

# Geometry, Geometry, Geometry

Jamie Nagle

University of Colorado, Boulder, USA

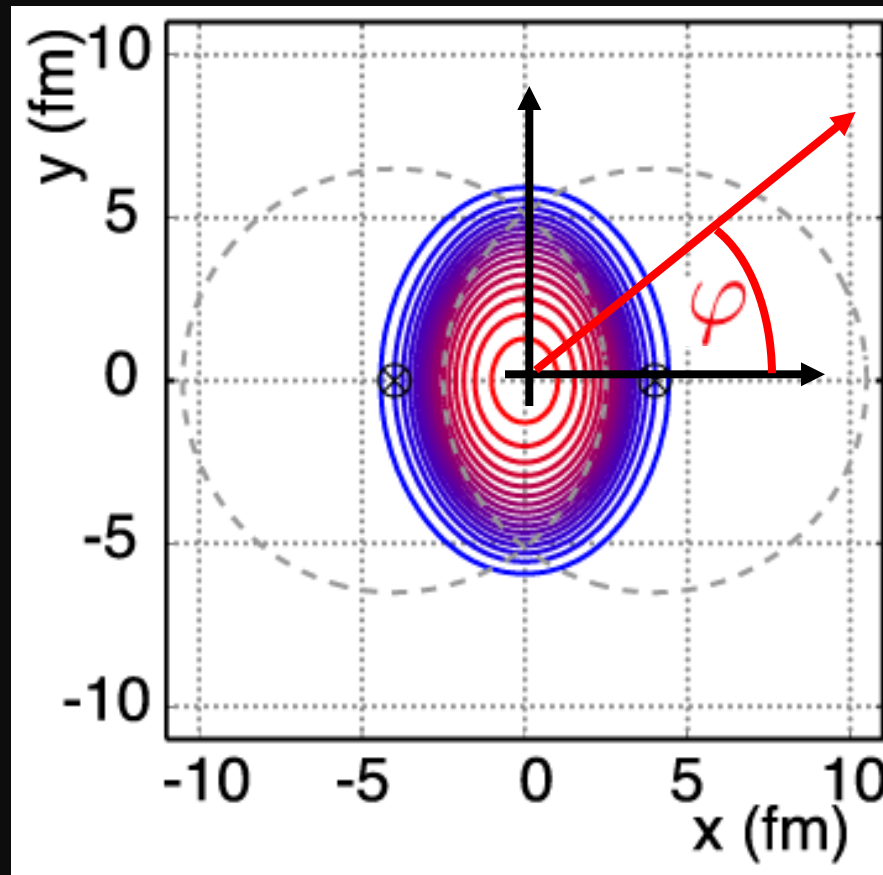


## Heavy ions: Experiments Confront Theory

7-9 November 2011 *Niels Bohr Institute*

Europe/Copenhagen timezone

# Simple Geometry

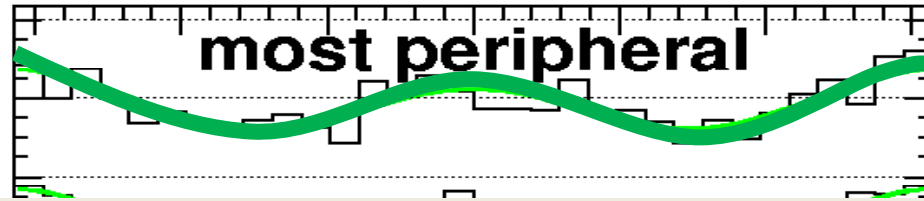


$$\frac{dN}{p_T dp_T dy d\varphi}(p_T, \varphi; b) = \frac{dN}{2\pi p_T dp_T dy} (1 + \underline{2v_2}(p_T; b) \cos(2\varphi) + \dots)$$

$v_2$  = “elliptic flow”

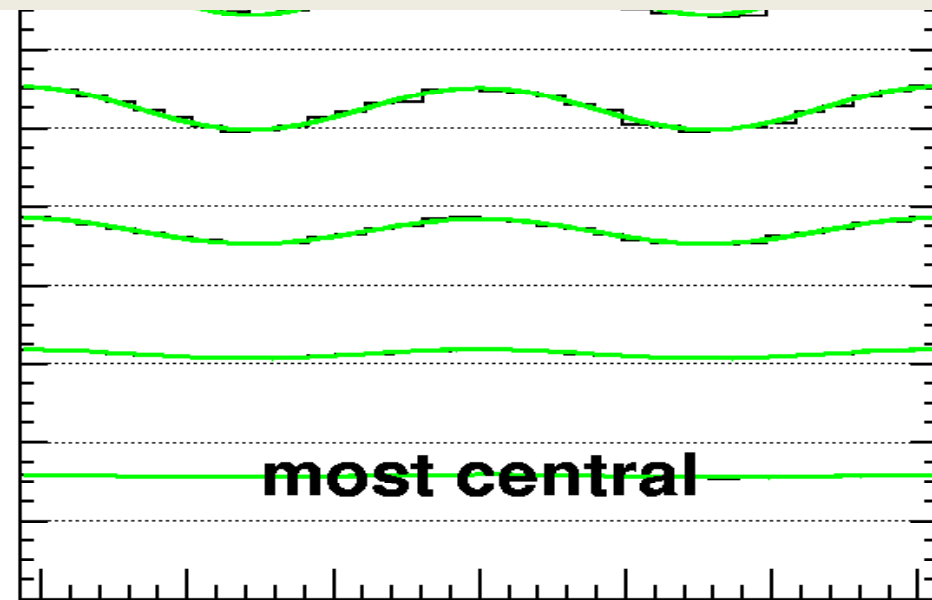
# Large Elliptic Flow

y scale



Literally “seen” in first days at RHIC

$dN/d(\phi - \Psi_2)$



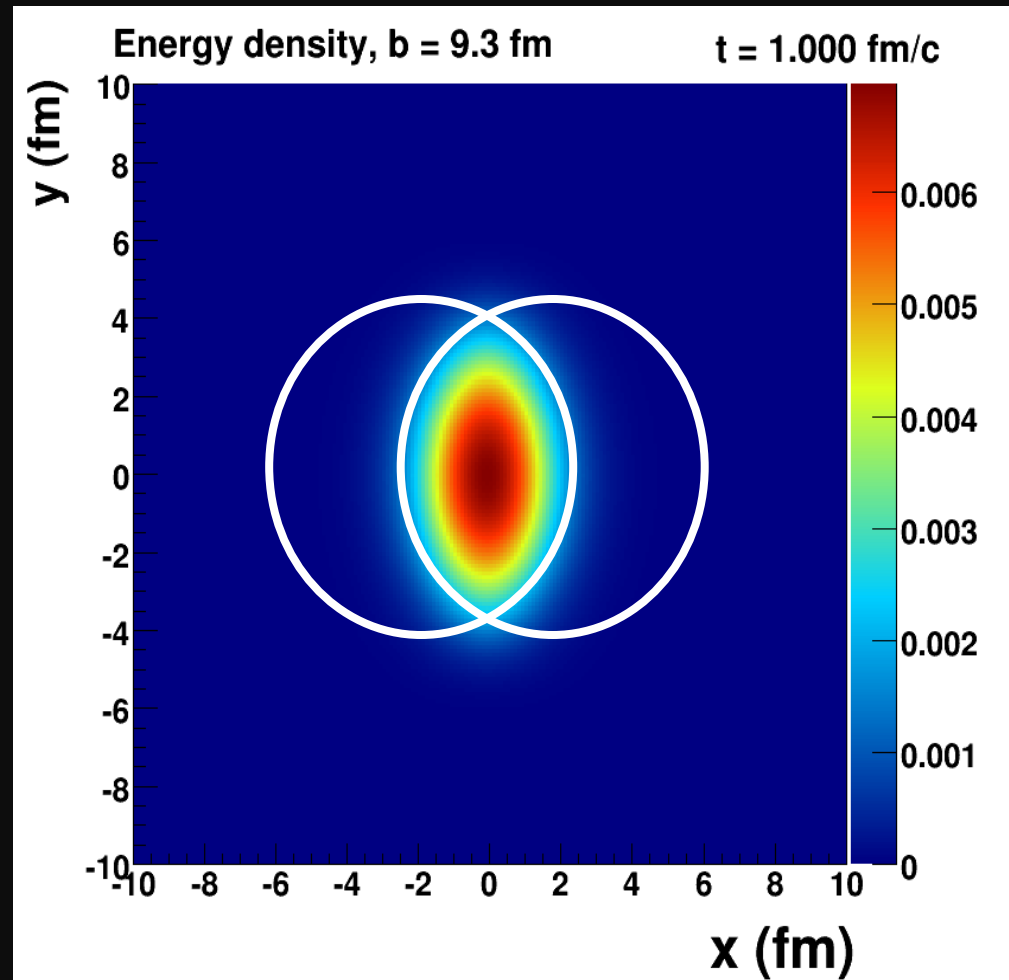
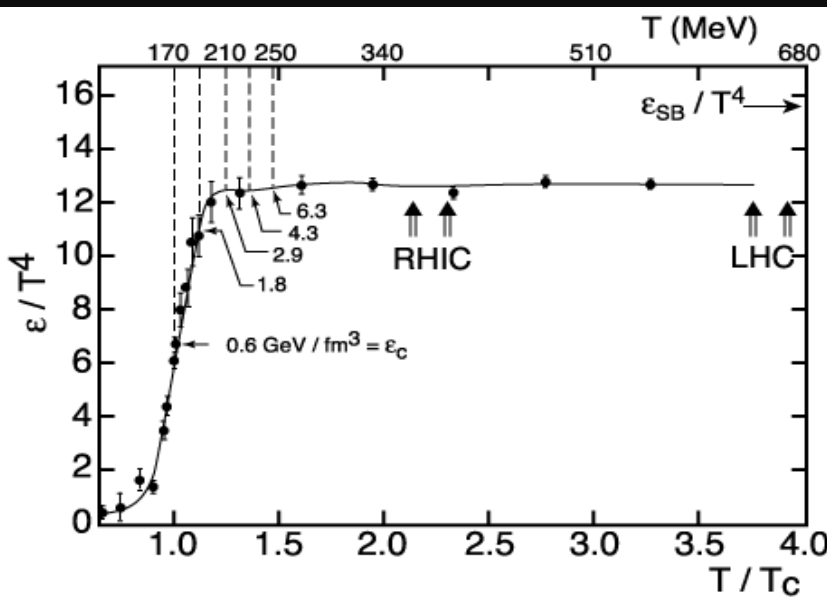
-3 -2 -1 0 1 2 3  
φ - Ψ₂ (radian)

# How Large is the Flow Really?

Assume early thermalization and run ideal hydrodynamics  
(i.e. no dissipation  $\rightarrow$  zero shear + zero bulk viscosity)

Key Inputs:

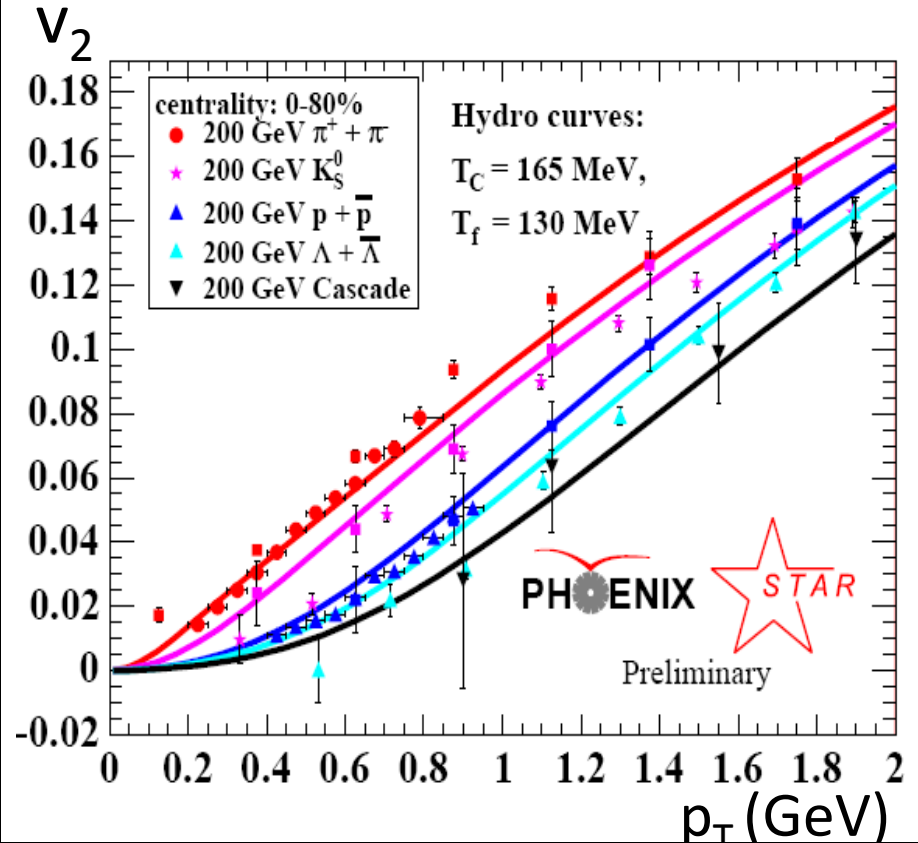
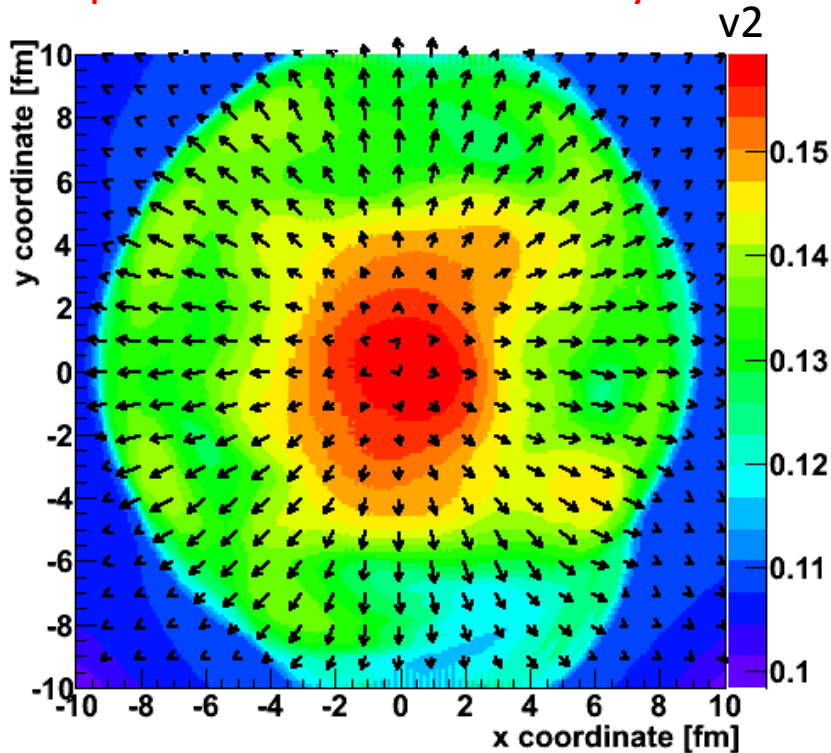
- *Initial Geometry*
- $\ell$ QCD Equation of State



# Fluid cells “freeze-out” below $T_f$

## Isotropic hadrons in cell rest frame, then boosted

Temperature Profile + Velocity Vectors

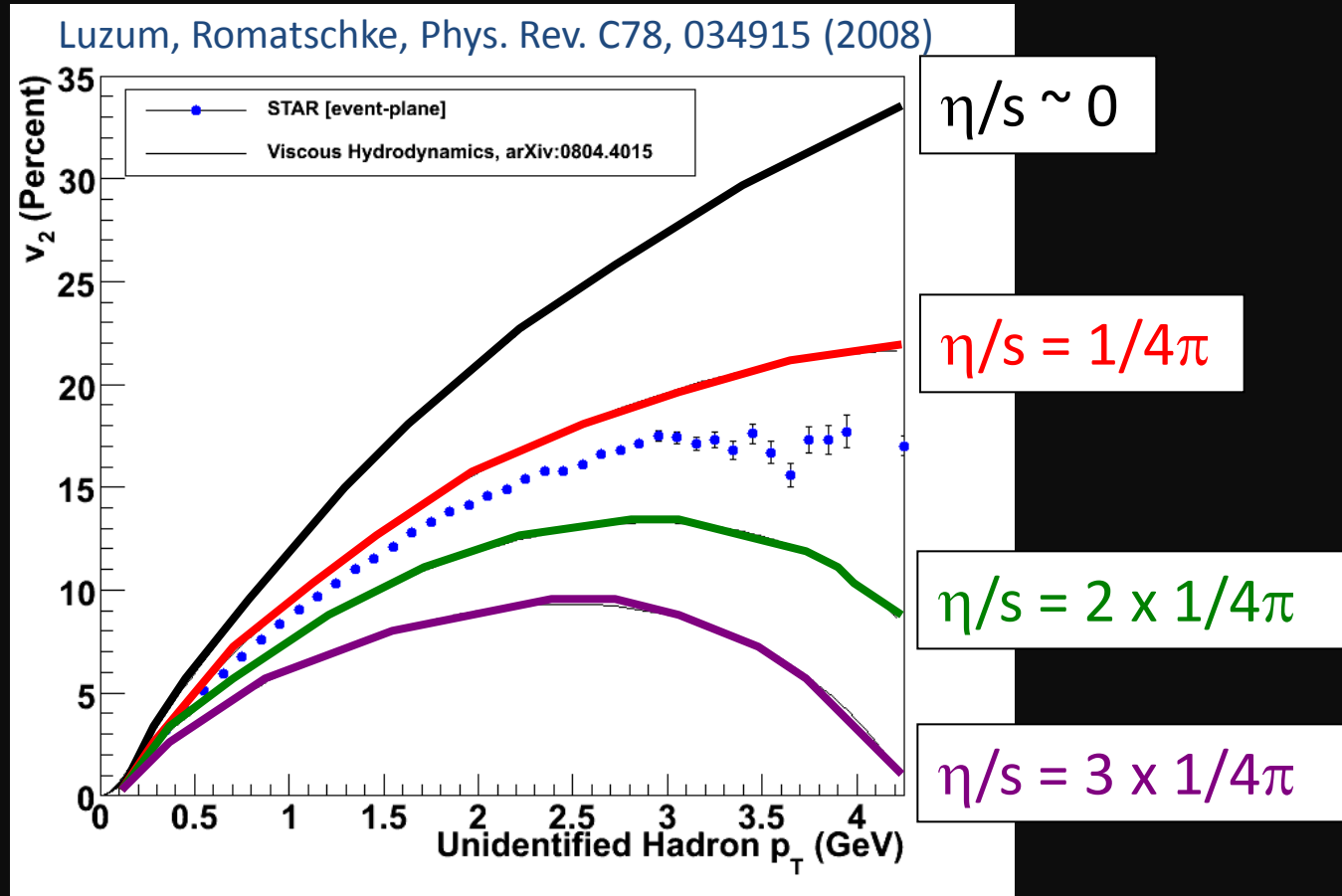


Characteristic flow pattern observed for  $p_T < 2 \text{ GeV}$

*Circa 2005: Quark-Gluon Plasma = Perfect Fluid*

# How to Quantify QGP $\eta/s$ ?

## Relativistic viscous hydrodynamics compared to data



$$\left(\frac{\eta}{s}\right) / \left(\frac{1}{4\pi}\right) = 1.3 \pm 1.3 \text{ (theory)} \pm 1.0 \text{ (experiment)}$$

# Very close to the bound!

$$\left(\frac{\eta}{s}\right) / \left(\frac{1}{4\pi}\right) = \underline{1.3 \pm 1.3 \text{ (theory)}} \pm \underline{1.0 \text{ (experiment)}}$$

## What dominated these uncertainties a few years ago?

**Different ways of measuring  $v_2$  gave  $\pm 20\%$  variations.**

Now resolved that the methods have different sensitivities to “nonflow” and fluctuations.

For example Ollitrault, Poskanzer, Voloshin Phys. Rev. C80, 014904 (2009).

In ideal hydrodynamics  $v_2 \propto \varepsilon_2$  (initial eccentricity)

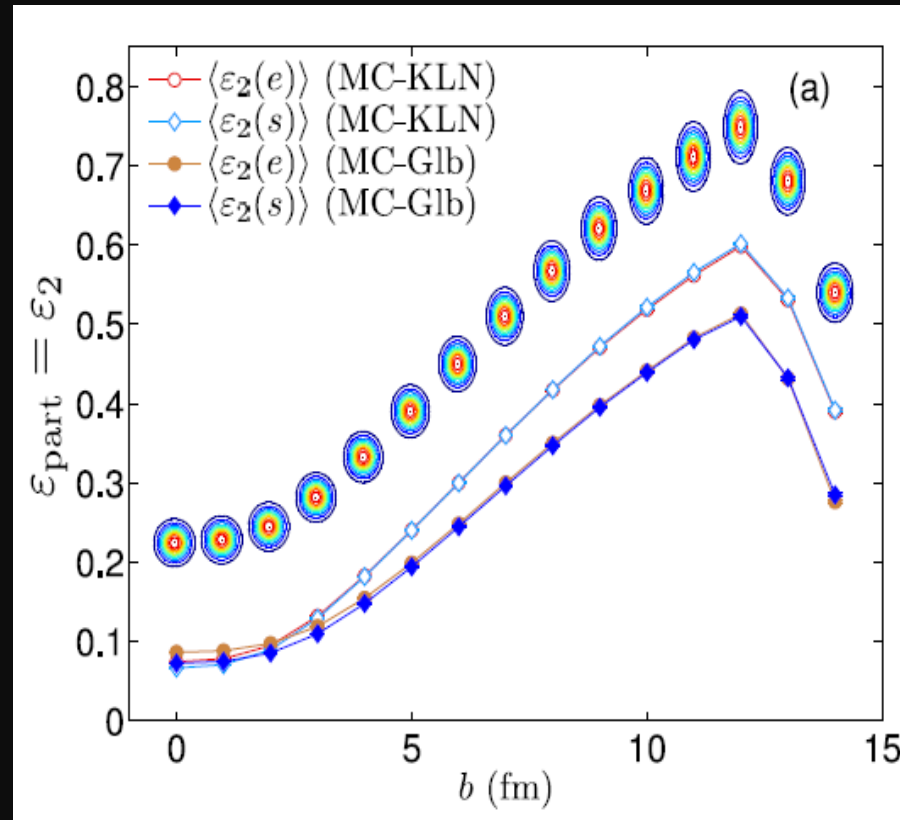
**Uncertainty on  $\eta/s$  from initial geometry was  $\sim 100\%$ .**

**Other sources subdominant:**

hadronization, EOS, pre-flow, ... (worth re-examination)



# Initial Condition Uncertainty



Elliptic Flow reasonably described by either:

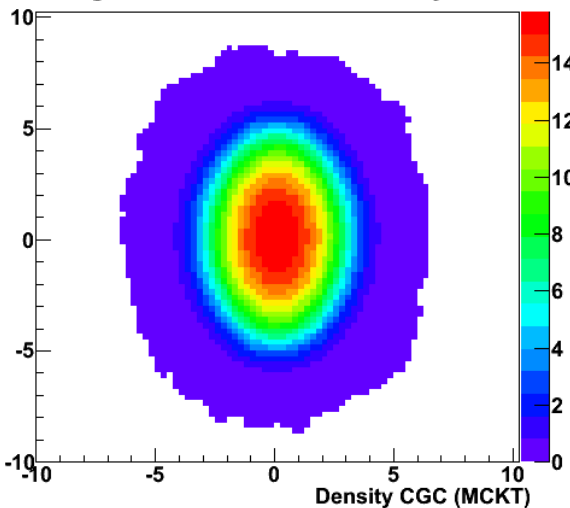
- A) Smaller eccentricity (Glauber) + Less Dissipation ( $\eta/s \approx 1/4\pi$ )
- B) Larger eccentricity (Gluon Saturation) + More Dissipation ( $\eta/s \approx 2/4\pi$ )

Neither **A** nor **B**, and yet these give a range of uncertainty

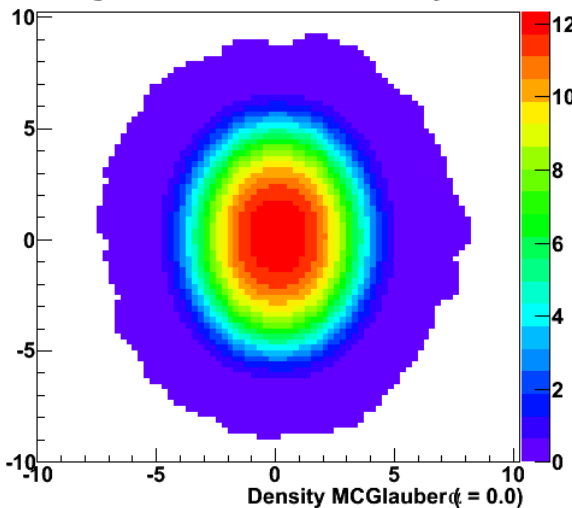


# Examine $b = 7$ fm case

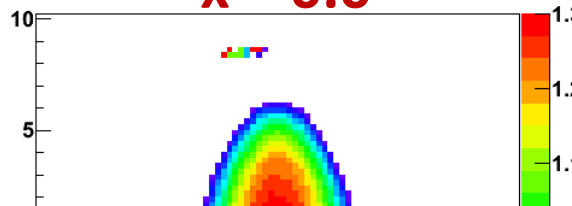
AuAu@200 GeV  $b=7.0$  Eccentricity = 0.377



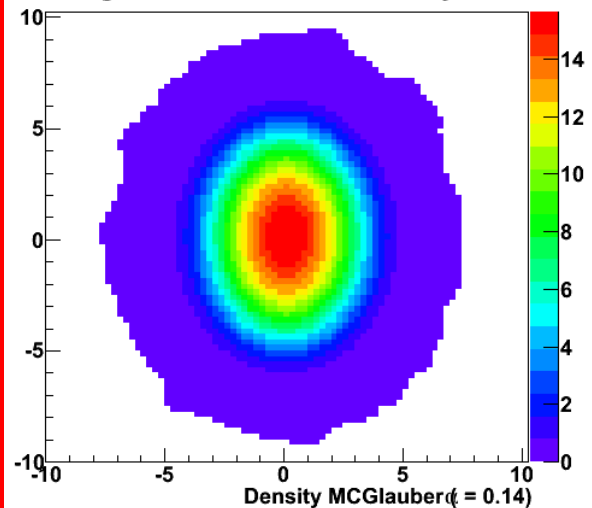
AuAu@200 GeV  $b=7.0$  Eccentricity = 0.278



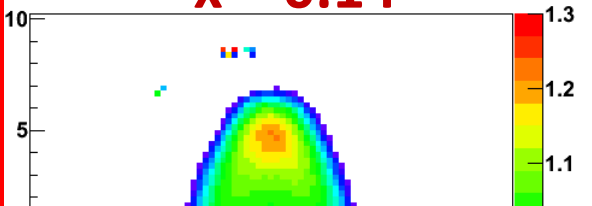
$x = 0.0$



AuAu@200 GeV  $b=7.0$  Eccentricity = 0.294



$x = 0.14$



Just saying “Glauber” is not enough (different variants).

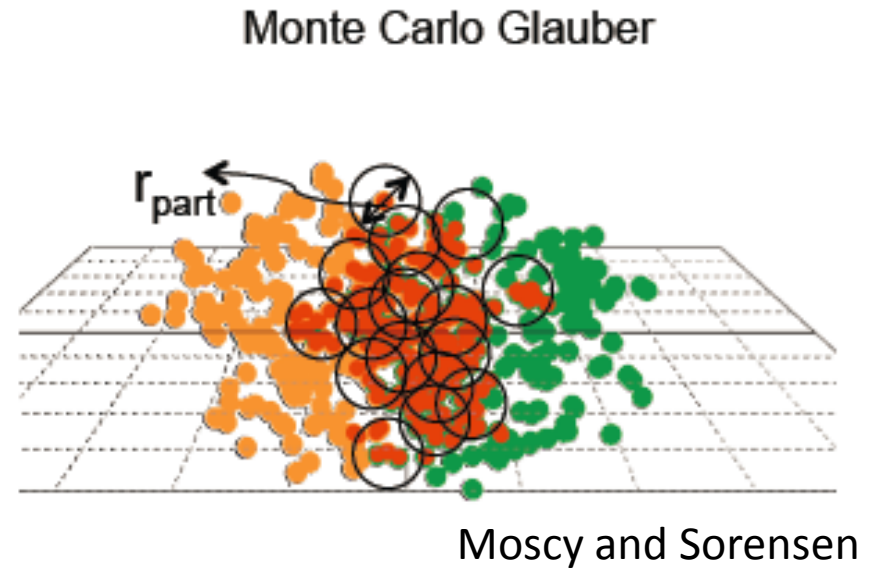
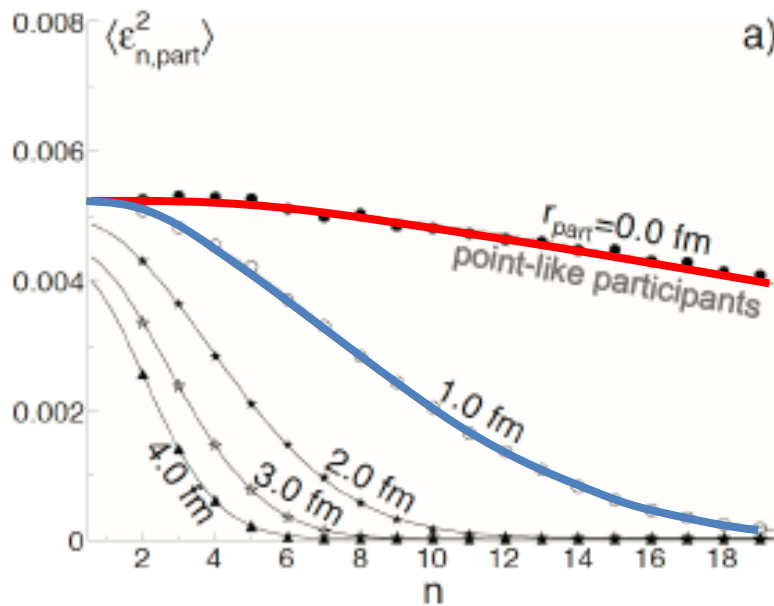
$$\frac{dN_{ch}}{d\eta} = n_{pp} \left[ (1 - x) \frac{N_{part}}{2} + x N_{coll} \right]$$

Ratio CGC / Glauber: 0.0

Ratio CGC / Glauber: 0.14

MC Glauber variations have 26 and 22% lower eccentricities than MCKT (with saturation effects).

# Spatial Moments Initial Smearing

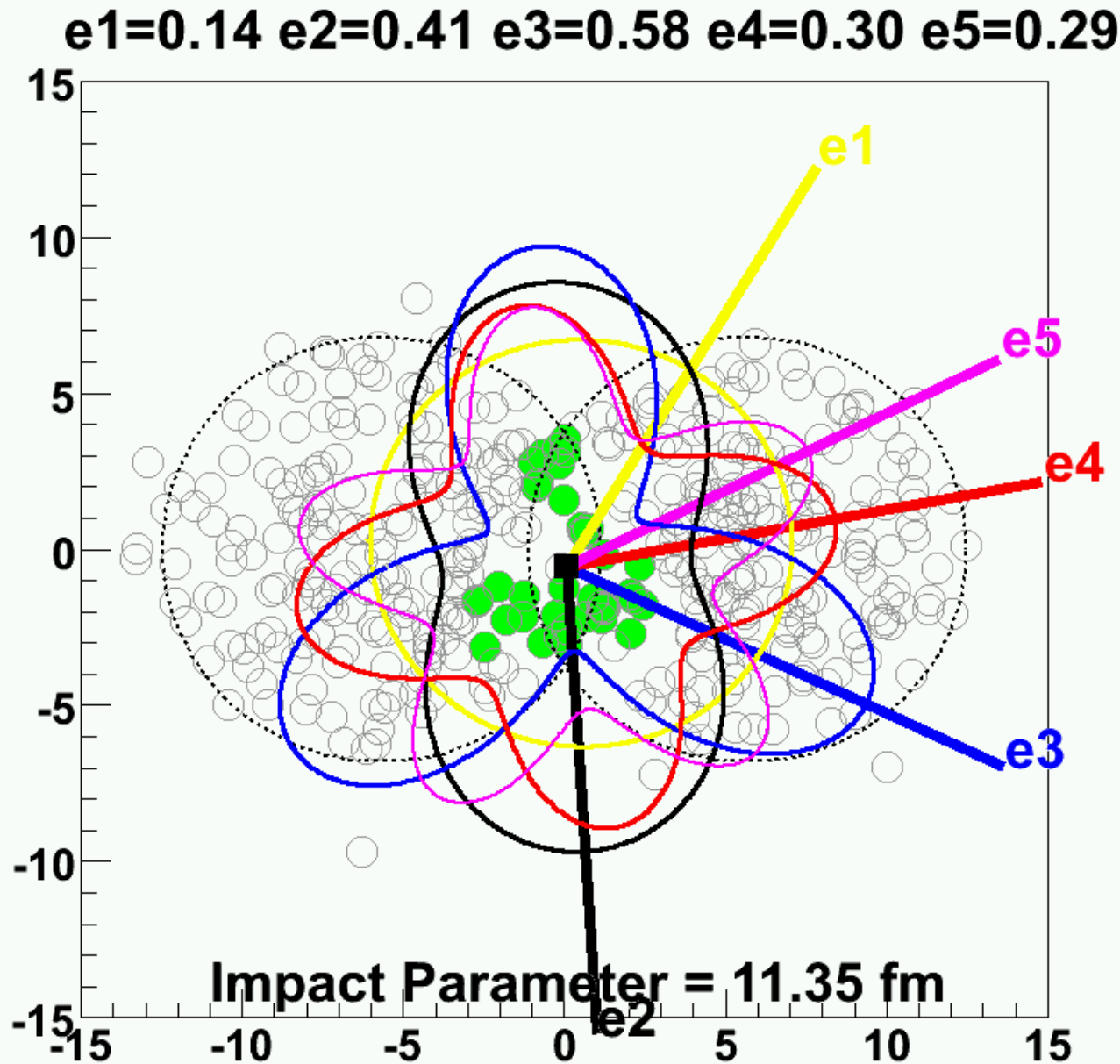


The initial smearing for the starting distribution suppresses the higher moments almost like a Gaussian drop off. Spatial smearing could be in initial state (e.g. size of nucleon) or during evolution.

For a 1 fm smearing this effect is modest for the  $n=2,3,4$  moments.

	<u>Au-Au 30-40% Central</u>		
	Point-like	Smear (1 fm radius)	
$\langle \epsilon_2 \rangle$	0.359	0.346	-4%
$\langle \epsilon_3 \rangle$	0.197	0.185	-6%
$\langle \epsilon_4 \rangle$	0.197	0.179	-9%

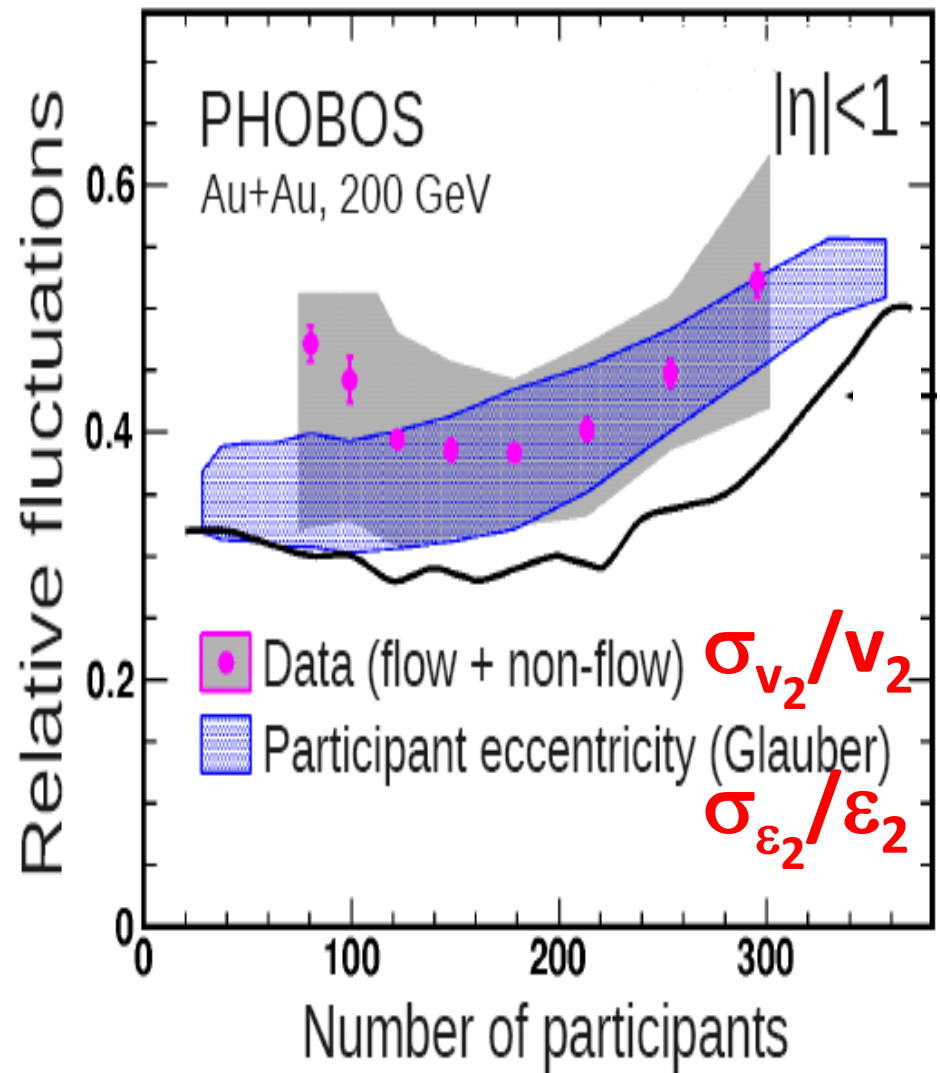
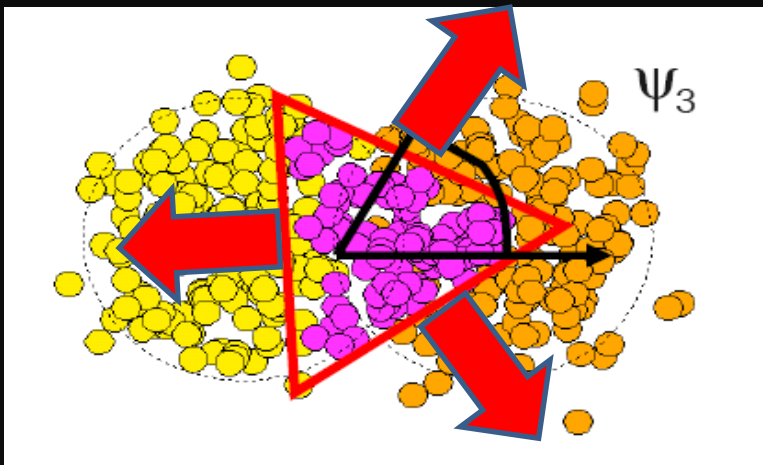
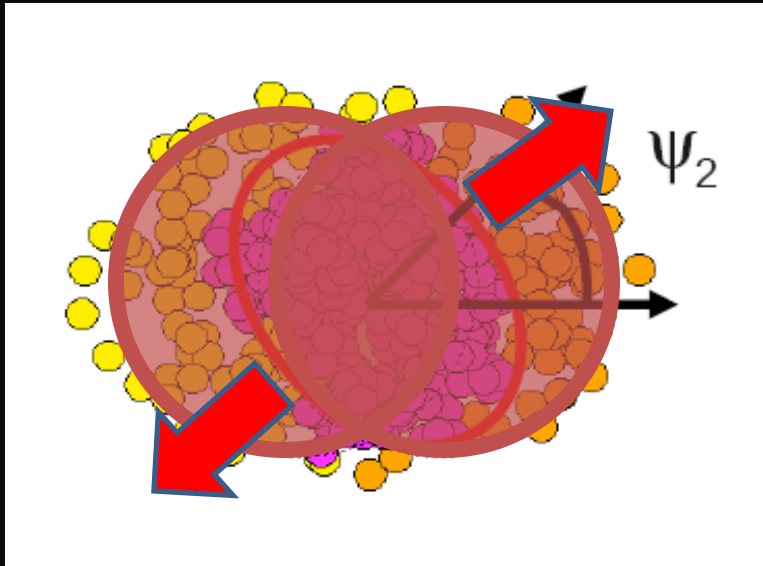
# Realization → Lumpy Initial Conditions



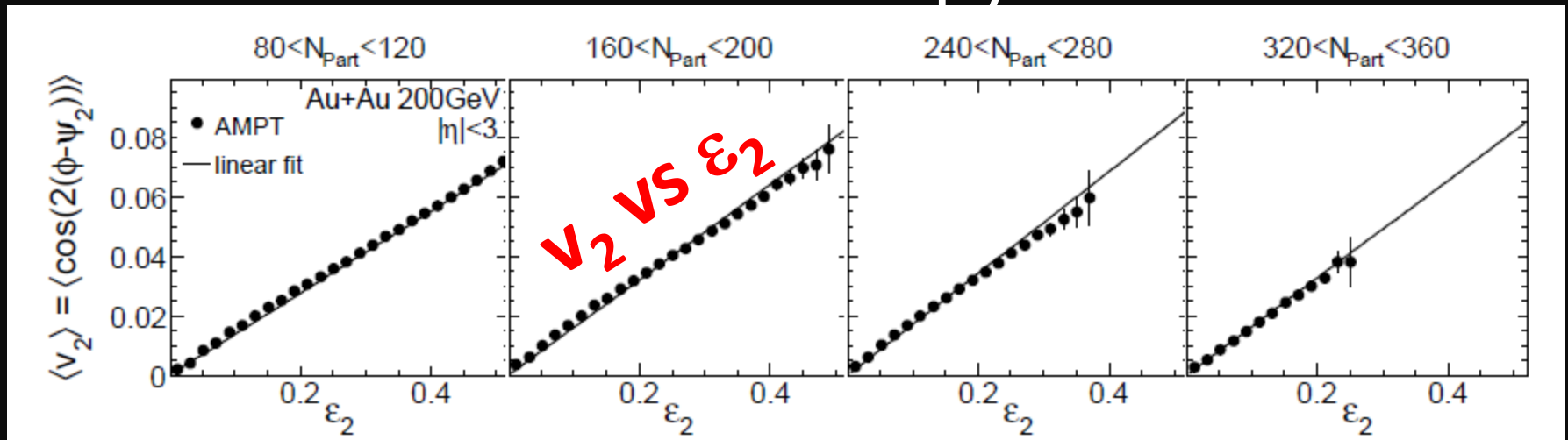
# Fluctuations Dominate

$$\sigma(v_2)/v_2$$

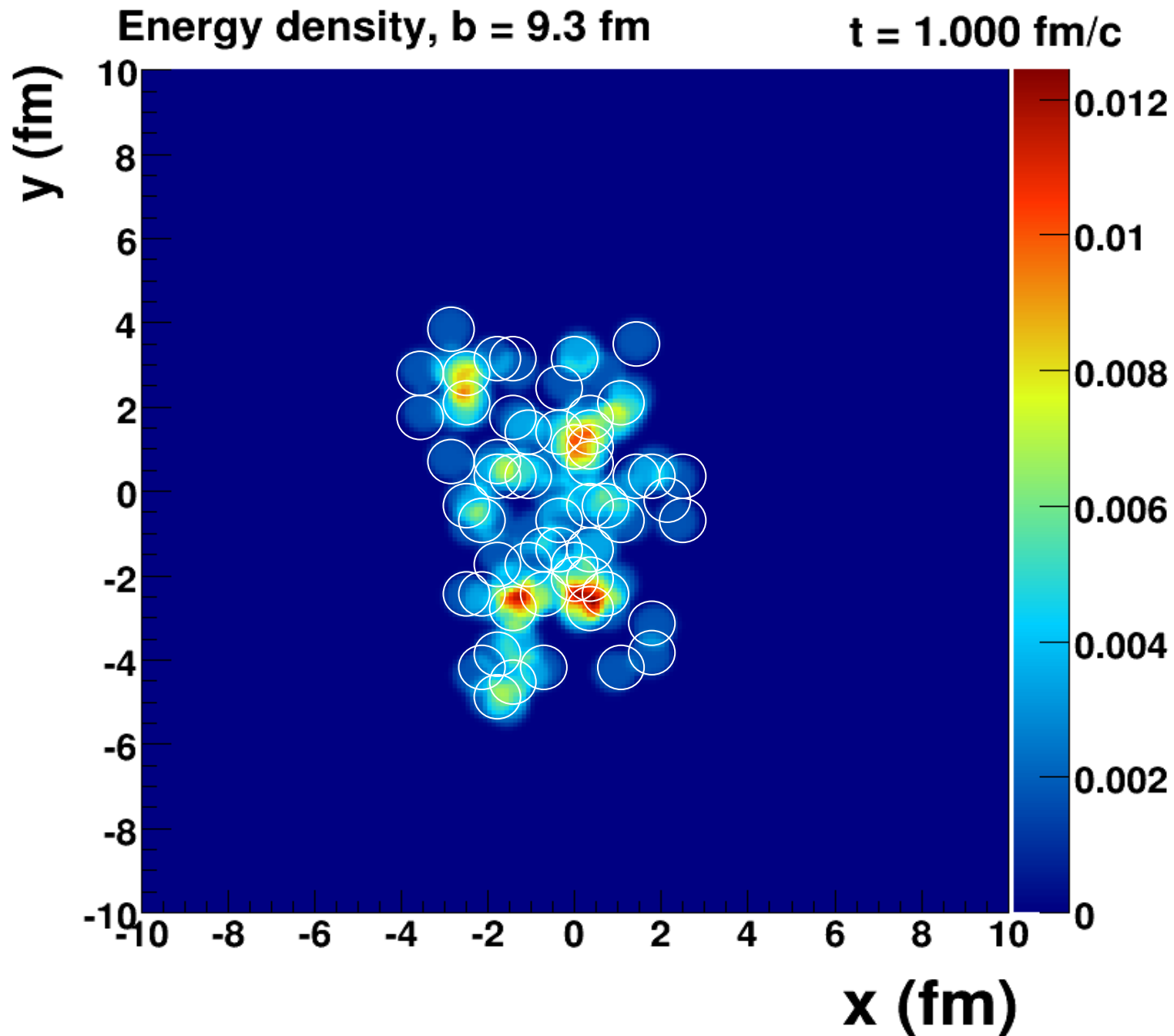
Flow dictated by  
nucleon geometry



# Spatial moments translate into momentum anisotropy moments



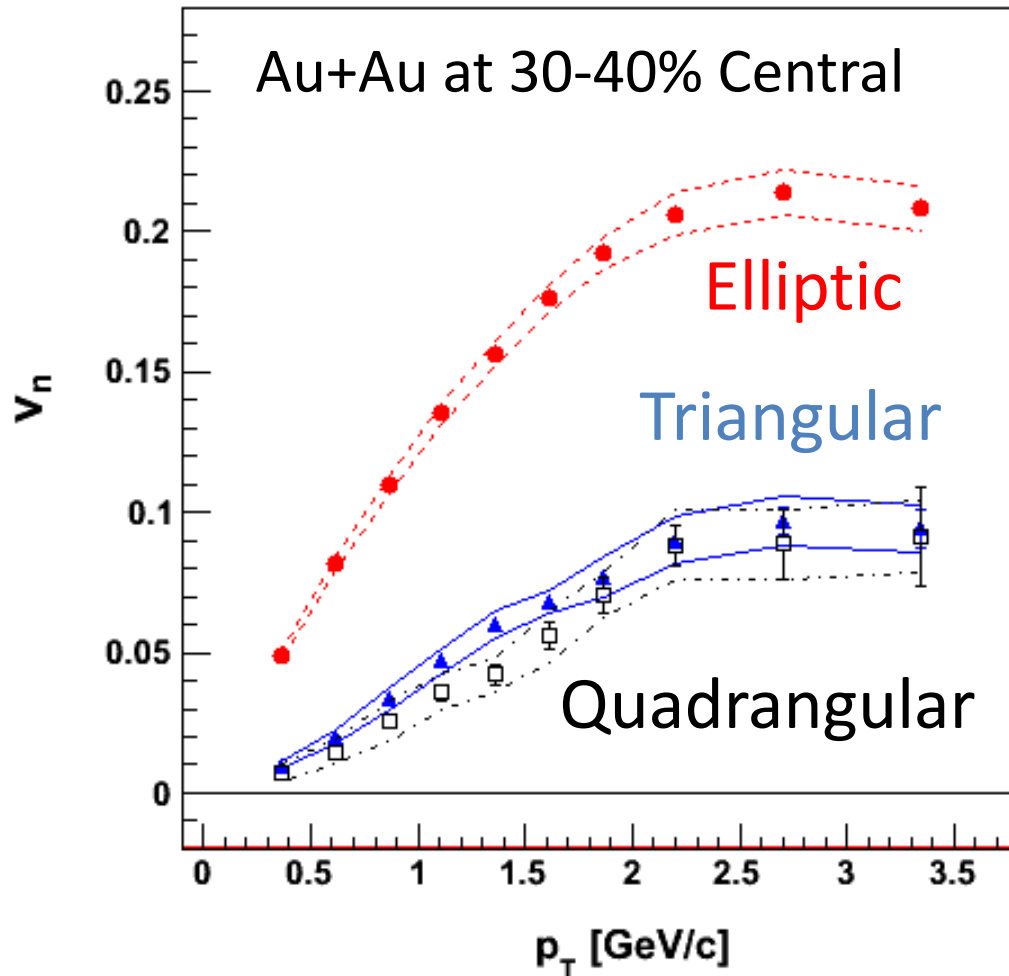
$v_3$  is as irrefutable as  $v_2$  from  $\epsilon_2$ .  
Now quantify implications.



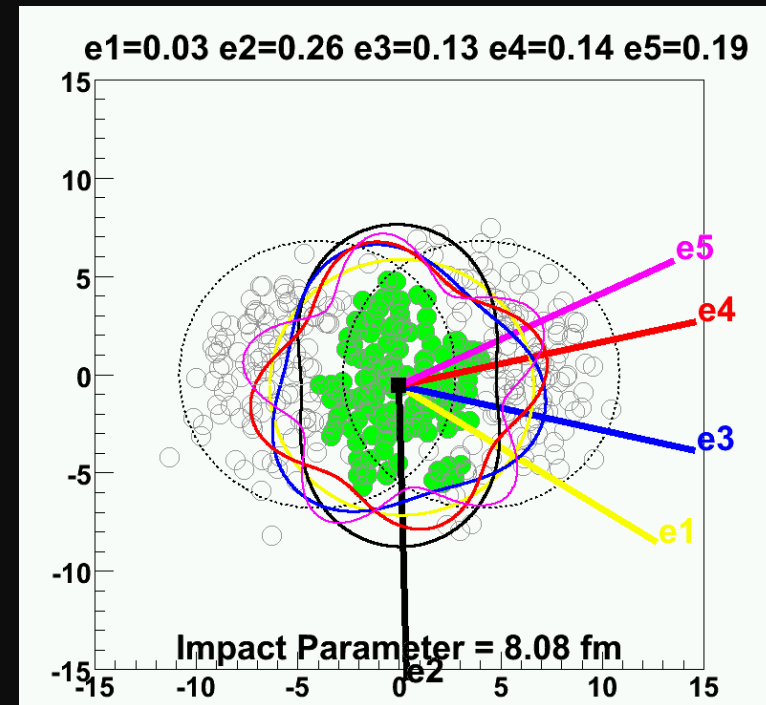
Romatschke=viscous hydrodynamics, McCumber=lumpy conditions + animation

# PHENIX Experiment

- $v_2 \{\Phi_2 \text{ forw.}\eta\}$
- ▲—  $v_3 \{\Phi_3 \text{ forw.}\eta\}$
- $v_4 \{\Phi_4 \text{ forw.}\eta\}$



Systematic uncertainties defined by the variations with  $\Phi_n$  from different  $\Delta\eta$  and from different methods.

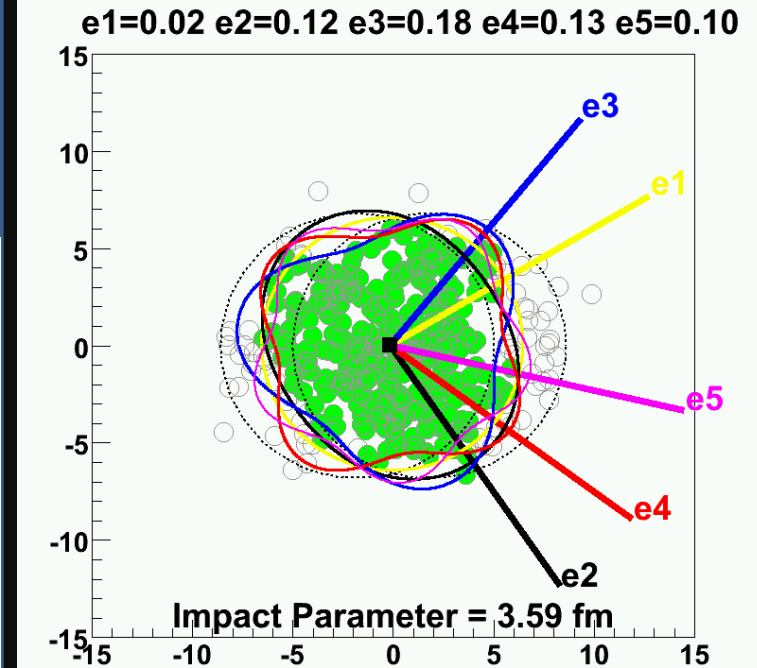
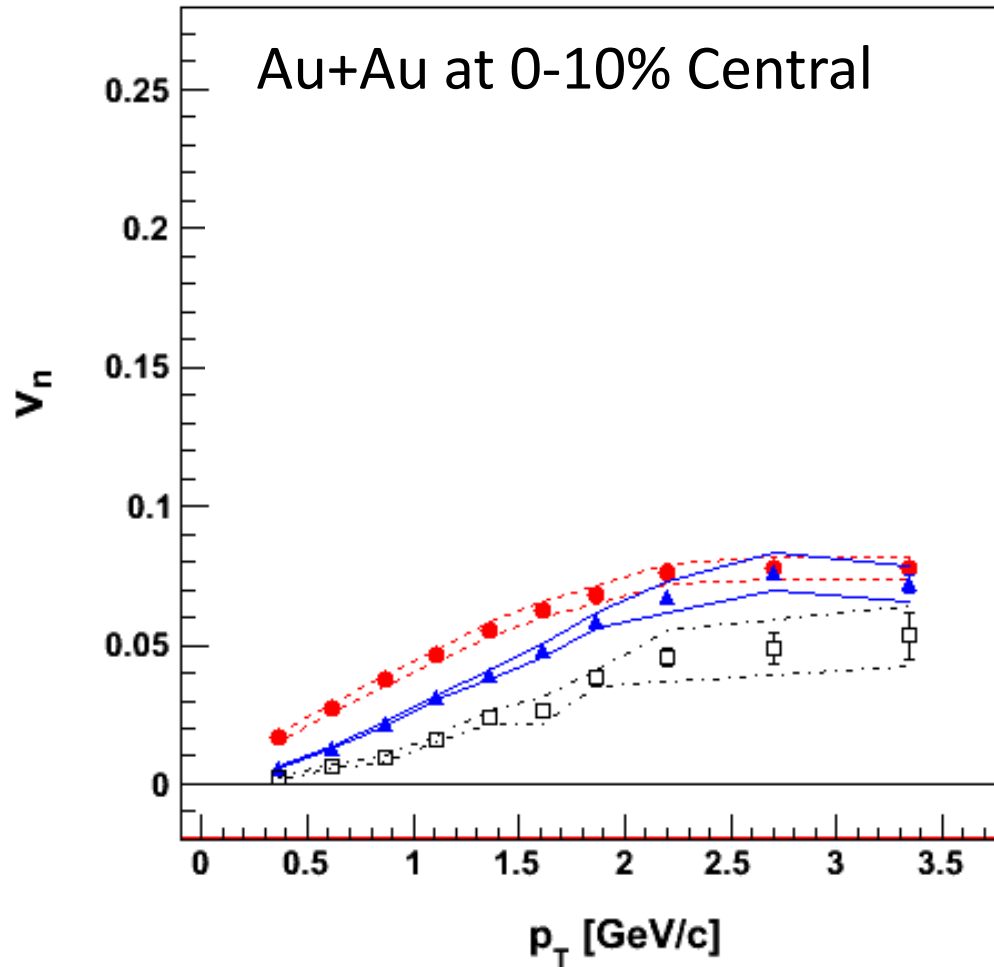


$$\varepsilon_2 \approx 2 \times \varepsilon_3 \approx 2 \times \varepsilon_4$$



# PHENIX Experiment

- $v_2 \{\Phi_2 \text{ forw.}\eta\}$
- ▲—  $v_3 \{\Phi_3 \text{ forw.}\eta\}$
- - □ - -  $v_4 \{\Phi_4 \text{ forw.}\eta\}$



$$\varepsilon_2 \approx \varepsilon_3 \approx \varepsilon_4$$

How many  
moments are  
important at RHIC?  
5,6,7<sup>th</sup>

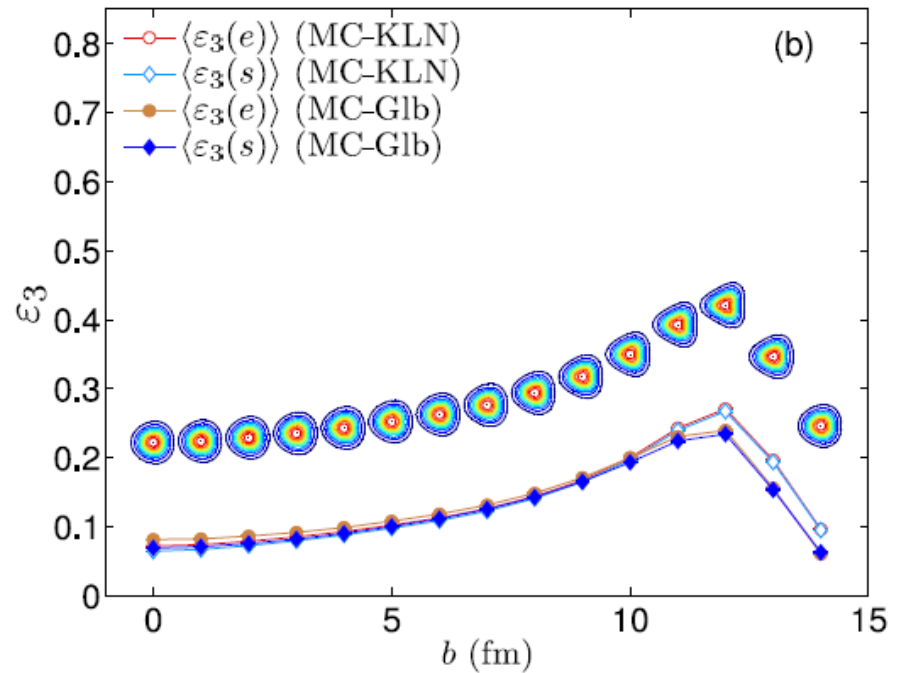
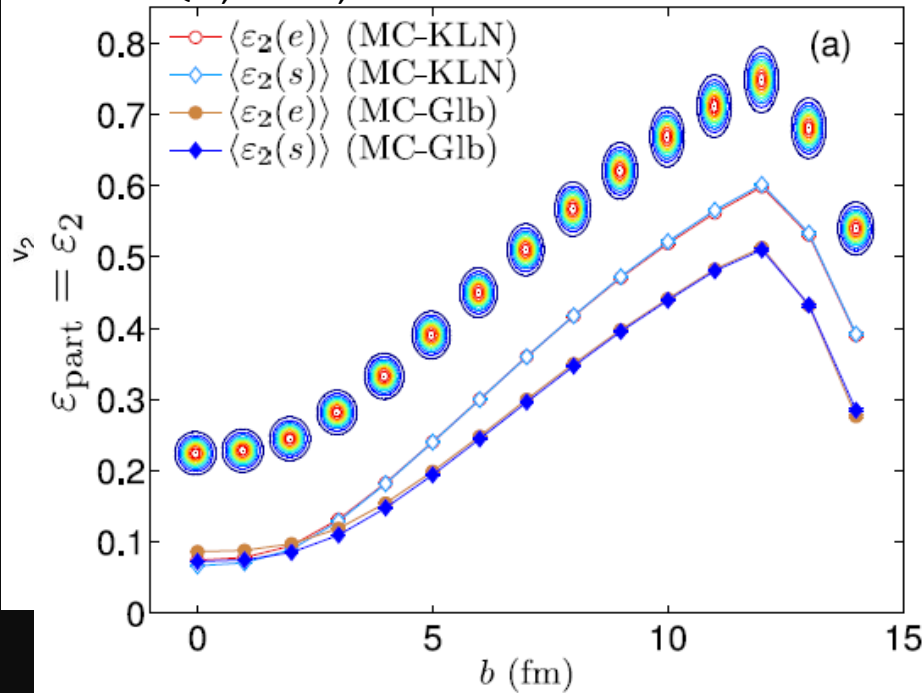
# Initial Condition Constraints

Glauber and CGC similar spatial triangularity  $\varepsilon_3$

Thus CGC with larger  $\eta/s$  gives smaller  $v_3$

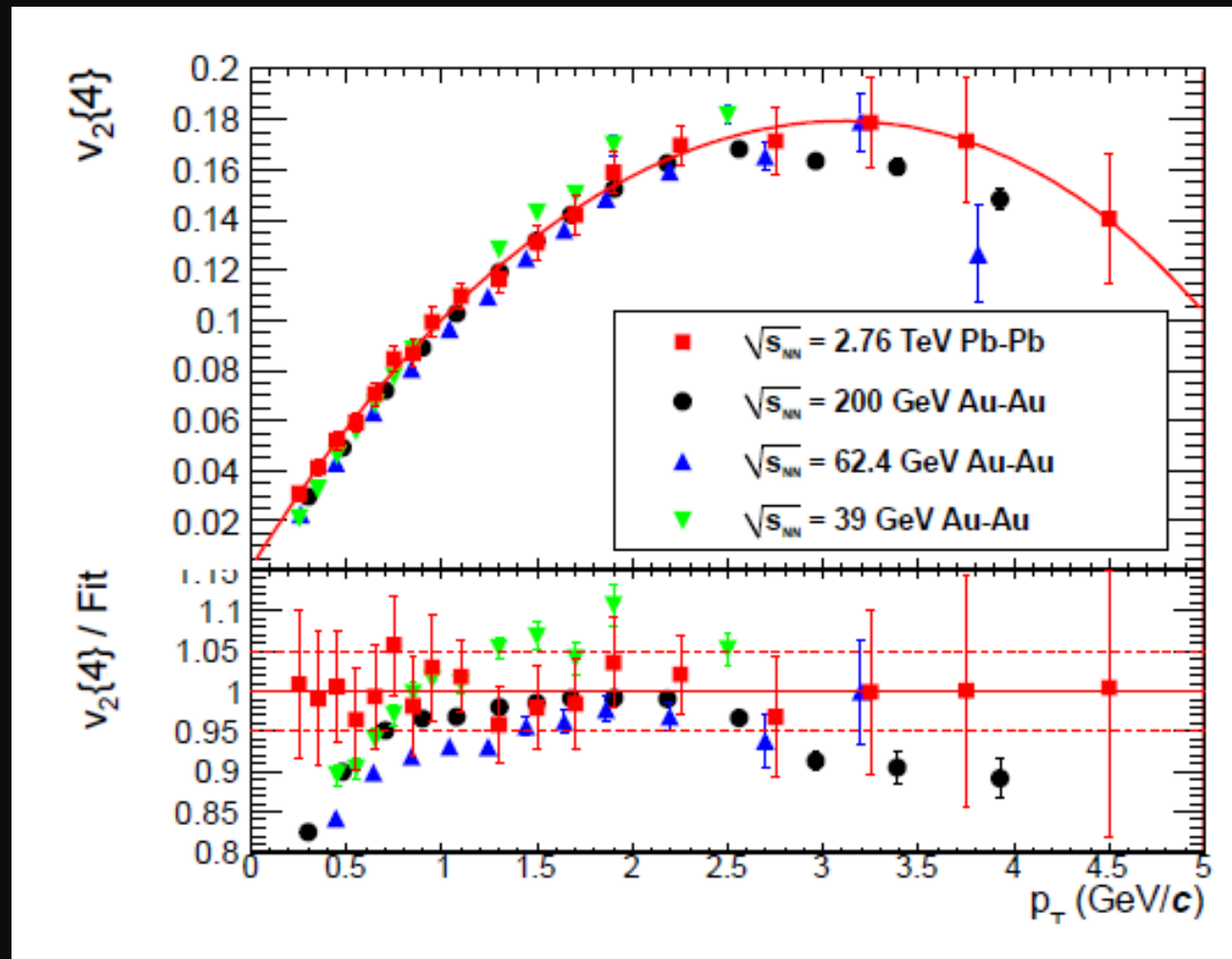
*Real Prediction!*

Qiu, Heinz, arXiv:1104.0650v2



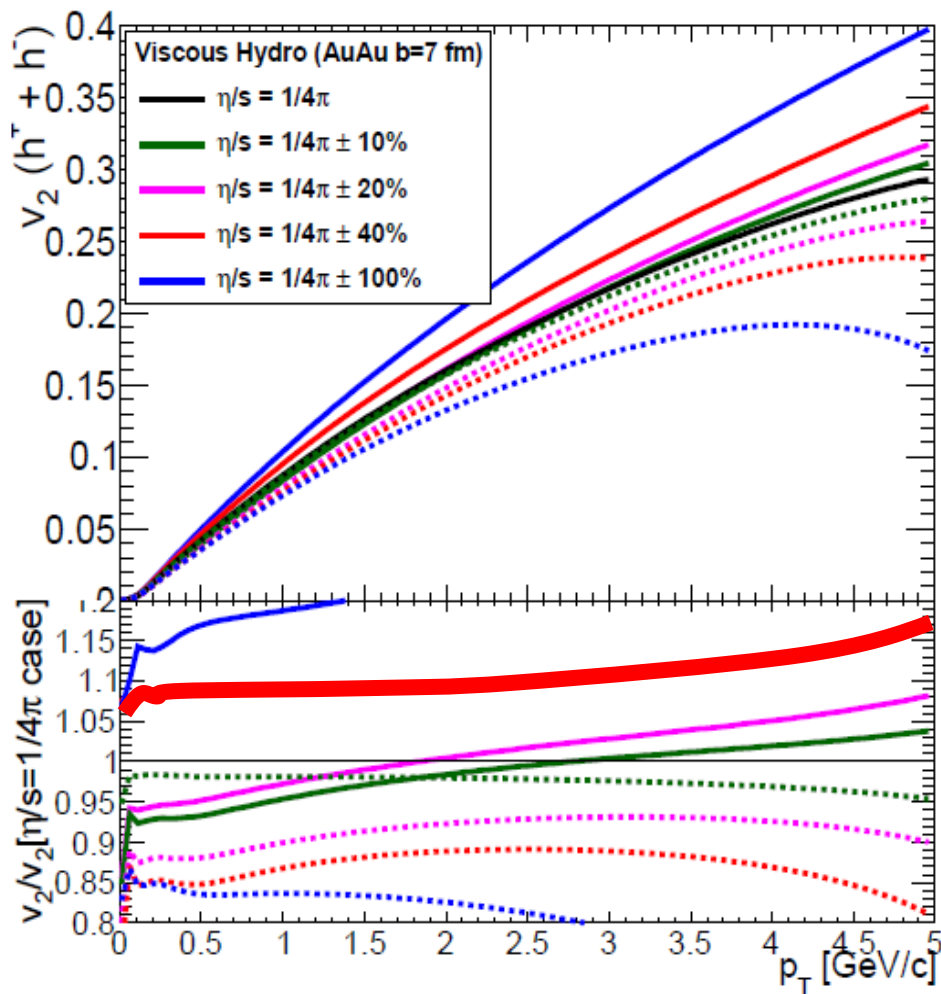
Needs theory calculation with full event-by-event fluctuations and confirmation at LHC energies.

# How does flow change with 14 x energy?



What is required to get this remarkable agreement?  
*JN, Bearden, Zajc, arXiv:1102.0680*

# Change in $\eta/s$ to change $v_2$ by 5%?



$\Delta(\eta/s) \sim 20\% \rightarrow \Delta v_2 \sim 5\%$

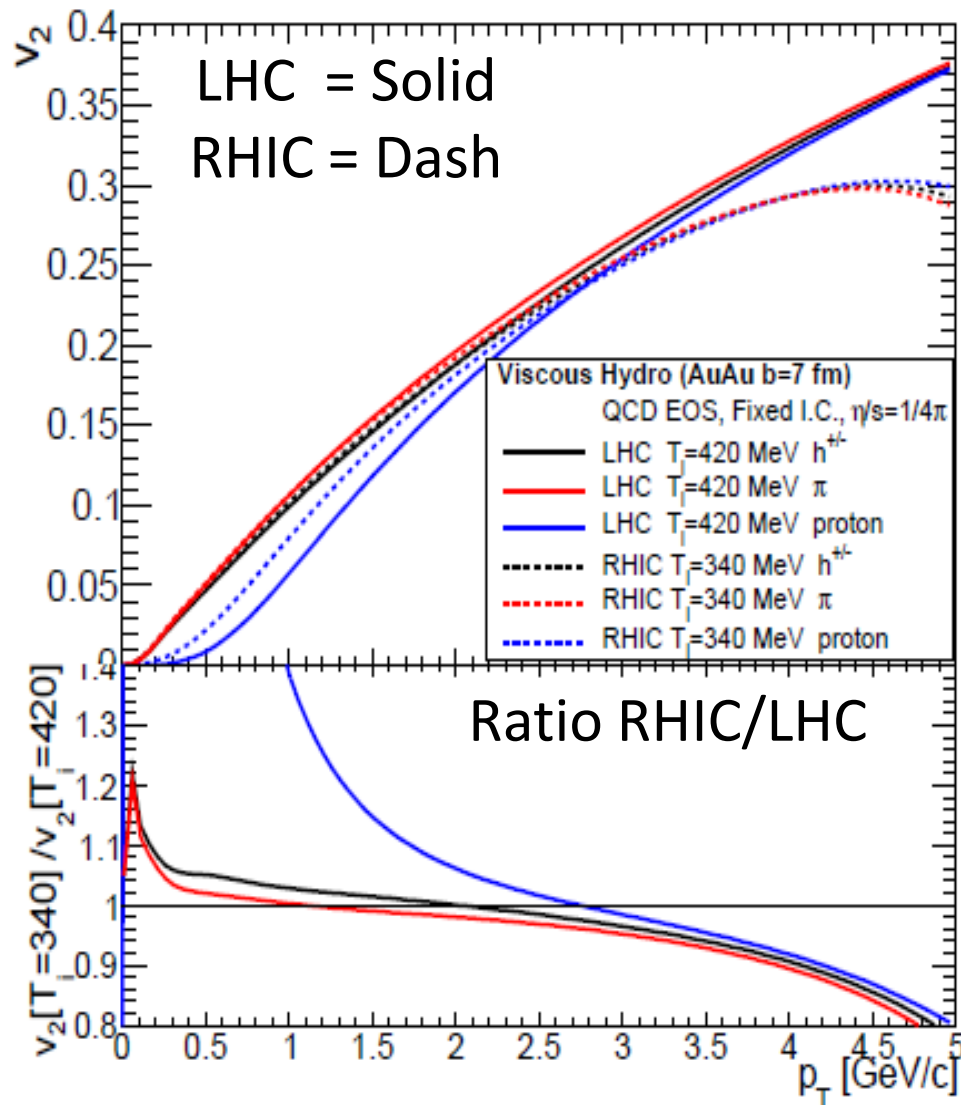
$\Delta(\eta/s) \sim 40\% \rightarrow \Delta v_2 \sim 10\%$

$\Delta(\eta/s) \sim 100\% \rightarrow \Delta v_2 \sim 25\%$

Always take ratios.

Low  $p_T$  turn out to be  
very sensitive  
and where  
experiments have  
smallest uncertainties

# $T_i = 420$ MeV (LHC) versus $T_i = 340$ MeV (RHIC)?



Black = all hadrons

Very similar  $v_2(p_T)$  except  
larger viscous effects for  
 $p_T > 3$  GeV/c

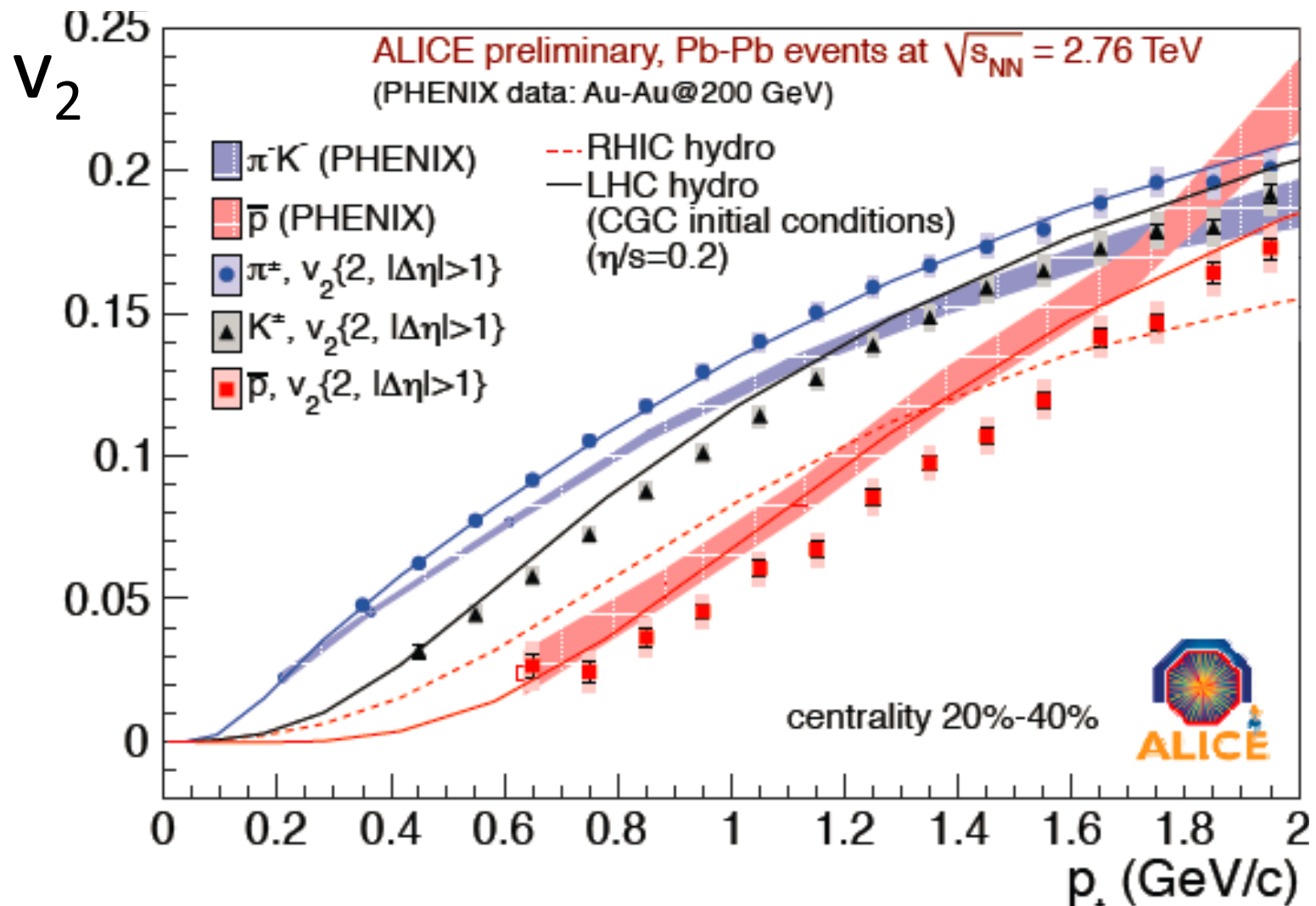
Blue = protons

Large difference for  $v_2(p_T)$   
due to larger radial boost  
at LHC temperatures.

Solid prediction.

Previously noted with ideal  
hydrodynamics by Kestin, Heinz

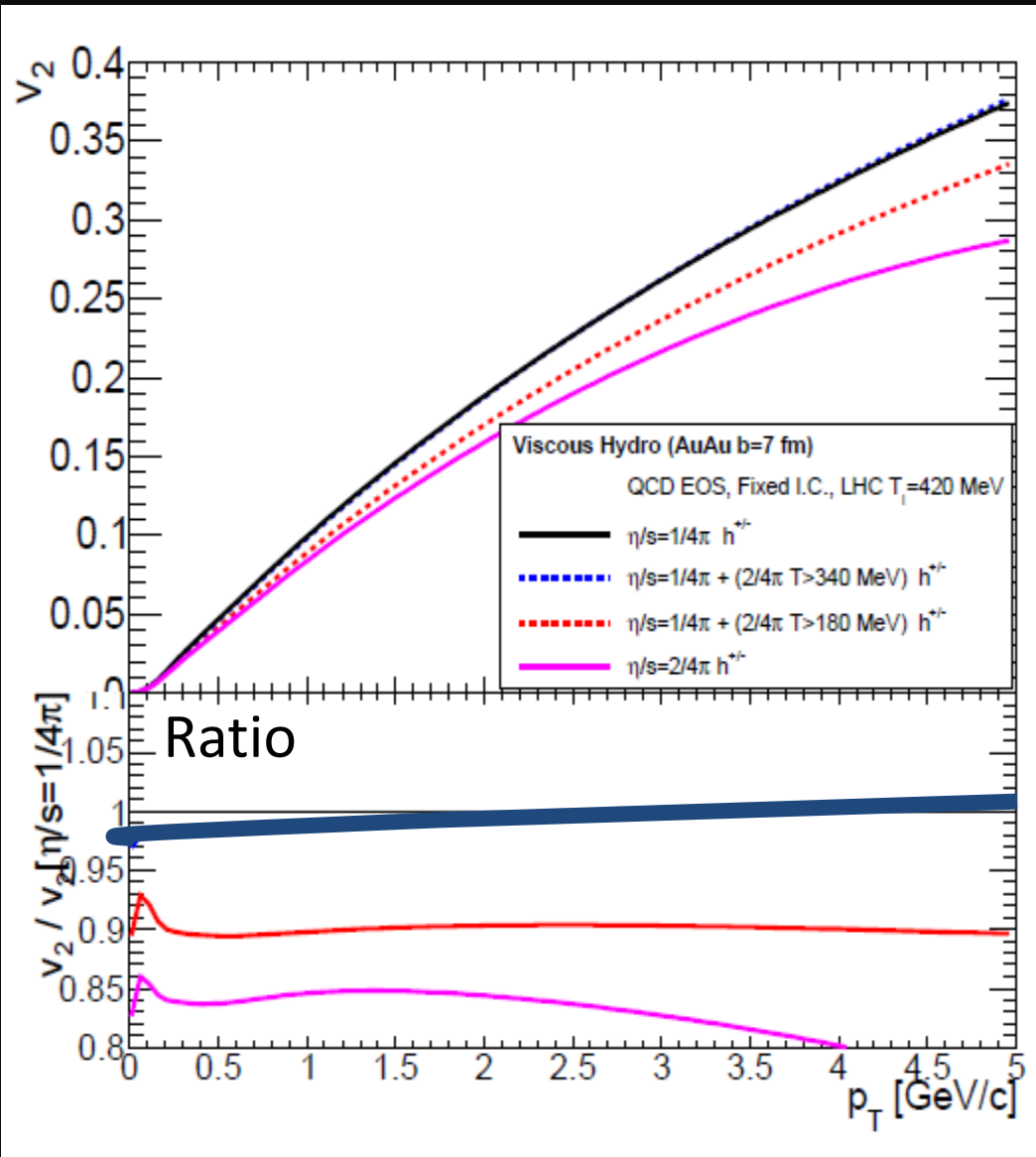
# General Feature Confirmed in ALICE Data



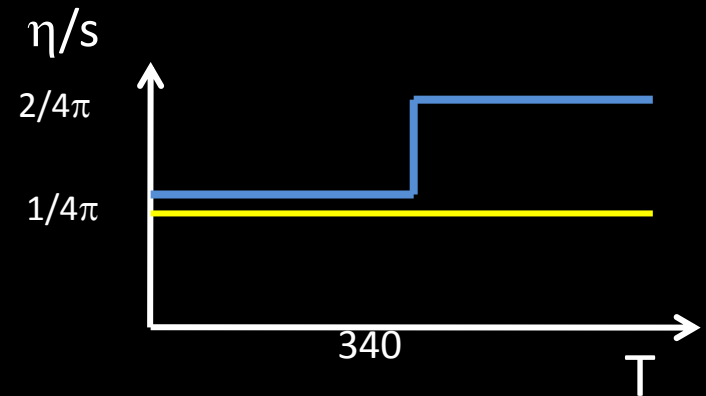
Exact radial boost still not matching in detail  
Try Hydro + Cascade afterburner

# What if $\eta/s$ is larger for $T > 340$ MeV?

*(just the range sampled at the LHC in early times)?*



Consider case I:

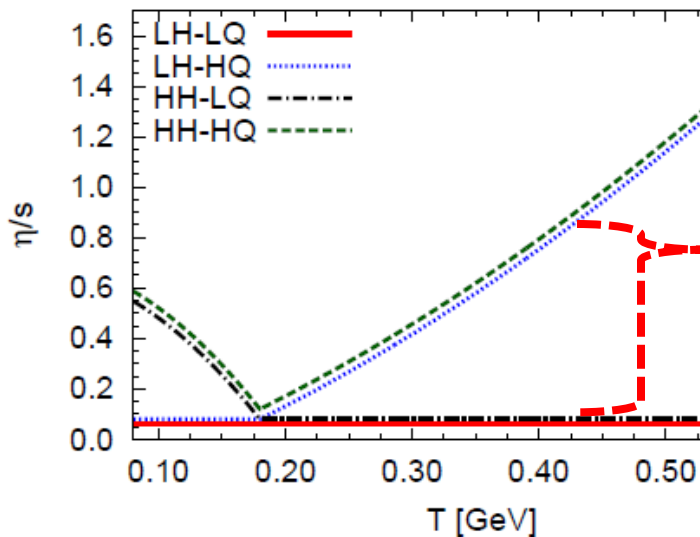


No change in  $v_2(p_T)$ !

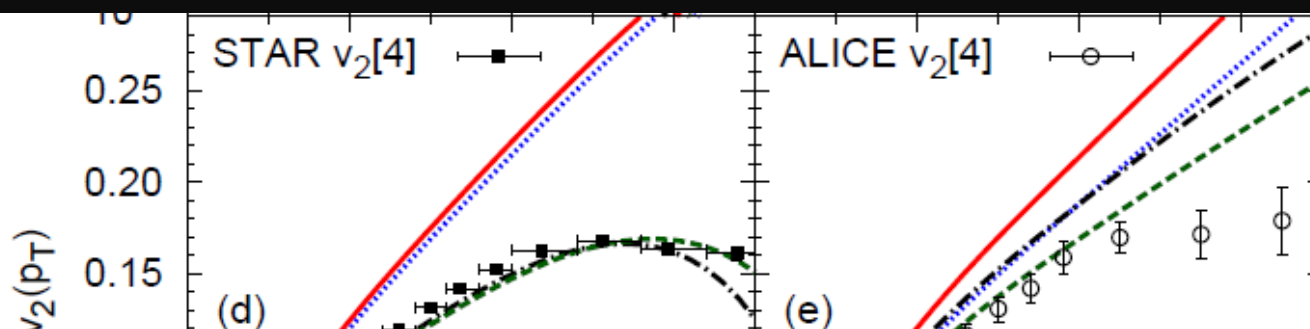
Earliest LHC time has no big impact on  $\eta/s$ .



# Recent study of $\eta/s$ (T) – arXiv:1101.2442

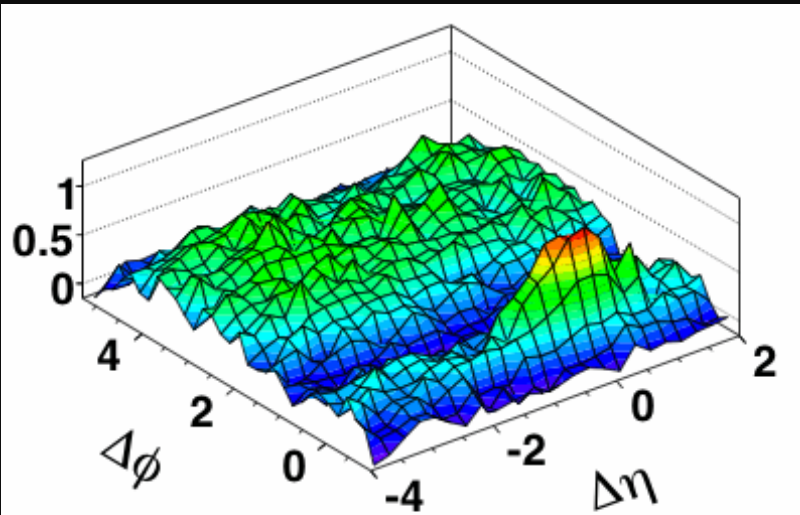


Factor 10 increase in  $\eta/s$  results in 15% reduction in  $v_2$  at  $p_T = 2$  GeV/c  
Seems consistent with my finding that a factor of 2 for  $T > 340$  MeV has almost no effect.



‘RHIC  $v_2$  is dominated by  $\eta/s$  below  $T_c$  and LHC  $v_2$  by  $\eta/s$  above  $T_c$ ’  
***seems an unlikely fine tuning problem***

# “Ridge” and “Shoulders”

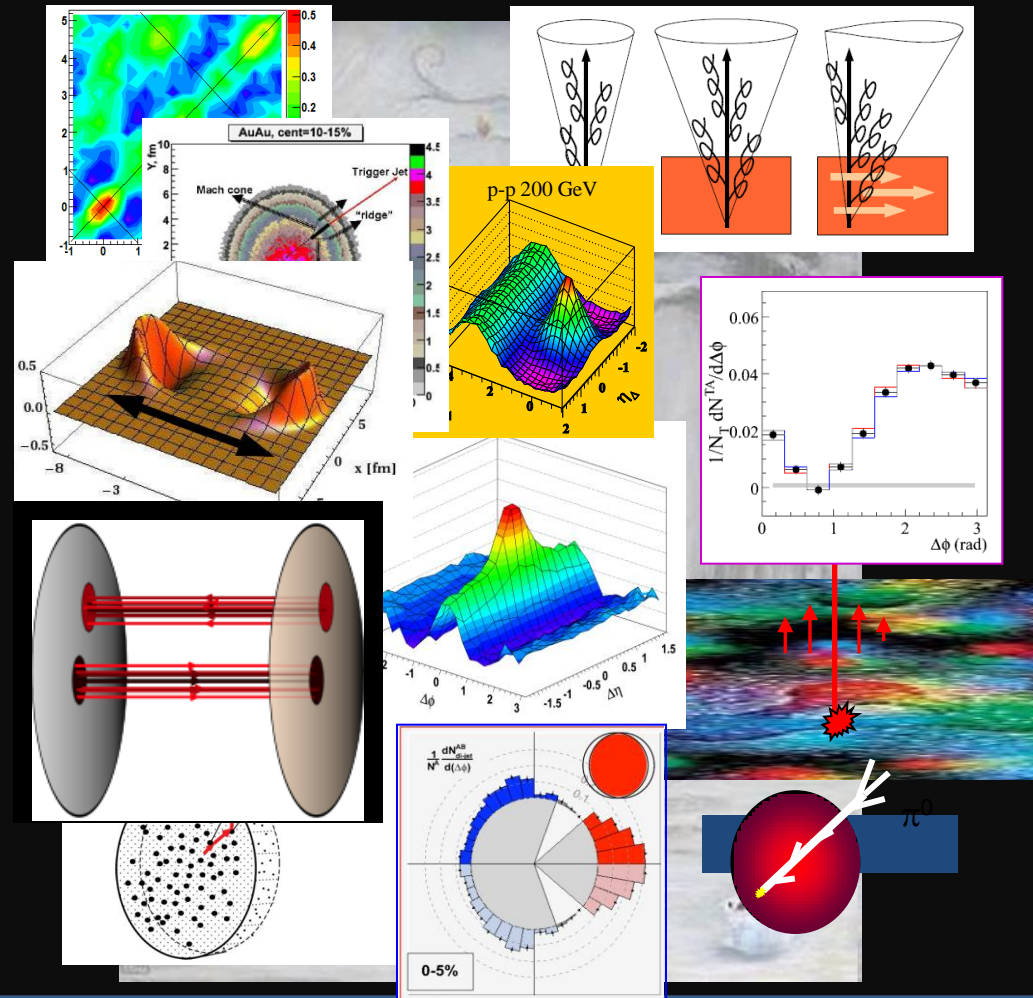


(b) Au+Au 0-30% (PHOBOS)

Two years ago (QM09)  
I gave a talk with my  
prediction.

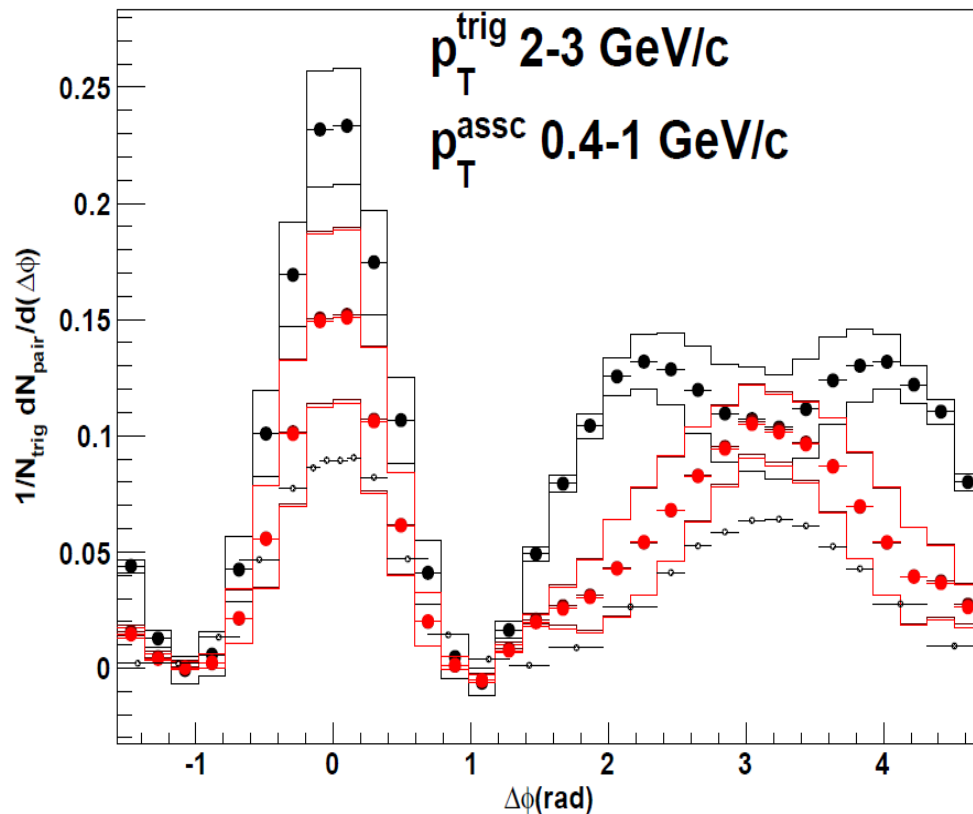
The “death” of the ridge  
and shoulders  
(*aka* shock response)

Features in two-particle  
correlations that have  
generated a lot of excitement.



# Re-Evaluation of Shock Response

PHENIX Au+Au 0-20% Central



## Black points

PHENIX published with  $v_2$  background modulation and ZYAM.

## Red points

(NOT PHENIX OFFICIAL)

Use AMPT  $v_3/v_2$  ratio and include  $v_2$  and  $v_3$  background modulation.  
Calculation by A. Adare

Dominant ridge and shoulders will be gone.

Detailed careful analysis needs  $v_1 \dots v_5$  and method checks.

# LHC Results

“Ridge” and “Shoulders” seen at LHC as well...

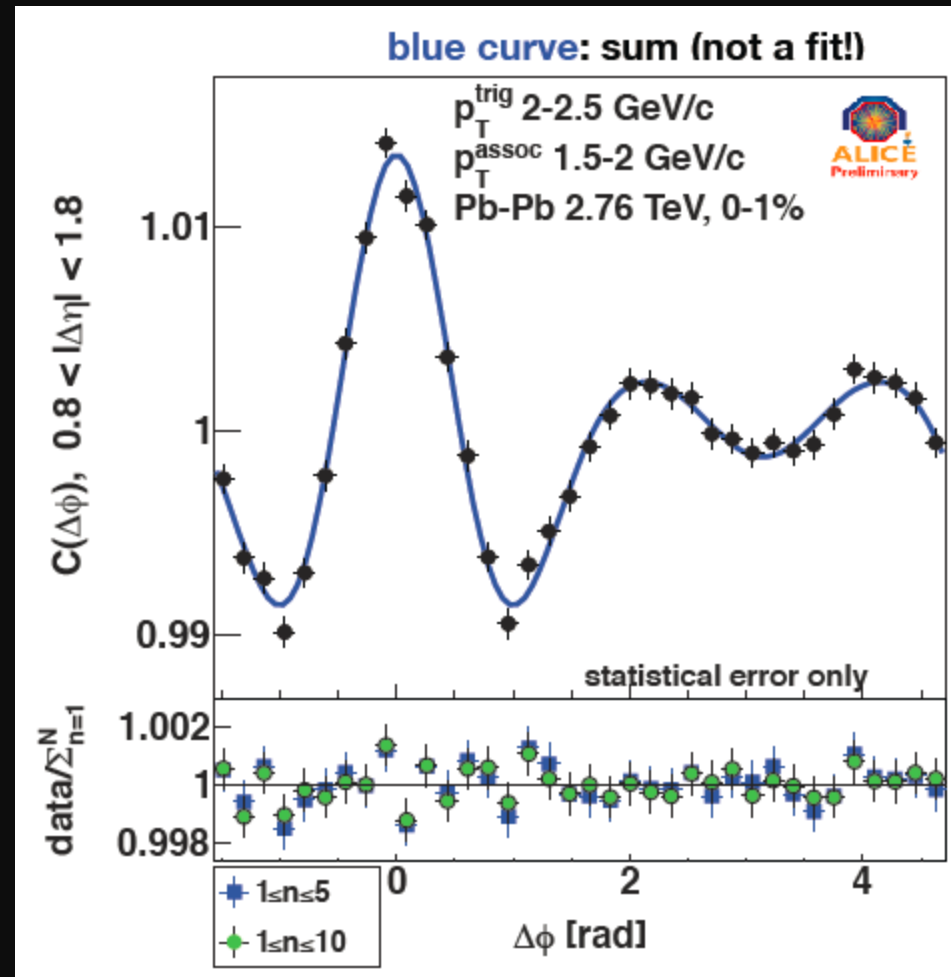
ALICE: arXiv:1105.3865

Fourier

Decomposition  
characterizes the  
distribution

Are the Fourier  
Coefficient just the  
higher order flow  
moments?

Key Tests



# Sound Mode Commentary

**Viscous Horizon:** “Its verbal definition is that it separates the wavelengths of sound which *are* and *are not* dissipated by viscosity effects.” i.e. damped to “un-observably small magnitude” [Shuryak].

Smooth medium and fluctuations

$$T_{\mu\nu} = \tilde{T}_{\mu\nu} + \delta T_{\mu\nu}$$

Dispersion relation from shear viscosity

$$\omega = c_s k - \frac{i}{2} \frac{4\eta}{3s} \frac{k^2}{T}$$

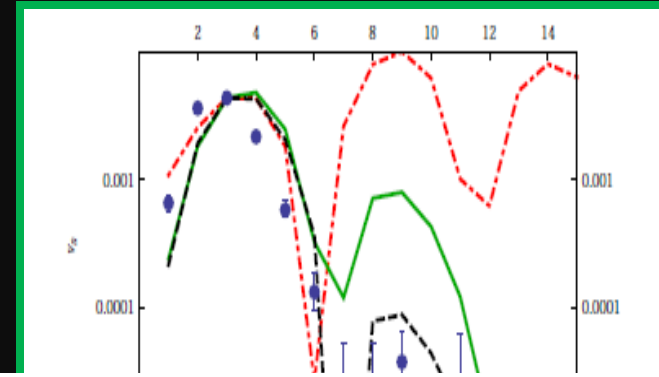
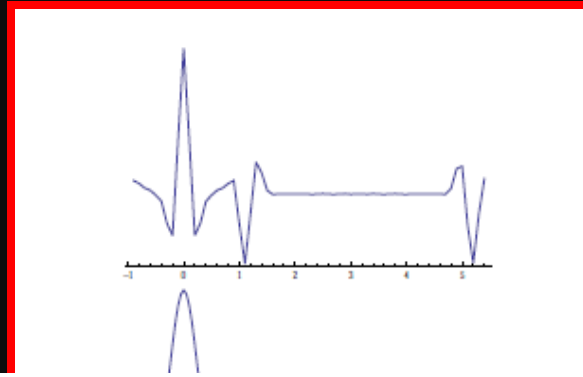
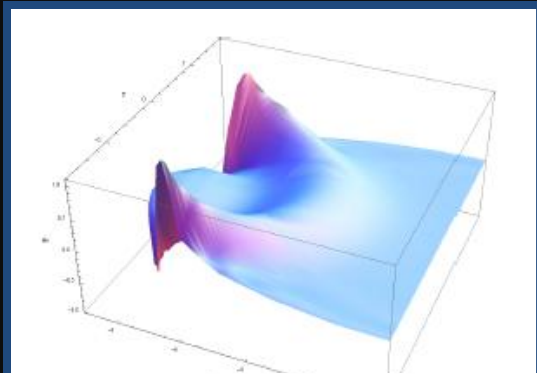
Damping of fluctuations  
depends on wavenumber (k),  
time (t),  $\eta/s$  and Temperature.

$$\delta T_{\mu\nu}(t) = \exp\left(-\frac{2}{3} \frac{\eta}{s} \frac{k^2 t}{T}\right) \delta T_{\mu\nu}(0)$$

Need more formal definition of viscous horizon,  
e.g. when wavelength mode damped by 1/e.

“Un-observably small” depends on experimental sensitivity.

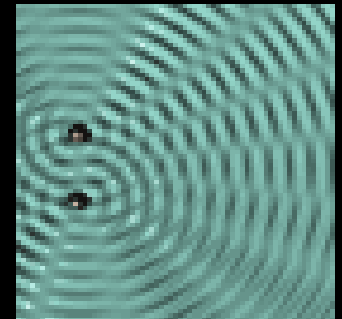
One then evolves with hydrodynamic equations this “hot spot”.  
At the freeze-out surface, apply Cooper-Frye hadronization and  
calculate angular distribution of particles.  
Finally, one decomposes the angular distribution into harmonics ( $n$ ).



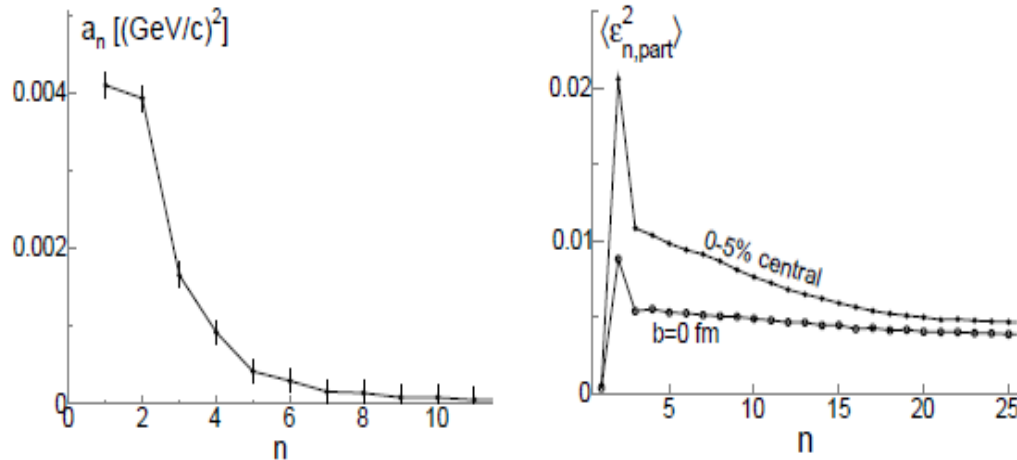
Structure due to detailed interaction of hot spot waves  
with freeze-out surface and cut-off effects.

Needs double check on robustness of higher moment  
structures with realistic geometry.

Also, check of how lumpy medium  
impacts propagation of state  
(not small perturbations).



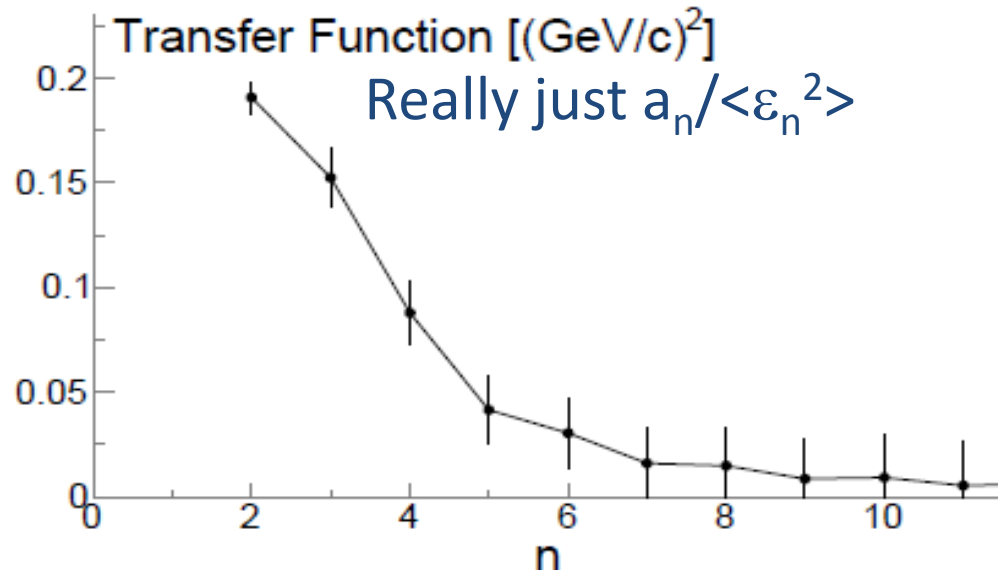
Mocsy, Sorensen use STAR published  $p_T$ - $p_T$  and  $\Delta\phi$  correlator ( $a_n$ ) which in the limit of just number correlations  $\approx v_n^2$ .



They argue that the transfer function goes to zero when  
 $\ell(\text{mean free path}) \approx$   
 $\lambda$  (wavelength of the mode) =  
 $2\pi R/n$

They choose  $n=5$  as the cut-off (viscous horizon) and  $\langle R \rangle = 3$  fm (average radial position of participants) and get  $\ell=3.5$  fm.

They state plugging into kinetic formula gives  $\eta/s \sim 5 \times 1/4\pi$ .



<http://arxiv.org/abs/arXiv:1008.3381> (Mocsy, Sorensen)

**Lots of factors of 2**  
**floating around**



Shuryak's damping equation  
by wavenumber (k)

$$\delta T_{\mu\nu}(t) = \exp\left(-\frac{2}{3} \frac{\eta}{s} \frac{k^2 t}{T}\right) \delta T_{\mu\nu}(0)$$

They attempt to relate k to angular  
harmonic moment n  
However, k is mode in radial  
direction and n is angular  
correlation for momentum vectors.

$$n\lambda = 2\pi R$$

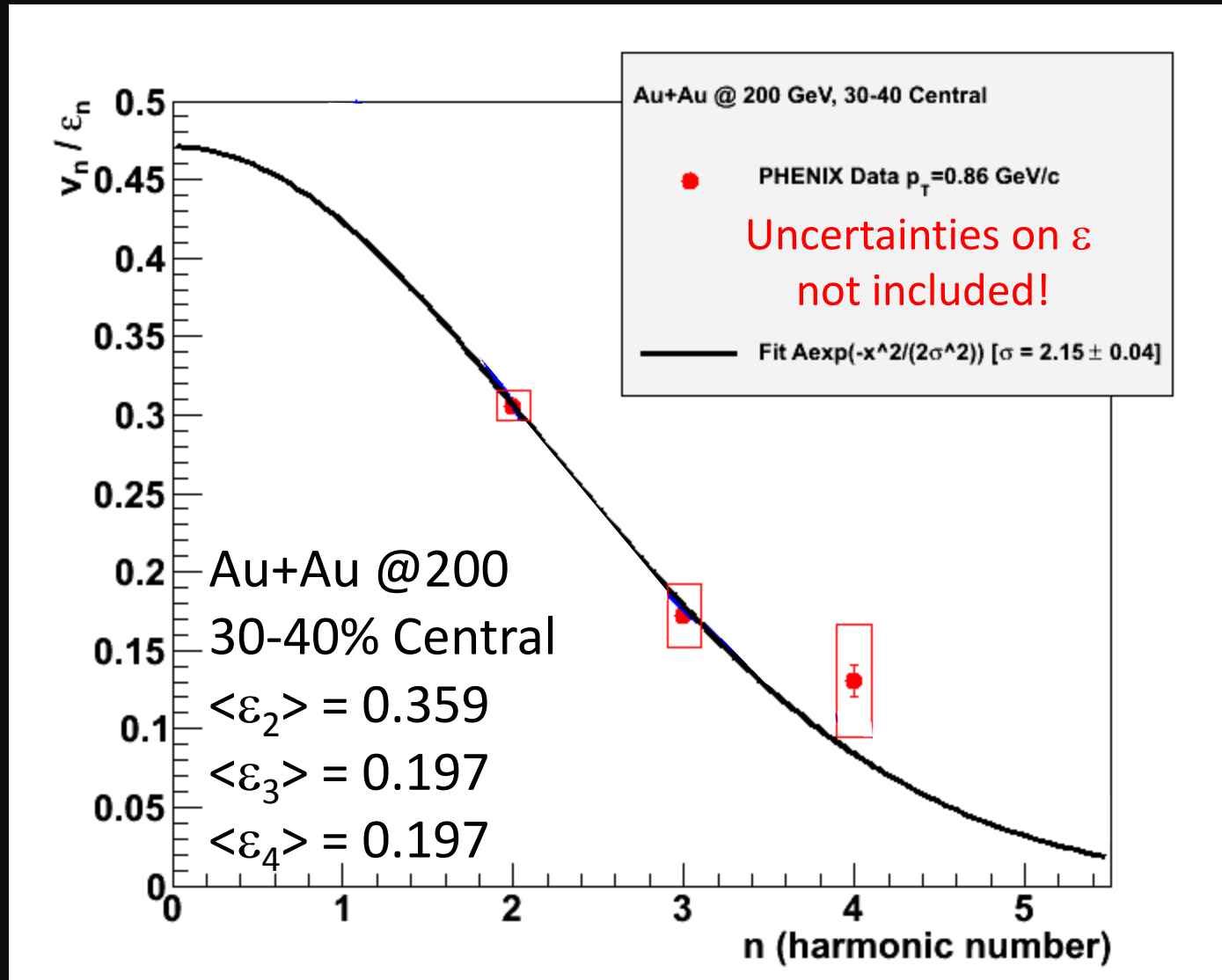
**I believe this is  
simply incorrect.**

$$k = \frac{2\pi n}{\lambda} = \frac{n}{R}$$

$$\frac{\delta T_{\mu\nu}(t)}{\delta T_{\mu\nu}(0)} = \exp\left(-\frac{2}{3} \frac{\eta}{s} \frac{n^2}{R^2} \frac{t}{3T}\right)$$

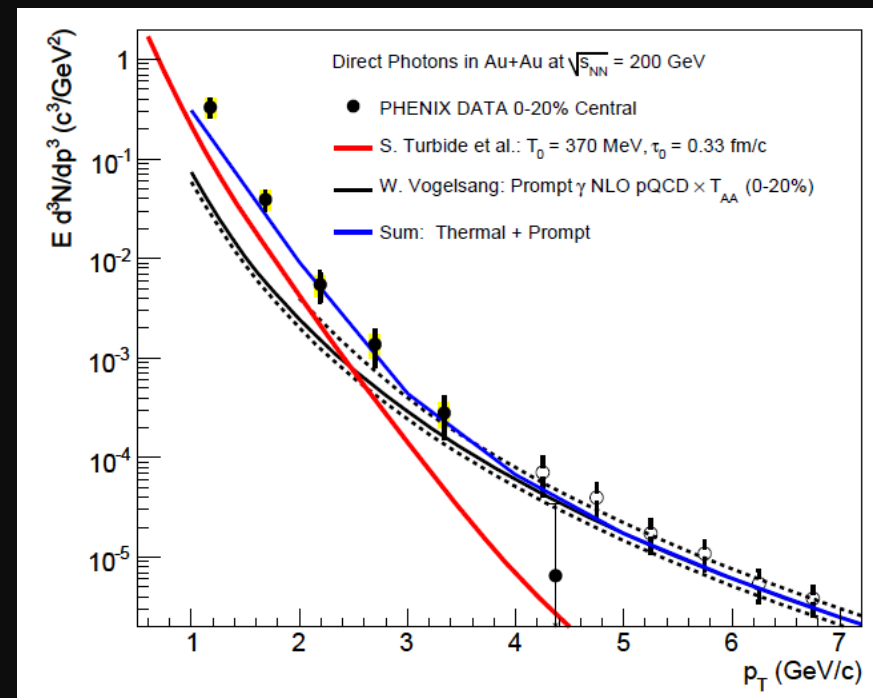
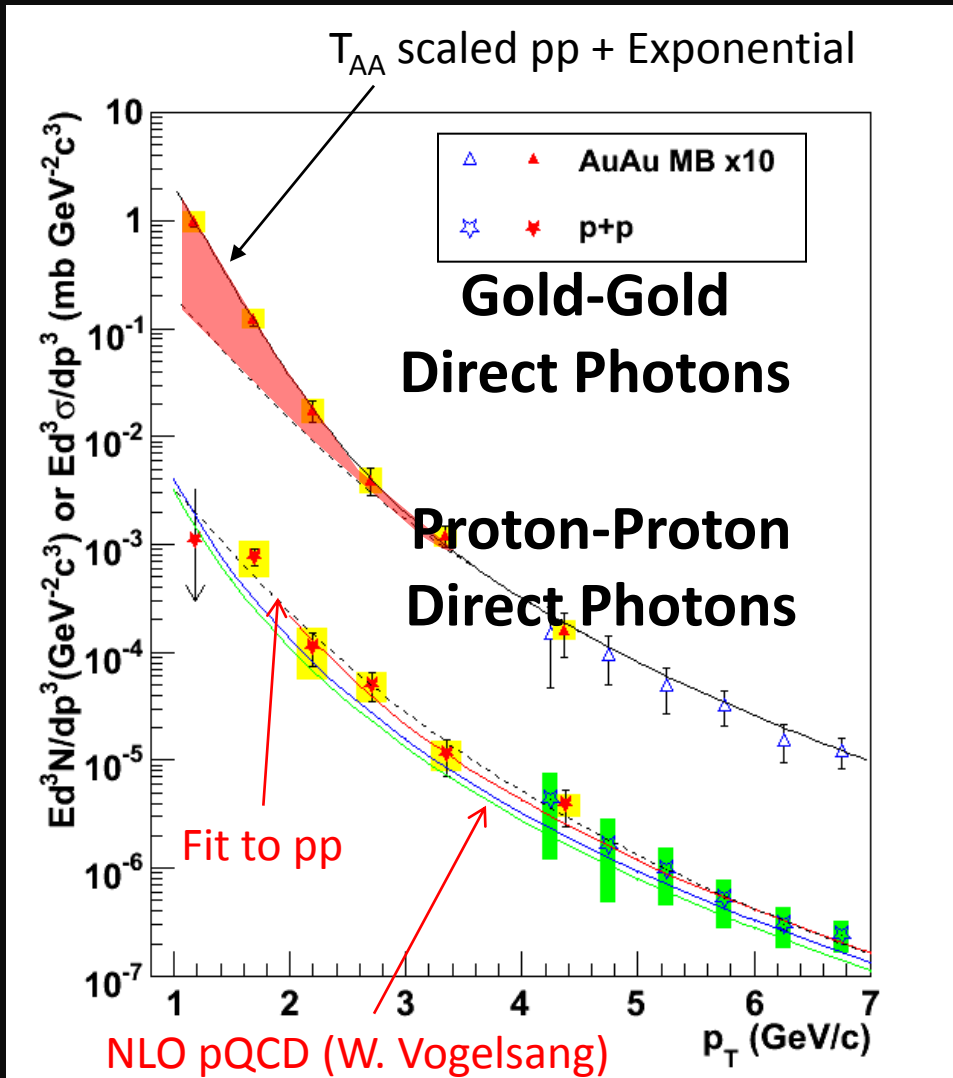
$$\frac{\delta T_{\mu\nu}(t)}{\delta T_{\mu\nu}(0)} \propto \frac{v_n}{\varepsilon_n} \propto \exp\left(-\frac{2}{3} \frac{\eta}{s} \frac{n^2}{R^2} \frac{t}{3T}\right)$$

Interesting feature of larger damping for higher moments.



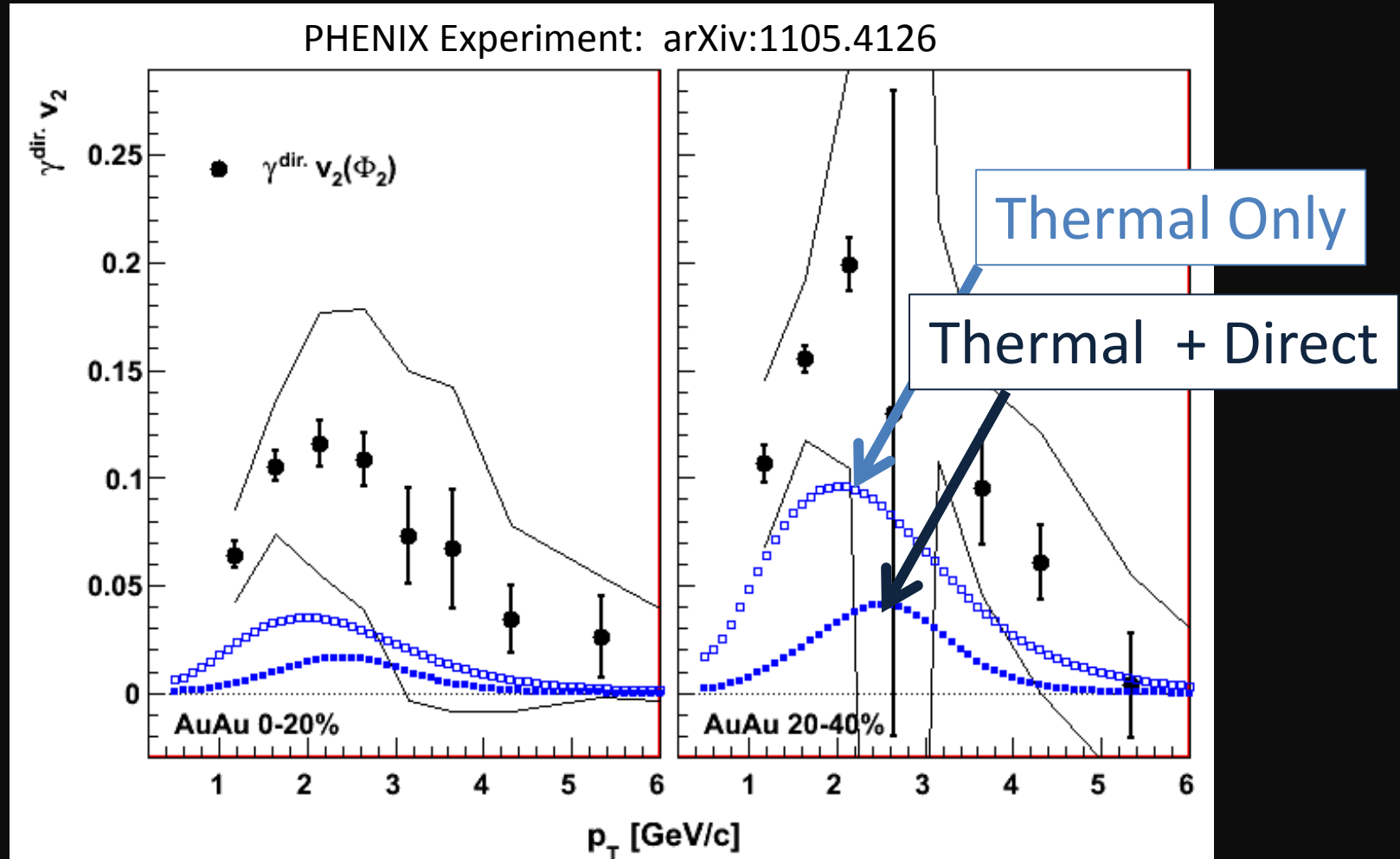
However, I believe full viscous hydrodynamics needed to relate to  $\eta/s$  (*no skipping steps*).

# Thermal Photon Emission



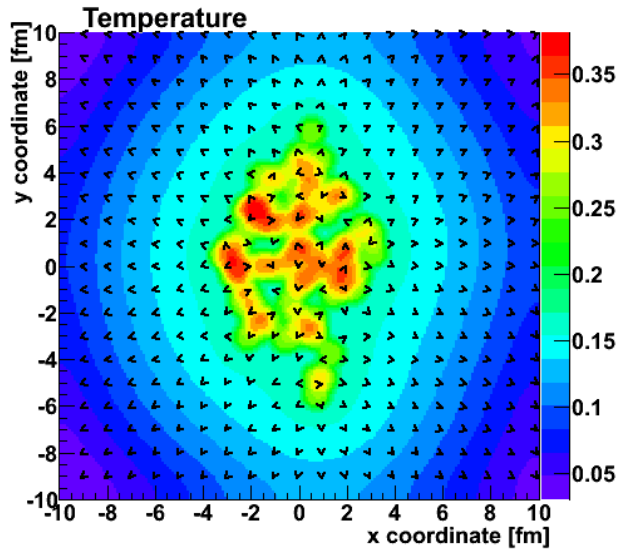
Specific prediction  
for flow ( $v_2$ ) of  
thermal photons

# Challenge of Direct Photon Flow

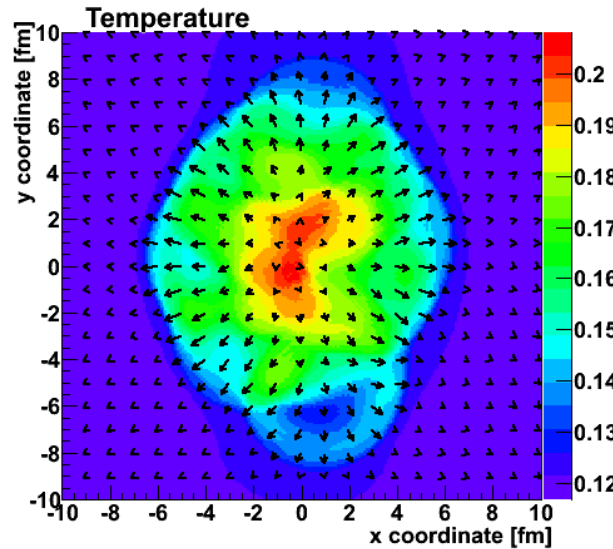


# Hydrodynamic Calculation

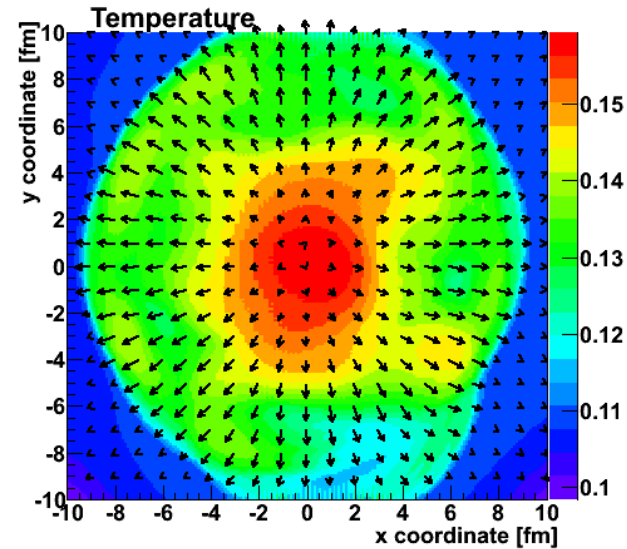
*KZ Mendoza, M. McCumber, JN, P. Romatschke*



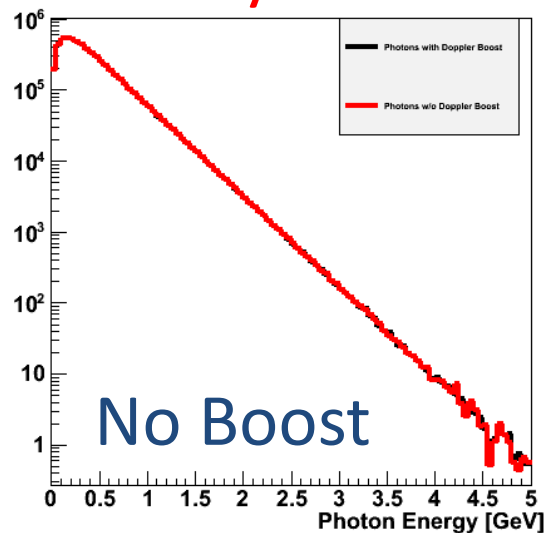
Early Time



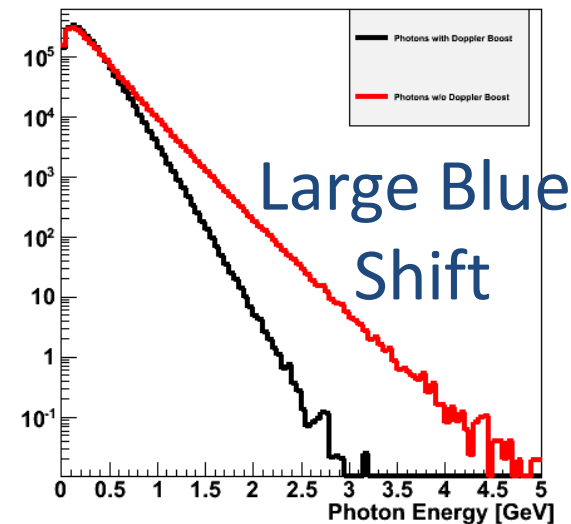
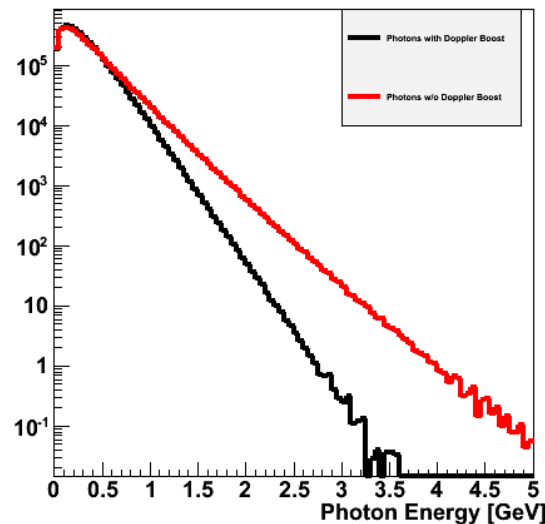
Middle Time



Late Time



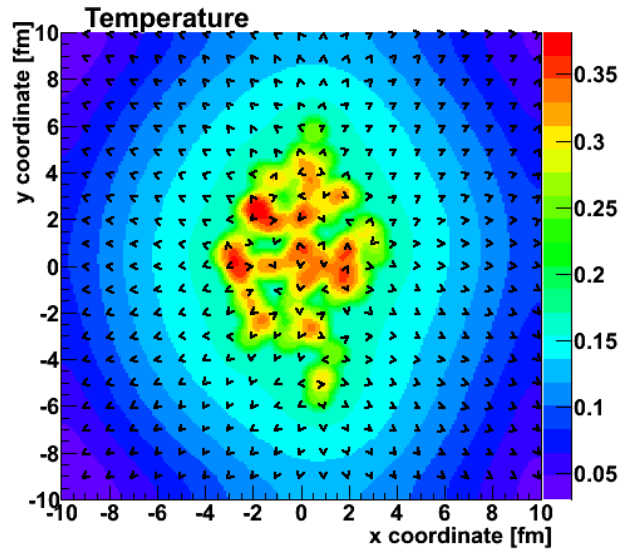
No Boost



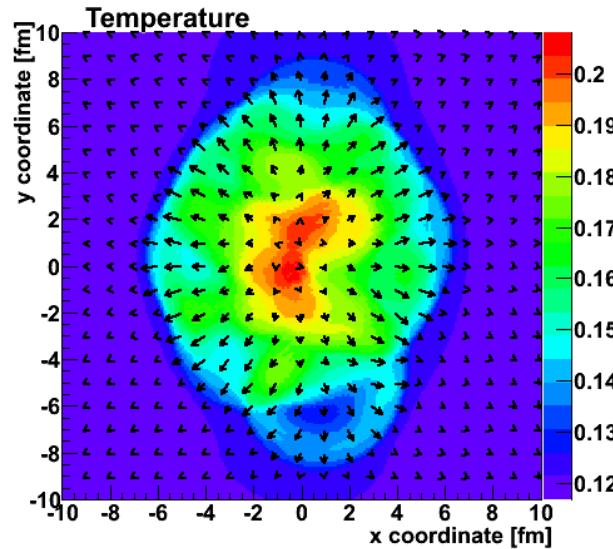
Large Blue Shift

# Hydrodynamic Calculation

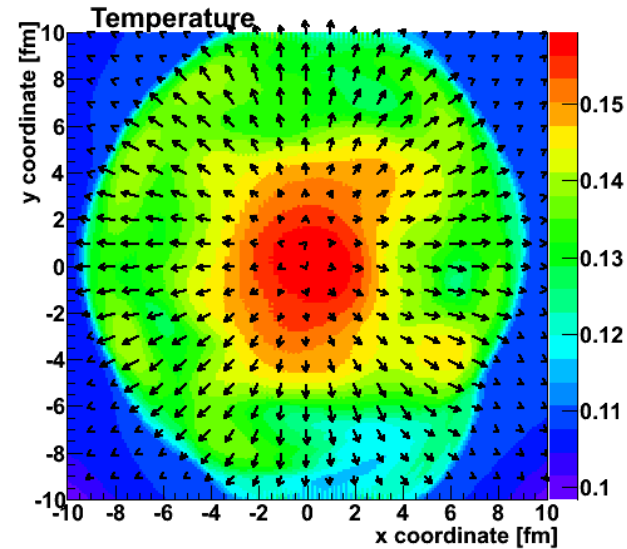
*KZ Mendoza, M. McCumber, JN, P. Romatschke*



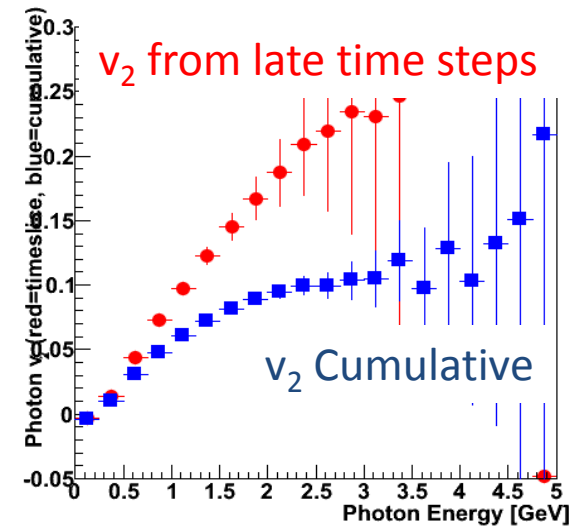
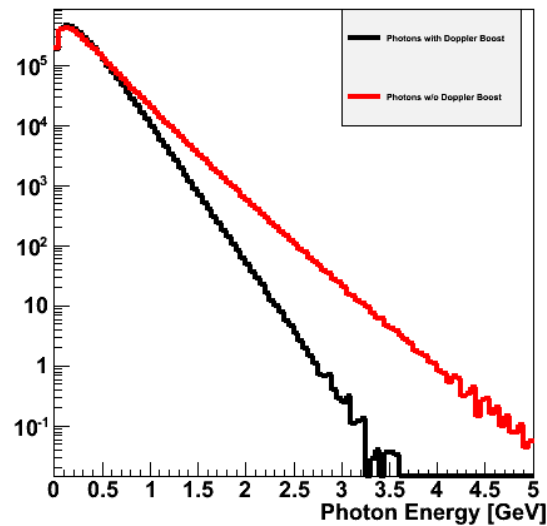
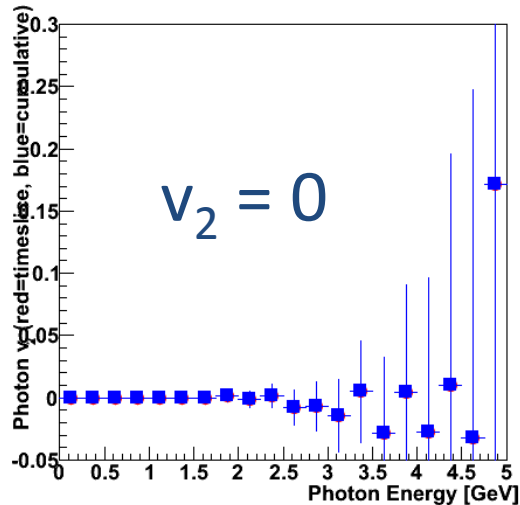
Early Time



Middle Time



Late Time



# Summary

- Enormous progress in this area in the last one year
- Dramatic reduction in  $\eta/s$  uncertainties around the corner
  - Exotic shock response and ridge features gone
- Key is now to understand how localized energy is so quickly transformed into effective heating  
(nicely related to jet quenching puzzle)
  - Major upgrades at RHIC planned.  
Critical for complementing the studies at LHC and really measuring the full excitation function



# PHENIX Decadal Plan

## Major Upgrade Proposed Extending into EIC Era

### The PHENIX Experiment at RHIC

*Decadal Plan 2011–2020*

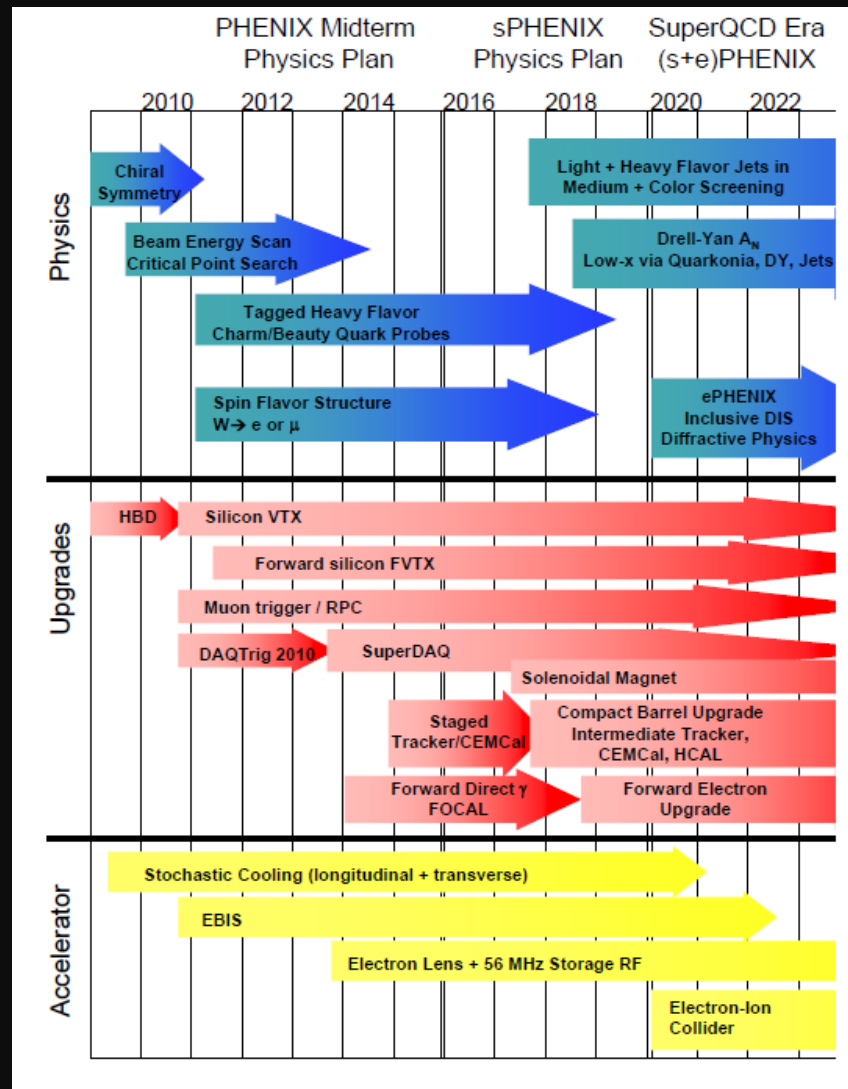
Brookhaven National Laboratory

Relativistic Heavy Ion Collider

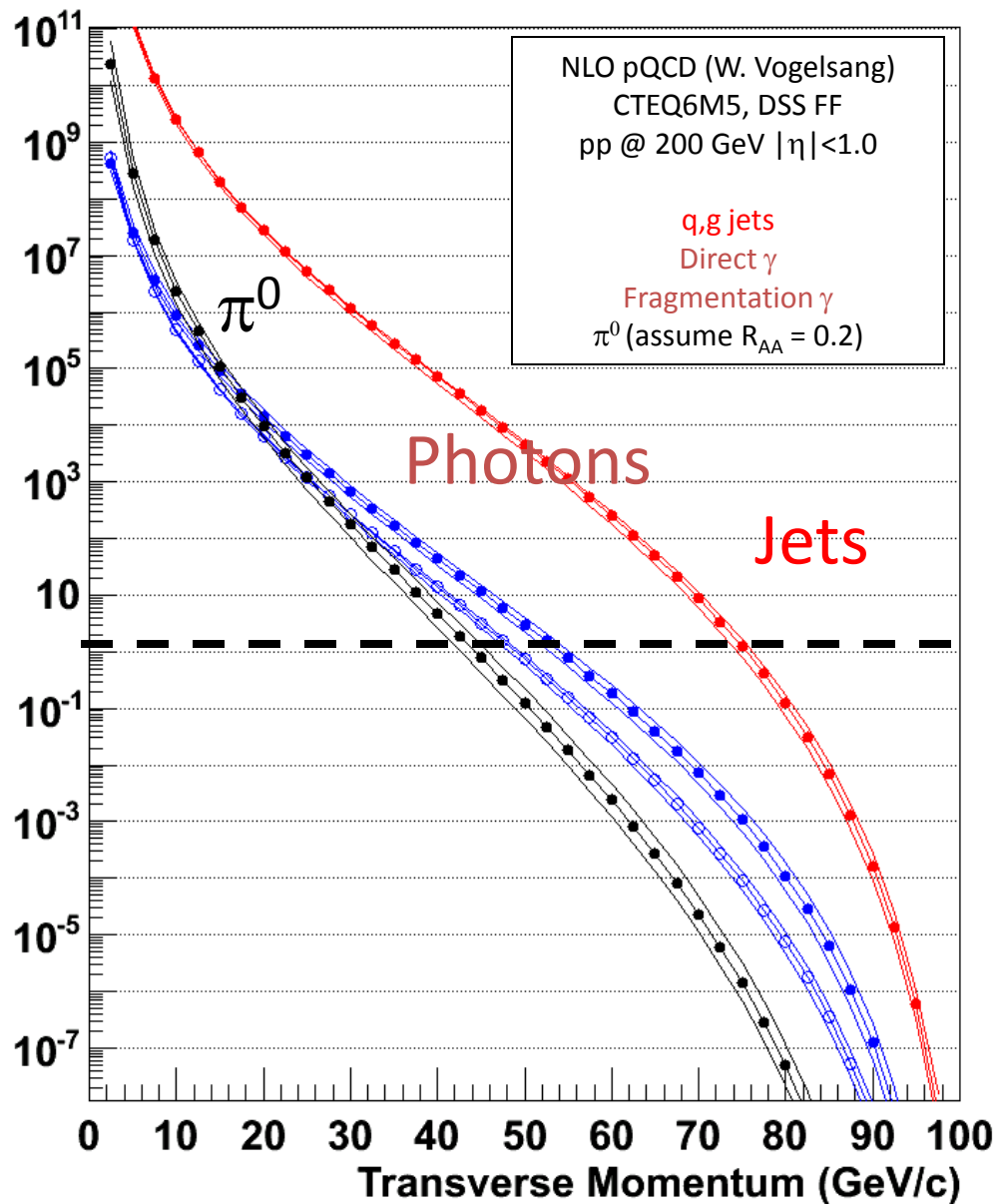
October, 2010



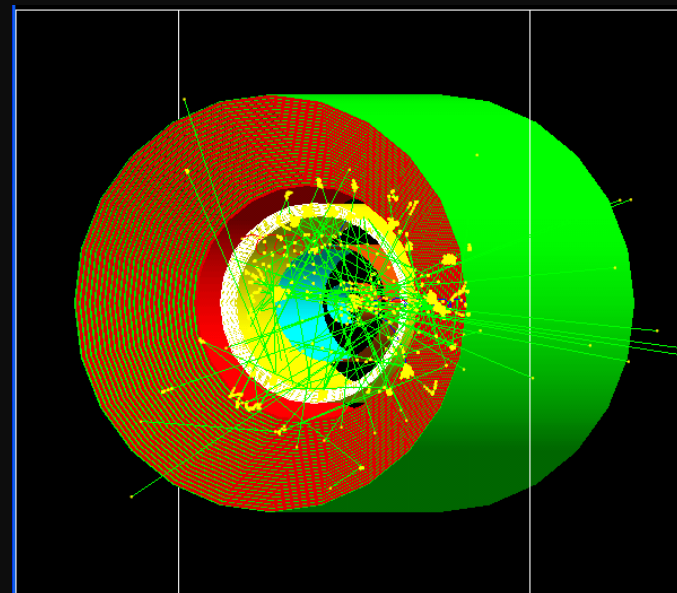
We are interested in  
feedback, suggestions,  
involvement.



Counts per 2.5 GeV bin in 50B AuAu Events



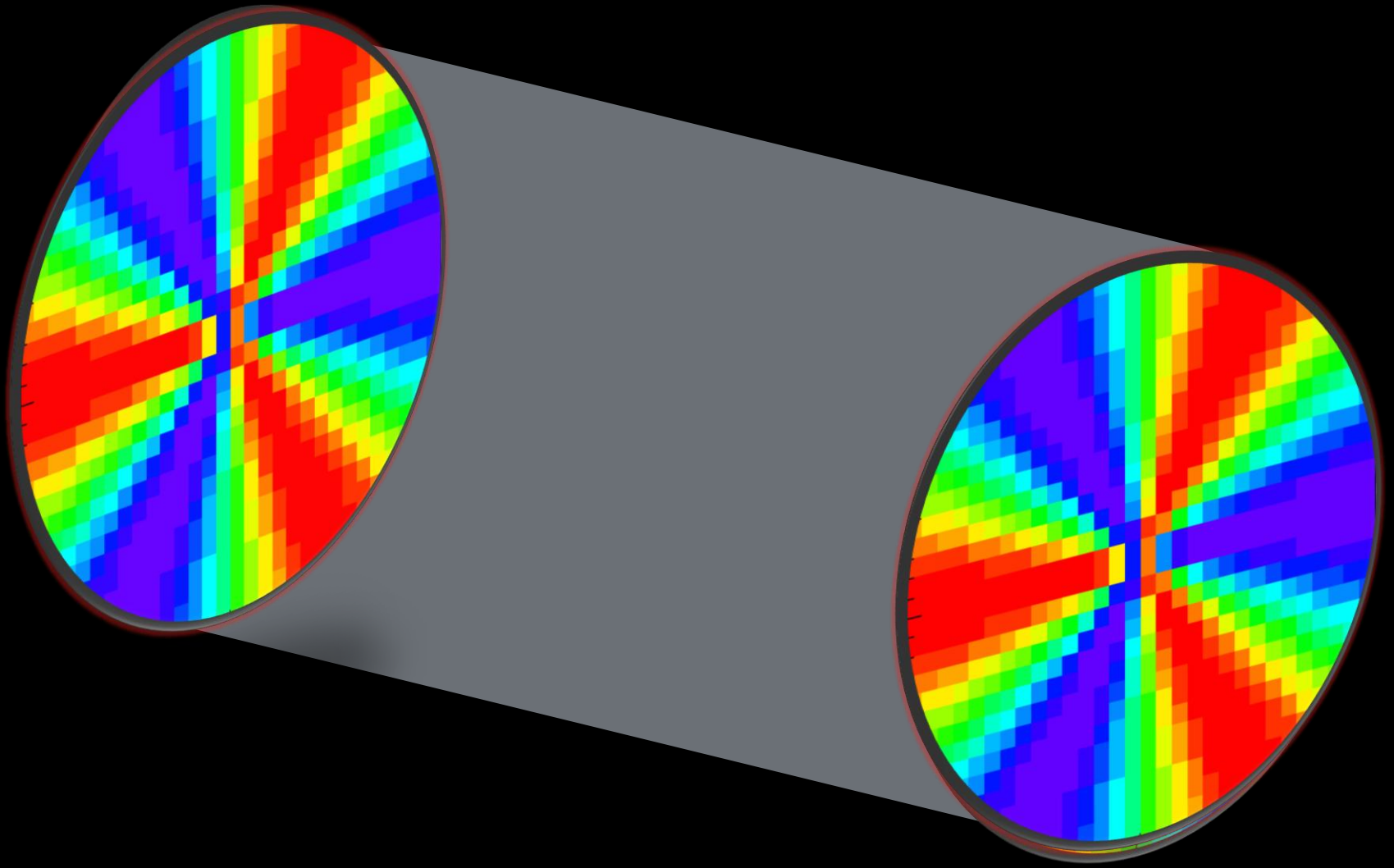
sPHENIX proposal  
for a hermetic  
EMCal/HCAL jet  
detector at RHIC



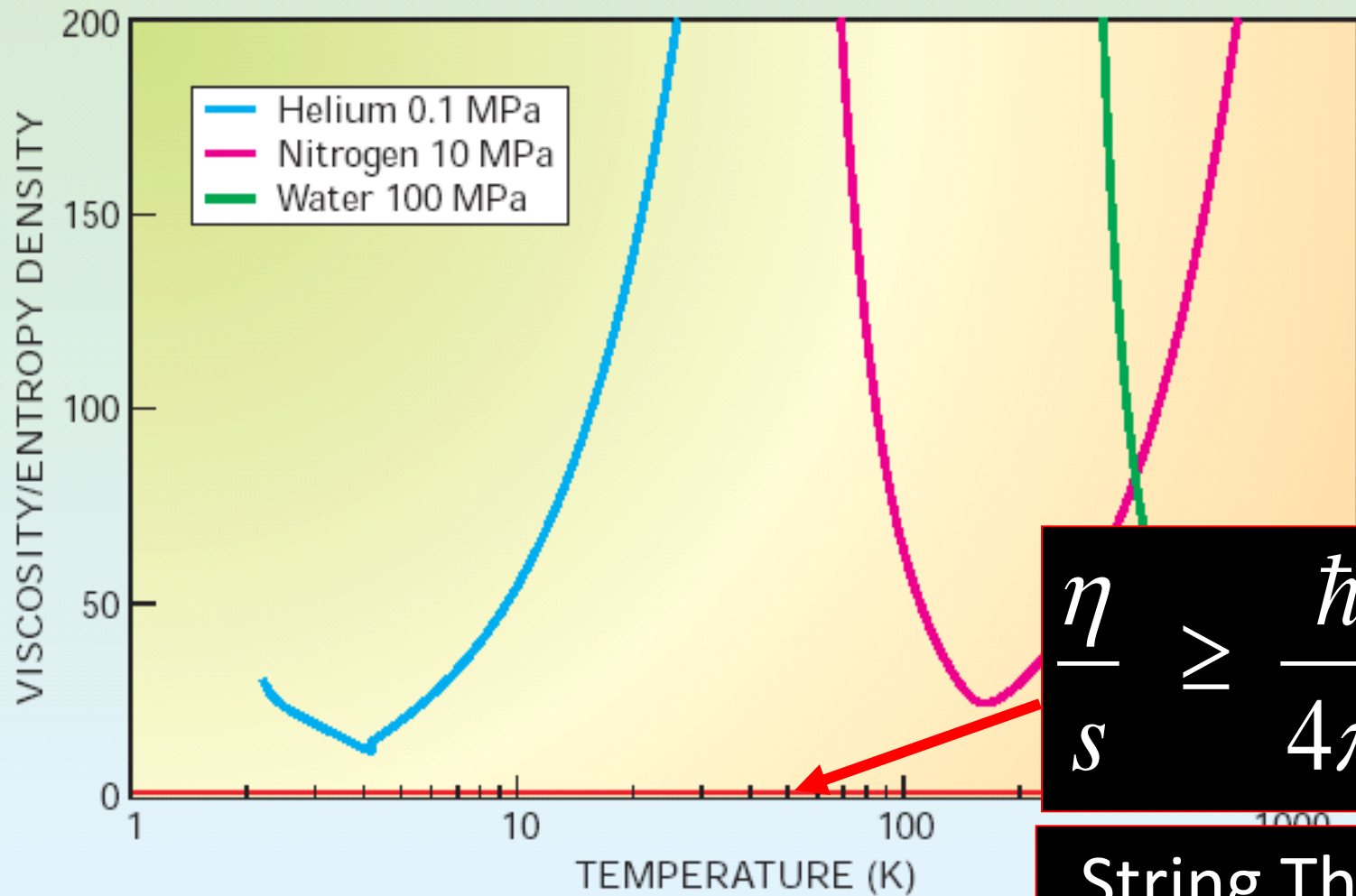
# Extras

# Measuring $v_3$

First, can detectors separated by  $\Delta\eta = 2$  or even  $\Delta\eta = 6$  measure event-by-event the 3<sup>rd</sup> order participant plane?



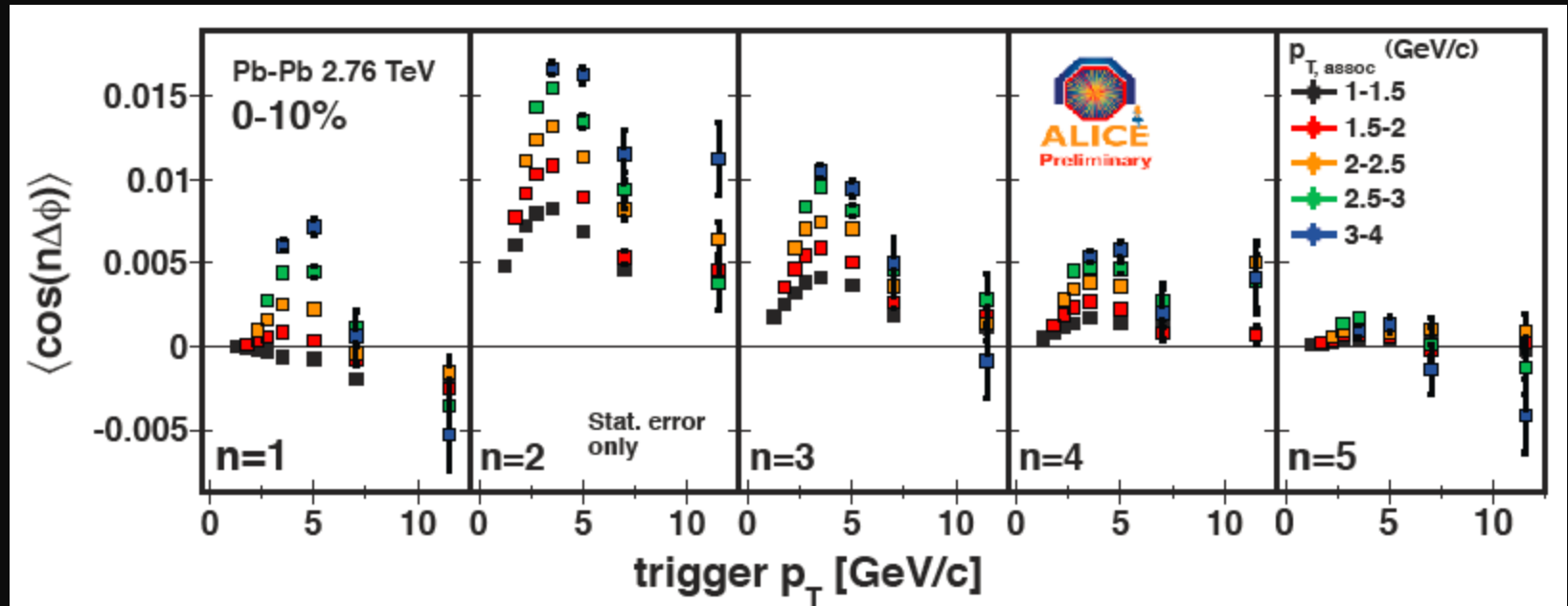
# Viscosity Roadmap



$$\frac{\eta}{s} \geq \frac{\hbar}{4\pi k_B}$$

String Theory  
Lowest Bound!

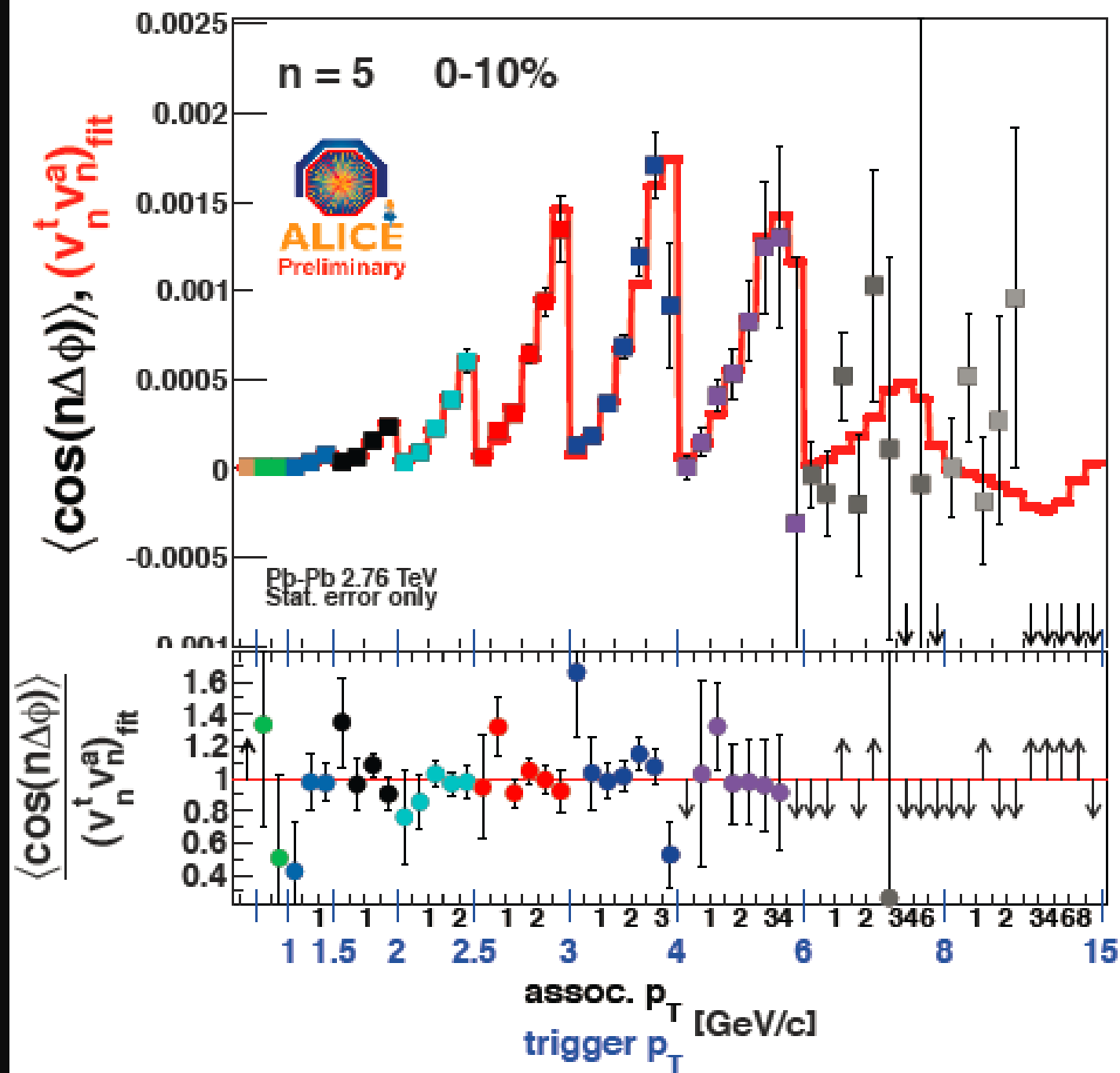
# Tests and model comparisons can be done on Coefficients



If entire 2-particle correlation from bulk flow then:

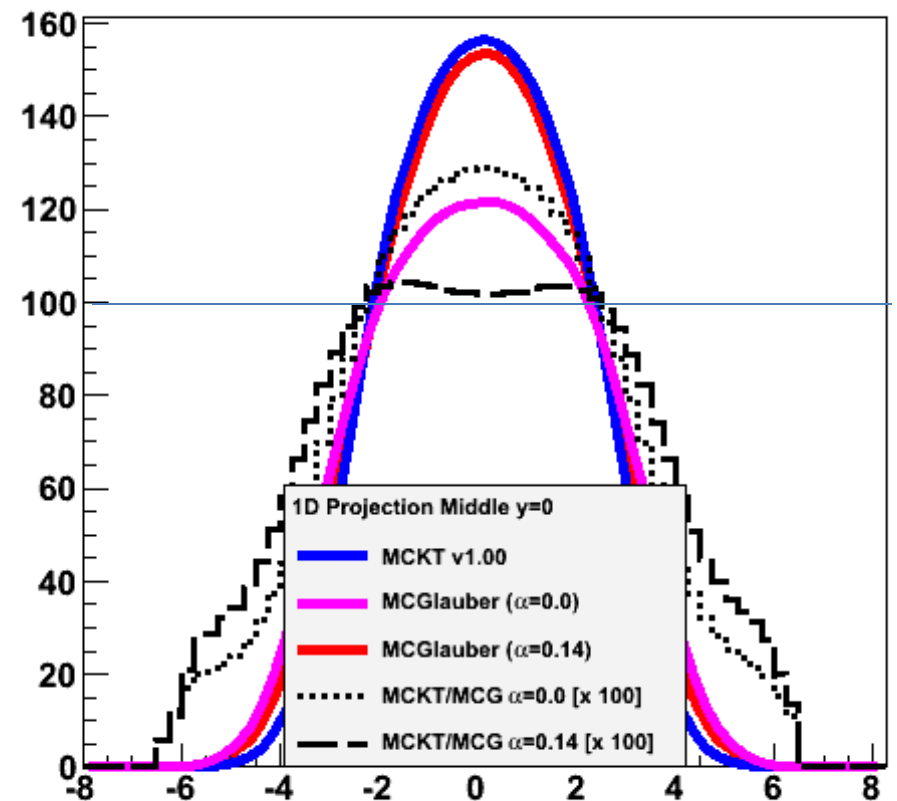
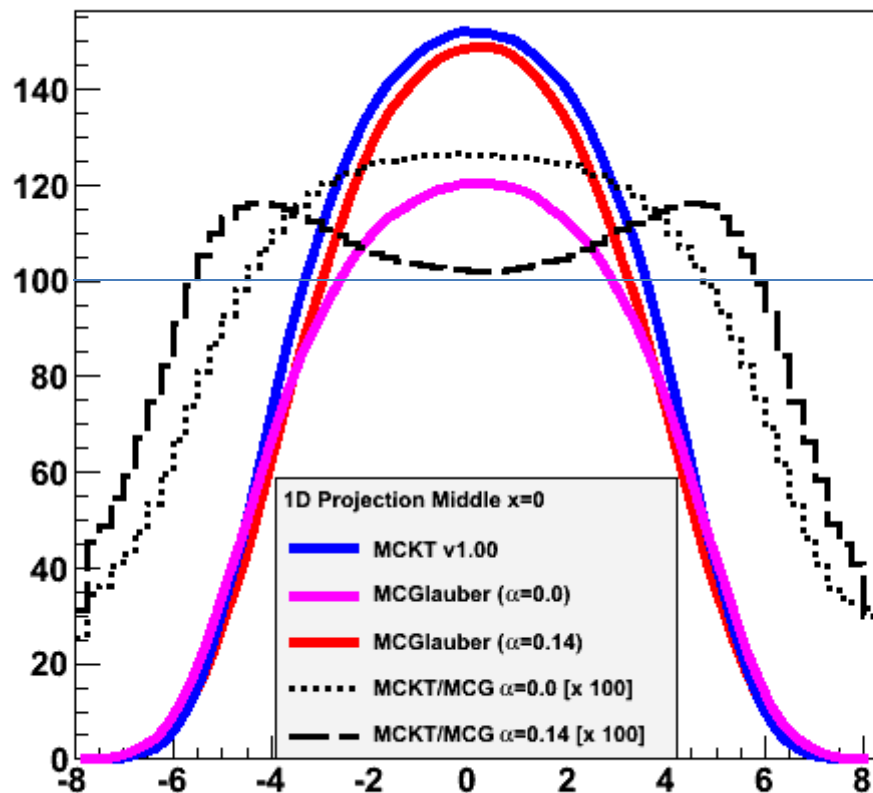
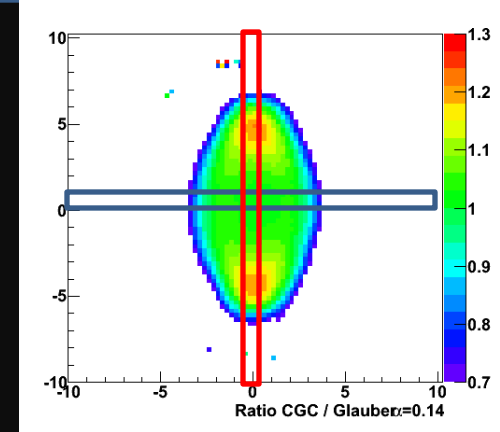
$$\text{Fourier Coefficient } v_{n\Delta}(p_{T1}, p_{T2}) = v_n(p_{T1}) \times v_n(p_{T2})$$

Deviations indicate important non-flow effects  
(jets and medium response for example)



# Ratio Density MCKT (Saturation) / Glauber

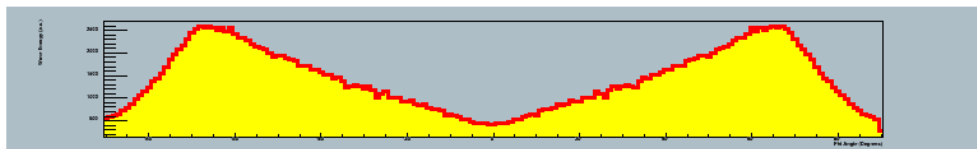
With the normalization, MCKT has a larger entropy density in the middle of the almond than both Glauber cases.





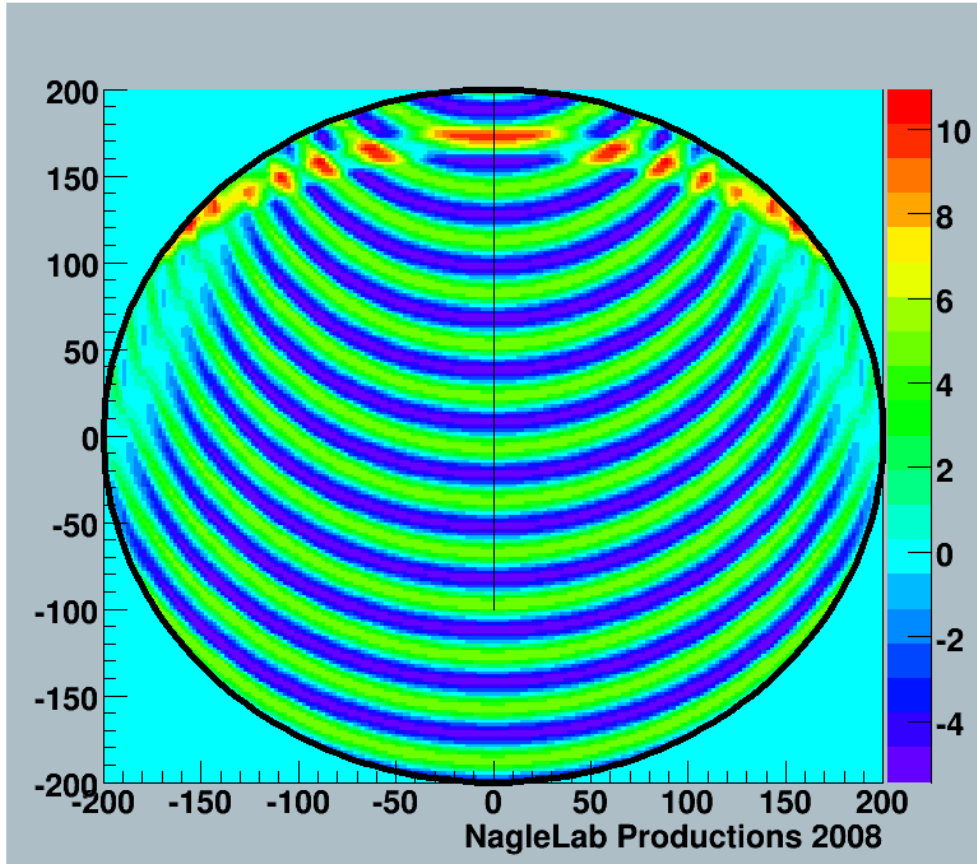
# Perfect Fluid Response

Super-sonic quark traversing medium  
results in shock wave response

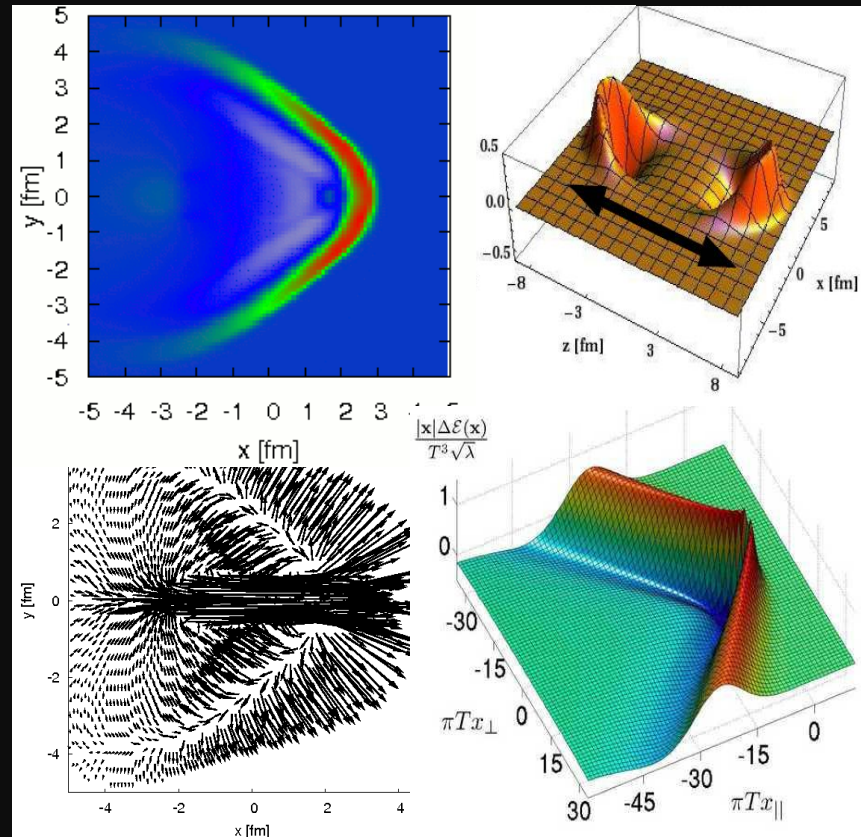


$\phi = -90^\circ$

$\phi = +90^\circ$



Observed in the data?



Lots of theory papers!