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# Angular Correlations at the LHC

$$C(\Delta\phi\Delta\eta) \equiv \frac{N_{\rm mixed}}{N_{\rm same}} \frac{{\rm d}^2 N_{\rm same}/{\rm d}\Delta\phi{\rm d}\Delta\eta}{{\rm d}^2 N_{\rm mixed}/{\rm d}\Delta\phi{\rm d}\Delta\eta}$$

Contributions to the two-particle ΔΦ, Δη angular correlation come from anisotropic flow, jets, resonances, HBT, etc



## Angular Correlations



For very peripheral collisions or when triggered with a high-p<sub>t</sub> charged particle the dominant contribution to two particle angular correlations is due to jet-correlations More central heavy-ion collisions look very very different!





#### measure anisotropic flow

$$\langle v_n \rangle = \langle \langle e^{in(\phi_1 - \Psi_n)} \rangle \rangle$$

 since the common symmetry plane cannot be measured event-by-event, we measure quantities which do not depend on it's orientation: multi-particle azimuthal correlations

$$\langle \langle e^{in(\phi_1 - \phi_2)} \rangle \rangle = \langle \langle e^{in(\phi_1 - \Psi_n - (\phi_2 - \Psi_n))} \rangle \rangle$$
  
=  $\langle \langle e^{in(\phi_1 - \Psi_n)} \rangle \langle e^{-in(\phi_2 - \Psi_n)} \rangle \rangle$   
=  $\langle v_n^2 \rangle$ 

 assuming that <u>only</u> correlations with the symmetry plane are present - not a very good assumption (jets, resonances, etc)!

# Can we isolate the flow?

- flow is a collective effect
  - flow "factorizes"
    - $<v_iv_j>/<v_j> = <v_iv_k>/<v_k>$ 
      - correlate particles separated in rapidity
      - correlate particles separated in pt
  - multi-particle correlation
    - build cumulants

#### **Collective Motion**

A+A independent # of p+p





P+P

A+A collective behavior

tests if we are dealing with a common symmetry plane!

# multi-particle correlations

 for detectors with uniform acceptance 2<sup>nd</sup> and 4<sup>th</sup> cumulant are given by:

> Borghini, Dihn and Ollitrault, PRC 64, 054901 (2001)

$$c_n\{2\} \equiv \left\langle \left\langle e^{in(\phi_1 - \phi_2)} \right\rangle \right\rangle = v_n^2 + \delta_2$$

$$c_n\{4\} \equiv \left\langle \left\langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \right\rangle \right\rangle - 2 \left\langle \left\langle e^{in(\phi_1 - \phi_2)} \right\rangle \right\rangle^2$$

$$= v_n^4 + 4v_n^2 \delta_2 + 2\delta_2^2 - 2(v_n^2 + \delta_2)^2$$

$$= -v_n^4$$

we got rid of two particle nonflow correlations! we can remove nonflow order by order

## v<sub>2</sub> from cumulants



cumulants show behavior as expected when correlations are dominated by collective flow

# What about p-p?



the 2 and 4-particle correlations decrease with increasing number of particles produced which is typical behavior of correlations involving only few particles the 4-particle cumulant,  $QC{4}$ , even has the wrong sign compared to true flow!

# What about p-p?



remember  $QC{2} = v^2, QC{4} = -v^4$ 

models like PYTHIA and PHOJET have the correct sign for the correlations and capture the general trends observed in p-p

no evidence for elliptic flow in this multiplicity range

The Perfect Liquid



The system produced at the LHC behaves as a very low viscosity fluid (a perfect fluid), constraints dependence of  $\eta$ /s versus temperature

# Flow Analysis Methods

flow analysis methods have different sensitivity to nonflow and fluctuations

I focus on the cumulants

Borghini, Dihn and Ollitrault, PRC 64, 054901 (2001) Bilandzic, Snellings and Voloshin, PRC 83, 044913 (2011)

$$v_n^2 \{2\} = \bar{v}_n^2 + \sigma_v^2 + \delta$$
$$v_n^2 \{4\} = \bar{v}_n^2 - \sigma_v^2$$
$$v_n^2 \{6\} = \bar{v}_n^2 - \sigma_v^2$$
$$v_n^2 \{8\} = \bar{v}_n^2 - \sigma_v^2$$

excellent opportunity to study flow fluctuations and from these get a handle on initial conditions!

#### Flow Fluctuations

when (2-particle) nonflow is corrected for or negligible!

in limit of "small" (not necessarily Gaussian) fluctuations

 $v_n^2 \{2\} = \bar{v}_n^2 + \sigma_v^2$  $v_n^2 \{4\} = \bar{v}_n^2 - \sigma_v^2$  $v_n^2 \{2\} + v_n^2 \{4\} = 2\bar{v}_n^2$  $v_n^2 \{2\} - v_n^2 \{4\} = 2\sigma_v^2$ 

$$v_n\{4\} = 0$$
$$v_n\{2\} = \frac{2}{\sqrt{\pi}}\bar{v}_n$$

## v2 versus centrality in ALICE



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Two particle  $v_2$  estimates depend on  $\Delta \eta$ Higher order cumulant  $v_2$ estimates are consistent within uncertainties Two particle  $v_2$  estimates are corrected for nonflow based on HIJING The estimated nonflow correction for  $\Delta \eta > 1$  is included in the systematic uncertainty

#### v<sub>2</sub> Fluctuations



tune initial conditions? what are the other contributions? need to compare to full model calculations!

### v<sub>2</sub> as function of p<sub>t</sub>

![](_page_15_Figure_1.jpeg)

Elliptic flow as function of transverse momentum does not change much from RHIC to LHC energies, can we understand that?

![](_page_16_Figure_0.jpeg)

## v<sub>2</sub> for identified particles

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![](_page_17_Figure_2.jpeg)

at small  $(m_t-m_0)/n_q$  the scaling in the data resemble the scaling as observed in hydrodynamics

![](_page_17_Figure_4.jpeg)

## Initial conditions and vn

G. Qin, H. Petersen, S. Bass, and B. Muller

![](_page_18_Figure_2.jpeg)

![](_page_18_Figure_3.jpeg)

$$\frac{2\pi}{N}\frac{dN}{d\phi} = 1 + \sum_{n=2,4,6,\dots}^{\infty} 2v_n \cos n(\phi - \Psi_R) \qquad \qquad \frac{2\pi}{N}\frac{dN}{d\phi} = 1 + \sum_{n=1}^{\infty} 2v_n \cos n(\phi - \Psi_n)$$

initial spatial geometry not a smooth almond (for which all odd harmonics are zero due to reflection symmetry) may give rise to higher odd harmonics which give additional independent constraint on the initial conditions!!!

## Shear Viscosity

#### Music, Sangyong Jeon

![](_page_19_Figure_2.jpeg)

#### initial conditions

#### ideal hydro $\eta/s=0$ viscous hydro $\eta/s=0.16$

![](_page_19_Figure_5.jpeg)

Larger η/s clearly smoothes the distributions and suppresses the higher harmonics (e.g. v<sub>3</sub>)

Hydro: Alver, Gombeaud, Luzum & Ollitrault, Phys. Rev. C82 (2010) 20

![](_page_20_Figure_0.jpeg)

![](_page_20_Figure_1.jpeg)

u<sub>1</sub> > u<sub>2</sub> > u<sub>3</sub> shear viscosity will make them equal and destroy the elliptic flow v<sub>2</sub> higher harmonics represent smaller differences which get destroyed more easily, and which, if measurable, makes them more sensitive probes to η/s

# Triangular Flow

 $\dots v_3$  Glauber  $\eta$ /s=0.08 We observe significant  $v_3$  which 0.1 .....v<sub>3</sub> CGC η/s=0.16 compared to  $v_2$  has a different centrality dependence The centrality dependence and 0.05 magnitude are similar to predictions for MC Glauber with  $\eta/s=0.08$  but above MC-⊓¥ KLN CGC with  $\eta/s=0.16$ 0 20 10 0

![](_page_21_Figure_2.jpeg)

ALICE Collaboration, arXiv:1105.3865 accepted for PRL

The  $v_3$  with respect to the reaction plane determined in the ZDC and with the  $v_2$  participant plane is consistent with zero as expected if  $v_3$  is due to fluctuations of the initial eccentricity

The  $v_3{2}$  is about two times larger than  $v_3{4}$  which is also consistent with expectations based on initial eccentricity fluctuations

# Triangular Flow

![](_page_22_Figure_1.jpeg)

The behavior of  $v_3$  as function of  $p_t$  for pions, kaons and protons shows the same features as observed for  $v_2$ (the mass splitting, the crossing of the pions with protons at intermediate  $p_t$ )

# Geometry and Harmonics

![](_page_23_Figure_1.jpeg)

G-L Ma and X-N Wang, arXiv:1011.5249v2

For central collisions at intermediate  $p_t$  the higher harmonics  $v_3$  and  $v_4$  cross  $v_2$  and become the dominant harmonics

For more central collisions this occurs already at lower pt

![](_page_23_Figure_5.jpeg)

#### Geometry

![](_page_24_Figure_1.jpeg)

$$C(\Delta\phi) \equiv \frac{N_{\text{mixed}}}{N_{\text{same}}} \frac{\mathrm{d}N_{\text{same}}/\mathrm{d}\Delta\phi}{\mathrm{d}N_{\text{mixed}}/\mathrm{d}\Delta\phi}$$

![](_page_24_Figure_3.jpeg)

We observe a doubly-peaked structure in the azimuthal correlation function opposite to the trigger particle before the subtraction of v<sub>2</sub>

The red line shows the sum of the measured anisotropic flow Fourier coefficients. The flow coefficients give a natural explanation of the observed correlation structure known as the Mach cone and ridge

## Conclusions

- Anisotropic flow measurements provided strong constraints on the properties of hot and dense matter produced at RHIC and LHC energies and have lead to the new paradigm of the QGP as the so called perfect liquid
  - At the LHC we observe even stronger flow than at RHIC which is expected for almost perfect fluid behavior
- The first measurements of v<sub>3</sub>, v<sub>4</sub> and v<sub>5</sub> have recently been made at RHIC and the LHC and indicate that these flow coefficients behave as expected from fluctuations of the initial spatial eccentricity (geometry!) and a created system which has a small  $\eta/s$ 
  - provide new strong experimental constraints on η/s and initial conditions

# Thanks