A BPS Road to Holography: Decoupling Limits and Non-Lorentzian Geometries

14Danish Quantum Field Theory Meeting, SDU, Odense, August 13, 2025 Niels Obers (Niels Bohr Institute)





based on work:

2410.03591 (JHEP), 2311.10564 (PRL), 2502.20310 (JHEP)

(Blair, Lahnsteiner, NO, Yan)

2506.19720 (Blair, NO, Yan)

2501.10178 (JHEP) (Harmark, Lahnsteiner, NO)

& earlier papers with Harmark, Hartong, and Bidussi/Menculini/Oling/Yan

Introduction/Motivation

- D-branes & their BPS nature underly major advances in ST
 - AdS/CFT correspondence (dual descriptions of D-branes)
 - matrix theory proposal (large N limit of D-brane wv. theory)
- → decoupling limits probe rich physics
 - simplification by removing part of spectrum
 - access to non-perturbative regimes

revisit decoupling limits in light of recent advances

- intruiging relations to non-Lorentzian corners of string theory e.g. NRST as a simpler corner of (relativistic) ST

Gomis,Ooguri(2000),Danielsson,Guijosa,Kruczenski(2000) Harmark,Hartong,NO(2017)/Bergshoeff,Gomis,Yan(2018) /& many papers since

- new boostrap techniques/study of amplitudes in BFSS

Han, Hartnoll, Kruthoff (2020) / Dorey, Mouland, Zhao (2022) / Mouland (2023), Tropper, Wang (2023) / Herderschee, Maldacena (2023) / Komatsu et al (2024)...

Questions

1. What are the guiding principles for mapping out self-consistent decoupling limits in ST?

Guided by a BPS road..

- 2. In the context of holography, what is the role of the tendimensional non-Lorentzian (NL) geometry coupled to matrix theory on the D-branes (or NRST for F-strings)?

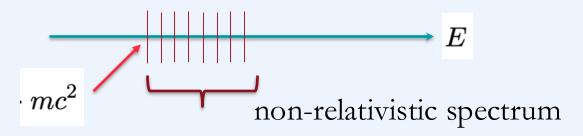
 To Holography ...
- 3. How is the ten-dimensional bulk geometry generated intrinsically? And back via TTbar

related work on holography and NL geometries:

Avila, Guijosa, Olmedo (2023) Lambert, Smith (2024 (3)) Fontanella, Nieto Garcia (2024 (3)) Guijosa (2025) Bergshoeff, Lambert, Smith (2025) Bergshoeff, Grosvenor, Romano, Yan (2025)

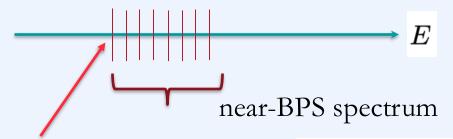
Part 1: BPS decoupling limits

Anatomy of BPS decoupling limits



$$E = \sqrt{(mc^2)^2 + (pc)^2} - \Phi q = \sqrt{(mc^2)^2 + (pc)^2} - 1 \cdot mc^2 \simeq \frac{p^2}{2m}$$

- same for near-BPS limit



$$E = \sqrt{M_1^2 + M^2} - A_1 M_1 \simeq \frac{M^2}{2M_1}$$
 $M_1 \to \omega^2, M \to \omega, \omega \to \infty$, $A_1 \to 1$

- scale to supersymmetric background, cancel divergence with critical electric field

Matrix Theory: A BPS perspective

Matrix theory arises from BPS decoupling limit of D0-brane wv.

→ relies on fact that tension=charge & fine tuning RR gauge field

$$S_{\rm D0} = -\frac{1}{g_{\rm s} \sqrt{\alpha'}} \int d\tau \sqrt{-\dot{X}^{\mu} \dot{X}_{\mu}} + \frac{1}{\sqrt{\alpha'}} \int C^{(1)}, \qquad \mu = 0, \dots, 9.$$

take limit:

$$g_{\mathrm{s}} o \omega^{-3/2} \, g_{\mathrm{s}} \,, \qquad X^0 o \sqrt{\omega} \, X^0 \,, \qquad X^i o rac{X^i}{\sqrt{\omega}} \,, \qquad C^{(1)} o \omega^2 \, g_{\mathrm{s}}^{-1} \, \mathrm{d} X^0 \,,$$

$$\omega \, o \, \infty \,.$$

gives NR action:
$$S = \frac{1}{g_{\rm s} \sqrt{\alpha'}} \int \mathrm{d}\tau \left(\frac{1}{2} \, \dot{X}^i \, \dot{X}^i + 2 \, \psi^{\intercal} \, \dot{\psi} \right),$$

BFSS

non-abelian (& SUSY) generalization is BFSS:

$$S_{\mathrm{BFSS}} = \frac{1}{R} \int \mathrm{d}\tau \, \mathrm{tr} \left[\tfrac{1}{2} \dot{X}^i \, \dot{X}^i + \tfrac{1}{4} \big[X^i \, , \, X^j \big] \big[X_i \, , \, X_j \big] + 2 \Big(\psi^\intercal \, \dot{\psi} - \psi^\intercal \, \gamma^i \big[\psi \, , \, X^i \big] \Big) \right],$$

- BFSS (1996) conjecture: N,R to infinity = M-theory on flat spacetime
- Susskind (1997) fixed N: DLCQ of M-theory

-Seiberg/Sen (1997): light like circle as infinite boost of spacelike circle



Prelude: Non-Lorentzian geometry

start with flat relativistic metric

$$\mathrm{d} s^2 = \omega \, \mathrm{d} x^A \, \mathrm{d} x^B \, \eta_{AB} + \omega^{-1} \, \mathrm{d} x^{A'} \, \mathrm{d} x^{A'} \,,$$

A = 0..p longitudinal directions A' = p+1,...9-p transverse directions

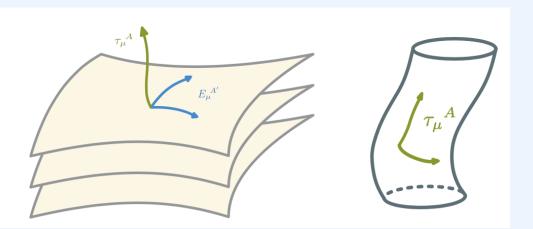
$$\omega \to \infty$$
 geometry \to non-Lorentzian: p-brane Newton-Cartan (NC) geometry

 \rightarrow local SO(1,p) and SO(9-p) geometry & p-brane boost symmetry

$$\delta_{
m G} x^A = 0 \,, \qquad \delta_{
m G} x^{A'} = \Lambda^{A'}{}_A \, x^A \,,$$

generalize to curved bgrs. $dx^A \to \tau_\mu{}^A dx^\mu \qquad dx^{A'} \to E_\mu{}^{A'} dx^\mu$

$$\mathrm{d}x^{A'} \to E_{\mu}{}^{A'} \, \mathrm{d}x^{\mu}$$



$$au_{\mu
u}= au_{\mu}{}^A\, au_{
u}{}^B\,\eta_{AB}\,,$$
 $E_{\mu
u}=E_{\mu}{}^{A'}\,E_{
u}{}^{A'}\,.$

Matrix 0-brane theory (M0T)

DLCQ M-Theory

lightlike compactification M0T: BFSS Matrix theory

can go beyond the limit in flat space:

curved space of M0T

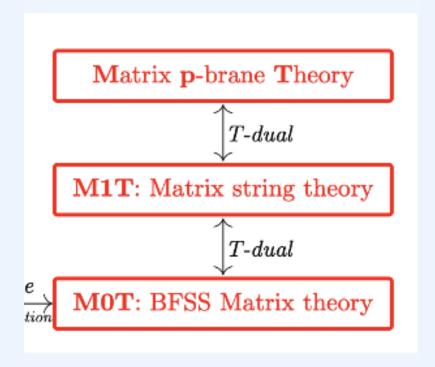
$$G_{\mu\nu} = \omega \, \tau_{\mu\nu} + \omega^{-1} \, E_{\mu\nu} \,, \qquad \Phi = \varphi - \frac{3}{2} \, \ln \omega \,, \qquad B^{(2)} = b^{(2)} \,,$$
 $C^{(1)} = \omega^2 \, e^{-\varphi} \, \tau^0 + c^{(1)} \,, \qquad C^{(q)} = c^{(q)} \quad \text{if } q \neq 1 \,.$ $\omega \to \infty$.

Matrix theory of D0-branes (M0T) couples to non-Lorentzian 10D geometry: M0T target space = Newton-Cartan geometry + other string fields

$$S_{\rm D0}^{\rm M0T} = \frac{1}{2\sqrt{\alpha'}} \int {\rm d}\tau \, e^{-\varphi} \, \frac{\dot{X}^{\mu} \, \dot{X}^{\nu} \, E_{\mu\nu}}{\dot{X}^{\mu} \, \tau_{\mu}{}^{0}} + \frac{1}{\sqrt{\alpha'}} \int c^{(1)} \, . \label{eq:SD0T}$$

- as opposed to earlier work: now applied to the full type II string theory containing all possible extended objects and in general backgrounds

MpT in curved spacetime



via T-duality → MpT = near-BPS decoupling limit of Dp-branes

MpT target space = p-brane NC geometry + other string fields

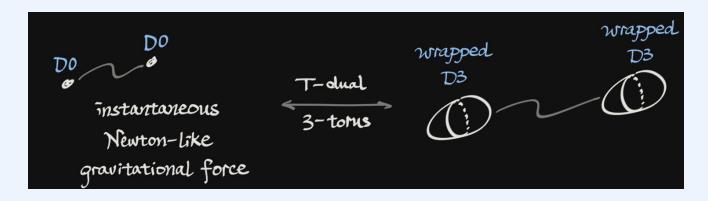
Dp-branes in MpT

can study wv. actions of generic Dq-branes and F-string in MpT

- form of actions related to BPS structure
- e.g. Dp-branes fundamental dof in MpT:

$$S_{{
m D}p}^{{
m M}p{
m T}} = -rac{T_p}{2} \int {
m d}^{p+1} \sigma \, e^{-arphi} \, \sqrt{- au} \, au^{lphaeta} \, E_{lphaeta} + T_p \int c^{(p+1)} \, .$$

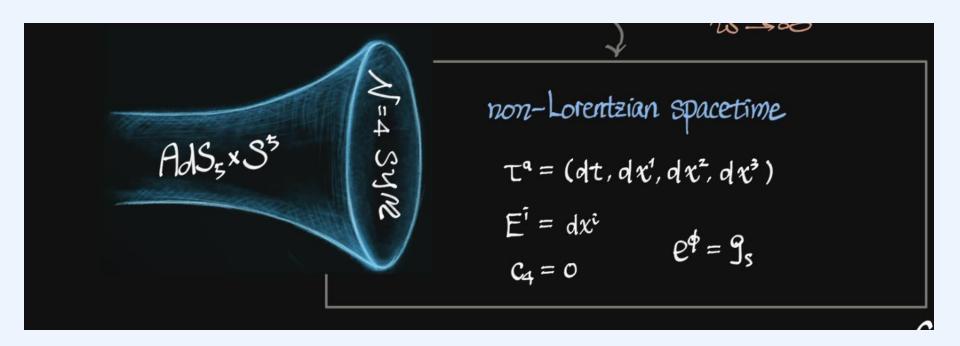
- wrapped Dp-branes behave like NR particles, experiencing Newton-like forces



- especially interesting is p=3:

M3T \rightarrow N=4 SYM sees 10D non-Lorentzian spacetime

Part 11: Application to Holography in string theory



BPS decoupling = near-horizon in bulk

- BPS decoupling limit at asymptotic infinity
 - → N=4 SYM coupled to non-Lorentzian M3T geometry

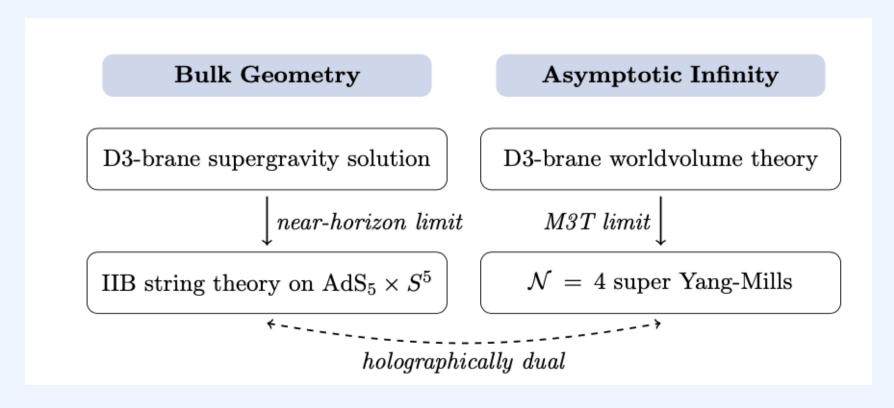
$$egin{aligned} \left(X^0, \, \cdots, \, X^p
ight) &= \omega^{rac{1}{2}} \left(t \, , \, x^1, \, \cdots, \, x^p
ight), \ \\ \left(X^{p+1}, \, \cdots, \, X^9
ight) &= \omega^{-rac{1}{2}} \left(x^{p+1}, \, \cdots, \, x^9
ight), \ \\ G_{
m s} &= \omega^{rac{p-3}{2}} \, g_{
m s} \end{aligned}$$

$$C^{(p+1)} = \omega^2 g_{\mathrm{s}}^{-1} \, \mathrm{d}t \wedge \cdots \wedge \mathrm{d}x^p \,.$$

• same asymptotic BPS decoupling limit, applied to D3-brane geometry generates the near-horizon limit!

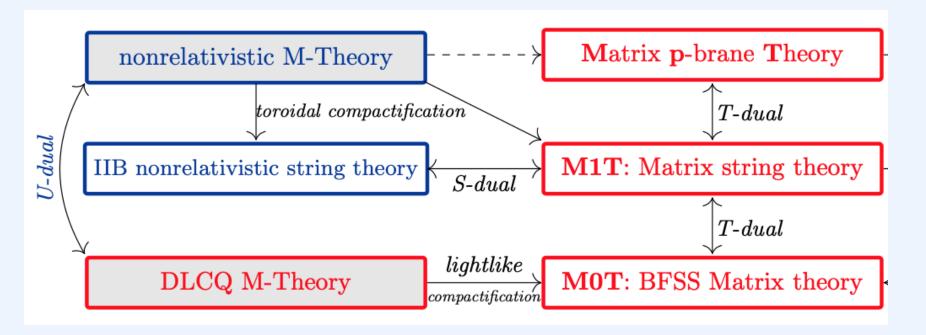
$$\mathrm{d} s^2 = \Omega(r)\,\mathrm{d} x^A\,\mathrm{d} x^B\,\eta_{AB} + \frac{\mathrm{d} r^2 + r^2\,\mathrm{d}\Omega_{8-p}^2}{\Omega(r)}\,, \qquad \quad \Omega(r) = \left(\frac{r}{\ell}\right)^{\frac{7-p}{2}},$$

AdS/CFT



- applies to general Dp-brane solutions as well as F-string (IMSY)
- near-horizon geometries themselves asymptotically approach an MpT limit
 - yields back NL geometry seen by SYM at asymptotic infinity
- geometrically explains: relationship between matrix theory & AdS/CFT decoupling

Connections to NR string theory/M-theory



- near crtical D1-string dual to near critical F1-string (=NRST)
- M-theory uplift (NR M-theory): near critical M2-brane
 - wrapped membranes WM2
 - open M2-branes on M5 (OM) (studied in 2000s)



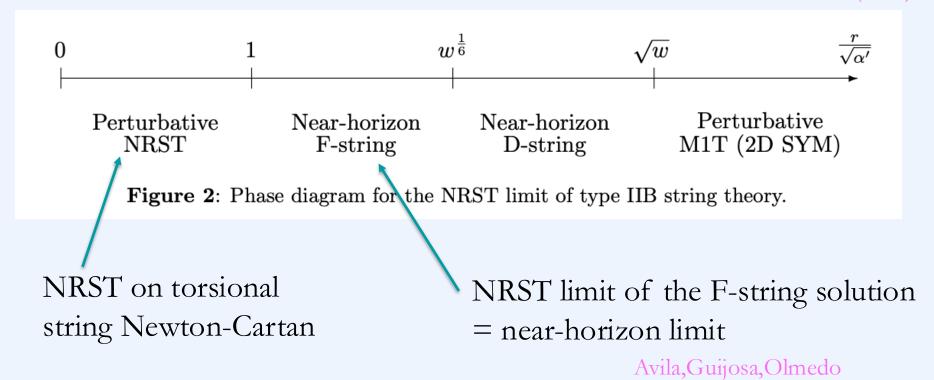
Blair, Gallegos, Zinnato (2021) Ebert, Sun, Yan (2021) Ebert, Yan (2023)

Lambert, Smith (2024)

see also:

Example of NRST in holography

Harmark, Lahnsteiner, NO (2025)



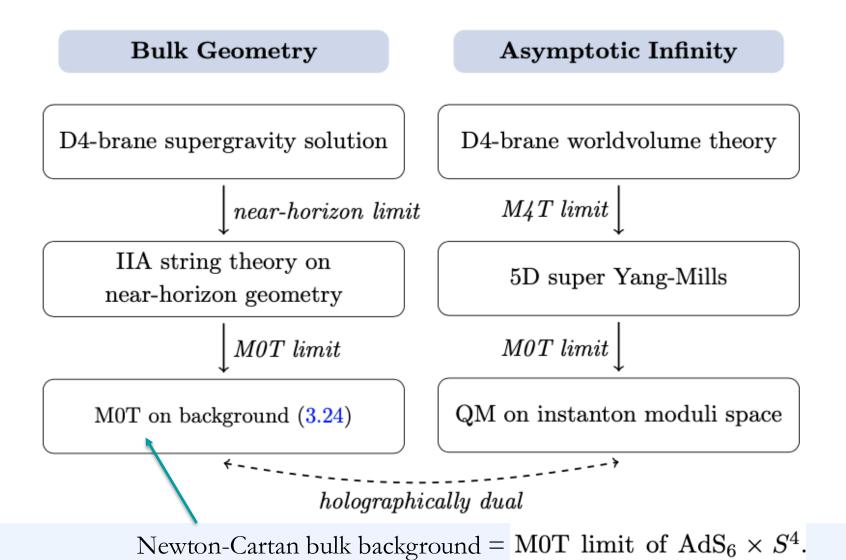
matrix string theory: 2nd quantization of NRST describes multi string states, joining/splitting via interaction term

Dijkgraaf, Verlinde, Verlinde/Motl

Landscape of holographic duals

- can use this perspective to generate
 new examples of holographic bulk geometries
- → can do asymptotic MpT limit and apply to bulk Dq-brane with (p not q)
 - involves double BPS decoupling limit zooming in on intersecting background bound states

Example of multi-critical BPS limit



many more duals can be generated: involving NL geometries in the bulk!

see also: Lambert, Smith (2024)

M-theory uplift and U-duality orbits

can unify the BPS deocupling limit using M-theory and U-dualities

- M0T BPS decoupling limit = DLCQ in M-theory
 - single BPS decoupling limits lie in U-duality orbit for ½ BPS (one DLCQ)
 - double BPS decoupling limits lie in U-duality orbit for ½ BPS (two DLCQ's)

Blair, Lahnsteiner, NO, Yan (2023) Gomis, Yan (2023) Dijkgraaf, de Boer, Harmark, NO (unpublished)

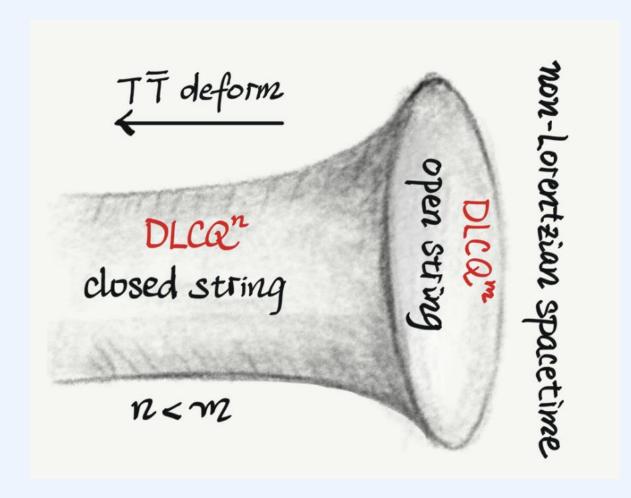
Holography as DLCQⁿ/DLCQ^m correspondence

→ holographic dulaities unified as DLCQⁿ/DLCQ^m with m > n

- AdS_5/CFT_4 = DLCQ^0/DLCQ^1:
 DLCQ at asympt. inf = near-horizon
 similarly for the other Dp-branes and NRST
- $AdS_3/CFT_2 = DLCQ^0/DLCQ^2$

- novel holographic dualities with NL geometry in bulk: DLCQ^1/DLCQ^2

Part III: Undoing the BPS limit → TTbar like deformation



Generating near-horizon bulk geometry

- intrinsic perspective on relation between asymptotic infinity and bulk geometry
 - can we `invert' the BPS decoupling limit?

Yes: related to TTbar deformation

$$rac{\partial \mathcal{L}(m{t})}{\partial m{t}} \sim \det T_{lphaeta}(m{t}) \,,$$

$$T_{\alpha\beta}(\boldsymbol{t}) = -\frac{2}{\sqrt{-\det \tau}} \frac{\partial \mathcal{L}(\boldsymbol{t})}{\partial \tau^{\alpha\beta}}$$

NRST can be TTbar deformed back to relativistic string theory

Blair (2020)

New Dp-brane TTbar

- → via duality web: new Dp-brane TTbar deformations
- deformations that induce a flow from SYM to DBI action (so far: abelian)

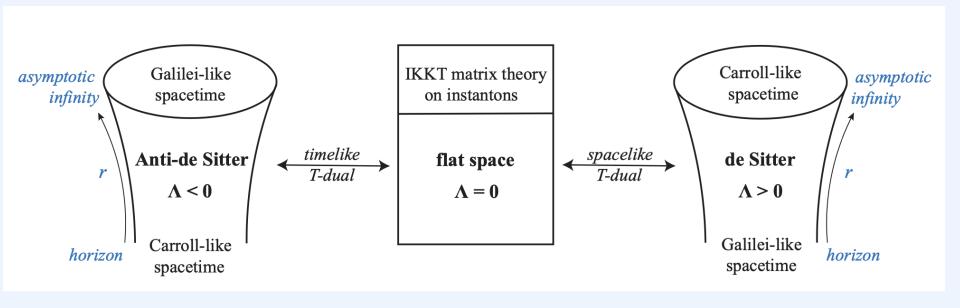
- universal result for (p+1)-dimensional free scalar field theory to DNG

$$\frac{\partial \mathcal{L}}{\partial t} = \frac{\sqrt{-\det \tau}}{2t^2} \left\{ \operatorname{tr} (\mathbb{1} - t \mathcal{T}) - (p-1) \left[\det (\mathbb{1} - t \mathcal{T}) \right]^{\frac{1}{p-1}} - 2 \right\}$$

- also expressions including worldvolume gauge fields
- checks with various results from other methods

Carroll geometry meets dS via holography

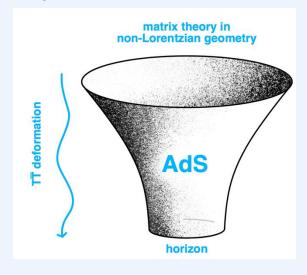
Blair, NO, Yan (2025)



- extends AdS case to (hypotethical) dS holography in string theory
- elaborates on work of Hull (1998): E-branes in type IIB*
- flat space corresp. of IKKT = T-dual to geometry with (conformal) dS
- geometry at asymptotic infinity is Carroll-like

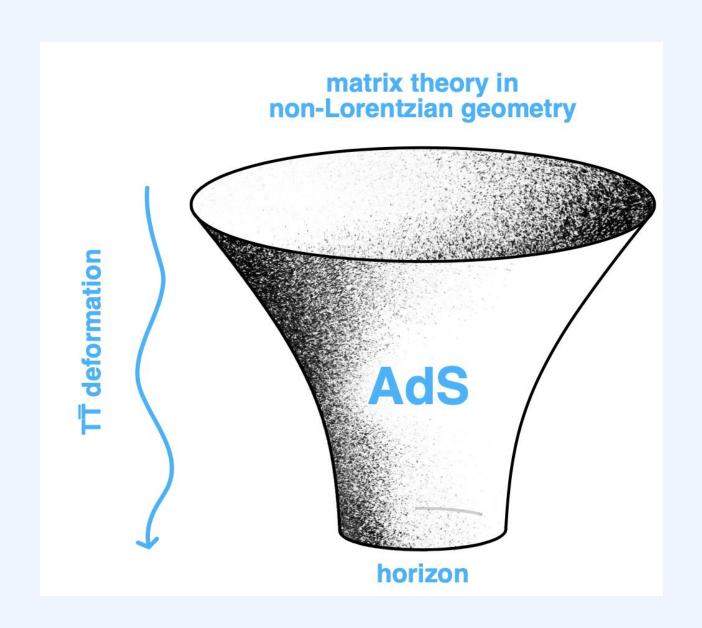
Outlook

- correspondence between SUGRA and matrix QM: revisit
- elements of the AdS/CFT correspondence: deeper insight into relationships between brane configs and possible decoupling limits



- -non-Lorentzian holography: solutions, EOMs, torsions constraints
- p-brane generalisations of TTbar deformation
- algebraic aspects of decoupling (novel IW contactions of superalgebra)
- matrix valued TSNC geometry? (in 2nd quantized NRST)
- extensions: tensionless, Carrollian/dS, heterotic

The end



Extra slides after this

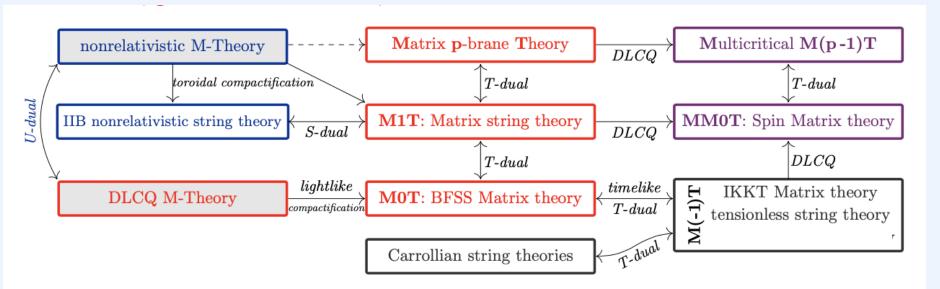
Unified web of non-Lorentzian string theories

obtained from insights using

Blair, Lahnsteiner, NO, Yan (2023) Gomis, Yan (2023)

- non-relativistic string theory on curved spactimes
- applying string theory U-dualities
- anatomy of BPS mass formulae and decoupling limits

de Boer, Dijkgraaf, Harmark, NO (2000s, unpublished)



each node describes decoupled theory including actions for:

- fundamental (light) degrees of freedom coupling to an appropriate (non-Lorentzian) target spacetime
- other (heavy) probe objects in the theory

Lessons from the web

- near-BPS limits in ST can be viewed as natural generalizations of non-relativistic point particle limit
- novel perspective on Matrix theory/NRST via non-Lorentzian
 backgrounds → new insights into holography
- new corners of the duality maze (via solution generating techniques): IKKT, Carrollian string theory,..
- multi-critical limits & relation to spin-Matrix theory
- limits and geometries can also be understood from novel IW contractions of the M-theory superalgebra

Remarks on MOT

objects in M0T

- D0-branes: BFSS theory (on curved Newton-Cartan bgr)
- D2-branes: NCYSM
- F1-string: novel action (target space Galilean/world-sheet Carrollian boost)

important/relevant later: emergent dilation symmetry for (sufficiently) well-behaved function Delta

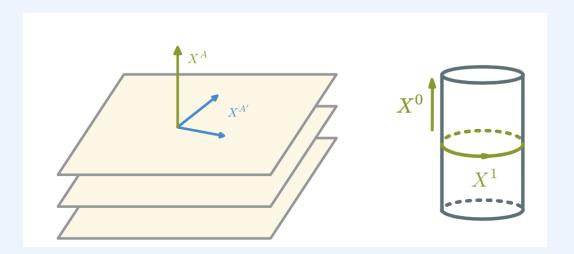
;
$$\omega o \omega \, \Delta^{-1}$$

$$au^0 o \Delta^{rac{1}{2}} \, au^0 \,, \qquad E^i o \Delta^{-rac{1}{2}} \, E^i, \qquad e^arphi o \Delta^{-rac{3}{2}} \, e^arphi \,.$$

underlying symmetry from (two) algebraic perspectives: extension of super-Bargmann altebra with extra bosonic charges can also see as D0-brane adapted IW contraction of relativistic IIA superalgebra

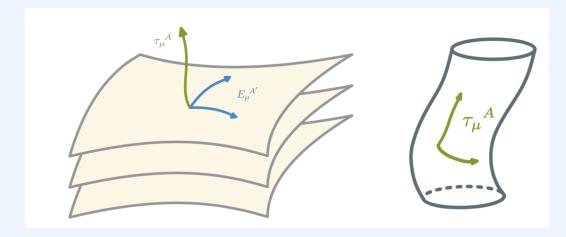
How does MpT relate to NR string theory (analogous decoupling limit)?

Gomis/Ooguri NR string lives in flat space



figures from review on NRST Gerben Oling & Ziqi Yan (2202.12698)

General target space probed by NR strings



multi-critical examples - NL geometry in bulk

- multicritical limit adapted to D0-D4 → bulk M0T geometry

$$au_{\mu\nu} \,\mathrm{d}\mathtt{z}^{\mu} \,\mathrm{d}\mathtt{z}^{
u} \equiv au^{A} \, au^{B} \,\eta_{AB} = -\left(\frac{\mathbb{\Gamma}}{\widetilde{\ell}}\right)^{\frac{3}{2}} \,\mathrm{d}\mathtt{z}^{0} \,\mathrm{d}\mathtt{z}^{0},$$

$$E_{\mu\nu}\,\mathrm{d}\mathbf{z}^{\mu}\,\mathrm{d}\mathbf{z}^{\nu} \equiv E^{A'}E^{A'} = \left(\frac{\widetilde{\ell}}{\mathbb{r}}\right)^{\frac{3}{2}} \left[\left(\frac{\mathbb{r}}{\widetilde{\ell}}\right)^{3} \mathrm{d}\mathbf{z}^{I} \mathrm{d}\mathbf{z}^{I} + \mathrm{d}\mathbb{r}^{2} + \mathbb{r}^{2}\,\mathrm{d}\Omega_{4}^{2} \right]$$

M0T limit of $AdS_6 \times S^4$.

multi-critical examples - NL geometry in bulk

- multicritical limit adapted to D1-Dp → bulk torsional string Newton-Cartan

$$\tau_{\mu\nu}dx^{\mu}dx^{\nu} = -\left(\frac{r}{\ell}\right)^{(6-p)/2}dt^2 + \left(\frac{r}{\ell}\right)^{-(6-p)/2}dv^2$$

$$E_{\mu\nu}dx^{\mu}dx^{\nu} = \left(\frac{r}{\ell}\right)^{(6-p)/2} \sum_{a=1}^{p} (dy^{a})^{2} + \left(\frac{r}{\ell}\right)^{-(6-p)/2} \sum_{i=1}^{8-p} dr^{i} dr^{i} ,$$

$$a_{(p+1)} = \left(\frac{r}{\ell}\right)^{6-p} dt \wedge dy^1 \wedge \dots \wedge dy^p , \quad g_s e^{\varphi} = g_s \left(\frac{r}{\ell}\right)^{-(6-p)(3-p)/4}$$

$$\ell^{6-p} = \frac{(2\pi\sqrt{\alpha'})^{7-p}g_s N}{(6-p)\Omega_{7-p}R_v}$$

- torsion constraints from SUSY?
- can also consider NS5-F1 \rightarrow reproduces solition that is earlier found

Duality asymmetry of NL geometries

Blair, Lahnsteiner, NO, Yan (2025)

- rich structure:
- T- and S-duality transformations exhibit novel asymmetric properties:

depending on choice of transformation & value of the background fields, codimension of foliation structure of NL geometry can change/stay same)

- duality asymmetry between NL geometries for Matrix theory underlying Morita equivalence in matrix theory and non-commutative YM
- Hashimoto-Itzhaki/Maldacena-Russo holography for noncommutative YM with B-field & generalizations

Nodal Riemann spheres

interesting degenerate topologies of 2D surfaces show up both on world-sheetand target space

Gomis, Yan (2023)

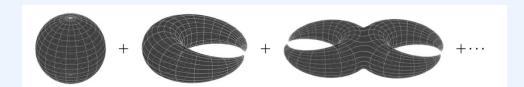
- F-string in M0T:

modular parameter

 $\tau \to i \infty$.

genus expansion (cf. ambitwistor)

 $\underline{pinching},$



Geyer, Mason, Monteiro, Tourkine (2015)

- pinched torus in target space:

