

#### LUND UNIVERSITY

# Jets in a quark-gluon plasma

#### Konrad Tywoniuk

Y. Mehtar-Tani, C.A. Salgado, KT PRL 106 (2011) 122002 Y. Mehtar-Tani, C.A. Salgado, KT PLB 707 (2011) 156 Y. Mehtar-Tani, KT arXiv:1105:1346 [hep-ph] Y. Mehtar-Tani, C.A. Salgado, KT arXiv:1112.5031 [hep-ph]

Spåtind 2012, 2-7 January 2012





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Large time domain for pQCD:  $\frac{1}{\sqrt{s}} < t < \frac{\sqrt{s}}{\Lambda_{\text{OCD}}^2}$ 

#### **QCD COHERENCE IN VACUUM**

[Dokshitzer, Fadin, Khoze, Troyan, Lipatov, Bassetto, Mueller, Ciafaloni, Marchesini....]



- leading singularities:

 $\propto \frac{d\omega_i}{\omega_i} \frac{d\theta_i}{\theta_i} \Theta \left(\theta_{i-1} - \theta_i\right)$ 

TASSO Collaboration, Z. Phys. C 47 (1990) 187 OPAL Collaboration, Phys. Lett. B 247 (1990) 617

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- included in most MC generators: PYTHIA, HERWIG
- soft & collinear divergences
- interferences ⇒ angular ordering

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CMS Experiment at LHC, CERN Data recorded: Mon Nov 8 11:30:53 2010 CEST Run/Event: 150431 / 630470 Lumi section: 173

100 N



#### MEDIUM MODIFIES THE JET EVOLUTION!





#### Fundamental building block of the QCD cascade!

$$\langle dN_q \rangle_{\varphi} = \frac{\alpha_s C_F}{\pi} \frac{d\omega}{\omega} \frac{d\theta}{\theta} \Theta(\cos\theta - \cos\theta_{q\bar{q}})$$

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Fundamental building block of the QCD cascade!

- Reason: emissions at large angles are sensitive to the total charge of the emitting system
- The antenna provides a nice laboratory!
- Question: how will the antenna radiation pattern look like if it were to traverse a quark-gluon medium?

#### ANTENNA SETUP IN MEDIUM

Mehtar-Tani, Salgado, KT PRL 106 (2011) 122002 Mehtar-Tani, Salgado, KT arXiv:1102.4317 [hep-ph]

 eikonal approximation for fixed opening angle of the pair
 medium is modeled as a classical

background field

$$J_q^{(0)}(x) = g\delta^{(3)}\left(\vec{x} - \frac{\vec{p}}{E}t\right)\Theta(t)Q_q$$



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g\*,γ\*

x+=0

 $x^+=L^+$ 

9

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 $\vartheta_{q\overline{q}}$ 

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Classical Yang-Mills eq:  $[D_{\mu}, F^{\mu\nu}] = J^{\nu}$ ,  $[D_{\mu}, J^{\mu}] = 0$ 

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Classical Yang-Mills eq:  $[D_{\mu}, F^{\mu\nu}] = J^{\nu}$ ,  $[D_{\mu}, J^{\mu}] = 0$ Linear response:  $\Box A^{i} - 2ig [A^{-}_{med}, \partial^{+}A^{i}] = -\frac{\partial^{i}}{\partial^{+}}J^{+} + J^{i}$ 

Gelis, Mehtar-Tani (2005), Mehtar-Tani (2007)

### **ANTENNA IN MEDIUM**

Y. Mehtar-Tani, KT arXiv: 1105.1346 [hep-ph], E. lancu, J. Casalderrey-Solana arXiv: 1105.1760 [hep-ph]

$$\begin{aligned} \text{Multiple scattering} &\Rightarrow \text{effective propagators:} \\ \mathcal{J} = \text{Re} \left\{ \int_{0}^{\infty} dy'^{+} \int_{0}^{y'^{+}} dy^{+} (1 - \Delta_{\text{med}}(y^{+}, 0)) \\ &\times \int d^{2} \boldsymbol{z} \exp \left[ -i \bar{\boldsymbol{\kappa}} \cdot \boldsymbol{z} - \frac{1}{2} \int_{y'^{+}}^{\infty} d\xi \, n(\xi) \sigma(\boldsymbol{z}) + i \frac{k^{+}}{2} \delta \boldsymbol{n}^{2} y^{+} \right] \quad |\delta n| \simeq \theta_{q\bar{q}} \\ &\times \left( \partial_{y} - i k^{+} \, \delta \boldsymbol{n} \right) \cdot \partial_{z} \, \mathcal{K}(y'^{+}, \boldsymbol{z}; y^{+}, \boldsymbol{y} \, |k^{+}) |_{\boldsymbol{y} = \delta \boldsymbol{n} y^{+}} \right\} + \text{sym.} \end{aligned}$$

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Describes Brownian motion through medium potential...

$$\mathcal{K}\left(y^{\prime+}, \boldsymbol{z}; y^{+}, \boldsymbol{y} | k^{+}\right) = \int \mathcal{D}[\boldsymbol{r}] \exp\left[\int_{y^{+}}^{y^{\prime+}} d\xi \left(i\frac{k^{+}}{2}\dot{\boldsymbol{r}}^{2}(\xi)\right) \cdot \frac{1}{2}n(\xi)\sigma(\boldsymbol{r})\right)\right]$$
$$\sigma(\mathbf{r}) = 2\alpha_{S}C_{A} \int \frac{d^{2}\mathbf{q}}{(2\pi)^{2}} \mathcal{V}^{2}(\mathbf{q}) \left[1 - \cos(\mathbf{r} \cdot \mathbf{q})\right]$$
$$\mathcal{V}(\boldsymbol{q}) = \frac{1}{\boldsymbol{q}^{2} + m_{D}^{2}}$$

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### THE SOFT LIMIT

Considering soft gluon emissions: only the quarks interact!

$$J_q(x) = g U_p(x^+, 0) \,\delta^{(3)}(\vec{x} - \frac{\vec{p}}{E}t)\Theta(t) \,Q_q$$

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Wilson line along the trajectory:

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 $\Delta_{\rm med} = 1 - \frac{1}{N_c^2 - 1} \langle \operatorname{Tr} U_p(x^+, 0) U_{\bar{p}}^{\dagger}(x^+, 0) \rangle$ 

$$\approx 1 - e^{-\frac{1}{12}\hat{q}\,\theta_{q\bar{q}}^2 L^3}$$

- the decoherence parameter!



**q**: medium transport coefficient



 $\Theta_{n=1}$ 



 $\Delta_{\rm med} \rightarrow 0$  **Coherence** 



 $\Theta_{qar{q}}$ 

 $\Delta_{\rm med} \rightarrow 0$  **Coherence** 

 $\Delta_{\text{med}} \rightarrow 1$  **Decoherence** 

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#### **ONSET OF DECOHERENCE** - FINITE ENERGIES



 $\Delta_{\rm med} \rightarrow 1$ 





 $\Delta_{\rm med} \rightarrow 1$ 

- independent spectrum appears
- interferences are destroyed
- jet broadening

Independent component: Baier, Dokshitzer, Mueller, Peigne, Schiff (1997-2001), Zakharov (1996), Wiedemann (2000), Gyulassy, Levai, Vitev (2001-2002)

 $L = 10 \text{ fm}, \hat{q} = 5 \text{ GeV}^2/\text{ fm}$ 



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# CONCLUSIONS

 $\star$  copious jets in heavy-ion collisions at the LHC \* medium induces soft radiation at large angles  $\Rightarrow$  onset of decoherence  $\star$  a two scale problem:  $r_{\perp}$  vs.  $Q_s$  $\Rightarrow$  jet probes medium, and vice versa \* the radiation pattern off an antenna  $\Rightarrow$  building block for jet calculus in medium

#### **MEDIUM-INDUCED RADIATION**

Baier, Dokshitzer, Mueller, Peigne, Schiff (1997-2001), Zakharov (1996), Wiedemann (2000), Gyulassy, Levai, Vitev (2001-2002)

- Energy loss:  $\Delta E \simeq \frac{\alpha_s C_R}{2\pi} \hat{q} L^2$ Broadening:  $k_{\perp}^2 \simeq \hat{q} L \propto \frac{\Delta E}{L}$
- emitted off a single emitter
- gluon interaction  $\Rightarrow$  k<sub>⊥</sub>-broadening
- transport parameter:  $\hat{q} = m_D^2 / \lambda$
- infrared & collinear safe spectrum
- energy loss distribution:  $P(\Delta E)$
- need more emitters to see coherence!





### ANGULAR SPECTRUM

![](_page_37_Figure_1.jpeg)

![](_page_37_Figure_2.jpeg)

 $1/\theta$ 

Y. Mehtar-Tani, C.A. Salgado, KT arXiv:1112.5031 [hep-ph]

![](_page_38_Figure_0.jpeg)

# **ENERGY SPECTRUM**

![](_page_39_Figure_1.jpeg)

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