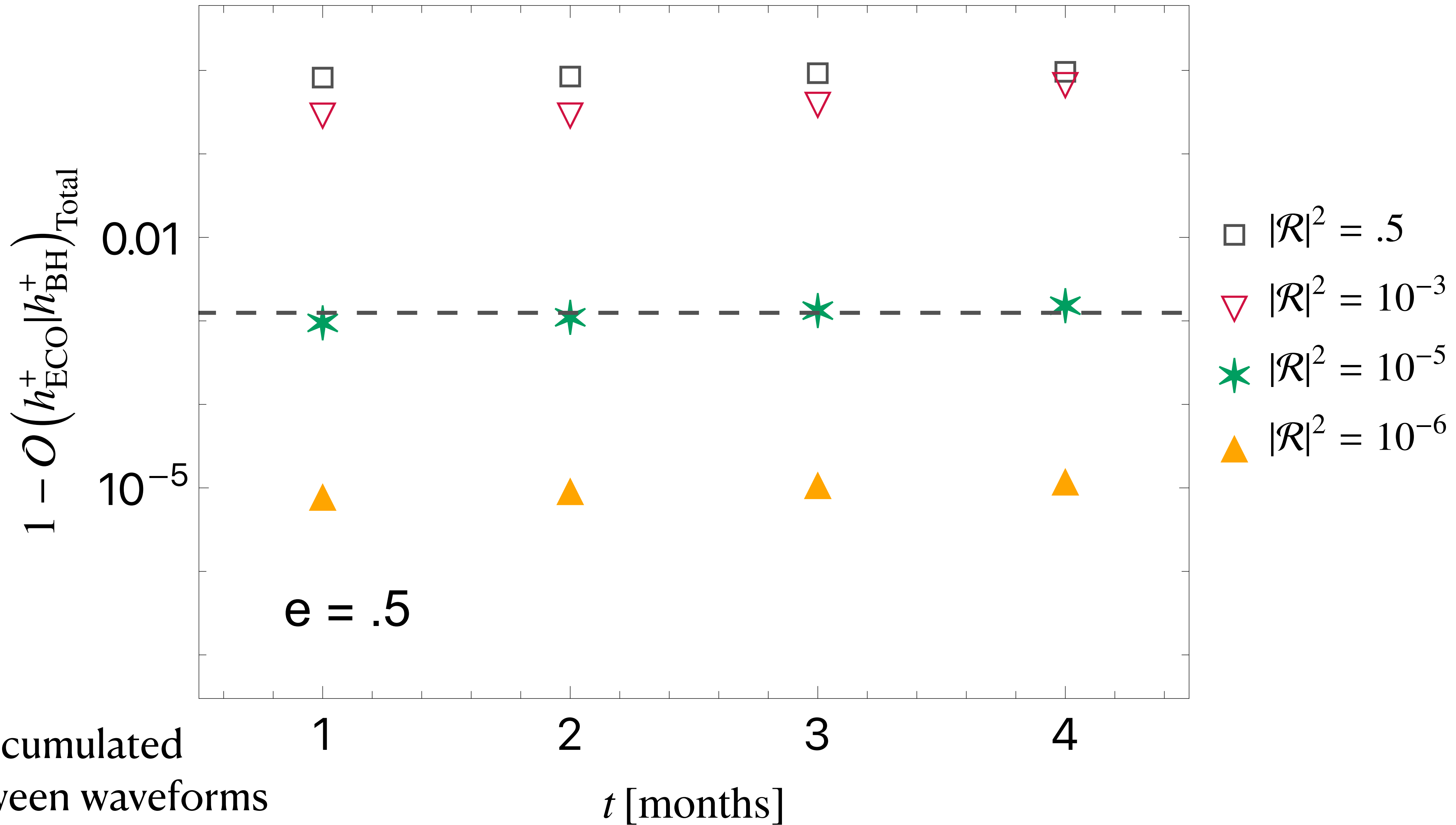
- $\delta\Phi = 2\delta\Phi_\phi + \delta\Phi_r$



- In preparation

Usefulness of TH

- $\dot{E}_{\text{ECO}} = (1 - |\mathcal{R}|^2)\dot{E}_H + \mathcal{O}(\epsilon)$ SD, PRD.102.064040 , Maggio+ PRD104 (2021) 10, 104026
- \mathcal{R} is the reflectivity of the ECO (QBH).
- position of the reflective surface $r_s = r_H(1 + \epsilon)$.
- Measuring $|\mathcal{R}|^2$ tests the "Blackness" of the hole and ϵ tests the horizon position. SD, S. Bose, PRD99,084001 (2019), Maselli+, PRL120,081101(2018)



- Image is the accumulated mismatch between waveforms with $|\mathcal{R}|^2 \neq 0$ and BH. with time. In preparation

Take Home

- TH in eccentric case computed analytically.
- Using eccentricity evolution frequency evolution is computed.
- In EMRI, final orbital quantities can change due to **modified TH**, implying **changed inspiral**.
- TH induces **large dephasing**.
- In EMRI TH can lead to significant dephasing, resulting in constraining $|\mathcal{R}|^2 \sim 10^{-5} - 10^{-6}$.