

Observational cosmology with X-ray luminous galaxy clusters

First Lecture: Cosmology, Clusters and Gas Mass Fraction

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Nordic Winter School, Gausdal

General outline for the two lectures

- Overview of current cosmological experiments and results
- Detailed examples of data analyses
- Paradigms and approaches in present-day cosmology including concrete examples

Observational cosmology

Underlying standpoints of data analysis to take home from the lectures:

- Deep understanding of **astrophysical processes** and **objects**
- Careful **design** of observations
- Justify and continuously revised **assumptions**
- Account properly for **covariances** between parameters, instrumental and astrophysical **systematic uncertainties** and **biases**
- **Simultaneous** fits of all the relevant **astrophysical** and **cosmological** parameters

Structure of the lectures

First Lecture:

- Introduction to cosmological analyses
- Expansion history measurements and cosmological results
- Detailed data analysis: gas mass fraction experiment

Second Lecture:

- Cosmological models and modeling
- Growth history measurements and cosmological results
- Detailed data analysis: cluster abundance experiment

Recent discoveries and current results

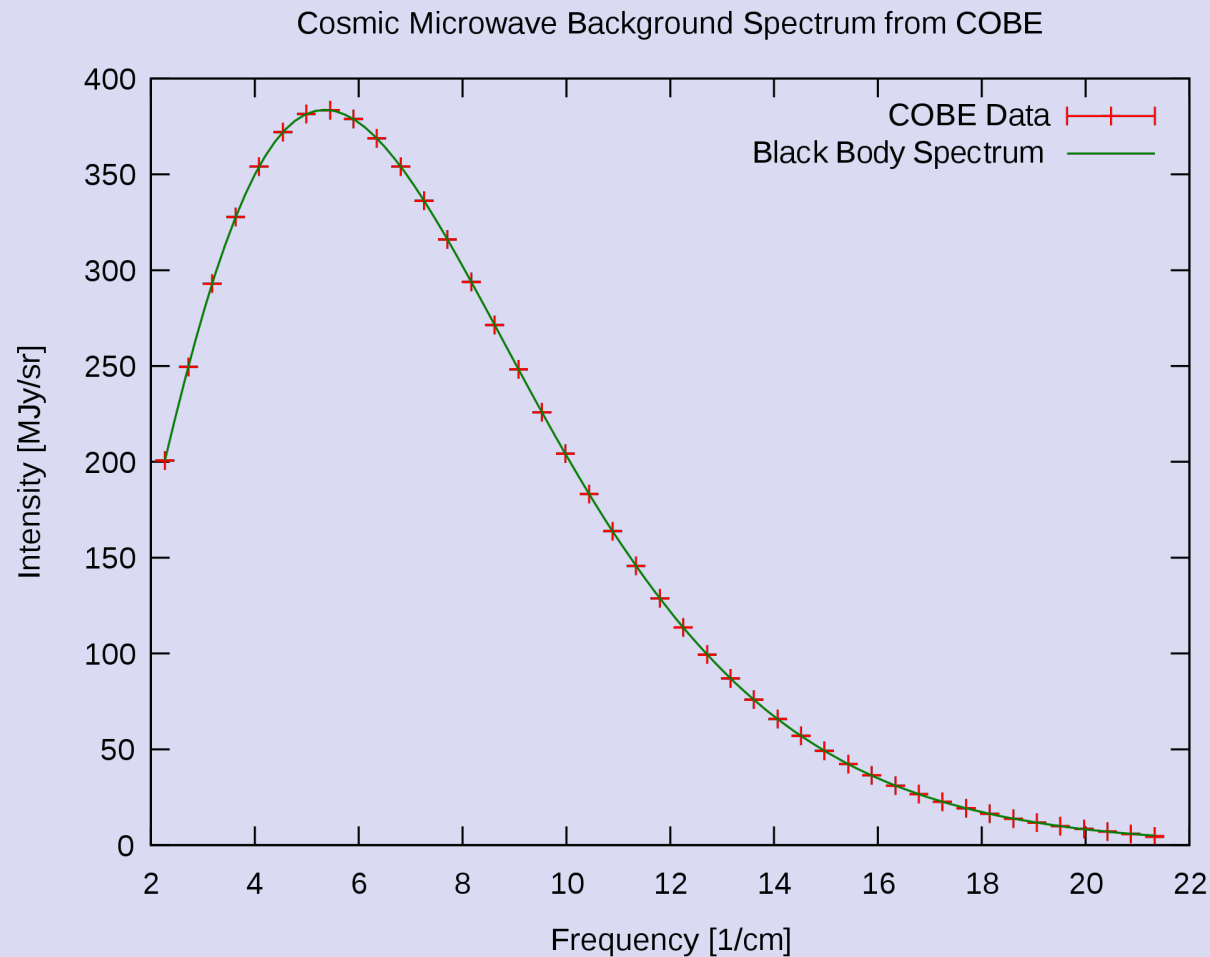
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Cosmic Microwave Background (CMB)

Nobel Prize in Physics 1978: Arno Penzias & Robert Wilson (CMB discovery in 1965)
[Pyotr Kapitsa (Low-temperature physics)]

Nobel Prize in Physics 2006: John Mather & George Smoot (CMB blackbody and anisotropy)



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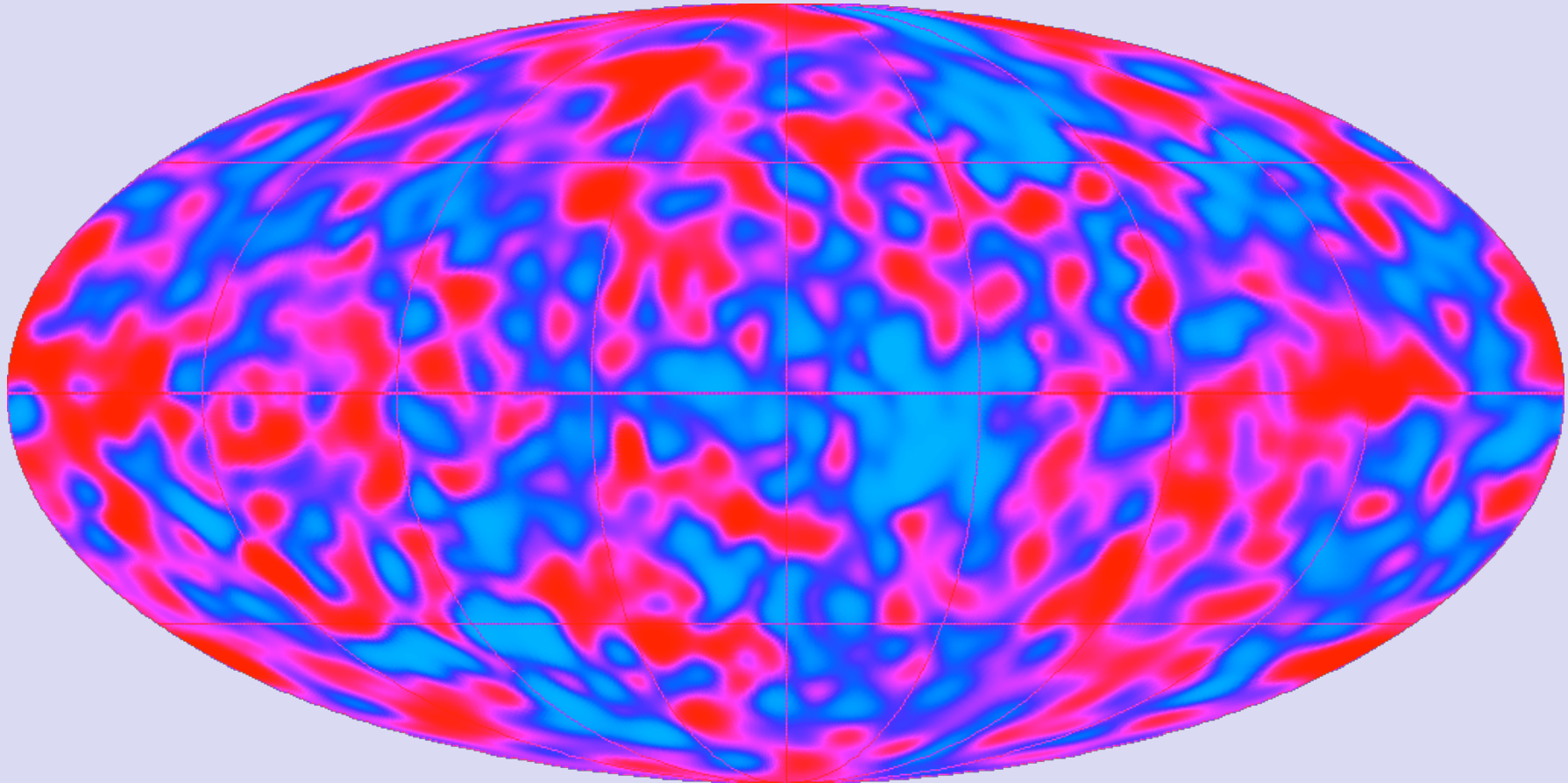
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From NASA's COBE satellite

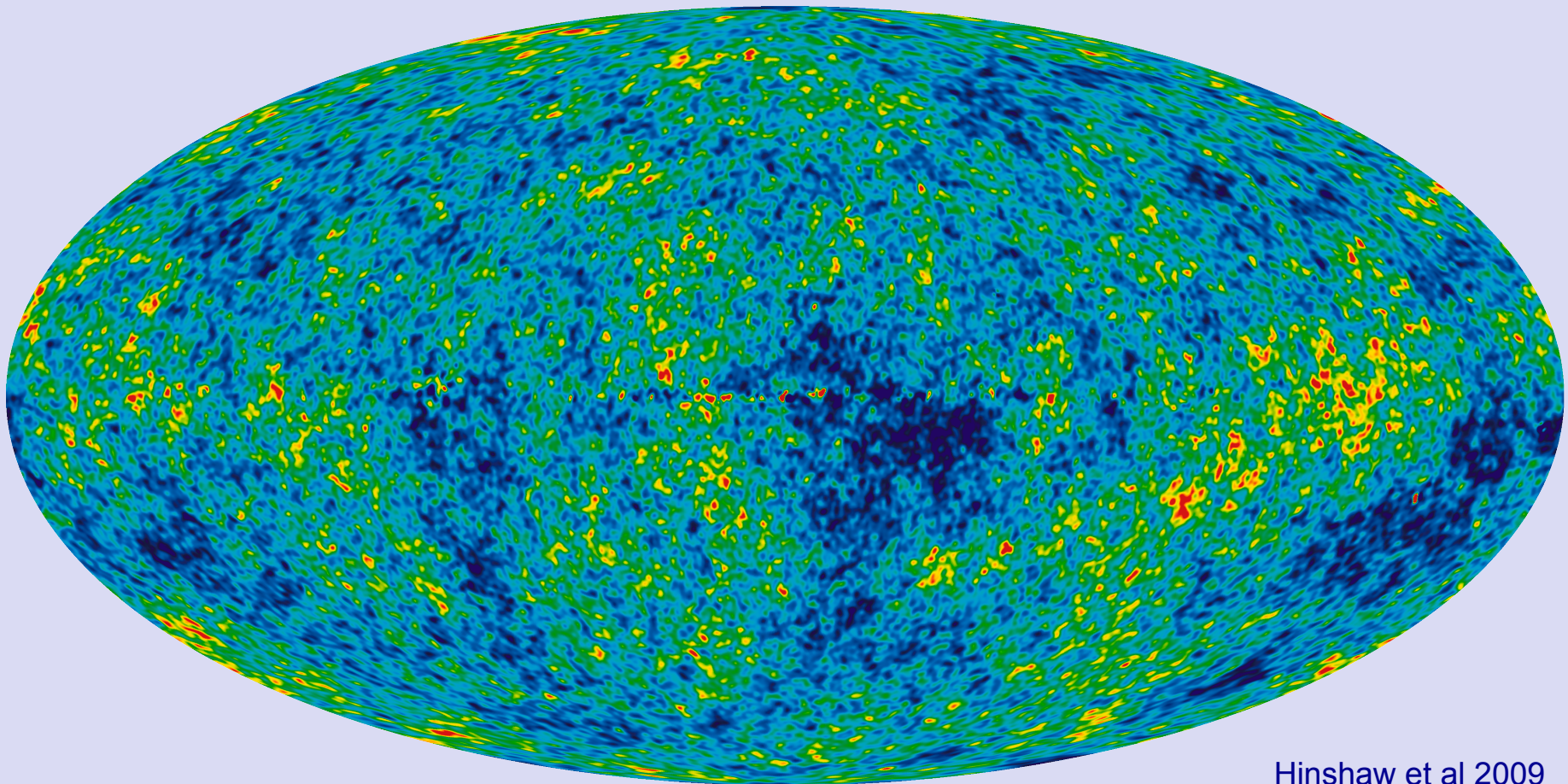


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Cosmic Microwave Background (CMB)

Measurements from current NASA's WMAP satellite (five-years)



Hinshaw et al 2009

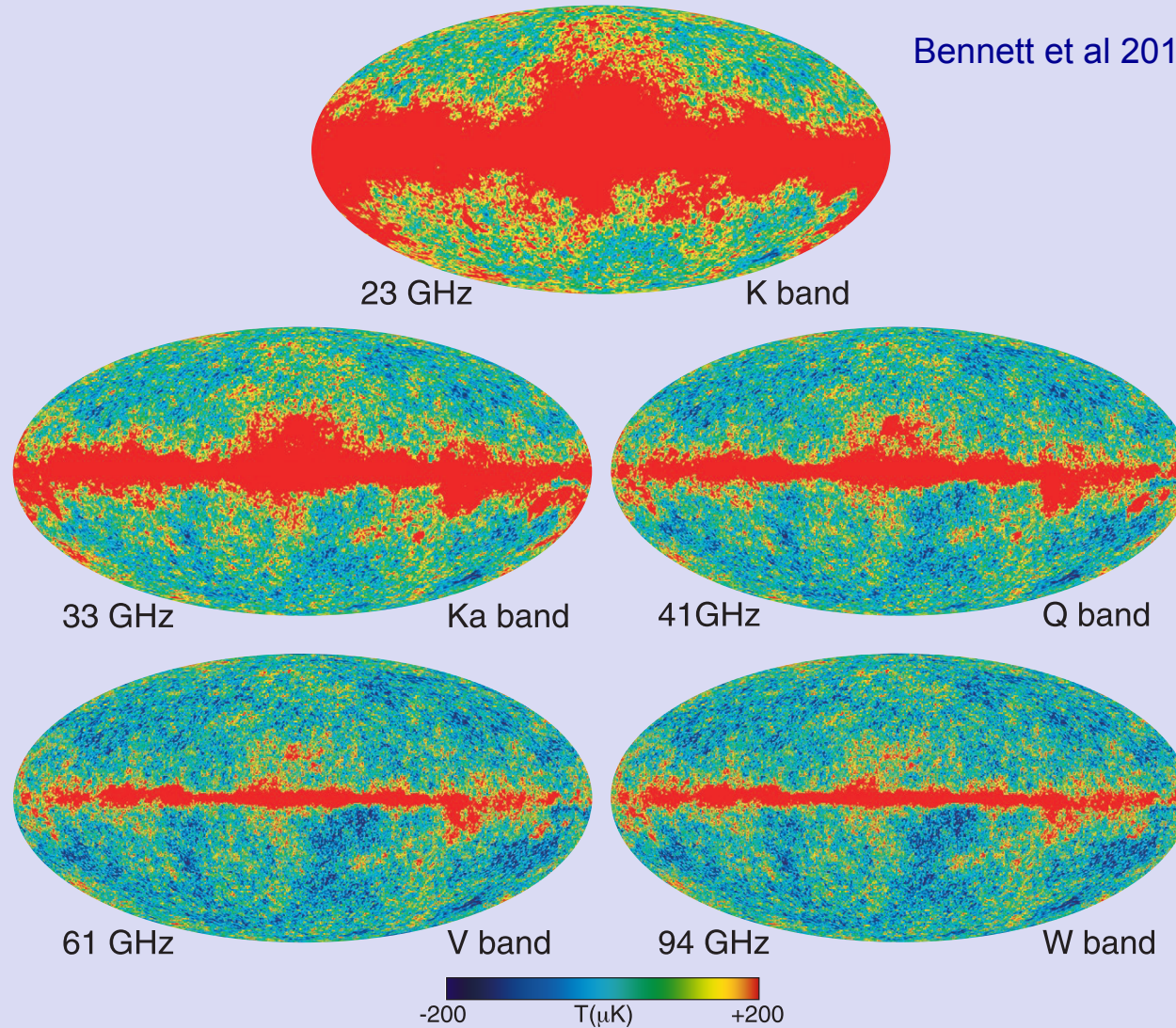
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Cosmic Microwave Background (CMB)

Final WMAP (nine-years) measurements

Bennett et al 2012

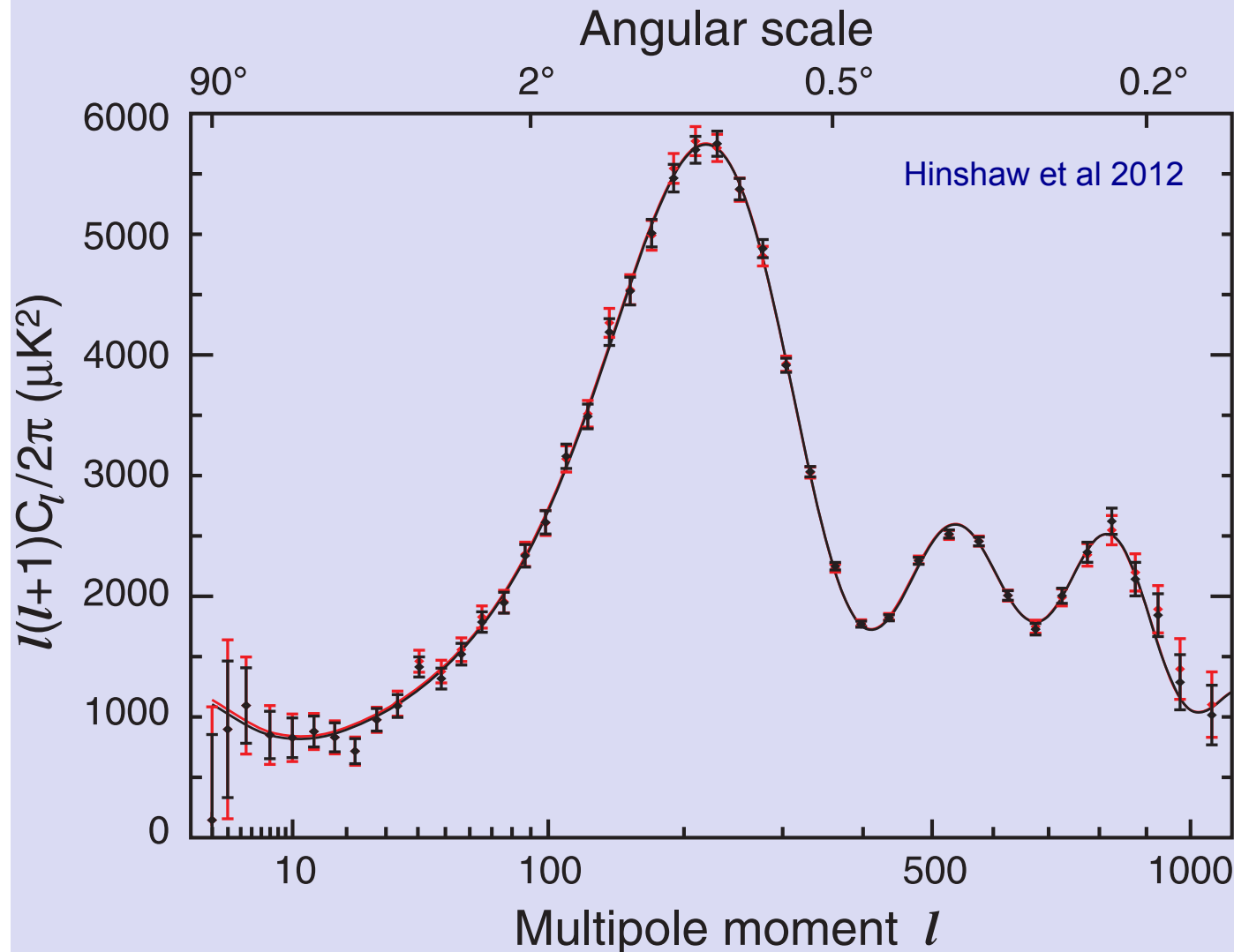


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Cosmic Microwave Background (CMB)

Final WMAP (nine-years) measurements



Red symbols (binned data) and line (best fit):
Using the previous WMAP MASTER-based power spectrum

Black symbols (binned data) and line (best fit):
Using the new WMAP C-1-weighted power spectrum

Up to 5% difference in the vicinity of $l \sim 50$. This affects mainly n_s .

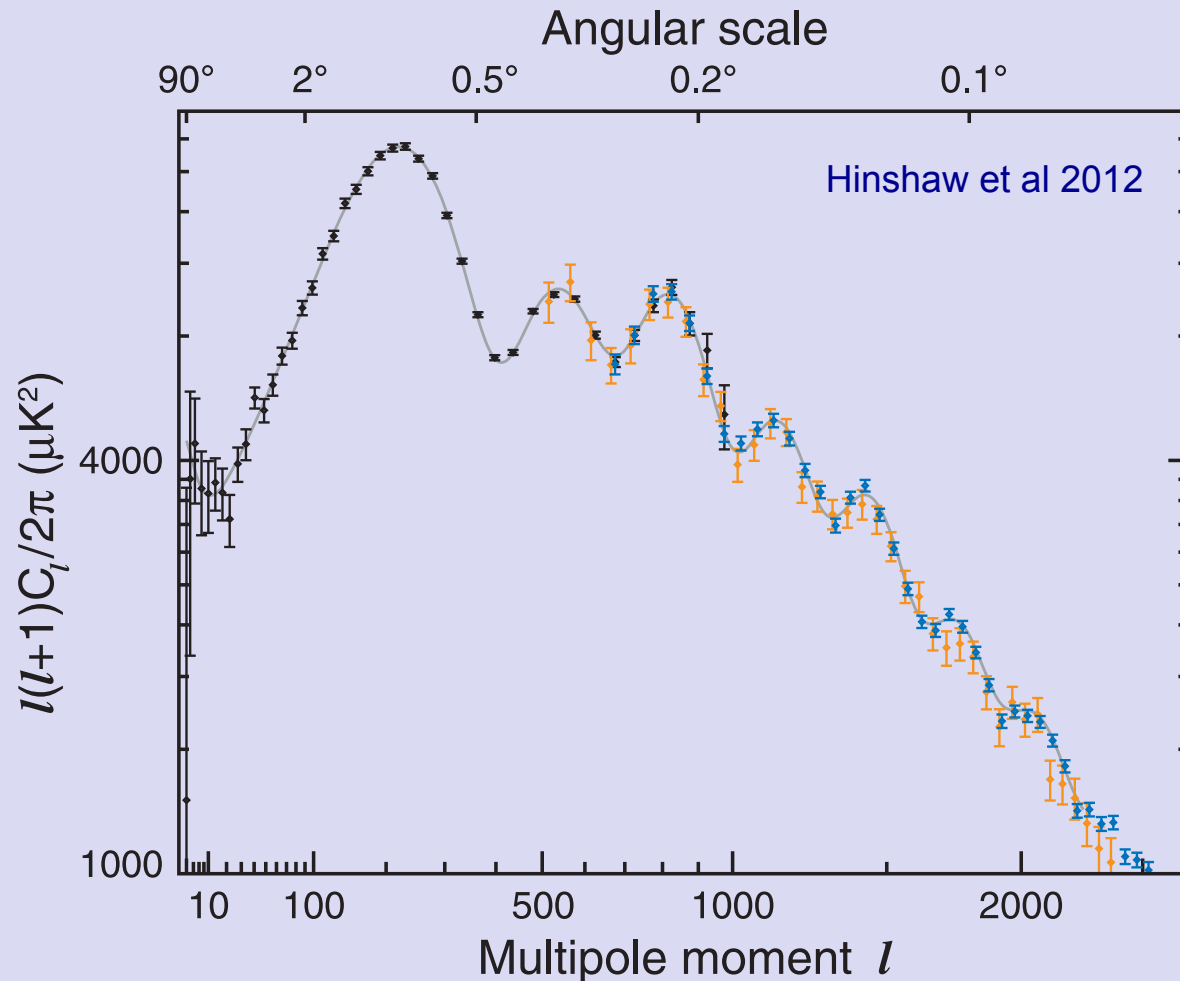
Overall, WMAP9 wrt WMAP7 improve the parameters an average of $\sim 10\%$, with almost 20% for $\Omega_c h^2$ and Ω_Λ , with about 10% from the new power spectrum.

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Cosmic Microwave Background (CMB)

Final WMAP (nine-years) measurements



Next: results from ESA's Planck satellite are coming soon this year...

Discovery of cosmic acceleration

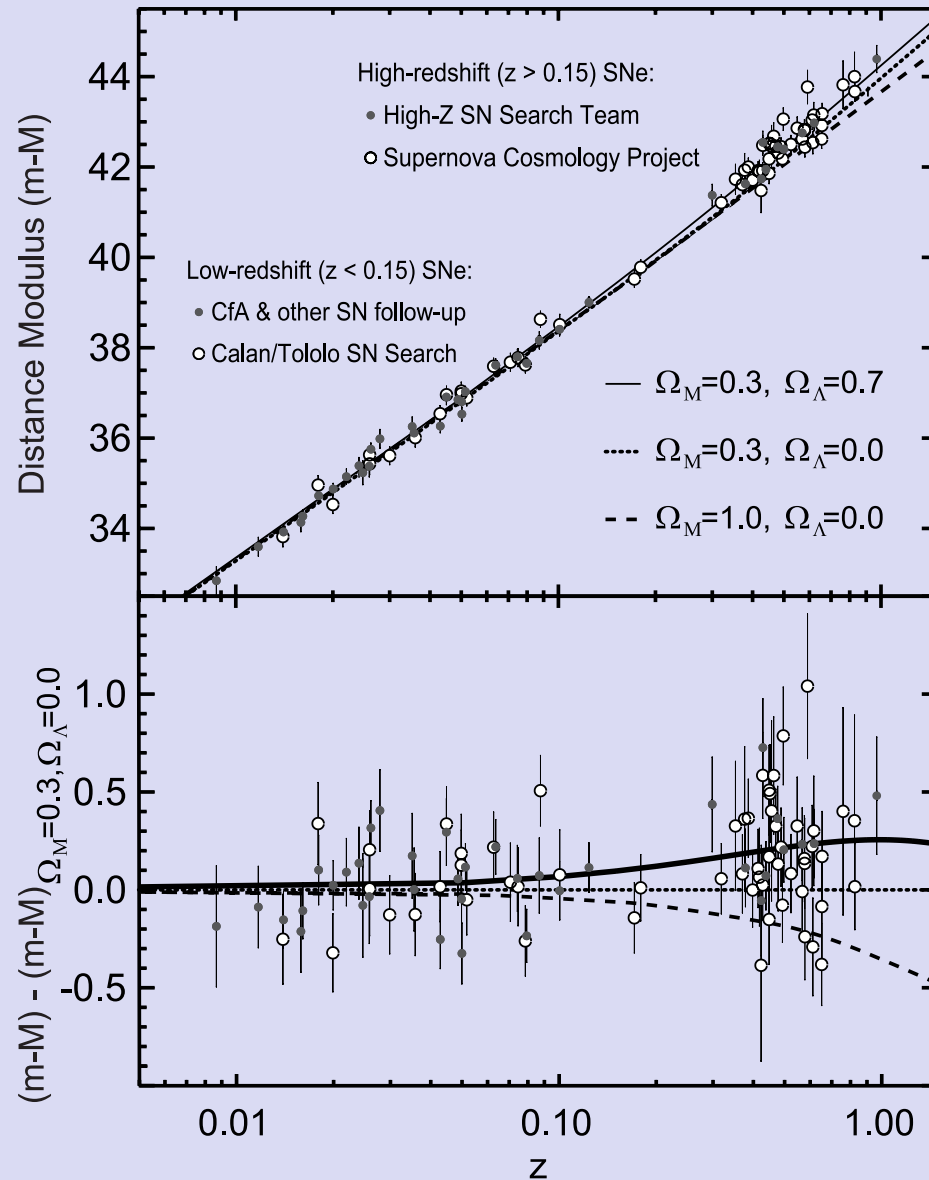


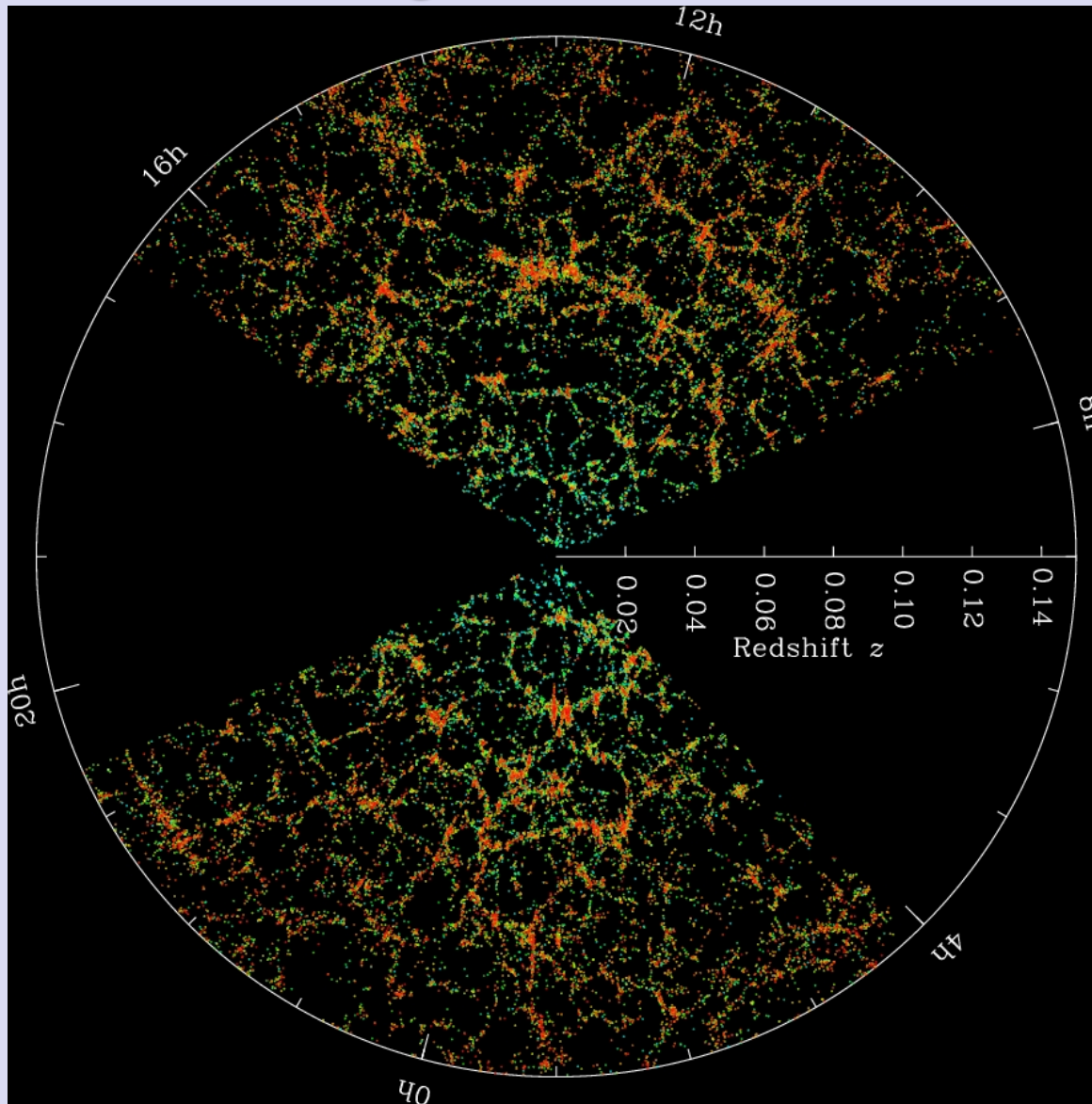
Figure from the dark energy review of
Frieman, Turner & Huterer, 2008,
ARA&A., 46, 385

-High-Z SN Search Team (HZT):
Riess et al 1998

-Supernova Cosmology Project (SCP):
Perlmutter et al et al. 1999

Nobel Prize Award in Physics 2011:
Saul Perlmutter (SCP)
Brian Schmidt & Adam Riess (HZT)

Large scale distribution of galaxies



Sloan Digital Sky Survey
(from the SDSS website)

Slice of a 3D map of galaxies

Galaxies are colored according
to the ages of their stars: the
redder, the more strongly cluster
since are made of older stars.

Outer circle: about two billion
light years old

>930000 galaxies

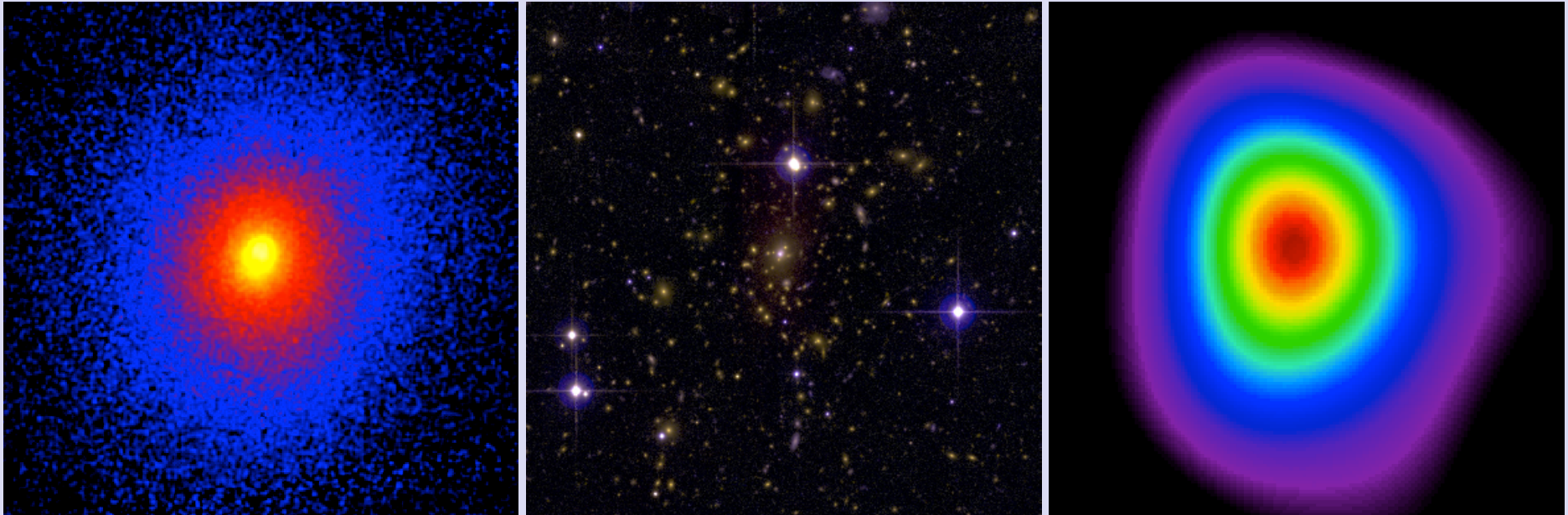
Cluster cosmology

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Cluster cosmology

Figure from Allen, Evrard & Mantz 11 (credits X-ray/Mantz; Optical/von der Linden et al; SZ/Marrone)



X-ray

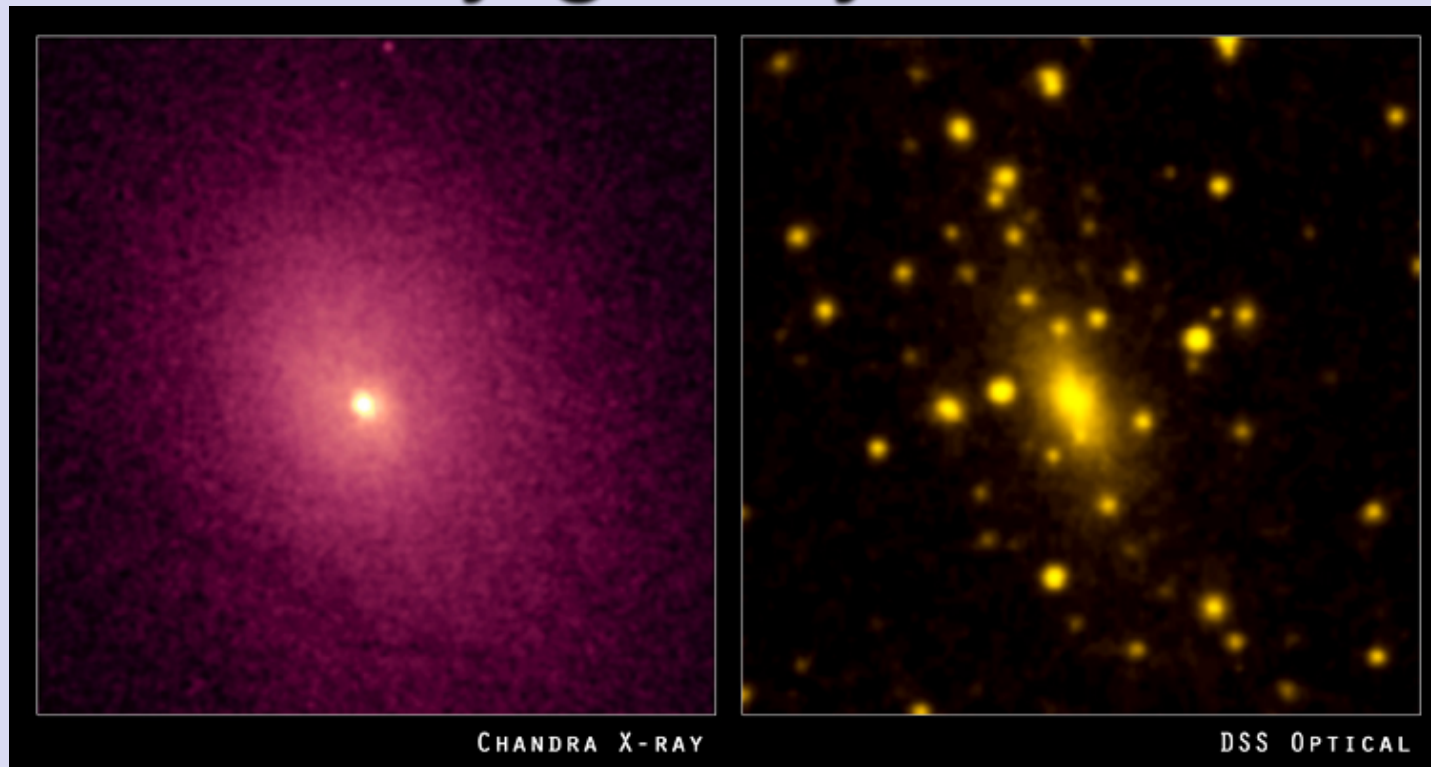
Optical

Millimeter (SZ)

- Images of galaxy cluster Abell 1835 in different wavelength
- Cosmology with galaxy clusters using X-ray observations:
 - Gas mass fraction
 - Abundance of clusters and their observable-mass relations

Cluster cosmology review: Allen, Evrard & Mantz, 2011, ARA&A, 49, 409

X-ray galaxy cluster



Galaxy cluster Abell 2029:

Thousands of galaxies optical (right panel)

Hot (multimillion Kelvin degrees) gas (left panel).

Dark matter (only gravitational interaction) $>10^{15}$ solar masses.

In 1933 Fritz Zwicky proposed “missing matter” (dark matter) in clusters of galaxies
(Note: Central enormous elliptically shaped galaxy)

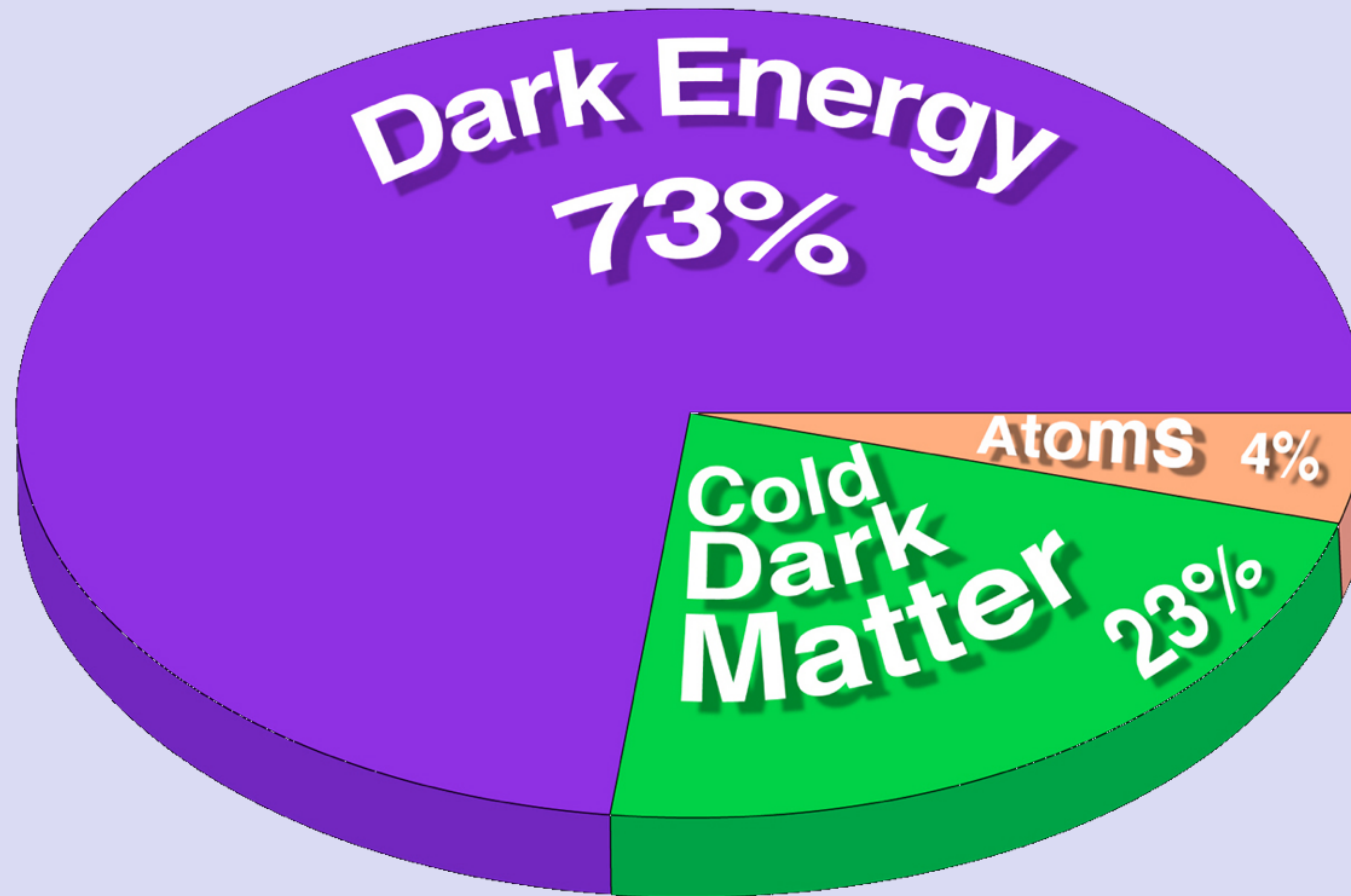
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Basic cosmology

Peebles; Luchin & Matarrese; Peacock;
Dodelson; Weinberg ; Mukhanov; etc.

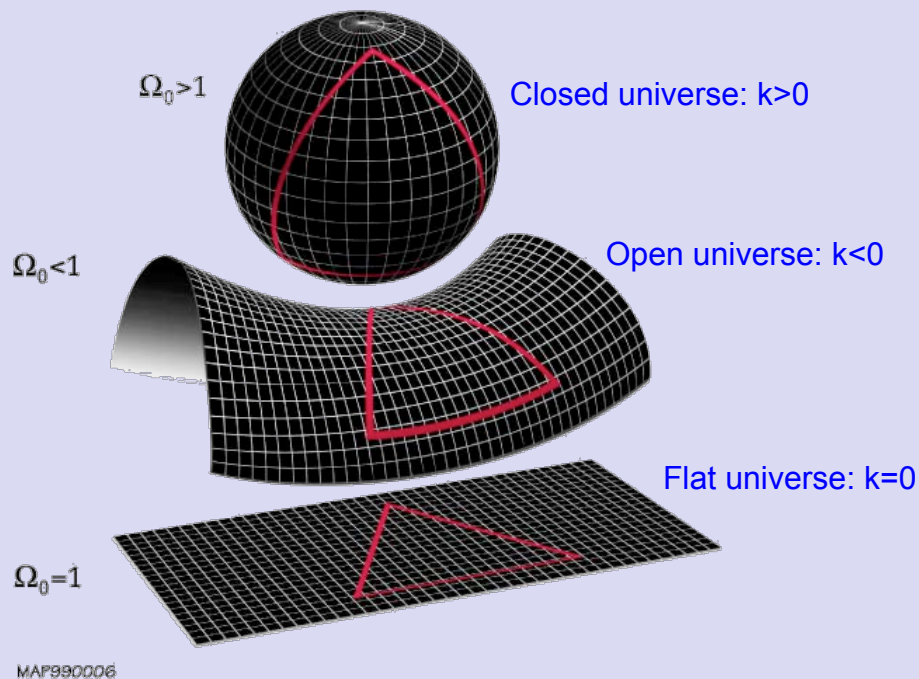
Cosmic pie: energy density of the Universe



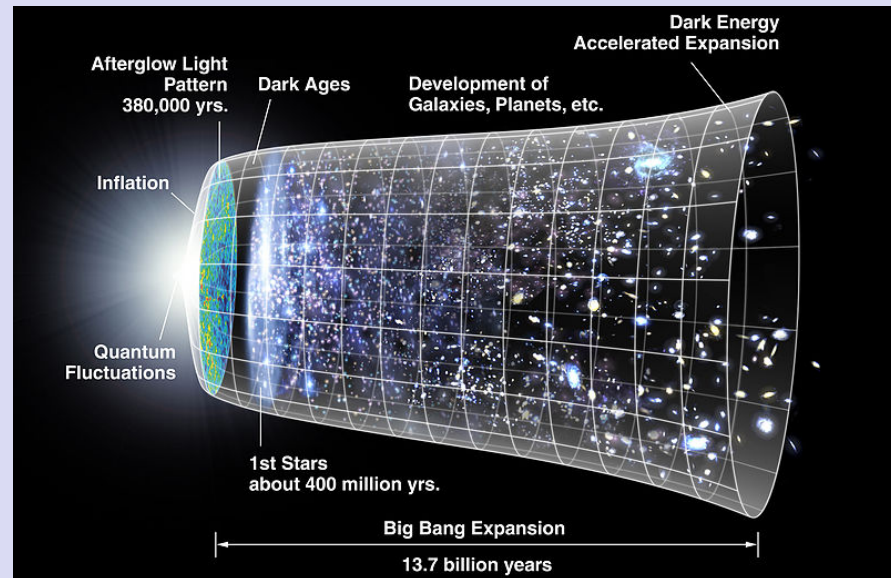
Space-time metric of the Universe

Friedmann-Lemaitre-Robertson-Walker (FLRW) metric: for an **isotropic** and **homogenous** space-time (**current data** indicate that at large scale these assumptions are nearly valid)

$$ds^2 = dt^2 - a^2(t) \left[dr^2 / (1 - kr^2) + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right]$$



Expansion history $a(t)$



Gravity and energy density

Einstein's equations

$$G^{\mu\nu} + \Lambda g^{\mu\nu} = 8\pi T^{\mu\nu}$$

$$c = G = 1$$

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} R g^{\mu\nu}$$

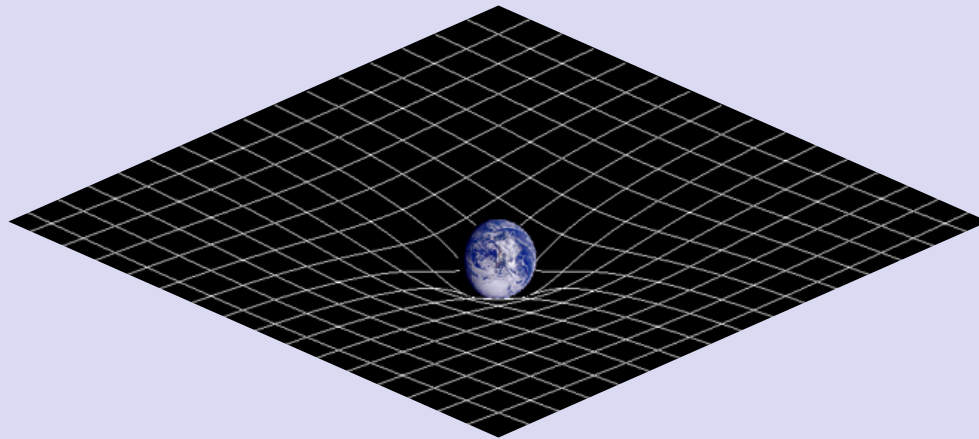
$R^{\mu\nu}$ Ricci tensor

R Ricci scalar

$G^{\mu\nu}$ metric tensor

Interaction between gravity and energy-matter

John Wheeler: Matter tells space how to curve and space tells matter how to move



Useful reference for perturbation theory: Ma & Bertschinger 1995, ApJ, 455, 7

Cosmic energy content and expansion

Example of stress-energy tensor:

$$T^{\mu\nu} = \begin{pmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

Fluid in a thermodynamic equilibrium

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G\rho}{3} - \frac{k}{a^2} + \frac{\Lambda}{3}$$
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3}$$

Friedmann equations: key equations in cosmology; Einstein field equations for the FRWL metric

$$\Lambda = 8\pi G\rho_{\text{VAC}} = -8\pi Gp_{\text{VAC}}$$

Energy density of the vacuum

Dark energy review: Frieman, Turner & Huterer 2008, ARA&A., 46, 385

Cosmic energy content and expansion

$$H^2 = \frac{8\pi G}{3} \rho - \frac{k}{a^2}$$

Friedmann equation

ρ sum of the energy densities of matter, dark energy, radiation

$$w = \frac{p_{de}}{\rho_{de}}$$

Dark energy
equation of state

$$a(t) = \frac{1}{1+z}$$

$a(t)$ scale factor
 z redshift

$$E(a) = \left[\Omega_m a^{-3} + \Omega_{de} a^{-3(1+w)} + \Omega_k a^{-2} \right]^{1/2}$$

Evolution parameter
 $E(a) = H(a)/H_0$

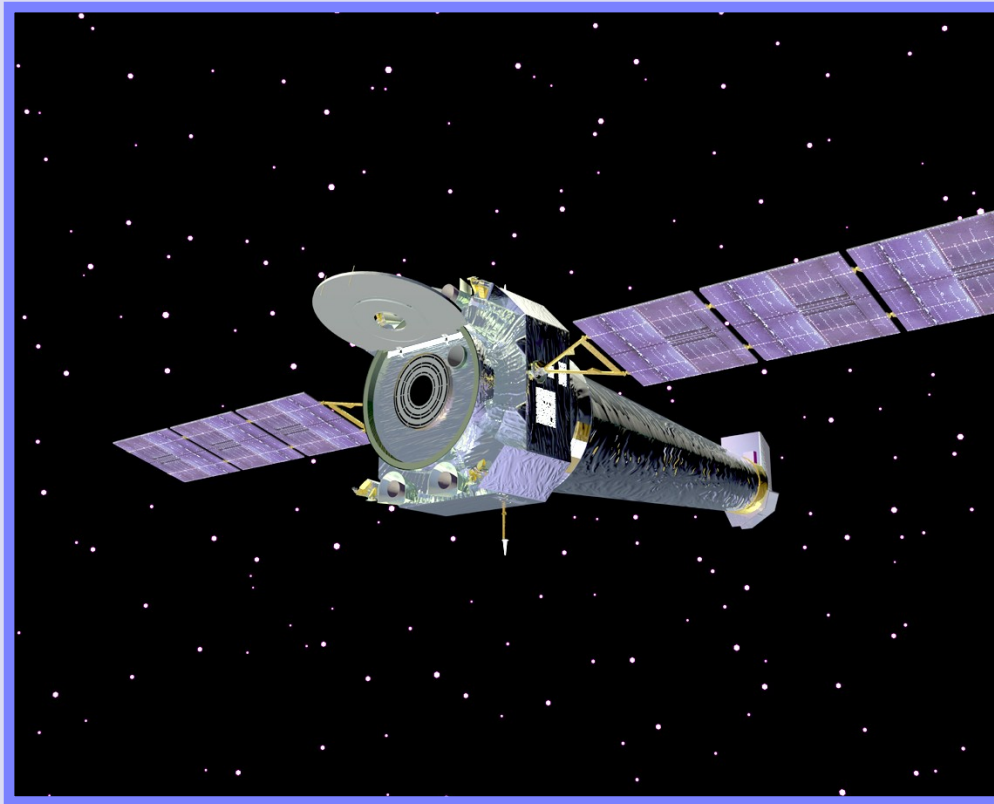
$$E(z) = \sqrt{\Omega_m (1+z)^3 + \Omega_{de} f(z) + \Omega_k (1+z)^2}$$

- i) flat Λ CDM $w=-1, \Omega_k=0$
- ii) flat w CDM w constant, $\Omega_k=0$
- iii) non-flat Λ CDM $w=-1, \Omega_k$ constant

The $f_{\text{gas}}(z)$ experiment

e.g. Allen et al 02, 04, 08; Ettori et al 03, 09;
Rapetti et al. 05, 07, 08; LaRoque et al 06

Chandra X-ray Observatory



Cluster cosmology revolutionized

First opportunity to carry out:

- > Detailed **spatially-resolved** and
- > **X-ray spectroscopy of galaxy clusters.**

Technical details for **ACIS** instrument (X-ray CCDs):

- Field of view **16x16 arcmin²**
- Good spectral resolution **~100eV** over 0.5-8 keV range.
- Exquisite spatial resolution (**0.5 arcsec** FWHM).

Measuring comic matter content using the gas mass fraction of X-ray luminous galaxy clusters

Consider a spherical region of observed angular radius θ within which the gas mass fraction is measured

$$R = \theta d_A$$

R physical size

$$L_x = 4\pi d_L^2 F_x$$

L_x X-ray Luminosity of the region
 F_x detected flux

$$d_L = d_A (1 + z)^2$$

d_L luminosity distance
 d_A angular diameter distance

Since the X-ray emission is mainly due to collisional processes (bremsstrahlung and emission line) and is optically thin

$$L_x \propto n^2 V$$

n mean matter density of colliding gas particles
 V volume of the emitting region: $V = 4\pi(\theta d_A)^3 / 3$

Measuring cosmic matter content using the gas mass fraction of X-ray luminous galaxy clusters

$$n \propto \frac{d_L}{d_A^{3/2}} \longrightarrow M_{gas} \propto nV \propto d_L d_A^{3/2}$$

M_{gas} observed gas mass within the measurement radius

$$M_{tot} \propto d_A$$

M_{tot} total mass determined by X-ray data assuming **hydrostatic equilibrium** (see below)

$$M_{gas} \text{ gas mass} \propto d_A(z)^{2.5} \text{ (X-ray Luminosity)}$$

$$M_{tot} \text{ total cluster mass} \propto d_A(z) \text{ (primarily X-ray Temperature)}$$

$$f_{gas} = \frac{M_{gas}}{M_{tot}} \propto d_L d_A^{1/2}$$

f_{gas} gas mass fraction

Measuring cosmic matter content using the gas mass fraction of X-ray luminous galaxy clusters

$$f_{gas} = \frac{M_{gas}}{M_{tot}} \propto d_A(z)^{1.5}$$

$$s = f_{stars} / f_{gas} = (0.16 \pm 0.05) h_{70}^{0.5}$$
 Baryonic mass fraction in stars

Lin & Mohr 04, Fukugita et al 98, White et al 93

$$f_{baryon} = f_{stars} + f_{gas} = f_{gas} (1 + s)$$
 Baryon mass fraction

Measuring cosmic matter content using the gas mass fraction of X-ray luminous galaxy clusters

$$f_{\text{baryon}} = b \frac{\Omega_b}{\Omega_m}$$

The matter content of rich clusters of galaxies is **expected to provide an almost fair sample of the matter content of the Universe** (White & Frenk 91, White et al. 93, Eke et al. 98).

b, the bias factor accounts for the relatively small amount of gas expelled when clusters form.

$$\Omega_m = \frac{b\Omega_b}{f_{\text{gas}}(1 + s)}$$

+HST+BBNS priors when clusters alone or +CMB data

Measuring cosmic acceleration using the gas mass fraction of X-ray luminous galaxy clusters

$$f_{gas}^{ref}(z) = \frac{b(z)\gamma K}{1+s(z)} \left(\frac{\Omega_b}{\Omega_m} \right) \varepsilon(\theta) \left[\frac{d_A^{ref}(z)}{d_A^{mod}(z;\theta)} \right]^{3/2}$$

Apparent evolution of the gas mass fraction

$$\varepsilon(\theta) = \left[\frac{H^{mod}(z;\theta) d_A^{mod}(z;\theta)}{H^{ref}(z) d_A^{ref}(z)} \right]^\eta$$

Small angular correction

$$\eta = 0.214 \pm 0.022 \quad \text{Measured from the data profiles}$$

Measuring cosmic acceleration using the gas mass fraction of X-ray luminous galaxy clusters

Small angular correction that accounts for the angle subtended at the measurement radius r_{2500} as the underlying cosmology varies

$$\varepsilon(\theta) = \left(\frac{\theta_{2500}^{ref}}{\theta_{2500}^{mod}} \right)^\eta$$

For each cluster the measured f_{gas} value at r_{2500} corresponds to a fixed angle θ_{2500}^{ref} for the reference cosmology that is slightly different from that θ_{2500}^{mod} for the test cosmology.

$$M_{2500} \propto 4\pi r_{2500}^3 \rho_{crit} / 3$$

Mass at the measurement radius r_{2500} , for which the density is 2500 the critical density

$$\rho_{crit} = 3H(z)^2 / 8\pi G$$

ρ_{crit} critical density

Measuring cosmic acceleration using the gas mass fraction of X-ray luminous galaxy clusters

Assuming **hydrostatic equilibrium** (HSE) in the intracluster medium (ICM) and **spherical symmetry** we can calculate the mass within a given radius using the following expression (Sarasin 1988)

$$M(r) = -\frac{rkT(r)}{G\mu m_p} \left[\frac{d \ln n}{d \ln r} + \frac{d \ln T}{d \ln r} \right]$$

$n(r)$ is the gas density
 $T(r)$ ICM temperature
 k Boltzmann constant
 μm_p mean molecular weight

Measuring **cluster masses** is one of the cornerstones of cluster cosmology. Under those assumptions we can measure the total mass from density $n(r)$ and temperature $T(r)$ profiles obtained from X-ray data. **Note also that $M(r)$ depends more strongly on $T(r)$ than $n(r)$.**

Given that the temperature, and temperature and density gradients, in the region of θ_{2500} are likely to be constant, we have

$$M_{2500} \propto r_{2500}$$

Measuring cosmic acceleration using the gas mass fraction of X-ray luminous galaxy clusters

$$\left. \begin{aligned} M_{2500} &\propto r_{2500} \\ M_{2500} &\propto 4\pi r_{2500}^3 \rho_{crit} / 3 \\ \rho_{crit} &= 3H(z)^2 / 8\pi G \end{aligned} \right\} r_{2500} \propto H(z)^{-1}$$

$$r_{2500} \propto H(z)^{-1} \longrightarrow \theta_{2500} = r_{2500} / d_A \propto [H(z)d_A]^{-1}$$

Angle spanned by r_{2500} at redshift z

$$\varepsilon(\theta) \equiv \frac{\varepsilon(\theta)^{\text{mod}}}{\varepsilon(\theta)^{\text{ref}}} = \left(\frac{\theta_{2500}^{\text{ref}}}{\theta_{2500}^{\text{mod}}} \right)^\eta \approx \left[\frac{H^{\text{mod}}(z; \theta) d_A^{\text{mod}}(z; \theta)}{H^{\text{ref}}(z) d_A^{\text{ref}}(z)} \right]^\eta$$

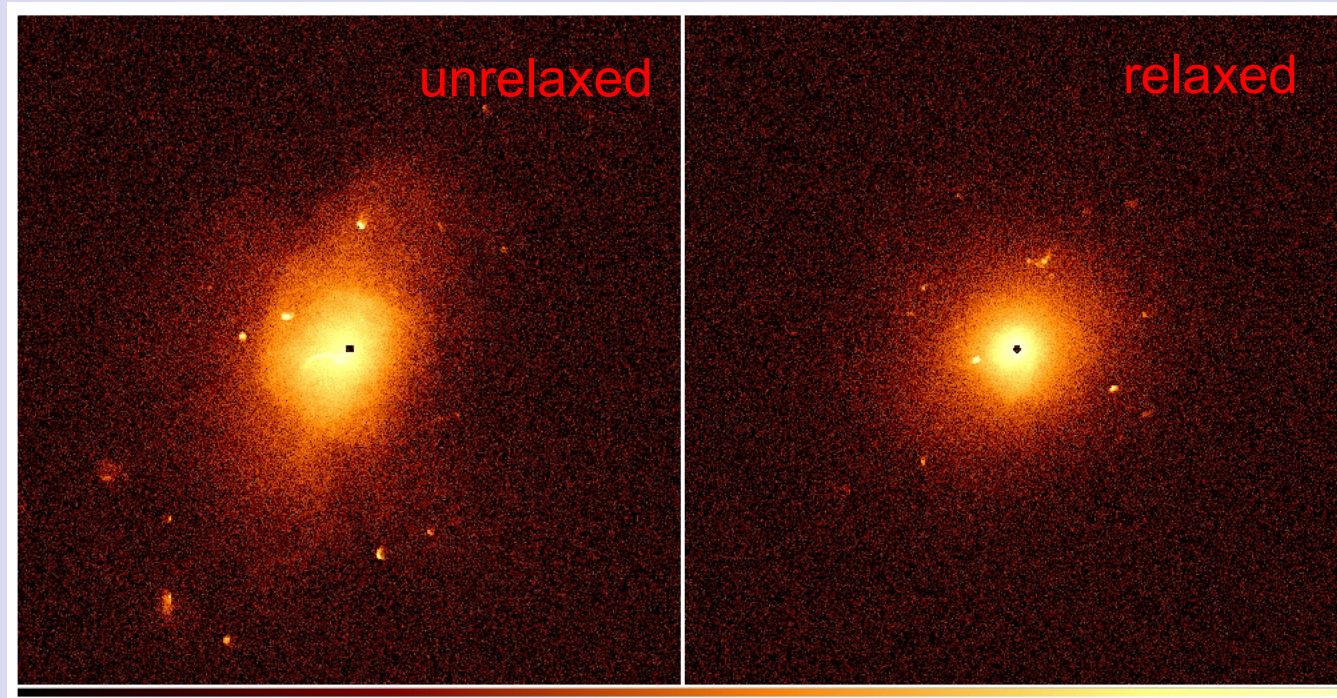
X-ray galaxy cluster cosmology: How well can we measure cluster mass?

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Hydro dynamical simulations of X-ray galaxy clusters

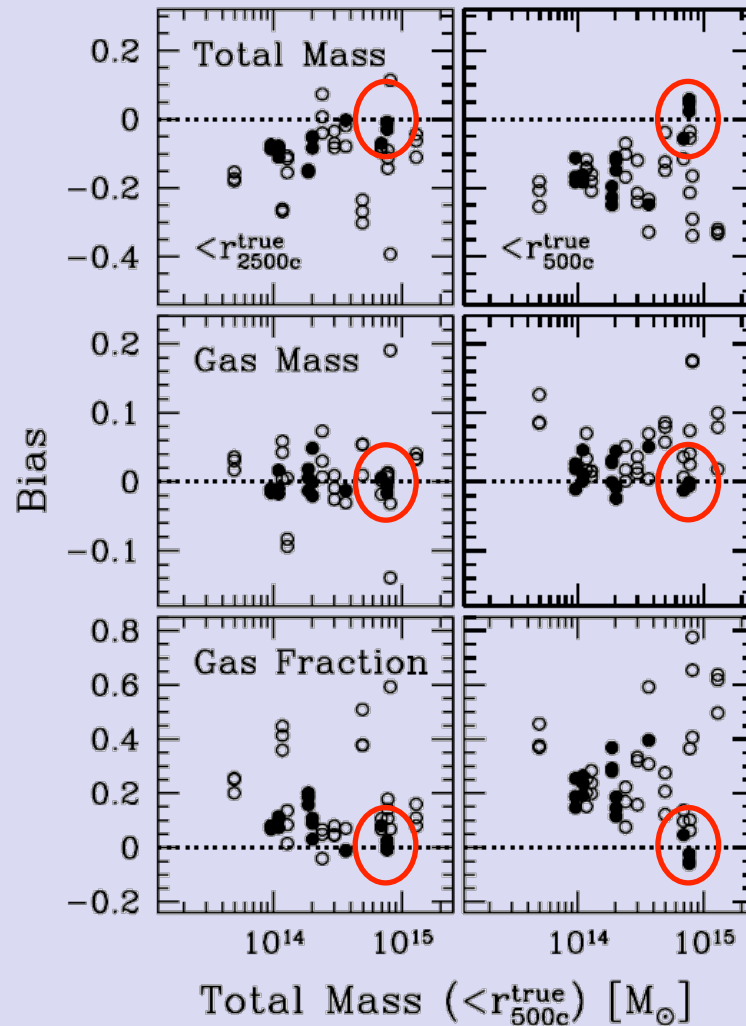
Nagai, Kravtsov, Vikhlinin 06



Very good news for X-ray galaxy cluster cosmology from the recent simulations: **systematics are relatively small and can be quantified.**

How accurately can we measure the mass?

Nagai, Kravtsov, Vikhlinin 06



For the largest, hottest ($kT > 5\text{keV}$), relaxed clusters (selection based on **X-ray morphology**) we currently expect to measure:

a) X-ray gas mass to $\sim 1\%$ accuracy.

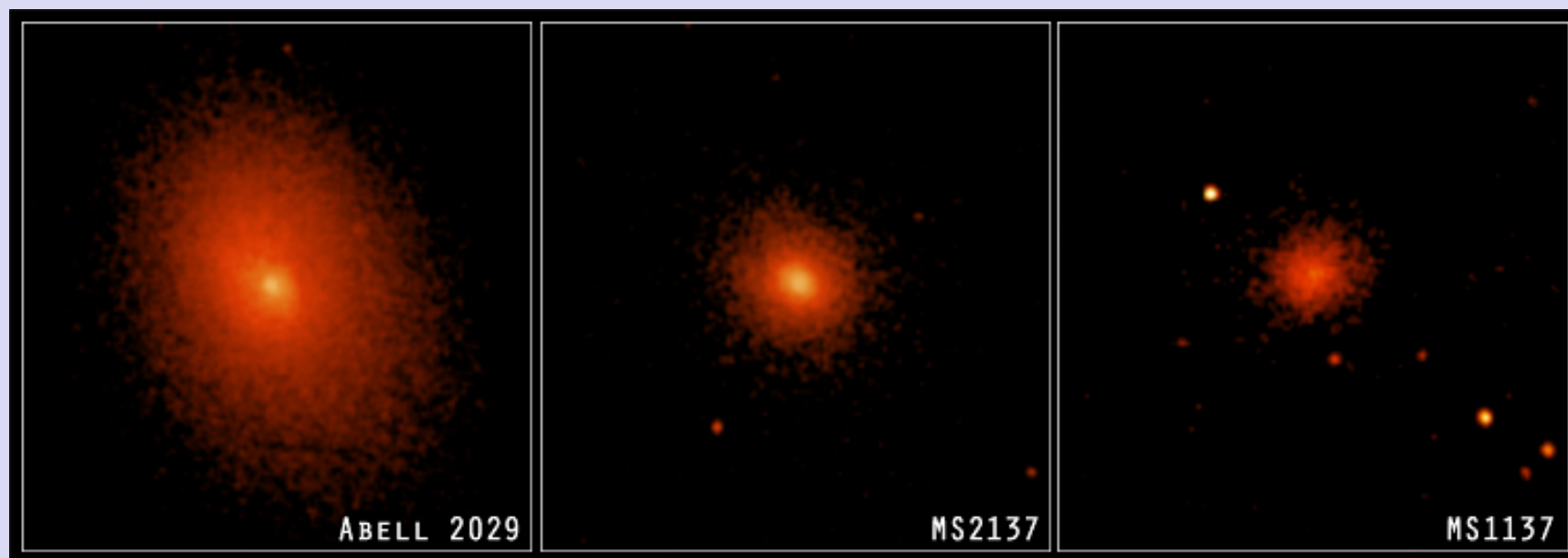
b) Total mass to few % accuracy (both bias and scatter).

Largest, relaxed clusters (filled points) inside red circles.

X-ray galaxy cluster data

It is crucial to use only **dynamically relaxed clusters**:

Regular X-ray morphology, low ellipticities, minimal centroid variation, sharp central brightness peaks centered on their dominant elliptical galaxies.



Measuring cosmic acceleration using the gas mass fraction of X-ray luminous galaxy clusters

$$F(d_A^{\text{mod}}) \equiv \frac{d_A^{\text{mod}}(z; \theta)^{3/2}}{\varepsilon(\theta)^{\text{mod}}} = \frac{b(z) \gamma K \left(\frac{\Omega_b}{\Omega_m} \right)}{1 + s(z)} \left[\frac{d_A^{\text{ref}}(z)^{3/2}}{\varepsilon(\theta)^{\text{ref}} f_{\text{gas}}^{\text{ref}}(z)} \right]$$

Angular diameter distance measurement for the fgas experiment

Compare to supernova measurements used as standard candles:

$$\mu^{\text{th}}(z) = 5 \log_{10} d_L(z; \theta) + \mu_0$$

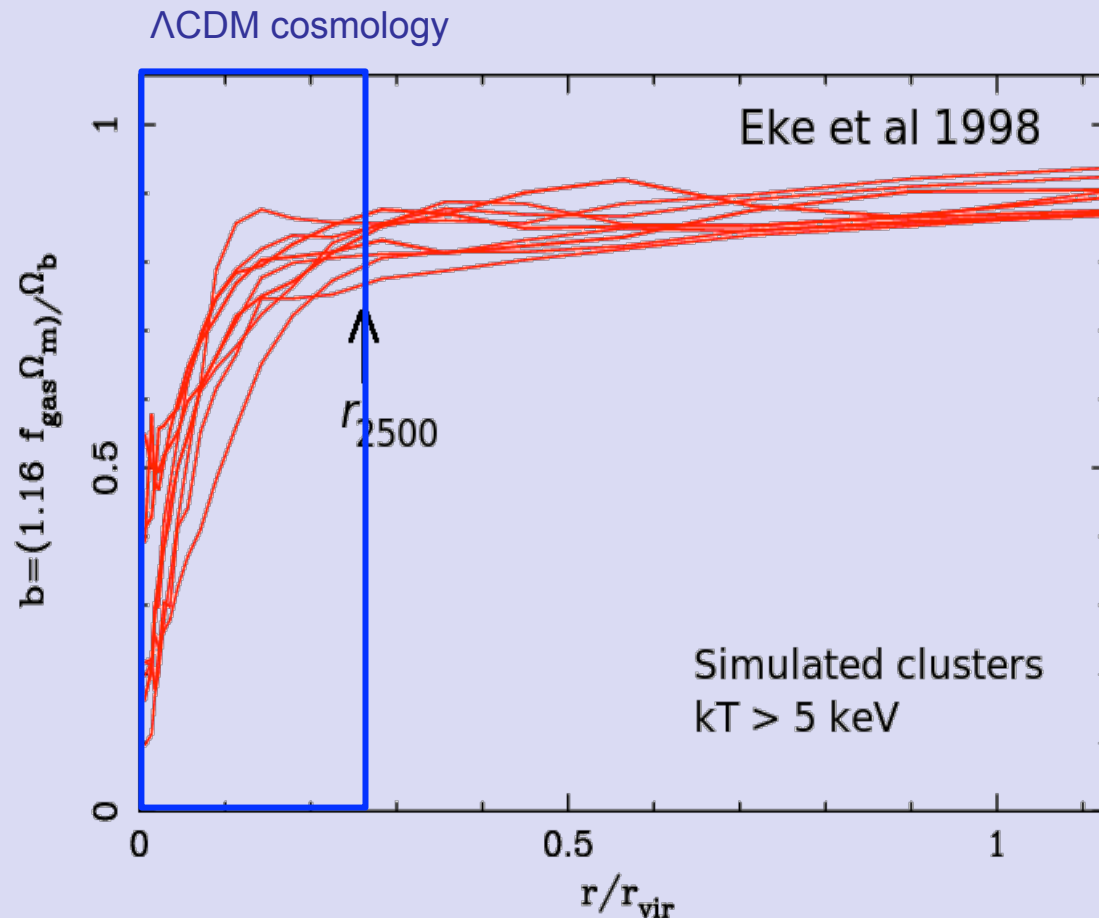
Apparent magnitude

$$d_L(z; \theta) = \frac{c(1+z)}{H_0} \int_0^z \frac{dz}{E(z; \theta)}$$

Luminosity distance

Gas mass: simulations

Low scatter total mass proxy



Simulations indicate low cluster-to-cluster scatter

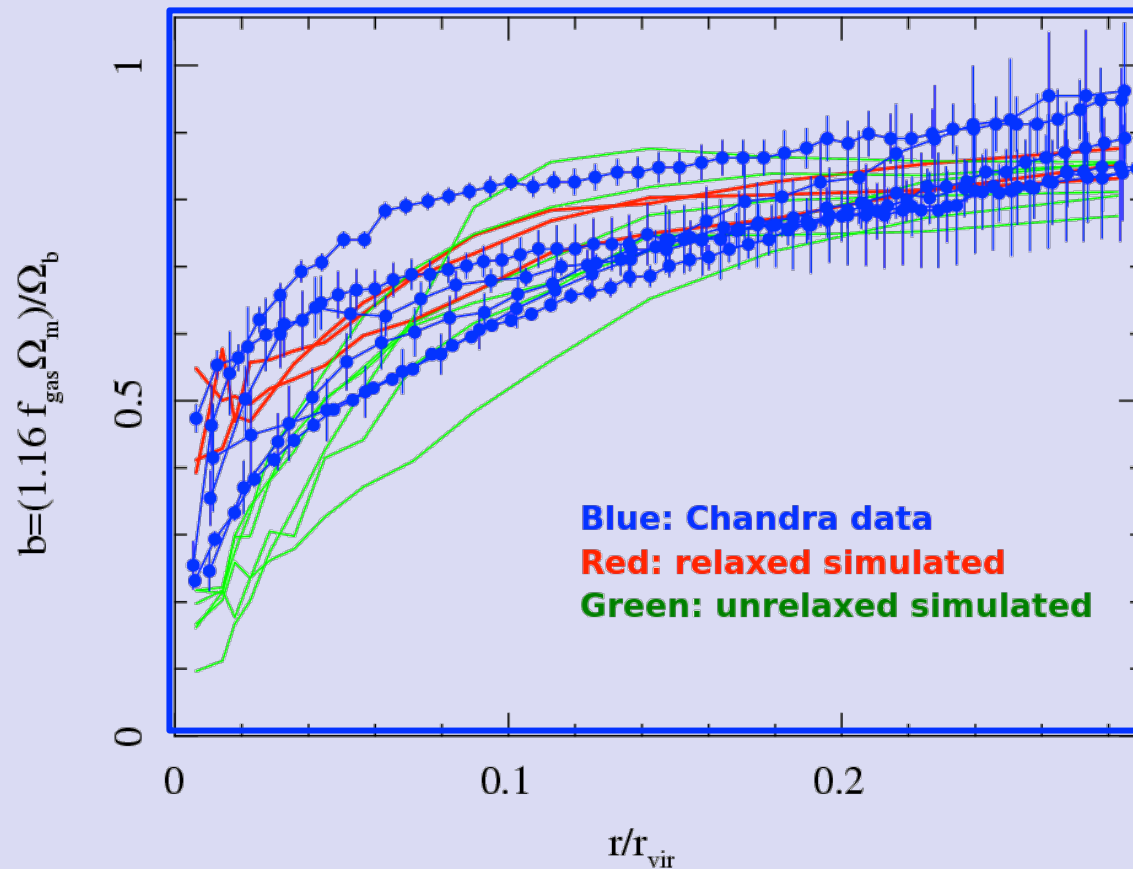
Gas-mass low-scatter total-mass proxy through f_{gas}

Simulations indicate that the baryon mass fraction in clusters is slightly lower than the mean value for the Universe as a whole. Some gas is lifted beyond the virial radius by shocks (e.g. Evrard et al 90, Thomas & Couchman 92, Navarro & White 93; NFW 95 etc, Kay et al 04, Ettori et al 06, Crain et al 06, Nagai et al 07).

Gas mass: data

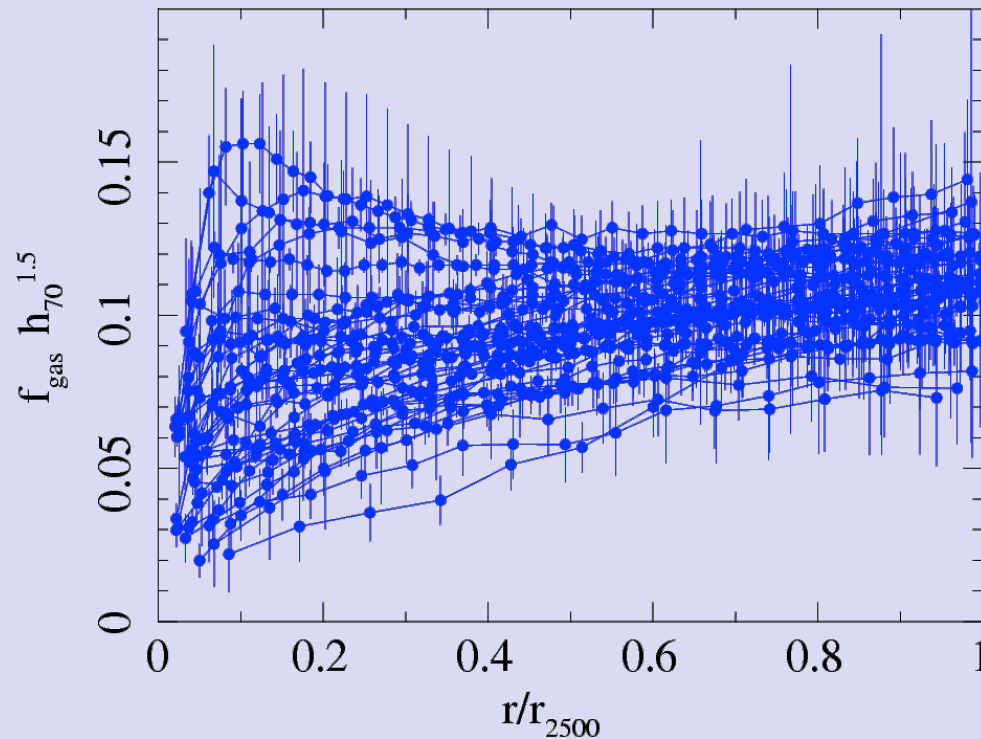
Low scatter total mass proxy

Λ CDM ($\Omega_m=0.3$, $\Omega_\Lambda=0.7$)



Undetected systematic scatter when weighted mean scatter is ~5% in distance

Chandra $f_{\text{gas}}(r)$ data: 42 relaxed clusters with redshifts $0.06 < z < 1.07$



Fitting a constant value at r_{2500} :

$$f_{\text{gas}}(r_{2500}) = (0.1104 \pm 0.0016) h_{70}^{-1.5}$$

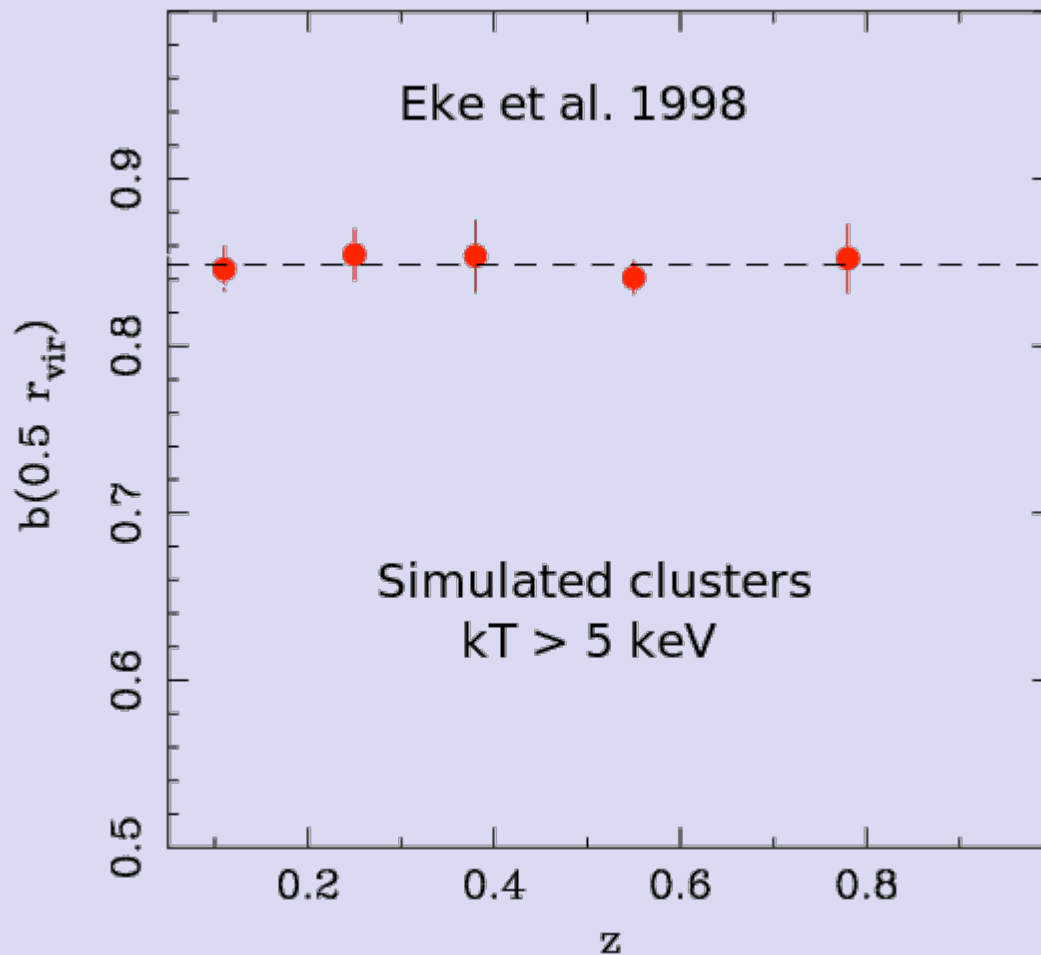
Fitting a power law $(0.7-1.2)r_{2500}$:

$$f_{\text{gas}}(r_{2500}) = (0.1105 \pm 0.0005) (r/r_{2500})^{0.214 \pm 0.022}$$

Assuming hydrostatical equilibrium and spherical symmetry (only relaxed clusters).

Gas mass: simulations

Low scatter total mass proxy



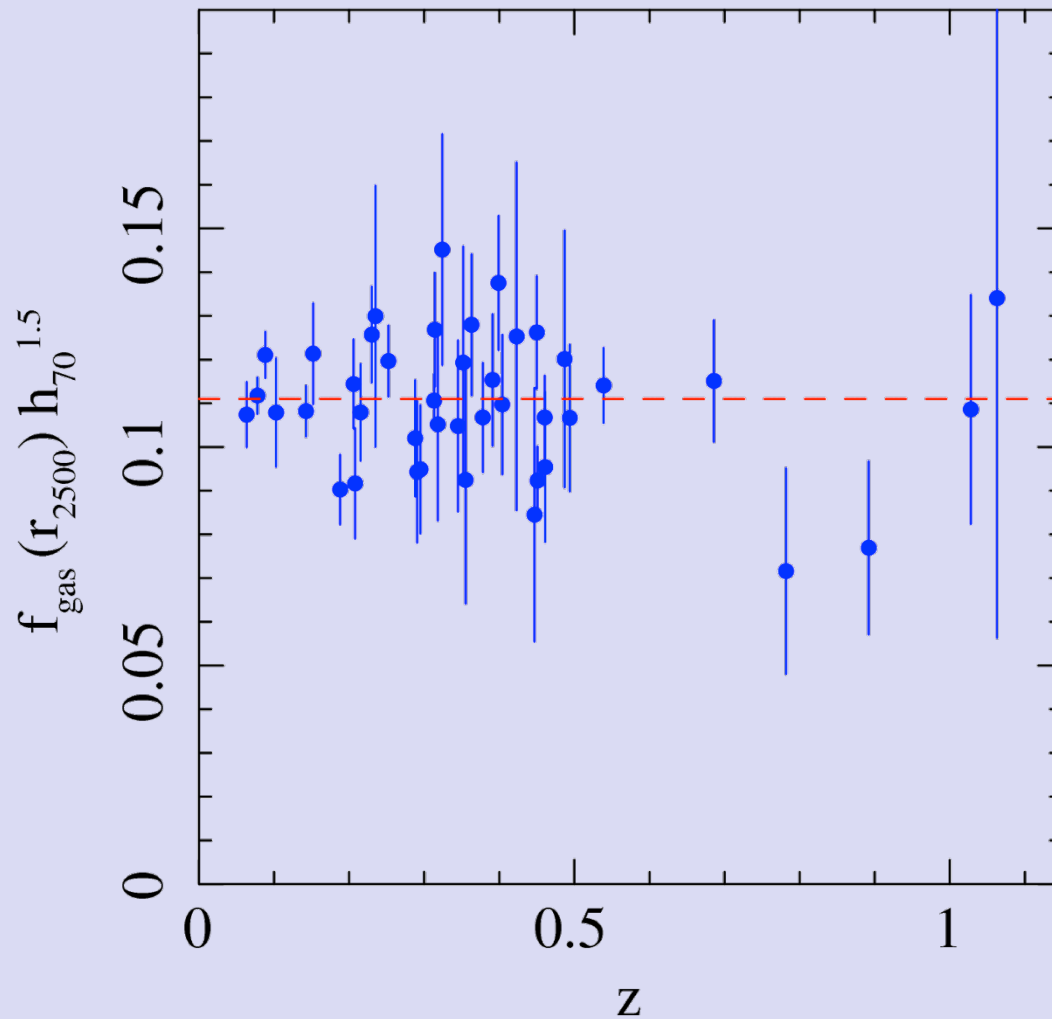
Tens of simulated clusters,
 $0 < z < 0.8$, $kT > 5 \text{ keV}$

Minimal evolution at r_{2500}

Gas-mass low-scatter total-mass proxy through f_{gas}

Gas mass fraction: data

Low scatter total mass proxy



Λ CDM ($\Omega_m=0.3$, $\Omega_\Lambda=0.7$)

42 dynamically relax clusters

Hot $kT > 5\text{keV}$

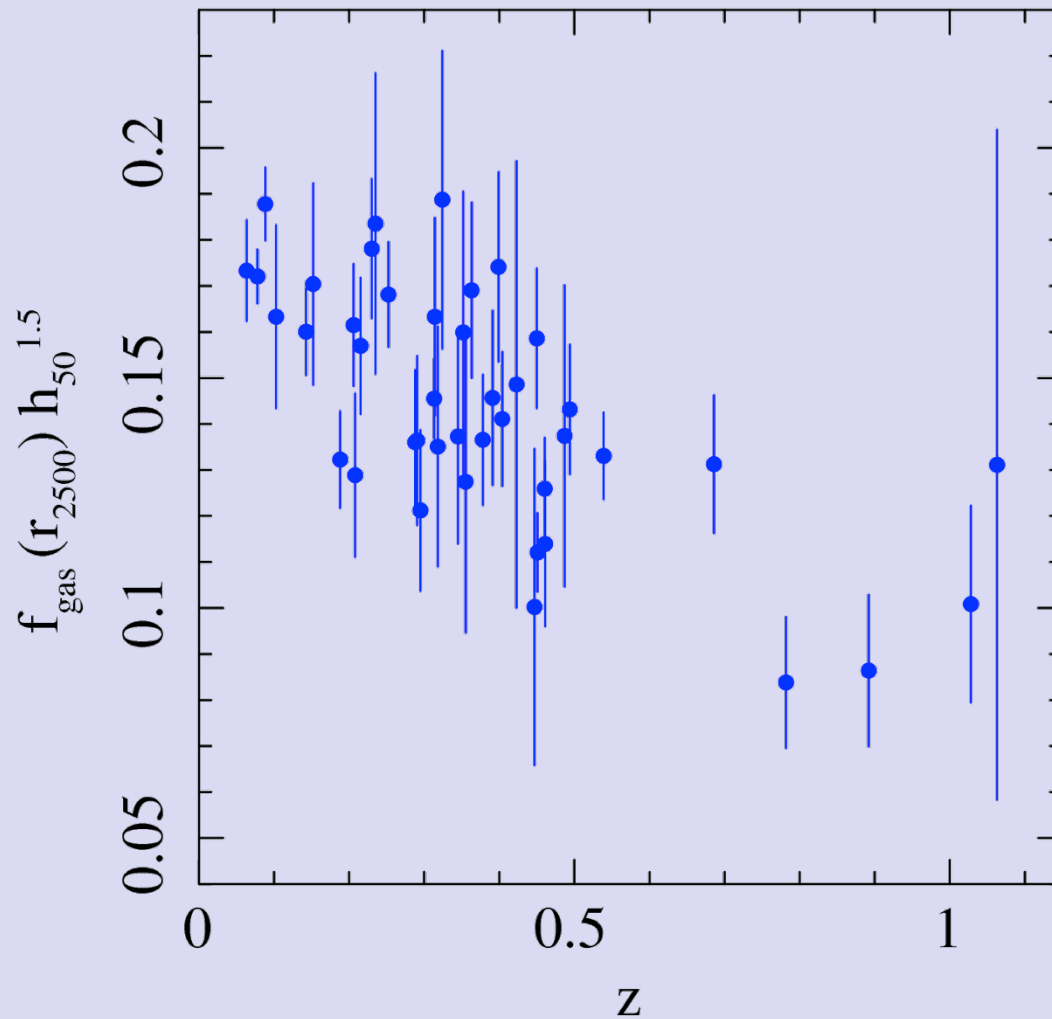
$0.06 < z < 1.07$

Scatter $< \sim 10\%$ in f_{gas}

Undetected systematic scatter
when weighted mean scatter
is $\sim 5\%$ in distance

Gas mass fraction: data

Low scatter total mass proxy



SCDM ($\Omega_m=1.0$, $\Omega_\Lambda=0.0$)

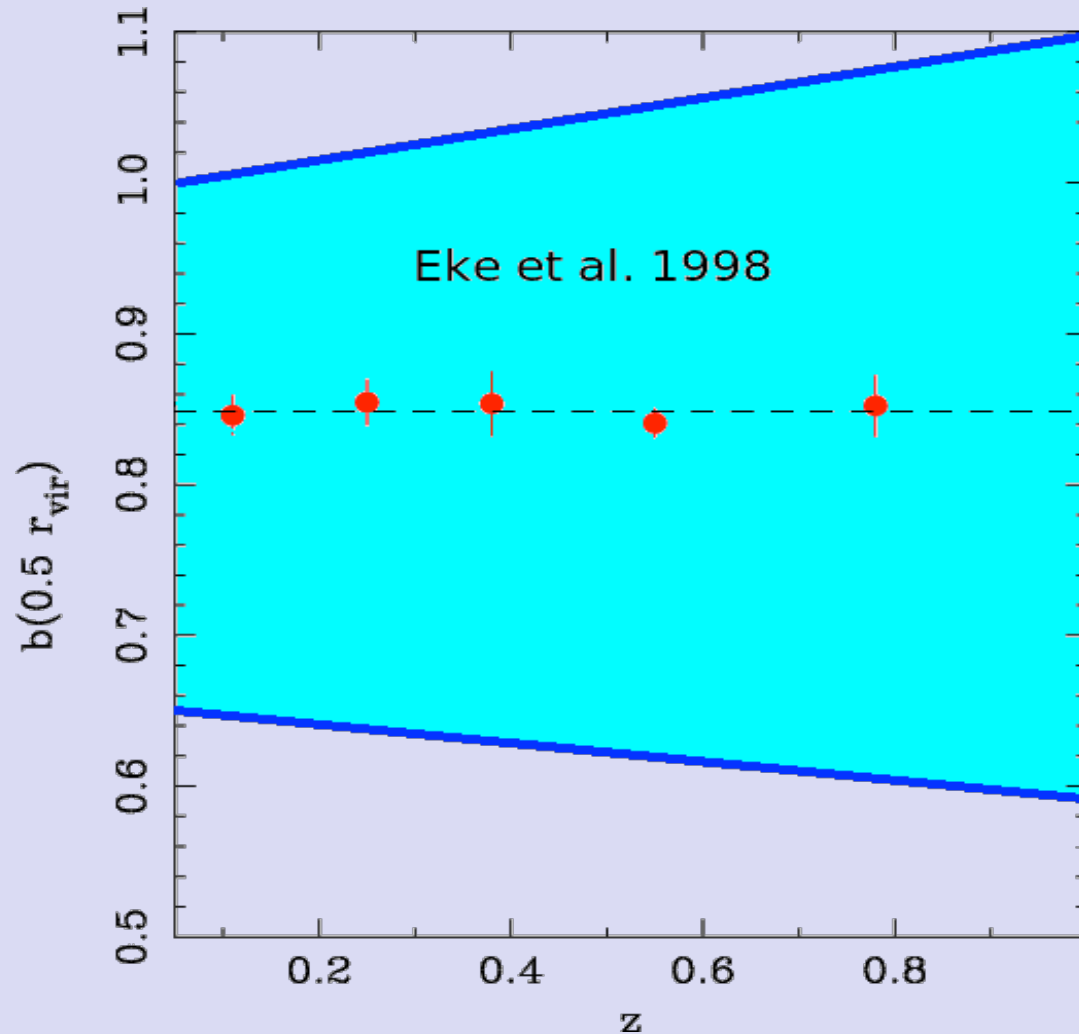
42 dynamically relax clusters

Hot $kT > 5\text{keV}$

$0.06 < z < 1.07$

Scatter $< \sim 10\%$ in f_{gas}

Allowances for systematic uncertainties



1) Gas depletion (simulation physics)

$$b(z) = b_0(1 + \alpha_b z)$$

normalization:

20% uniform prior

$$0.65 < b_0 < 1.0$$

evolution:

10% at $z=1$ uniform prior

$$-0.1 < \alpha_b < 0.1$$

Allowances for systematic uncertainties

2) Instrument calibration and modelling (gas clumping, etc.)

1.0 ± 0.1 , 10% Gaussian prior on K

3) Baryonic mass in stars

$$s(z) = s_0(1 + \alpha_s z)$$

normalization s_0 : 30% Gaussian uncertainty (observational)

evolution $-0.2 < \alpha_s < 0.2$: 20% at $z=1$ uniform prior (observational)

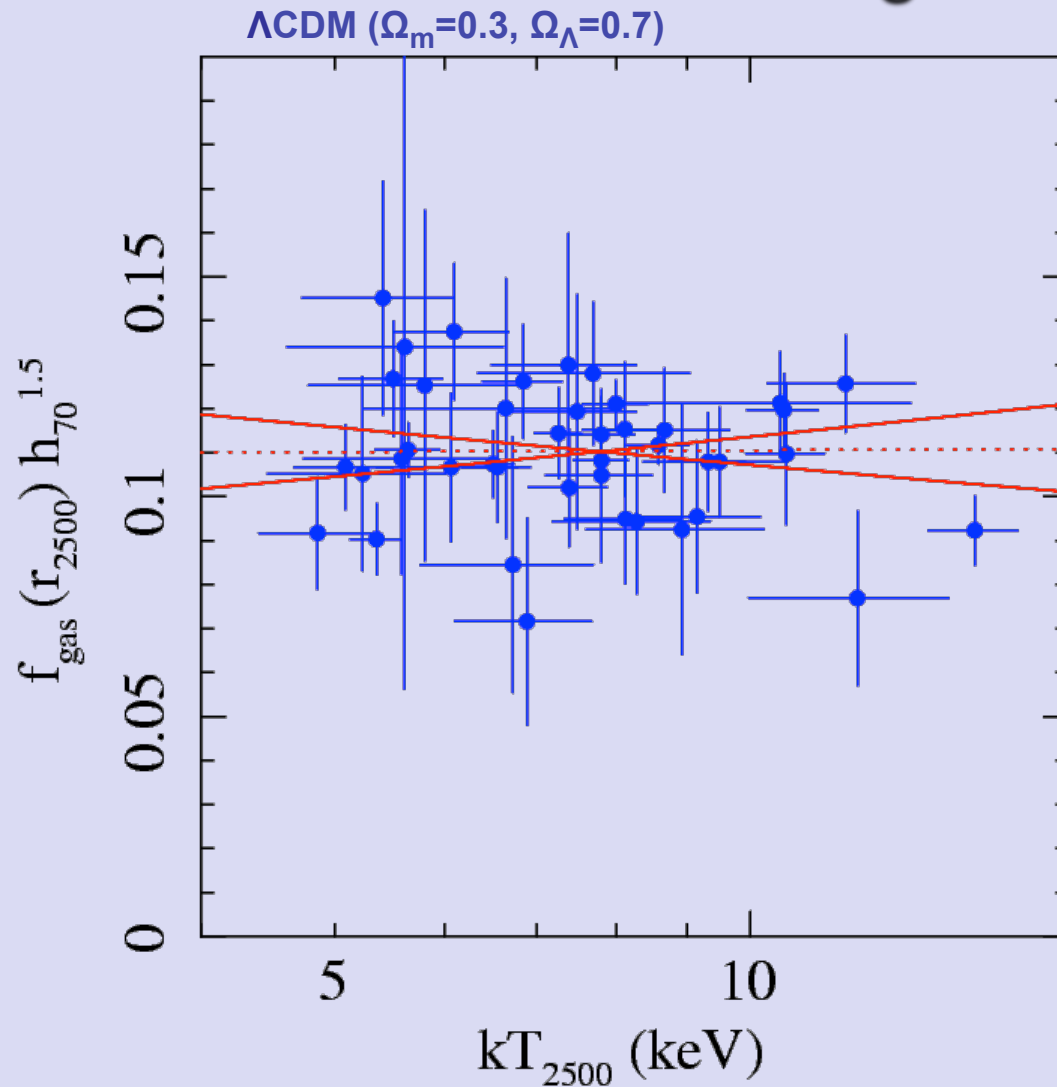
4) Non-thermal pressure support in gas: (primarily due to bulk motions)

$$\gamma = M_{\text{true}} / M_{\text{X-ray}}$$

$$1 < \gamma < 1.1$$

10% uniform (Nagai et al 07, Werner et al 09, Sanders et al 09)

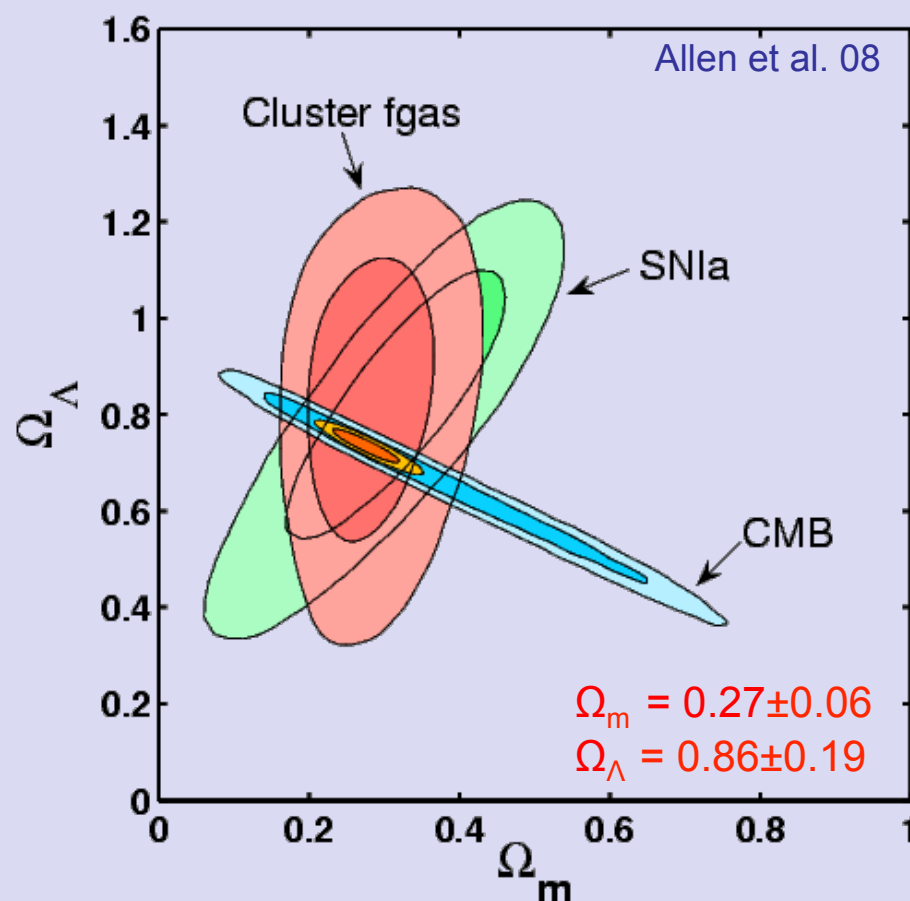
Chandra $f_{\text{gas}}(kT)$ data



Best fitting power law:
 $\alpha=0.005\pm0.058$ (solid lines 2σ limits).

f_{gas} essentially independent of temperature for the massive, dynamically relaxed clusters in the analysis.

Constraints on Λ CDM from three independent experiments



42 f_{gas} clusters (Allen et al 08)
including standard BBNS+HST priors
and full systematic allowances.

192 SNe Ia [Davis et al 07: Riess et al 07 (Gold sample), Wood-Vasey et al 07 (ESSENCE), Astier et al 06 (1st year SNLS)].

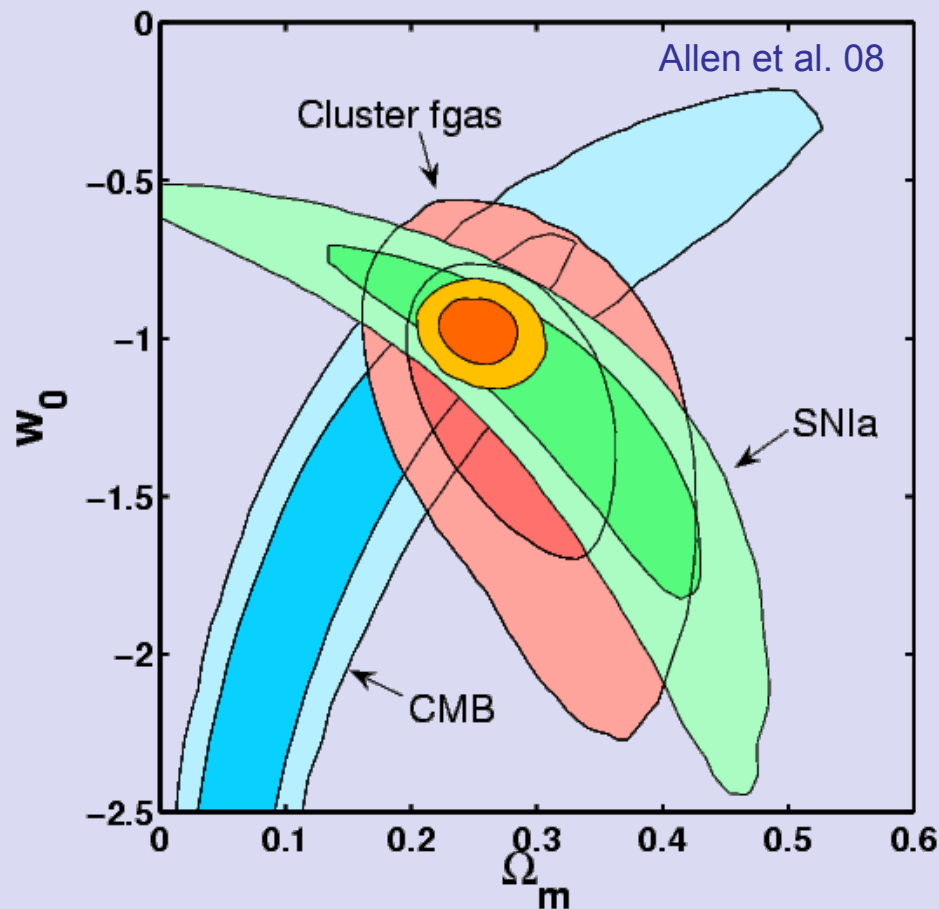
CMB data from WMAP3, CBI, Boomerang, ACBAR (prior $0.2 < h < 2.0$).

Combined constraints

$$\Omega_m = 0.275 \pm 0.033$$

$$\Omega_\Lambda = 0.735 \pm 0.023$$

Constraints on w CDM from three independent experiments



42 f_{gas} clusters (Allen et al 08)
including standard BBNS+HST priors
and full systematic allowances.

192 SNe Ia [Davis et al 07: Riess et al
07 (Gold sample), Wood-Vasey et al
07 (ESSENCE), Astier et al 06 (1st
year SNLS)].

CMB data from WMAP3, CBI,
Boomerang, ACBAR (prior
 $0.2 < h < 2.0$).

Combined constraints

$$\Omega_m = 0.253 \pm 0.021$$

$$w_0 = -0.98 \pm 0.07$$