

# Perturbative TQFTs

Konstantin Wernli

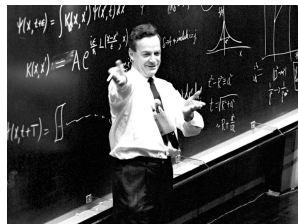
Centre for Quantum Mathematics, IMADA, SDU

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- 1 Introduction
- 2 Background
- 3 Perturbative TQFTs
- 4 Back to Chern-Simons and Evidence

In Quantum Mechanics, we have the *Feynman path integral*

$$\Psi(x_0, x_1) = \int_{\gamma: \substack{\gamma(0)=x_0 \\ \gamma(1)=x_1}} e^{\frac{i}{\hbar} S[\gamma]} D\gamma$$



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where  $\mathcal{L}: F_M \rightarrow \text{Dens}(M)$  is the Lagrangian density.

Classical physics are understood by solving the Euler-Lagrange equations,  $EL = \{\phi \in F_M, \delta S[\phi] = 0\}$ .



# Motivation: Functional integrals

In Quantum Field Theory, we are interested in the partition function

$$Z_M = \int_{\phi \in F_M} e^{\frac{i}{\hbar} S_M[\phi]} D\phi$$

and expectation values of observables

$$\langle O \rangle = \frac{1}{Z} \int_{\phi \in F_M} e^{\frac{i}{\hbar} S_M[\phi]} O(\phi) D\phi$$

Mathematical problem: Spaces  $F_M$  usually don't have sensible integration theories.

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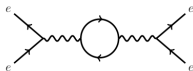
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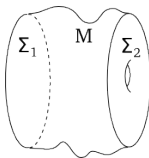
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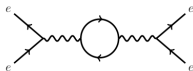


(b) Approach II

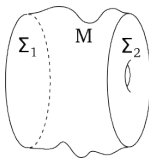
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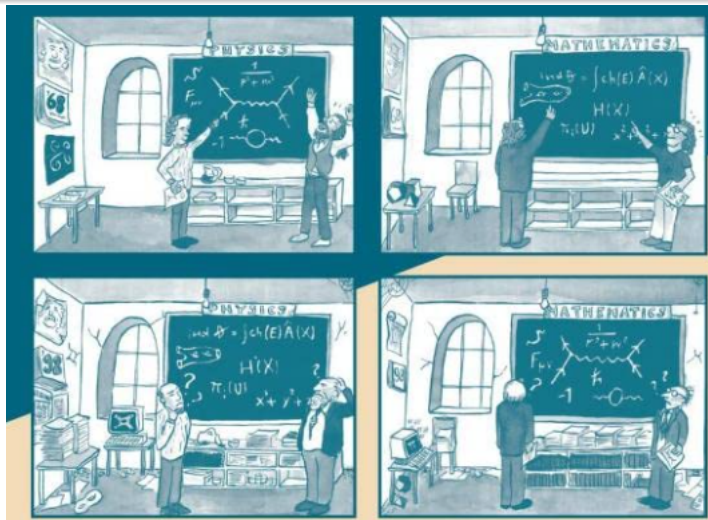


(a) Approach I



(b) Approach II

# Comic of the situation



# Chern-Simons theory: the drosophila melanogaster

Chern-Simons theory is the 3d gauge theory with

$$F_M = \Omega^1(M, \mathfrak{g}), \quad S_M = \int_M \frac{1}{2} \langle A, dA \rangle + \frac{1}{6} \langle A, [A, A] \rangle$$

where  $\mathfrak{g}, \langle \cdot, \cdot \rangle$  is a quadratic Lie algebra. The partition function

$$Z_M = \int_{F_M} e^{\frac{ik}{2\pi} S_{CS}(A)} DA$$

and Wilson loop observables for  $\gamma: S^1 \rightarrow M$

$$\langle O_\gamma \rangle_M = \frac{1}{Z} \int_{F_M} e^{\frac{ik}{2\pi} S_{CS}[A]} \left( \text{tr}_R P \exp \int_\gamma A \right)$$

were studied by Witten in 1989.



# Different Chern-Simons invariants

- Witten: For  $M = S^3$ ,  $\langle O_\gamma \rangle = J_{K(\gamma)}$  is the Jones polynomial of the knot  $K$ .
- Witten:  $Z_M$  are topological invariants of 3-manifolds.<sup>1</sup>
- Reshitikhin-Turaev described a *functorial* construction of similar invariants  $Z_{RT,M}^k$  based on quantum groups (Approach II).
- Fröhlich-King, Bar-Natan, Kontsevich, Axelrod-Singer, & many others studied the perturbative quantization  $Z_{CS,M}^{\text{pert}}$  of Chern-Simons (Approach I).

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<sup>1</sup>Witten considered framed 3-manifolds due to a 1-loop anomaly ▶

# Research Goals

Asymptotic Expansion Conjecture (AEC) [Witten, Andersen, ...]

Prove that as  $k \rightarrow \infty$ ,  $Z_{RT,M}^k \sim Z_{CS,M}^{\text{pert}}$ .

Wide Open Question (WOQ)

What is the analogous result for the Reshetikhin-Turaev functorial field theory?

Goal

Develop new (general) tools to combine perturbative QFT and functorial QFT to prove the WOQ and use WOQ to prove AEC.

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$$Z\left(\text{gluing diagram}\right) = \left\langle Z\left(\text{left boundary}\right), Z\left(\text{right boundary}\right) \right\rangle$$

Figure: Illustration of the gluing axiom

# Cobordism category

Reformulate in categorical language: Consider category **Cob** where

- Objects are  $d - 1$ -dimensional manifolds  $\Sigma$ ,
- Morphisms from  $\Sigma_1$  to  $\Sigma_2$  are  $d$ -manifolds  $\partial M$  with boundary  $\partial M = \Sigma_1 \sqcup \Sigma_2$ .

This category carries a symmetric monoidal structure (“tensor product”) given by the disjoint union.

## Definition

A TQFT is a symmetric monoidal functor

$$Z: \mathbf{Cob} \rightarrow \mathbf{Vect}.$$

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- The coefficients  $a_k(p)$  are given by sums over  $k$ -loop Feynman diagrams.



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- One then obtains a new symmetry, the BRST symmetry  $Q$  with  $Q^2 = 0$ .
- Only BRST cohomology is physically relevant.

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- Choose a gauge-fixing Lagrangian  $\mathcal{L} \subset \mathcal{F}_M$  such that  $\mathcal{S}_M$  restricted to  $\mathcal{L}$  has non-degenerate critical points.

# BV formalism II

Then one can prove that for a family of Lagrangians  $\mathcal{L}_t$  one has that

$$\frac{d}{dt} \int_{\mathcal{L}_t} e^{\frac{i}{\hbar} S_M} \mu^{\frac{1}{2}} = 0.$$

In particular, one has that

$$Z = \int_{F_M} e^{\frac{i}{\hbar} S_M[\phi]} \mathcal{D}[\phi] = \int_{\mathcal{L}} e^{\frac{i}{\hbar} S_M} \mu^{\frac{1}{2}}$$

when both sides are defined. Otherwise we use the perturbative expansion of the RHS as a definition.

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# What are classical TFTs?

For me, perturbative = semiclassical. Understand classical first.

- From a local  $d$ -dimensional field theory  $(F, S)$ , can extract a symplectic manifold  $F_\Sigma$  associated to a  $d - 1$ -dimensional manifold - the phase space.
- If  $\Sigma \subset \partial M$ , there is a map  $\pi: F_M \rightarrow F_\Sigma$  ("restriction")
- If  $M: \Sigma_1 \rightarrow \Sigma_2$  is a cobordism, then

$$(\pi_1 \times \pi_2)(EL) \subset F_{\Sigma_1} \times F_{\Sigma_2}$$

is a relation.

- In good cases, this is a canonical relation (lagrangian).

# The symplectic category

- This suggests “enhancing” the category of symplectic manifolds by allowing canonical relations as morphisms. Simply denote this category **Symp**.
- This idea was already studied Hörmander, Sniatycki-Tulczyjew, Weinstein, Guillemin-Sternberg... before TQFTs
- The main problem is that the composition of relations needs a transversality assumption
- Ignoring these problems, define classical TFTs as functor

$$F: \mathbf{Cob} \rightarrow \mathbf{Symp}$$



# Perturbative TQFTs - sketch of a definition

Let  $F: \mathbf{Cob} \rightarrow \mathbf{Symp}$  be a classical TFT.

## Data

A *perturbative TQFT quantizing  $F$*  assigns

- to  $d - 1$ -dimensional manifold  $\Sigma$  and a polarization  $\mathcal{P}$  of  $F_\Sigma$  a  $O_{F_\Sigma}[[\hbar]]$ -module  $\mathcal{H}_\Sigma^\mathcal{P}$
- to a  $d$ -dimensional manifold  $M$  with boundary an element  $\tilde{Z}_M \in H^{\text{top}}(EL_M) \otimes \mathcal{H}_{\partial M}^\mathcal{P}$ .

# Perturbative TQFTs - sketch of a definition

Here is an attempt to formalize the gluing axiom. Suppose  $M = M_1 \cup_{\Sigma} M_2$ .

Then we require that there is an open set of pairs of polarizations  $\{(\mathcal{P}_1, \mathcal{P}_2)\}$  of  $F_{\Sigma}$  with pairings

$$\langle \cdot, \cdot \rangle_{\mathcal{P}_1, \mathcal{P}_2, \Sigma} : \mathcal{H}_{\Sigma}^{\mathcal{P}_1} \times \mathcal{H}_{\Sigma}^{\mathcal{P}_2} \rightarrow \mathbb{C}[[\hbar]]$$

and

$$\pi_* \langle \tilde{Z}_{M_1}, \tilde{Z}_{M_2} \rangle = \tilde{Z}_M$$

where  $\pi : F_{M_1} \times F_{M_2} \rightarrow F_M$  and  $\pi_*$  denotes pushforward of forms.

# Perturbative TQFTs - remarks

- The module  $\mathcal{H}_\Sigma$  should be constructed as a “infinite level limit” of geometric quantization of  $F_\Sigma$ , the reduced phase space of the theory.
- The partition function is a top form on  $F_M$ . For a closed manifold  $M$ , the idea is that the integral of this top form

$$Z_M = \int_{EL_M} \tilde{Z}_M$$

(if it exists) is the formal power series describing the semiclassical asymptotics of the non-perturbative theory.

# BFV data

Goal: Extend BV formalism to spacetime manifolds with boundary.

To a  $d - 1$ -dimensional closed manifold  $\Sigma$  we associate a *BFV manifold* consisting of:

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Notice:  $\iota_{Q^{\partial}}\omega^{\partial} = dS^{\partial}$  for  $S^{\partial} = \iota_E \iota_{Q^{\partial}}\omega$  ( $E$  Euler vector field of  $M$ ).



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- a degree +1 vector field  $Q_M$ ,
- a surjective submersion  $\pi: \mathcal{F}_M \rightarrow \mathcal{F}_{\partial M}^{\partial}$  satisfying  $(Q_M)^2 = 0$ ,  $\delta\pi(Q_M) = Q_{\partial M}$  and

$$\iota_{Q_M}\omega + \delta\mathcal{S}_M = \pi^*\alpha^{\partial}.$$

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- Compatible with cutting and gluing of manifolds by taking fibered products  $\mathcal{F}_{M_1} \times_{\mathcal{F}_\Sigma} \mathcal{F}_{M_2}$
- Call corresponding functor  $\mathcal{F} : \mathbf{Cob} \rightarrow \mathbf{BFV}$  a *BFV functor*.

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- $\mathcal{H}_{\partial M}^{\mathcal{P}}$  is a certain “geometric quantisation” of the space of boundary fields.

# Perturbative Quantisation I: BV-BFV space of states

Let  $\mathcal{F}$  be a classical BFV functor, choose a polarisation  $\mathcal{P}$  of  $\mathcal{F}_{\Sigma}^{\partial}$ . BV-BFV space of states is a trivial vector bundle

$$\widehat{\mathcal{H}}_M^{\mathcal{P}} = \mathcal{M}_M \times (\mathcal{H}_{\partial M}^{\mathcal{P}}, \Omega^{\mathcal{P}})$$

where

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- $\mathcal{H}_{\partial M}^{\mathcal{P}}$  is a certain “geometric quantisation” of the space of boundary fields.
- $\Omega^{\mathcal{P}}$  is a coboundary operator which is a quantisation of  $S^{\partial}$ .

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## Conjecture

The  $(\hbar^2 \Delta + \Omega)$  cohomology gives a perturbative TQFT.

- 1 Introduction
- 2 Background
- 3 Perturbative TQFTs
- 4 Back to Chern-Simons and Evidence

# Reshetikhin-Turaev theories

- Ⓐ Mathematically, RT invariants  $Z_{RT}^k$  of a 3-manifold  $M = S_L^3$  are given by certain combinations of colored Jones polynomials of  $L$ , depending on a level  $k \in \mathbb{N}$ .
- Ⓑ By now it is understood that there is a 3-2-1 extended TQFT producing the RT invariants with  $Z_{RT}^k(S^1) = \mathbb{C}$ , a modular tensor category  $\mathbf{C}$ .
- Ⓒ Physically, these TQFTs correspond to Anyon models and form the basis for topological quantum computation

# The main conjecture

## Conjecture - How to prove WOQ

- There is a perturbative TQFT  $Z_{RT}^\infty$  such that

$$Z_{RT,M}^k \sim_{k \rightarrow \infty} Z_{RT,M}^\infty$$

- This perturbative TQFT coincides with the perturbative TQFT defined by the  $(\hbar^2 \Delta + \Omega)$ -cohomology of the perturbative quantization of the Chern-Simons BFV functor

Let me present some evidence of this fact.

## Some evidence for part A

- Jeffrey, Freed-Gompf, Rozansky, Andersen and many others worked on the asymptotic behaviour of  $Z_{RT}^k$ .
- In particular, Andersen-Hempel and Andersen-Mistegaard show that on finite order mapping tori the asymptotics are given by integrals of top forms over the moduli space of flat connections
- For manifolds with boundary, Andersen and many collaborators show that the state spaces of  $Z_{RT}^k$  can be computed via geometric quantization of moduli spaces of flat connections at level  $k$ .
- This is evidence that one can indeed understand the asymptotics of RT as a perturbative TQFT, but no results for manifolds with boundary exist.



# Evidence for part B

- Rozansky proved a certain class of manifolds  $M$  that

$$Z_{RT,M}^k \sim_{k \rightarrow \infty} \sum_{x \in S_{CS}(M)} e^{ikx} T_x \left( 1 + \sum_{l=1}^{\infty} r_{x,l} k^{-l} \right)$$

(the AEC at first order)

- For  $x = 0$ ,  $r_0$  can be identified with the LMO invariant

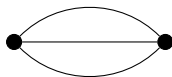
This gives a conjectural chain of equalities

$$(Z_{RT}^{\infty})_{x=0} = Z_{LMO} \stackrel{?}{=} Z_{KKTl} = Z_{BC} \stackrel{?}{=} Z_{BV}^{pert}$$

# Theta invariants

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At level 1 (2-loops) the first equality is known, and the invariant known as theta (or sunset) invariant:

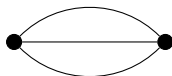


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## Theorem

[Kuperberg-Thurston '99, Lescop 2000's] The theta invariant equals the Casson invariant up to a framing correction.

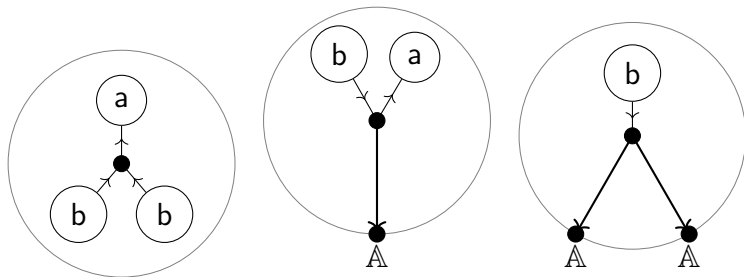
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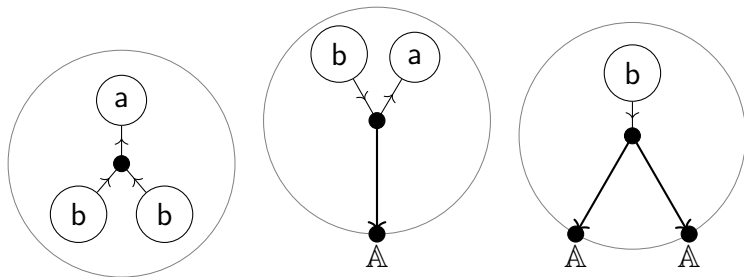
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Rules: e.g. Last one gives

$$\psi_{\Gamma} = g_i^{jk} \int_{C_2(\partial M)} (\pi_{1,*} b_1^i \eta_{12} \eta_{13}) \mathbb{A}_j \mathbb{A}_k$$



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- $(D \times S^1) \cup_{\varphi} (D \times S^1) \cong L_{p,q}$
- In particular, diffeomorphism type independent of  $m, n$  and only depends on  $q \bmod p$ .

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**Theorem [W., Cattaneo-Mnev-W.]**

The theta invariant (2-loop contribution to the state) of  $L_{(p,q)}$  is

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Agrees with Results by Kuperberg-Thurston-Lescop since the Casson(-Walker-Lescop) invariant of  $L(p, q)$  is  $s(q, p)$ .