Outline	Introduction	Model	Results	Summary

## Backreaction in Swiss Cheese models

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# Winter school for particle physics and cosmology, Skeikampen, 4.1.2013

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## Outline

- Introduction and motivation
- Model construction
- Results
- Summary

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► The Einstein equation is

$$G_{\mu\nu}=T_{\mu\nu}$$

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and the Einstein tensor is non-linear,

$$\langle G_{\mu\nu}(g_{\alpha\beta}) \rangle \neq G_{\mu\nu}(\langle g_{\alpha\beta} \rangle).$$

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- In cosmology we use (almost) always averaged Friedmann-Robertson-Walker metric.
- ► Is this correct?

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## Backreaction

The Hamiltonian constraint between expansion rate θ, energy density ρ, curvature <sup>(3)</sup>R and shear σ is

$$\frac{1}{3}\theta^2 = 8\pi G_{\rm N}\rho - \frac{1}{2}{}^{(3)}R + \sigma^2$$

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► Averaging both sides we get one of the Buchert equations

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► Like Friedmann's equation, except an extra term

$$\mathcal{Q}\equivrac{2}{3}\left(\langle heta^2
angle-\langle heta
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ight)-2\langle \sigma^2
angle$$

appears due to averaging

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Lemaître-Tolman-Bondi metric:

 $ds^{2} = -dt^{2} + X^{2}(t, r)dr^{2} + R^{2}(r, t)(d\theta^{2} + \sin^{2}\theta d\phi^{2})$ 

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- $\blacktriangleright \ \rightarrow \text{Darmois junction conditions}$

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LTB solution describes pressureless fluid

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- LTB solution describes pressureless fluid
- Expansion rate varies with r
- If the inside of an LTB region expands faster than outside, at some point the radial dust shells will overlap, generating a shell crossing singularity
- This combined with Darmois junction conditions restricts the models a lot

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One can show that under reasonably physical conditions, backreaction is always small in Swiss Cheese models

• The space has a regular symmetry center at r = 0

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- The areal radius is monotonic,  $\frac{\partial R(t,r)}{\partial r} \ge 0$ .

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- ► There is a big bang singularity at t = t<sub>B</sub>(r) ≤ 0, and there are no other singularities at least until time t = t<sub>0</sub>, with t<sub>0</sub> - t<sub>B</sub>(r) ~ t<sub>0</sub>.

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- The space has a regular symmetry center at r = 0
- The areal radius is monotonic,  $\frac{\partial R(t,r)}{\partial r} \ge 0$ .
- There is a big bang singularity at  $t = t_{\rm B}(r) \leq 0$ , and there are no other singularities at least until time  $t = t_0$ , with  $t_0 - t_{\rm B}(r) \sim t_0$ .
- The spacetime matches smoothly to a FRW dust universe at the boundary.
- The areal radius at the matching surface is small compared to the spacetime curvature radius of the background FRW universe.

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Our model				

▶ What happens when you break  $\frac{\partial R(t,r)}{\partial r} \ge 0$ ?

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- ▶ This induces a 2+1-dimensional thin singular shell on the surface

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## Our model

- What happens when you break  $\frac{\partial R(t,r)}{\partial r} \ge 0$ ?
- When  $\frac{\partial R(t,r)}{\partial r} = 0$ , there is a discontinuity in the extrinsic curvature.
- This induces a 2+1-dimensional thin singular shell on the surface
- Can we suspend our disbelief and construct a model with significant backreaction?

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## Not really

Having significant backreaction requires collapsing regions as

$$\mathcal{Q}\equivrac{2}{3}\left(\langle heta^2
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ight)-2\langle \sigma^2
angle$$

Collapsing regions require pressure to keep them from collapsing into singularities.

 $\Omega_{\mathcal{Q}} \lesssim 1\%$ 

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 It is possible to construct inhomogeneous spacetimes that solve Einstein equations exactly.

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- However, the junction conditions and having no pressure restrict backreaction to always be small in these models if the inhomogeneities are small.

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- It is possible to construct inhomogeneous spacetimes that solve Einstein equations exactly.
  - However, the junction conditions and having no pressure restrict backreaction to always be small in these models if the inhomogeneities are small.
  - Even if we break one of the junction conditions, we still could not have large backreaction
  - It appears likely that more complicated models are needed for estimation of backreaction.