

Learning Standard Model structure from LHC data with Riemannian Flow Matching

Based on [2607.16144]

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in collaboration with Midori Kato, Inar Timiryasov, and Oleg Ruchayskiy

Conference/Workshop

HAMLET – PHYSICS



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Website of the paper: <https://hep-ssl-webapp.pages.dev>



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2 Method

3 Results

4 Conclusions

5 Backup slides

Introduction



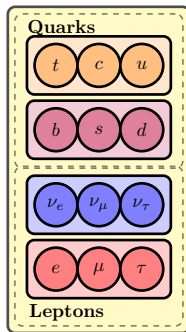
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Introduction: particle physics

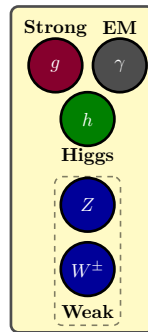
Particle physics deals with **fundamental particles**. The framework we use for particle physics is based on relativistic quantum mechanics (quantum field theory).

Interactions between particles are described by the **Standard Model**. Every interaction and particle the SM has predicted we have found, the last piece of the puzzle was the discovery of the Higgs boson.

Fermions



Bosons

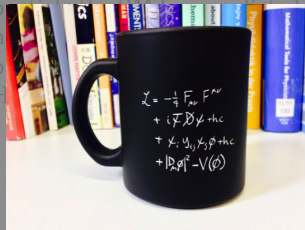


Introduction: particle physics

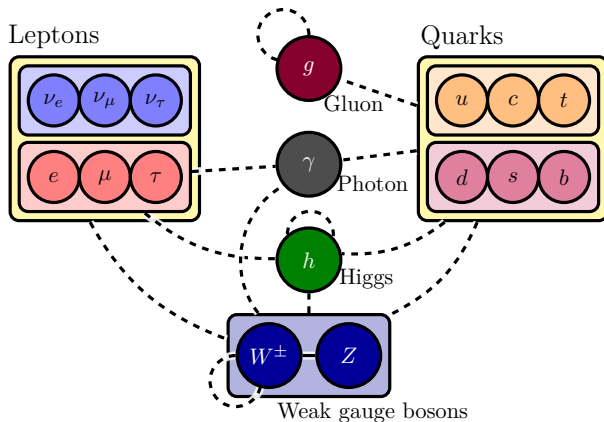
Particle physics deals with

$$\begin{aligned}\mathcal{L}_{\text{SM}} = & -\frac{1}{2} \text{tr}[\mathbf{W}_{\mu\nu} \mathbf{W}^{\mu\nu}] - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{2} \text{tr}[\mathbf{G}_{\mu\nu} \mathbf{G}^{\mu\nu}] + \sum_{\Psi} \bar{\Psi} \not{D} \Psi + |D_{\mu} H|^2 \\ & - \left(\sum \bar{L}_L \cdot H Y_{\ell} \ell_R + \sum \bar{Q}_L \cdot H Y_D D_R + \sum \bar{Q}_L \cdot \tilde{H} Y_U U_R + \text{H.c.} \right) \\ & + \mu^2 |H|^2 - \lambda |H|^4,\end{aligned}$$

Small enough to fit in a mug!



Introduction: particle physics



LHC experiment and ATLAS detector

The **Large Hadron Collider** is a 27 km radius proton – proton collider. It operates at a center of mass energy of 14 TeV.

It records tens of petabytes of data per year. It is currently on shutdown, but the next phase, the high luminosity LHC will record much (much!) more data.

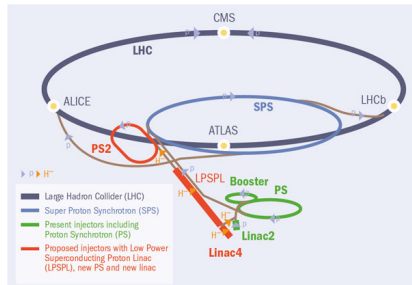


Image from CERN Courier

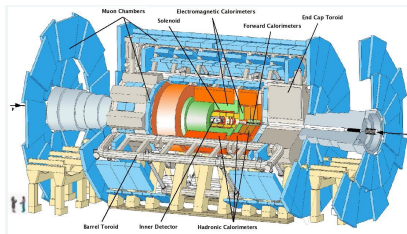


Image from UC Irvine

ATLAS

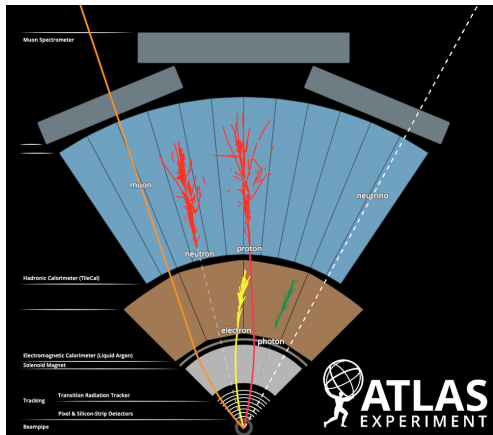
- A general-purpose detector
- Some objectives include: precision measurement, discovery of new physics...
- It's like an onion (or ogre), it has layers

ATLAS Data Collection

Again, it's like an onion (or an ogre)

ATLAS detects:

- **Electrons:** interact in the tracker and ECAL
- **Muons:** travels the whole detector
- **Jets:** interact until the hadronic calorimeter.
Types of *jets*:
 - **Small-R jets**
 - **Large-R jets**
 - **Hadronic τ**
- **Photons:** only interact ECAL



ATLAS released some data from Run 2, for anyone to access it, and do any analysis with it. They claim it is mostly for education purposes, but researchers have done analyses on them (like us!)

Dataset	Selection requirement	Size [GB]	Events	Frac. [%]
Full release	No additional selection	1200	5,770,706,892	100.00
1LMET30	≥ 1 lepton, $p_T > 7$ GeV; $E_T^{\text{miss}} > 30$ GeV	133.4	605,490,333	10.49
2to4lep	2–4 leptons, all with $p_T > 7$ GeV	74.2	283,951,733	4.92
2bjets	2 jets $p_T > 20$ GeV, ≥ 1 b -tagged	36.2	141,259,550	2.45
3J1LMET30	≥ 3 jets $p_T > 20$, ≥ 1 lepton $p_T > 3$, $E_T^{\text{miss}} > 30$ GeV	36.6	76,407,431	1.32
2muons	≥ 2 muons $p_T > 10$ GeV	18.3	72,563,712	1.26
3lep	≥ 3 leptons $p_T > 7$ GeV	5.5	17,193,618	0.30
exactly3lep	exactly 3 leptons $p_T > 7$ GeV	4.8	15,677,770	0.27
2J2LMET30	≥ 2 jets $p_T > 20$, ≥ 2 leptons $p_T > 7$, $E_T^{\text{miss}} > 30$ GeV	1.9	6,242,521	0.11
4lep	≥ 4 leptons $p_T > 7$ GeV	0.627	1,515,848	0.03
exactly4lep	exactly 4 leptons $p_T > 7$ GeV	0.555	1,362,611	0.02

You can access the datasets in [this link](#). We worked with the 2to4lep and 1L30MET datasets.

Working with these datasets was physics motivated

- The 2to4lep dataset contains known resonances in the invariant mass spectrum: $\Upsilon, J/\psi, Z$
- The 1L30MET dataset contains the cleanest channel for W and top quark

Generative models of continuous data

As of 2026, the dominant paradigm for generation of continuous data is **transport-based generative modeling**. These models construct a continuous path between two distributions (e.g. from noise to data).

Conditional flow matching

Model learns a velocity field u_t^θ whose flow transports the noise distribution into the data distribution:

$$dX_t = u_t^\theta(X_t) dt, \quad X_0 \sim p_{\text{init}}, \quad 0 \leq t \leq 1,$$

so that $X_1 \sim p_{\text{data}}$.

Our loss is:

$$\mathcal{L}_{\text{CFM}}(\theta) = \mathbb{E}_{t,z,x} \left\| u_t^\theta(x) - u_t^{\text{target}}(x | z) \right\|^2$$

Example of a conditional path: $x_t = (1 - t)\epsilon + tz$ where $\epsilon \sim \mathcal{N}(0, I)$. The target velocity is $u_t(x | z) = z - \epsilon$.

Generative models of continuous data

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Applications in image generation:

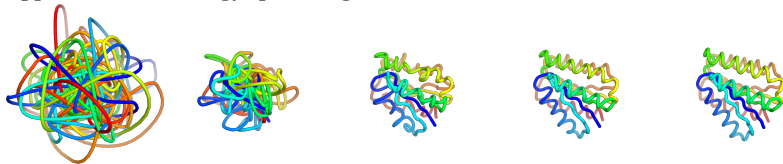


Images generated with FLUX.2 [klein] 4B

Generative models of continuous data

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Applications in biology: protein generation



Protein generated with FoldFlow

Our motivation and novelty

Open question:

Can a single generative model successfully capture physics recorded at the Large Hadron Collider?

Why should we care?

- The next LHC runs will record an insurmountable amount of data, and we must take as much advantage of it
- Our model can provide a further step towards bigger analyses on LHC data

Method



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Generative models for collider physics

Our goal: **generate collider physics events at reconstruction level**, using the ATLAS Open Data datasets as training examples. **Motivation:** can a single generative model learn to predict SM histograms?

Our first approach was: the model sees the kinematics distributions from ATLAS data $z = (p_T, \eta, \phi, E) \in \mathbb{R}^4$, and learns the velocity field that transports noise to ATLAS data.

		p_T	η	ϕ	E
Event 1	μ^-	45.16	0.71	1.79	57.20
	μ^+	35.80	2.41	3.10	201.55
	j	54.64	-0.05	-0.76	55.08
Event 2	μ^-	84.77	0.30	0.00	88.62
	μ^+	51.58	0.32	-1.44	54.24
	j	92.70	-2.37	2.04	502.28
	j	46.80	-1.09	-2.63	78.03
Event 3	μ^-	32.41	0.03	1.54	32.42
	μ^+	68.60	0.21	-0.97	70.13
	γ	45.31	-0.11	-3.09	45.57

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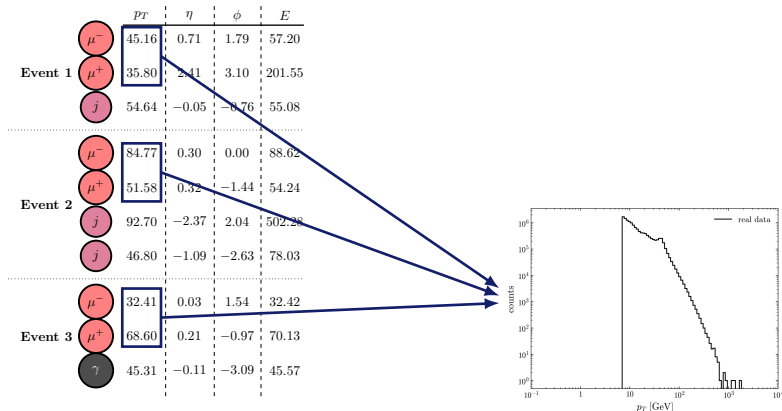
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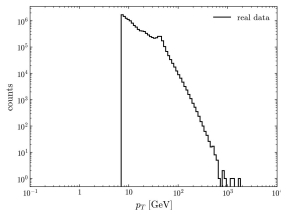
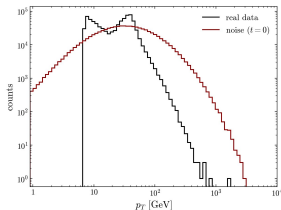


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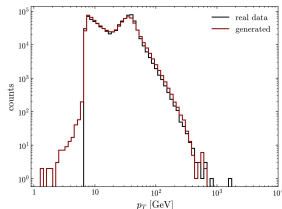
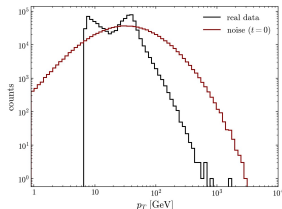


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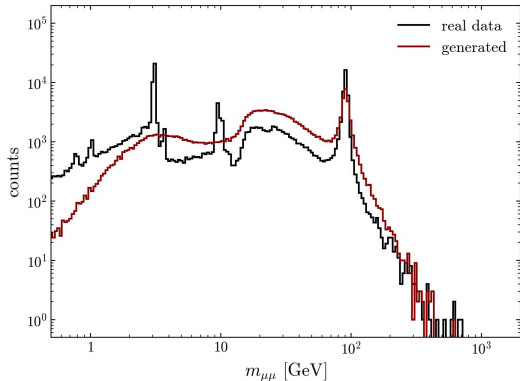


Generative models for collider physics

Consider now the invariant mass spectrum of opposite sign-same flavor muons:

$$m_{\mu\mu}^2 = 2 m_\mu^2 + 2(E_1 E_2 - \vec{p}_1 \cdot \vec{p}_2),$$

encodes known physics: peaks at $m_{\mu\mu} = m_Z, m_{J/\Psi}, \dots$. The same model failed to replicate the whole spectrum



Riemannian Flow Matching

We want to change our approach so that it keeps all of its particles on-shell:

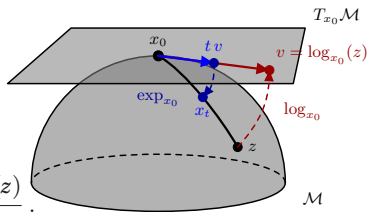
$$E^2 = |\vec{p}|^2 + m^2,$$

all particles must follow this relationship. We can enforce it by using **Riemannian Flow Matching**.

Given a Riemannian manifold $\mathcal{M}$ with metric g , where we want the noise, data, and the path in-between to be in. In this case, the velocity vector, $u_t(x) \in T\mathcal{M}$ in the tangent space of the manifold $\mathcal{M}$.

The path from noise to data now must follow **geodesic paths**. The geodesics connecting ϵ and z , and the velocity vector that transports between both; can be expressed in terms of the exp and log mappings

$$x_t = \exp_{\epsilon}(t \log_{\epsilon}(z)) , \quad u_t(x | z) = \frac{\log_{x_t}(z)}{1 - t} .$$



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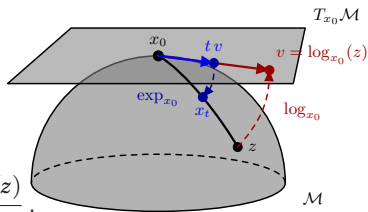
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The path from noise to data now must follow **geodesic paths**. The geodesics connecting ϵ and z and the velocity

Contrast with vanilla CFM:

$$x_t = (1 - t)\epsilon + tz, u_t = z - \epsilon$$

between terms of the exp and log mappings



$$x_t = \exp_{\epsilon}(t \log_{\epsilon}(z)) , \quad u_t(x | z) = \frac{\log_{x_t}(z)}{1 - t} .$$

Manifold choice

The data we have is written in terms of (p_T, η, ϕ, E, m) . Converting it to Euclidean momentum:

$$z = (E, \vec{p}, m),$$

This is then encoded to two types of manifolds:

Massless manifold (e, μ, γ)

$$\mathcal{M}_{\text{massless}} = \mathbb{R} \times S^2,$$
$$z = (y_E, \hat{n}),$$

Massive manifold (jets, τ)

$$\mathcal{M}_{\text{massive}} = \mathbb{R} \times S^2 \times \mathbb{R},$$
$$z = (y_E, \hat{n}, z_m),$$

Chart	Encode	Decode
Direction (S^2)	$\hat{n} = \mathbf{p}/ \mathbf{p} $	$\mathbf{p} = \mathbf{p} \hat{n}$
Log-energy	$y_E = \log(E/E_{\text{ref}}^k)$	$E = E_{\text{ref}}^k \exp(y_E)$
Boost-rapidity χ	$\chi = \text{arccosh}(E/m)$	$E = m \cosh \chi$
On-shell mass*	$m = \sqrt{E^2 - \mathbf{p} ^2}$, floored	$ \mathbf{p} = \sqrt{E^2 - m^2}$
Log-mass	$z_m = \log(m/m_{\text{floor}}^k)$	$m = m_{\text{floor}}^k \exp(z_m)$

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Explicit form for $\exp_z$ and $\log_z$ maps for S^2 are shown in backup slides

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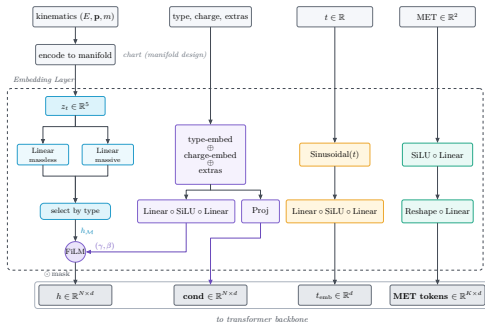
Model architecture: embedding

Input:

- Kinematics to generate: $(E, \vec{p}, m)$
- Conditioned variables: particle-type, MET, t , detector-features

Output:

- Kinematics embedded in higher dimension: h
- Higher dimensional conditions variables: t_{emb} , cond , MET tokens



Model architecture: backbone model

Input: output of the embedding layer

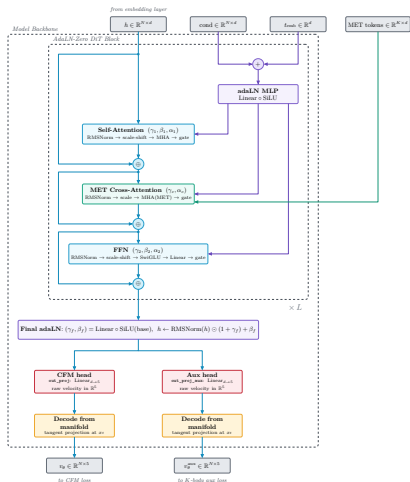
Backbone layer: six transformer adaLN-Zero layers:

- attention
- MET-cross attention
- feed-forward network

each one wrapped in a residual connection (ResNet-style)

Output: two velocity field:
 $u(t), u(t)^{\text{aux}} \in \mathbb{R}^5$

Two losses: main CFM loss, and an auxiliary loss that focuses on K -invariant masses



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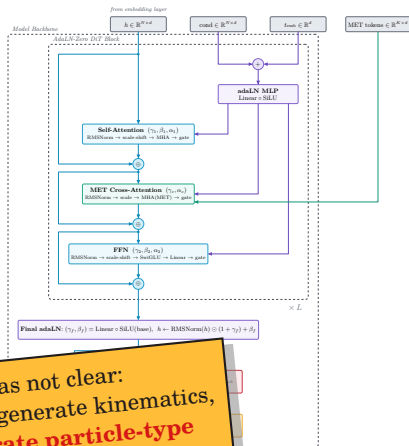
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In case it was not clear:
the model will only generate kinematics,
**it will not generate particle-type
or particle multiplicities**



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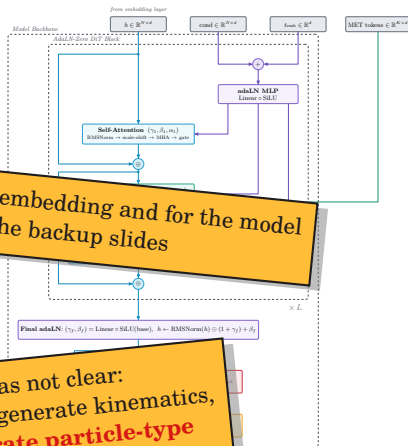
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Relevant functions for both embedding and for the model are defined in the backup slides

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Loss designs

CFM loss: modification from original CFM loss

$$\mathcal{L}_{\text{CFM}} = \mathbb{E} \left[q_i^{\text{eff}} \cdot H_{\delta} (u_i(x_t, t) - u_i(x_t, t \mid z)) \right]$$

H_{δ} : Huber loss

q_i^{eff} : quality-weighting

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Defined in backup slides

Auxiliary loss: ensures that the model becomes aware of the invariant mass

$$\mathcal{L}_{\text{aux}} = \sum_K \left\langle q_{\mathcal{I}} w_I^* \left[\mathbb{I}(t > t_{\text{gate}}) H_\delta(\Delta_{\mathcal{I}}^{\text{pos}}) + H_\delta(\Delta_I^{\text{vel}}) \right] \right\rangle ,$$

with

$$\Delta_I^{\text{pos}} = \log m_K^{\text{pred}}(\mathcal{I}) - \log m_K^{\text{truth}}(\mathcal{I}) ,$$

$$\Delta_I^{\text{vel}} = \log \dot{m}_K^{\text{pred}}(\mathcal{I}) - \log \dot{m}_K^{\text{truth}}(\mathcal{I}) ,$$

where $m_K^2 = (E_1 + \dots + E_K)^2 - (\vec{p}_1 + \dots + \vec{p}_K)^2$. Predicted kinematics are computed from *one Euler step*. w^* is as KDE distribution-matching weight (defined in backup slides)

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Teaches the model where the invariant mass should arrive

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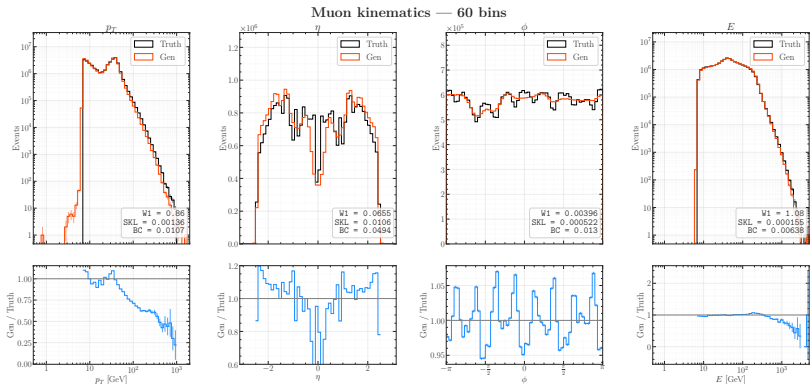
Teaches the direction of travel of the invariant mass

Results



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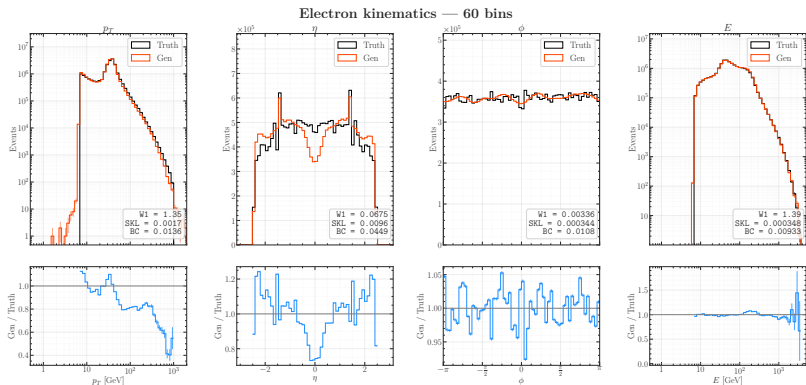
Single-particle kinematics (for example muon)



We can immediately notice

- $p_T < 7$ GeV cut not learned (hard to learn strict cuts like this with any generative model)
- $\eta = 0$ deficit: collider artifact (for some reason also replicated for other particles despite not present in data)
- ϕ relative uniformity is learned

Single-particle kinematics (for example electron)

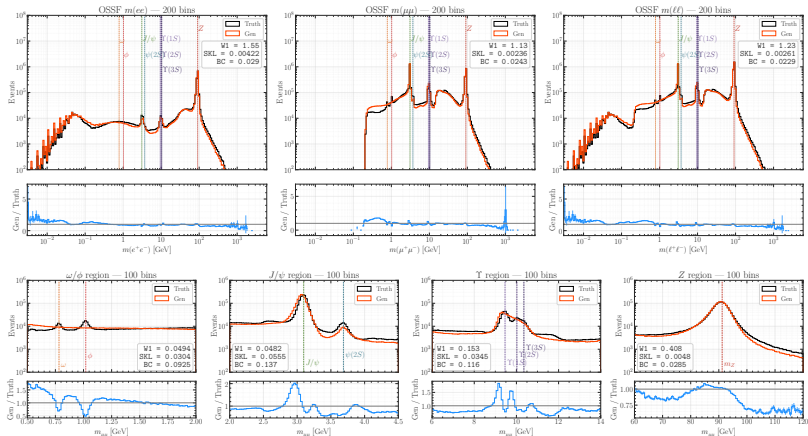


We can immediately notice

- $p_T < 7$ GeV cut not learned (hard to learn strict cuts like this with any generative model)
- $\eta = 0$ deficit, unknown reason
- ϕ relative uniformity is learned

Dilepton spectrum

Dilepton invariant masses and dimuon resonance zooms



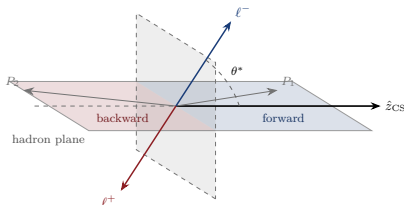
Plot of $m_{\ell\ell}^2 = 2(E_1 E_2 - \vec{p}_1 \cdot \vec{p}_2)$.

The model could reproduce very accurately several peaks in the invariant mass spectrum: $Z, J/\psi, \Psi(2S)$. It merged the Υ resonances.

Measurement of the Weinberg Angle

Forward-backwards asymmetry:

$$A_{\text{FB}} = \frac{N_F - N_B}{N_F + N_B}$$



N_F : number of ℓ^- in the **forward** direction

N_B : number of ℓ^- in the **backward** direction

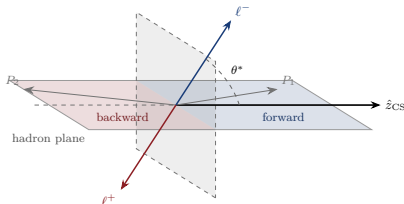
Measurement of the Weinberg Angle

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N_B : number of ℓ^- in the **backward direction**



Theory predicts A_{FB} depends on the $\sin \theta_W$, the sine of the Weinberg Angle.

$$A_{FB}^q = \frac{3(1 - 4 \sin^2 \theta_W)(1 - 4|Q_q| \sin^2 \theta_W)}{[1 + (1 - 4 \sin^2 \theta_W)][1 + (1 - 4|Q_q| \sin^2 \theta_W)^2]},$$

where Q_q is the charge of the initial quark.

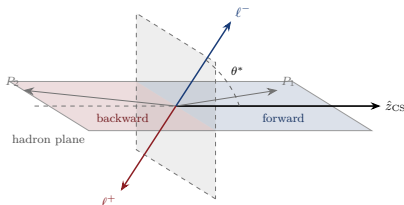
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where Q_q is the charge of the initial quark.

The model was never encouraged to reproduce the forward-backward asymmetry spectrum. If it could reproduce it, it is a small sign it learned the physics.

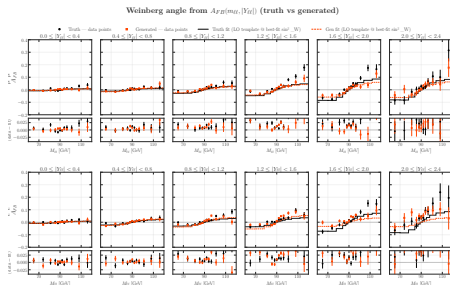
Measurement of the Weinberg Angle

Forward-backward asymmetry histograms at different total rapidities, $Y_{\ell\ell}$ and invariant masses $m_{\ell\ell}$:

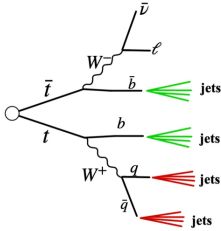
We do a fit using the leading-order template fits of the A_{FB} . We find from the fits that $\sin^2 \theta_W$ are

Source	$\sin^2 \theta_W$	σ_{fit}	σ_{gen}
$\mu^+\mu^-$ truth	0.23690	± 0.00084	—
$\mu^+\mu^-$ generated	0.23657	± 0.00082	± 0.00055
e^+e^- truth	0.23792	± 0.00103	—
e^+e^- generated	0.23627	± 0.00092	± 0.00073
PDG eff. leptonic	0.23122	—	—
CMS Run-1 NLO	0.23140	—	—

The “generation uncertainty”, σ_{gen} was computed using five different seeds during generation.

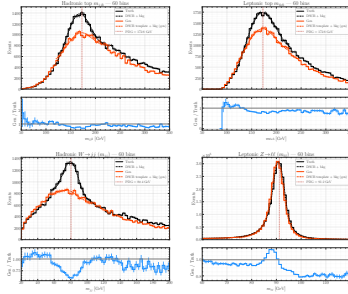


Fitting heavy resonances



Semileptonic $t\bar{t}$ process:
 $t\bar{t} \rightarrow WWbb \rightarrow jj\ell\nu bb$.

Heavy-particle mass closure: top, W, and Z



Observable	Truth μ	Gen μ	PDG
Hadronic top m_{jjb}	165.1 ± 1.2	$167.5 \pm 2.3 \pm 3.6$	172.76
Leptonic top $m_{\ell\nu b}$	166.9 ± 1.3	$166.8 \pm 1.6 \pm 1.9$	172.76
Hadronic W m_{jj}	80.0 ± 0.4	$77.3 \pm 1.8 \pm 3.5$	80.38
Leptonic Z $m_{\ell\ell}$	90.9 ± 0.1	$90.5 \pm 0.1 \pm 0.05$	91.19

Conclusions



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Conclusions

What we did:

- We built ShellFlow, a generative model whose only physical priors are the **on-shell condition** of particles, and knowledge of **invariant masses**.
- ShellFlow managed to capture a substantial part of the Standard Model structure that can be seen in colliders. This includes non-trivial **multi-particle correlations**, positions of heavy-resonances in different multi-particle spectrum of particles, and a reconstruction of the **Weinberg angle** that agrees with PDG values.

Limitations:

Not all the particle correlations are perfect, known issues include: deficit at $\eta = 0$ for all particles, no low-mass resonances (ω, ϕ), no peaks from 4-lepton invariant mass (Z, H), and W, t peaks are not sharp

What is next:

Train with MC data, mechanistic interpretability analysis (can we find *physical* features inside ShellFlow?), changing ShellFlow to also predict particle multiplicity or PID,...

Acknowledgments

I would like to thank our co-authors, since none of this work would have been possible without either of them

- Midori Kato
- Inar Timiryasov
- Oleg Ruchayskiy

We thank **NVIDIA** for providing us access to 8 A100s used during the writing of this project.

We acknowledge the work of the **ATLAS Collaboration** to record or simulate, reconstruct, and distribute the Open Data used in this paper, and to develop and support the software with which it was analysed.

Backup slides



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Collider coordinate system

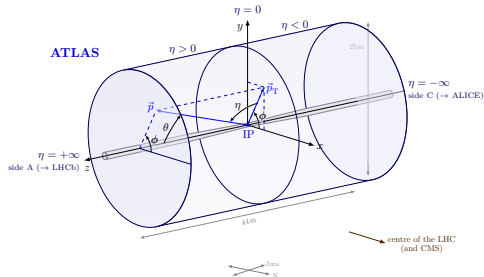
Detector geometry is cylindrically symmetric about the beam, so we work in $(\theta, \phi, p_T, \eta)$ derived from the 4-momentum (p_x, p_y, p_z, E) :

$$\theta = \arctan \frac{\sqrt{p_x^2 + p_y^2}}{p_z}$$

$$\phi = \arctan \frac{p_y}{p_x}$$

$$p_T = \sqrt{p_x^2 + p_y^2}$$

$$\eta = -\ln \tan \frac{\theta}{2} = \frac{1}{2} \ln \frac{|\mathbf{p}| + p_z}{|\mathbf{p}| - p_z}$$



Geodesics of different manifolds

Spheres

For $x, y \in S^{d-1}$, the exp and log maps have the form:

$$\log_x(y) = \frac{\theta}{\sin \theta} (y + \langle x, y \rangle x), \quad \exp_x(v) = \cos \|v\| x + \frac{\sin \|v\|}{\|v\|} v,$$

where $\theta = \arccos(\langle x, y \rangle)$ is the angle between x and y , $\langle x, y \rangle = \sum_i x_i y_i$ is the usual dot product, and $\|v\| = \langle v, v \rangle$.

Hyperboloids

For $x, y \in \mathbb{H}^d$, the exp and log maps have the form:

$$\log_x(y) = \frac{\chi}{\sinh \chi} (y - \langle x, y \rangle_{\mathbb{M}} x), \quad \exp_x(v) = \cosh \|v\|_{\mathbb{M}} x + \frac{\sinh \|v\|_{\mathbb{M}}}{\|v\|_{\mathbb{M}}} v,$$

where $\theta = \operatorname{arccosh}(\langle x, y \rangle_{\mathbb{M}})$ is the hyperbolic angle between x and y , $\langle x, y \rangle_{\mathbb{M}} = \sum_{i,j} g_{ij} x_i y_j = x_0 y_0 - \sum_{i>0} x_i y_i$ is the Minkowski product, and $\|v\|_{\mathbb{M}} = \langle v, v \rangle_{\mathbb{M}}$.

Here we can see why we chose not to embed the data into a hyperbolic manifold, cosh and sinh would present instabilities during training as they are not bounded functions, unlike cos or sin.

Definition of functions in model layers

Embedding layer functions

$$\begin{aligned}\text{SiLU}(x) &= x \sigma(x), & \text{FiLM}(h; \gamma, \beta) &= (1 + \gamma) \odot h + \beta, \\ \text{Sinusoidal}(t)_{2k} &= \sin(1000 t \omega_k), & \text{Sinusoidal}(t)_{2k+1} &= \cos(1000 t \omega_k), \\ \omega_k &= 10000^{-2k/d},\end{aligned}$$

where $\sigma(x)$ is a sigmoid function, and $\odot$ is the Hadamard matrix product.

Model functions

$$\text{FFN}(\text{SwiGLU}) = [\text{SiLU}(x W_1) \odot x W_3] W_2.$$

Training details

Hyperparameter	Value
<i>Architecture</i>	
DiT blocks L	6
model width d_{model}	128
attention heads	4
feed-forward expansion ratio	4
MET cross-attention tokens	4
dropout	0
trainable parameters	3.0×10^6
<i>Optimiser (AdamW)</i>	
peak learning rate	6×10^{-4}
weight decay	0.01
(β_1, β_2)	(0.9, 0.999)
ε	10^{-8}
gradient clip (global norm)	1.0
EMA decay ρ	0.9995
<i>Learning-rate schedule</i>	
type	cosine
warm-up	10^4 steps, linear
minimum learning rate	10^{-6}
epochs	30
total optimiser steps	717,180

Hyperparameter	Value
<i>Data and batching</i>	
batch size per GPU	4,096
global batch size	32,768
train/validation split	0.95 / 0.05
split seed	42
<i>Loss</i>	
Huber scale δ	1.0
K -body orders	$K \in \{2, 3, 4\}$, equal wt.
position gate t_{gate}	0.99
exponent α	1.0
sharpness k	5
truth-prior mix η	$1/3$ ($w_{\text{max}}^* = 3$)
quality floor $q_{\text{floor}}^{\text{cfm}}$	0.5
<i>Flow-time sampling</i>	
distribution	$t \sim \mathcal{U}[t_{\text{min}}, t_{\text{max}}]$
$(t_{\text{min}}, t_{\text{max}})$	(0.01, $1 - 10^{-6}$)
<i>Compute</i>	
GPUs	8, data-parallel
numerical precision	bf16-mixed
compilation	torch.compile

Hardware and software specifications

Component	Specification
GPU	8× NVIDIA A100 (80 GB), NVLink
CPU	2× AMD EPYC 7J13 (128 cores)
System memory	2 TB
Storage	19 TB NVMe RAID array
Software	PyTorch 2.10 (CUDA 12.8), Linux

Functions in loss functions

Huber loss:

$$H_\delta(x) = \begin{cases} \frac{1}{2}x^2, & |x| < \delta, \\ \delta \cdot (|x| - \frac{1}{2}\delta), & \text{otherwise,} \end{cases}$$

Quality weighting:

$$q_i^{\text{eff}} = q_{\text{floor}}^{\text{cfm}} + (1 - q_{\text{floor}}^{\text{cfm}}) q_i, \quad q_i^{\text{eff}} \in [q_{\text{floor}}^{\text{cfm}}, 1],$$

with

$$q_k = \begin{cases} \sigma(\alpha_{\text{iso}}^{\text{lep}}(c_{\text{iso}}^{\text{lep}} - r_{\text{iso}})) \sigma(\alpha_{\text{ip}}^{\text{lep}}(c_{\text{ip}}^{\text{lep}} - |d_0/\sigma_{d_0}|)) & k \in \{e, \mu\}, \\ \sigma(\alpha_{\text{iso}}^\gamma(c_{\text{iso}}^\gamma - r_{\text{iso}}^\gamma)) & k = \gamma, \\ \sigma(\alpha_{\text{jvt}}(\text{JVT} - c_{\text{jvt}})) & k = \text{jet}, \\ \sigma(\alpha_{\text{rnn}}(s_{\text{rnn}} - c_{\text{rnn}})) & k = \tau_{\text{had}}, \\ \sigma(\alpha_{D_2}(c_{D_2} - D_2)) & k = \text{large-}R \text{ jet}, \\ 1 & \text{otherwise (MET, padding).} \end{cases}$$

with $\sigma(x) = (1 + e^{-x})^{-1}$, $r_{\text{iso}} = \min(\text{ptvarcone30}, \text{topoetcone20})/p_T$ and $r_{\text{iso}}^\gamma = \min(\text{topoetcone40}, \text{ptcone20})/p_T$. The values of c and α are shown in the next slide.

More on quality features

Type	Feature	cut c	slope α
Leptons (e, μ)	$\min(\text{ptvarcone30}, \text{topoetcone20})/p_T$	0.15	10.0
	$ \text{d0sig} $	3.00	1.00
Photon	$\min(\text{topoetcone40}, \text{ptcone20})/p_T$	0.15	10.0
Jet	jvt	0.50	5.00
Tau	RNNJetScore	0.50	5.00
Large- R jet	D2	1.50	5.00

Functions in loss functions

KDE distribution-matching weight:

$$w^*(r) = \left[(1 - \eta) r^k + \eta \right]^{-\alpha}, \quad r = \frac{\hat{p}_{\text{gen}}}{\hat{p}_{\text{truth}}},$$

gaussian KDE of the K -body mass density,

$$\hat{p}(\log m) = \frac{1}{h} \sum \kappa \left(\frac{\log m - \log m_j}{h} \right), \quad \kappa(u) = \frac{1}{\sqrt{2\pi}} \exp(-u^2/2),$$

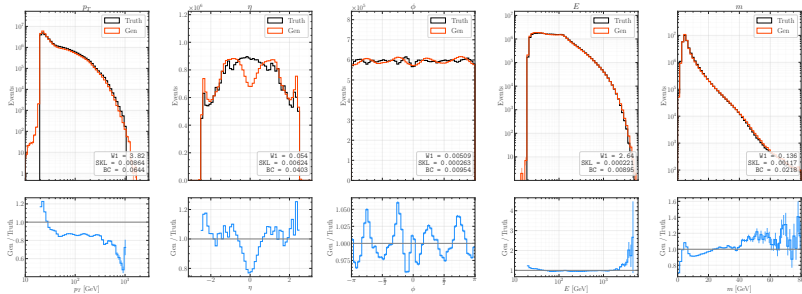
Evaluation metrics

Metrics evaluation distances between distributions x, y used in evaluating the quality of generation:

W_1	Wasserstein distance	$\hat{W}_1 = \frac{1}{n} \sum_{i=1}^N x_i - y_i $
SKL	Symmetric KL divergence	$SKL = \frac{1}{2} \sum_{i=1}^n x_i \log \frac{x_i}{y_i} + y_i \log \frac{y_i}{x_i}$
BC	Bray-Curtis Dissimilarity	$BC = \frac{\sum_{i=1}^n x_i - y_i }{\sum_{i=1}^n x_i + y_i}$

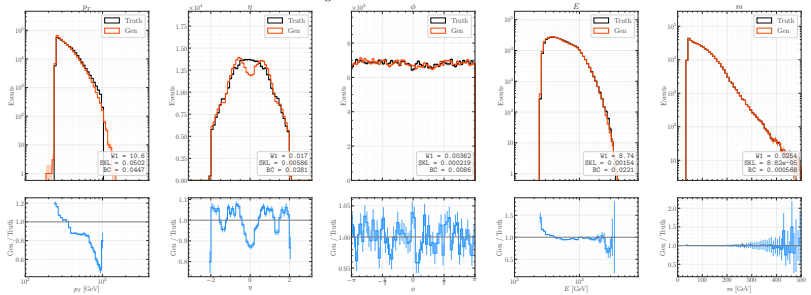
Jet kinematics

Jet kinematics — 60 bins



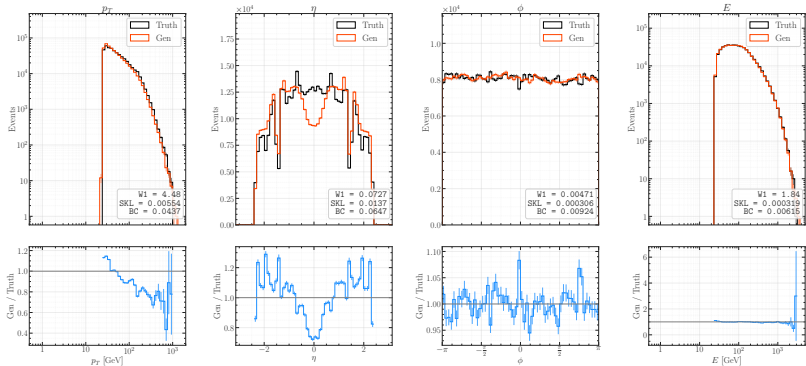
Large- R jet kinematics

Large- R Jet kinematics — 60 bins



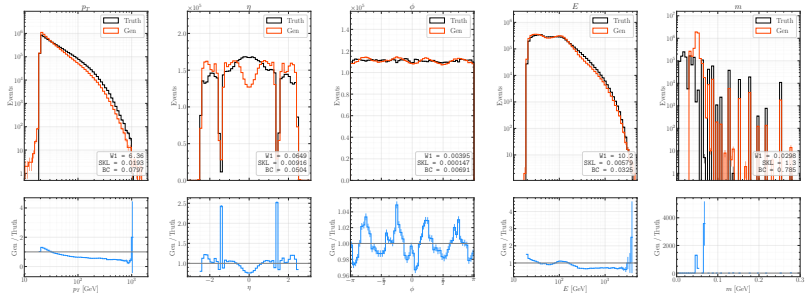
Photon kinematics

Photon kinematics — 60 bins



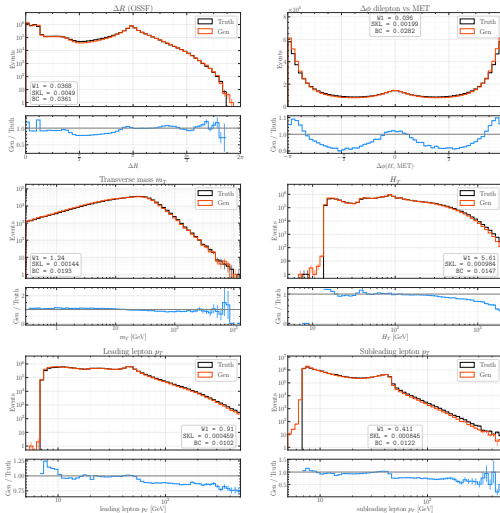
Tau kinematics

Tau kinematics — 60 bins



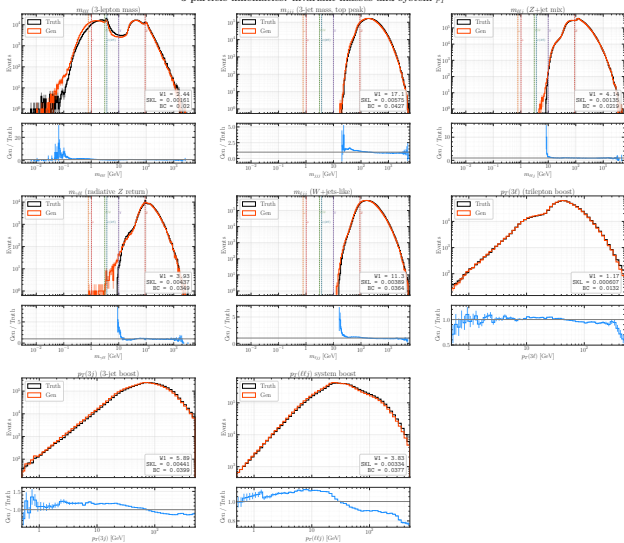
Additional evaluation metrics

Inter-particle kinematics: ΔR , $\Delta\phi$, m_T , H_T , lepton p_T — 60 bins



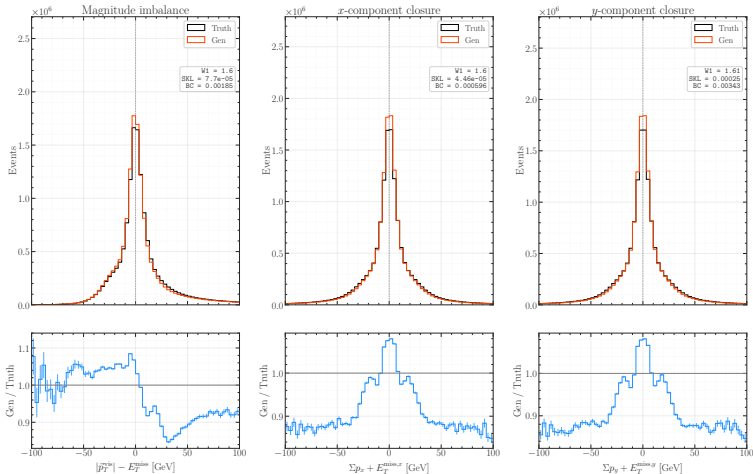
Three-particle features

3-particle kinematics: invariant masses and system p_T



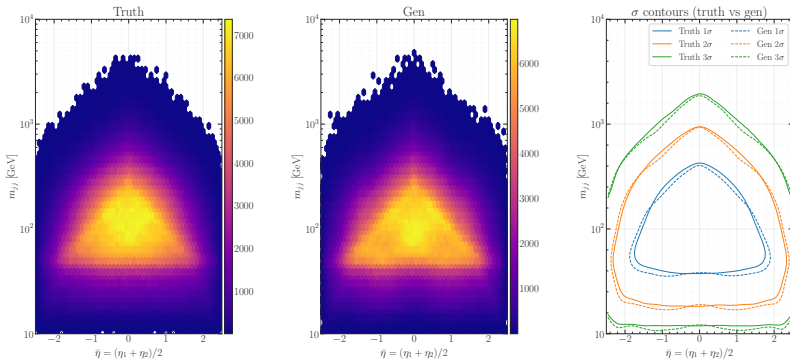
Transverse momentum conservation

Transverse momentum conservation — 60 bins



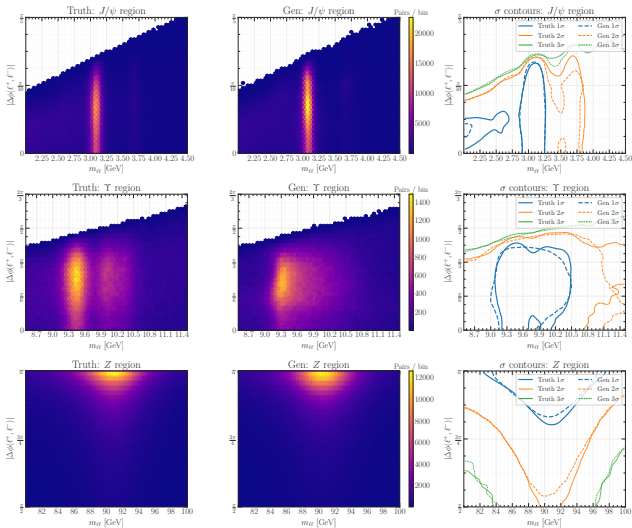
Leading di-jet kinematics

Leading dijet kinematics: m_{jj} vs $\bar{\eta}$ — 60 bins



OSSF dilepton pairs

OSSF dilepton: $m_{\ell\ell}$ vs $|\Delta\phi|$ (truth vs generated)



Four-lepton invariant mass

Higgs $m_{4\ell}$ in the $H \rightarrow ZZ^* \rightarrow 4\ell$ selection

