Detection and modelling of eccentric intermediate mass ratio in-spirals

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With

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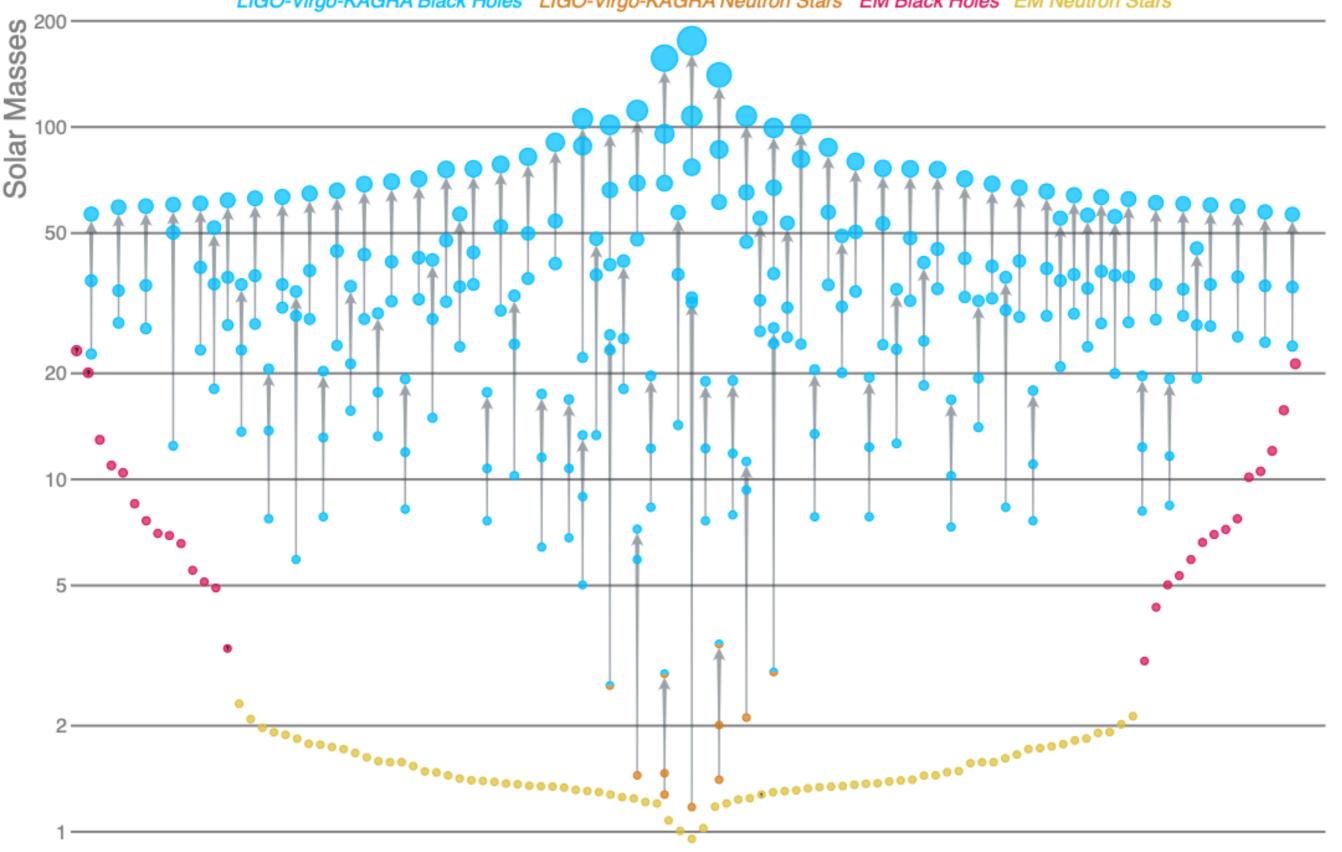




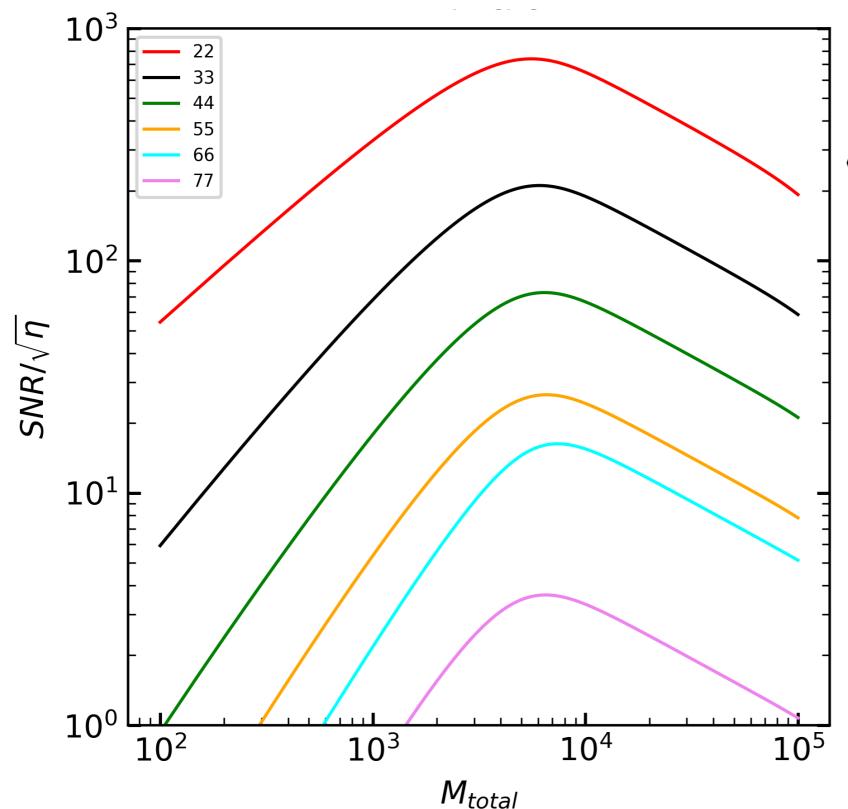


Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes LIGO-Virgo-KAGRA Neutron Stars EM Black Holes EM Neutron Stars



Detectability of a circular binary with at least one IMBH merging in DECIGO-B band



$$\langle \rho^2 \rangle = \frac{3}{20} \sum_{n} \int_{f_{\text{low}}}^{f_{\text{up}}} \frac{h_{c,n}^2(f_n)}{f_n^2 S_h(f_n)} df$$

$$h_{c,n} = (\pi d)^{-1} \sqrt{2\dot{E}_n/\dot{f}_n}$$

Scaling

$$1/\sqrt{\eta}$$

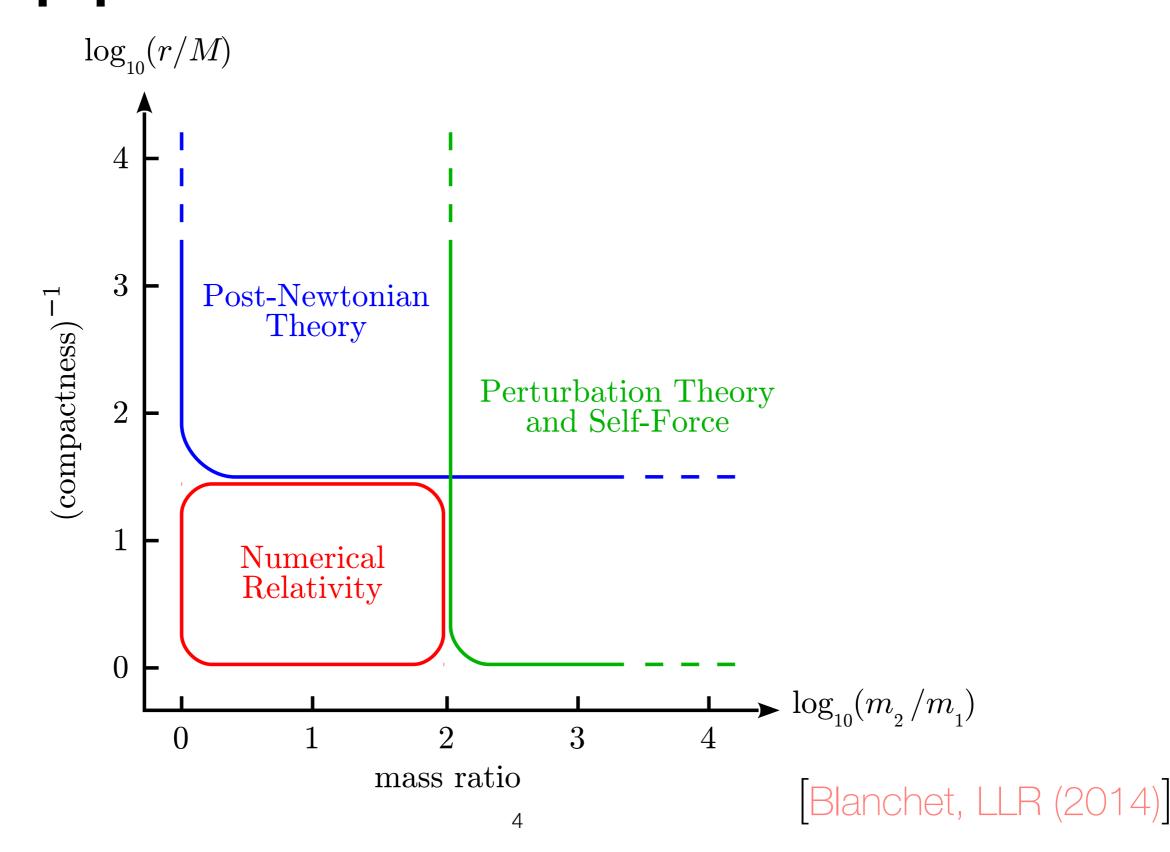
$$= 2 (\eta = 0.25)$$

$$\sim 3 (\eta = 0.1)$$

$$\sim 10 (\eta = 0.01)$$

in prep.

Approximation schemes



Outline

- PN and BHP solutions for ICBs in circular orbits
- Eccentric corrections and hybridisation
- Model systematics: impact of mass-ratio,
 eccentricity and higher order BHP terms

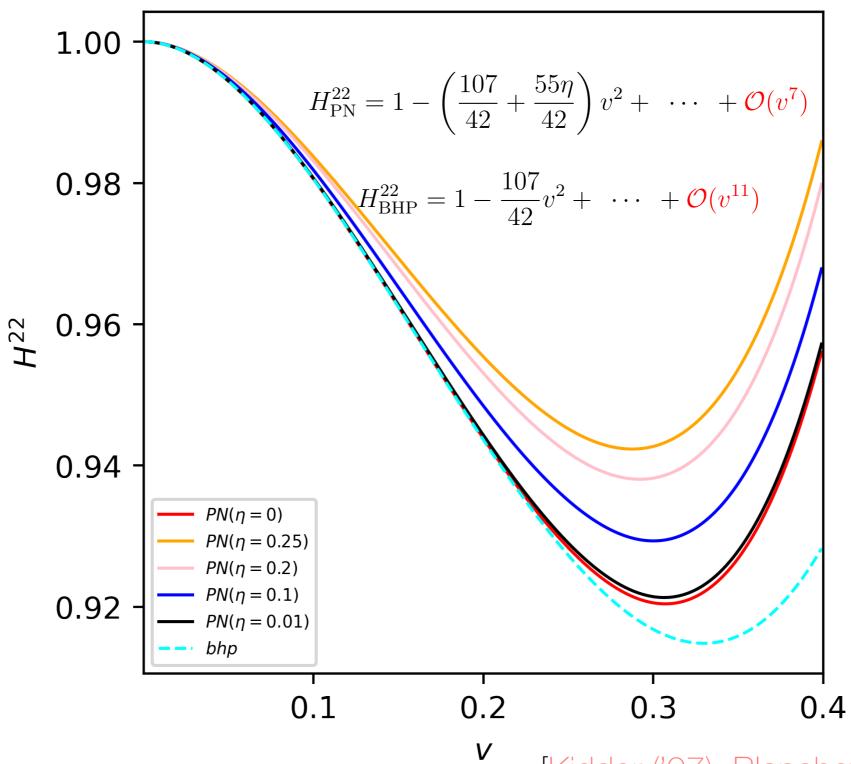
Approximation schemes

- Taylor series expansions in small parameters such as internal source velocities
- Point particles, multipolar expansions etc.
- High order computations (nontrivial due to IR divergences)

Detector	LIGO/Virgo		ET		LISA		
Masses (M_{\odot})	1.4×1.4	10 × 10	1.4×1.4	500×500	$10^5 \times 10^5$	$10^7 \times 10^7$	
PN order	cumulative number of cycles						
Newtonian	2 562.599	95.502	744 401.36	37.90	28 095.39	9.534	
1PN	143.453	17.879	4 433.85	9.60	618.31	3.386	
1.5PN	-94.817	-20.797	-1005.78	-12.63	-265.70	-5.181	
2PN	5.811	2.124	23.94	1.44	11.35	0.677	
2.5PN	-8.105	-4.604	-17.01	-3.42	-12.47	-1.821	
3PN	1.858	1.731	2.69	1.43	2.59	0.876	
3.5PN	-0.627	-0.689	-0.93	-0.59	-0.91	-0.383	
4PN	-0.107	-0.064	-0.12	-0.04	-0.12	-0.013	
4.5PN	0.098	0.118	0.14	0.10	0.14	0.065	

[Order counting terminology: $nPN := V^{2n}$]

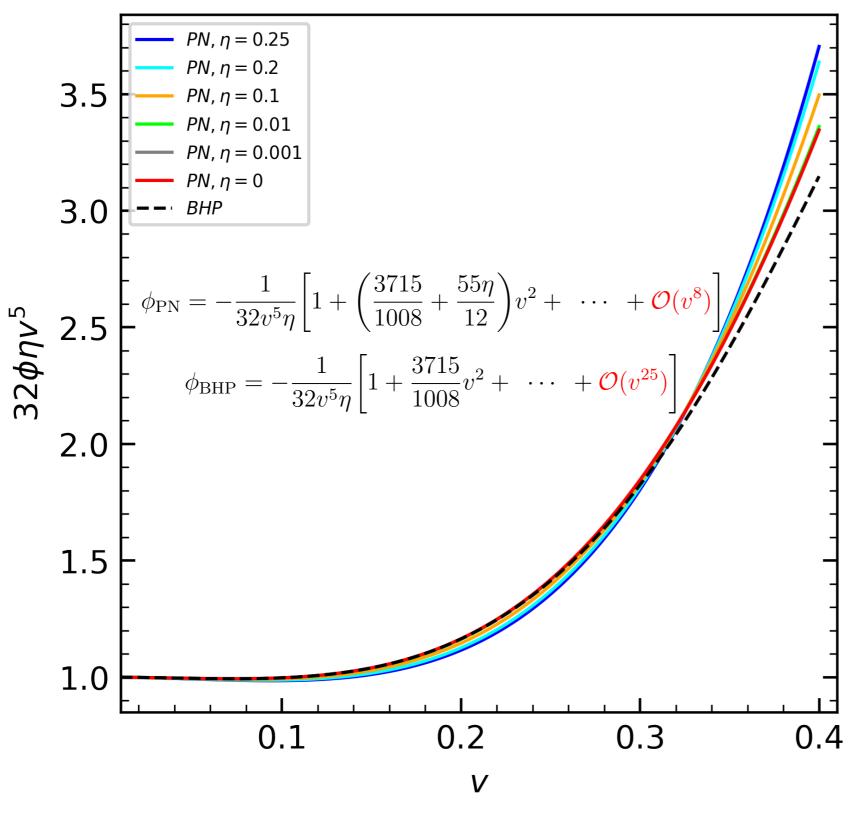
PN-BHP comparisons



 $O(V^{2n})$: terms beyond nPN and ignored]

[Kidder ('07), Blanchet + ('08)] [Fujita-lyer ('10)]

PN-BHP comparisons



Scaling

$$1/(32\eta v^5)|_{v=0.4}$$
 $\sim 10 \, (\eta = 0.25)$
 $\sim 30 \, (\eta = 0.1)$
 $\sim 300 \, (\eta = 0.01)$

[Blanchet + ('02)] [Fujita ('12)]

Eccentric Case

PN-Phase

$$\phi_{\text{PN}} = -\frac{1}{32v^5\eta} \left\{ 1 + \left(\frac{3715}{1008} + \frac{55\eta}{12} \right) v^2 + \dots + \mathcal{O}(v^8) \right.$$

$$\left. - \frac{785}{272} e_0^2 \left(\frac{v_0}{v} \right)^{19/3} \left[1 + \left(\frac{6955261}{2215584} + \frac{436441\eta}{79128} \right) v^2 + \left(\frac{2833}{1008} - \frac{197\eta}{36} \right) v_0^2 + \dots + \mathcal{O}(v^7) \right] \right.$$

$$\left. + \mathcal{O}(e_0^4) \right\}$$

PN-BHP-Phase

$$\begin{split} \phi_{\text{PN-BHP}} &= -\frac{1}{32v^5\eta} \Bigg\{ 1 + \left(\frac{3715}{1008} + \frac{55\eta}{12} \right) v^2 + \cdots \right. \\ &+ \frac{2554404624135128353}{2214468081745920} v^8 + \cdots \right. \\ &+ \mathcal{O}(v^{25}) \\ &- \frac{785}{272} e_0^2 \left(\frac{v_0}{v} \right)^{19/3} \Bigg[1 + \left(\frac{6955261}{2215584} + \frac{436441\eta}{79128} \right) v^2 + \left(\frac{2833}{1008} - \frac{197\eta}{36} \right) v_0^2 + \cdots \\ &+ (\cdots) v^8 + (\cdots) v^7 v_0 + \cdots \right. \\ &+ (\cdots) v^7 v_0 + \cdots \right. \\ &+ \left. (\cdots) v^7 v_0 + \cdots \right. \\ \\ &+ \left. (\cdots) v^7 v_0 + \cdots \right. \\$$

Impact of mass ratio on # GW cycles: PN contribution

- System: M=1000 Msun, q=200
- \blacksquare DECIGO-B: f_{low} =0.01Hz, $f_{up} = f_{lso}$
- PN phase model: 3.5PN (circular) 3PN (Eccentric)
- BHP phase model: 12PN (circular) & 5PN (Eccentric)

PN Order ↓	$(\Delta N_{cycle}^{PN(e_0^{10})} - \Delta N_{cycle}^{BHP(e_0^{10})})$	$(\Delta N_{cycle}^{PN(e_0^{10})}$ - $\Delta N_{cycle}^{BHP(e_0^{10})})$	$(\Delta N_{cycle}^{PN(e_0^{10})} - \Delta N_{cycle}^{BHP(e_0^{10})})$	$(\Delta N_{cycle}^{PN(e_0^{10})} - \Delta N_{cycle}^{BHP(e_0^{10})})$
$e_0 \rightarrow$	$e_0 = 0$	$e_0 = 0.1$	$e_0 = 0.2$	$e_0 = 0.3$
0 PN	0	0	0	0
1 PN	293.963	$-23316.496 \times 10^{-5}$	-2.479	-9.671
1.5 PN	0	0	0	0
2 PN	4.351	-5023.803×10^{-5}	-0.223	-0.541
2.5 PN	0.515	-229.227×10^{-5}	-1028.718×10^{-5}	-2400.073×10^{-5}
3 PN	2.917	-279.143×10^{-5}	-1105.646×10^{-5}	-2393.290×10^{-5}
3.5 PN	-0.2	-	-	-
Total cycles accumulated	301.746	-0.288	-2.724	-10.260

[In prep.]

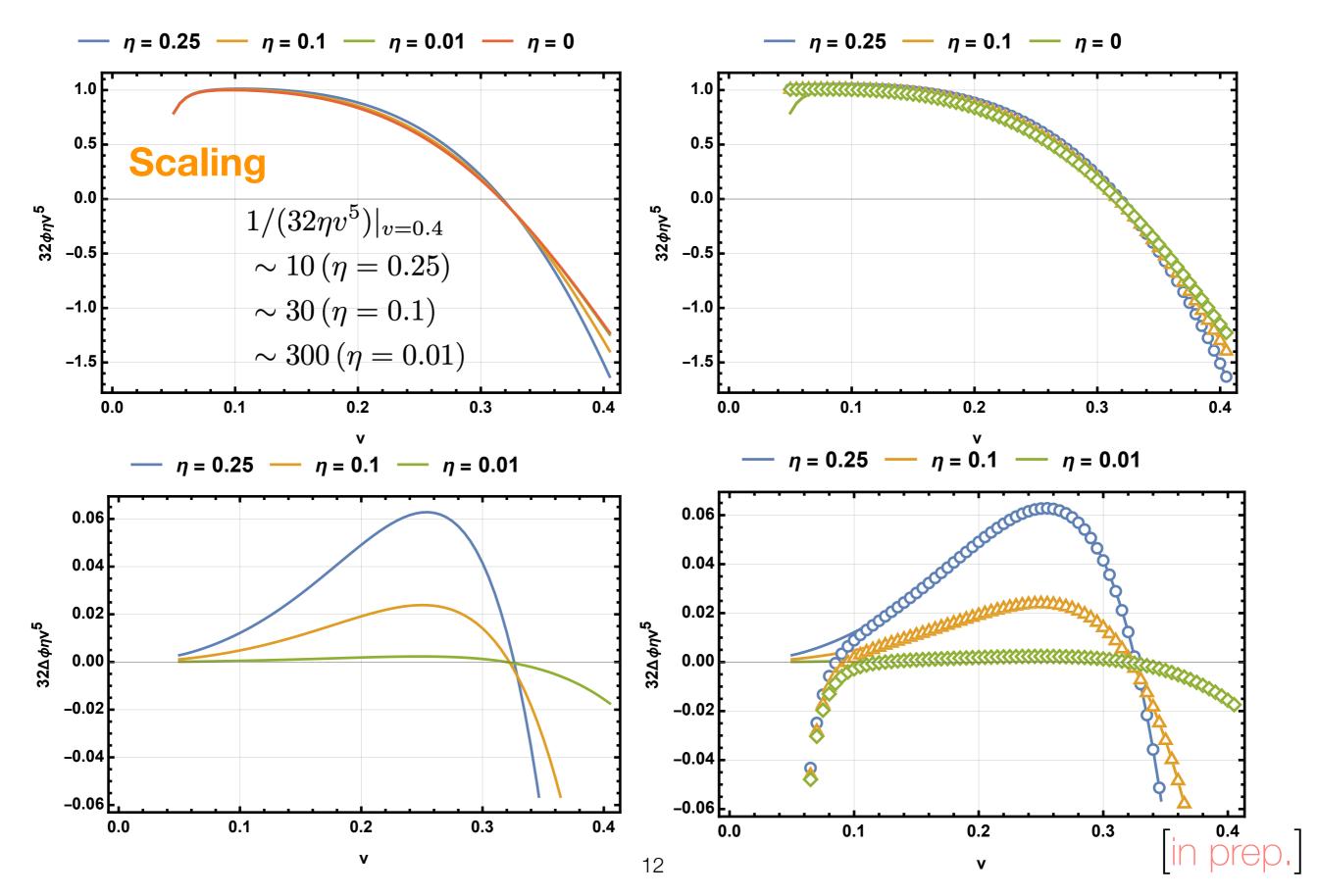
Impact of higher order terms in BHP theory on # GW cycles

PN Order ↓	$\Delta N_{cycle}^{BHP(e_0^{10})}$	$\Delta N_{cycle}^{BHP(e_0^{10})}$	$\Delta N_{cycle}^{BHP(e_0^{10})}$	$\Delta N_{cycle}^{BHP(e_0^{10})}$
$e_0 \rightarrow$	$e_0 = 0$	$e_0 = 0.1$	$e_0 = 0.2$	$e_0 = 0.3$
3.5 PN	-	-350.27×10^{-4}	-140.11×10^{-3}	-315.245×10^{-3}
4 PN	-63.011	423.11×10^{-4}	169.24×10^{-3}	380.80×10^{-3}
4.5 PN	47.008	-12.82×10^{-4}	-5.13×10^{-3}	-11.54×10^{-3}
5 PN	-51.388	13.39×10^{-4}	5.36×10^{-3}	12.05×10^{-3}
5.5PN-12PN	26.49	-	-	-
Total cycles accumulated	-40.90	73.41×10^{-4}	29.36×10^{-3}	66.07×10^{-3}

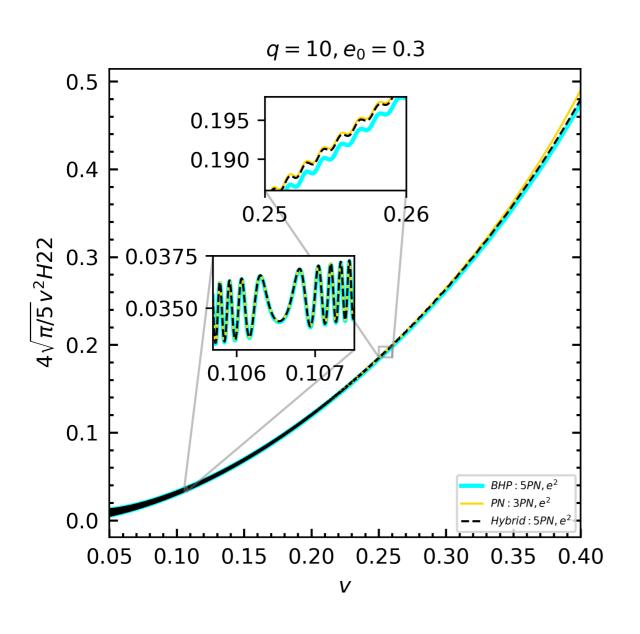
Likely to be more relevant for higher mass ratios

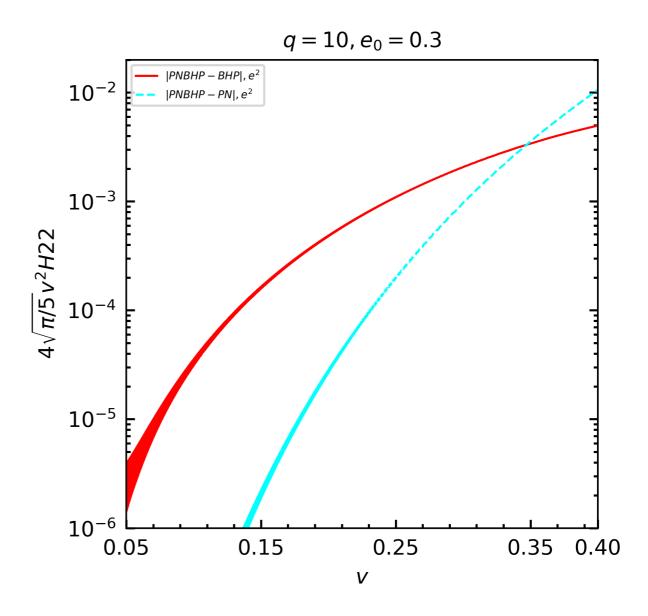
[In prep.]

phase comparisons (e~o.3)

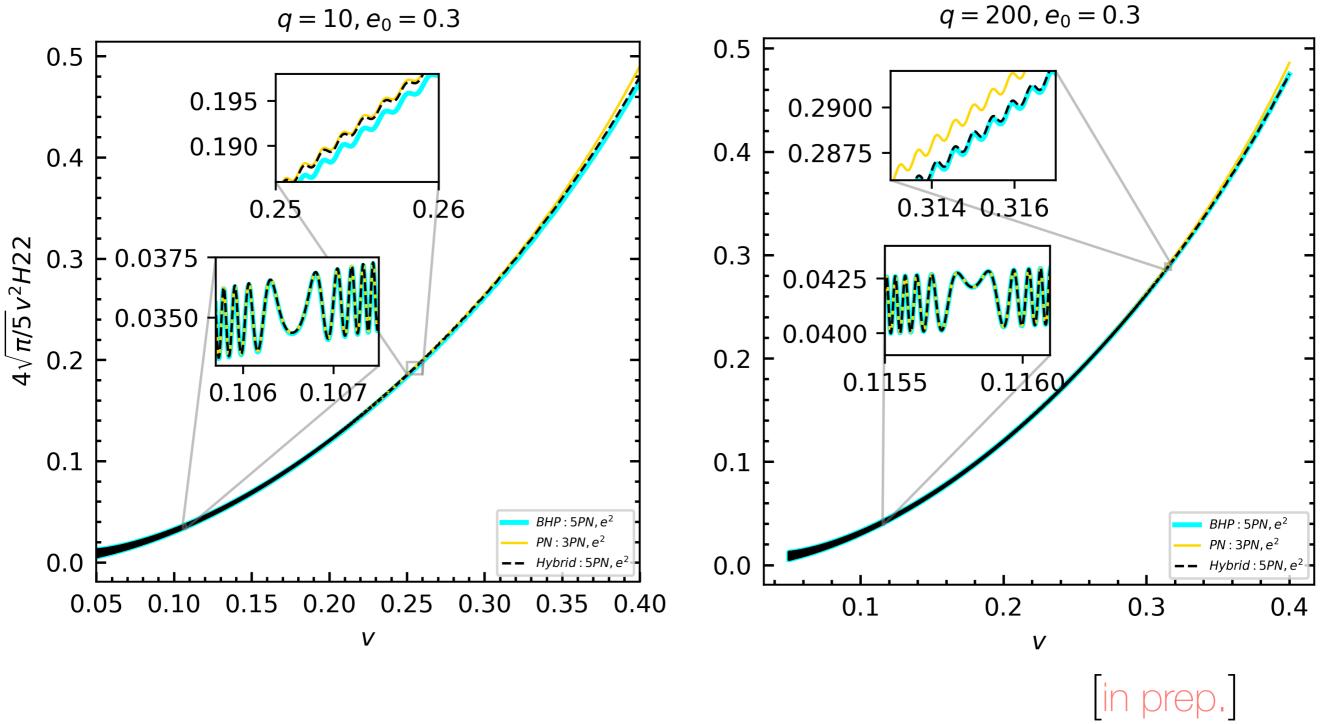


22 mode amp comparisons

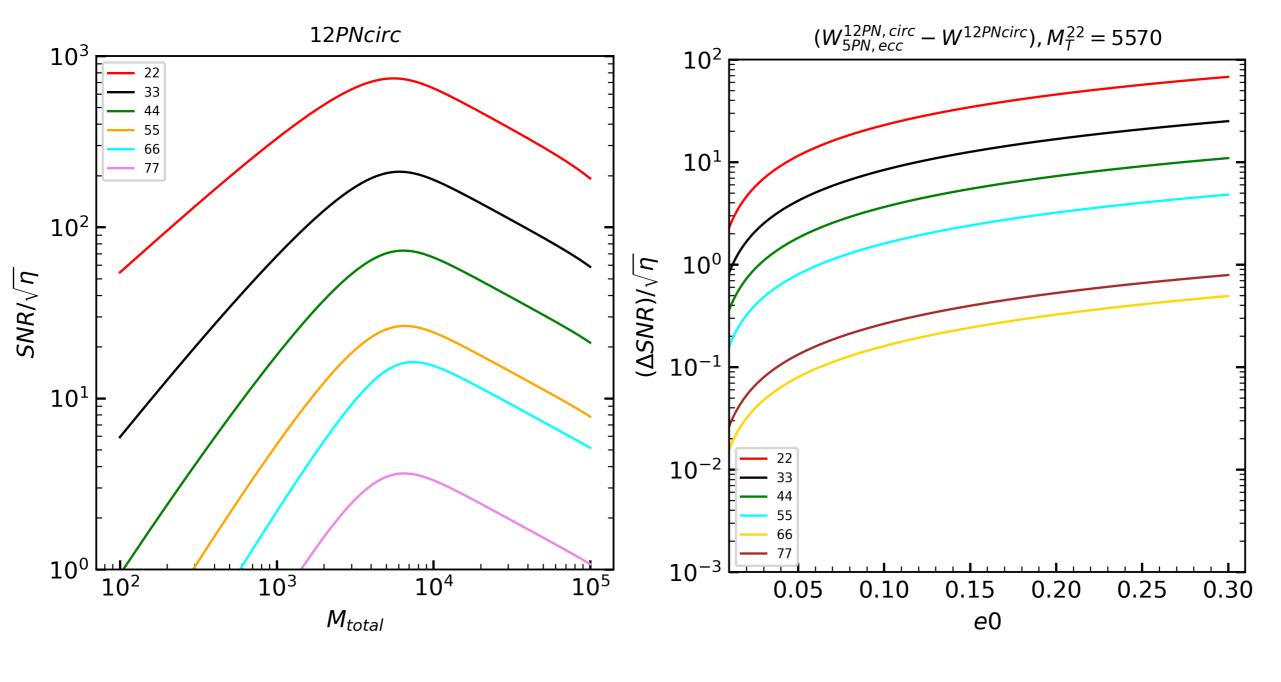




22 mode comparisons



Impact of eccentricity on SNR

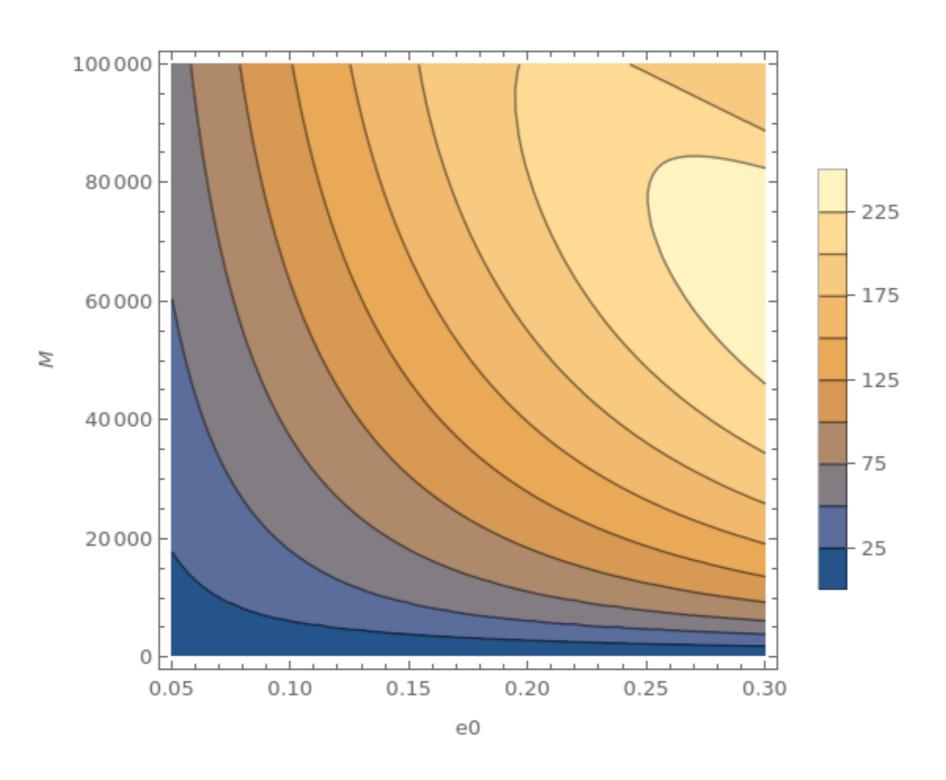


Summary

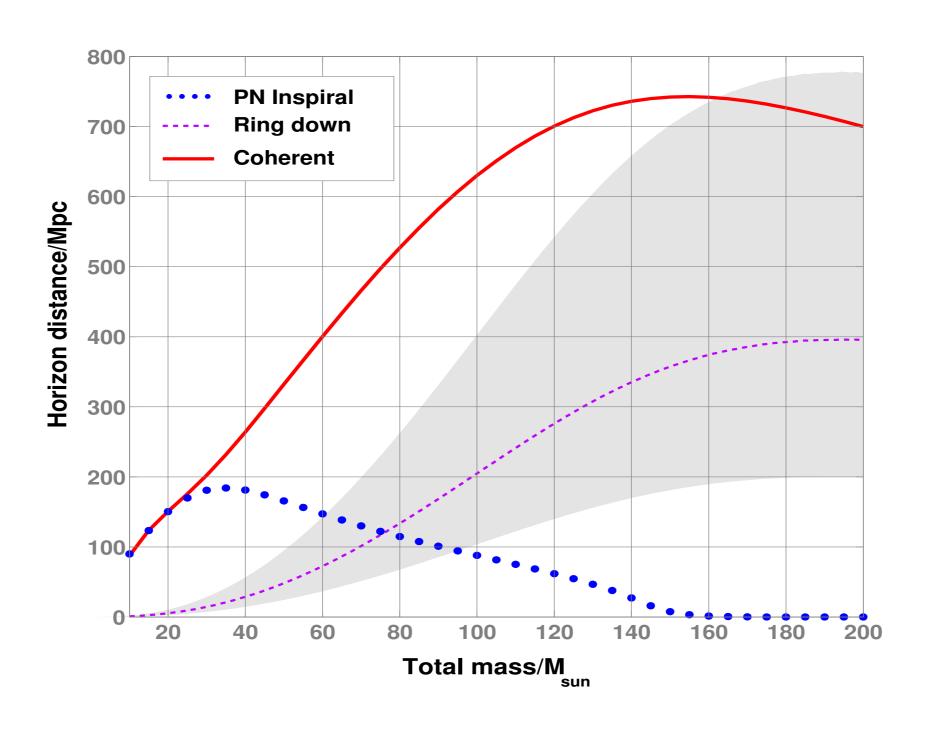
- Compact binaries involving at least one IMBH and a stellar or supermassive BH can be classified as IMRIs and should be detectable by detectors in space and on ground.
- Analytical solutions from perturbative approaches in GR may be combined to obtain models that may help detect and extract properties of these sources.
- Based on simple SNR estimates, once may be able to infer the presence of higher modes and/or eccentric nature of the binary.

additional slides

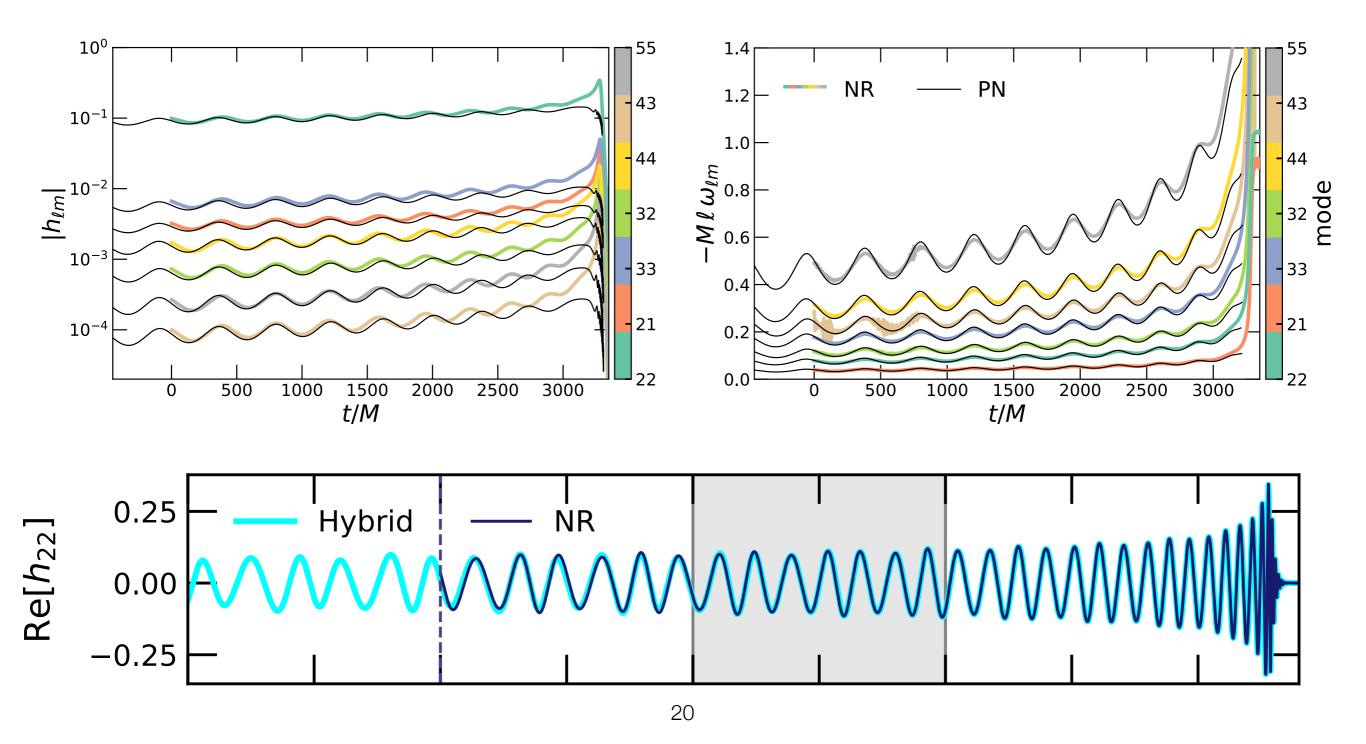
Gain in eccentricity



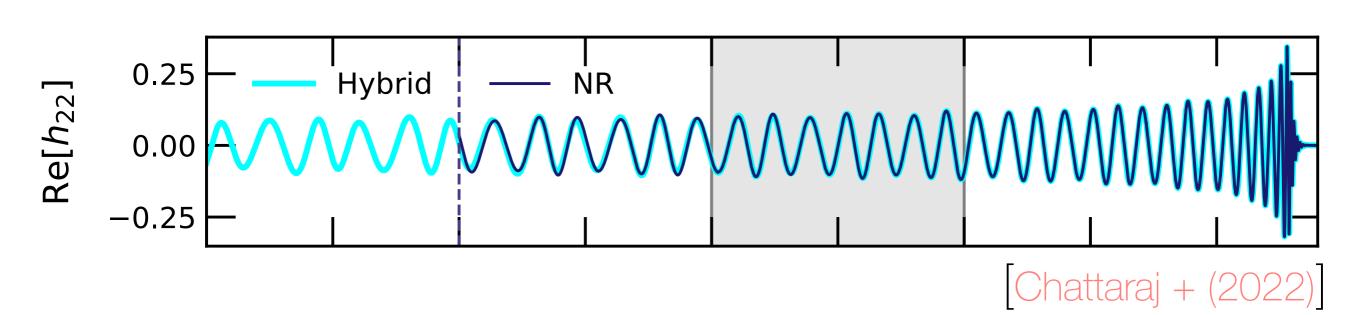
Coherent inputs



PN-NR comparisons



Construction of Target WF



$$\delta = \min_{t_0, \varphi_0, \psi} \int_{t_i}^{t_f} dt \sum_{\ell, m} \left| h_{\ell m}^{NR}(t - t_0) e^{i(m\varphi_0 + \psi)} - h_{\ell m}^{PN}(t) \right|$$

$$h_{\ell m}^{
m hyb}(t) \equiv \tau(t) \, h_{\ell m}^{
m NR}(t - t_0') \, e^{i(m\varphi_0' + \psi')} + (1 - \tau(t)) \, h_{\ell m}^{
m PN}(t)$$

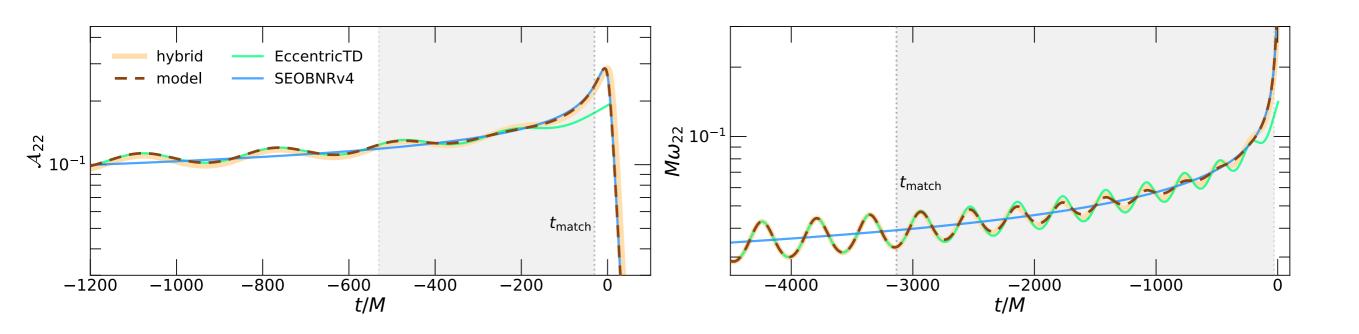
$$\tau(t) \equiv \begin{cases} 0 & \text{if } t < t_{i} \\ \frac{t - t_{i}}{t_{f} - t_{i}} & \text{if } t_{i} \leq t < t_{f} \\ 1 & \text{if } t_{f} \leq t. \end{cases}$$

A quadrupole mode model

$$\mathcal{A}_{22}^{\text{model}}(t) \equiv \tau_{\text{a}}(t) \, \mathcal{A}_{22}^{\text{IMR}}(t) + (1 - \tau_{\text{a}}(t)) \, \mathcal{A}_{22}^{\text{inspiral}}(t)$$

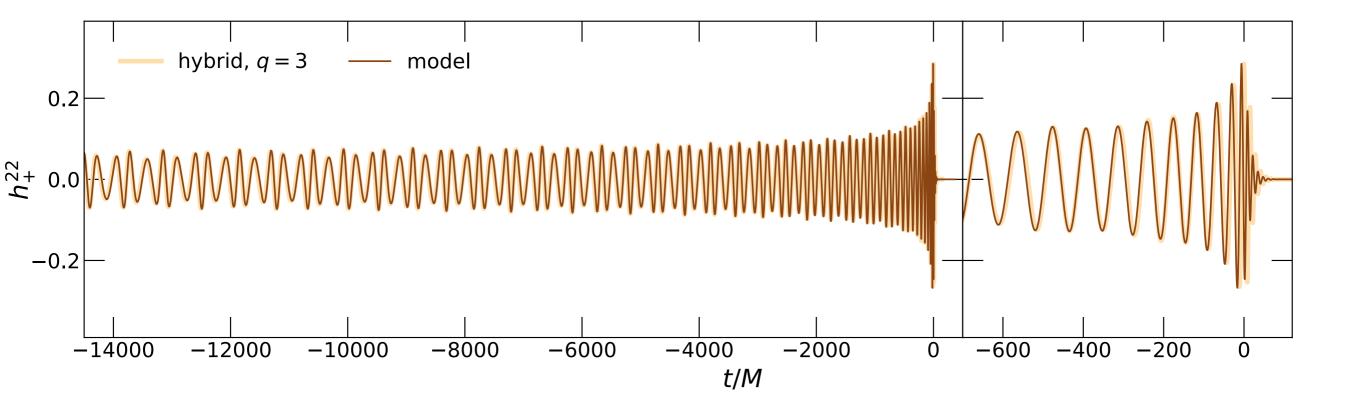
$$\omega_{22}^{\text{model}}(t) \equiv \tau_{\text{a}}(t) \,\omega_{22}^{\text{IMR}}(t) + (1 - \tau_{\text{a}}(t)) \,\omega_{22}^{\text{inspiral}}(t)$$

$$\tau_{a}(t) \equiv \begin{cases} 0 & \text{if } t < t_{i} \\ \frac{t - t_{i}}{t_{f} - t_{i}} & \text{if } t_{i} \leq t < t_{f} \\ 1 & \text{if } t_{f} \leq t. \end{cases}$$

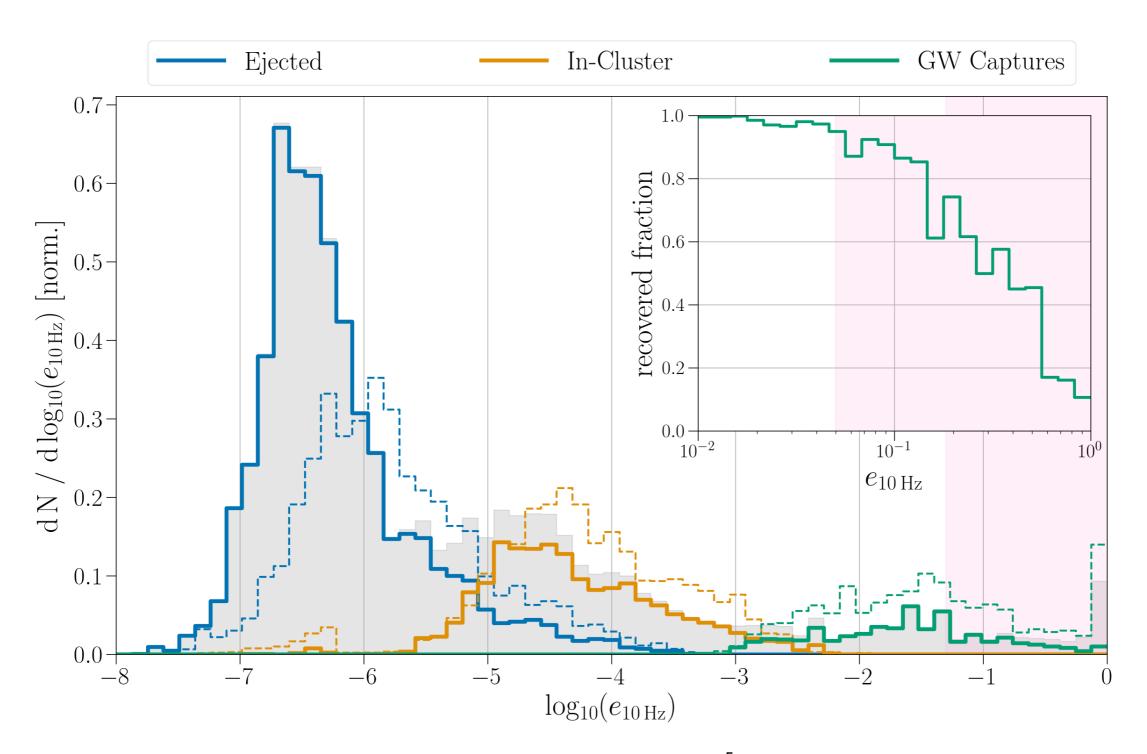


Chattaraj + (2022)

IMR construction



Eccentric Population



Zevin et al., ApJL, 921 L43 (2021)]