

Insights into the evolution of hadronic collisions with flow observables

NBI Heavy-ion Seminar, Copenhagen
Vytautas Vislavicius, Lund University, Sweden
February 9th, 2023



LUND
UNIVERSITY

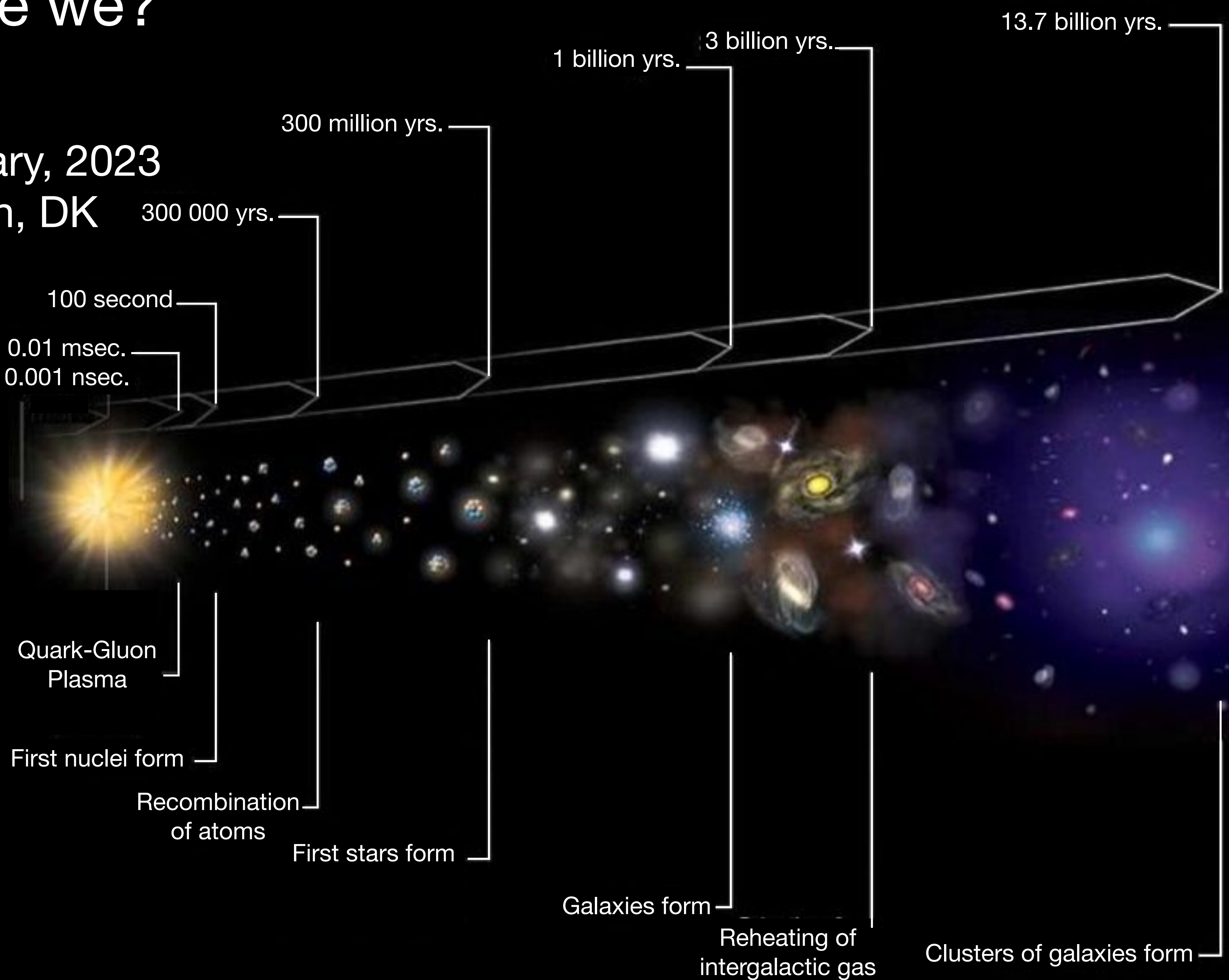
Where (and when) are we?

Where (and when) are we?

~11:20, Thursday, 9th February, 2023
Aud. A., NBI, Copenhagen, DK

Where (and when) are we?

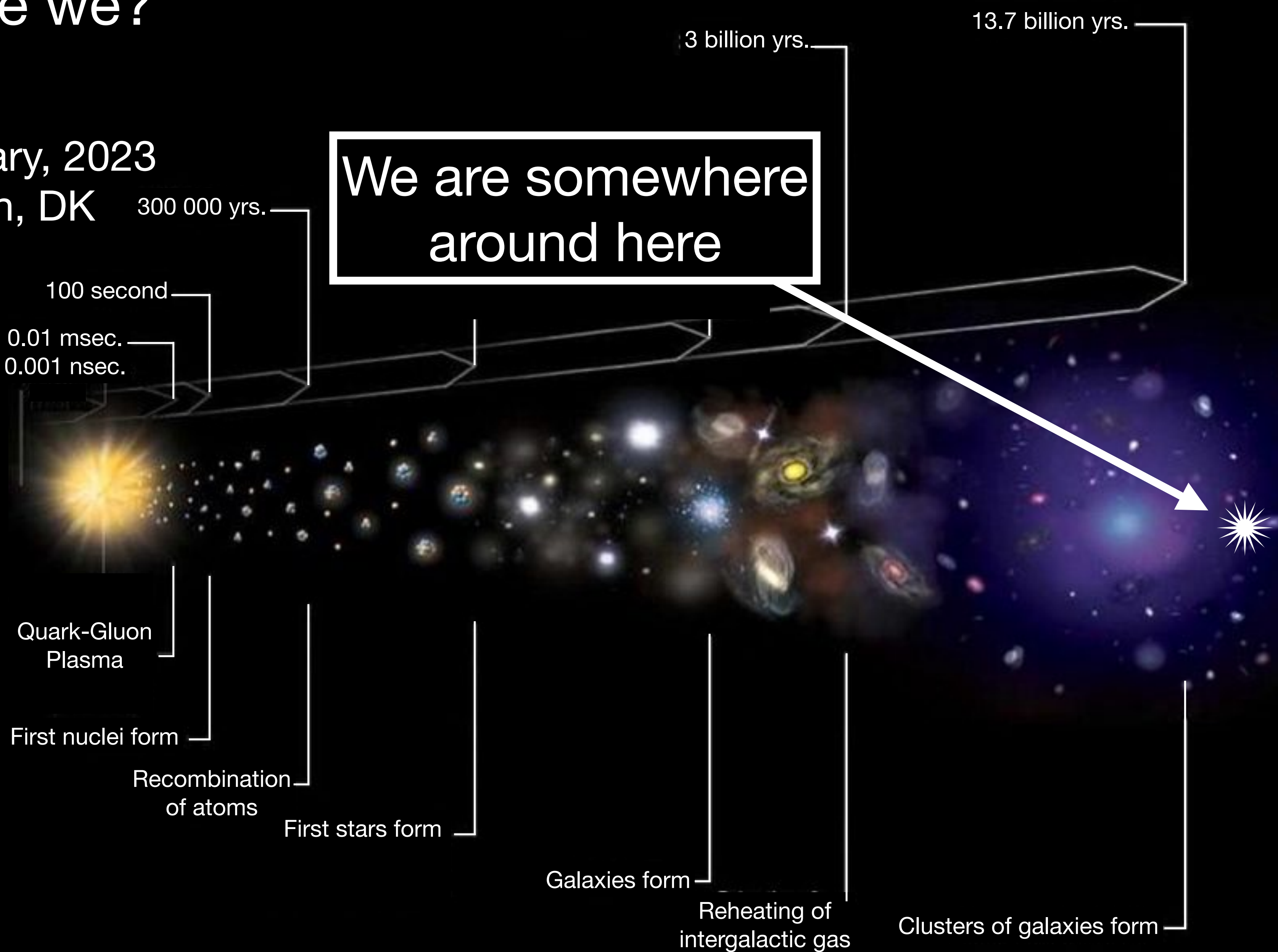
~11:20, Thursday, 9th February, 2023
Aud. A., NBI, Copenhagen, DK



Where (and when) are we?

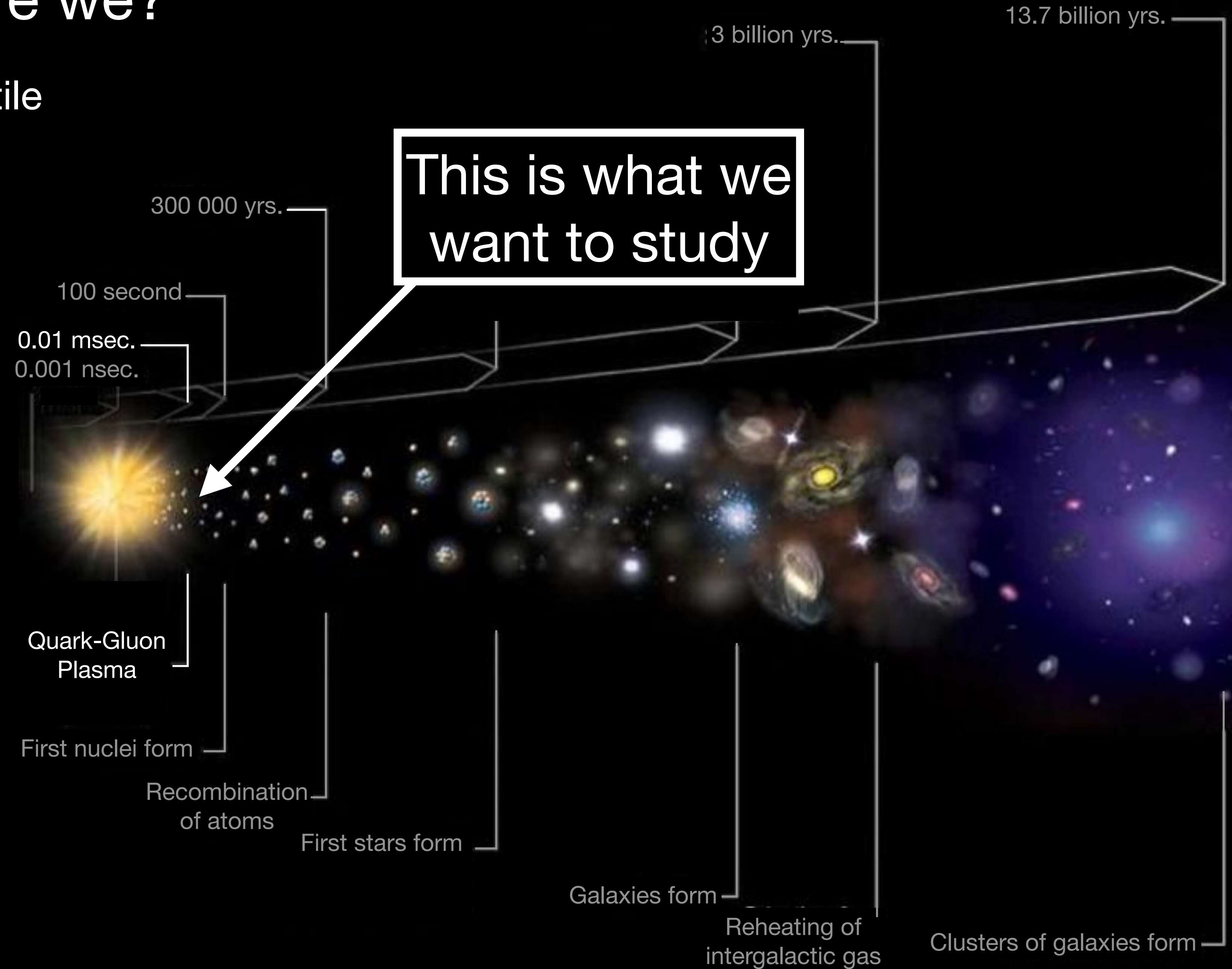
~11:20, Thursday, 9th February, 2023
Aud. A., NBI, Copenhagen, DK

We are somewhere around here



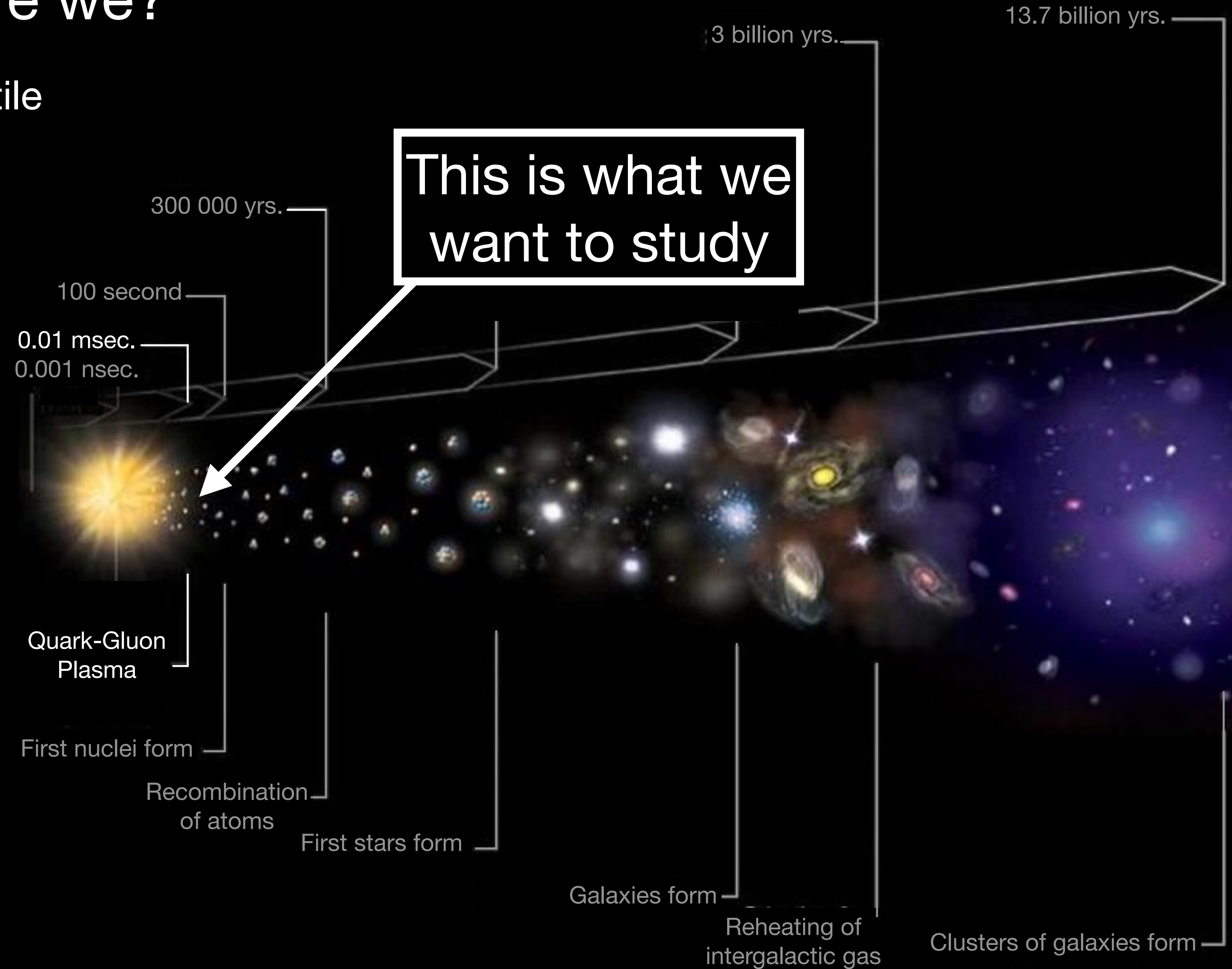
Where (and when) are we?

- Interested in extremely volatile processes that happened 13+ billion years ago



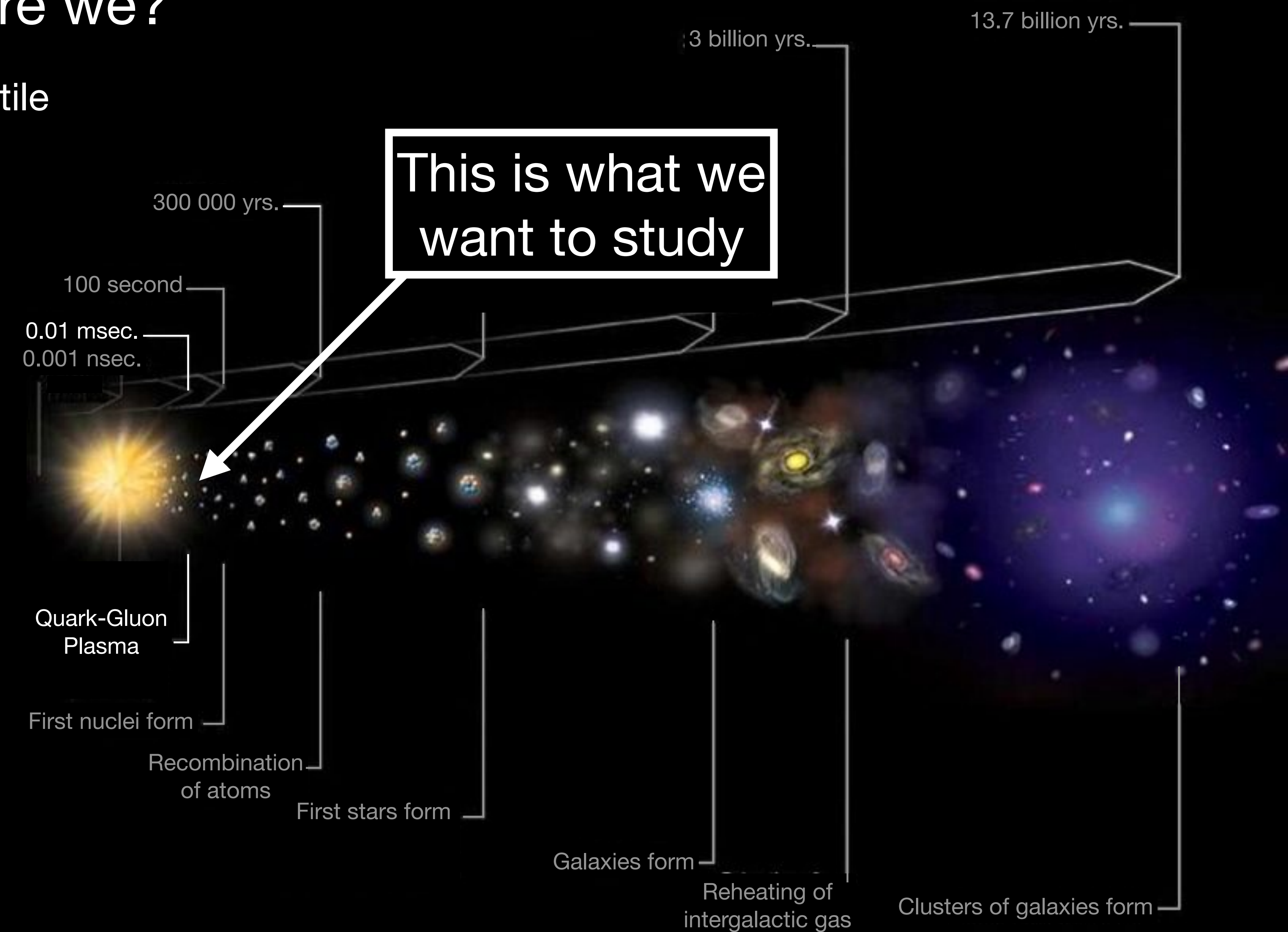
Where (and when) are we?

- Interested in extremely volatile processes that happened 13+ billion years ago
- Don't have means, knowledge, or resources to make new Big Bang(s)



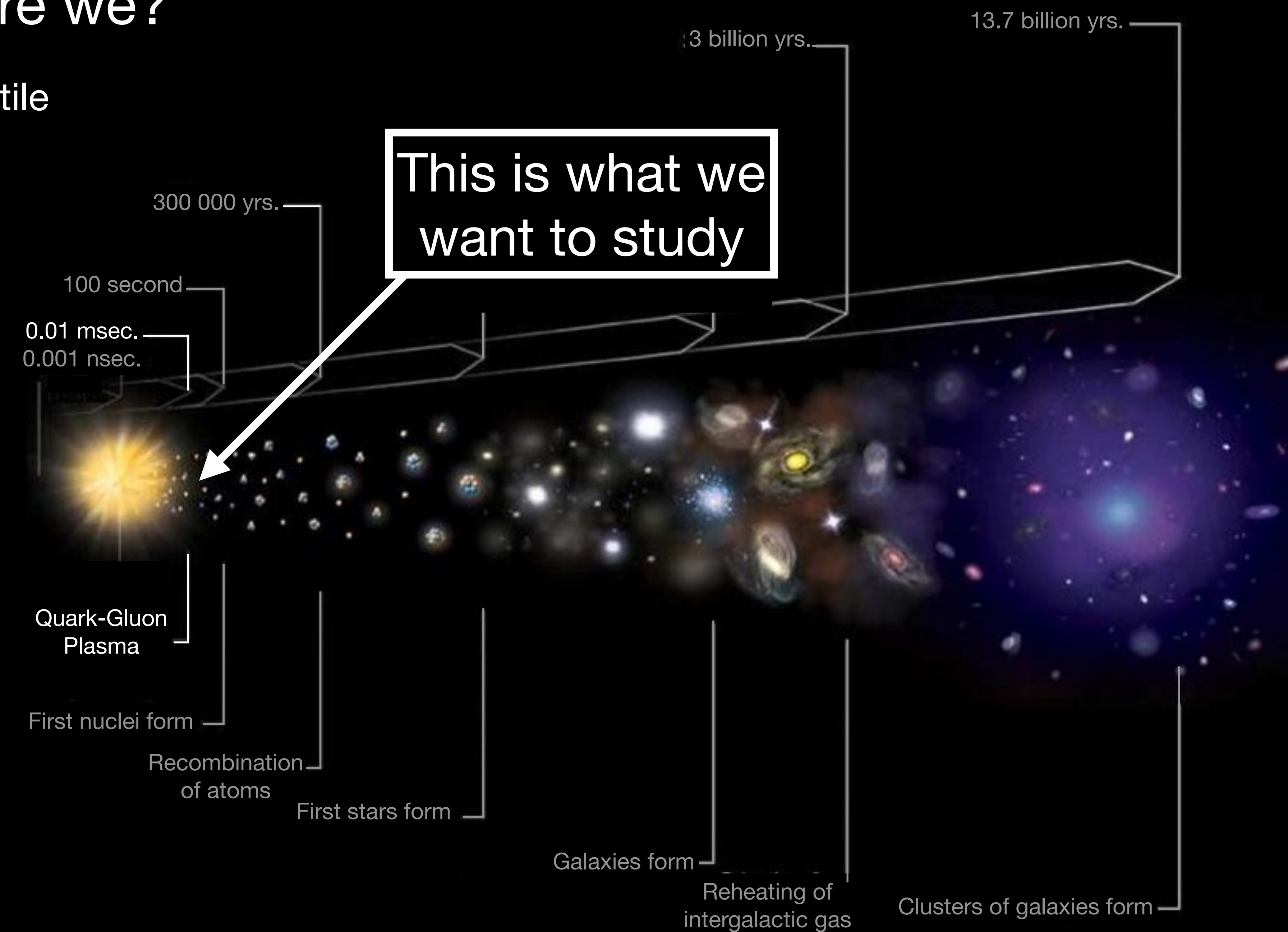
Where (and when) are we?

- Interested in extremely volatile processes that happened 13+ billion years ago
- Don't have means, knowledge, or resources to make new Big Bang(s)
- ...

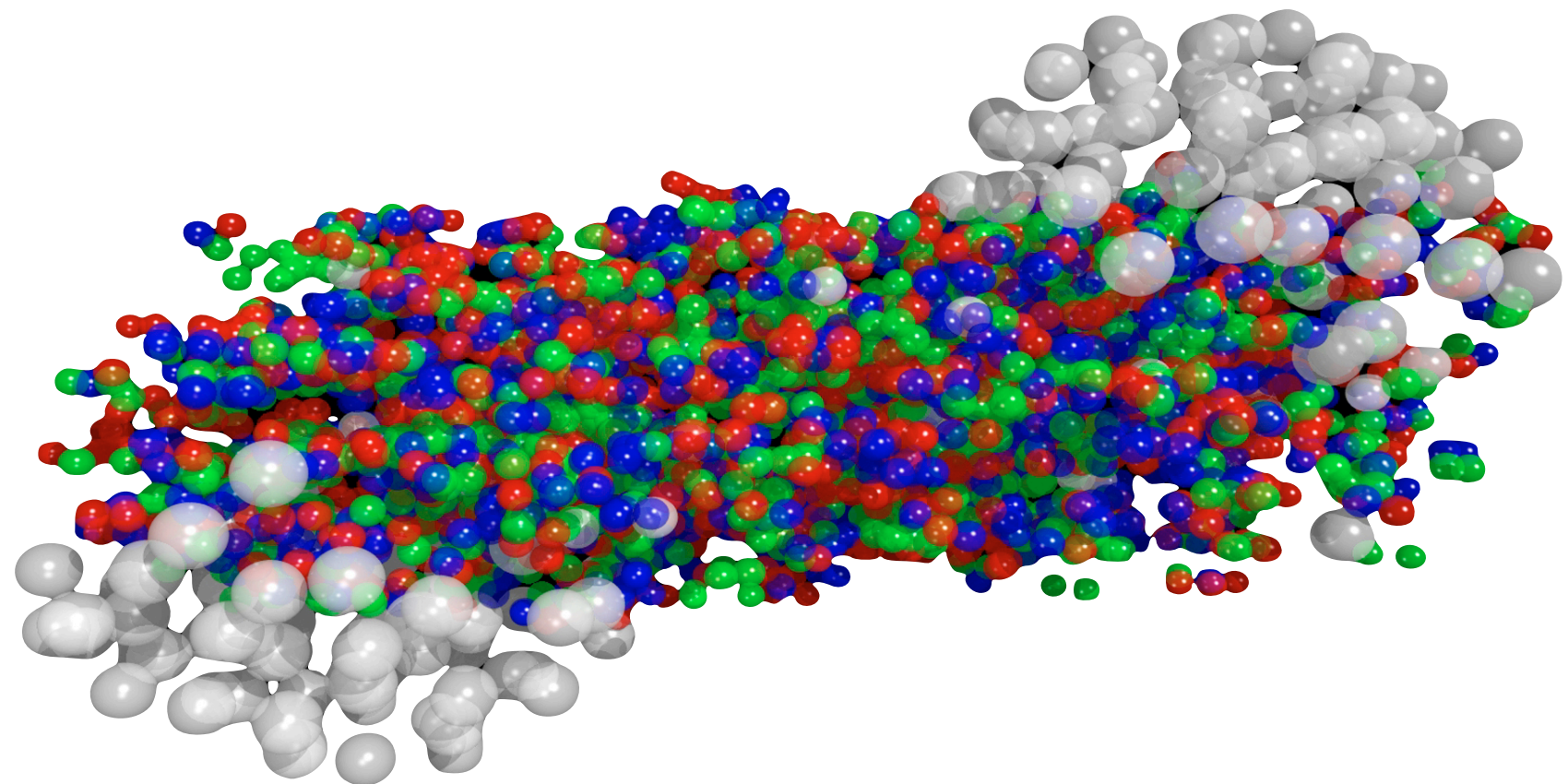


Where (and when) are we?

- Interested in extremely volatile processes that happened 13+ billion years ago
- Don't have means, knowledge, or resources to make new Big Bang(s)
- ...
- But we can produce small droplets of QGP by smashing heavy ions* at relativistic energies

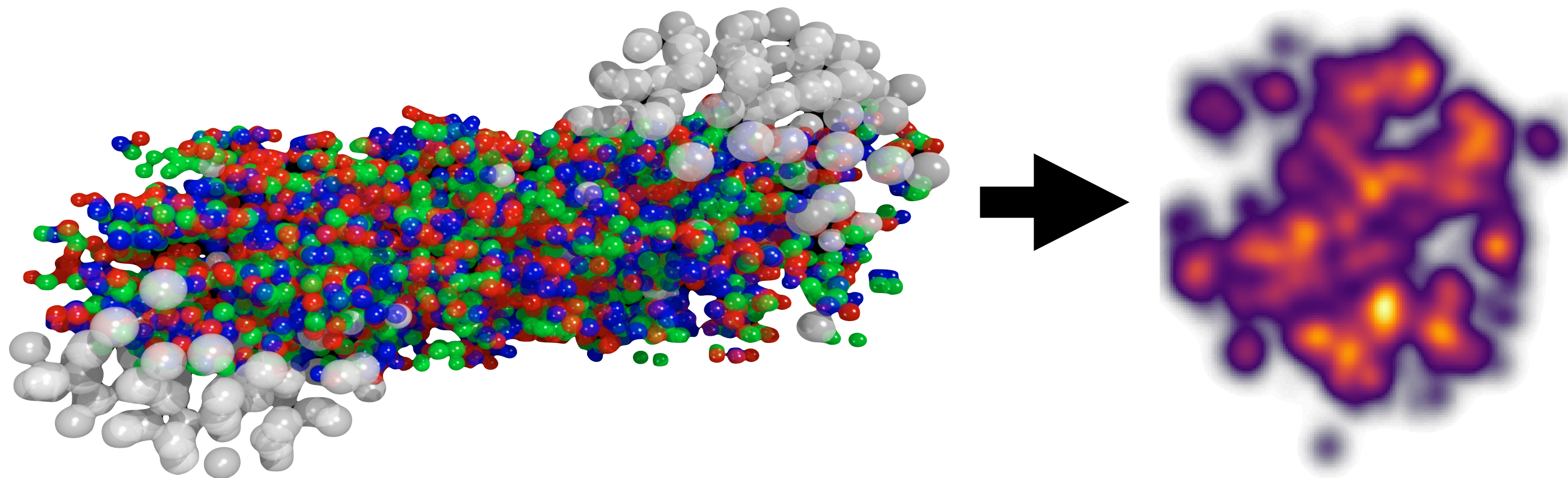


From collisions to measurements



From collisions to measurements

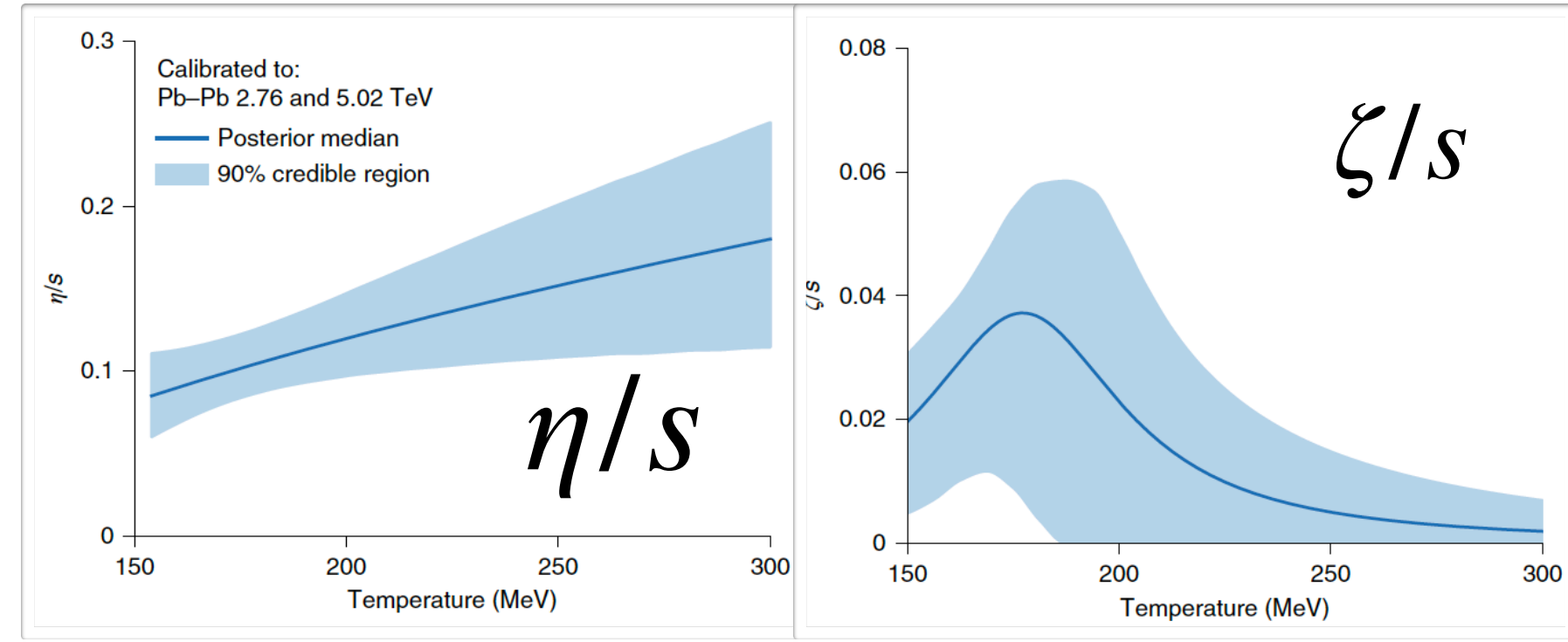
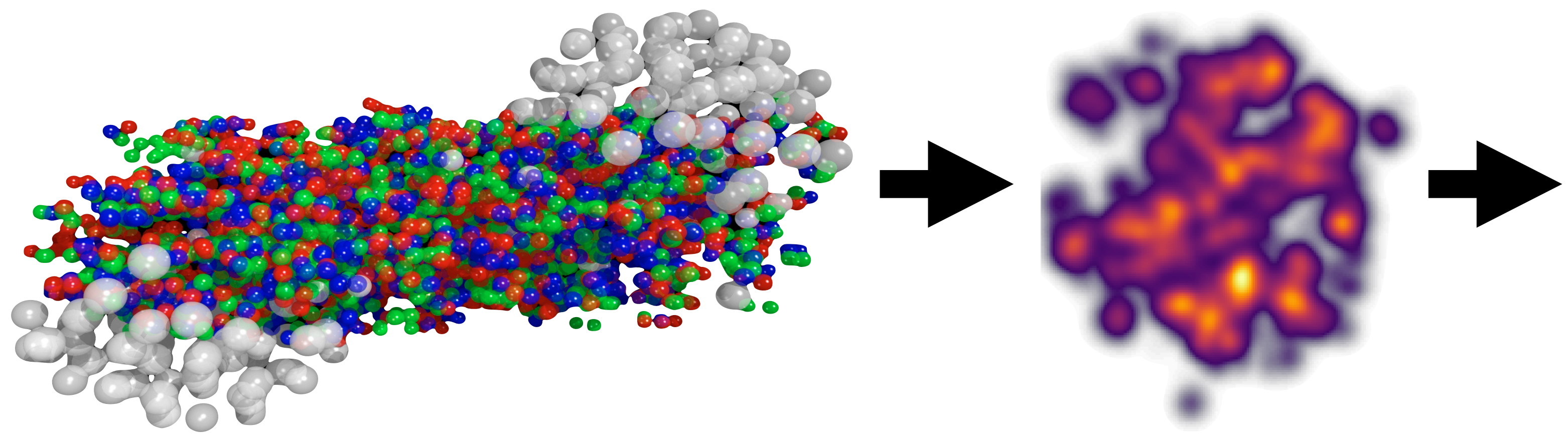
- Overlap between colliding nuclei:
⇒ Initial state, geometry & its fluctuations



From collisions to measurements

- Overlap between colliding nuclei:
⇒ Initial state, geometry & its fluctuations

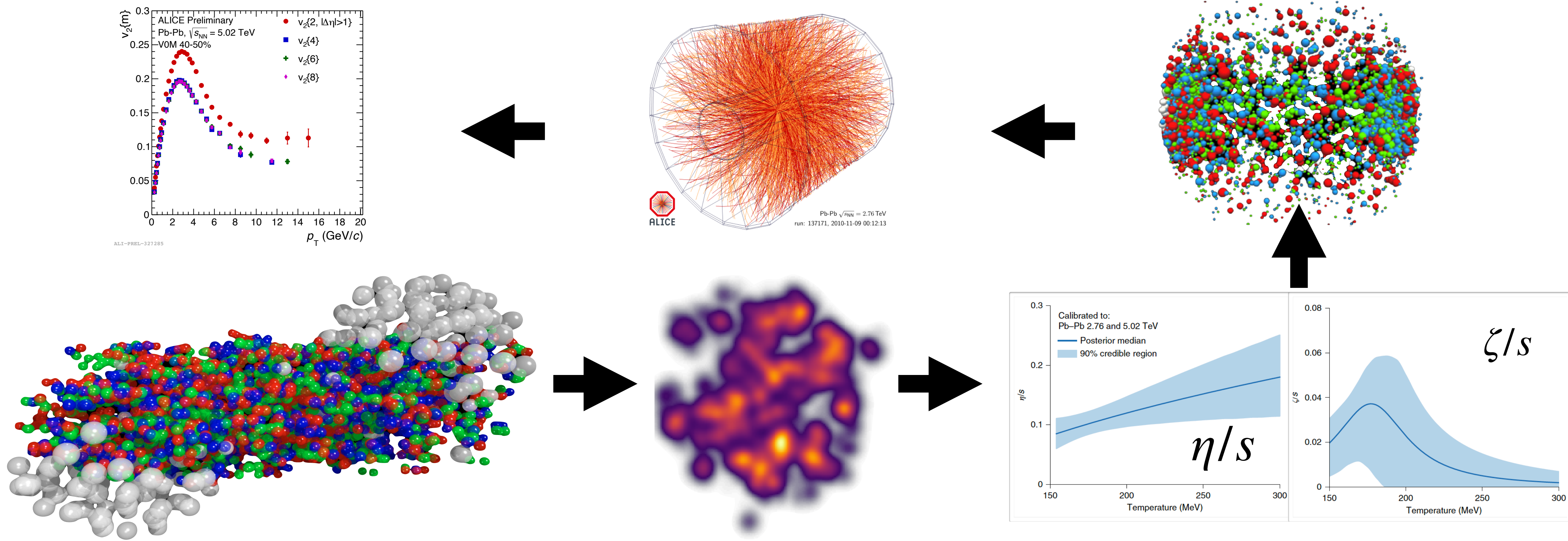
- Hydrodynamic expansion of QGP:
⇒ Radial and anisotropic flow, sensitive to initial state and properties of QGP



From collisions to measurements

- Overlap between colliding nuclei:
 ⇒ Initial state, geometry & its fluctuations

- Hydrodynamic expansion of QGP:
 ⇒ Radial and anisotropic flow, sensitive to initial state and properties of QGP

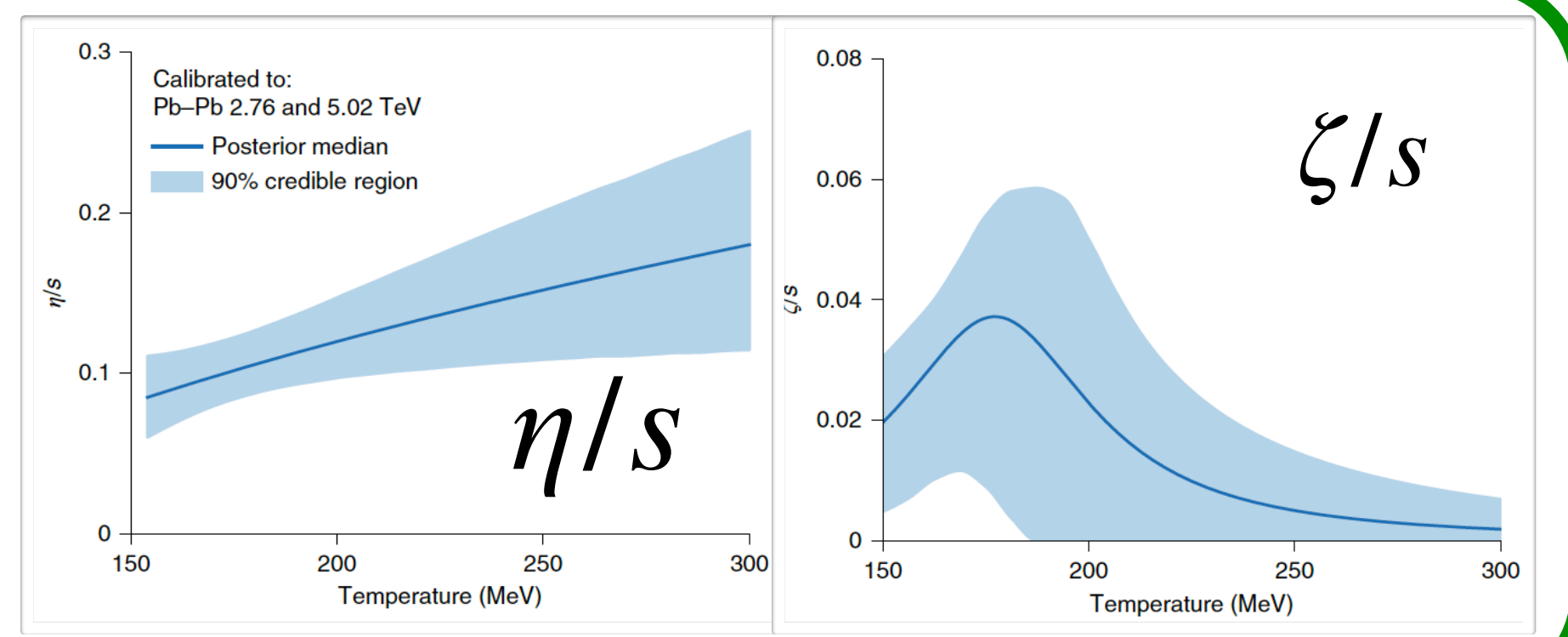
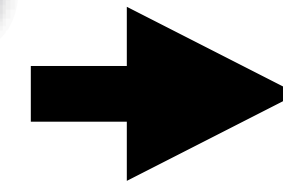
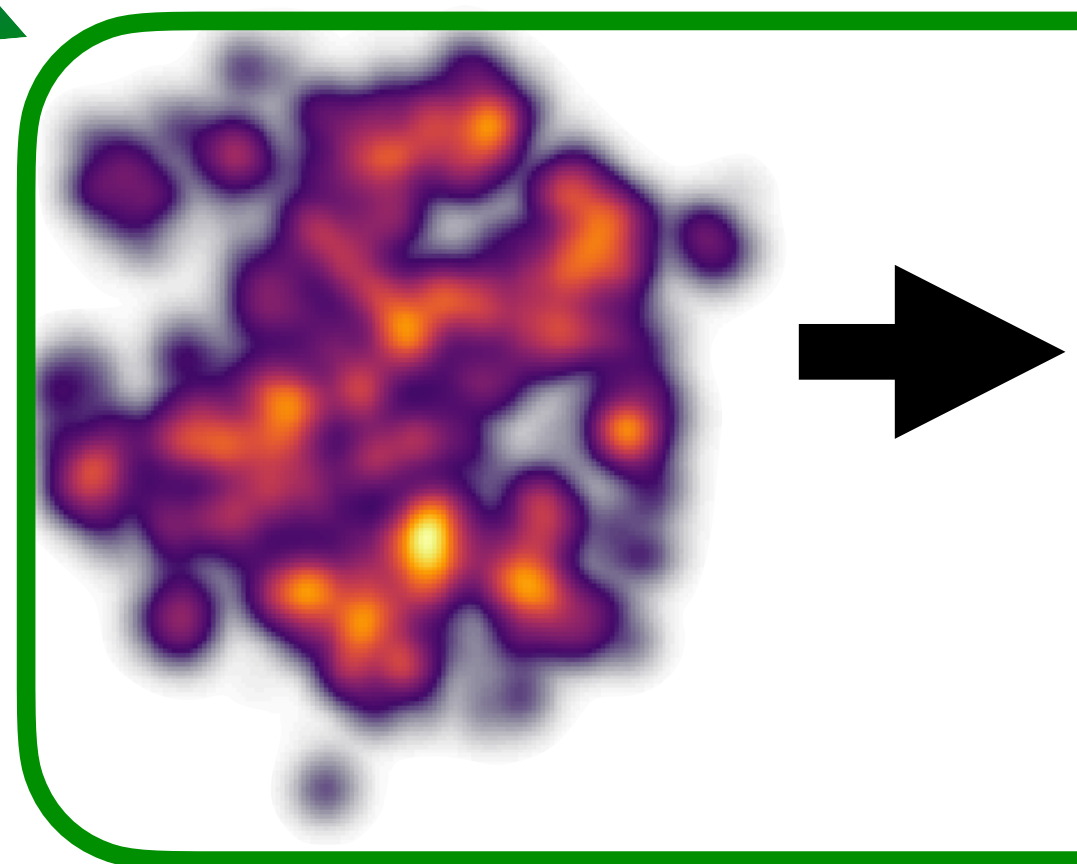
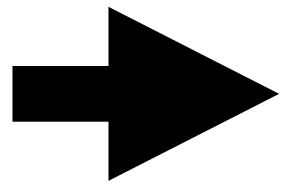
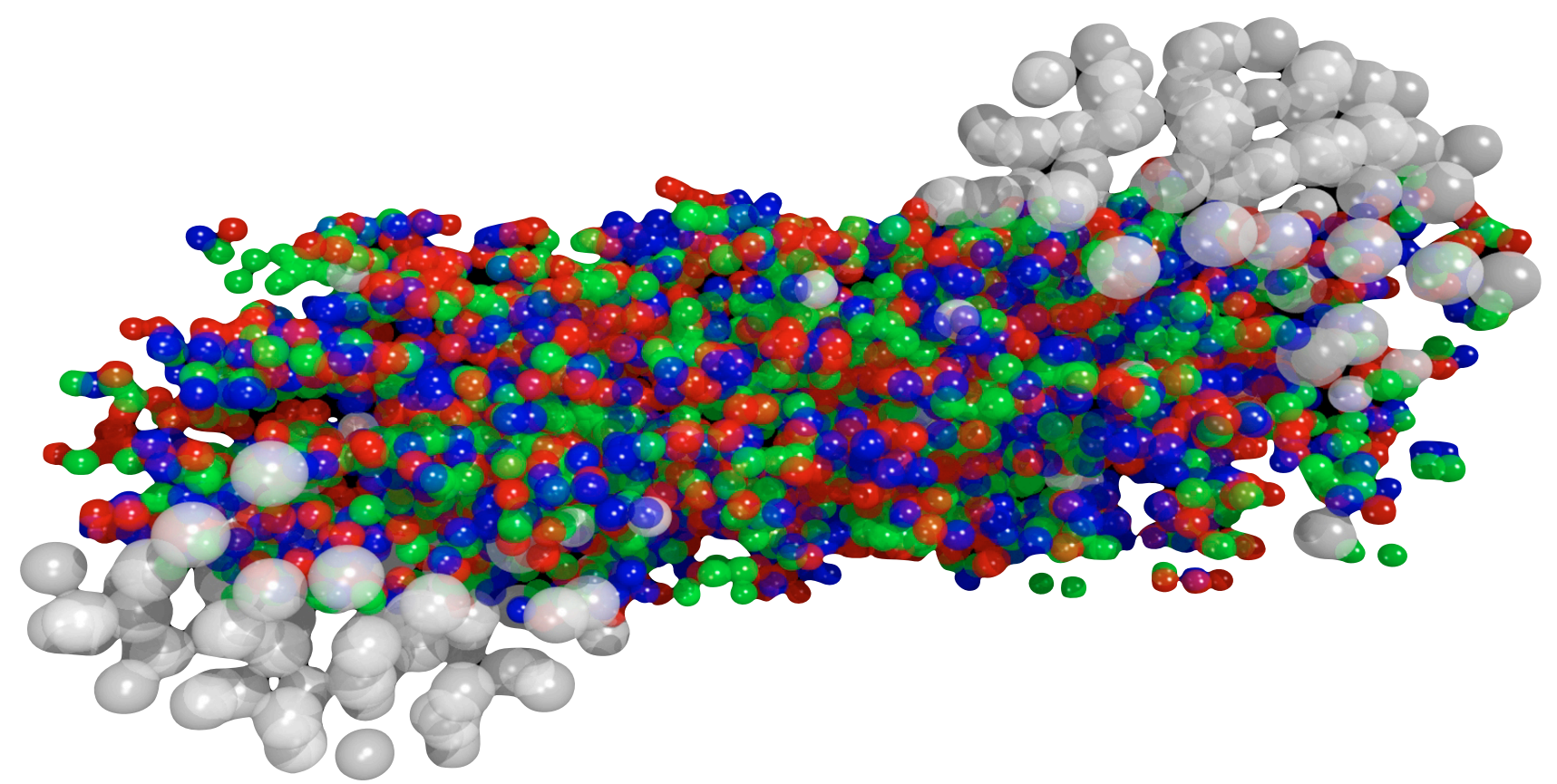
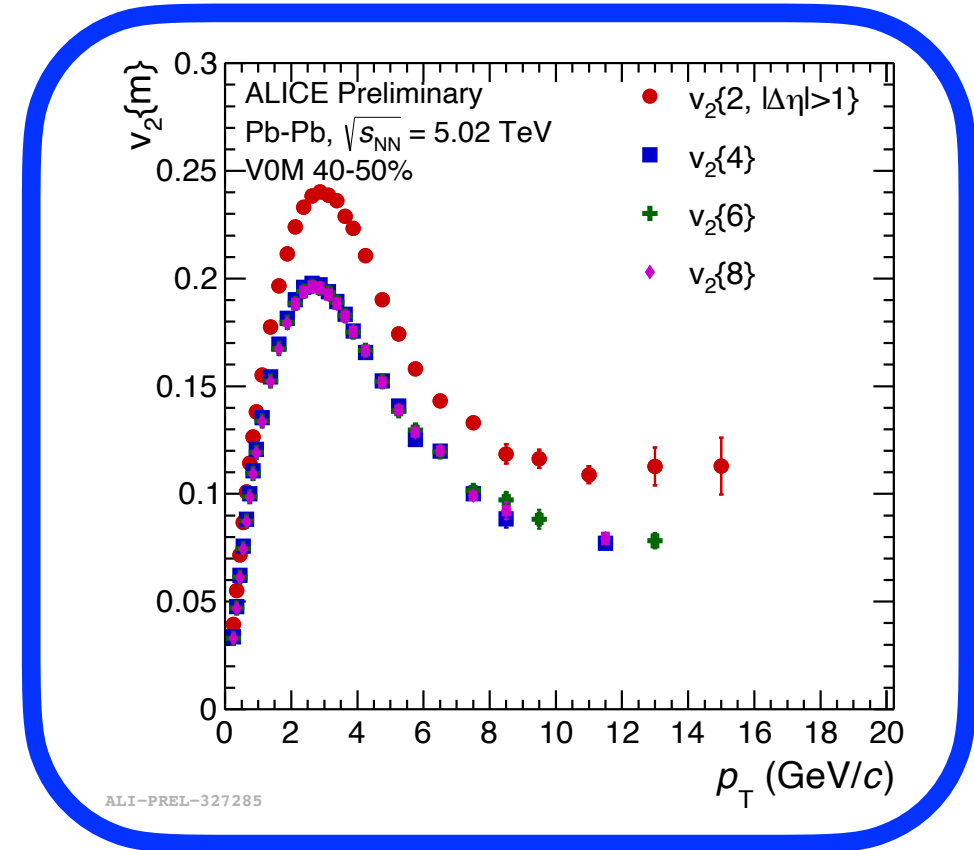


From collisions to measurements

- Overlap between colliding nuclei:
 ⇒ Initial state, geometry & its fluctuations

- Hydrodynamic expansion of QGP:
 ⇒ Radial and anisotropic flow, sensitive to initial state and properties of QGP

What can we learn about the **initial state and properties of QGP** from the measurements of **radial and anisotropic flow**?

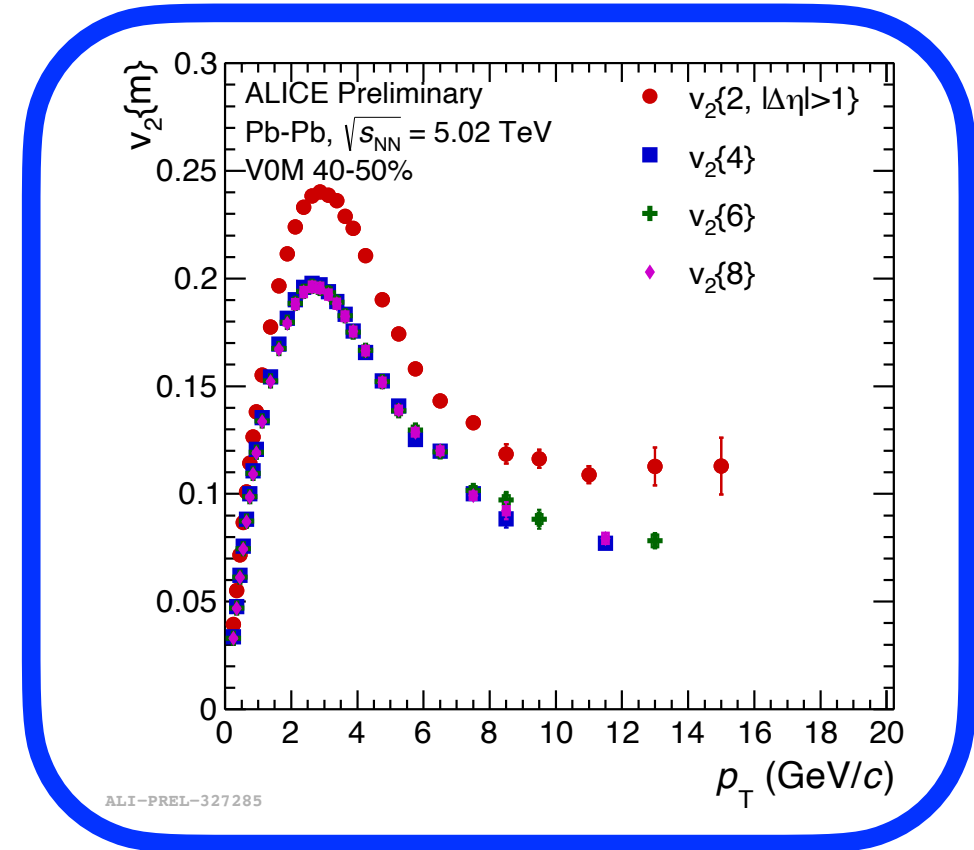


From collisions to measurements

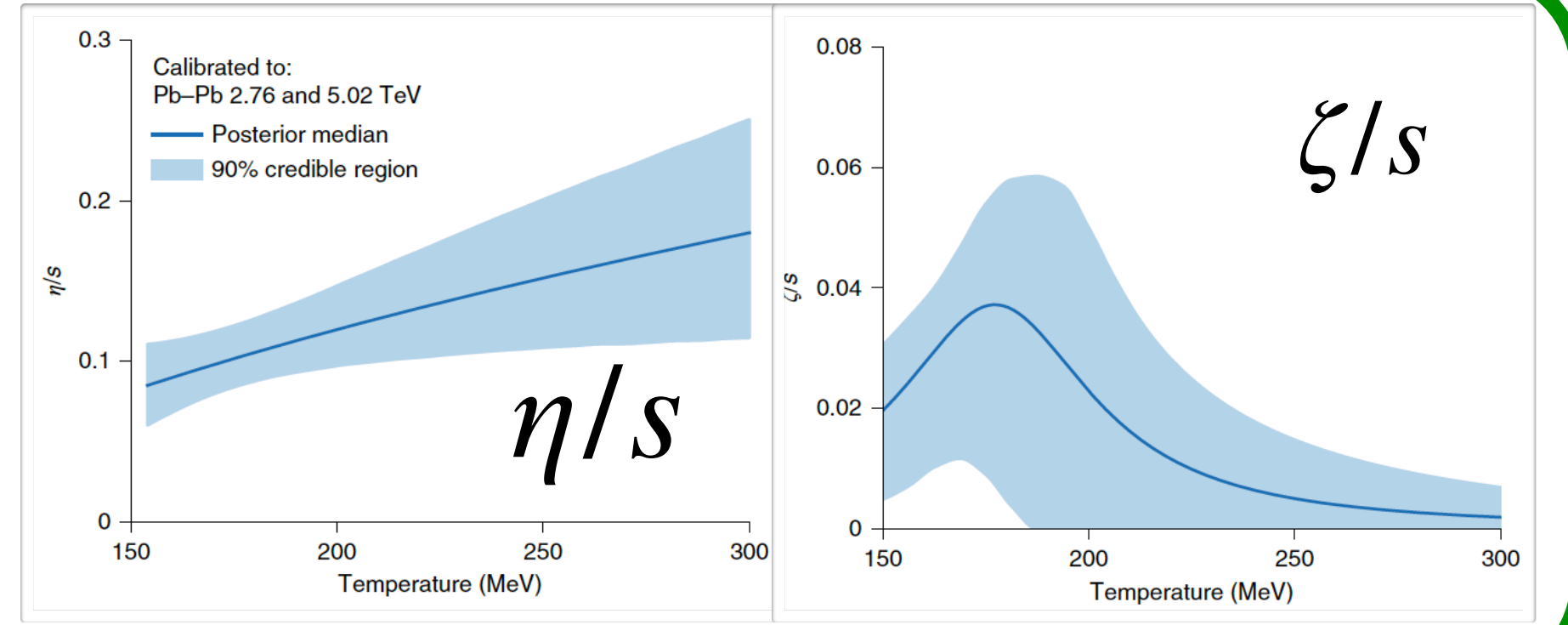
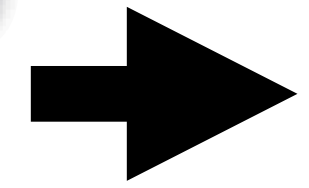
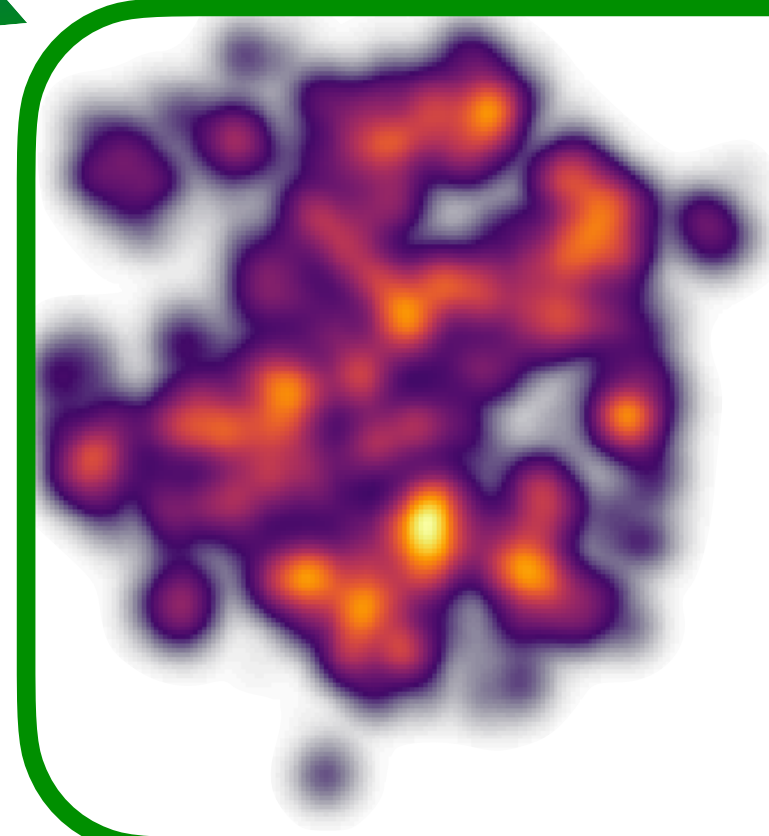
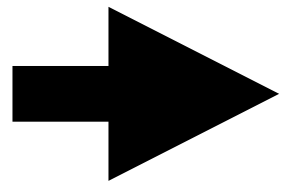
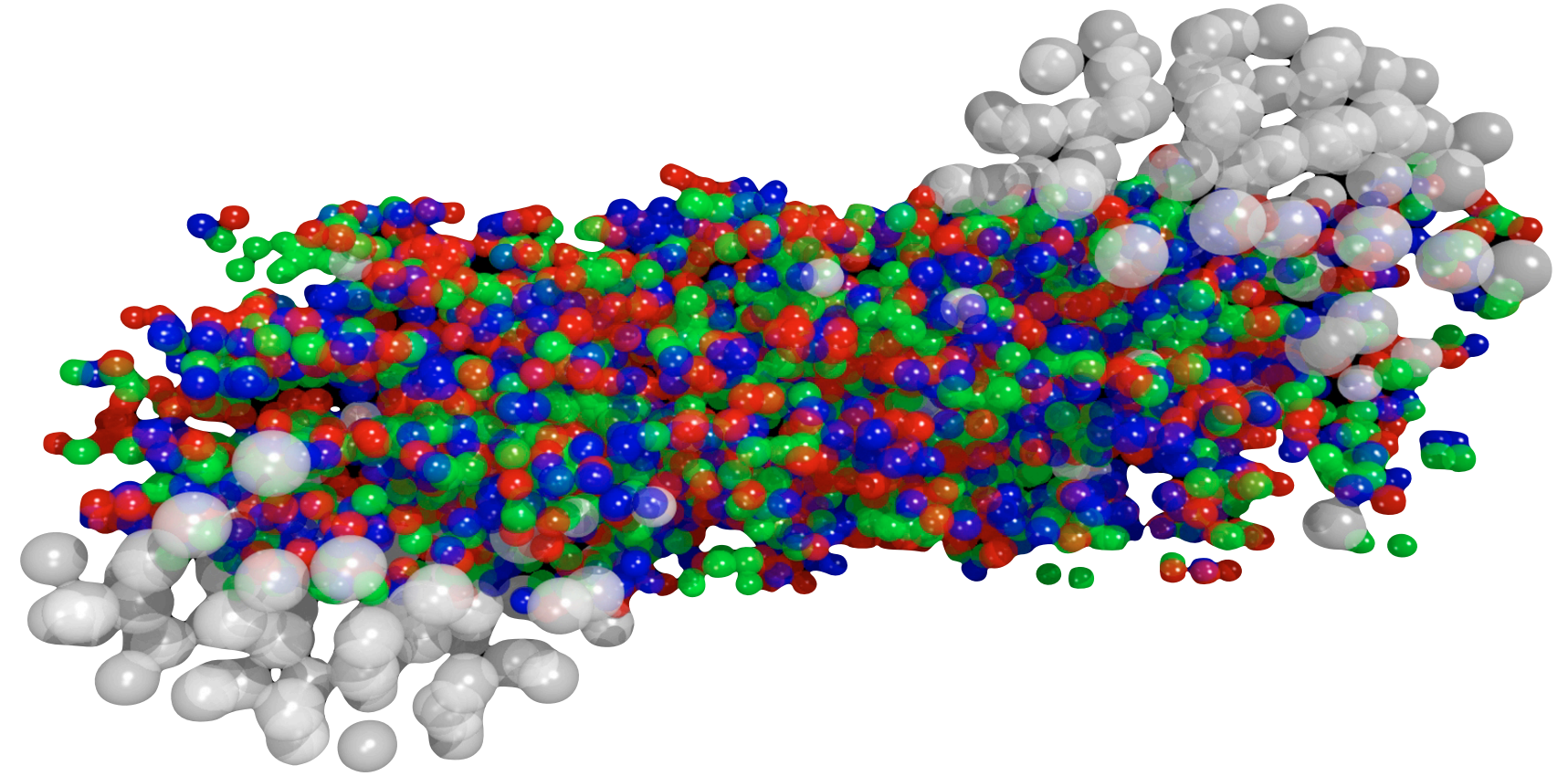
- Overlap between colliding nuclei:
 ⇒ Initial state, geometry & its fluctuations

- Hydrodynamic expansion of QGP:
 ⇒ Radial and anisotropic flow, sensitive to initial state and properties of QGP

What can we learn about the **initial state and properties of QGP** from the measurements of **radial and anisotropic flow**?

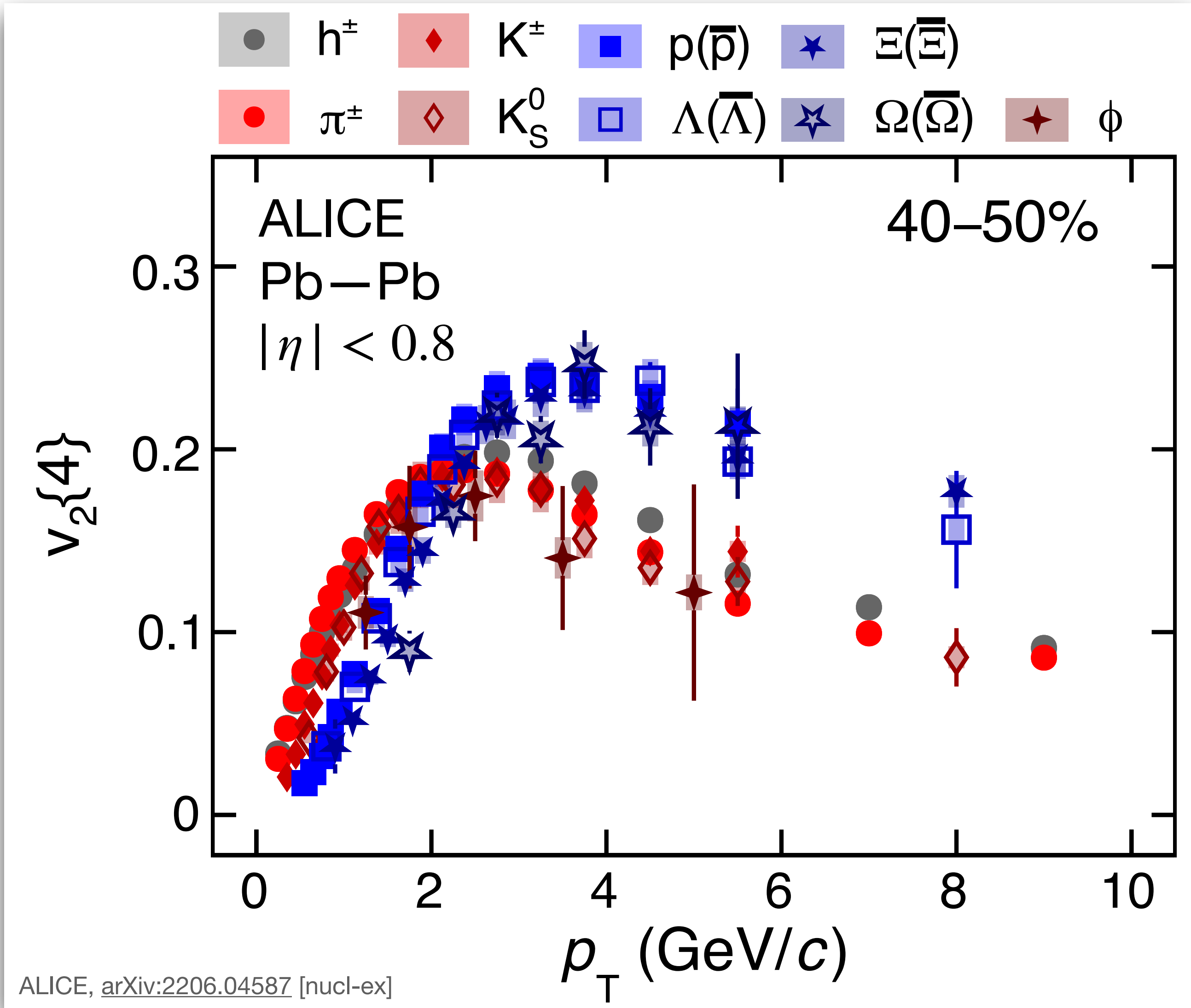


- Radial flow: boost of particles in radial direction by the expanding QGP
- Anisotropic flow: azimuthal asymmetries in particle production, eg elliptic flow v_2



v_2 of baryons and mesons in Pb–Pb

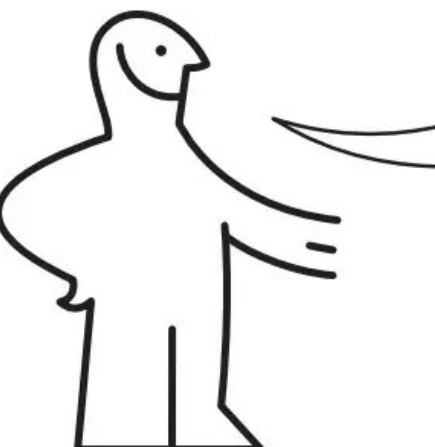
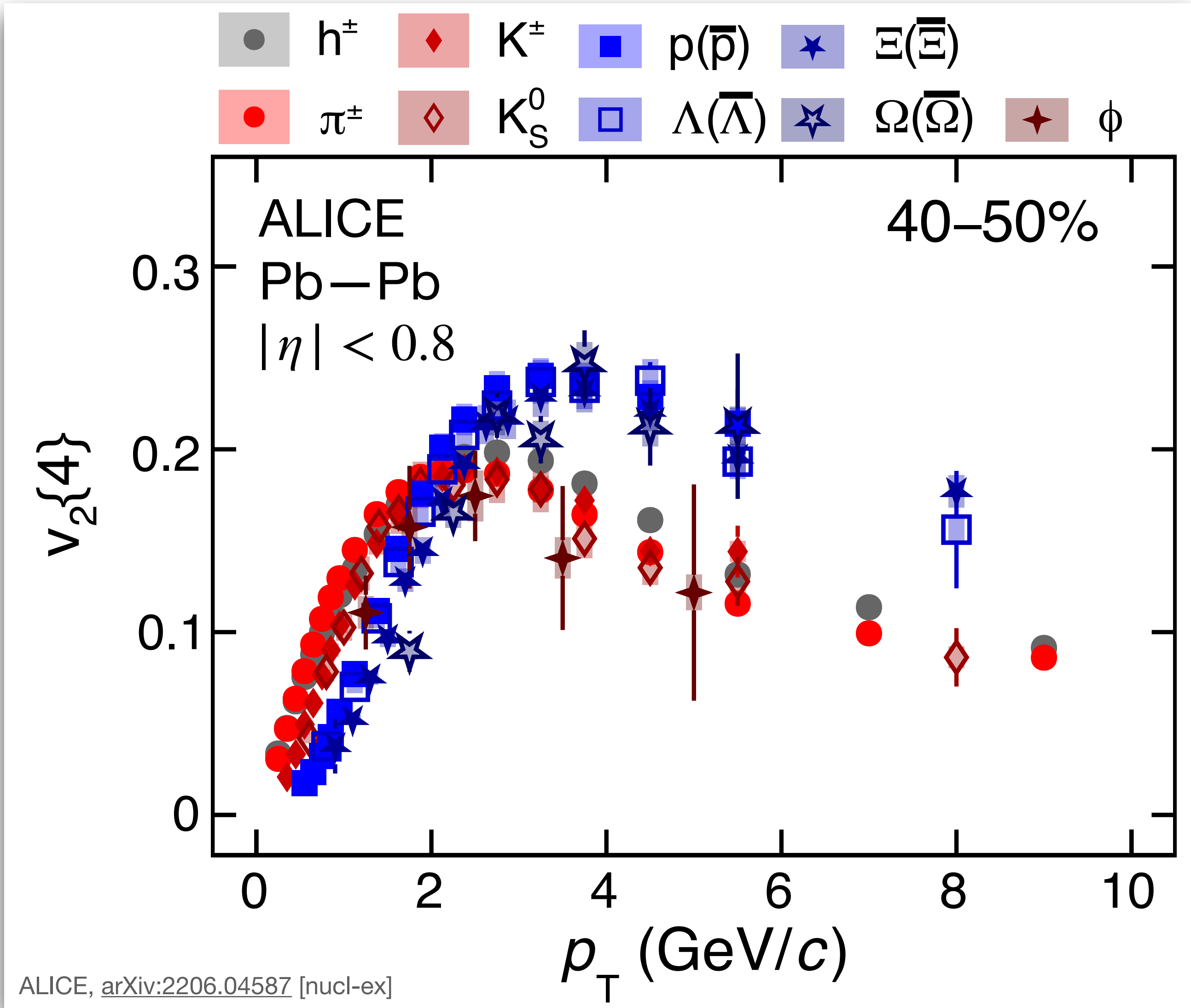
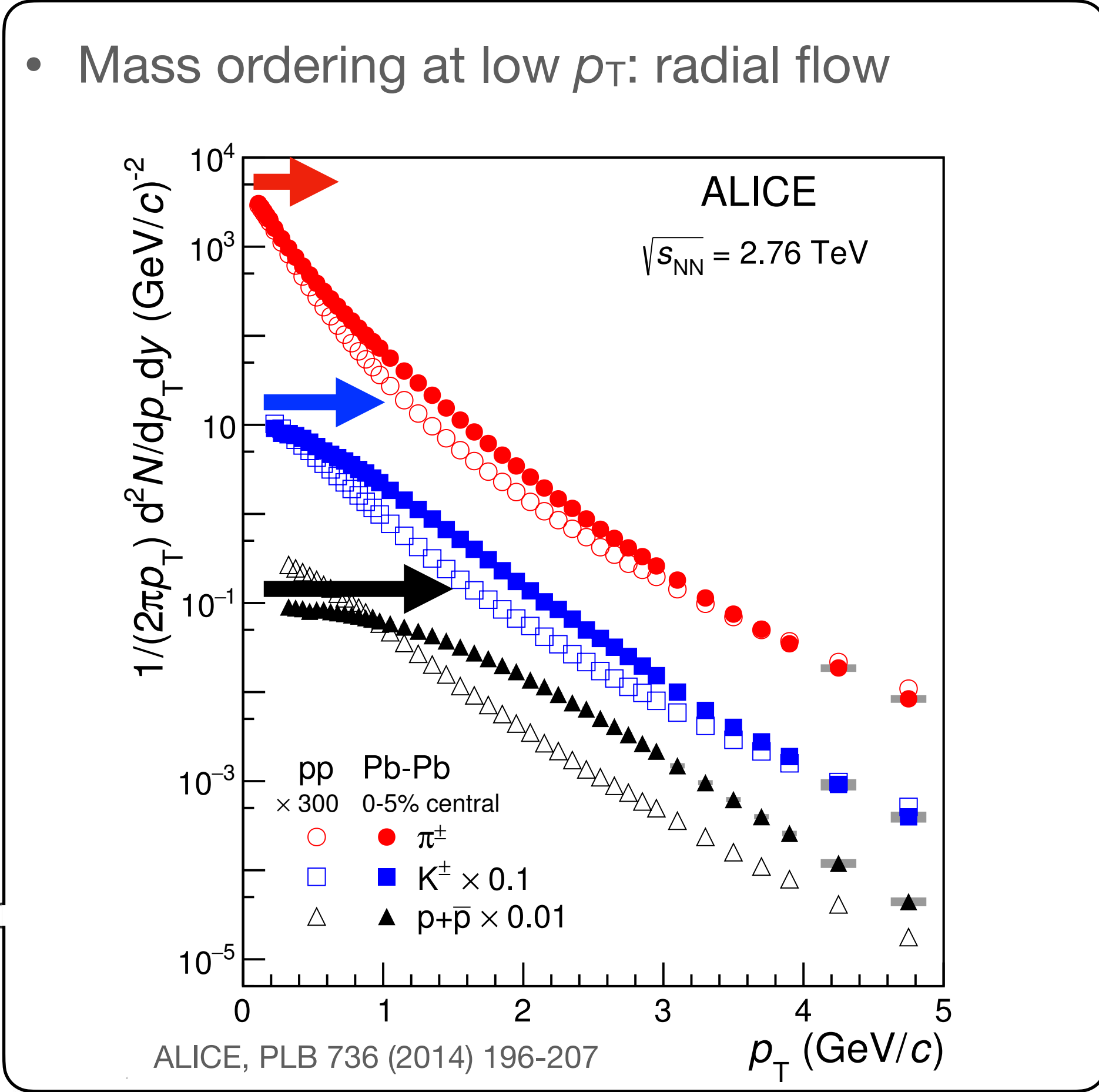
- Low p_T :
 - Inverse mass ordering \Rightarrow expected from hydrodynamics



ALICE, arXiv:2206.04587 [nucl-ex]

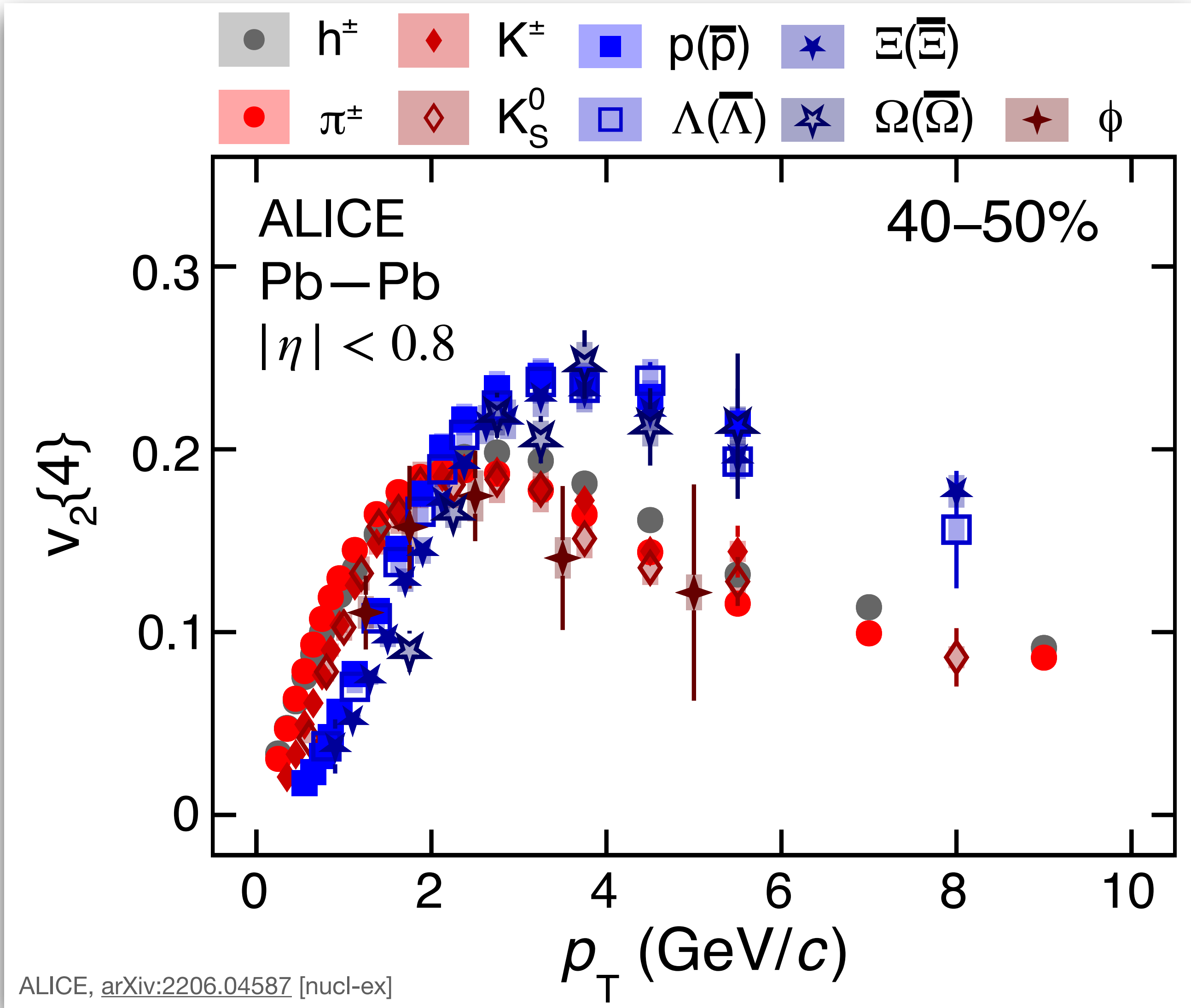
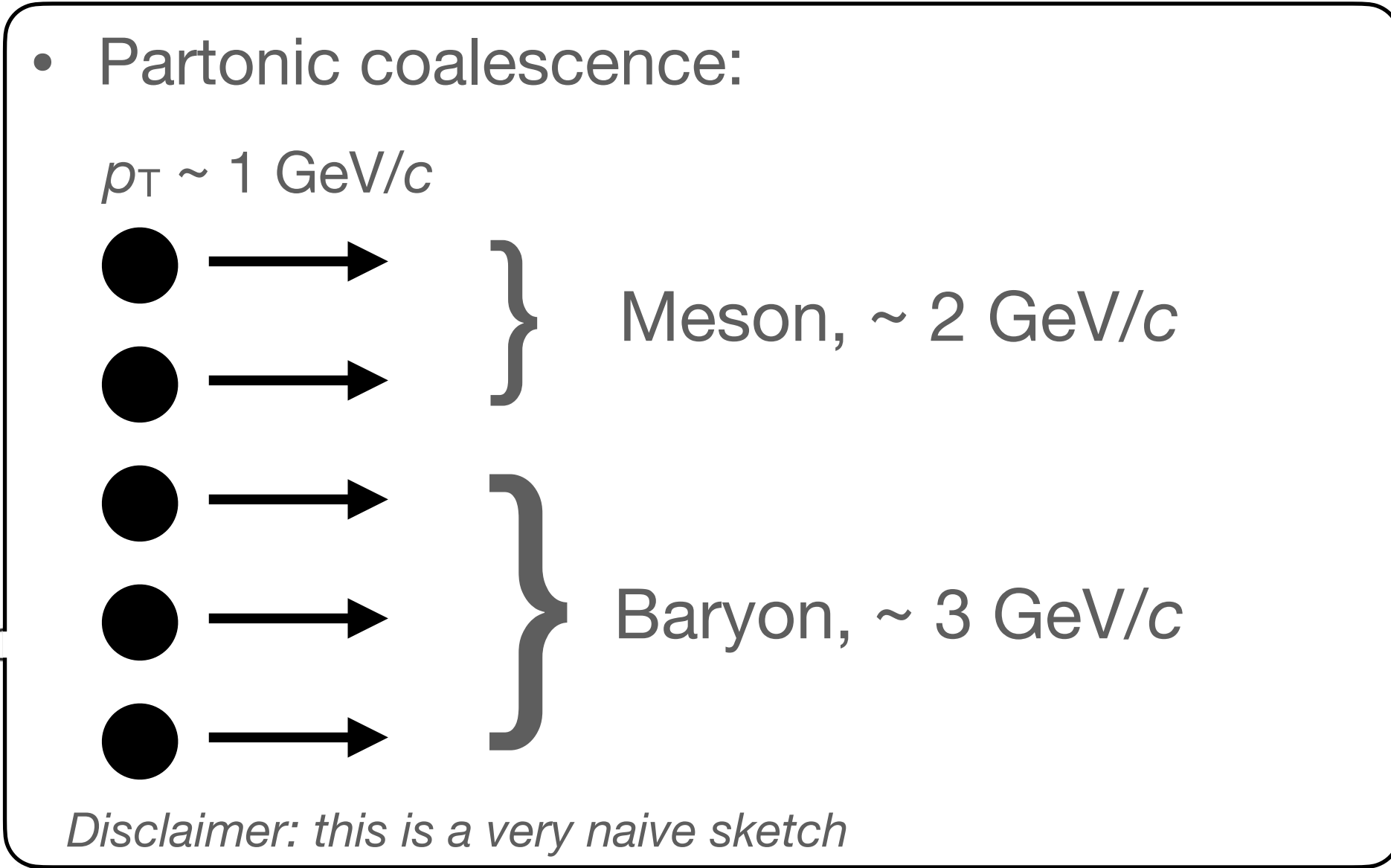
v_2 of baryons and mesons in Pb–Pb

- Low p_T :
 - Inverse mass ordering \Rightarrow expected from hydrodynamics



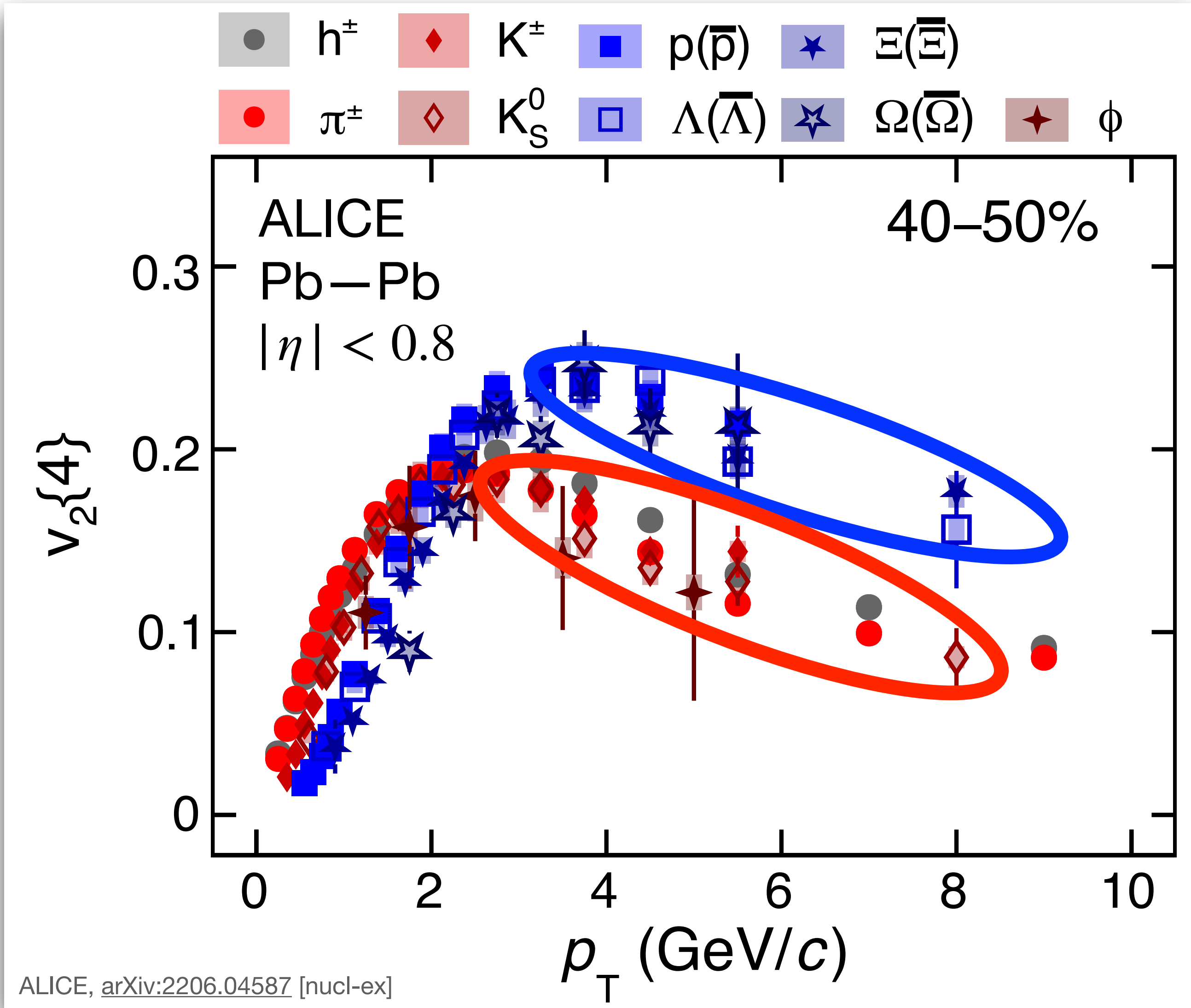
v_2 of baryons and mesons in Pb–Pb

- Low p_T :
 - Inverse mass ordering \Rightarrow expected from hydrodynamics
- Intermediate p_T :
 - Crossing point \Rightarrow radial flow and/or coalescence?



v_2 of baryons and mesons in Pb–Pb

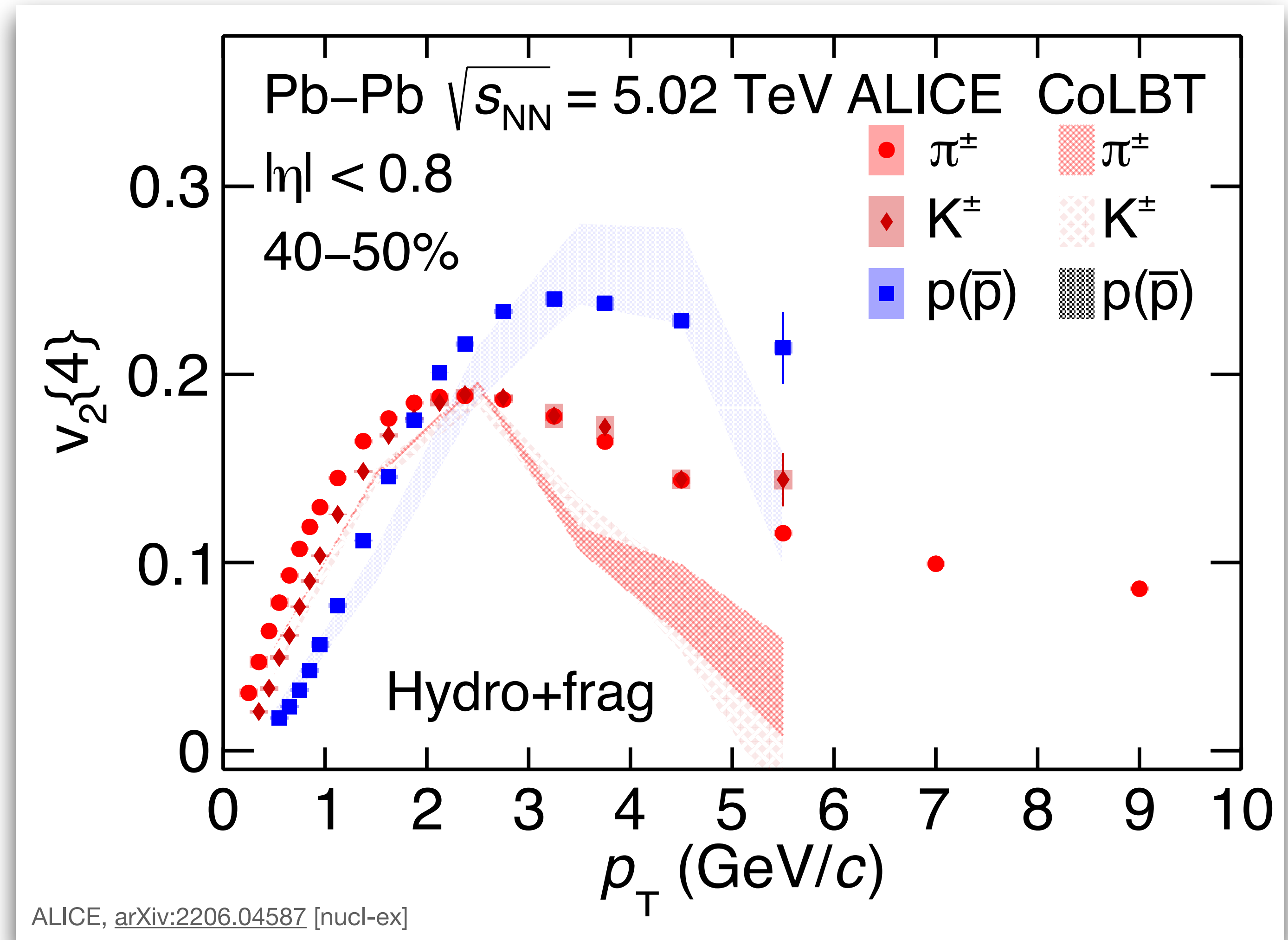
- Low p_T :
 - Inverse mass ordering \Rightarrow expected from hydrodynamics
- Intermediate p_T :
 - Crossing point \Rightarrow radial flow and/or coalescence?
 - **Baryon/meson** grouping \Rightarrow coalescence?
- * High p_T : jet quenching [1]



v_2 of baryons and mesons in Pb–Pb

comparison to hydro model

- Measurements of v_2 :
 - ⇒ Baryon/meson ordering and crossing
- Hydro + fragmentation:
 - ✗ Underestimates the data in most cases
 - ✓ Baryon/meson crossing predicted
 - ⇒ Arises from species-dependent p_T cut, where fragmentation dominates over hydro

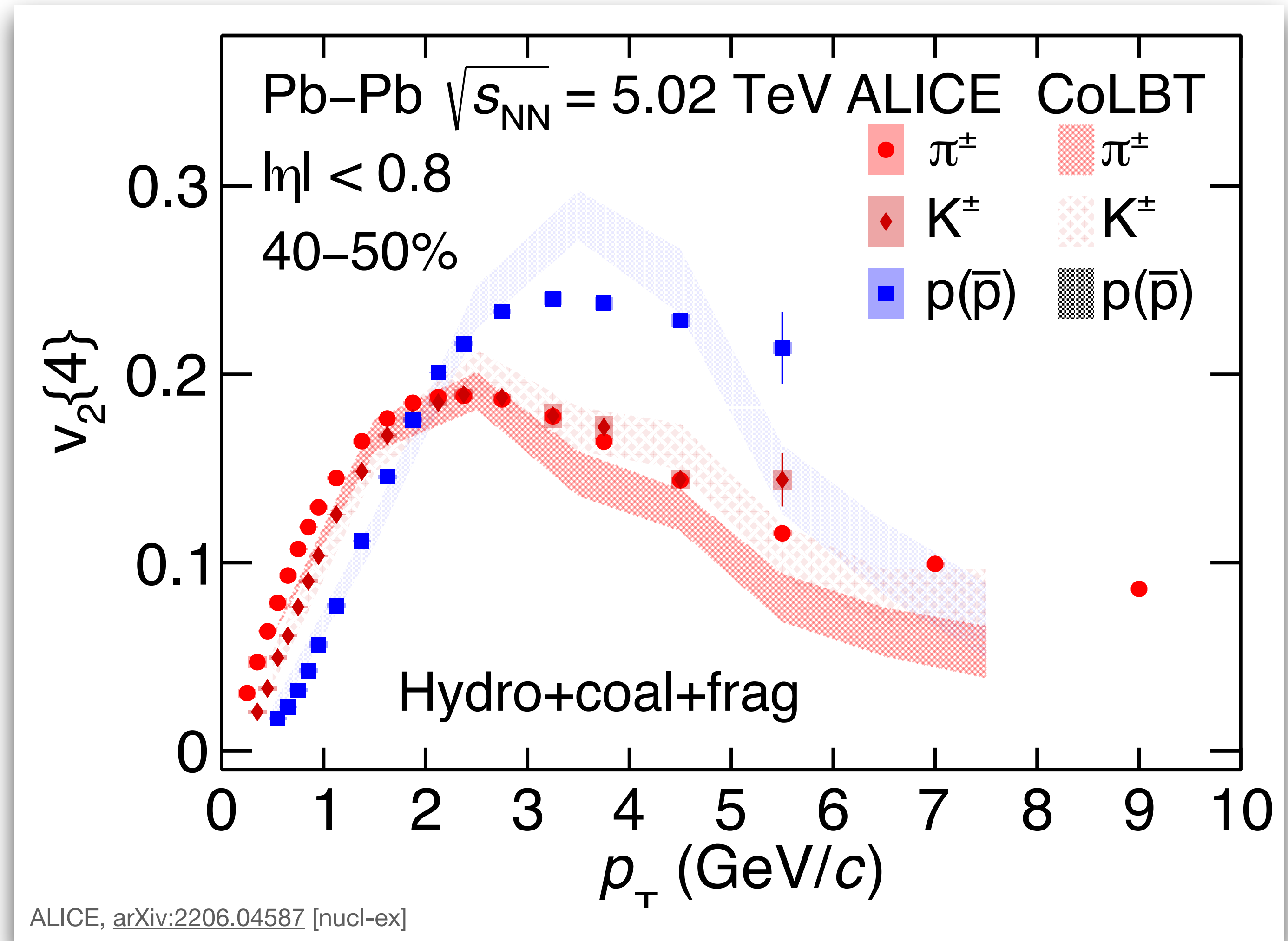


v_2 of baryons and mesons in Pb–Pb

comparison to hydro model

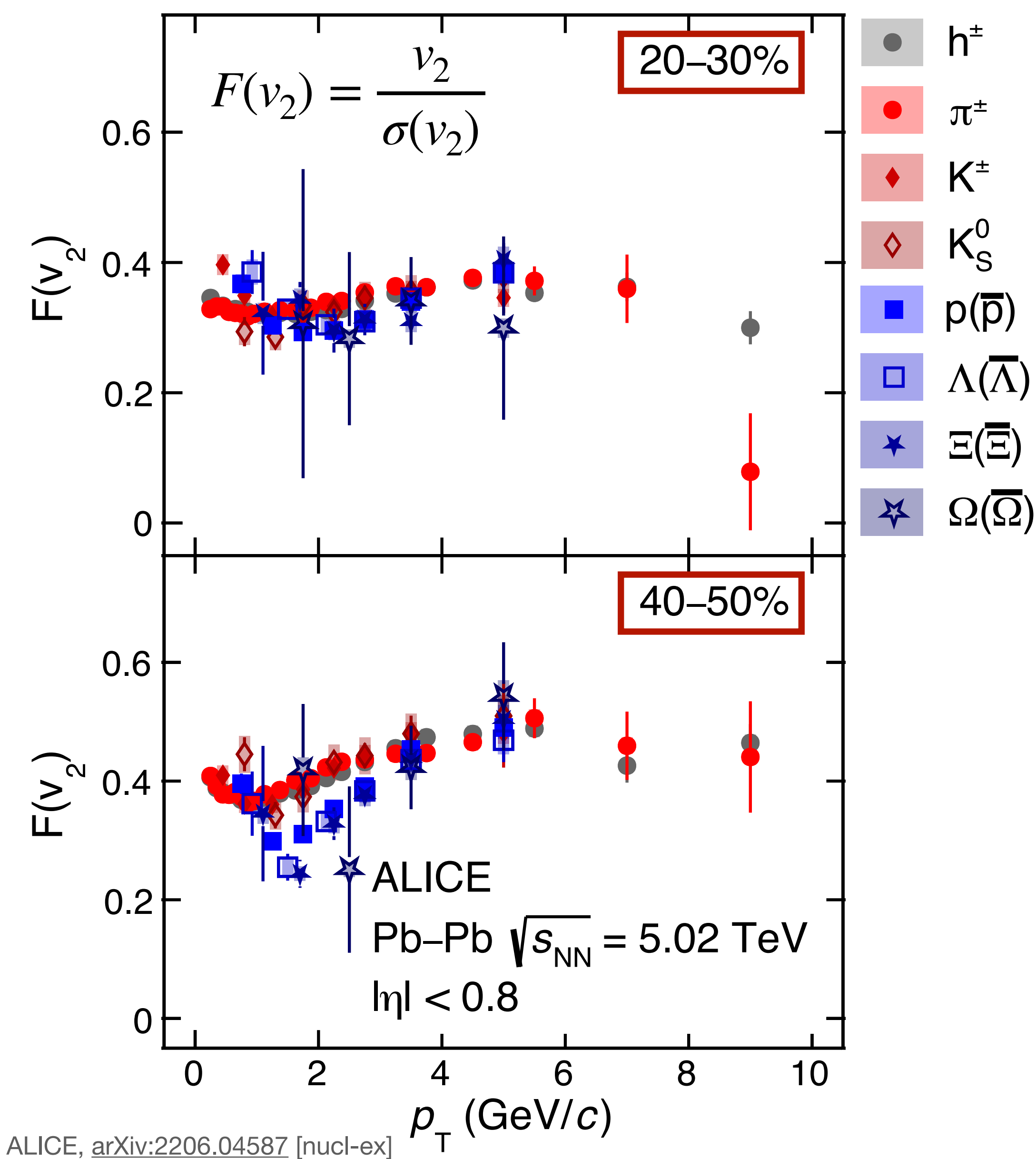
- Measurements of v_2 :
 - ⇒ Baryon/meson ordering and crossing
- Hydro + fragmentation:
 - ✗ Underestimates the data in most cases
 - ✓ Baryon/meson crossing predicted
 - ⇒ Arises from species-dependent p_T cut, where fragmentation dominates over hydro

- Hydro + coalescence + fragmentation:
 - ✓ Significantly better description of data
 - **But crossing is not unique to coalescence!**



v_2 fluctuations of baryons and mesons in Pb—Pb collisions

- Emerging p_T dependence from central to peripheral collisions
- Baryon/meson grouping in semi-central collisions
 - ⇒ Different from that observed for v_2
 - ⇒ Could point to a different origin of this observation



ALICE, arXiv:2206.04587 [nucl-ex]

v_2 fluctuations: skewness and kurtosis in Pb–Pb collisions

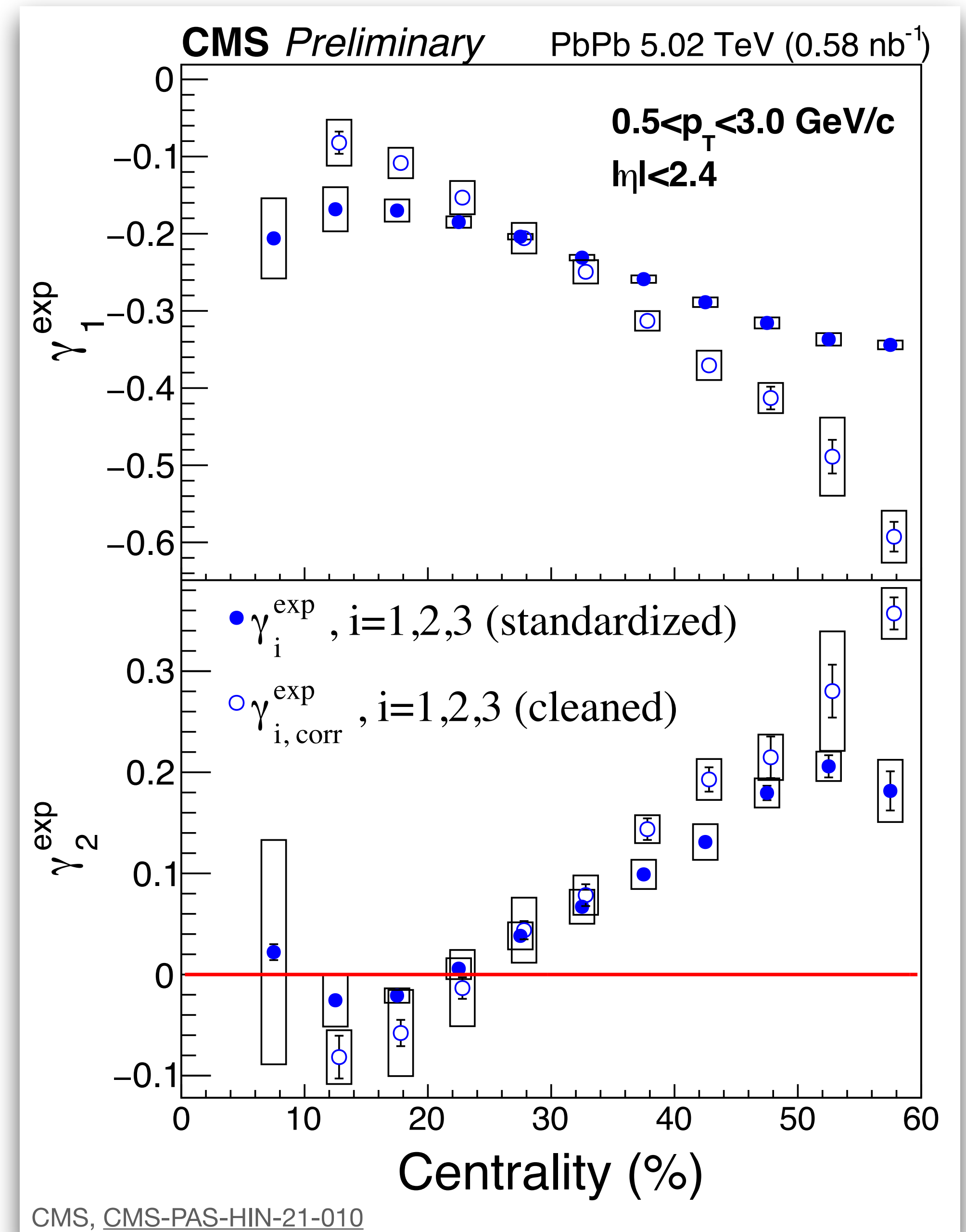
Measure v_2 with multiparticle cumulants:

⇒ Sensitive to underlying v_2 probability

density function (PDF) and thus **initial geometry**

- Skewness (γ_1) decreasing with centrality, PDF becoming less symmetric
- Kurtosis (γ_2) increasing with centrality, tails become “fatter”

Do γ_1, γ_2 probe initial geometry exclusively?



v_2 fluctuations: skewness and kurtosis in Pb–Pb collisions

Measure v_2 with multiparticle cumulants:

⇒ Sensitive to underlying v_2 probability

density function (PDF) and thus **initial geometry**

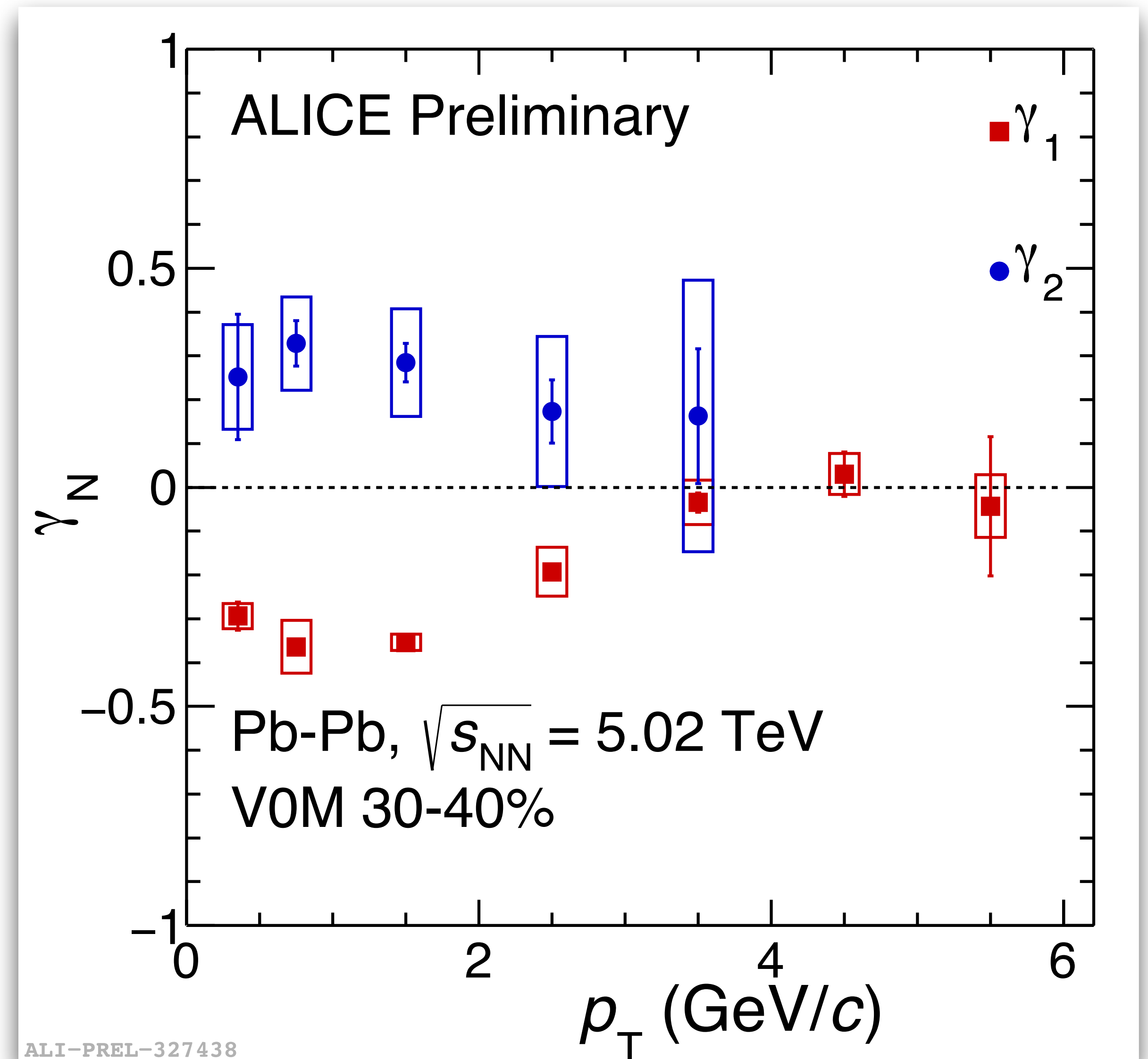
- Skewness (γ_1) decreasing with centrality, PDF becoming less symmetric
- Kurtosis (γ_2) increasing with centrality, tails become “fatter”

Do γ_1 , γ_2 probe initial geometry exclusively?

Not necessarily!

- Both γ_1 and γ_2 show evolution with p_T

⇒ Suggests that v_2 PDF is modified by the evolution of QGP



Is QGP produced exclusively in heavy-ion collisions?

“Standard” paradigm:

- QGP in heavy-ion collisions
- Cold nuclear matter in proton-ion collisions
- Reference in pp

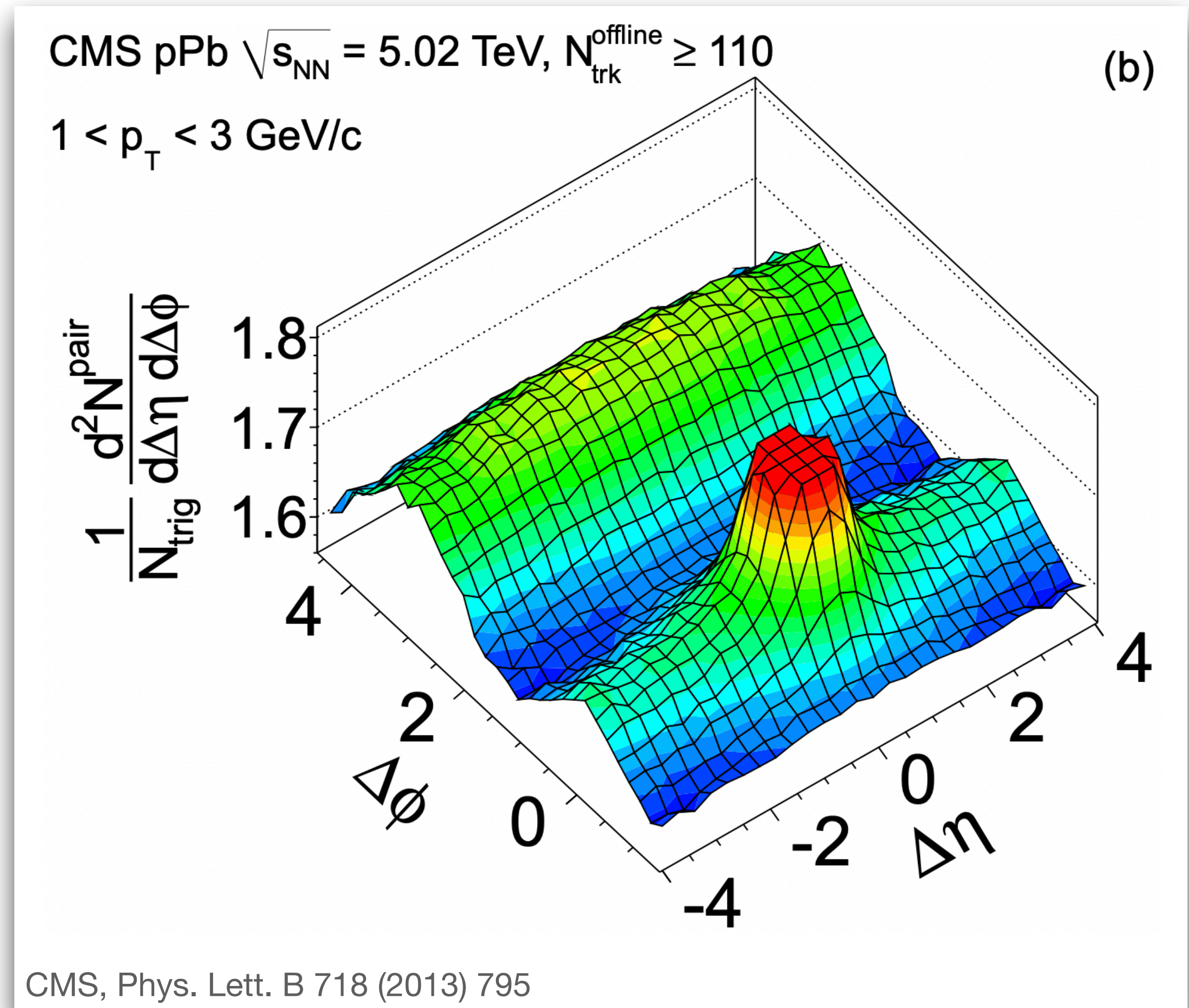
Is QGP produced exclusively in heavy-ion collisions?

“Standard” paradigm:

- QGP in heavy-ion collisions
- Cold nuclear matter in proton-ion collisions
- Reference in pp

Pandora’s box since 2013:

- Near-side ridge in p-Pb, reminiscent of QGP



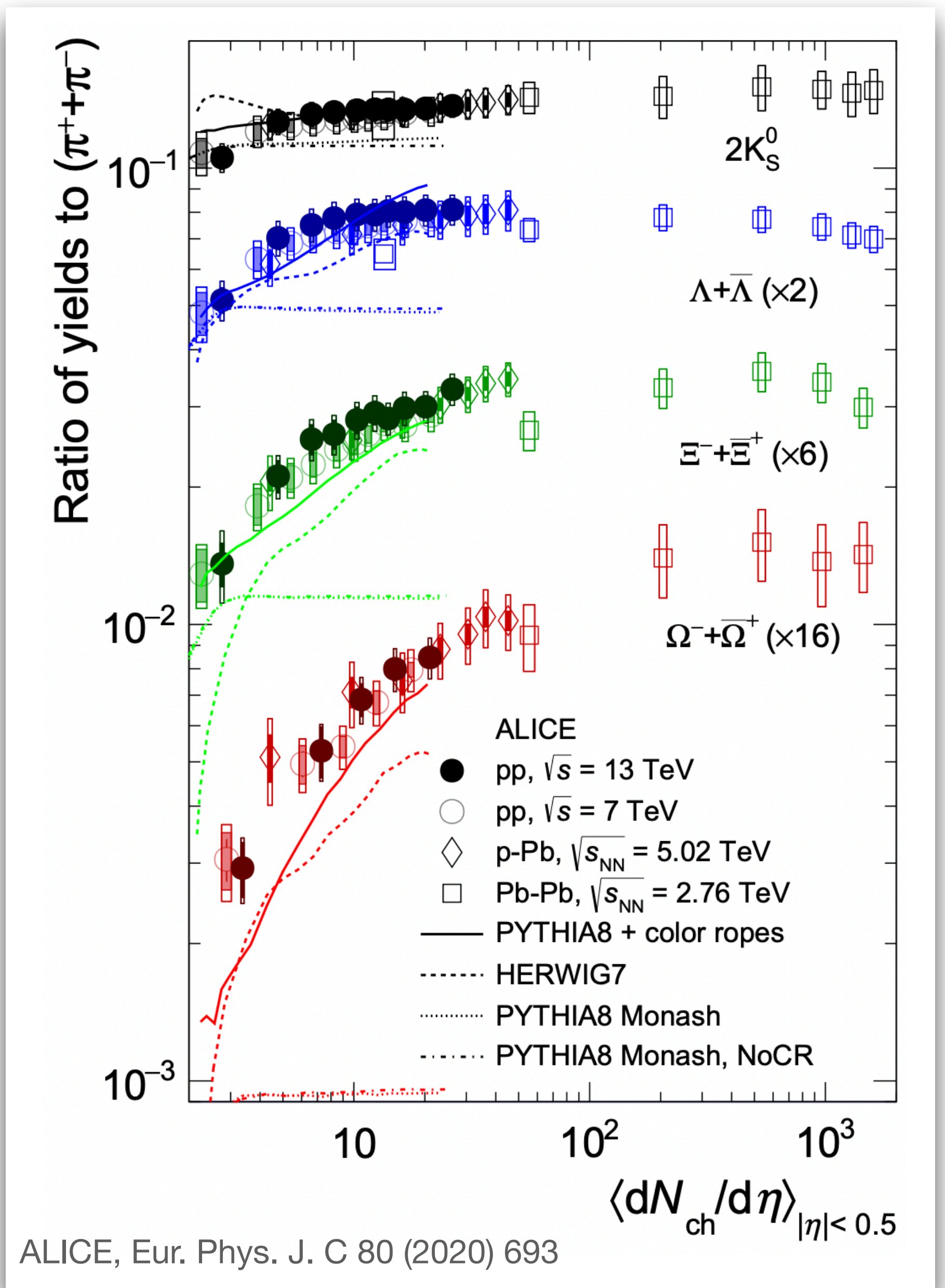
Is QGP produced exclusively in heavy-ion collisions?

“Standard” paradigm:

- QGP in heavy-ion collisions
- Cold nuclear matter in proton-ion collisions
- Reference in pp

Pandora’s box since 2013:

- Near-side ridge in p-Pb, reminiscent of QGP
- Particle production mechanism in pp and p–Pb similar to that in Pb–Pb



Is QGP produced exclusively in heavy-ion collisions?

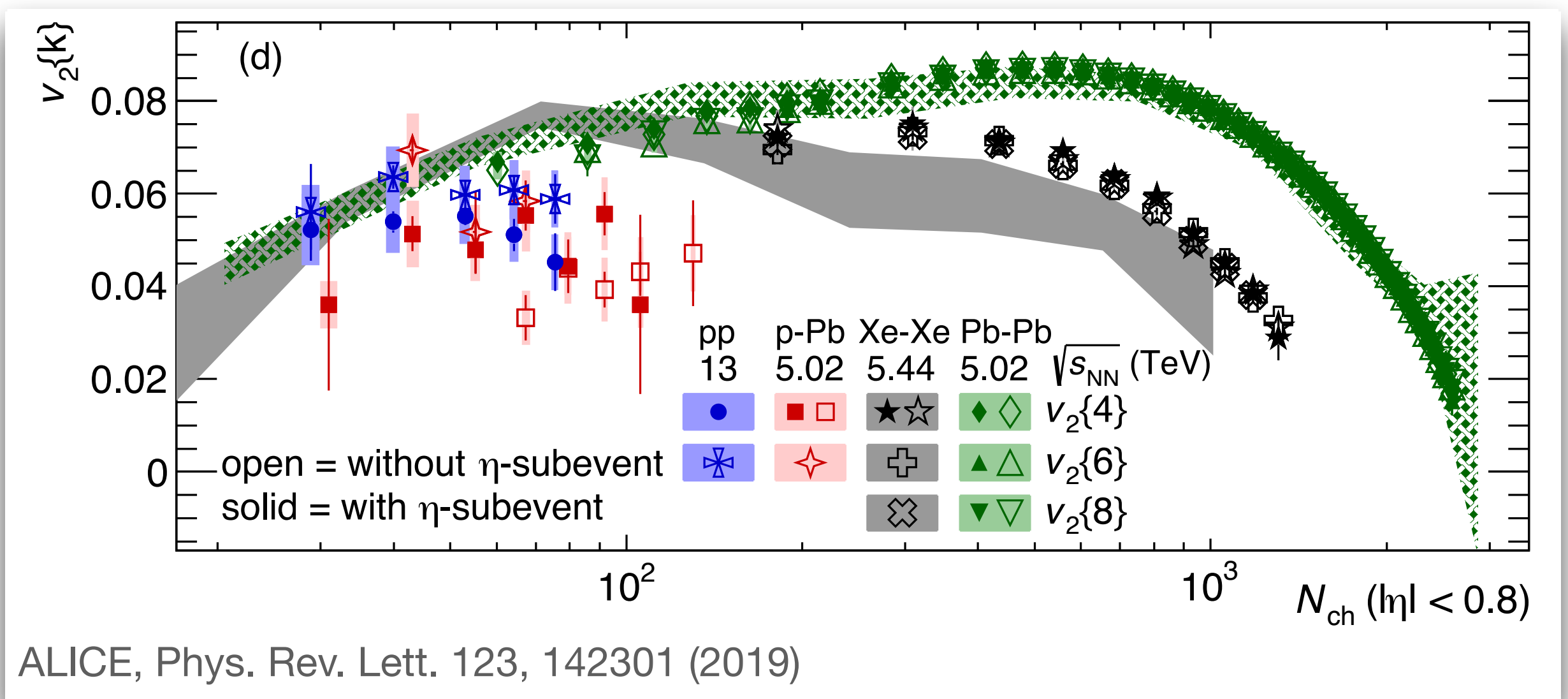
“Standard” paradigm:

- QGP in heavy-ion collisions
- Cold nuclear matter in proton-ion collisions
- Reference in pp

Pandora’s box since 2013:

- Near-side ridge in p-Pb, reminiscent of QGP
- Particle production mechanism in pp and p–Pb similar to that in Pb–Pb
- Non-zero v_2 coefficients

Is v_2 measured in pp and p–Pb collisions QGP-like?
Study it in p_T -differential way



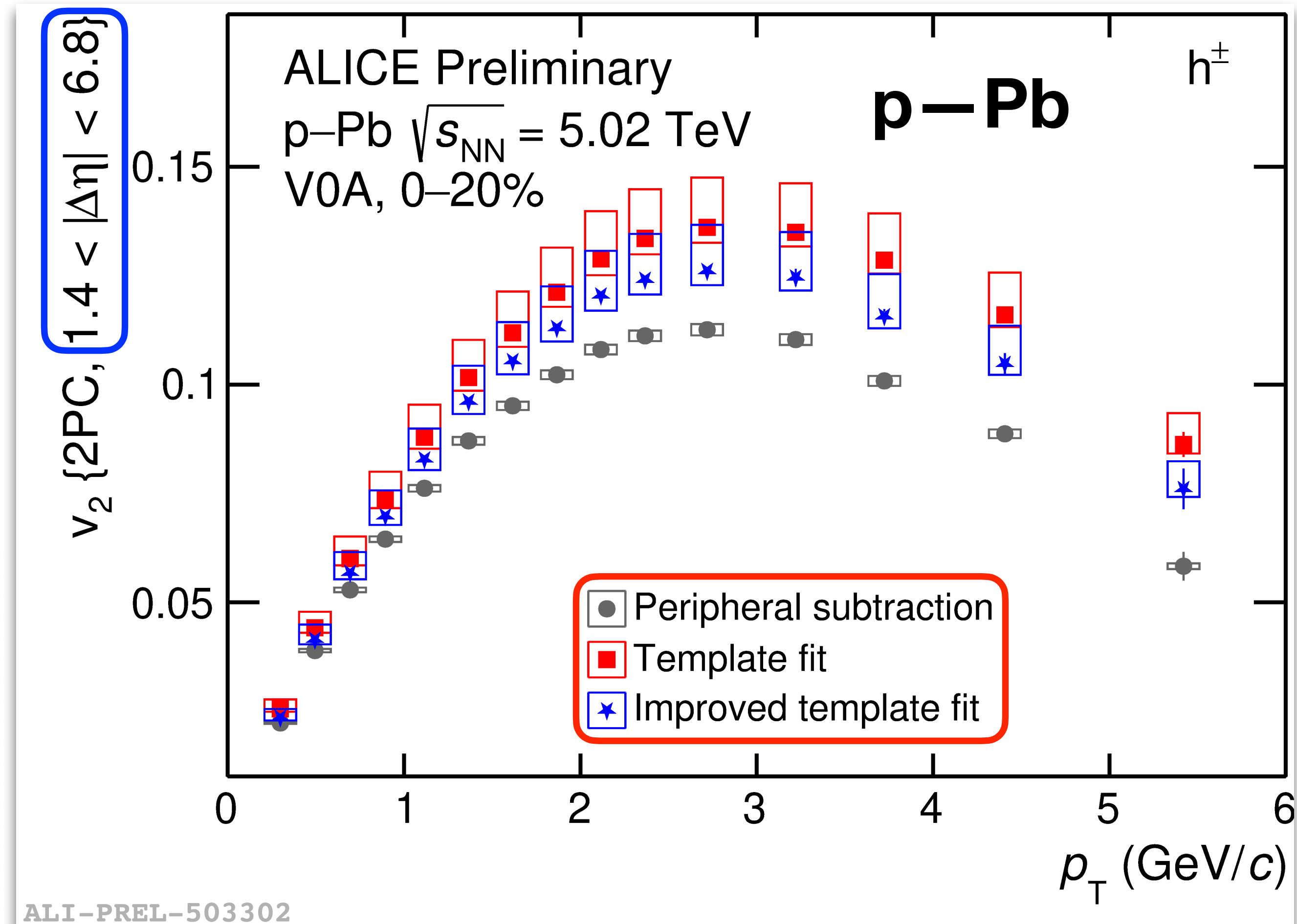
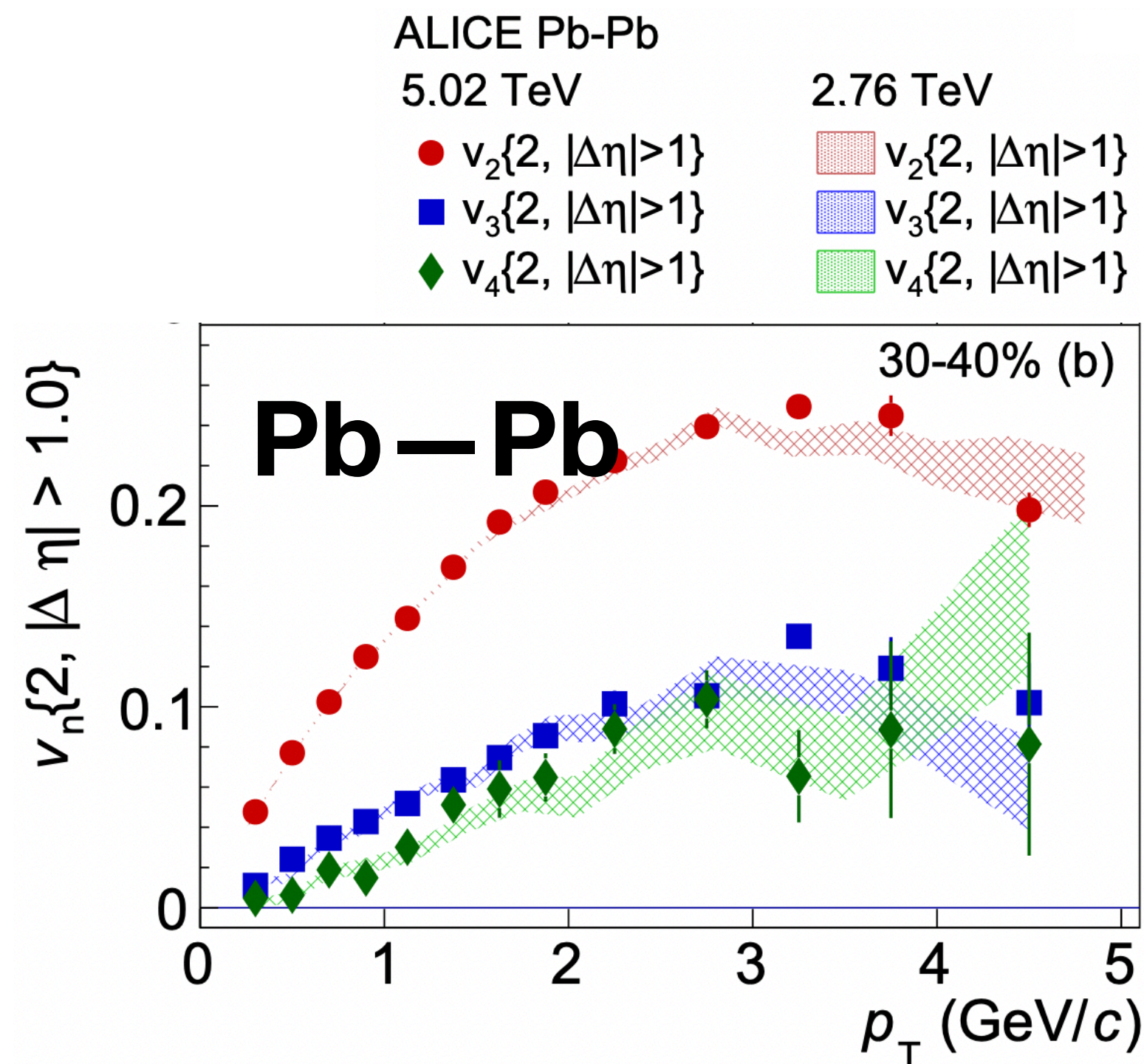
Anisotropic flow in p–Pb collisions for charged hadrons

- Small systems dominated by non-flow

Suppress by:

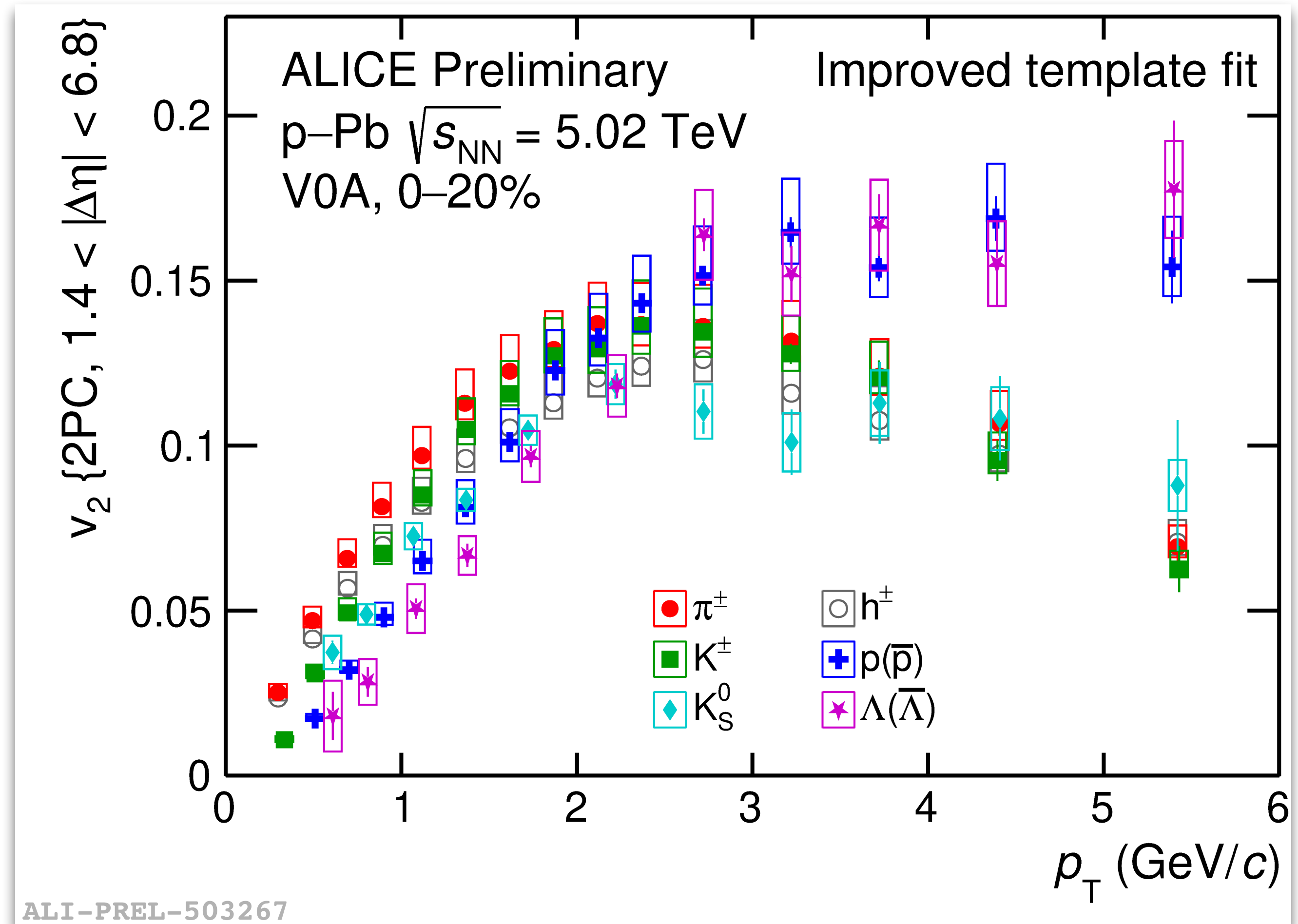
- Large $|\Delta\eta|$ gap
- Subtraction of low multiplicity collisions

- Qualitatively similar trends in Pb–Pb and p–Pb



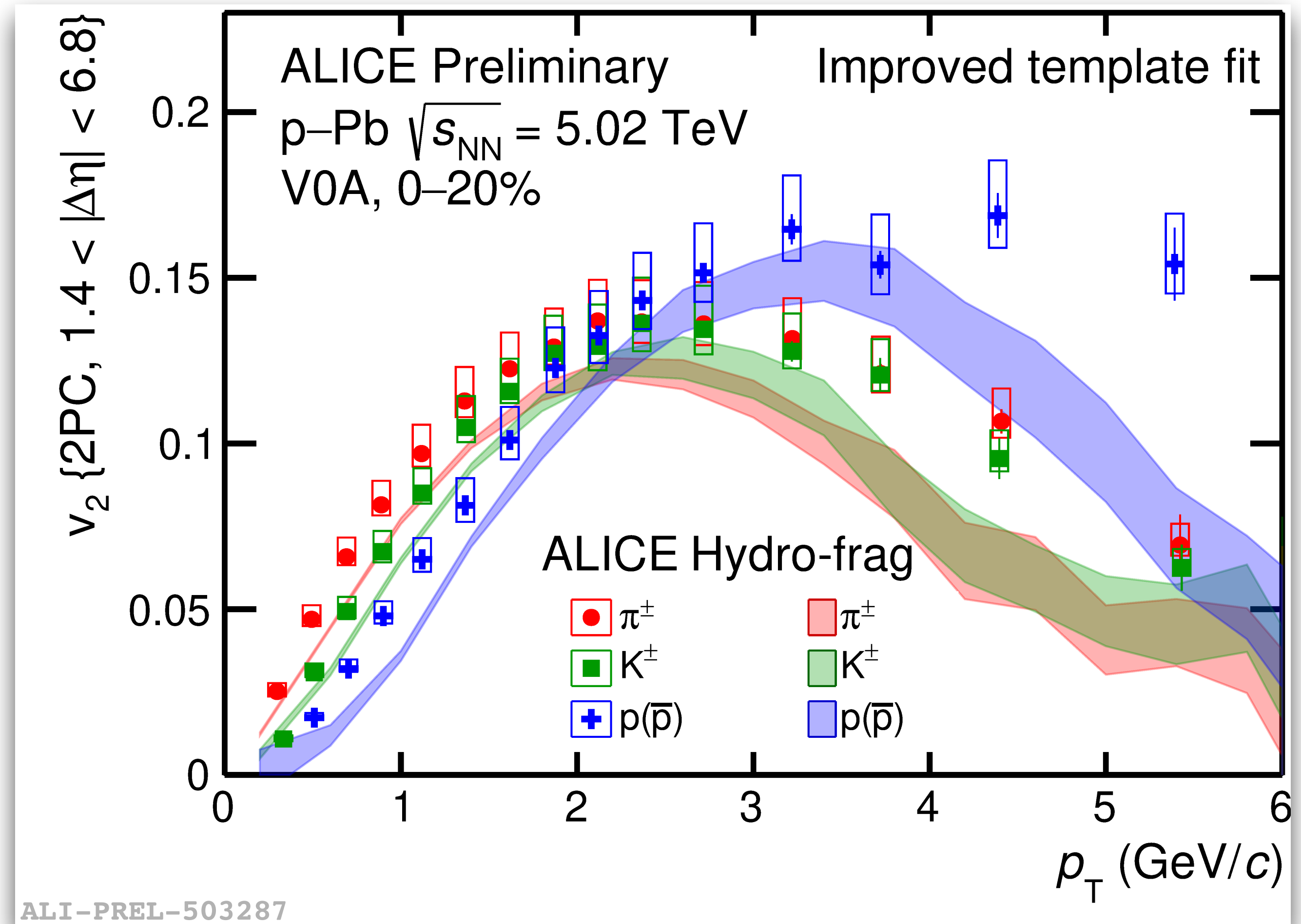
Anisotropic flow in p–Pb collisions for identified particles

- p_T -differential measurements of v_2 in p–Pb collisions for identified particles:
 - Inverse mass ordering at low p_T
 - Baryon/meson grouping at intermediate p_T with 3σ



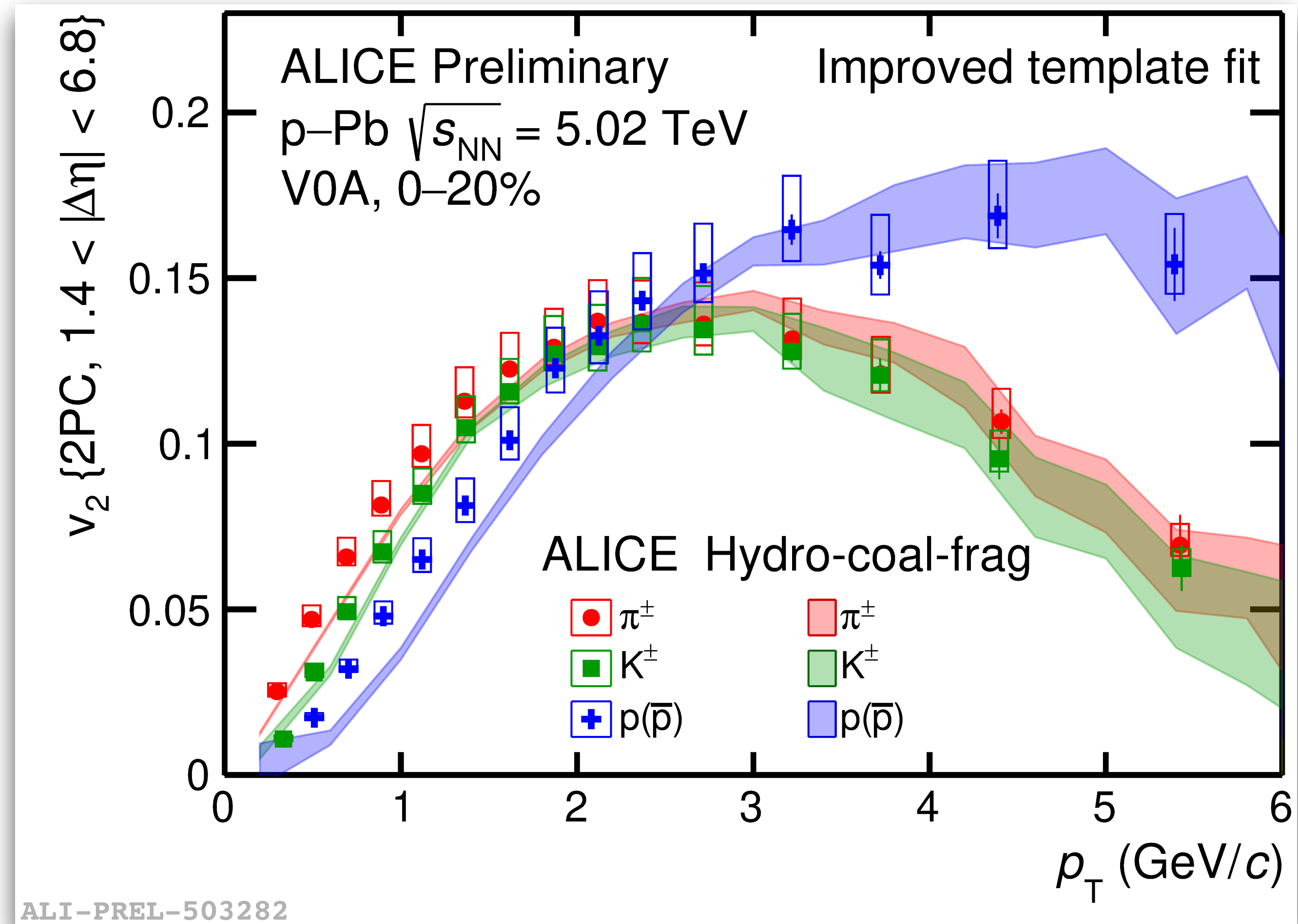
Anisotropic flow in p–Pb collisions for identified particles

- p_T -differential measurements of v_2 in p–Pb collisions for identified particles:
 - Inverse mass ordering at low p_T
 - Baryon/meson grouping at intermediate p_T with 3σ
- Comparison to models:
 - Hydro + fragmentation: largely underestimates data at intermediate p_T



Anisotropic flow in p–Pb collisions for identified particles

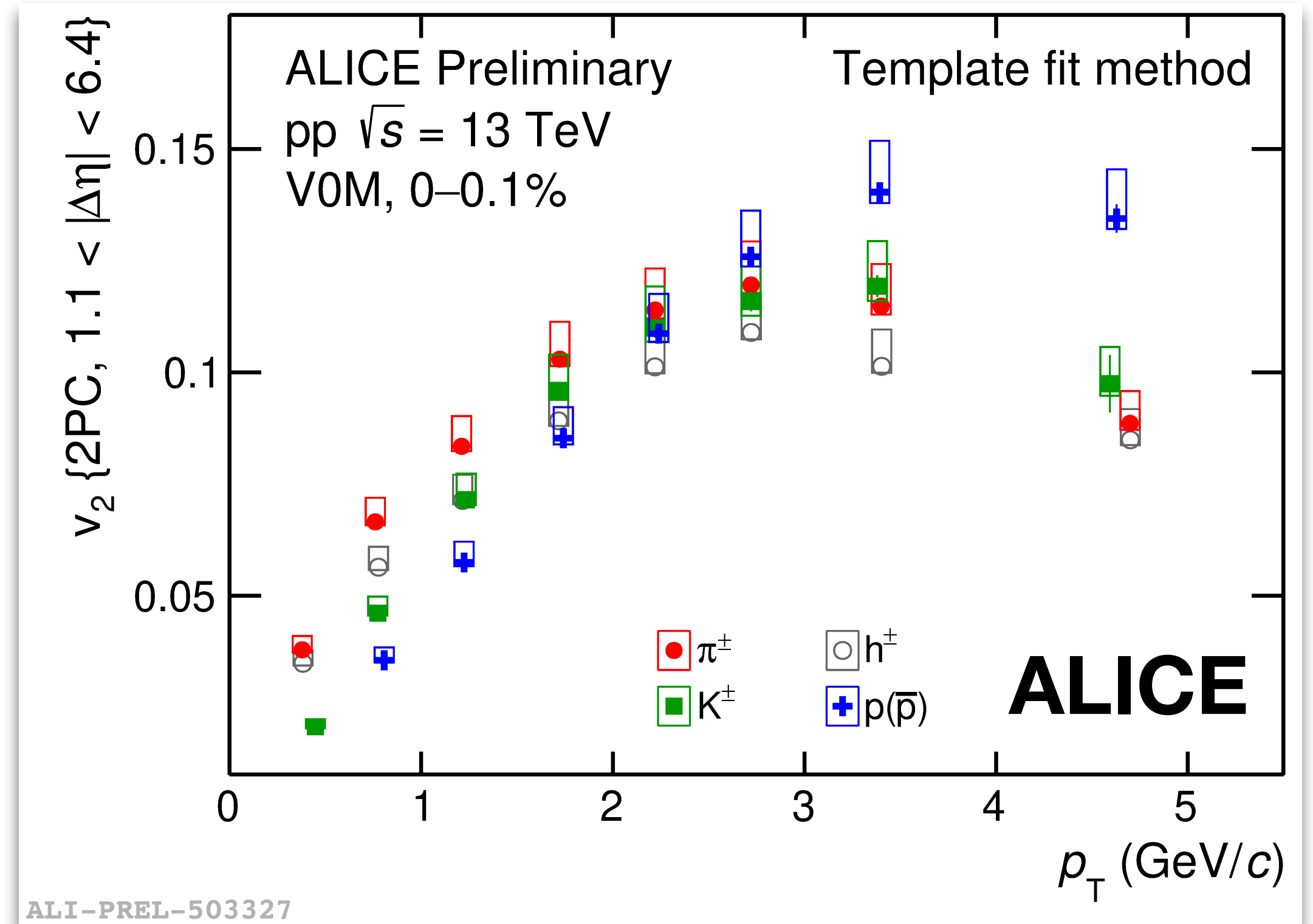
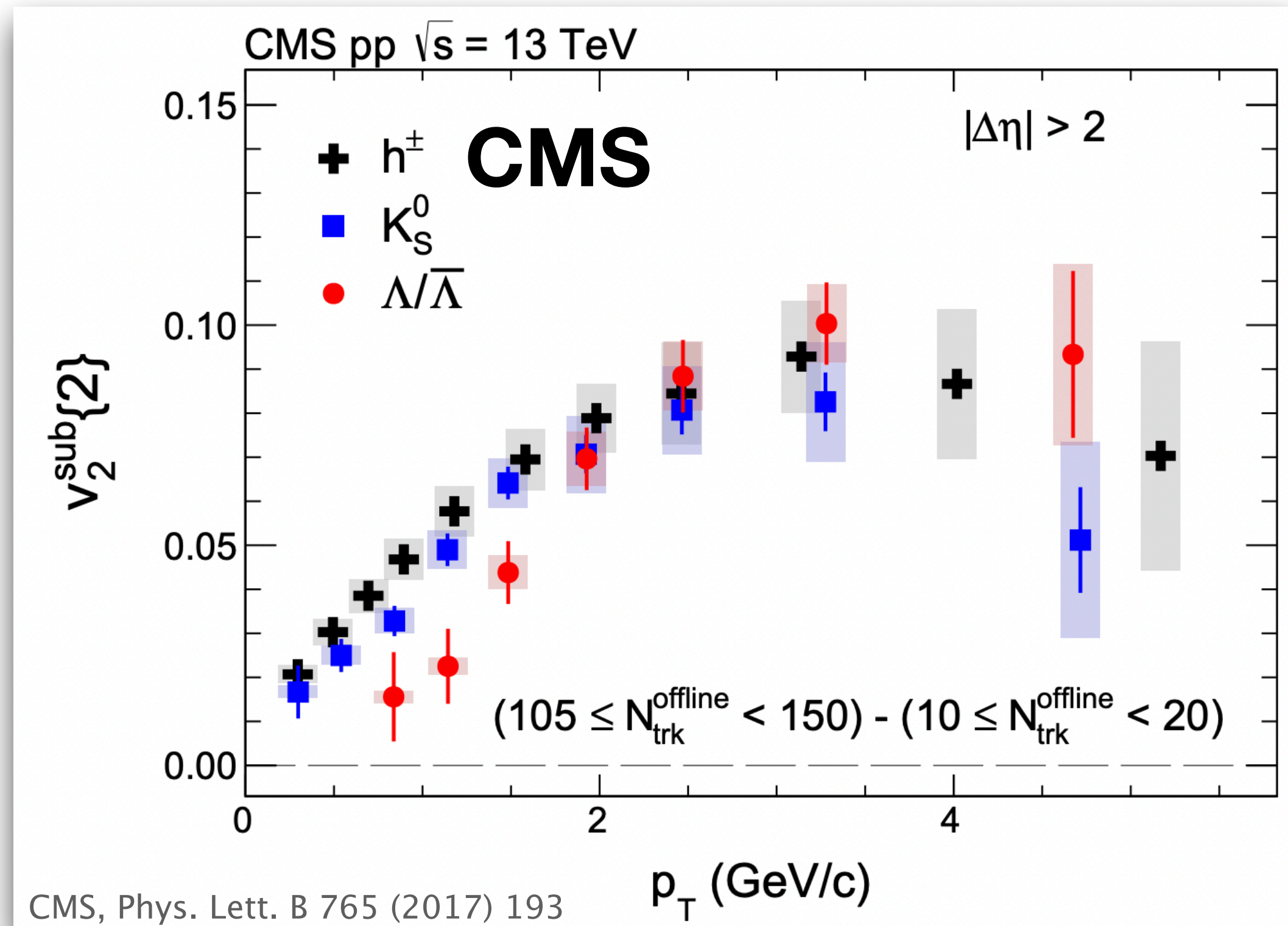
- p_T -differential measurements of v_2 in p–Pb collisions for identified particles:
 - Inverse mass ordering at low p_T
 - Baryon/meson grouping at intermediate p_T with 3σ
 - Comparison to models:
 - Hydro + fragmentation: largely underestimates data at intermediate p_T
 - Hydro + coalescence + fragmentation: good description of data at intermediate p_T
- ⇒ Presence of partonic collectivity in p–Pb



Anisotropic flow in pp collisions for identified particles

- Inverse mass ordering at low p_T
- Clear separation ($> 3\sigma$) between protons and π/K at intermediate p_T !

Hydrodynamics—like trends observed
in pp and p—Pb. Are they of the same origin?
⇒ Need to study the initial state



Correlation between $[p_T]$ and v_2

- Shape of the fireball: anisotropic flow, $\varepsilon_n \rightarrow v_n$
 - Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$
 - Initial state: geometry and fluctuations of **shape** and **size**
 - Final state: correlation between v_n and $[p_T]$
- ⇒ Study with Pearson correlation coefficient:

$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$

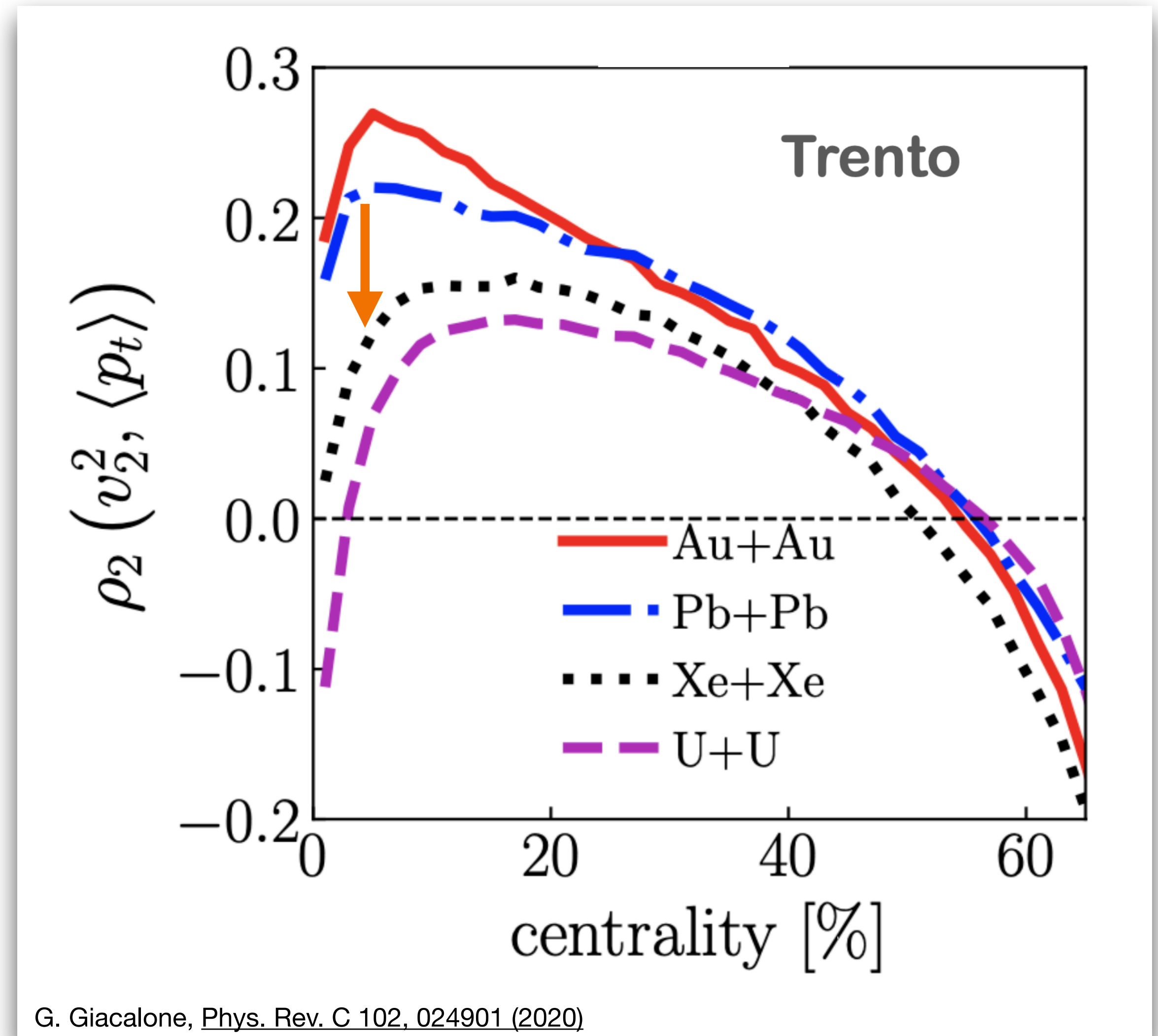
Correlation between $[p_T]$ and v_2 and deformation of nuclei

- Shape of the fireball: anisotropic flow, $\varepsilon_n \rightarrow v_n$
- Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$
- Initial state: geometry and fluctuations of **shape** and **size**
- Final state: correlation between v_n and $[p_T]$

⇒ Study with Pearson correlation coefficient:

$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$

⇒ ρ_2 is significantly smaller in central collisions of deformed **Xe** nuclei (deformation parameter $\beta_2 \approx 0.16$) compared to spherical **Pb** ($\beta_2 \approx 0$)



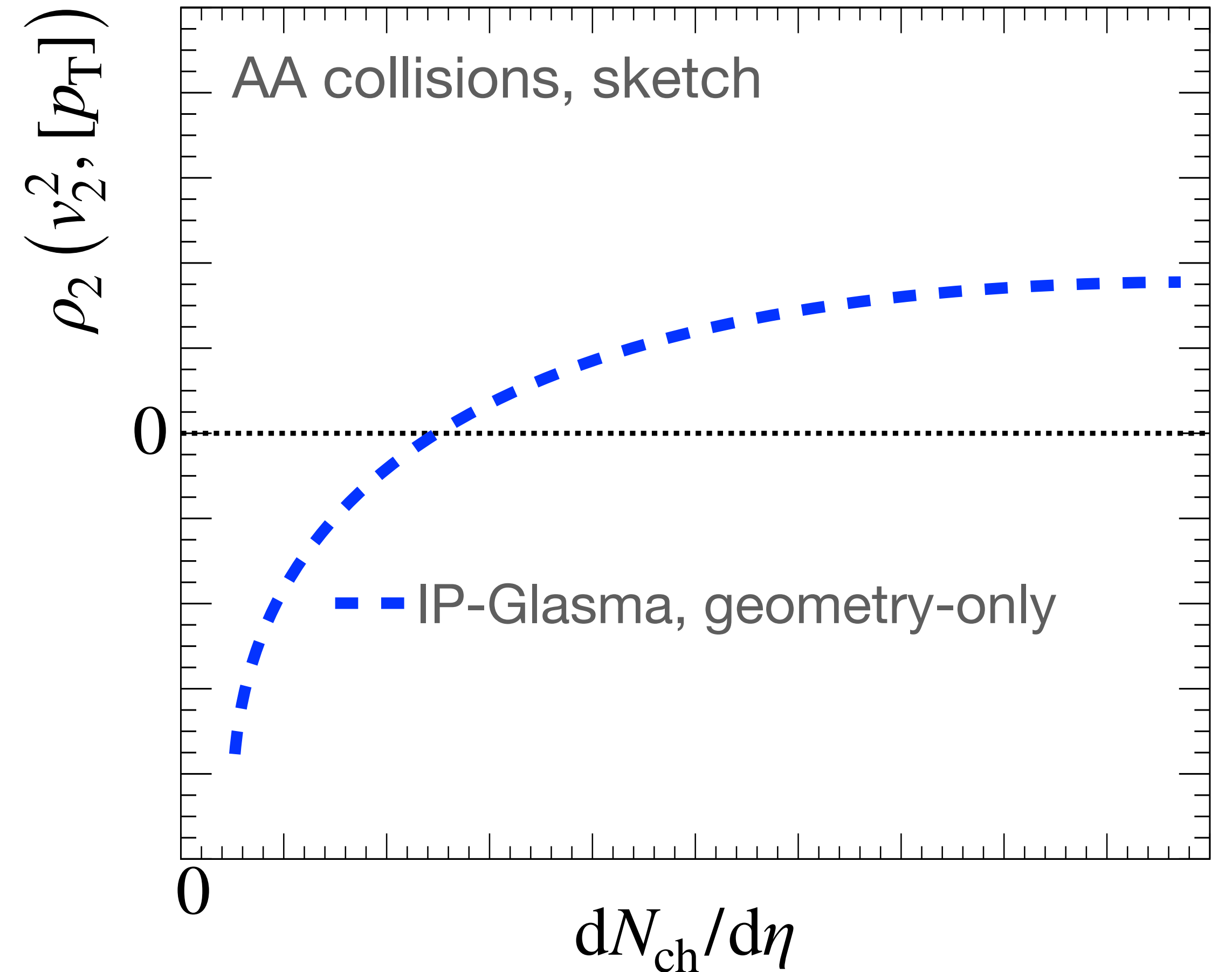
G. Giacalone, Phys. Rev. C 102, 024901 (2020)

Correlation between $[p_T]$ and v_2 at low multiplicity

- Shape of the fireball: anisotropic flow, $\varepsilon_n \rightarrow v_n$
- Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$
- Initial state: geometry and fluctuations of **shape** and **size**
- Final state: correlation between v_n and $[p_T]$

⇒ Study with Pearson correlation coefficient:

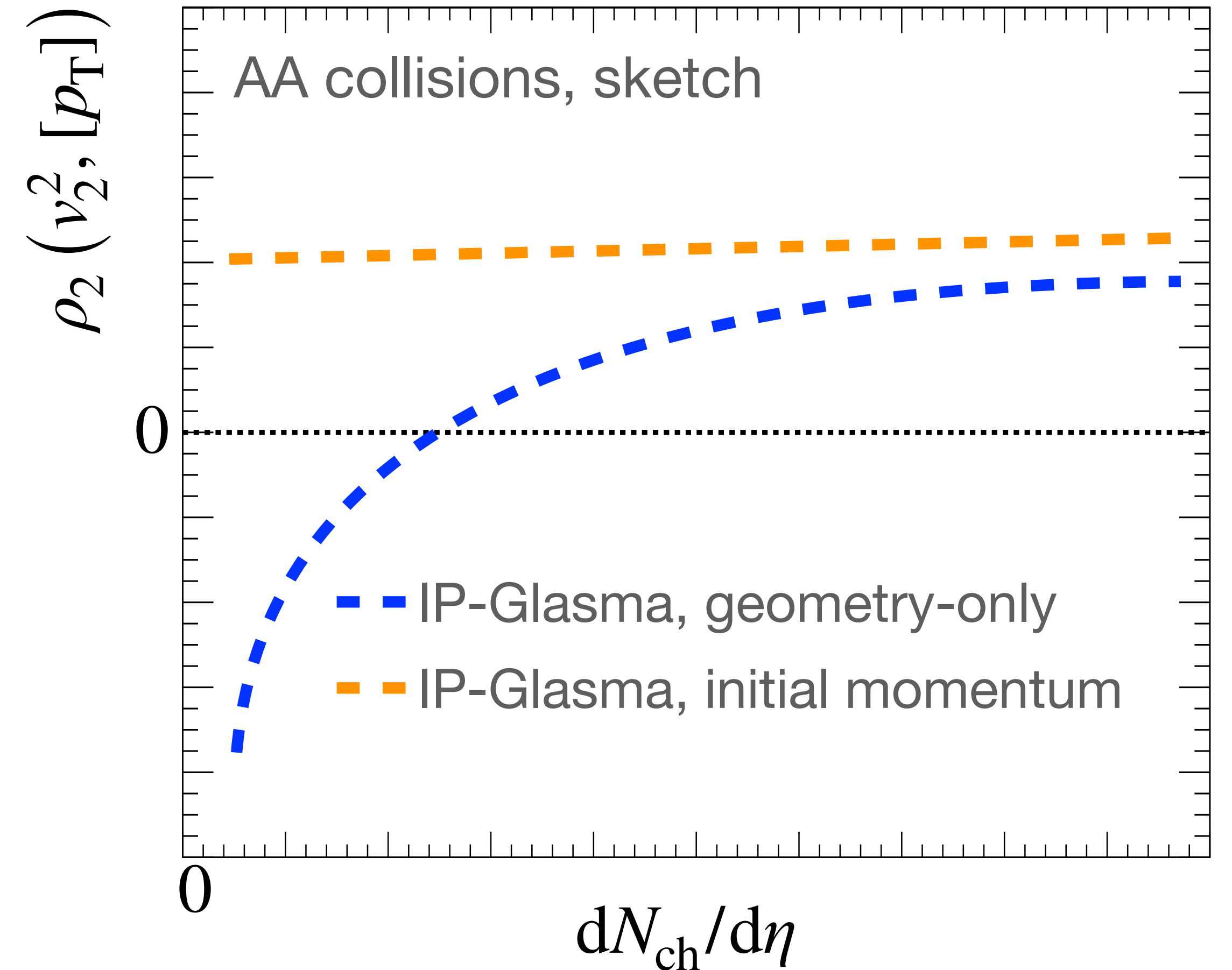
$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$



Correlation between $[p_T]$ and v_2 at low multiplicity

- Shape of the fireball: anisotropic flow, $\varepsilon_n \rightarrow v_n$
 - Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$
 - Initial state: geometry and fluctuations of **shape** and **size**
 - Final state: correlation between v_n and $[p_T]$
- ⇒ Study with Pearson correlation coefficient:

$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$



Correlation between $[p_T]$ and v_2 at low multiplicity

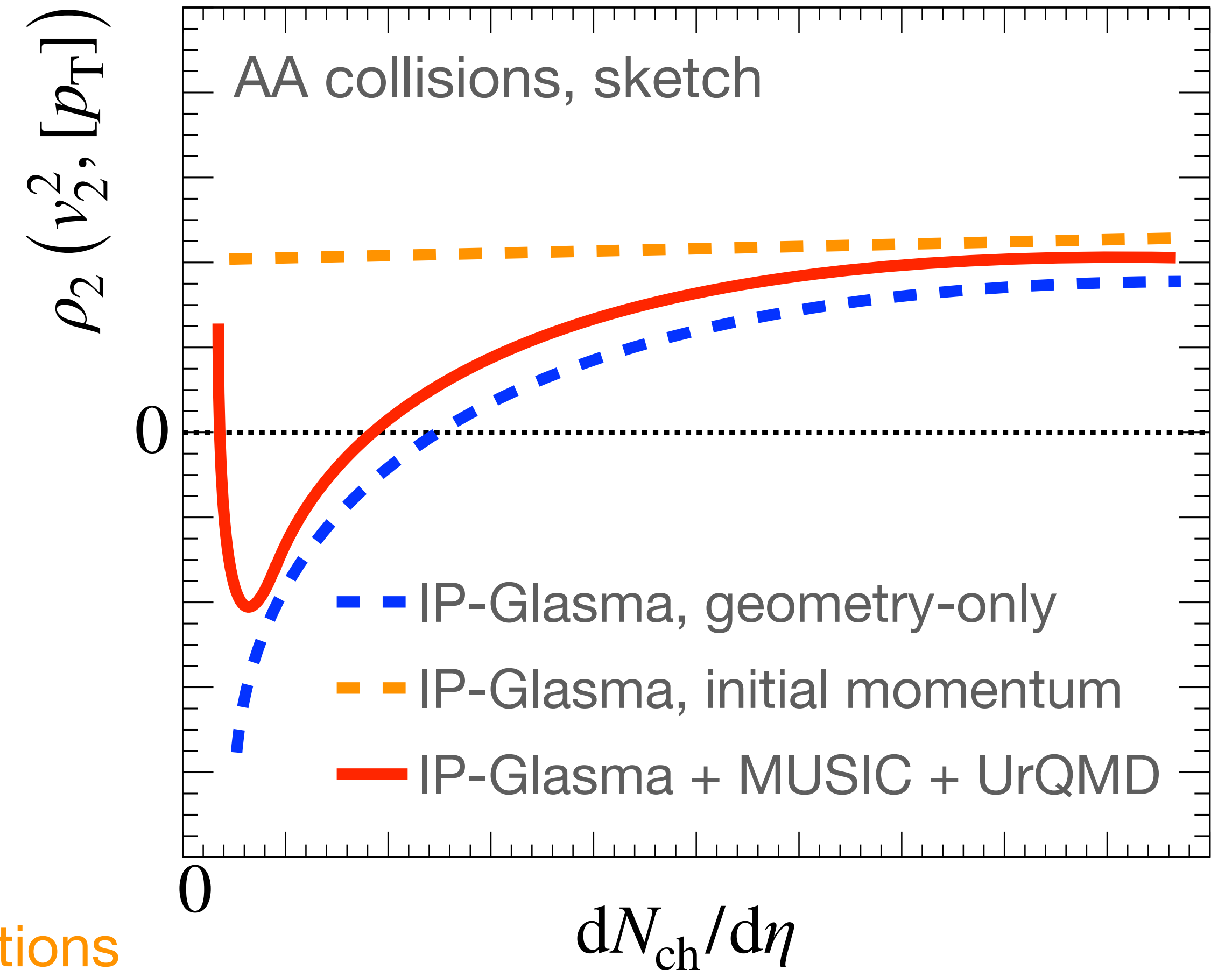
- Shape of the fireball: anisotropic flow, $\varepsilon_n \rightarrow v_n$
- Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$
- Initial state: geometry and fluctuations of **shape** and **size**
- Final state: correlation between v_n and $[p_T]$

⇒ Study with Pearson correlation coefficient:

$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$

Low multiplicity: **geometry** → **initial momentum correlations**

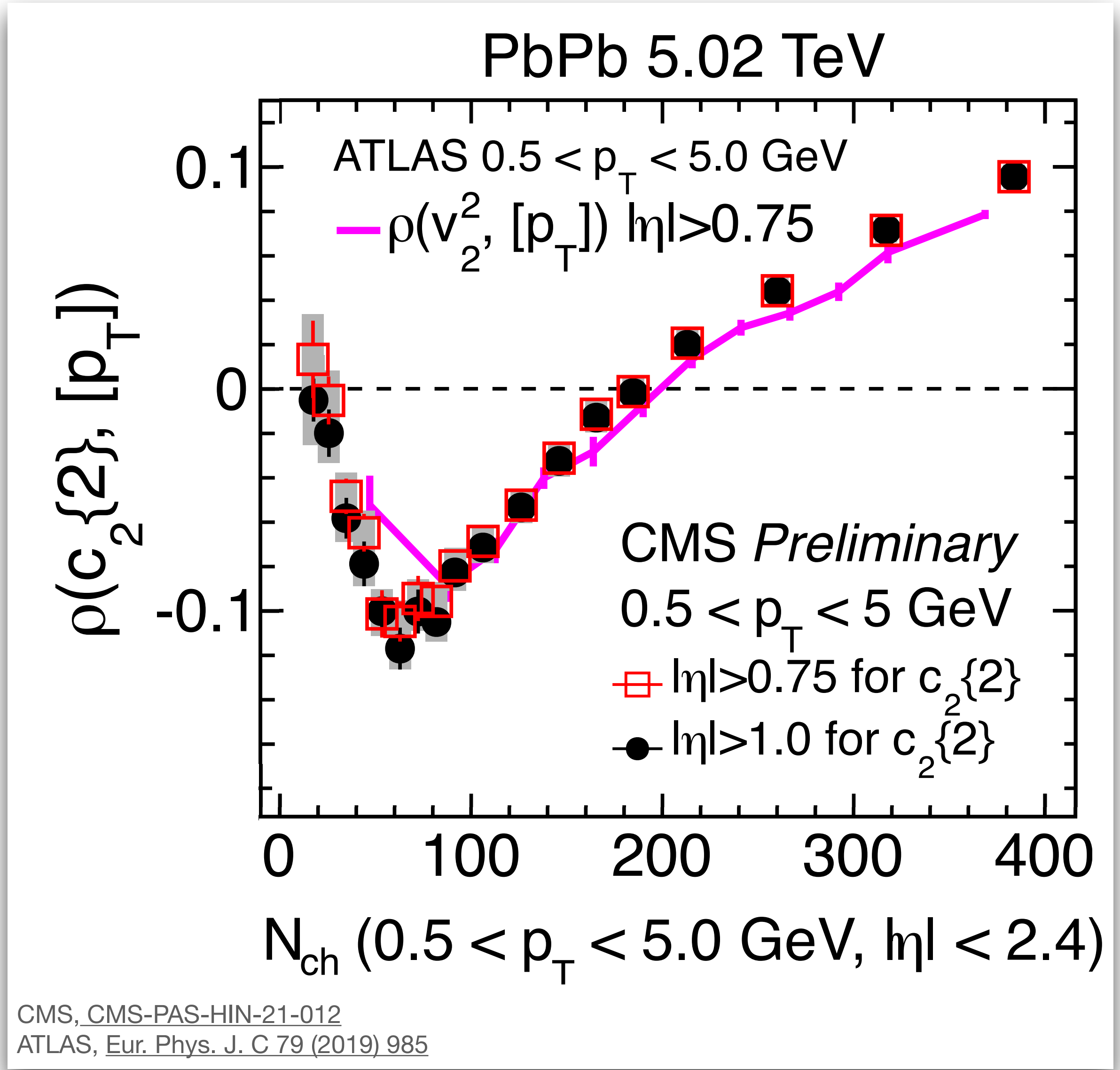
⇒ Change of slope sign → presence of CGC?



Correlation between $[p_T]$ and v_2 in Pb—Pb at low multiplicity

$\rho(v_2^2, [p_T])$ in Pb—Pb:

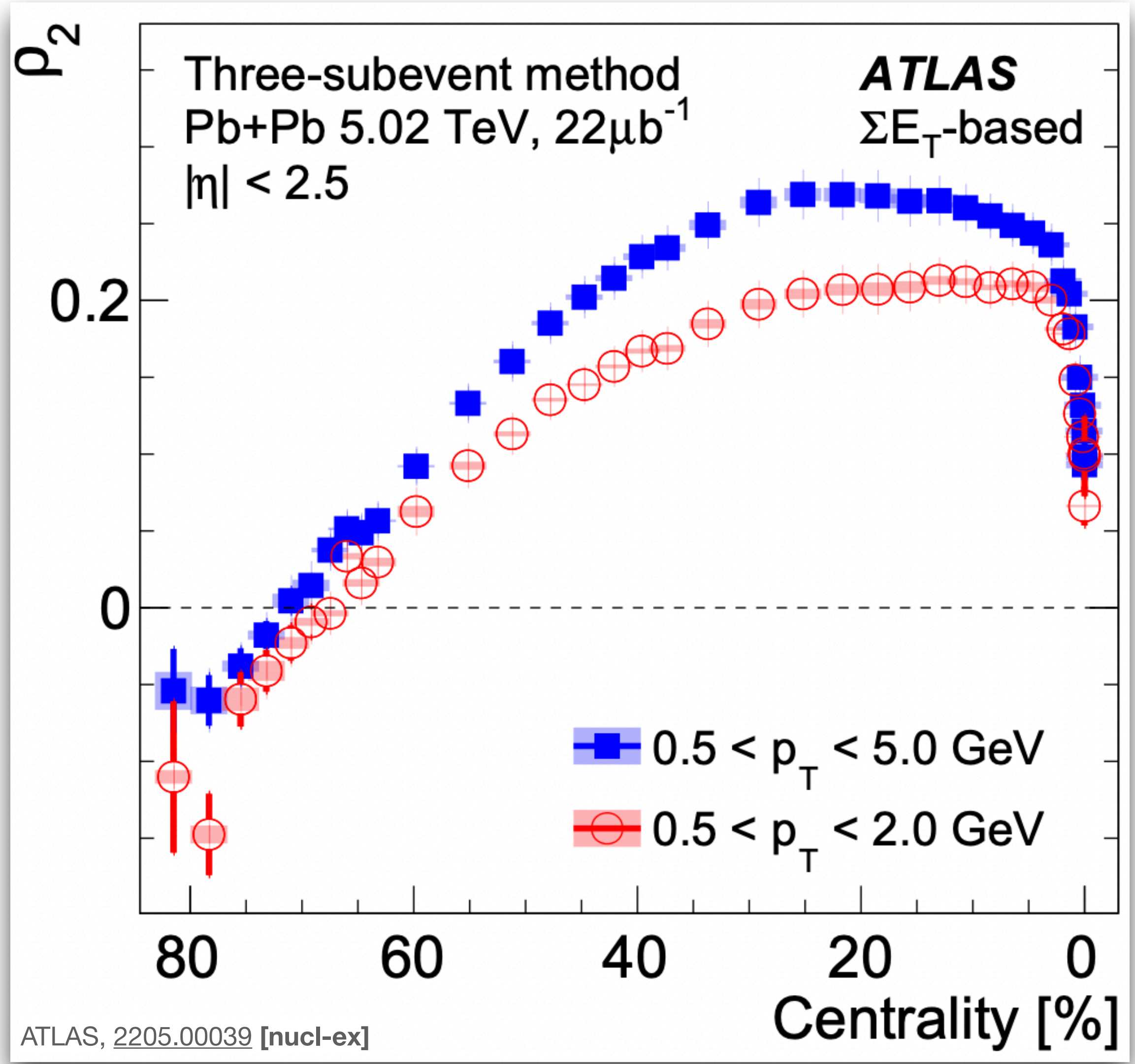
- Decreasing + increasing trend at low multiplicity



Correlation between $[p_T]$ and v_2 in Pb–Pb at low multiplicity

$\rho(v_2^2, [p_T])$ in Pb–Pb:

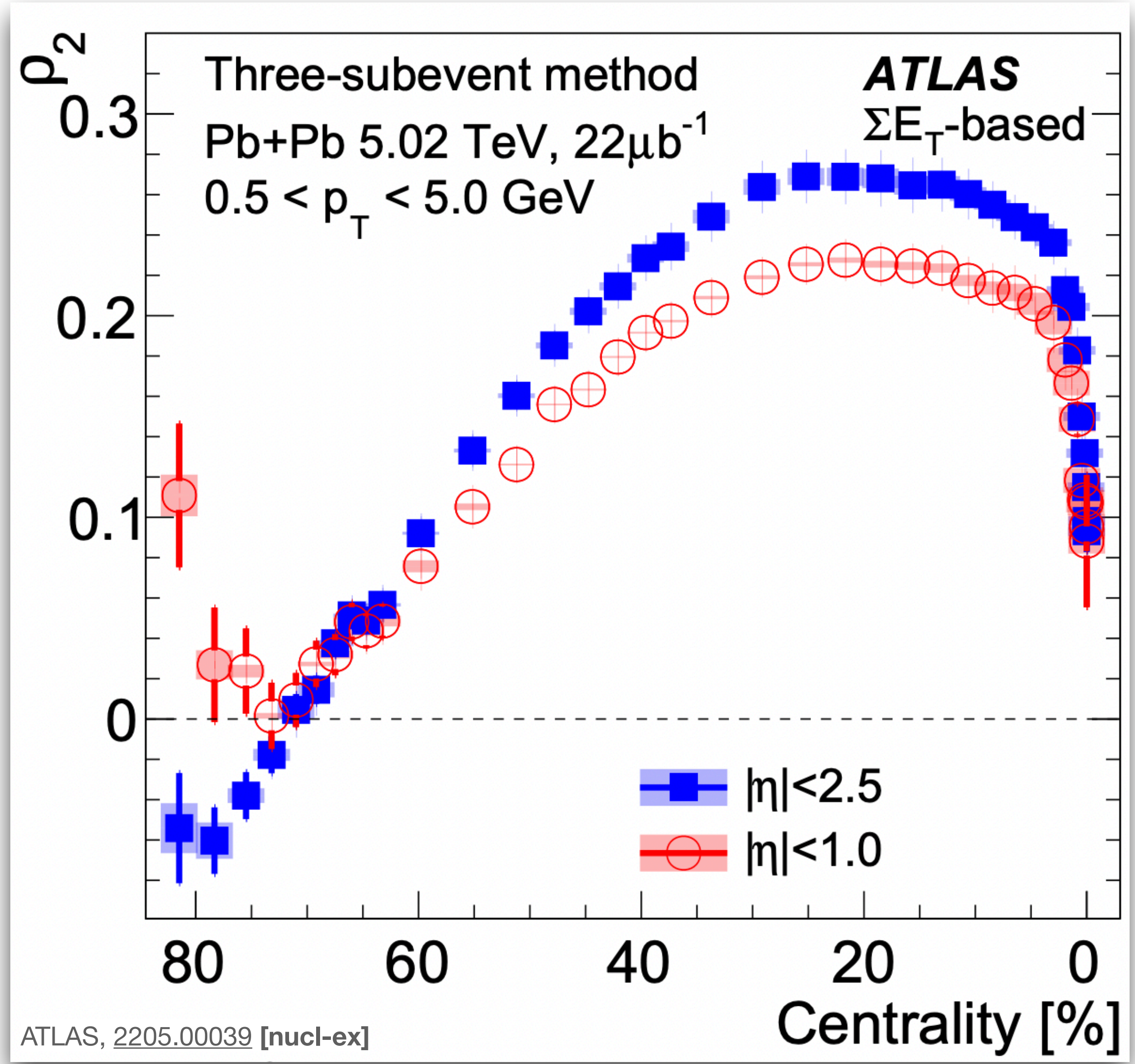
- Decreasing + increasing trend at low multiplicity
- Sensitive to p_T interval...



Correlation between $[p_T]$ and v_2 in Pb–Pb at low multiplicity

$\rho(v_2^2, [p_T])$ in Pb–Pb:

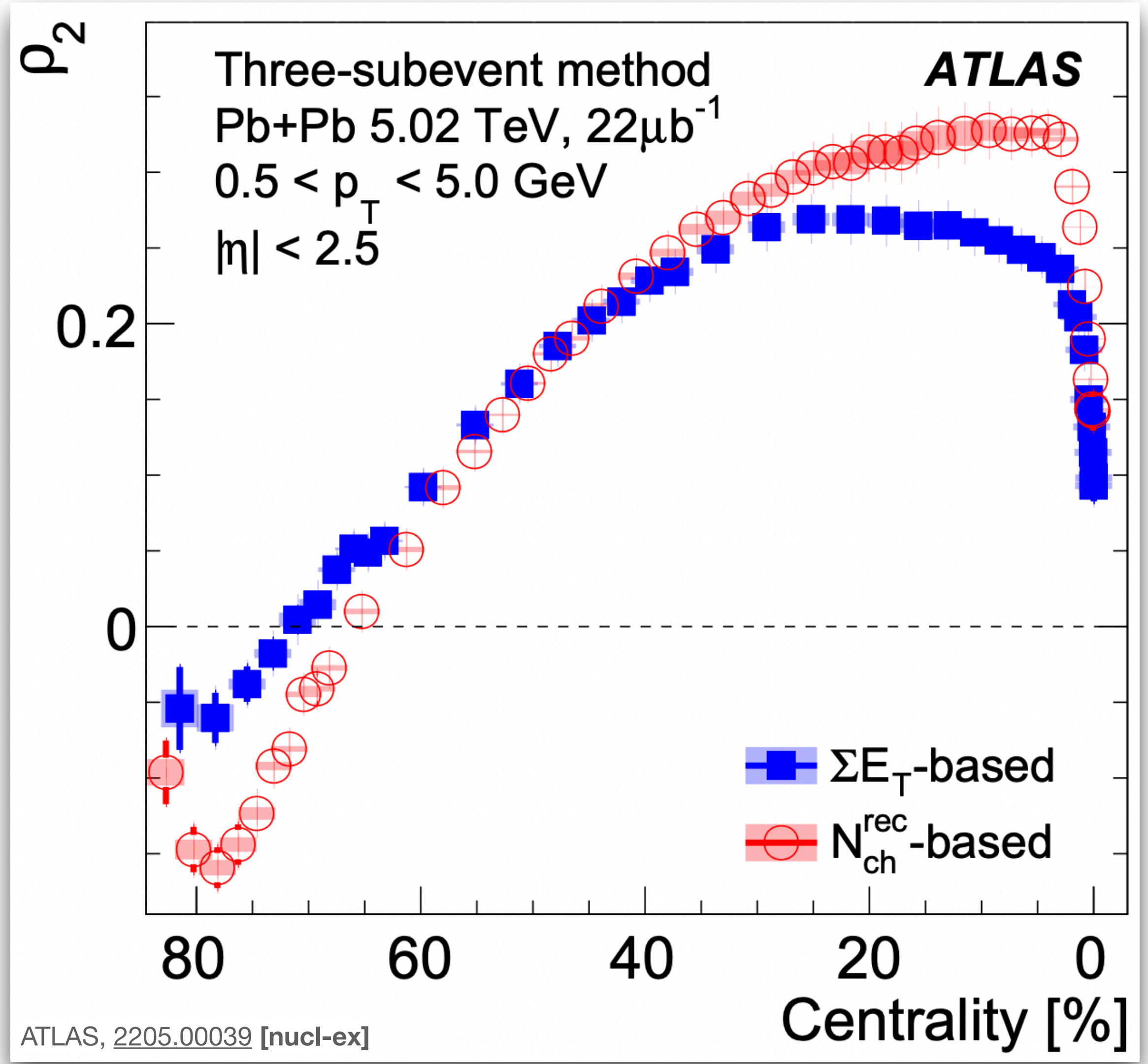
- Decreasing + increasing trend at low multiplicity
- Sensitive to p_T interval...
 - and pseudorapidity range...



Correlation between $[p_T]$ and v_2 in Pb–Pb at low multiplicity

$\rho(v_2^2, [p_T])$ in Pb–Pb:

- Decreasing + increasing trend at low multiplicity
- Sensitive to p_T interval...
 - and pseudorapidity range...
 - and even the multiplicity estimator



Correlation between $[p_T]$ and v_2 comparison to models

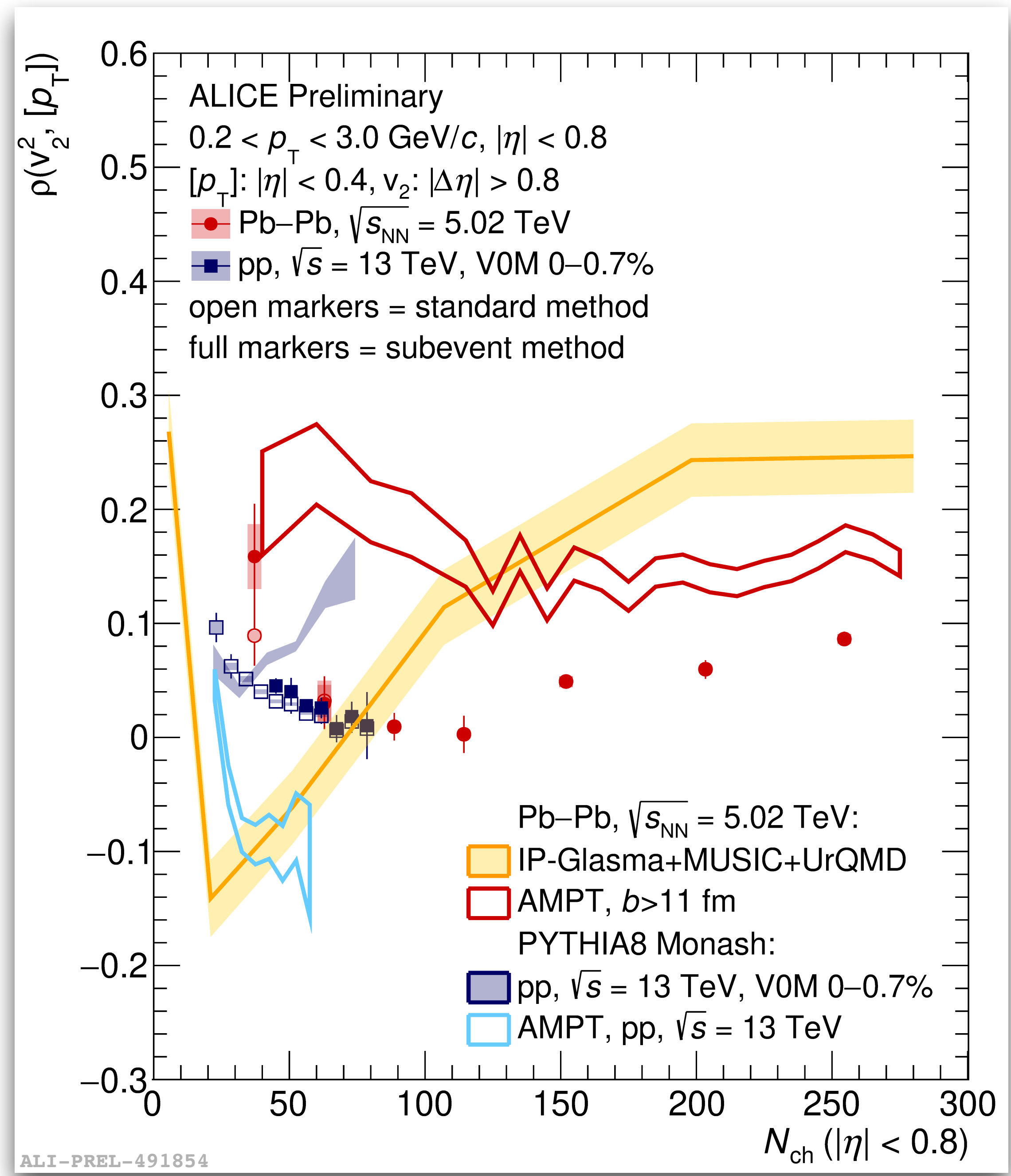
$\rho(v_2^2, [p_T])$ in Pb–Pb:

- IP-glasma+MUSIC+UrQMD:
 - Slope change around 20 charged tracks, significantly lower than in data
- AMPT:
 - Change of slope also observed, although at significantly higher N_{ch}

⇒ Slope change not exclusive to IP-Glasma

$\rho(v_2^2, [p_T])$ in pp:

- Consistent with Pb–Pb at similar N_{ch}
- Underestimated by AMPT, overestimated by PYTHIA

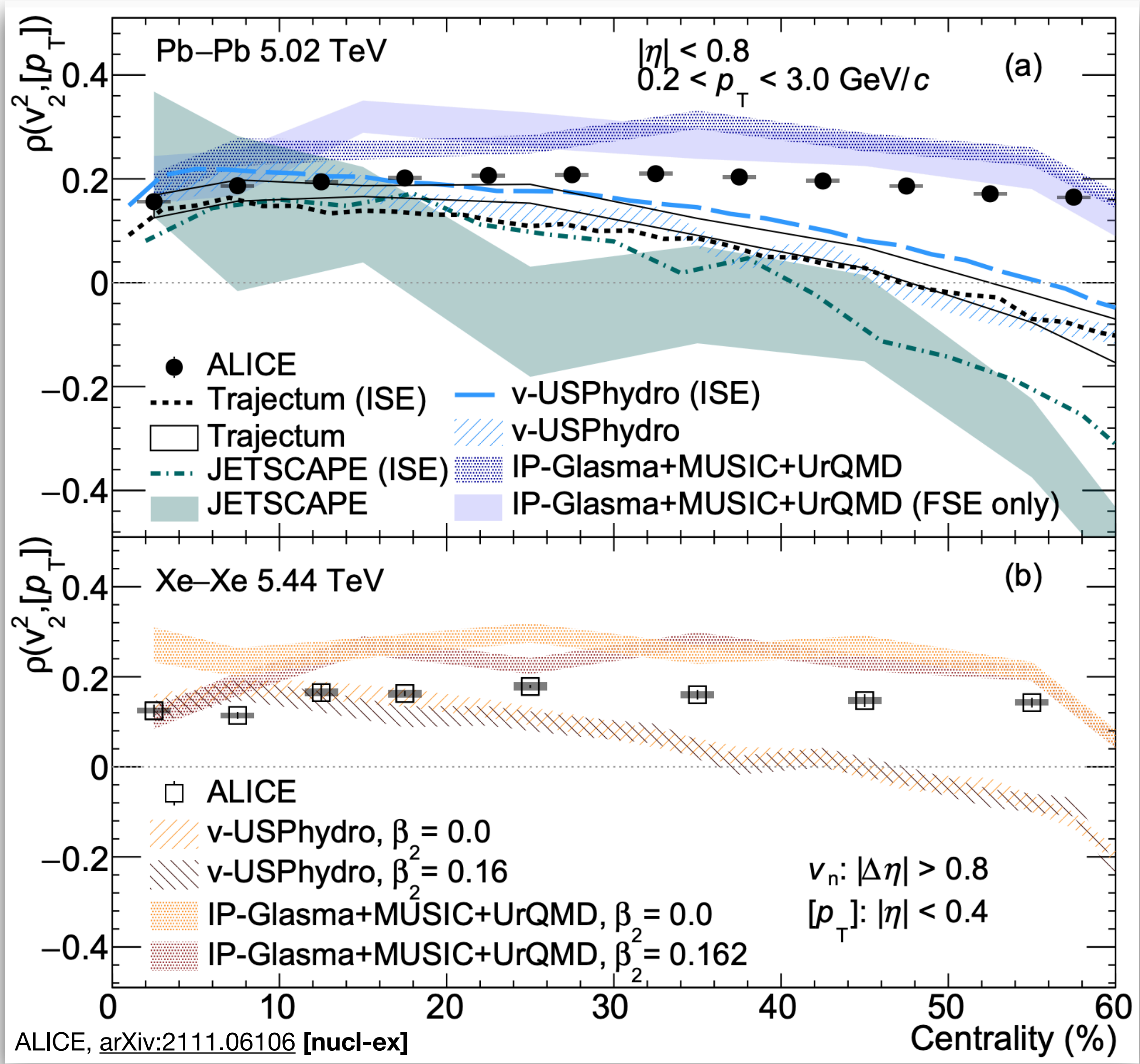


Correlation between $[p_T]$ and v_2 in Pb–Pb and Xe–Xe collisions

- ρ_2 slightly larger in Pb–Pb compared to Xe–Xe
- Comparison to models:
 - Below 20% centrality, all models provide a decent description

– More peripheral → best described by models with IP-Glasma

- Xe–Xe:
 - $\beta_2 = 0.162$ gives better description in most central collisions, similar to $\beta_2 = 0$ in more peripheral



Summary

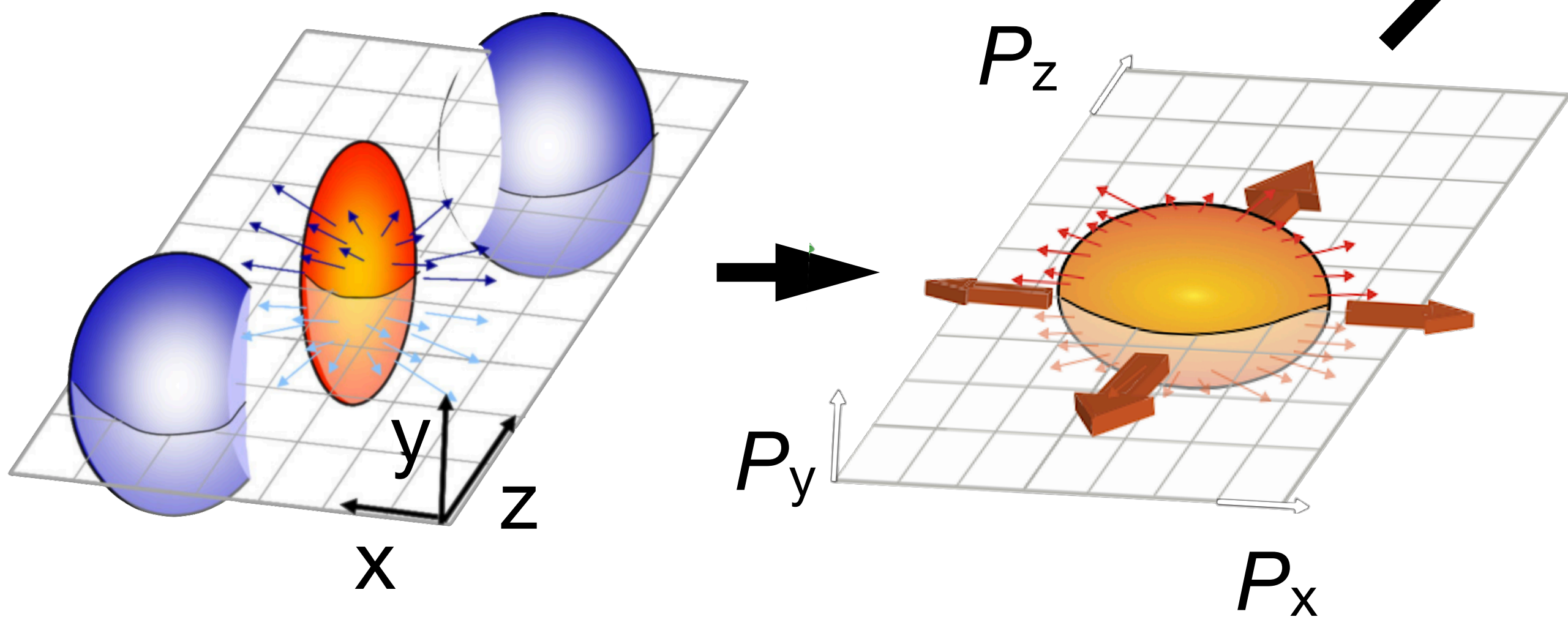
- p_T -differential v_2 of identified hadrons in pp and p–Pb collisions show remarkable similarities to Pb–Pb collisions
 - ⇒ Suggests dynamical evolution similar to that in Pb–Pb
- Relative flow fluctuations: emerging p_T dependence in peripheral collisions
- Higher moments of v_2 PDF: evolution with centrality and p_T suggests sensitivity to initial geometry and transport properties of QGP
- Correlations between $[p_T]$ and v_2 :
 - ⇒ Highly sensitive to kinematic cuts and multiplicity estimator
 - ⇒ Data better described by models with IP-Glasma in initial conditions
 - ⇒ Observed decreasing trend at small N_{ch} in Pb–Pb collisions
 - ⇒ Change of slope at low N_{ch} not unique to models with IP-Glasma

Backup

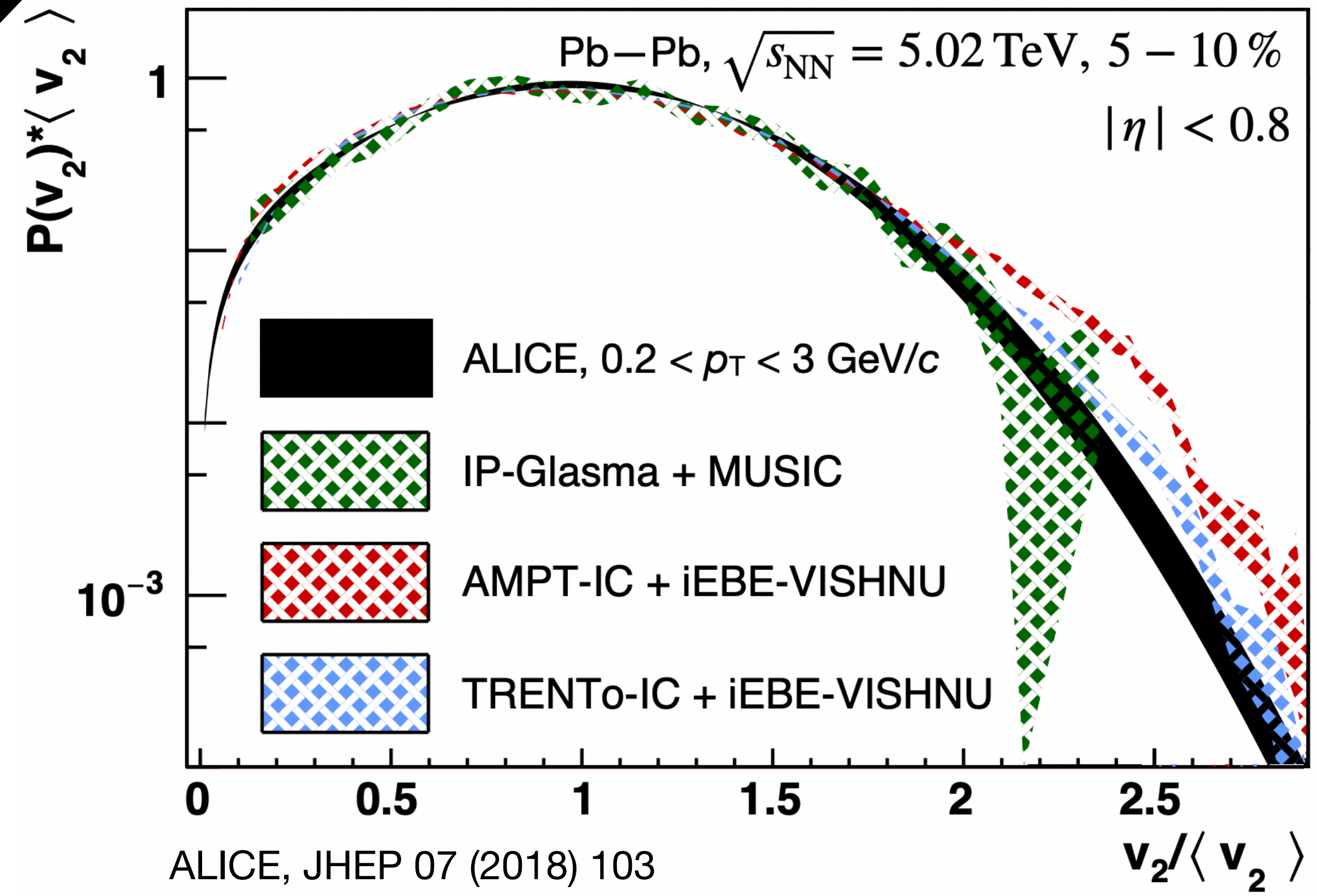
From collisions to measurement

- Overlap between colliding nuclei:
 ⇒ Initial state, geometry & its fluctuations
- Hydrodynamic expansion of QGP:
 ⇒ Radial and anisotropic flow, sensitive to initial state and properties of QGP

⇒ (Non-)Gaussian probability density function of v_n sensitive to initial state eccentricity ϵ_n
 ⇒ v_n fluctuations measured w.r.t. averaged Ψ_n



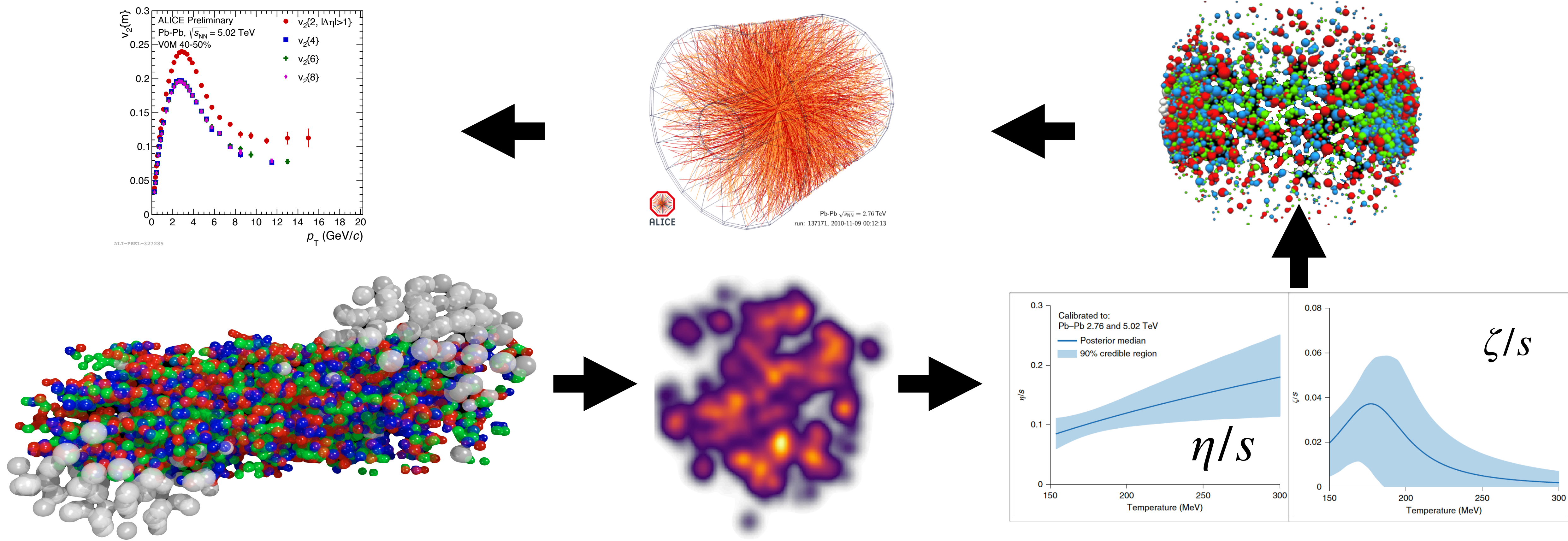
$$\frac{dN}{d\varphi} \propto 1 + 2 \sum_{n=1}^{\infty} v_n \cos n (\varphi - \Psi_n)$$



From collisions to measurements

- Overlap between colliding nuclei:
 ⇒ Initial state, geometry & its fluctuations

- Hydrodynamic expansion of QGP:
 ⇒ Radial and anisotropic flow, sensitive to initial state and properties of QGP



Correlation between $[p_T]$ and v_2 comparison to models

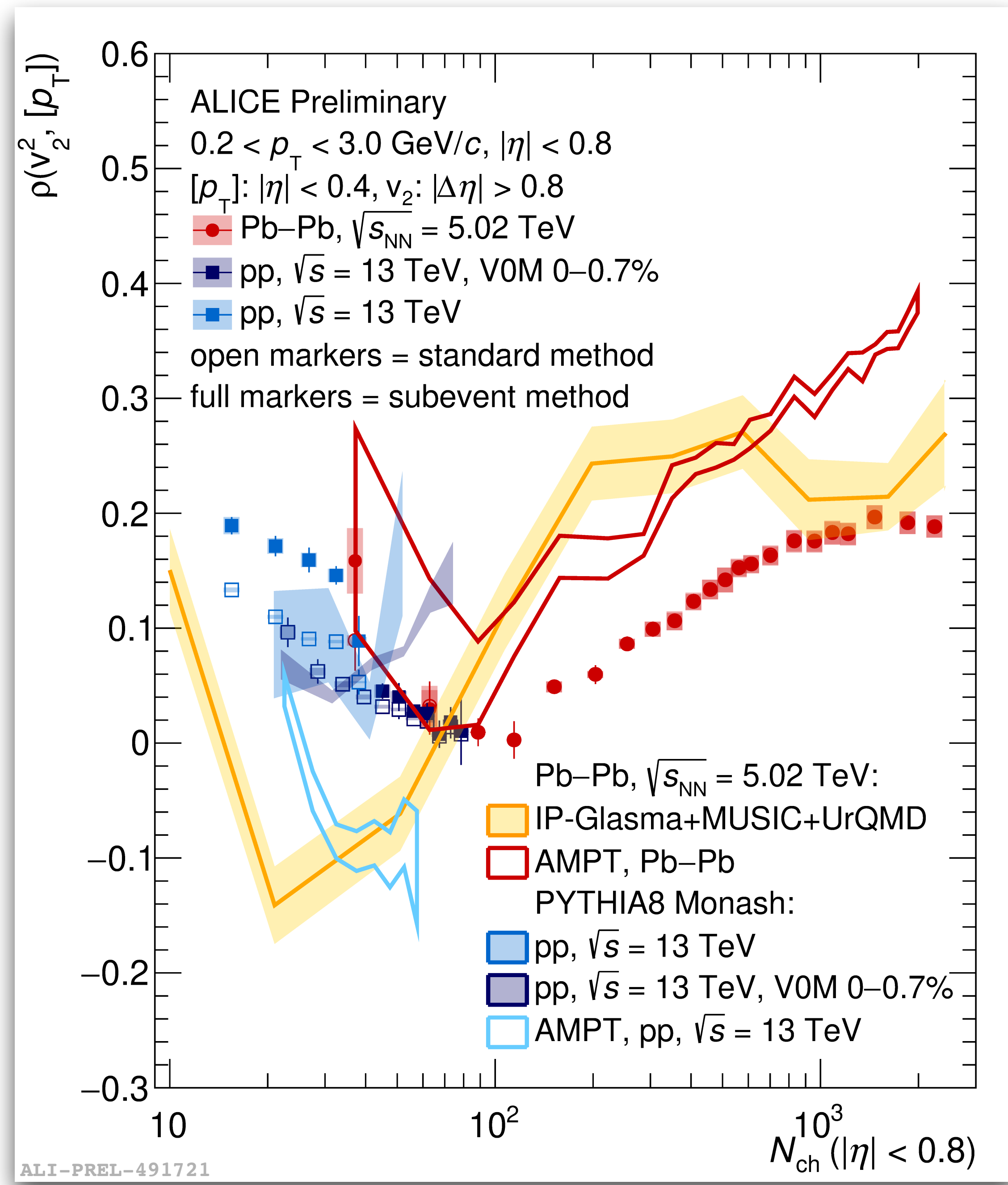
$\rho(v_2^2, [p_T])$ in Pb–Pb:

- IP-glasma+MUSIC+UrQMD:
 - Slope change around 20 charged tracks, significantly lower than in data
- AMPT:
 - Change of slope also observed, although at significantly higher N_{ch}

⇒ Slope change not exclusive to IP-Glasma

$\rho(v_2^2, [p_T])$ in pp:

- Consistent with Pb–Pb at similar N_{ch}
- Underestimated by AMPT, overestimated by PYTHIA



Flow vector fluctuations in Pb–Pb

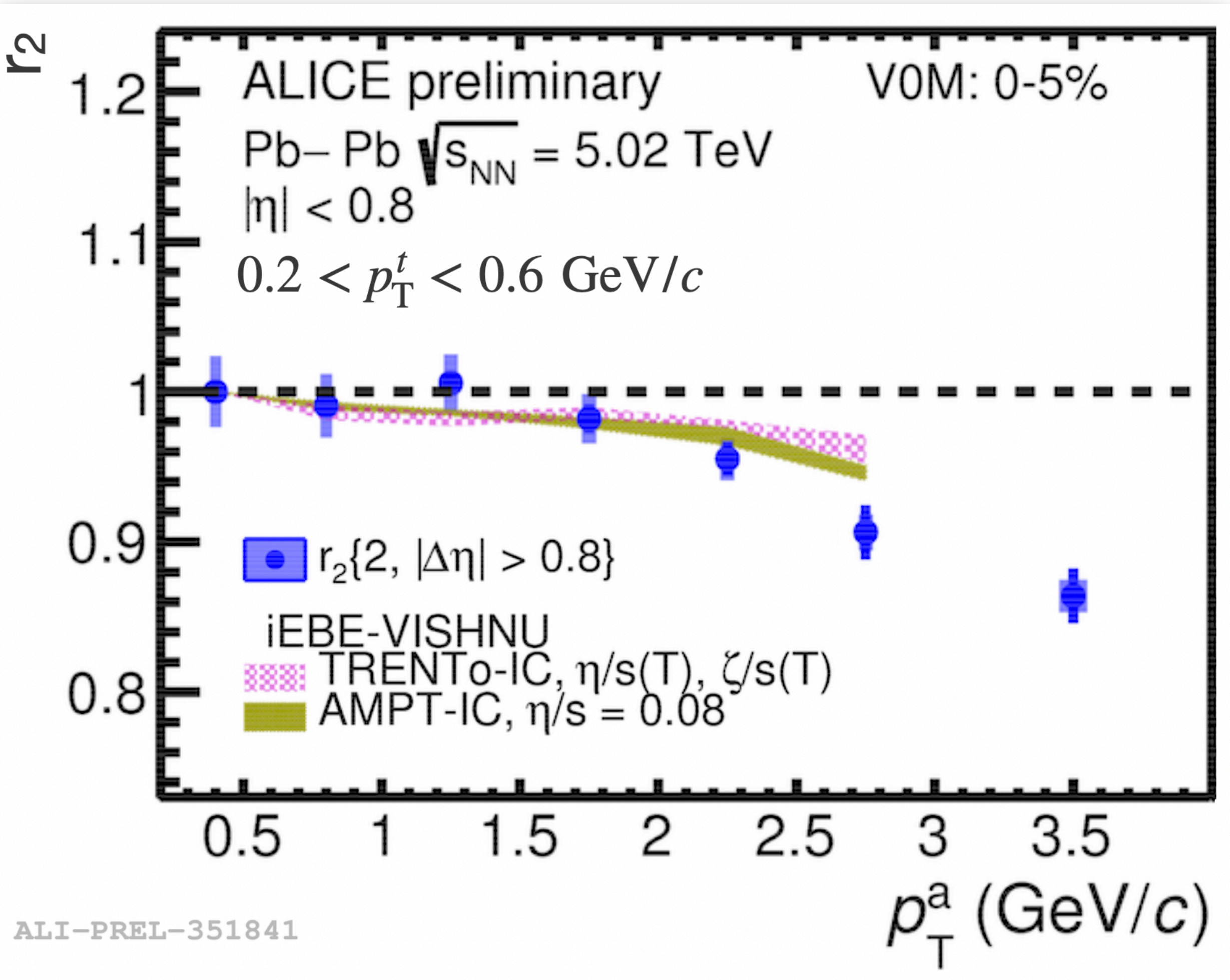
Flow factorisation ratio r_n =
$$\frac{\langle v_n^a v_n^t \cos[n (\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$$

(**t**rigger and **a**ssociated)

Flow vector fluctuations in Pb–Pb

Flow factorisation ratio $r_n = \frac{\langle v_n^a v_n^t \cos[n(\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$
 (trigger and associated)

Largest fluctuations in central Pb–Pb collisions at high $p_T \Rightarrow$ large event-by-event fluctuations in the initial state



Flow vector fluctuations in Pb–Pb

Flow factorisation ratio $r_n = \frac{\langle v_n^a v_n^t \cos[n(\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$
 (trigger and associated)

Largest fluctuations in central Pb–Pb collisions at high $p_T \Rightarrow$ large event-by-event fluctuations in the initial state

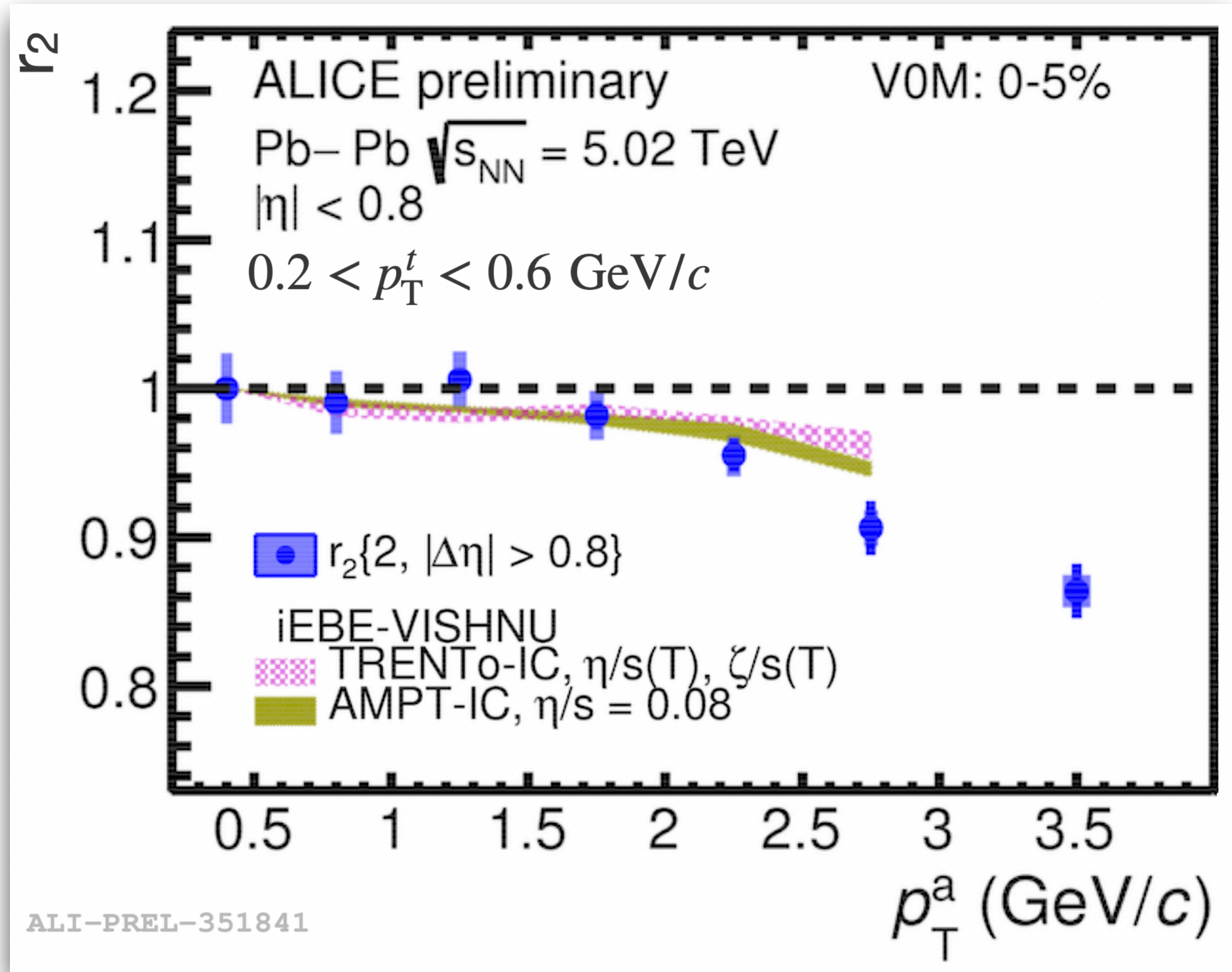
Deviations from $r_n = 1$ can be due to:

- *Flow angle fluctuations,*
- *Flow magnitude fluctuations,*

$$\langle \cos [n(\Psi_n^a - \Psi_n^t)] \rangle \neq 1$$

$$\langle v_n^a v_n^t \rangle \neq \sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}$$

- Cannot be measured directly, but upper/lower limits can be estimated



Flow vector fluctuations in Pb–Pb

Flow factorisation ratio $r_n = \frac{\langle v_n^a v_n^t \cos[n(\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$
 (trigger and associated)

Largest fluctuations in central Pb–Pb collisions at high $p_T \Rightarrow$ large event-by-event fluctuations in the initial state

Deviations from $r_n = 1$ can be due to:

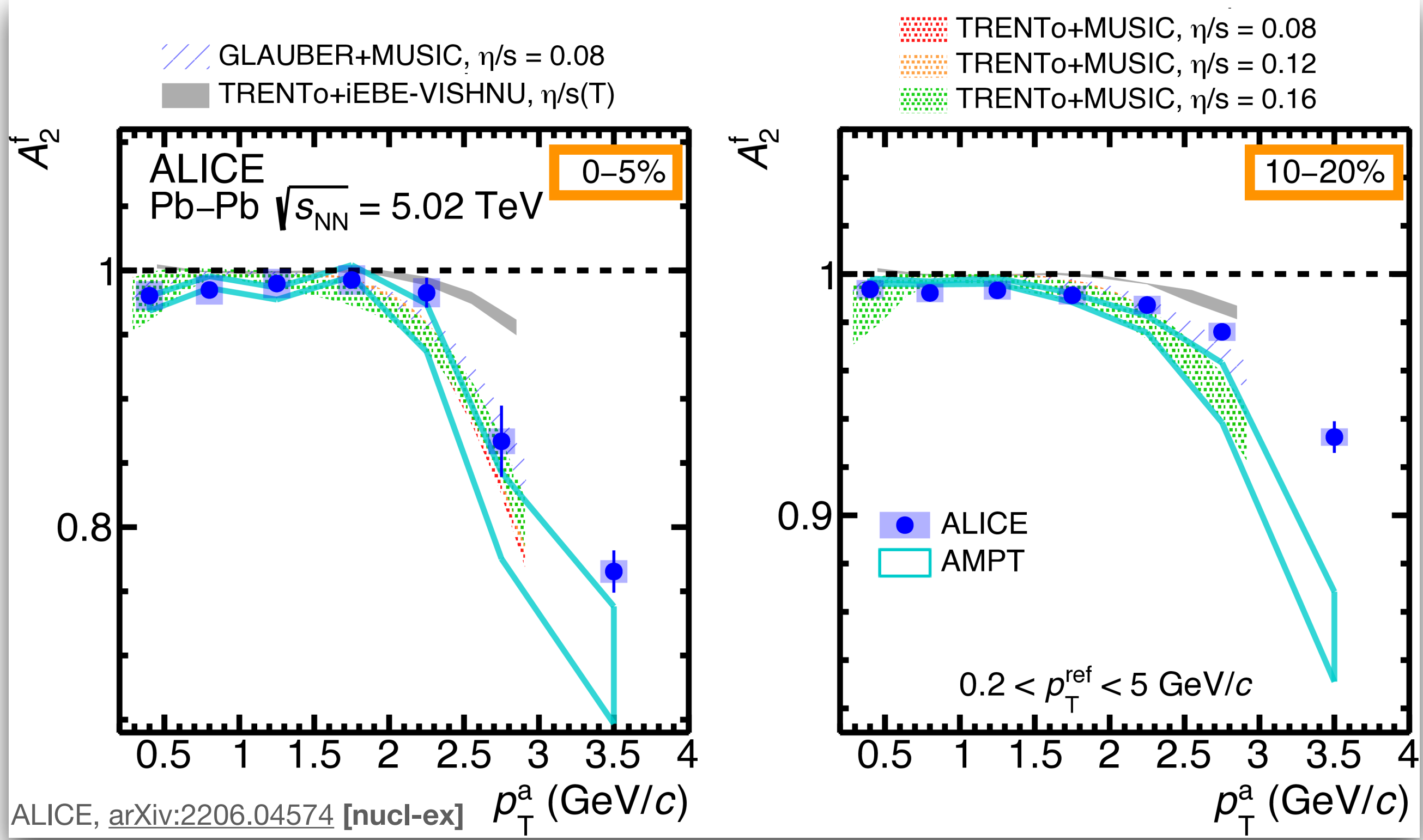
- *Flow angle fluctuations,*
 $\langle \cos [n(\Psi_n^a - \Psi_n^t)] \rangle \neq 1$

- *Flow magnitude fluctuations,*

$$\langle v_n^a v_n^t \rangle \neq \sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}$$

- Cannot be measured directly, but upper/lower limits can be estimated

Flow angle fluctuations



• Sensitive to fluctuations in initial state, little sensitivity to η/s

Flow vector fluctuations in Pb–Pb

Flow factorisation ratio $r_n = \frac{\langle v_n^a v_n^t \cos[n(\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$
 (trigger and associated)

Largest fluctuations in central Pb–Pb collisions at high $p_T \Rightarrow$ large event-by-event fluctuations in the initial state

Deviations from $r_n = 1$ can be due to:

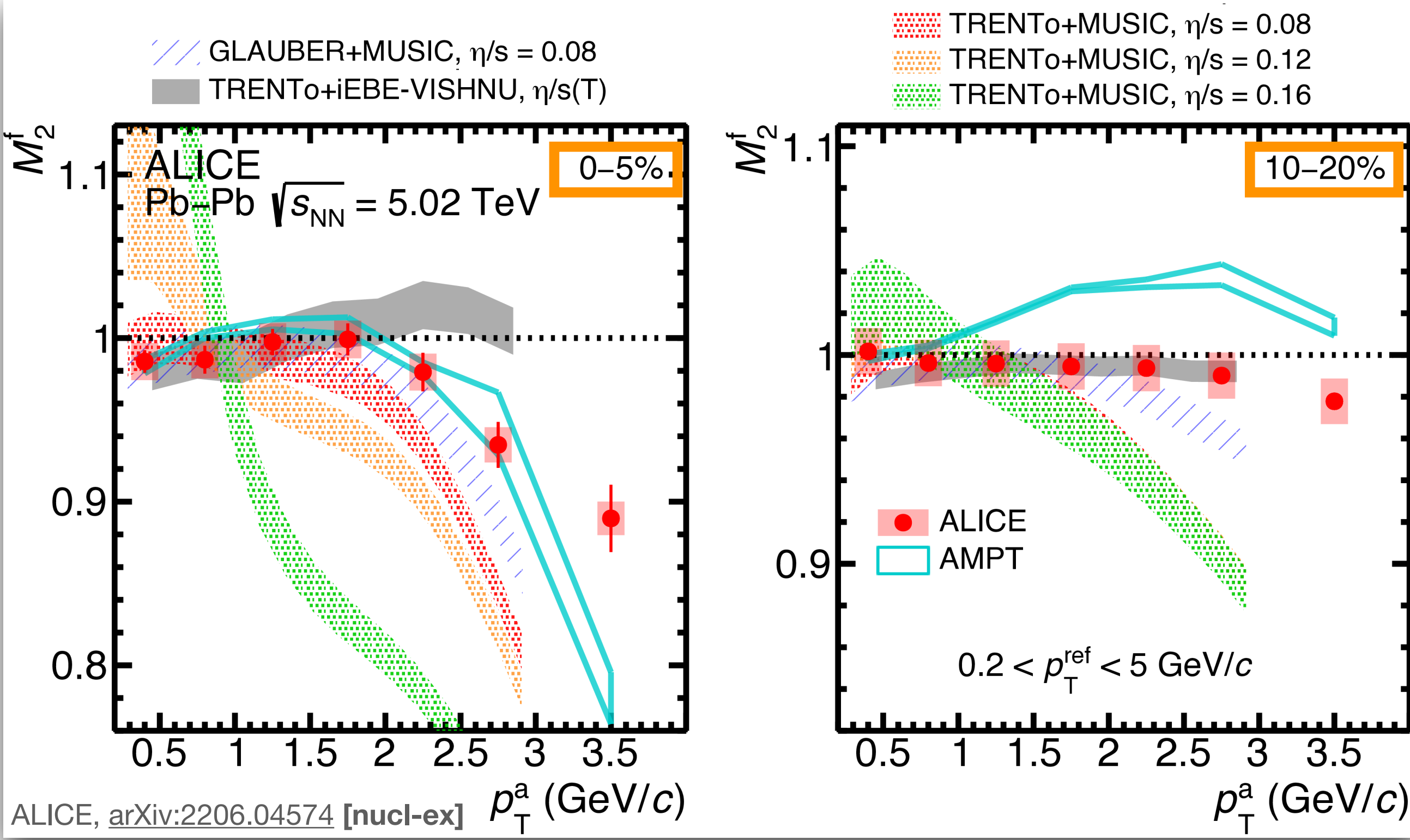
- *Flow angle fluctuations,*
 $\langle \cos [n(\Psi_n^a - \Psi_n^t)] \rangle \neq 1$

- *Flow magnitude fluctuations,*

$$\langle v_n^a v_n^t \rangle \neq \sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}$$

- Cannot be measured directly, but upper/lower limits can be estimated

Flow magnitude fluctuations



- Sensitive to fluctuations in initial state, little sensitivity to η/s
- Strong sensitivity to shear viscosity, but only in the most central collisions

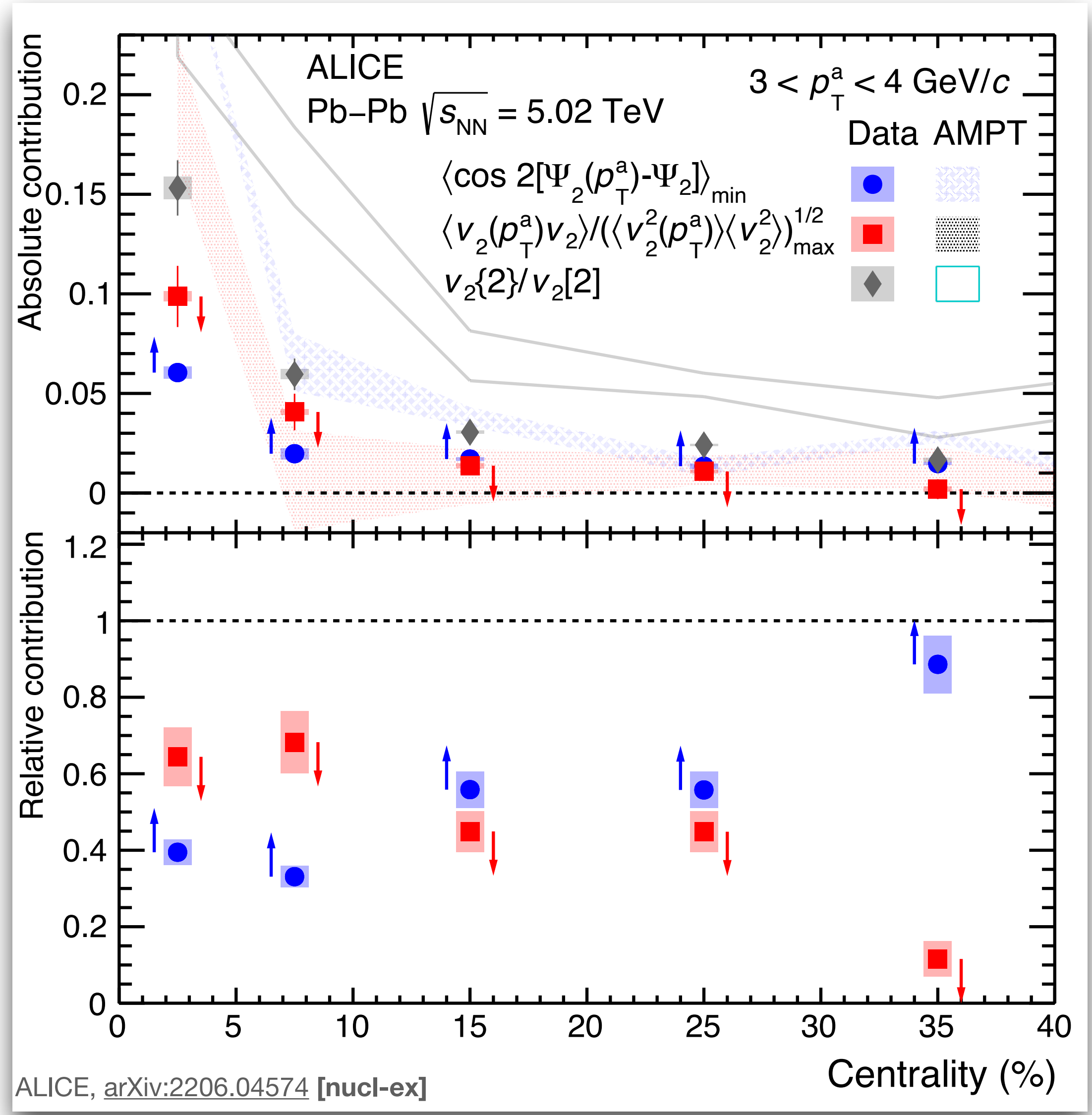
Flow vector fluctuation limits in Pb–Pb

Flow factorisation ratio $r_n = \frac{\langle v_n^a v_n^t \cos[n(\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$
 (trigger and associated)

Largest fluctuations in central Pb–Pb collisions at high $p_T \Rightarrow$ large event-by-event fluctuations in the initial state

- At least 40% of fluctuations in central collisions originate from **flow angle** fluctuations
- Above 30% centrality, **flow magnitude** fluctuations are suppressed

First measurement separating flow angle and magnitude fluctuations \Rightarrow Challenges the assumption of a common symmetry plane



Flow vector fluctuations in Pb–Pb

Define flow factorisation as
$$r_n = \frac{V_{n\Delta}(p_T^a, p_T^t)}{\sqrt{V_{n\Delta}(p_T^a, p_T^a) \cdot V_{n\Delta}(p_T^t, p_T^t)}} = \frac{\langle v_n^a v_n^t \cos[n(\Psi_n^a - \Psi_n^t)] \rangle}{\sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}}$$

Deviations from $r_n = 1$ can be due to:

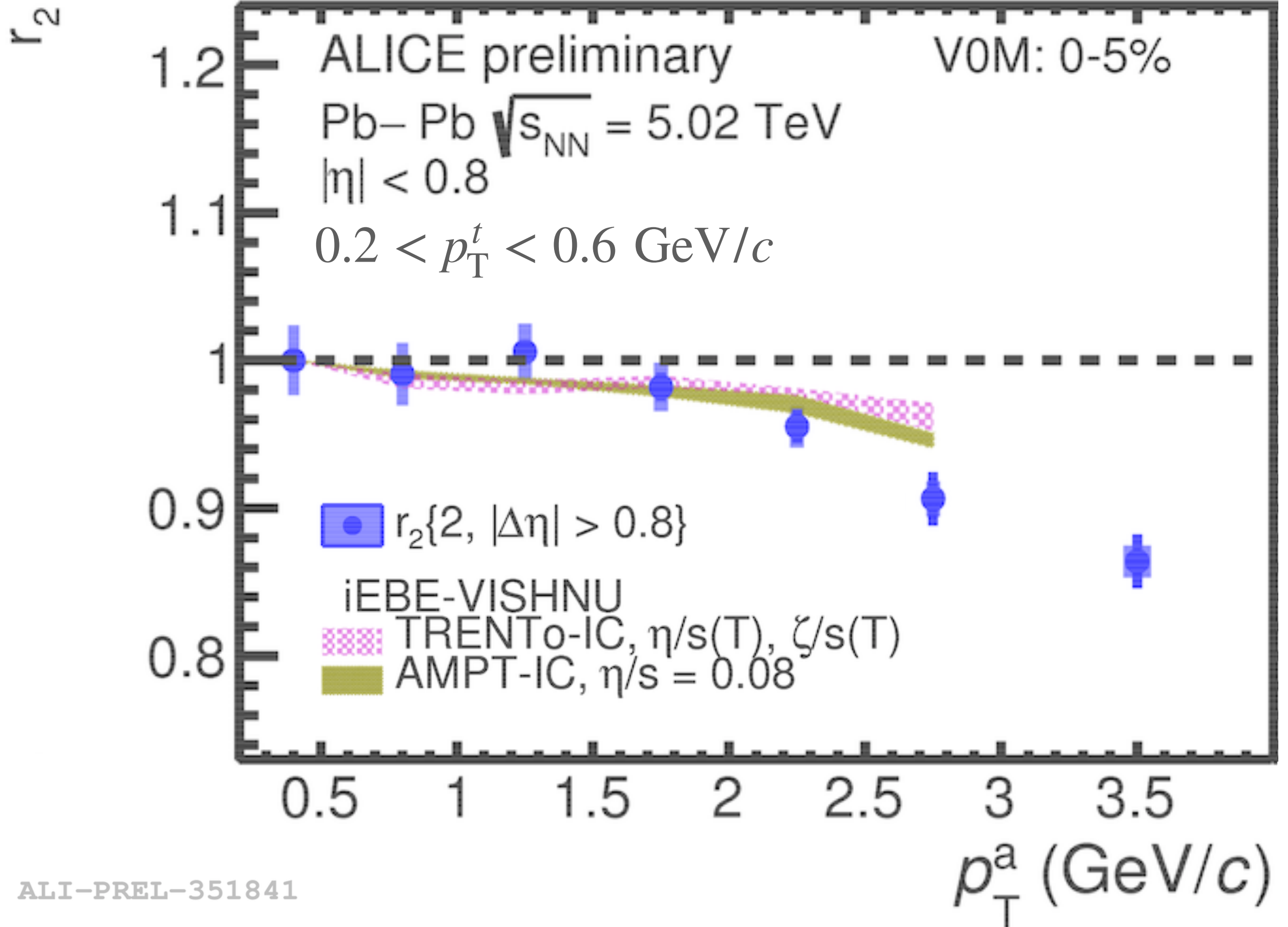
- *Flow magnitude fluctuations,*

$$\langle v_n^a v_n^t \rangle \neq \sqrt{\langle v_n^{a,2} \rangle \langle v_n^{t,2} \rangle}$$

- *Flow angle fluctuations,*

$$\langle \cos [n(\Psi_n^a - \Psi_n^t)] \rangle \neq 1$$

- Cannot measure directly, but *can* measure upper/lower limits!



ALI-PREL-351841

Correlation between $[p_T]$ and v_2

- **Shape of the fireball: anisotropic flow, $\varepsilon_n \rightarrow v_n$**
- **Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$**
- Initial state: geometry and fluctuations of **shape** and **size**
- Final state: correlation between v_n and $[p_T]$

Study with Pearson correlation coefficient:

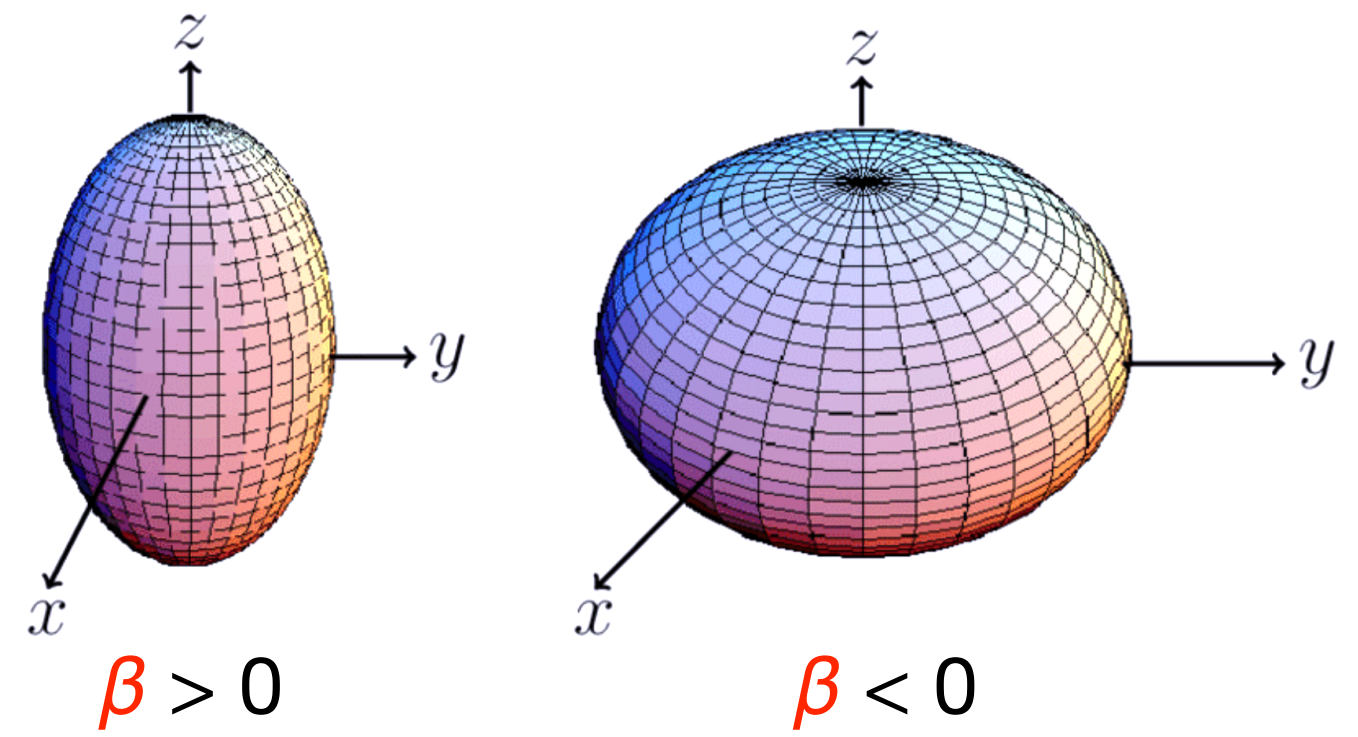
$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$

For deformed nuclei

Significantly smaller ρ_2 in central Xe—Xe, compared to Pb—Pb

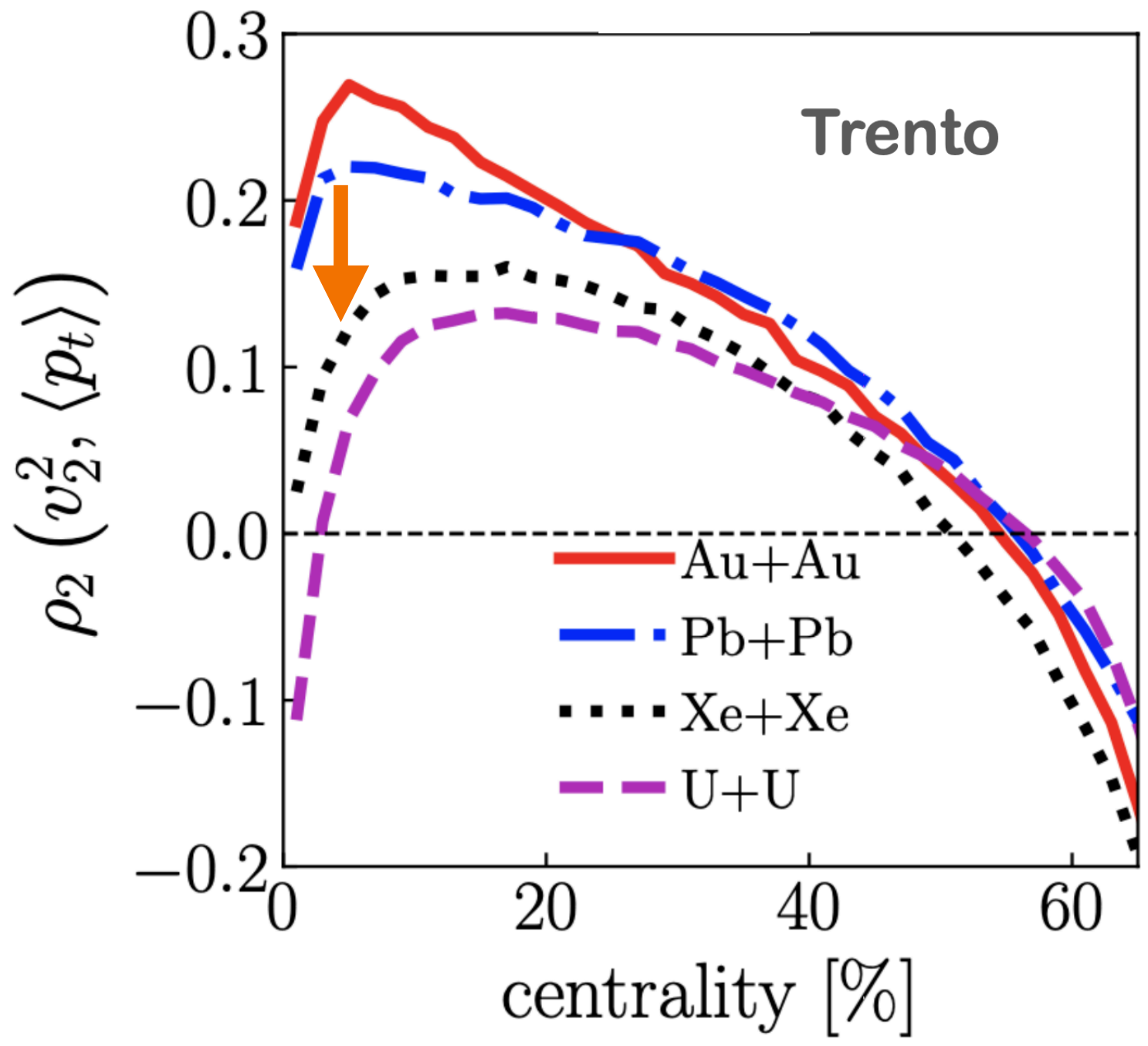
⇒ Deformation β reduces ρ_2

$$D_{WS}(r) = \frac{D_0}{1 + e^{\left(r - R_0(1 + \beta Y_{20}) \right) / a}}$$



Pb—Pb: $\beta \approx 0$
 Xe—Xe: $\beta \approx 0.16$

G. Giacalone, Phys. Rev. C 102, 024901 (2020)



Correlation between $[p_T]$ and v_2

- Shape of the fireball: anisotropic flow, $\epsilon_n \rightarrow v_n$
- Size of the fireball: radial flow, $[p_T], 1/R \rightarrow [p_T]$
- Initial state: geometry and fluctuations of shape and size
- Final state: correlation between v_n and $[p_T]$

For deformed nuclei

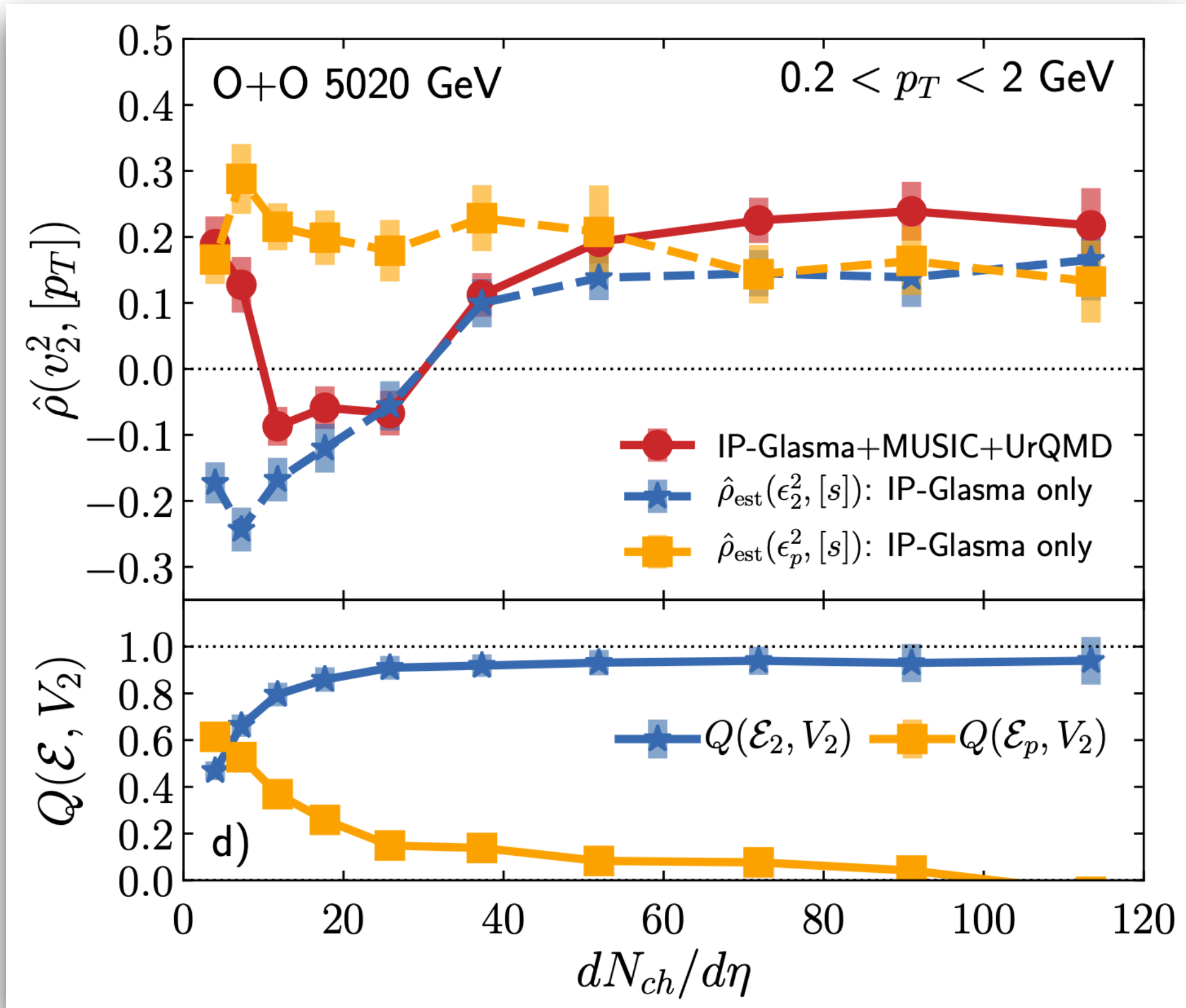
Significantly smaller ρ_2 in central Xe—Xe, compared to Pb—Pb
 \Rightarrow Deformation β reduces ρ_2

Probing the initial state

- Low multiplicity: geometry \rightarrow initial momentum correlations
 \Rightarrow Change of slope sign \rightarrow presence of CGC?

Study with Pearson correlation coefficient:

$$\rho_n \left(v_n^2, [p_T] \right) = \frac{\text{cov} \left(v_n^2, [p_T] \right)}{\sqrt{\text{var} \left(v_n^2 \right)} \sqrt{\text{var} \left([p_T] \right)}}$$



G. Giacalone, Phys. Rev. Lett. 125, 192301 (2020)