Jet Contribution to the γ-ray Flux in NGC 1068

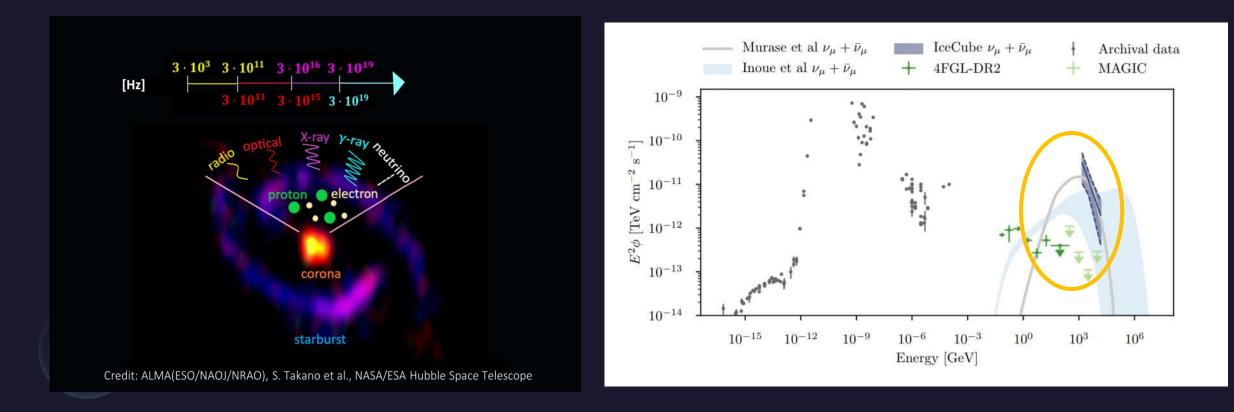
Silvia Salvatore

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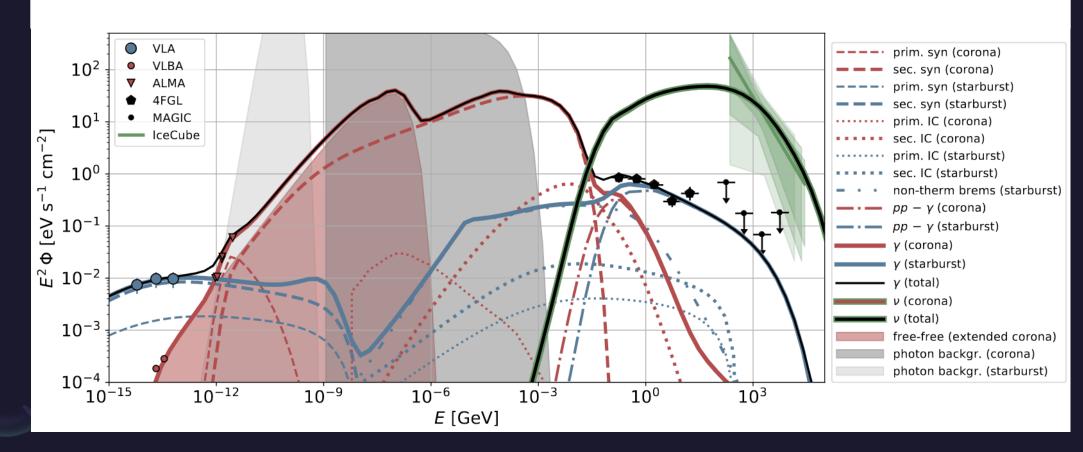


Why are neutrinos important?

- ALMA observations \rightarrow inner core and outer ring
- For energies between 100 GeV and 10 TeV \rightarrow significant difference in gamma-ray and neutrino fluxes

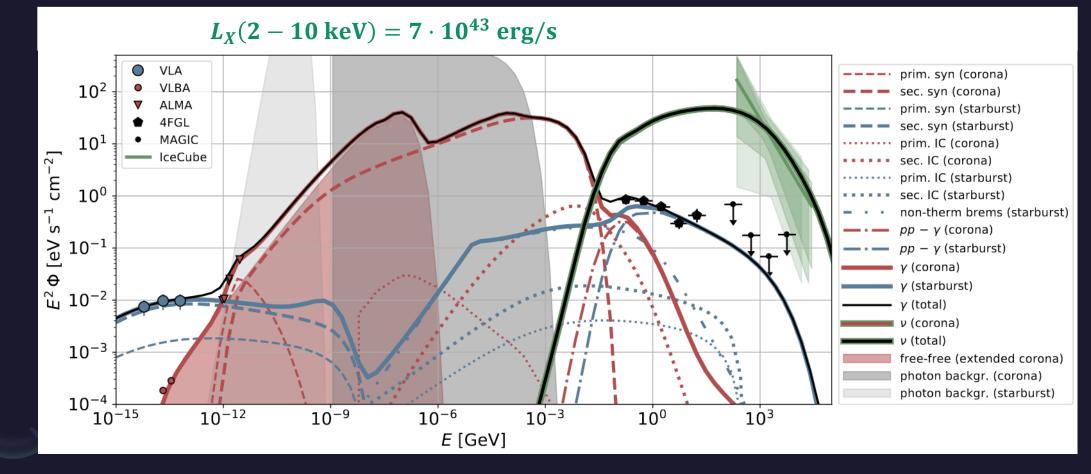


AGN-corona + starburst ring model



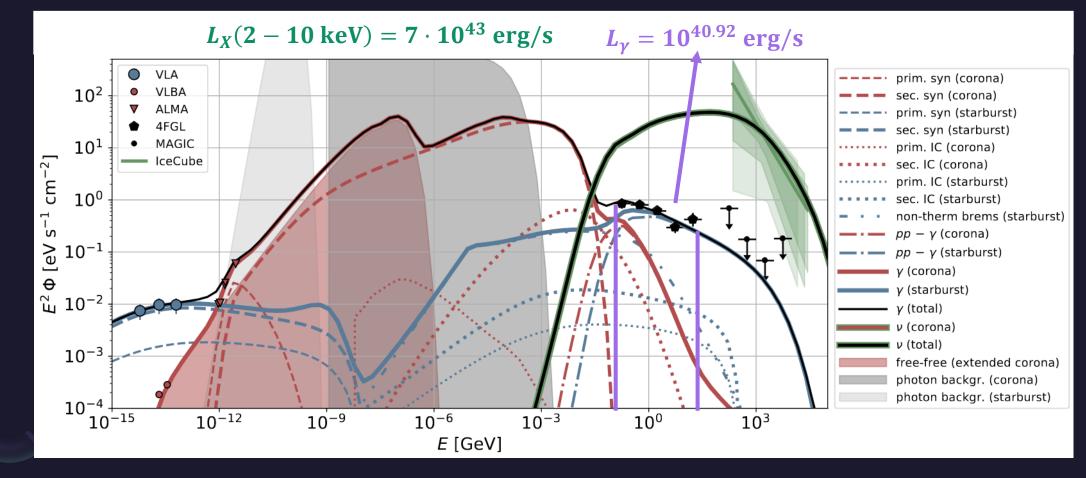
Eichmann,, **Salvatore**, ..., 2022

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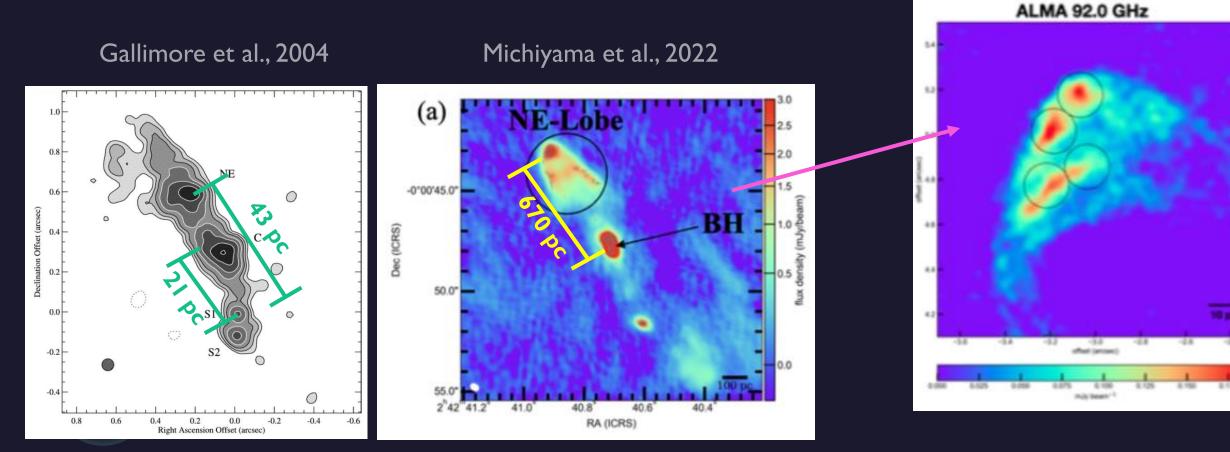
Introducing the jet

Radio data

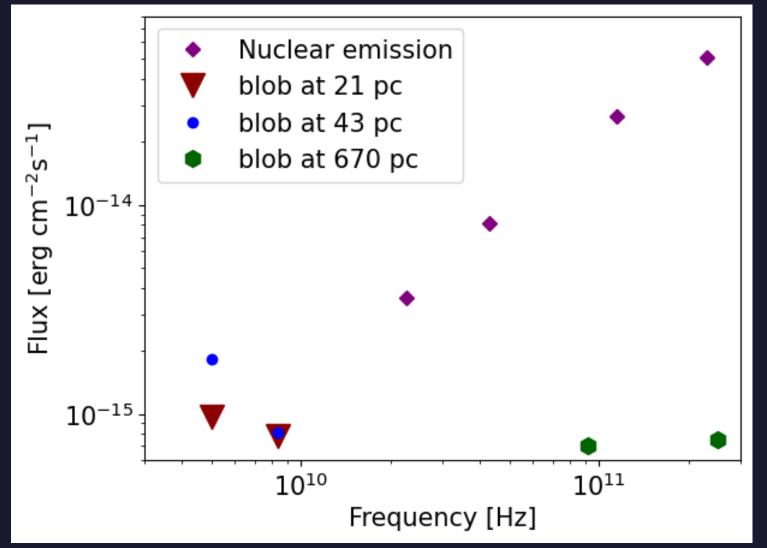
Michiyama et al., 2022 Gallimore et al., 2004 (a) 1.0 obe -0*00'45.0* Offset (arcsec) BH 0.4 Dec (ICRS 0.2 Declin 50.0" 0.0 -0.2 41.0 40.4 40.6 -0.4 0.8 0.4 0.2 0.0 -0.2 -0.6 RA (ICRS) Right Ascension Offset (arcsec)

Introducing the jet

Radio data



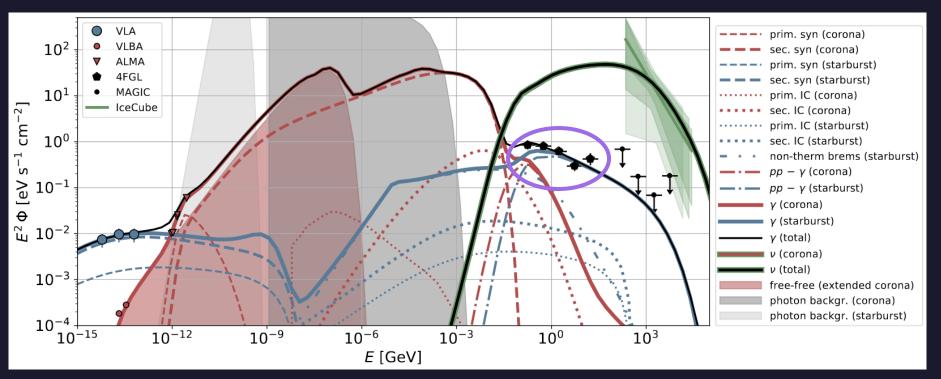
Radio data



How to produce high energy photons from these blobs?

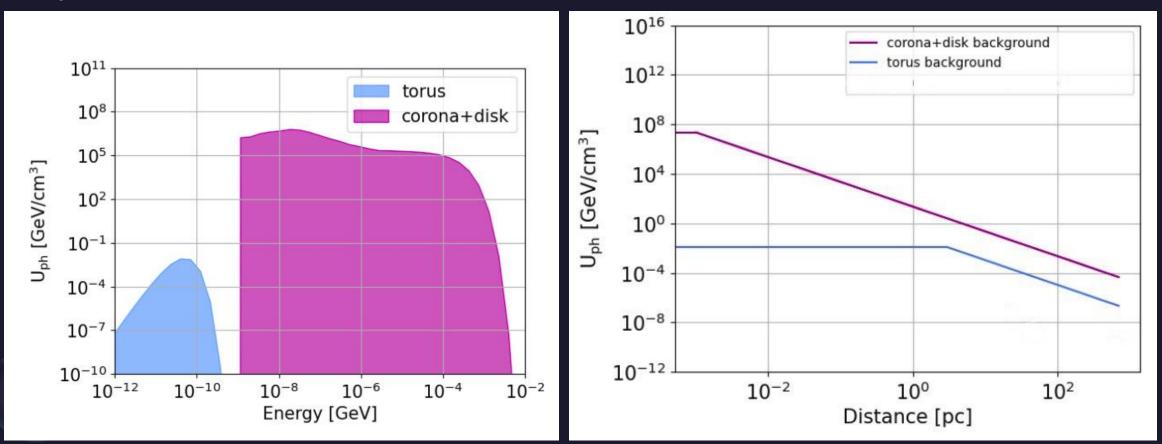
Possible γ -ray production scenarios:

- Inverse Compton
- $p\gamma$ interactions



Eichmann,, **Salvatore**, ..., 2022

• X-ray field from the **AGN-corona+disk** or IR field from the **torus**?



• X-ray photon field from the AGN-corona:

$$U_{COR} = \overline{U_{COR}} \left(\frac{d}{r_{COR}}\right)^{-2}$$

• Synchrotron photon field:

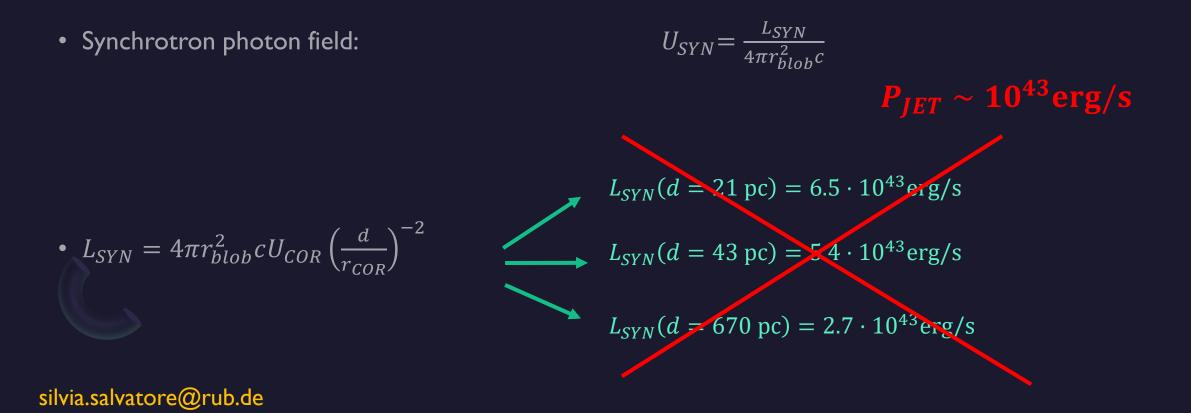
$$U_{SYN} = \frac{L_{SYN}}{4\pi r_{blob}^2 c}$$

•
$$L_{SYN} = 4\pi r_{blob}^2 c U_{COR} \left(\frac{d}{r_{COR}}\right)^{-2}$$

 $L_{SYN}(d = 21 \text{ pc}) = 6.5 \cdot 10^{43} \text{ erg/s}$
 $L_{SYN}(d = 43 \text{ pc}) = 5.4 \cdot 10^{43} \text{ erg/s}$
 $L_{SYN}(d = 670 \text{ pc}) = 2.7 \cdot 10^{43} \text{ erg/s}$

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$$L_{SYN}(d = 21 \,\mathrm{pc}) = 6.5 \cdot 10^{43} \,\mathrm{erg/s}$$

 $L_{SYN}(d = 43 \text{ pc}) = 54 \cdot 10^{43} \text{ erg/s}$

$$L_{SYN}(d = 670 \text{ pc}) = 2.7 \cdot 10^{43} \text{erg/s}$$

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• In the Thomson regime, assuming no absorption process is in action, this relation holds:

$$L_{IC} = L_{SYN} \frac{U_{RAD}}{U_B}$$

• From (AG Pacholczyk, JA Roberts, 1971):

minimum total energy stored in the source \rightarrow equipartition between particles and magnetic field energies:

$$E_B = \frac{3}{4}(1+k)E_e \longrightarrow B_{min} = (4.5)^{\frac{2}{7}}(1+k)^{\frac{2}{7}}c_{12}^{\frac{2}{7}}r_{blob}^{-\frac{6}{7}}L_{SYN}^{\frac{2}{7}}$$

$$L_{IC} = L_{SYN} \overline{U_{COR}} \left(\frac{d}{r_{COR}}\right)^{-2} \left(\frac{8\pi}{B_{min}^2}\right) \sim L_{\gamma} \sim 10^{40.92} \text{ erg/s}$$

3.9 \cdot 10⁷ GeV/cm³ 0.15 mpc

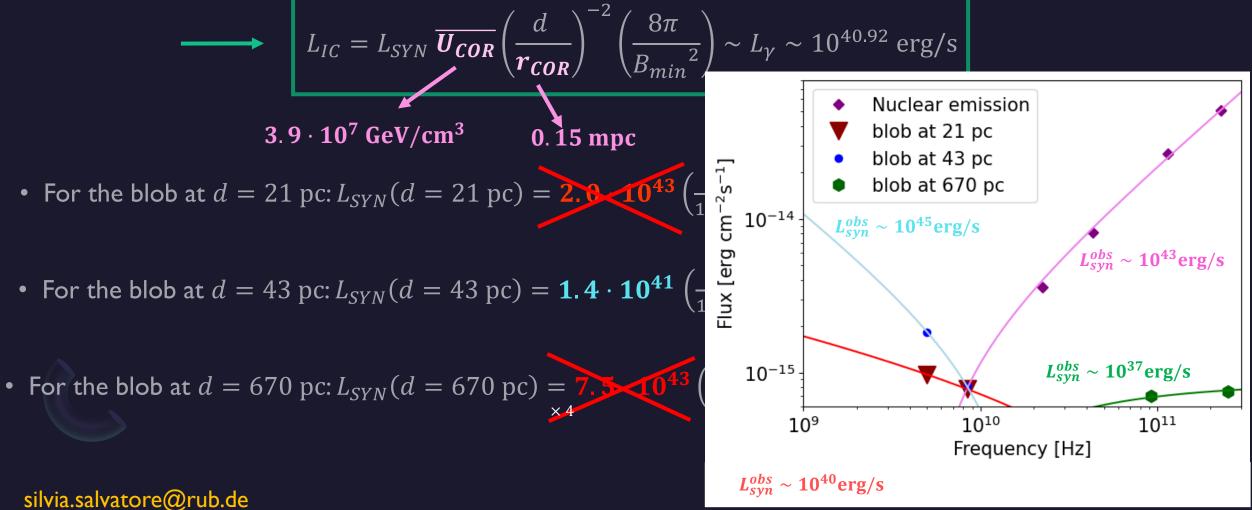
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$$3.9 \cdot 10^7 \text{ GeV/cm}^3 \qquad 0.15 \text{ mpc}$$

• For the blob at
$$d = 21 \text{ pc: } L_{SYN}(d = 21 \text{ pc}) = 2.0 \cdot 10^{43} \left(\frac{c_{12}}{1.3 \cdot 10^7}\right)^{\overline{3}} \left(\frac{r_{blob}}{2.5 \text{ pc}}\right)^{-4} \left(\frac{1+k}{101}\right)^{\overline{3}} \text{ erg/s}$$

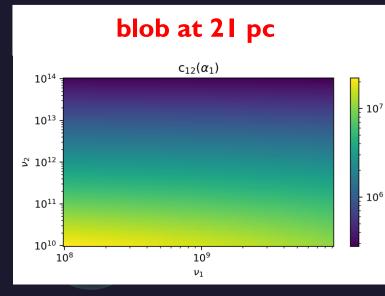
• For the blob at
$$d = 43 \text{ pc: } L_{SYN}(d = 43 \text{ pc}) = 1.4 \cdot 10^{41} \left(\frac{c_{12}}{1.3 \cdot 10^7}\right)^{\frac{4}{3}} \left(\frac{r_{blob}}{4.6 \text{ pc}}\right)^{-4} \left(\frac{1+k}{101}\right)^{\frac{4}{3}} \text{ erg/s}$$

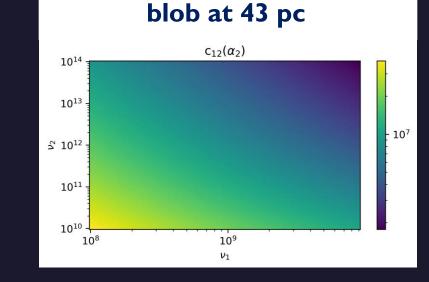
• For the blob at $d = 670 \text{ pc: } L_{SYN}(d = 670 \text{ pc}) = 7.5 \cdot 10^{43} \left(\frac{c_{12}}{3.9 \cdot 10^6}\right)^{\frac{4}{3}} \left(\frac{r_{blob}}{5 \text{ pc}}\right)^{-4} \left(\frac{1+k}{101}\right)^{\frac{4}{3}} \text{ erg/s}$



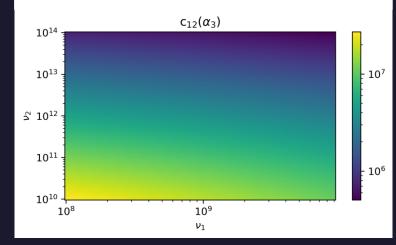
How robust is the dependence on c_{12} ?

$$c_{12=} c_2^{-1} c_1^{\frac{1}{2}} \frac{2\alpha - 2}{2\alpha - 1} \cdot \frac{\nu_1^{(1-2\alpha)/2} - \nu_2^{(1-2\alpha)/2}}{\nu_1^{1-\alpha} - \nu_2^{1-\alpha}} \left(c_{1=} \frac{3e}{4\pi m^3 c^5} \right) c_2 = \frac{2e^4}{3m^4 c^7}$$



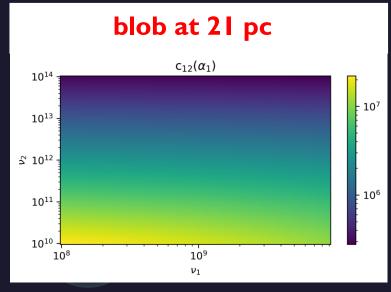


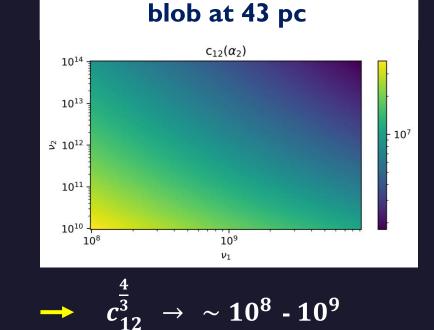


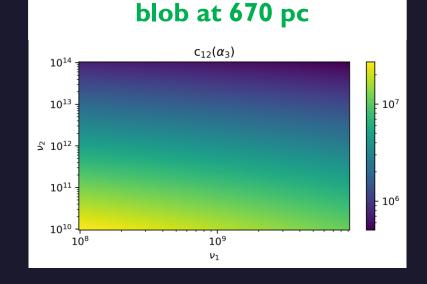


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How robust is the dependence on *k* ?

$$k \simeq \left(\frac{m_p}{m_e}\right)^{\frac{q-1}{2}}$$

• q: spectral index of injected CRs

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• If q ranges in [2, 2.5] \rightarrow k ranges in ~ [40, 300]
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• If q ranges in $[2, 2.5] \rightarrow k$ ranges in ~ [40, 300]

$$\rightarrow k^{\frac{4}{3}} \rightarrow \sim 10^{-1} - 1$$

Production of γ -rays from p γ interaction

$$L_{\gamma} = \int \mathrm{d}r^3 \int_{E_{\gamma,min}}^{E_{\gamma,max}} \mathrm{d}E_{\gamma} E_{\gamma} \sum_{i} \epsilon^i_{\pi\gamma}(E_{\gamma})$$

$$\epsilon_{\pi\gamma}^{i}\left(E_{\gamma}\right) = \frac{A_{i}n_{\mathrm{p}}\left(\gamma_{\mathrm{p}}\right)}{\gamma_{\mathrm{p}}^{2}} \int_{\frac{\varepsilon_{l}'}{2\gamma_{\mathrm{p}}}}^{\infty} \mathrm{d}\varepsilon \frac{n_{\mathrm{ph}}(\varepsilon)\left(\frac{r_{\mathrm{cor}}}{d}\right)^{2}}{\varepsilon^{2}} f\left(\gamma_{\mathrm{p}},\varepsilon\right)$$

$$n_{\rm p}(\gamma_{\rm p}) \equiv \frac{dN}{dV d\gamma_{\rm p}} = n_0 \left(\frac{\gamma_{\rm p}}{\gamma_{0,\rm p}}\right)^{-q_{\rm p}} \quad \text{for} \quad \gamma_{\rm p,min} < \gamma_{\rm p} < \gamma_{\rm p,max}$$

$$U_{\rm p}^{\rm tot} = n_0 \int_{\gamma_{\rm p,min}}^{\gamma_{\rm p,max}} \mathrm{d}\gamma_{\rm p} n_{\rm p} \left(\gamma_{\rm p}\right) \le \frac{P_{\rm jet}}{4\pi r_{\rm blob}^2 c E_{0,\rm p}}$$

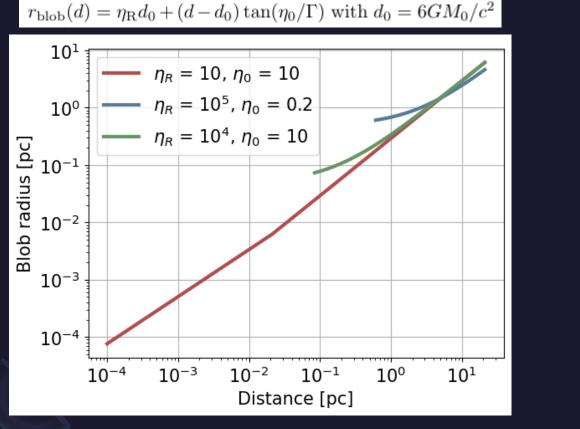
$$L_{\gamma} = V \int_{10^{-1} \text{GeV}}^{10^2 \text{GeV}} \mathrm{d}E_{\gamma} E_{\gamma} \sum_{i} \frac{A_{i} n_{0} \left(\frac{\gamma_{\mathrm{p}}}{\gamma_{0,\mathrm{p}}}\right)^{-q_{\mathrm{p}}}}{\gamma_{\mathrm{p}}^{2}} \int_{\frac{\varepsilon_{l}'}{2\gamma_{\mathrm{p}}}}^{\infty} \mathrm{d}\varepsilon \frac{n_{\mathrm{ph}}(\varepsilon) \left(\frac{r_{\mathrm{cor}}}{d}\right)^{2}}{\varepsilon^{2}} f(\gamma_{\mathrm{p}},\varepsilon) \,.$$

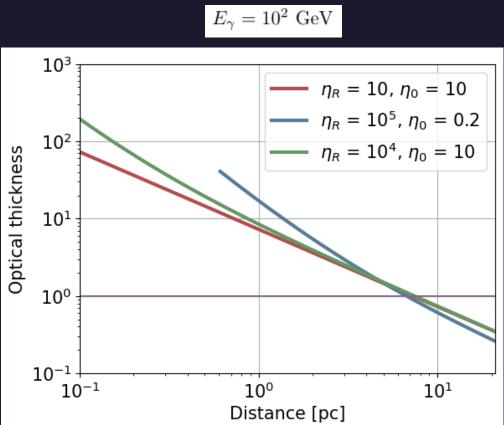
 $\gamma_{p,\mathrm{max}} \sim 1000$

 $P_{\rm jet} = 10^{43} \text{ erg/s}.$

 $\ll 10^{40.92} \text{ erg/s}$

Optical thickness





Conclusions

- The gamma-rays could be predicted by the jet component in the leptonic scenario
- Uncertainty on the radio data ______ uncertainty on the emitted luminosity

In particular: blob radius + spectral index

Lenain et al., 2008: Fermi-LAT gamma-rays predicted by the jet

blob radius of 6 pc \leftarrow 2.7 pc

background photon field: torus corona + disk

