

# Numerical Simulations of Relativistic Multifluids

I. Hawke

with thanks to Nils Andersson, Greg Comer, Cesar Lopez-Monsalvo and Lars Samuelsson

School of Mathematics,  
University of Southampton, UK

MICRA2009, 26 August 2009

## 1 Multifluids

- Introduction
- Framework
- Two-stream instabilities
- Nonlinear code
- Implementation issues
- 1d results
- 2d results

## 2 Summary

## **Punchline:**

Multifluid models may be useful or essential for representing physics such as superfluidity. However, it is computationally expensive and the additional instabilities are messy.

- Multiple particle species will be important in NSs.
- For certain phenomena (e.g. superfluids, heat conduction) a single ideal fluid approximation is insufficient.
- Multiple fluids occupying the same volume may give a sufficient model;
  - ▶ Standard hydrodynamic model of superfluids is inherently multifluid;
  - ▶ Some causal heat conduction models can be derived from multifluid approaches.

# Introduction

- Multiple particle species will be important in NSs.
- For certain phenomena (e.g. superfluids, heat conduction) a single ideal fluid approximation is insufficient.
- Multiple fluids occupying the same volume may give a sufficient model;
  - ▶ Standard hydrodynamic model of superfluids is inherently multifluid;
  - ▶ Some causal heat conduction models can be derived from multifluid approaches.

The key physics question is determining what the fluids represent.

That question is ignored here; instead look at numerical methods for a toy problem.

Carter introduced a Hamiltonian framework based on species number currents  $n_X^\nu$  and a master function (“energy”)  $\Lambda \equiv \Lambda(n_X^\alpha, n_\alpha^Y)$ .

- Action principle varying  $\Lambda$ ; clean coupling to GR.
- Conjugate momenta  $\mu_\nu^X$  need not be parallel to  $n_X^\nu$ ; *entrainment*.
- Covariant continuity and Euler equations follow:

$$\begin{aligned}\nabla_\mu n_X^\mu &= \Gamma^X, \\ 2n_X^\mu \nabla_{[\mu} \mu_{\nu]}^X &= f_\nu^X.\end{aligned}$$

- Typically assume inter-species forces balance,  $\sum_X f_\nu^X = 0$ .
- Even in simple case with no species production ( $\Gamma^X = 0$ ) and no forces the fluids can still interact through the entrainment.

# Carter Framework

Carter introduced a Hamiltonian framework based on species number currents  $n_X^\nu$  and a master function (“energy”)  $\Lambda \equiv \Lambda(n_X^\alpha, n_\alpha^Y)$ .

- Action principle varying  $\Lambda$ ; clean coupling to GR.
- Conjugate momenta  $\mu_\nu^X$  need not be parallel to  $n_X^\nu$ ; *entrainment*.
- Covariant continuity and Euler equations follow:

$$\begin{aligned}\nabla_\mu n_X^\mu &= \Gamma^X, \\ 2n_X^\mu \nabla_{[\mu} \mu_{\nu]}^X &= f_\nu^X.\end{aligned}$$

- Typically assume inter-species forces balance,  $\sum_X f_\nu^X = 0$ .
- Even in simple case with no species production ( $\Gamma^X = 0$ ) and no forces the fluids can still interact through the entrainment.

Carter introduced a Hamiltonian framework based on species number currents  $n_X^\nu$  and a master function (“energy”)  $\Lambda \equiv \Lambda(n_X^\alpha, n_\alpha^Y)$ .

- Action principle varying  $\Lambda$ ; clean coupling to GR.
- Conjugate momenta  $\mu_\nu^X$  need not be parallel to  $n_X^\nu$ ; *entrainment*.
- Covariant continuity and Euler equations follow:

$$\begin{aligned}\nabla_\mu n_X^\mu &= \Gamma^X, \\ 2n_X^\mu \nabla_{[\mu} \mu_{\nu]}^X &= f_\nu^X.\end{aligned}$$

- Typically assume inter-species forces balance,  $\sum_X f_\nu^X = 0$ .
- Even in simple case with no species production ( $\Gamma^X = 0$ ) and no forces the fluids can still interact through the entrainment.



Carter introduced a Hamiltonian framework based on species number currents  $n_X^\nu$  and a master function (“energy”)  $\Lambda \equiv \Lambda(n_X^\alpha, n_\alpha^Y)$ .

- Action principle varying  $\Lambda$ ; clean coupling to GR.
- Conjugate momenta  $\mu_\nu^X$  need not be parallel to  $n_X^\nu$ ; *entrainment*.
- Covariant continuity and Euler equations follow:

$$\begin{aligned}\nabla_\mu n_X^\mu &= \Gamma^X, \\ 2n_X^\mu \nabla_{[\mu} \mu_{\nu]}^X &= f_\nu^X.\end{aligned}$$

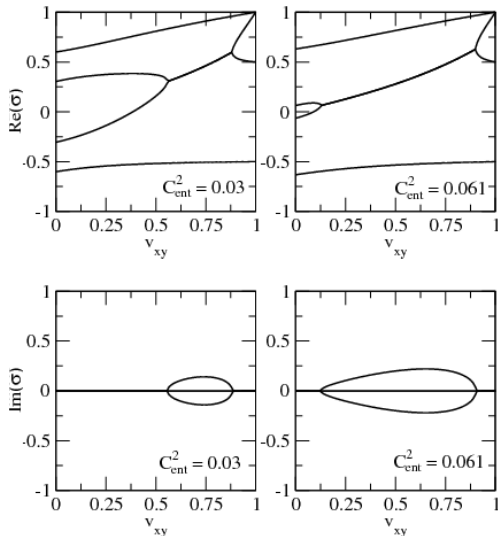
- Typically assume inter-species forces balance,  $\sum_X f_\nu^X = 0$ .
- Even in simple case with no species production ( $\Gamma^X = 0$ ) and no forces the fluids can still interact through the entrainment.

# Two-stream instabilities

Linear analysis shows instabilities when the *relative flow* is large and entrainment occurs.

Plane parallel waves on flat backgrounds without boundaries are checked.

The dispersion relation predicts high frequency modes will explode.



Samuelsson et al, arXiv:0906.4002

# Nonlinear code

As a first attempt we look at the simplest case:

- No species production or forces.
- “Generalized polytrope” EOS (Prix et al, PRD71 043005 (2005), Samuelsson et al, arXiv:0906.4002).
- Periodic boundaries, 1+1 or 2+1d, special relativity.
- Central differencing with Kreiss-Oliger dissipation.

The equations become

$$\partial_t n_X^t = -\partial_j n_X^j,$$
$$\partial_t \mu_i^X = \partial_i \mu_t^X + 2 \sum_j \frac{n_X^j}{n_X^t} \partial_{[ij} \mu_{j]}^X.$$

# Nonlinear code

As a first attempt we look at the simplest case:

- No species production or forces.
- “Generalized polytrope” EOS (Prix et al, PRD71 043005 (2005), Samuelsson et al, arXiv:0906.4002).
- Periodic boundaries, 1+1 or 2+1d, special relativity.
- Central differencing with Kreiss-Oliger dissipation.

The equations become

$$\begin{aligned}\partial_t n_X^t &= -\partial_j n_X^j, \\ \partial_t \mu_i^X &= \partial_i \mu_t^X + 2 \sum_j \frac{n_X^j}{n_X^t} \partial_{[i} \mu_{j]}^X.\end{aligned}$$

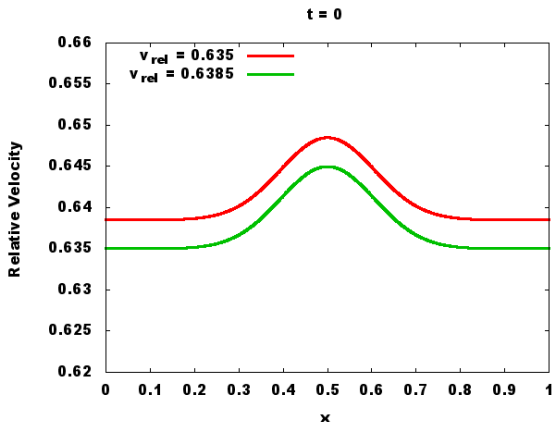
# Implementation issues

The evolved variables  $(n_X^t, \mu_i^X)$  do not immediately contain all the necessary information. All other components of  $n_X^\nu$  must be found from e.g. the master function.

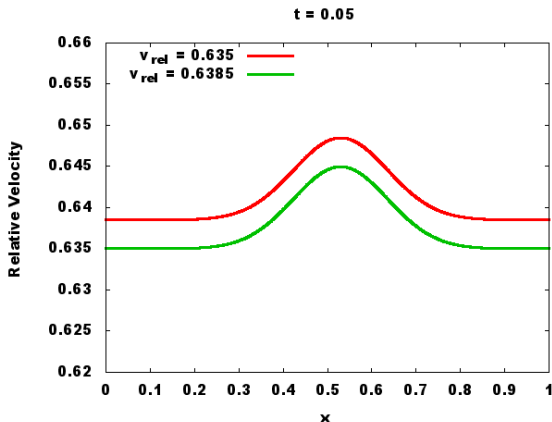
Currently the code

- 1 guesses the value of the scalars  $n_X^\alpha n_\alpha^Y$ ;
- 2 computes  $\Lambda$  and derivatives from these guesses;
- 3 uses the definition of  $\mu_i^X$  to find  $n_X^i$  by solving a linear system;
- 4 uses the resulting approximation to  $n_X^\nu$  to check the scalar guesses, implying a nonlinear root-find.

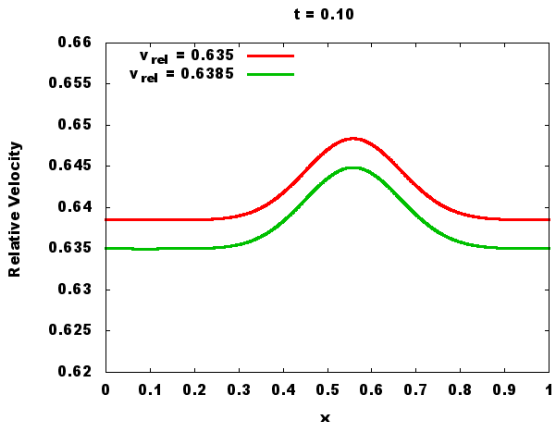
This “primitive variable recovery” requires over 90% of the run-time.



Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.

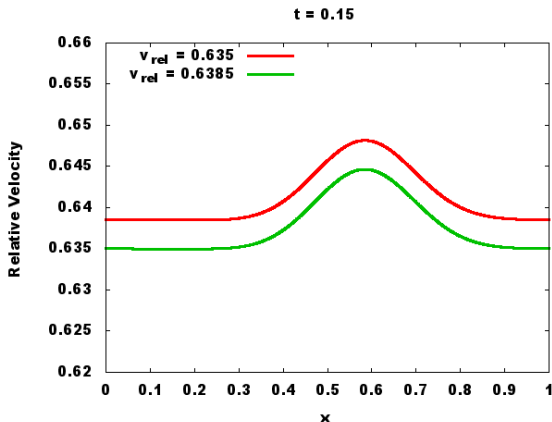


Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.

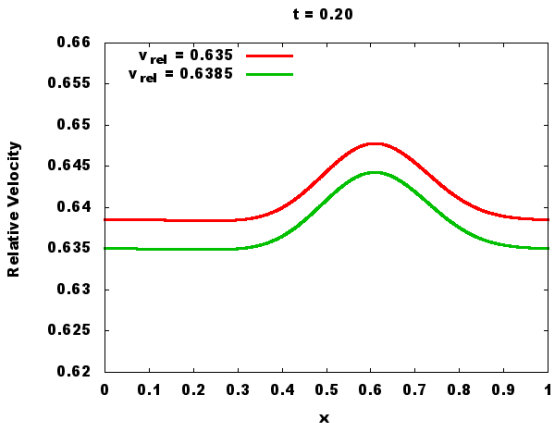


Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.

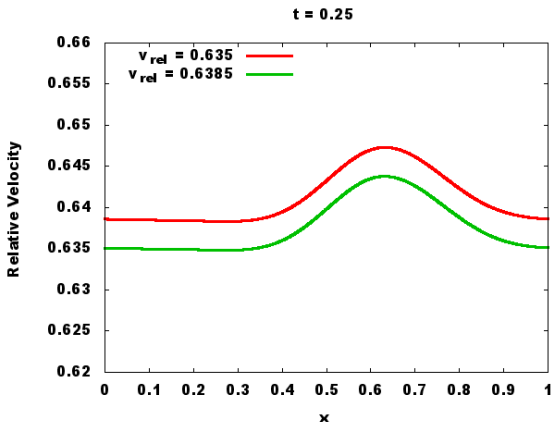




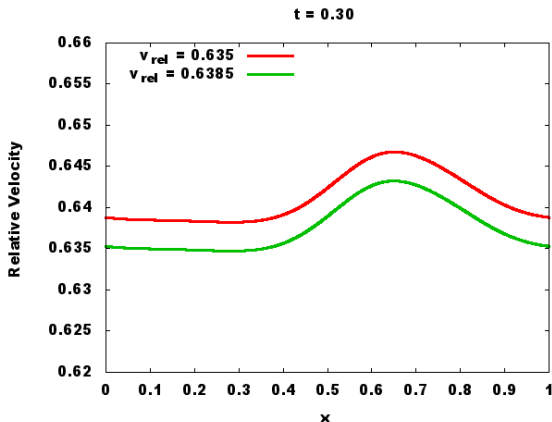
Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.



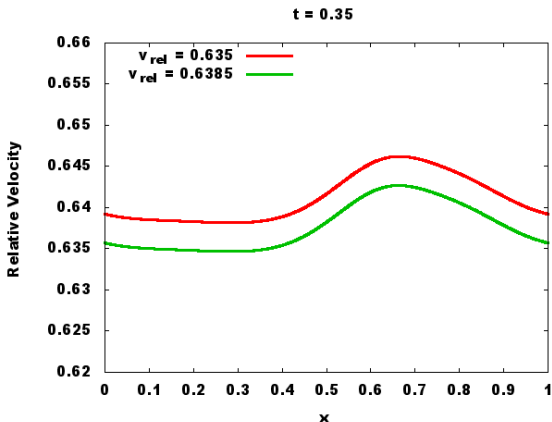
Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.



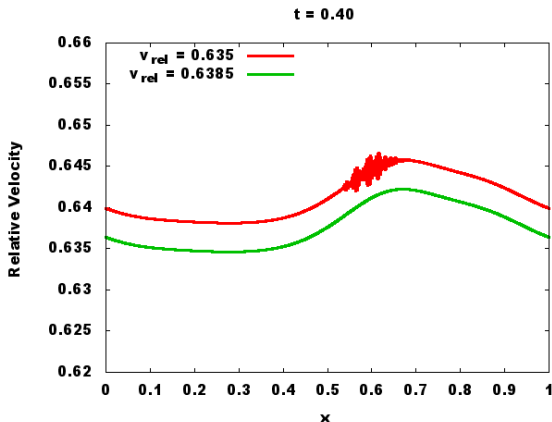
Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.



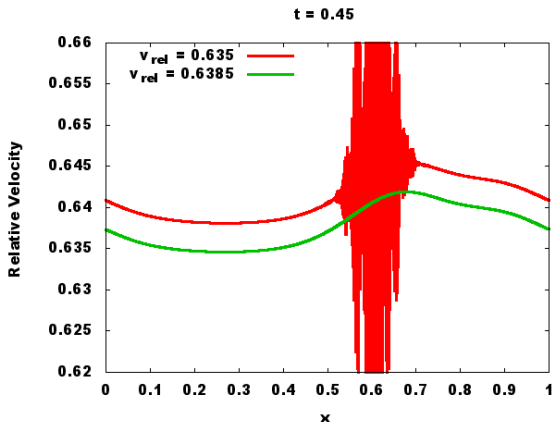
Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.



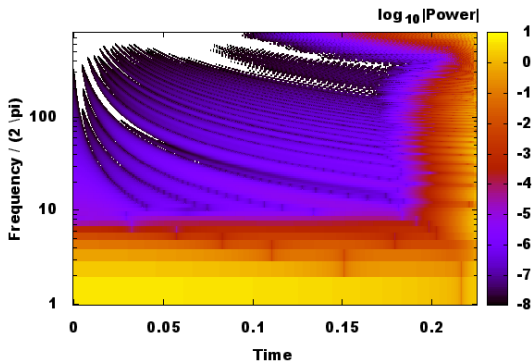
Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.



Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.

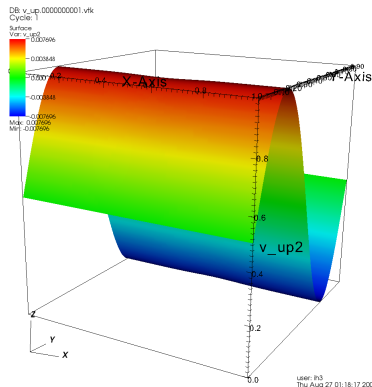
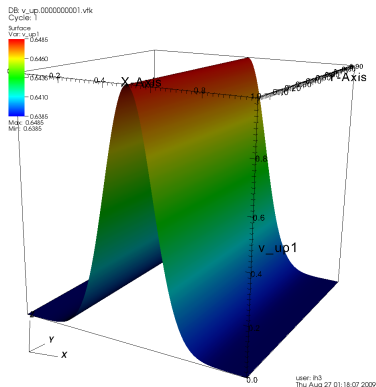


Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.

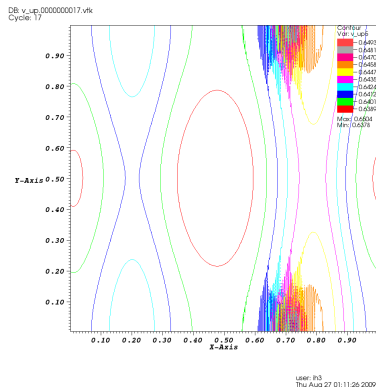
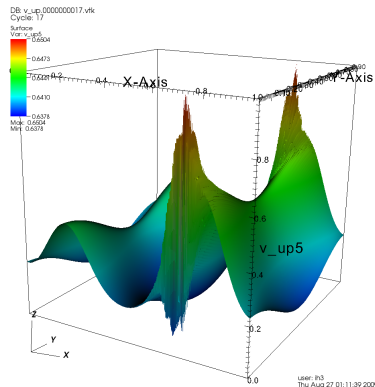


Near the instability limit the fully nonlinear evolutions reproduce the linear analysis. The instability cascades down from high frequency.





It is straightforward to extend to a 2d “shearing box”; perturbing the initial data in the other direction has no effect on the instability.



It is straightforward to extend to a 2d “shearing box”; perturbing the initial data in the other direction has no effect on the instability.

- Multiple (ideal) fluids give a simple framework for modelling complex interactions.
- General conservation law forms are not obvious.
- Numerically expensive to convert between required types of variables.
- Linear instabilities are clear in nonlinear evolutions.