

Finite temperature field theory and heavy ion collisions

Nordic Winter School on Particle Physics and
Cosmology 2013, Skeikampen, Norway

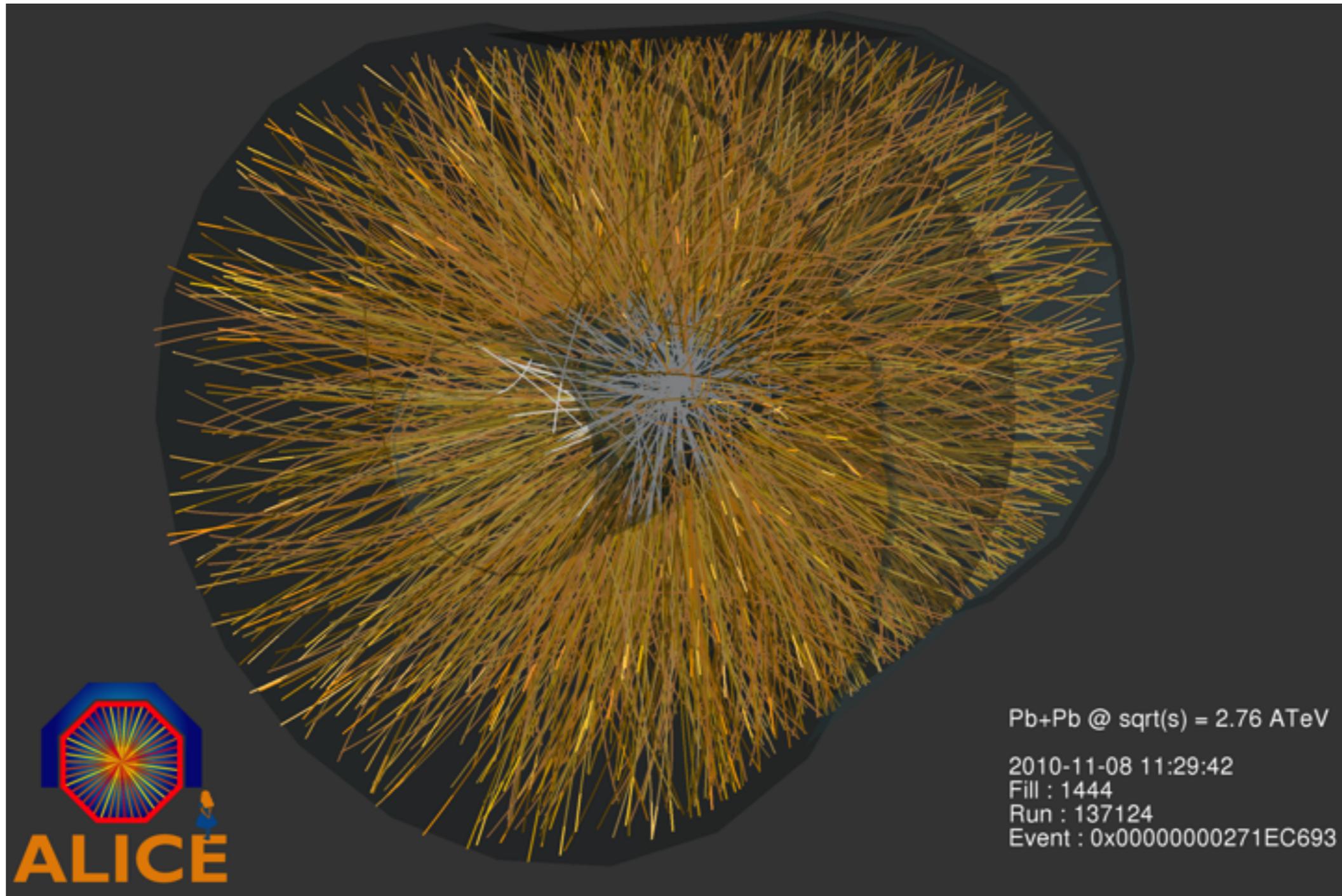
Lecture 1: Heavy ion collisions; hydrodynamics

Simon Caron-Huot
(NBIA, Copenhagen & IAS, Princeton)

Plan for lectures

- Heavy ion collisions: main results, hydrodynamics description
- Finite temperature field theory (Schwinger-Keldysh)
- Spectral functions, properties and scales of weakly coupled plasmas
- Energy loss in a plasma and hard probes

- I'll be talking about something *relatively* cold



$T \sim 400 \div 500 \text{ MeV}$ (LHC₃)

(~5 trillions K)

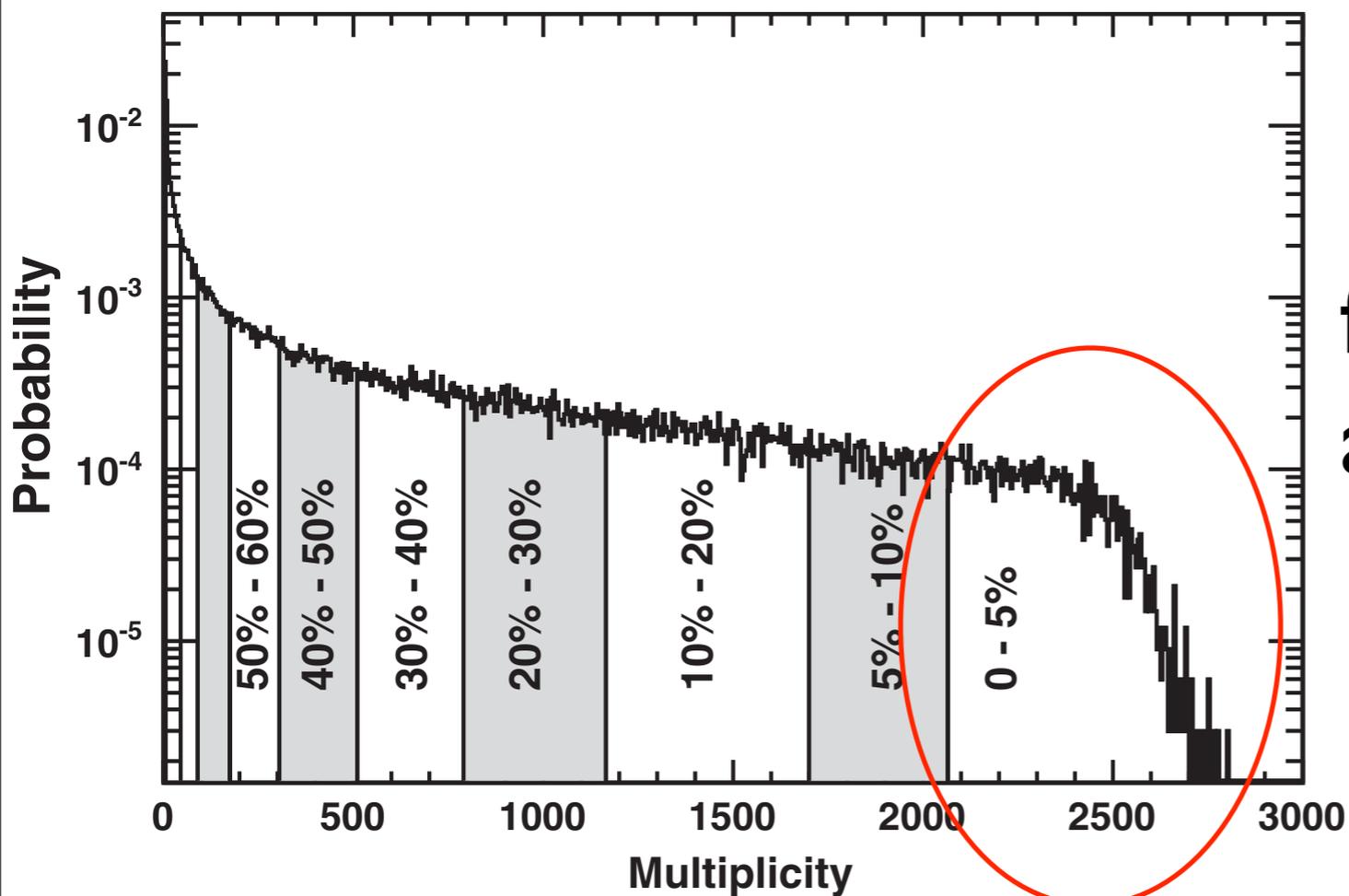
(ultrarelativistic) Heavy ion collisions

- Modern history:
 - RHIC (Brookhaven), Au+Au @ $\sqrt{s_{NN}}=200$ GeV
 - LHC, Pb+Pb @ $\sqrt{s_{NN}}=2.76$ TeV
- Main goal is to study QCD, the theory of the strong interactions, around and above the *deconfinement transition*
- First energies at which there is evidence for a new *deconfined* phase of matter

- Integrated luminosity $\sim 160 \mu\text{b}^{-1}$
(Atlas&CMS 2011, Alice is a bit lower)
- Total Pb-Pb cross-section $\sigma \sim 7\div 8 \text{ b}$
- Multiplicity is extremely well correlated with impact parameter

$dN_{\text{ch}}/dy \sim 1000$ (RHIC)
 ~ 1600 (LHC)

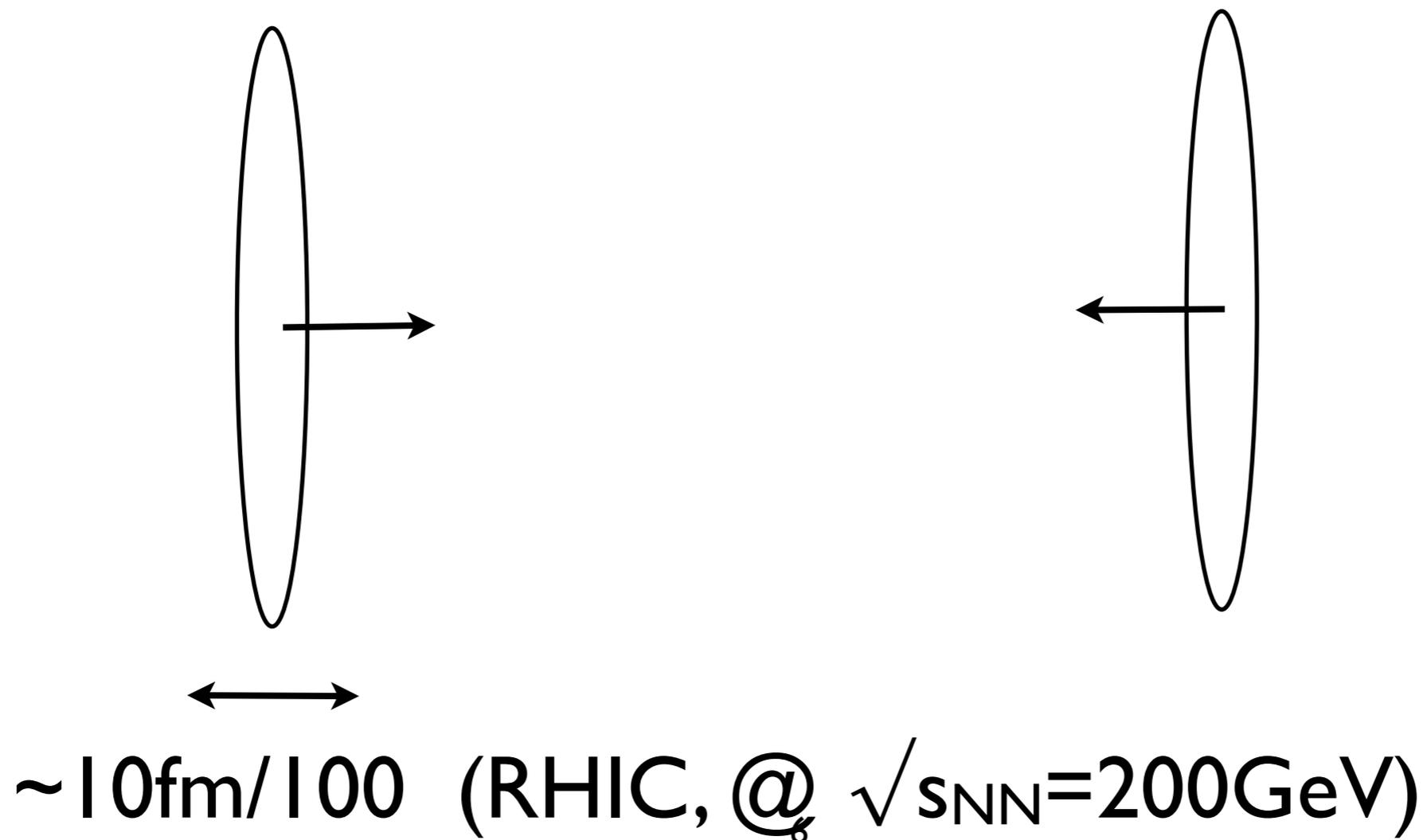
for 5% most central
at midrapidity



(Alice, PRL 252302)

Stages of a heavy ion collision

I. Lorentz-contracted 'pancakes' hit each other



1. A high density but out-of-equilibrium state is created; it expands longitudinally ($\tau \sim 0-0.5 \div 1 \text{ fm}/c$)

2. Rescattering becomes important: from here one assumes *local thermodynamic equilibrium*.

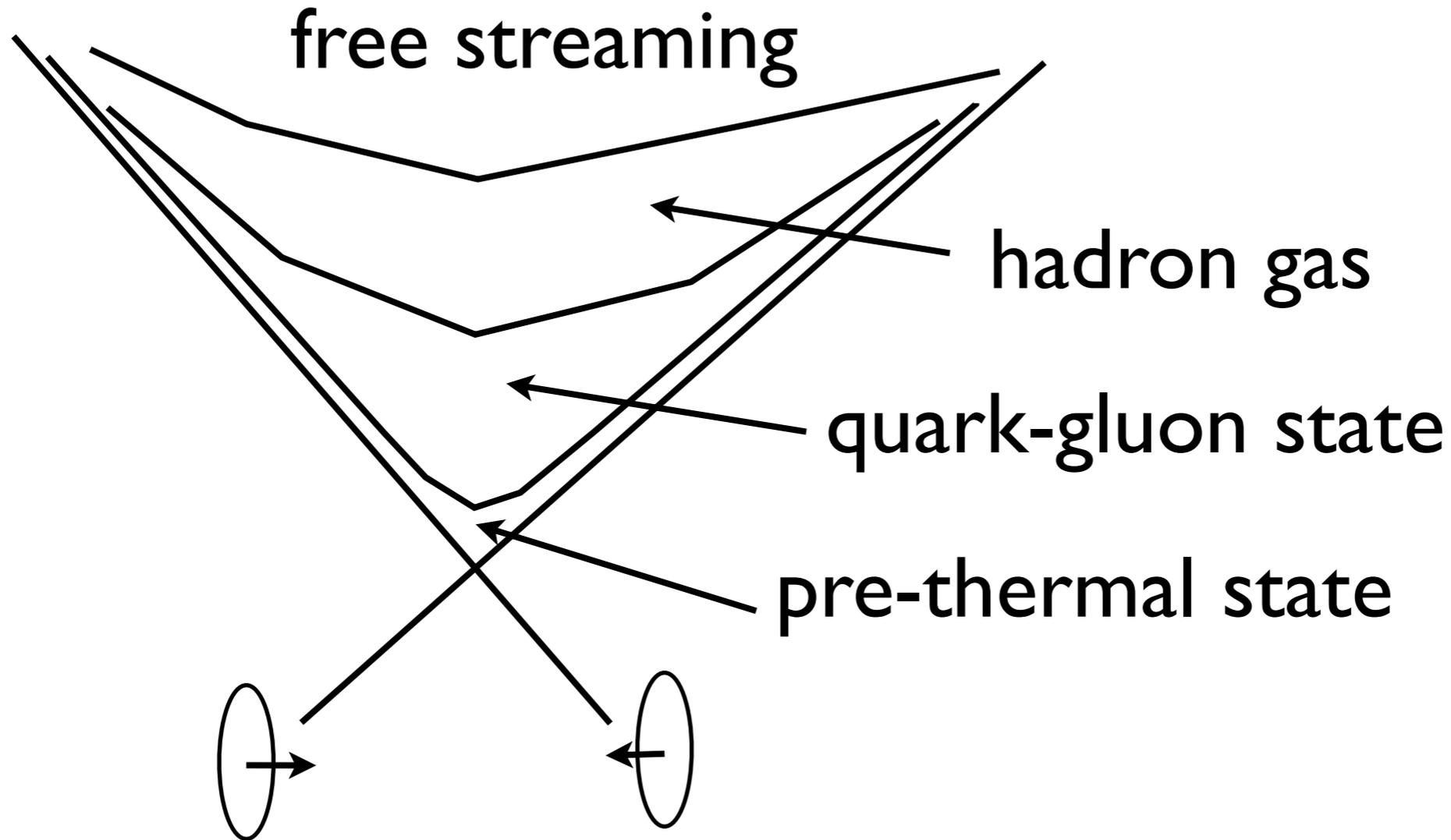
The system expands and cools hydrodynamically

3. T drops below $T_c \sim 170 \text{ MeV}$ ($\tau \sim 10 \text{ fm}/c$): *hadron gas*.

Around that time, *chemical freeze-out* occurs.

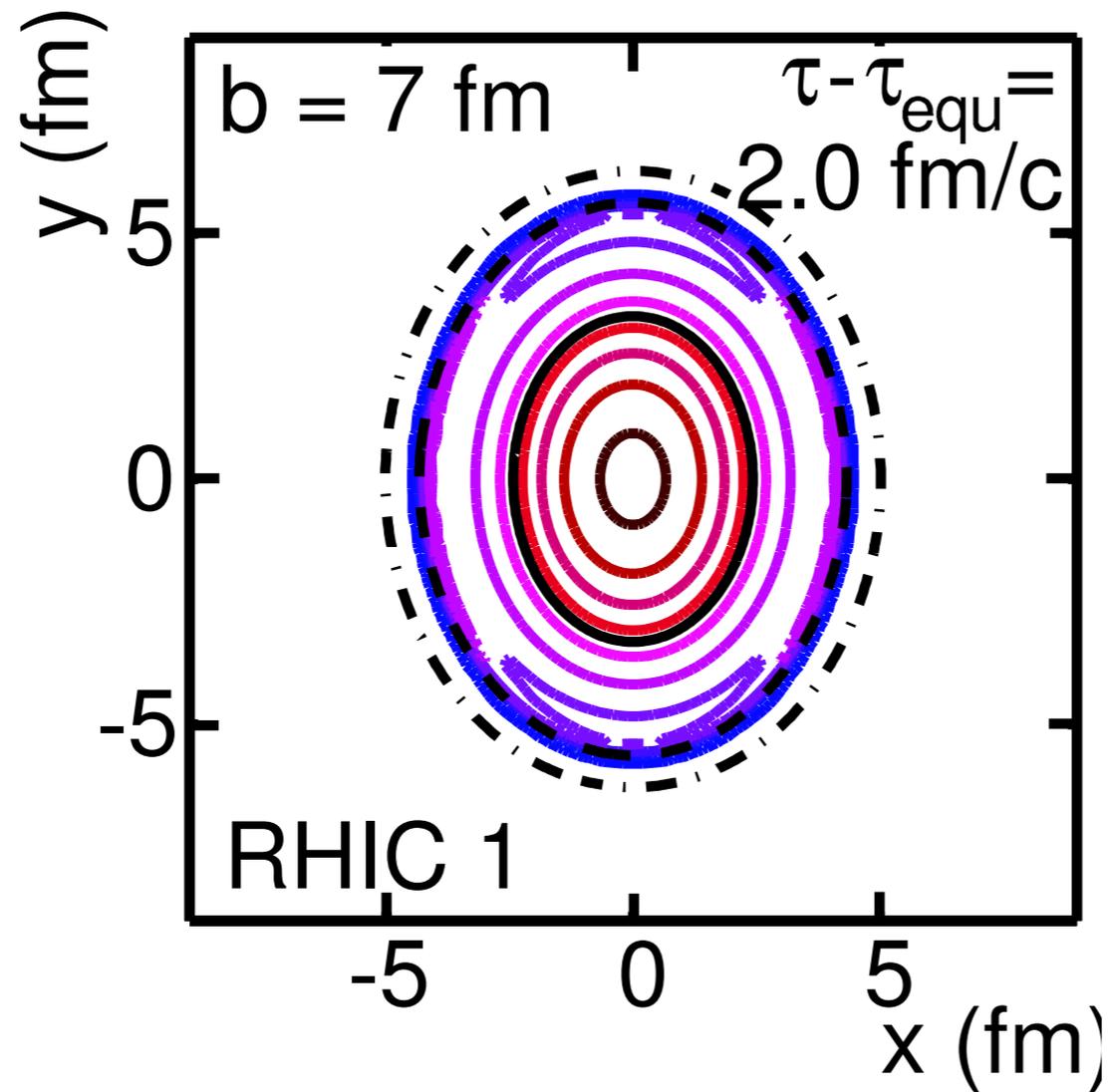
4. *Kinetic freeze-out* ($\tau \sim 20 \text{ fm}/c$): hadrons stream to the detector

Space-time diagram



Many stages/transitions seem prohibitively difficult to describe; fortunately, many of the details do not matter too much!

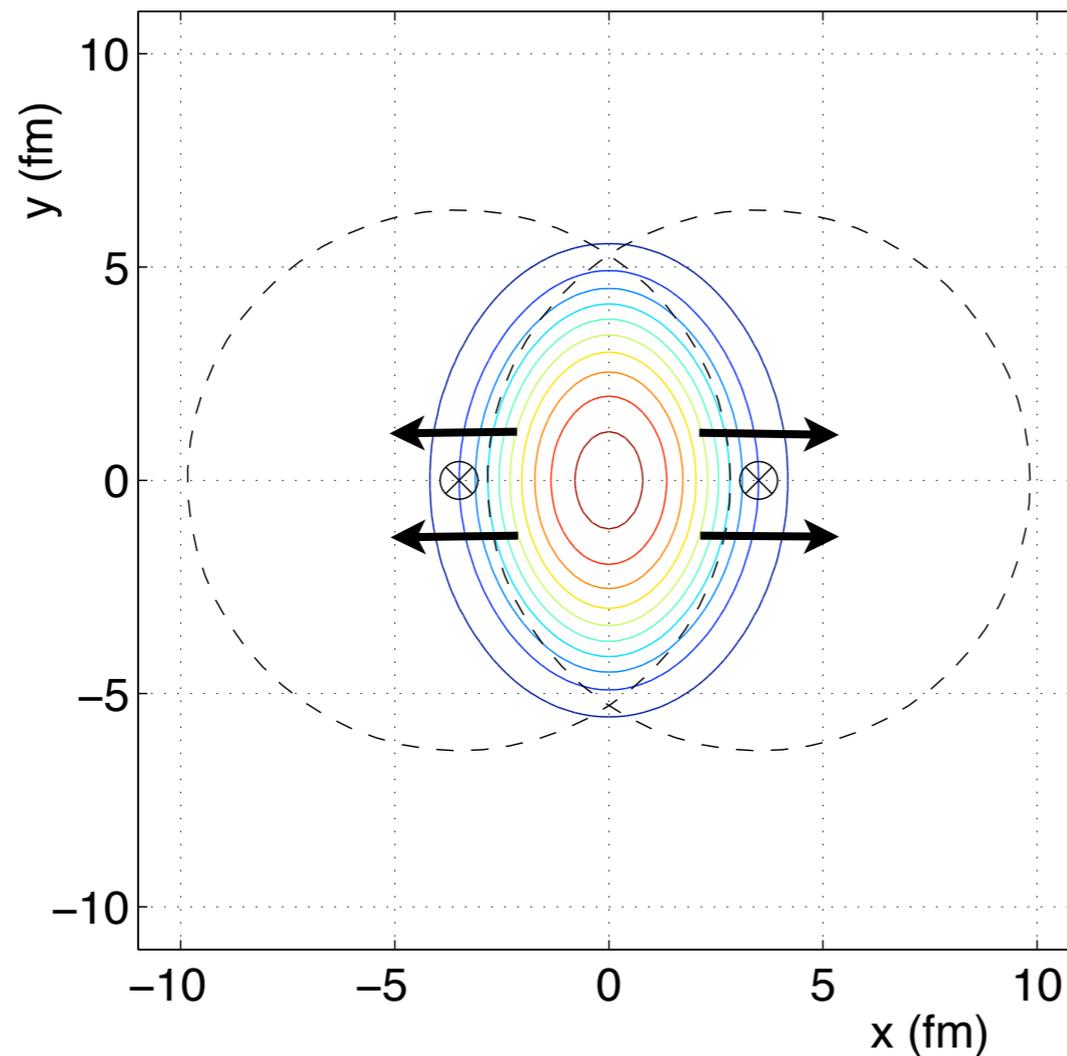
- What's the signature of rescattering?



(picture from U. Heinz's excellent review, 0901.4355)

Rescattering converts density gradients into *acceleration*
(‘fluid-like’ behavior)

- Radial acceleration per se is *not* observable
- However, initial conditions can be anisotropic



(again from U.Heinz's review)

- Beginning: *spacetime anisotropy*, mom. *isotropy*
- Final state: *momentum* \perp *space anisotropy*

Hydrodynamics, I

- Hydrodynamics used to be the ‘dynamics of water’
- In modern usage, it refers to the effective theory of [any] medium, on distances larger than a mean free path
- By definition, this describes *local thermodynamic equilibrium*
- There’s a hydrodynamics theory for any fluid; it depends on its IR degrees of freedom and is parametrized by a set of [Wilson] coefficients.

Hydrodynamics, II

- Start from equilibrium. Say equilibria are parametrized by a temperature T and a 4-velocity, $u_\mu u^\mu = -1$.

$$T^{\mu\nu} = \eta^{\mu\nu} p(T) + u^\mu u^\nu (\epsilon(T) + p(T))$$

- 4 parameters; system closed by conservation laws. For nonrelativistic u :

$$\partial_t \vec{v} + (\vec{v} \cdot \vec{\partial}) \vec{v} = -\frac{1}{\rho} \vec{\partial} p, \quad [\text{Euler's equations}]$$

$$\partial_t \rho + \rho \vec{\partial} \cdot \vec{v} + \vec{v} \cdot \vec{\partial} \rho = 0.$$

[notice p acts like a potential: ∂p drives acceleration]

Hydrodynamics, III

- If T, u are not constant (but still slowly varying on mfp. scale), there will be small corrections.
- Most general first order correction: (in local rest frame)

$$T_{ij} = \delta_{ij}p - \eta \left[\partial_i u_j + \partial_j u_i - \frac{2}{3} \delta_{ij} \partial_k u_k \right] - \zeta \delta_{ij} \partial_k u_k \quad (+\mathcal{O}(\partial^2))$$

- *shear* and *bulk* viscosities

Hydrodynamics, IV

- Linearized perturbations give *sound waves*, but also *diffusive modes*

$$\left(-i\omega + \frac{\eta}{\epsilon + p} k_z^2\right) u_x = 0$$

- These are acausal. This is not a conceptual problem, but *in practice*, causes numerical instabilities

Hydrodynamics, V

- The solution is to go to *second order*
- This removes the acausality and stabilize the dissipative effects
- Can be understand from Müller-Israel-Stewart theory

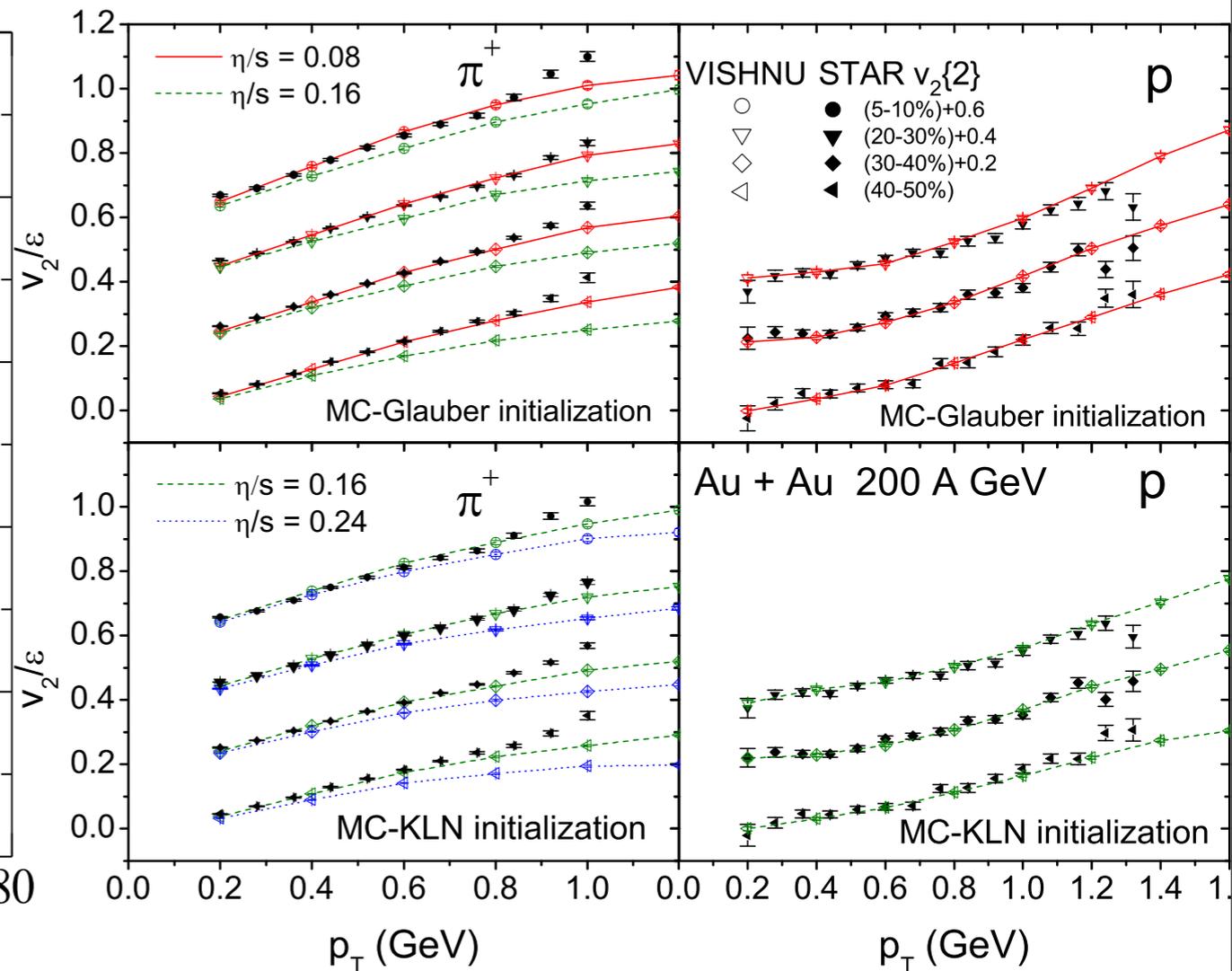
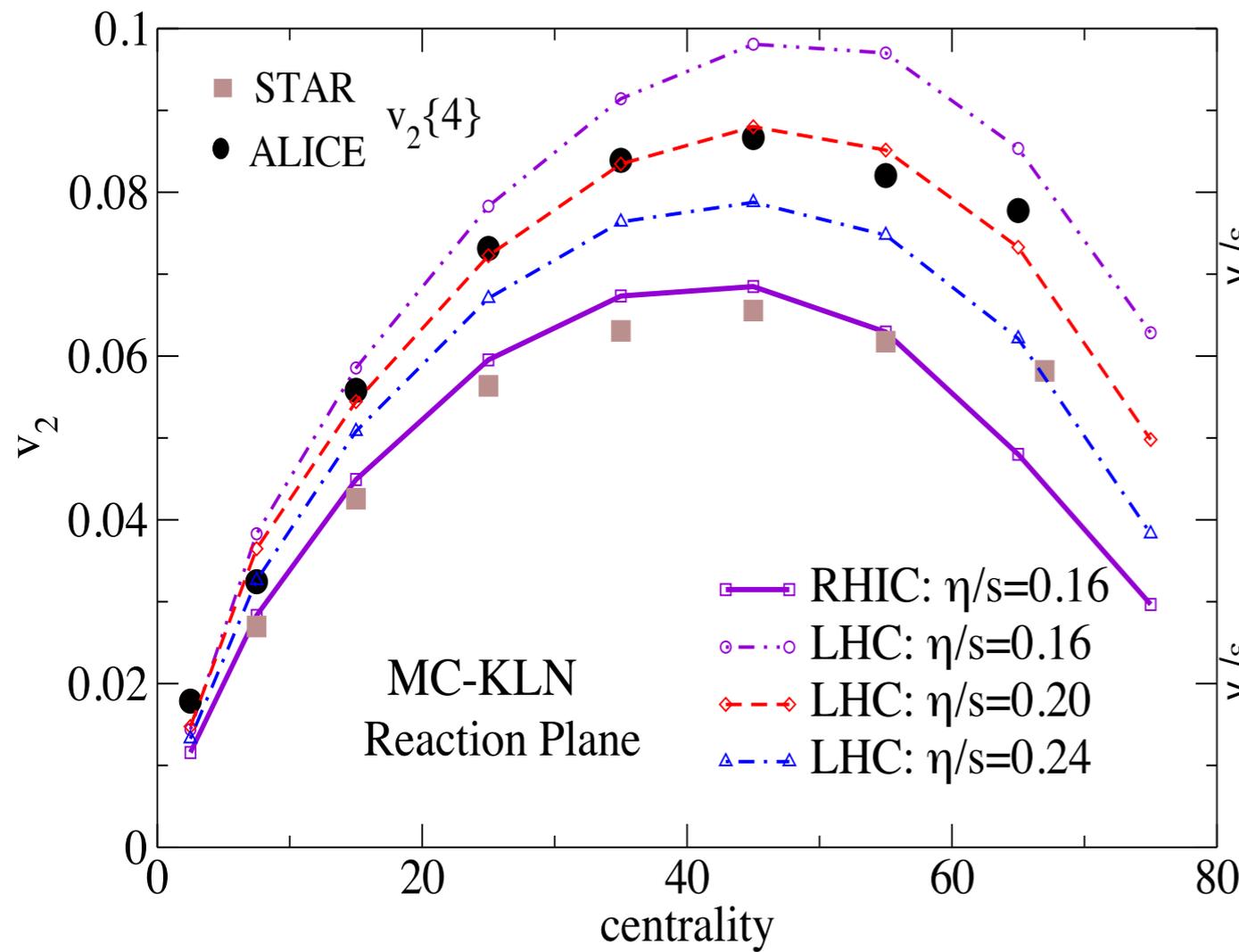
$$\pi_{ij} = -\eta[\partial_i u_j + \dots] \rightarrow -\eta[\partial_i u_j + \dots] - \tau_\pi u \cdot \partial \pi_{ij}$$

- The linearized shear mode becomes:

$$-i\omega + \frac{k_z^2}{1 - i\omega\tau_\pi} \frac{\eta}{\epsilon + p} = 0$$

Success of hydro

[plots from Heinz,Chen&Song | 108.5323]



$$\frac{dN}{d\phi} \propto 1 + \sum_{n=1} 2v_n(p_T) \cos(n(\phi - \phi_n))$$

present

consensus:

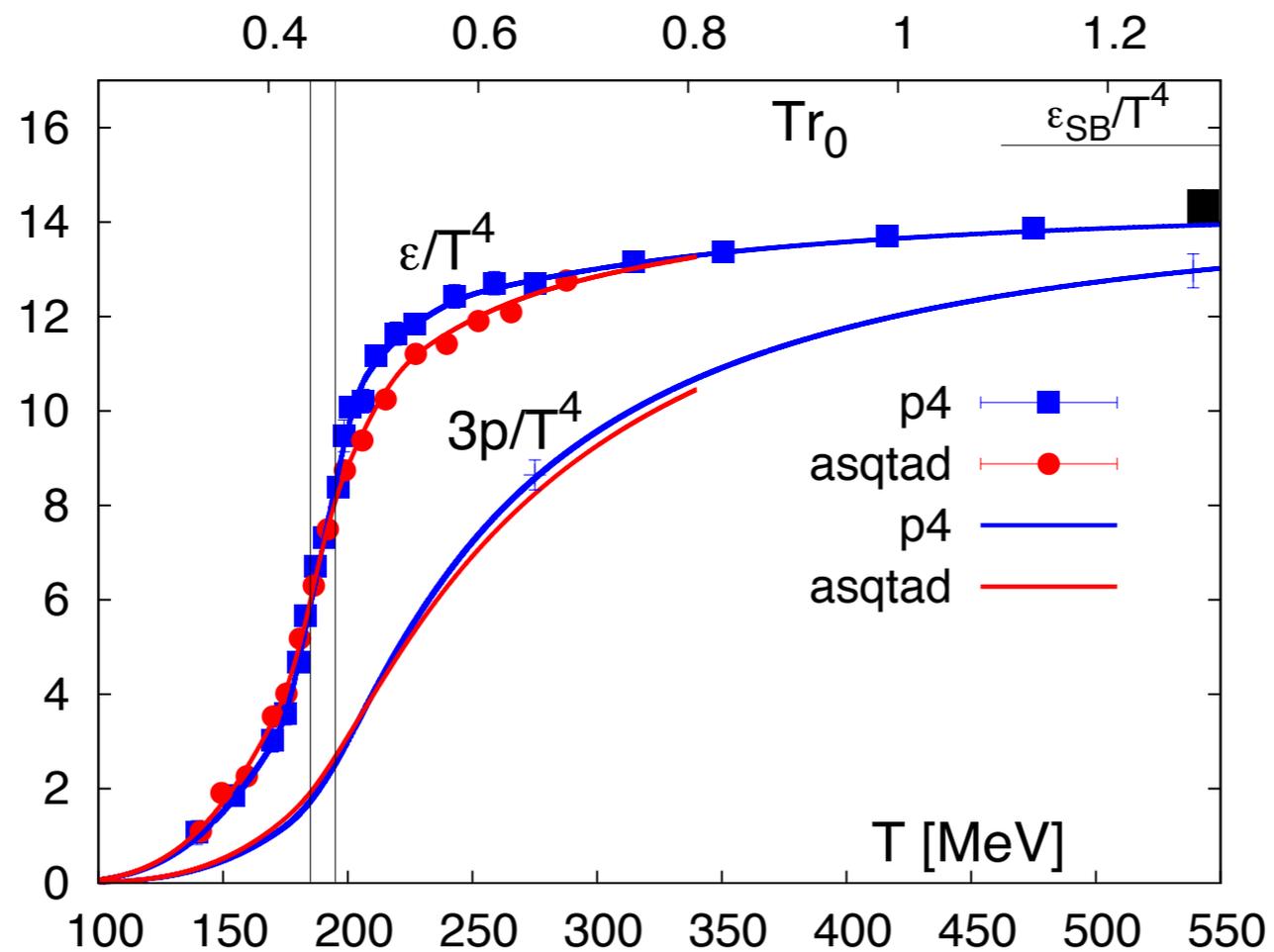
$$4\pi \frac{\eta}{s} < 2.5$$

References

- There are many good textbooks on finite T field theory. I recommend:
 - Kapusta and Gale,
Finite temperature field theory: principles and applications
 - M. LeBellac, *Thermal field theory*
 - A. Das, *Finite temperature field theory*
- I also found the following *lecture notes/reviews* helpful:
 - P. Romatschke, 0902.3663
 - B. Müller et al., 1202.3233
 - E. Iancu, 1205.0579
 - U. Heinz, 0901.4355

[extra:]

- The QCD equation of state (2+1 flavors)



(Karsch *et al*,
0903.4379)