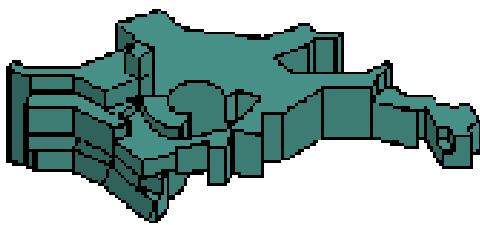


General relativistic hydrodynamics beyond bouncing polytropes



Pablo Cerdá-Durán
Max Planck Institut für Astrophysik
Garching bei München, Germany

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Punch line(s)

To model **rotating** compact objects
we need **MHD and relativity**

To model **MHD and relativity** we need
new physics and new methods

Outline

Motivation

GR (magneto)hydrodynamics

New problems

Conclusions

Motivation

Astrophysical scenarios

Collapse to black hole

Core collapse supernovae

Proto-neutron stars

Pulsar/Magnetar formation

Isolated neutron stars

Instabilities

Asteroseismology

Binary neutron stars

Motivation

Why relativistic?

Shock propagation $\sim 0.1c$

Low density regions

Relativistic jets

Newtonian: Alfvén speed $> c$

Why general relativistic?

Compact objects (BH, NS) = strong gravity

Black holes

Gravitational waves

Status

Gravity				
	Newtonian	Modified Newtonian (TOV-like potentials)	Approximately relativistic (PPN, CFC, CFC+)	Fully relativistic (ADM, BSSN, FCF)
Microphysics	Polytropes Ideal gas			
Tabulated EOS				
Simple neutrino treatment			Ongoing work	Ongoing work
Boltzmann neutrino transport	Ongoing work	Ongoing work	Ongoing work	Ongoing work

(adapted from the table by Harald Dimmelmeier)

Status

Magnetic fields

Ideal MHD

Resistive MHD

Beyond MHD (Hall effect,
ambipolar diffusion)

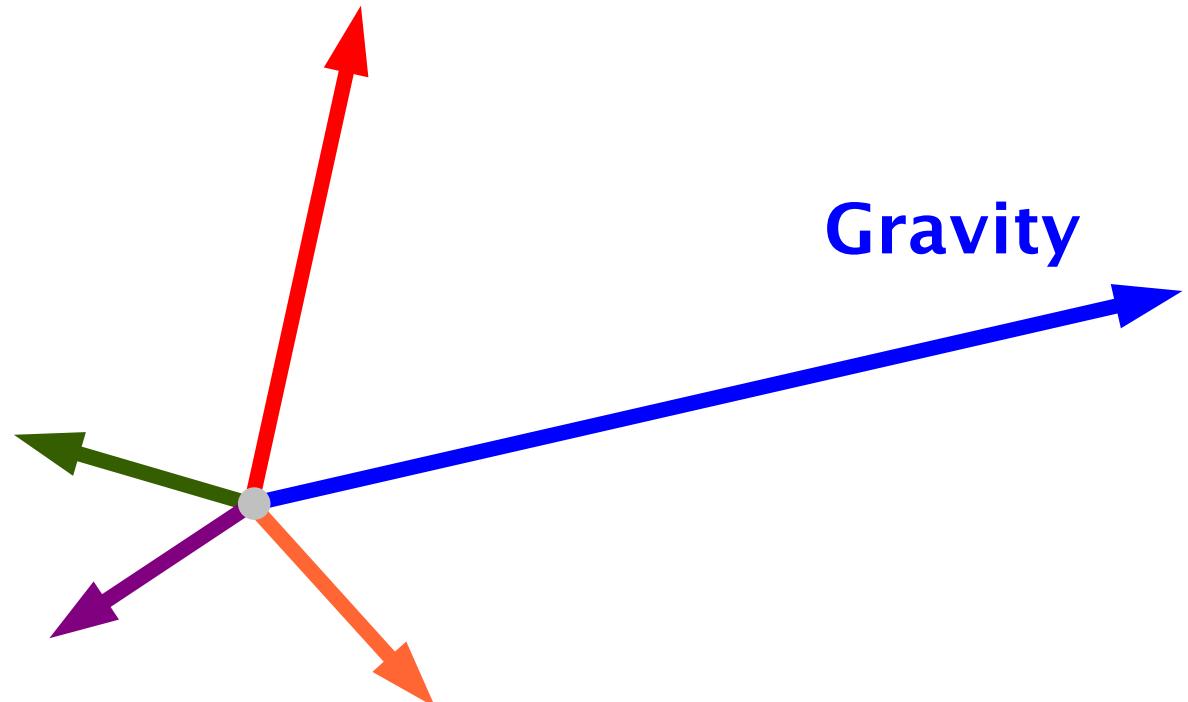
Small scales

Magneto-rotational instability
(MRI)

Turbulence and dynamo

Microphysics

Gravity



**Multi-D
performance**

1D

2D

3D

Framework

3+1 decomposition

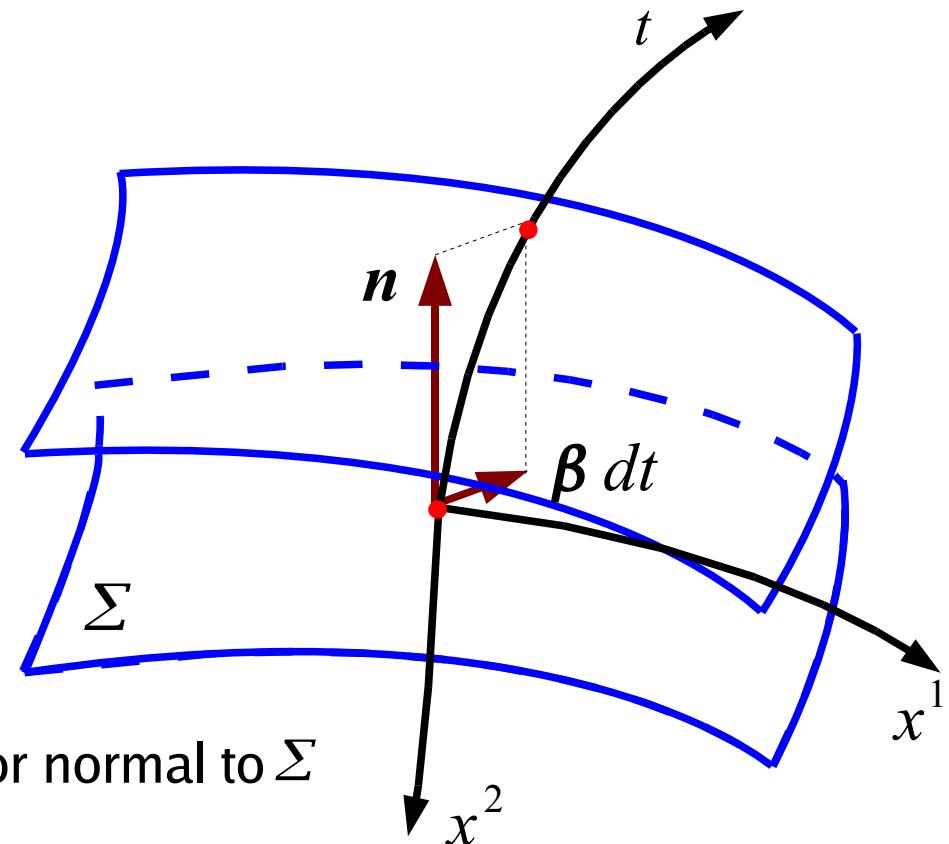
$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = (\beta^i \beta_i - \alpha^2) dt^2 + 2 \beta_i dx^i dt + \gamma_{ij} dx^i dx^j$$

α : lapse

ϕ : conformal factor

β^i : shift vector

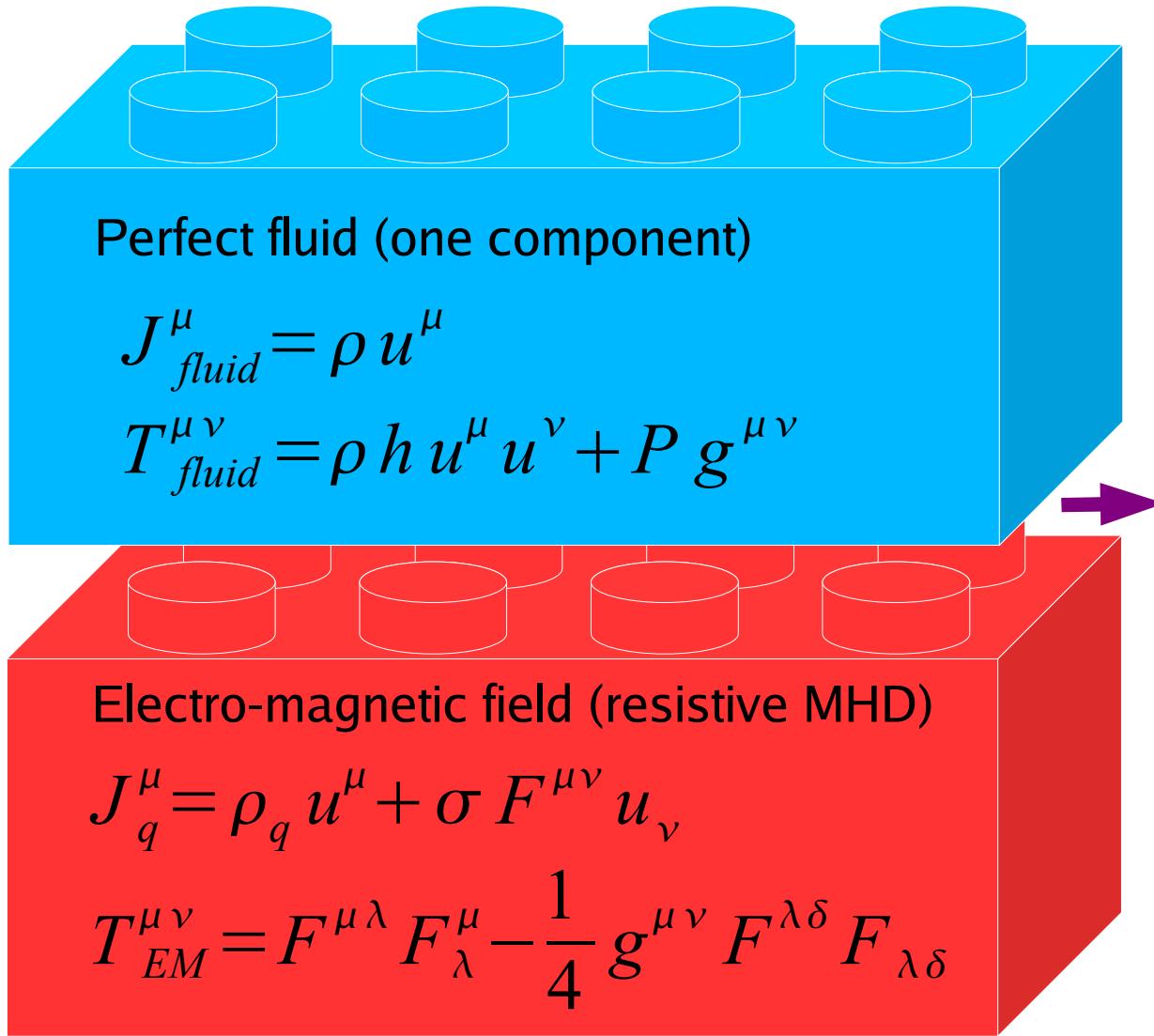
γ_{ij} : 3-metric



$n^\mu n_\mu = -1$: timelike 4-vector normal to Σ

$\perp^\mu_\nu \equiv g^\mu_\nu + n^\mu n_\nu$: projector onto Σ

Energy-momentum tensor



Energy-momentum tensor

$$T^{\mu\nu} = T_{fluid}^{\mu\nu} + T_{EM}^{\mu\nu} + \dots$$

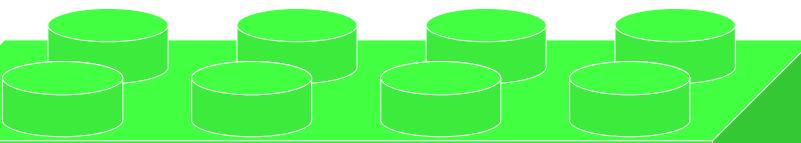
Currents

J_{fluid}^{μ} : Rest-mass current

J_q^{μ} : Charge current

Faraday tensor $F^{\mu\nu}$

Equations



Einstein's equations

$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$

Energy-momentum conservation
(Bianchi identities)

$$T^{\mu\nu}_{;\mu} = 0$$

Continuity (no particle creation)

$$J^\mu_{fluid;\mu} = 0$$

$$J^\mu_{q;\mu} = 0$$

Maxwell's equations

$$*F^{\mu\nu}_{;\nu} = 0$$

$$F^{\mu\nu}_{;\nu} = 4\pi J^\mu_q$$

Energy-momentum tensor

$$T^{\mu\nu} = T^{\mu\nu}_{fluid} + T^{\mu\nu}_{EM} + \dots$$

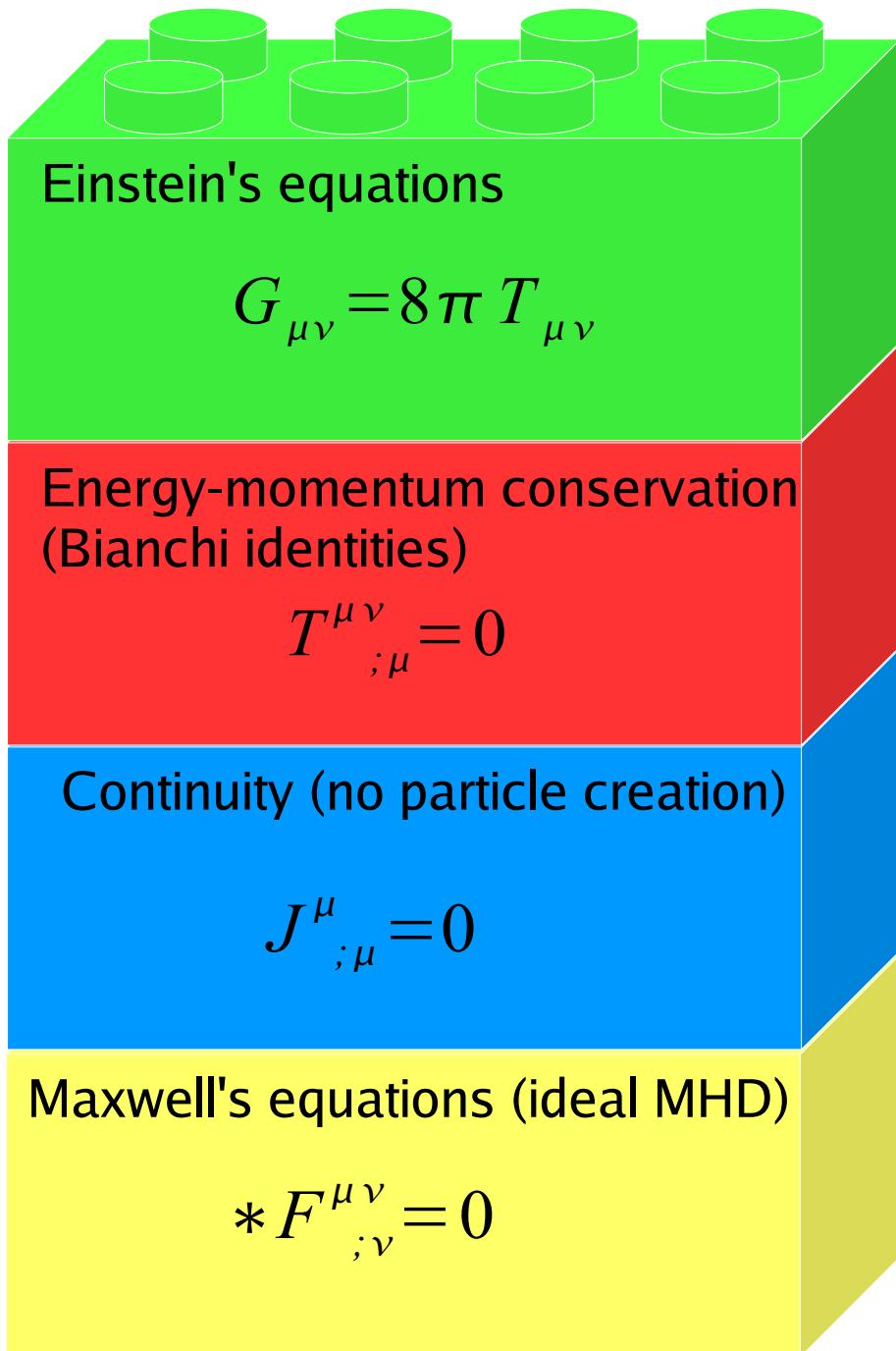
Currents

J^μ_{fluid} : Rest-mass current

J^μ_q : Charge current

Faraday tensor $F^{\mu\nu}$

Equations - Ideal MHD



Energy-momentum tensor

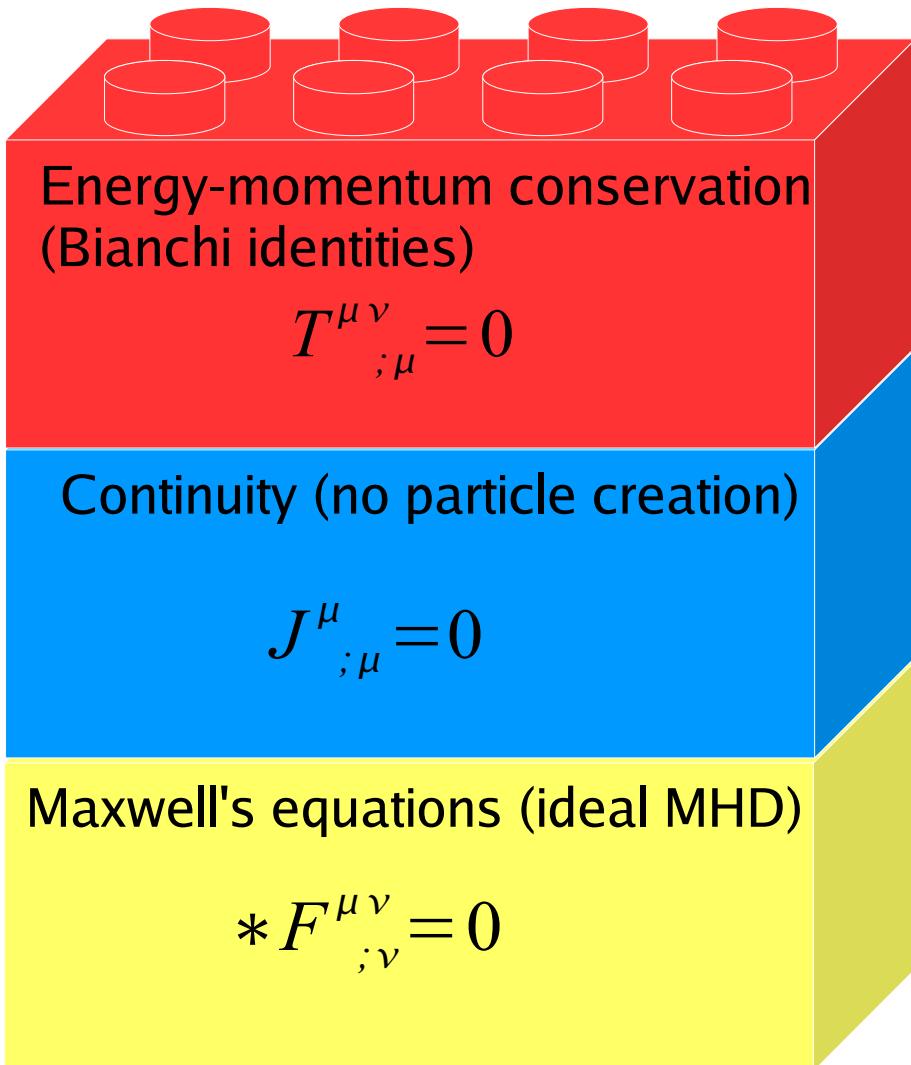
$$T^{\mu\nu} = T^{\mu\nu}_{fluid} + T^{\mu\nu}_{EM} + \dots$$

Currents

J^μ : Rest-mass current

Faraday tensor $F^{\mu\nu}$

Equations – Ideal MHD



Hyperbolic system of
conservation (balance) laws

$$\frac{\partial \hat{\mathbf{U}}}{\partial t} + \nabla_i \hat{\mathbf{F}}^i(\hat{\mathbf{U}}) = \hat{\mathbf{S}}(\hat{\mathbf{U}})$$

$$\nabla_k \hat{B}^k = 0$$

Conserved variables

$$\hat{\mathbf{U}} \equiv \sqrt{\bar{y}}(D, S_i, \tau, B^k)$$

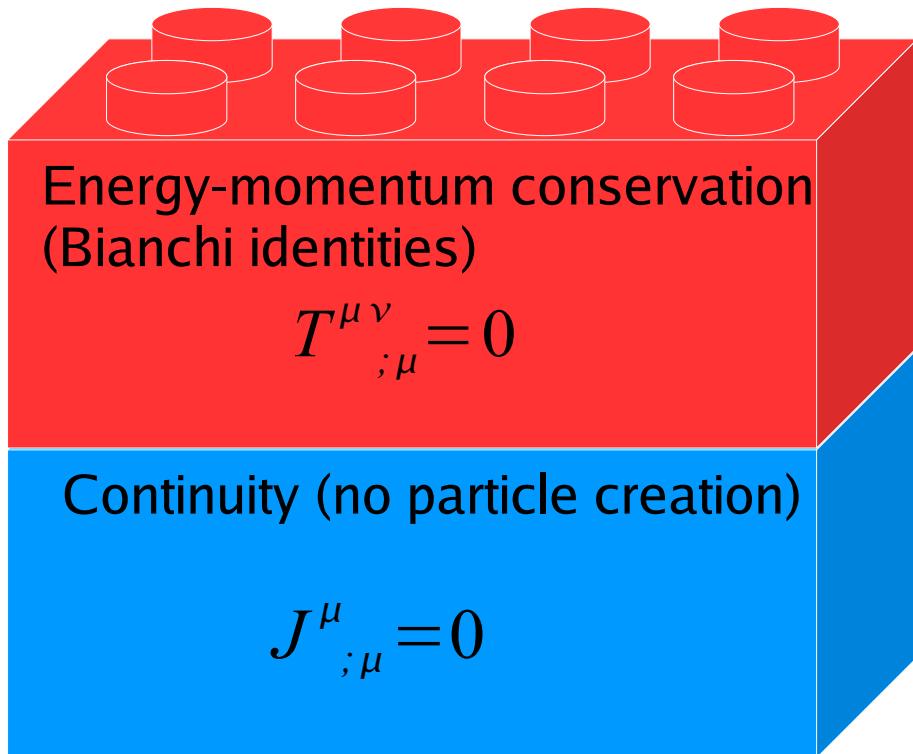
$$= (\hat{D}, \hat{S}_i, \hat{\tau}, \hat{B}^k)$$

$D \equiv -n_\mu J^\mu$: relativistic density

$S_i \equiv -\perp_i^\mu n^\nu T_{\mu\nu}$: momentum

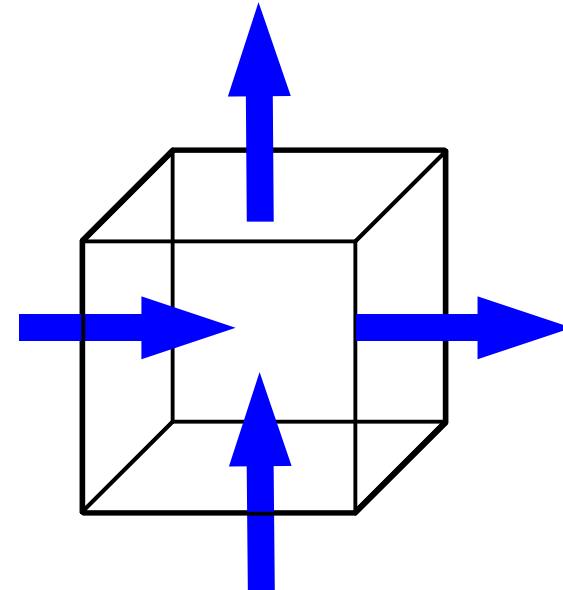
$\tau \equiv n^\mu n^\nu T_{\mu\nu} - D$: energy
(without density)

Equations – Ideal MHD



Conservation of
Number of particles
Momentum
Energy

$$\partial_t \int dV \hat{U}_{fluid} + \int dA \cdot \hat{F} = \int dV \hat{S}$$
$$\hat{U}_{fluid} \equiv (\hat{D}, \hat{S}_i, \hat{\tau})$$

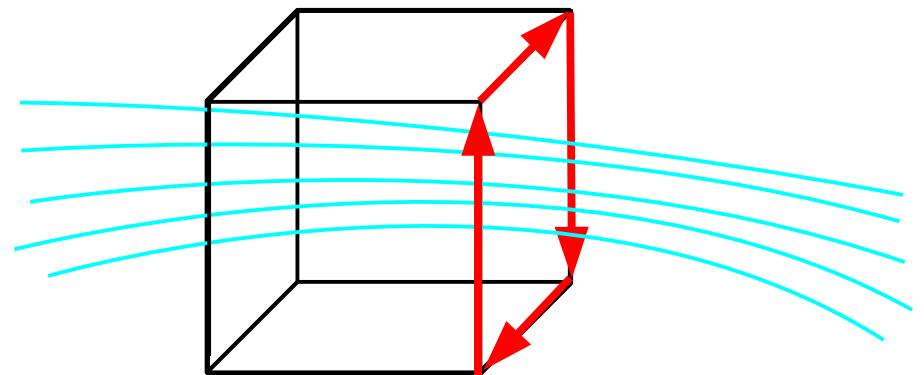


Equations – Ideal MHD

Conservation of magnetic flux

Induction equation $\longrightarrow \partial_t \int dA_k \hat{B}^k + \oint dl \cdot \hat{E} = 0$

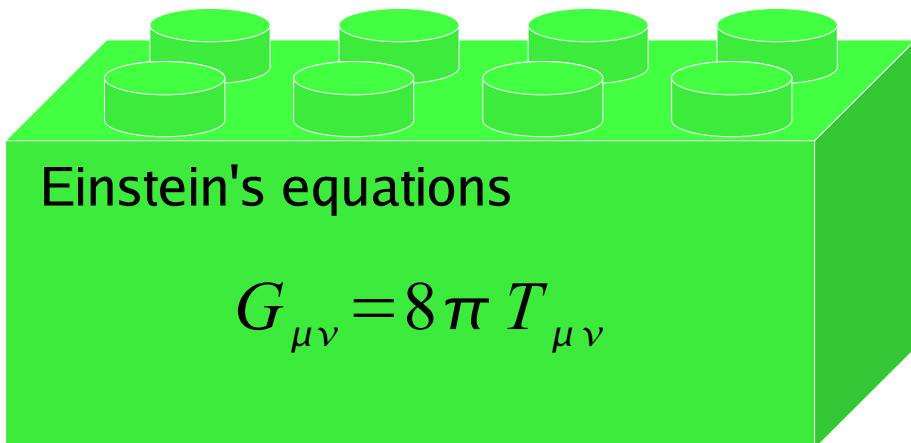
Solenoidal condition $\longrightarrow \oint dA_k \hat{B}^k = 0$



Maxwell's equations (ideal MHD)

$$*F^{\mu\nu}_{;\nu} = 0$$

Equations – space-time



$$\gamma_{ij} = \phi^4 f_{ij} + h_{ij}$$

$$\phi^4 \equiv \text{Tr}(\gamma_{ij})$$

10 variables (α , β^i , ϕ , h_{ij})
 - 4 gauge conditions
 —————
 6 degrees of freedom

Newtonian limit

$$\alpha^2 \simeq 1 - 2U$$

$$\phi^4 \simeq 1$$

$$\beta_i, h_{ij} \simeq 0$$

1 degree of freedom

U : Newtonian potential

Vacuum-Minkowski limit

$$\square h^{ij} = 0$$

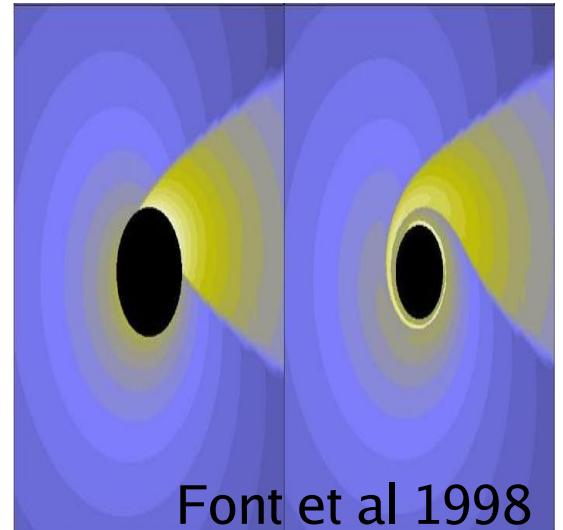
$$\phi, \alpha \simeq 1$$

$$\beta^i \simeq 0$$

2 degree of freedom

h^+ , h^\times : Gravitational waves

Frame-dragging



β_i

CoCoNuT code

General relativistic ideal MHD code

(Dimmelmeier et al. 2002, Cerdá-Durán et al 2008)



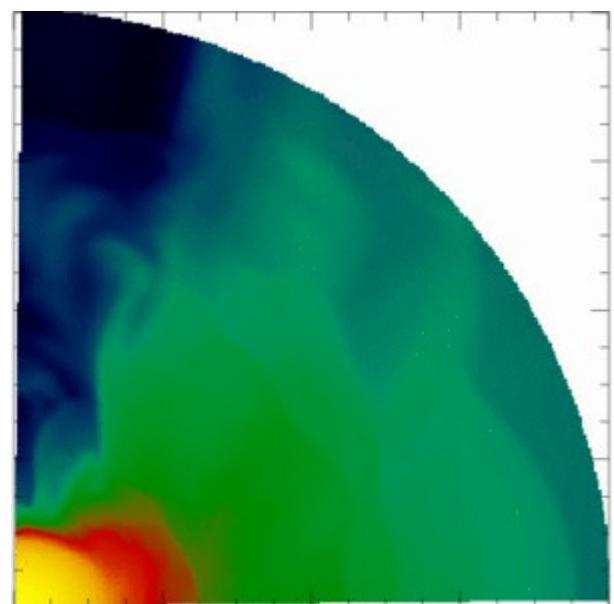
Spherical coordinates: 2D (MHD) and 3D (only HD)

Godunov-type schemes + CT scheme

General relativity: dynamic space-time in the CFC approx. (Isenberg 1979, Wilson 1989) with **spectral methods** (**LORENE**: www.lorene.obspm.fr) (Dimmelmeier et al. 2004)

Microphysics: finite temperature EOS (SHEN, LS) and deleptonization scheme (Liebendörfer 2005)

Applications: core collapse supernovae, isolated neutron star evolution and BH formation.



Some new problems

Microphysics/recovery

Multi-scale problem

3D-efficiency

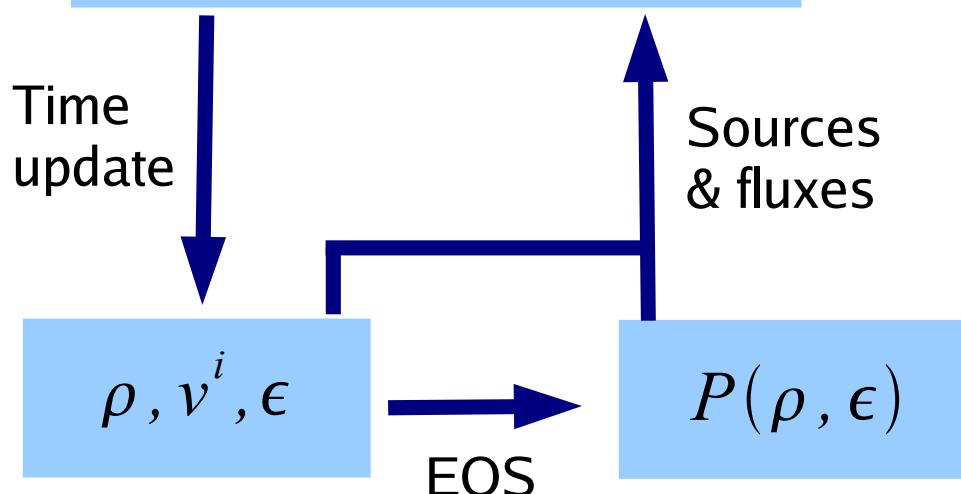
“realistic” MHD

Microphysics/recovery

Micromechanics/recovery

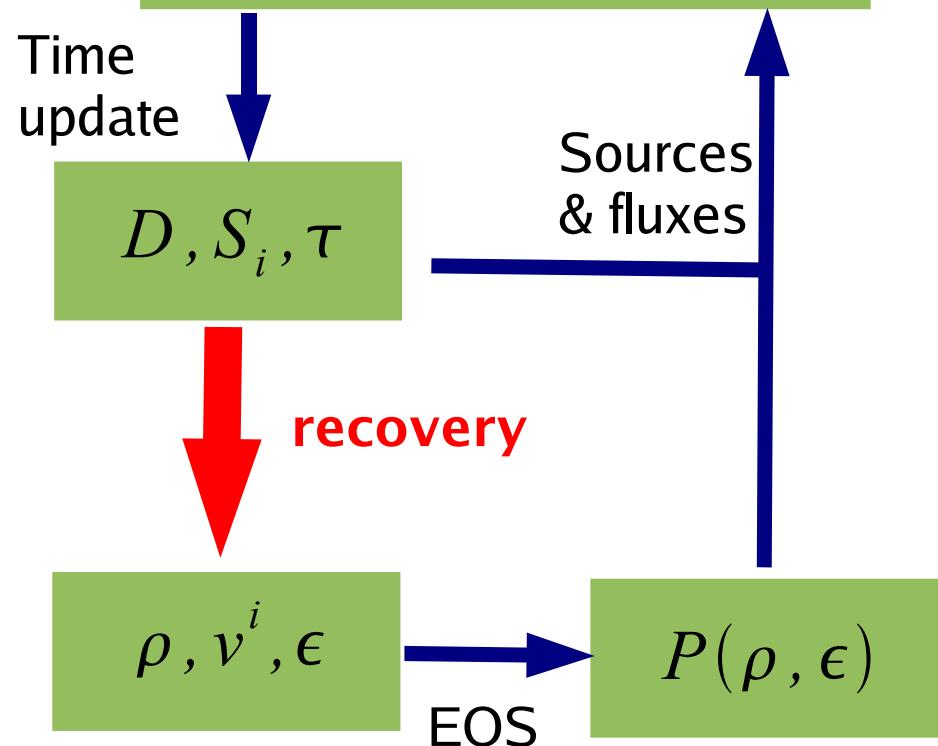
Newtonian

$$\begin{aligned}\partial_t \rho &= S_\rho(\rho, v^i) \\ \partial_t v^i &= S_{v^i}(\rho, v^i, P) \\ \partial_t \epsilon &= S_\epsilon(\rho, v^i, \epsilon, P)\end{aligned}$$



Relativistic

$$\begin{aligned}\partial_t D &= S_D(D, S_i) \\ \partial_t S_i &= S_{S_i}(D, S_i, P) \\ \partial_t \tau &= S_\tau(D, S_i, \tau, P)\end{aligned}$$



Micromechanics/recovery

GR-HD

$$\tau, D, S_i, Y_e$$

$$P - P(\rho, \epsilon, Y_e) = 0$$

Newton
-
Raphson

$$\epsilon = \epsilon(\rho, T, Y_e)$$

$$P$$

$$\rho, v^i, \epsilon$$

GR-MHD

$$\tau, D, S_i, B^k, Y_e$$

$$S^2 = \dots$$
$$\tau = \dots$$
$$\epsilon - \epsilon(\rho, T, Y_e) = 0$$

$$z \equiv \rho h W^2, W, T$$

$$\rho, v^i, \epsilon, B^k$$

Microphysics/recovery

Efficiency: number of iterations?

Use the previous time-step value

Robustness: always converge?

“Safe” recovery values (next slide)

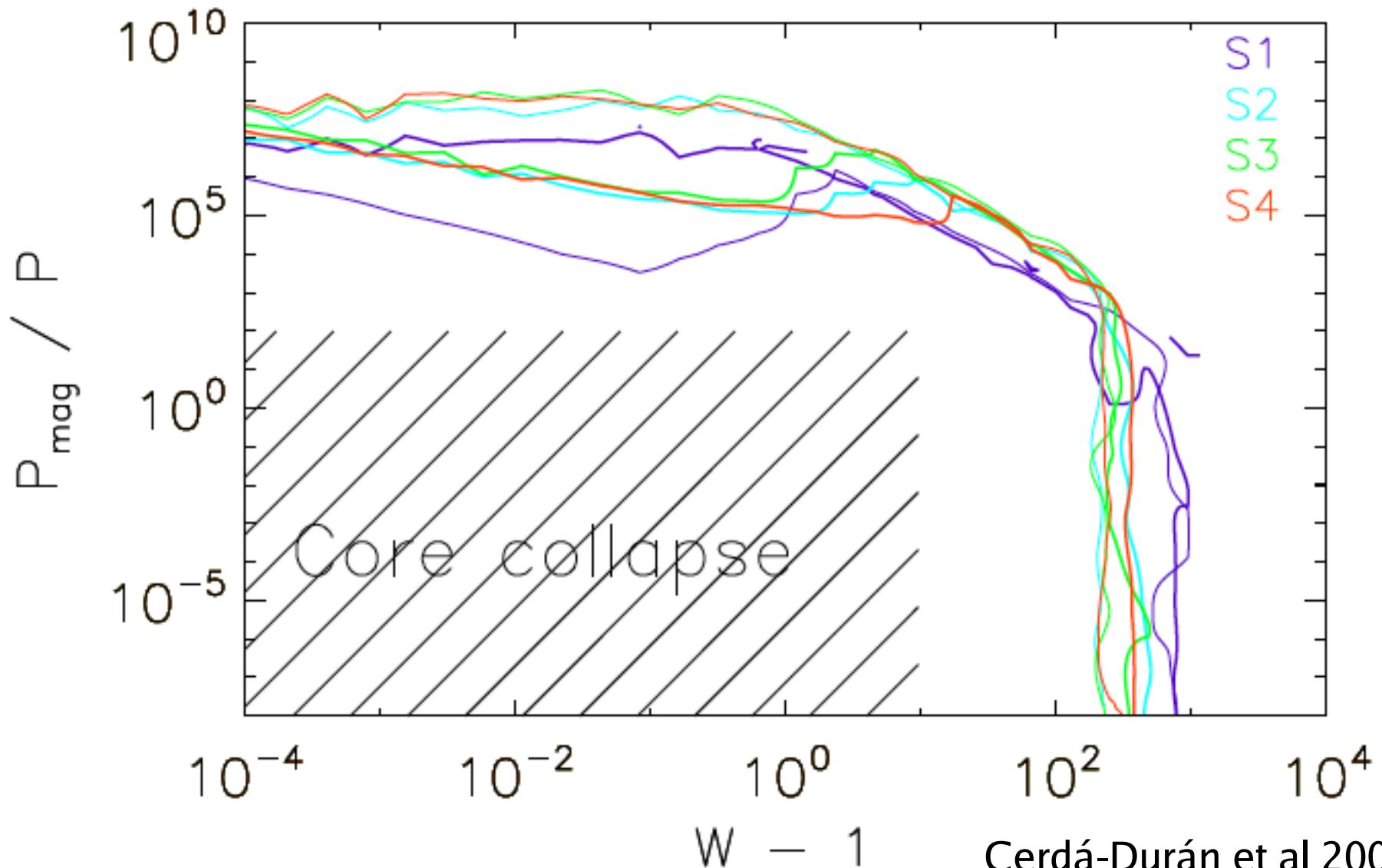
Parallel computation: can be a bottle neck!

Recovery: ~10–20% CPU time

Bad balance: Amdahl's law limits number of CPUs

Microphysics/recovery

“Safe” recovery values (SHEN EOS, collapse and bounce conditions)



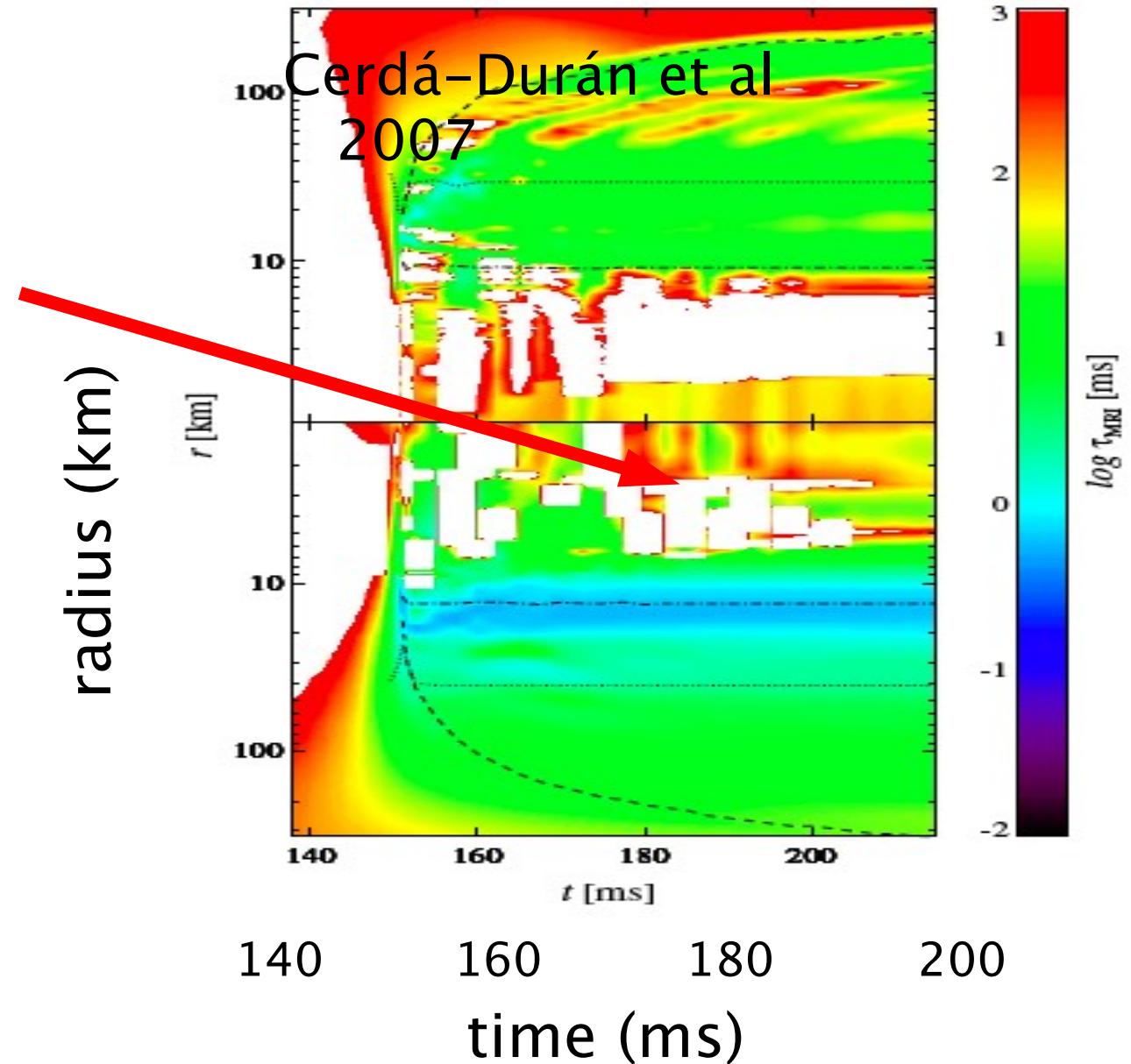
Cerdá-Durán et al 2008

Multi-scale problem

Multi-scale problem - MRI (small scales)

s20 model (Woosley et al 2002)+ rotation ($T/|W|=0.05$)

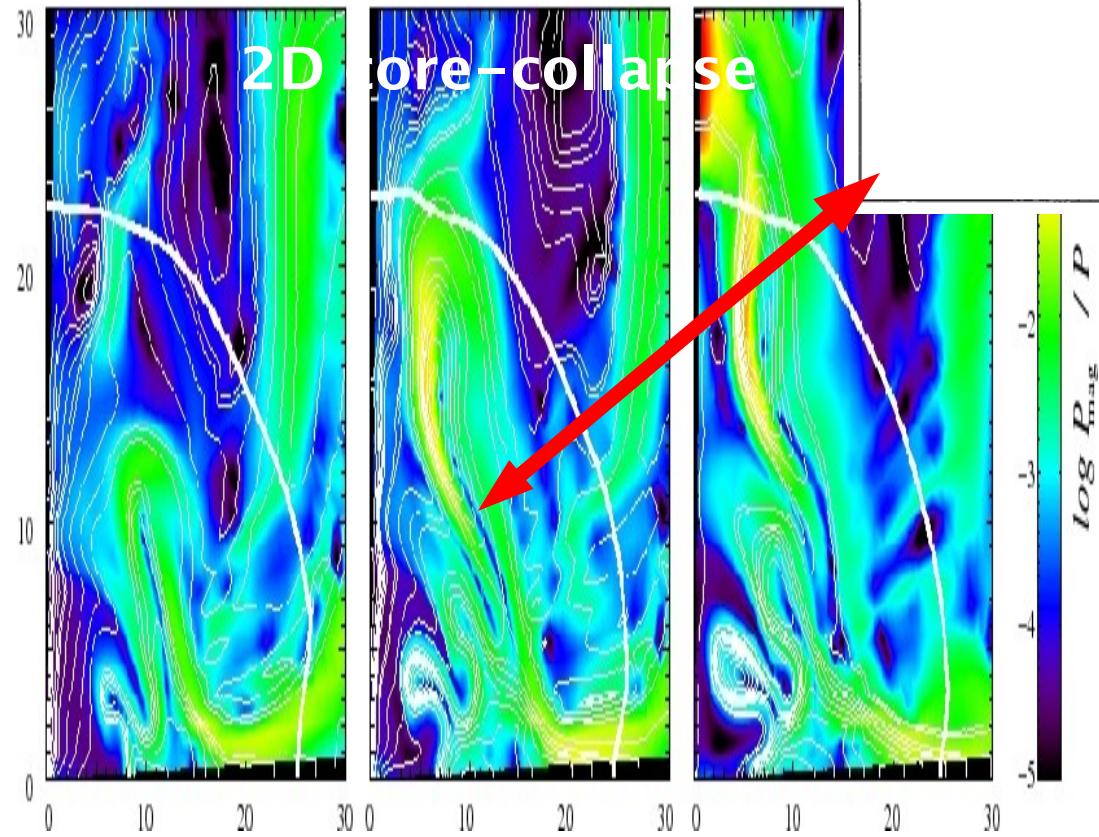
Unstable to
hydromagnetic
instabilities
(MRI, magneto-
convection)
Timescale ~ 1 ms!



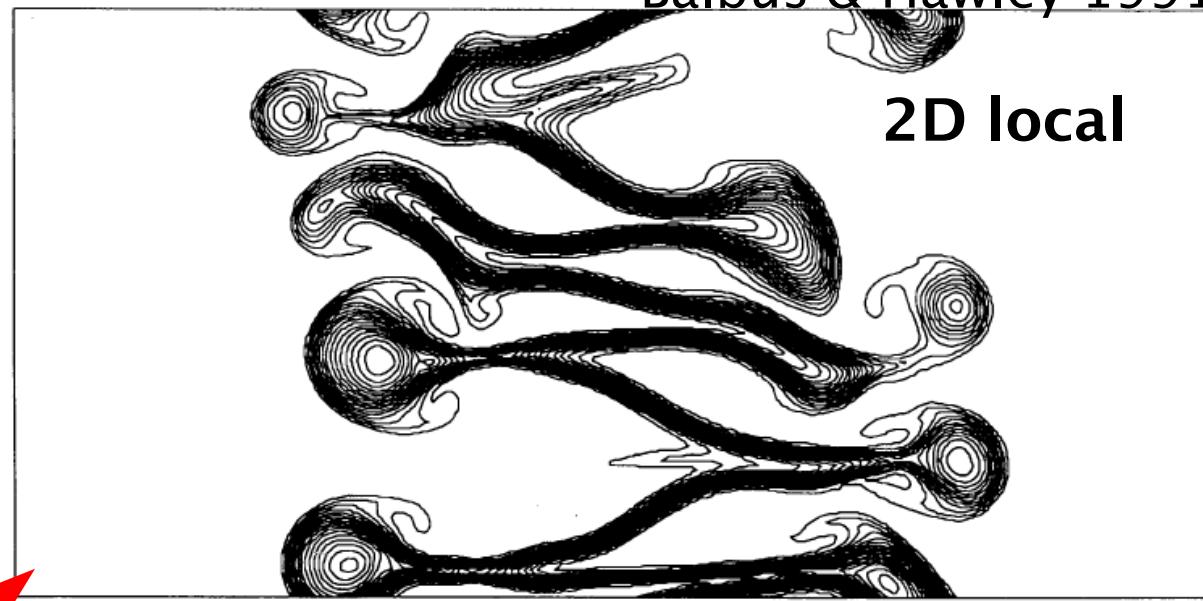
Multi-scale problem - MRI (small scales)

Balbus & Hawley 1991

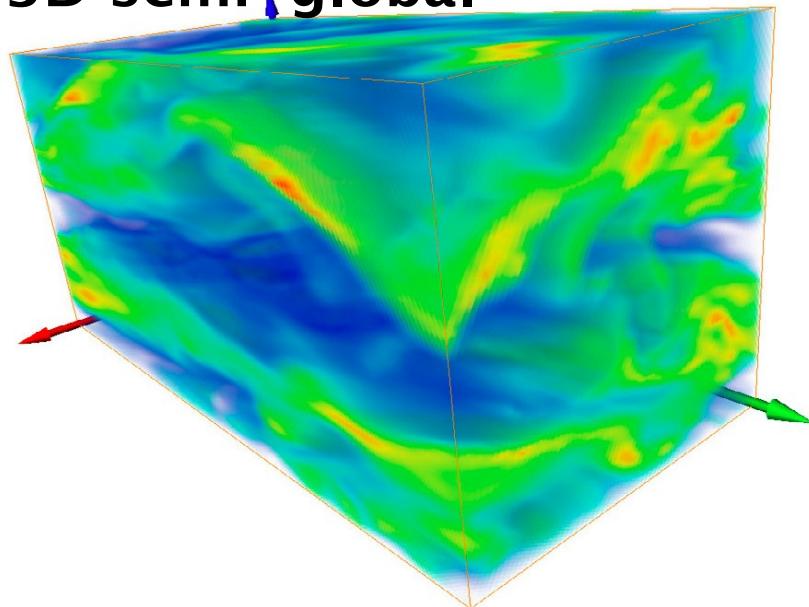
3D problem!



Cerdá-Durán et al 2008

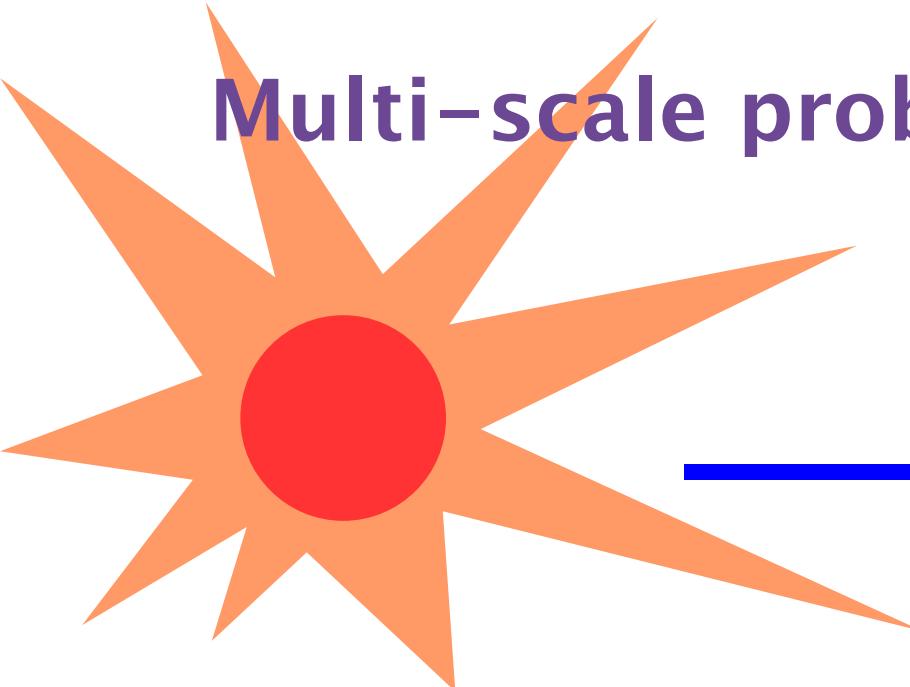


3D semi-global



Obergaulinger et al 2008

Multi-scale problem - long term evolution



Supernovae



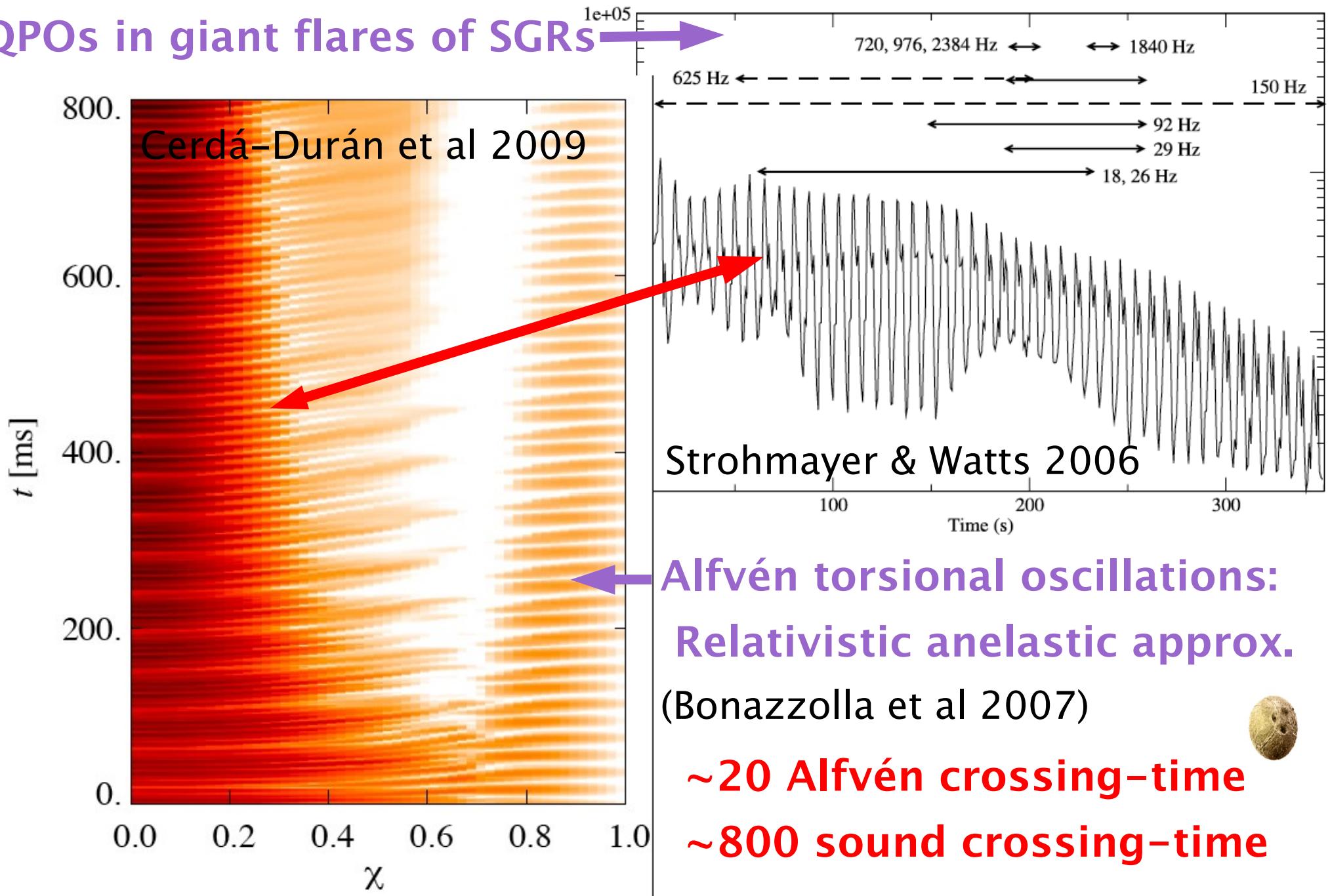
Pulsar
Magnetars
Strange stars
Other NS

**Problems: Dynamo, MRI, convection,
cooling ...**

Methods: Implicit?, anelastic? ...

Multi-scale problem - long term evolution

QPOs in giant flares of SGRs



3D efficiency

Efficient 3D

Cartesian coordinates

- + Easy to implement
- + Best for scenarios without any symmetry (e.g. binary)
- Not possible to recover spherical symmetry and difficult for axisymmetry (Cartoon method)
- Difficult to treat large domains (AMR, FMR)

Spherical polar coordinates

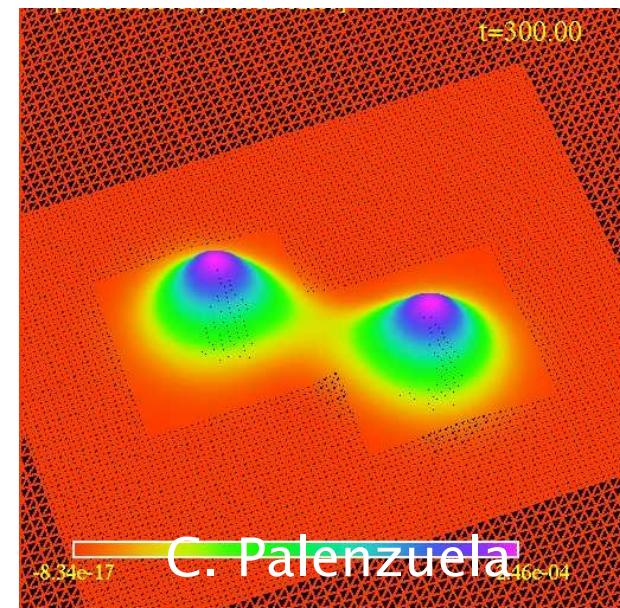
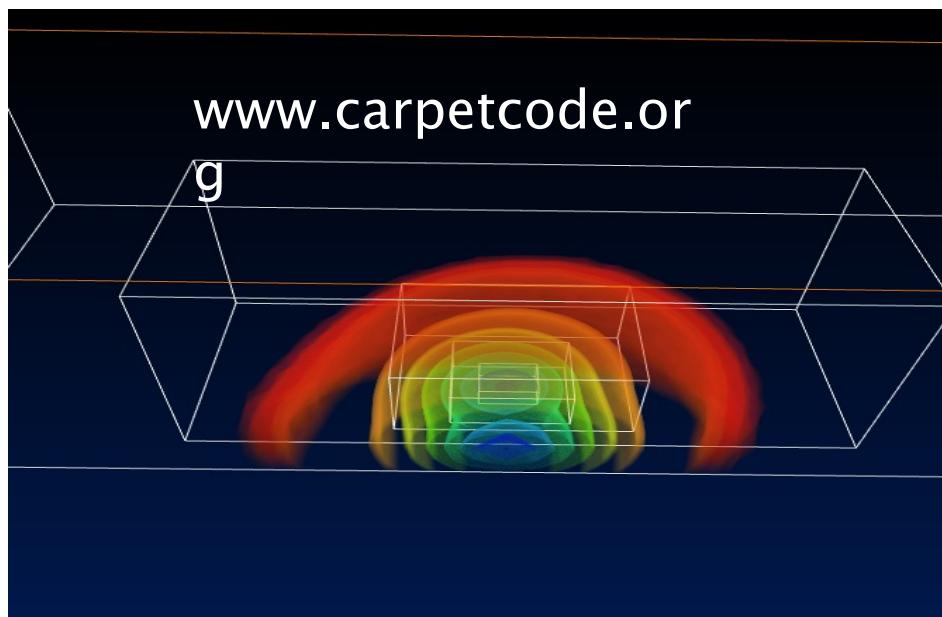
- Geometric terms everywhere!
- + Well adapted for quasi-spherical objects (isolated stars)
- + Spherical and axisymmetry easily imposed
- Severe time-step restrictions for explicit methods
- + Large radial domains and compactification

Adaptive Mesh refinement (AMR)

- Extremely complicated
- + Long expertise / available infrastructure
- + No time-step problem
- Non-adapted to the problem (Cartesian)
- MHD+AMR difficult

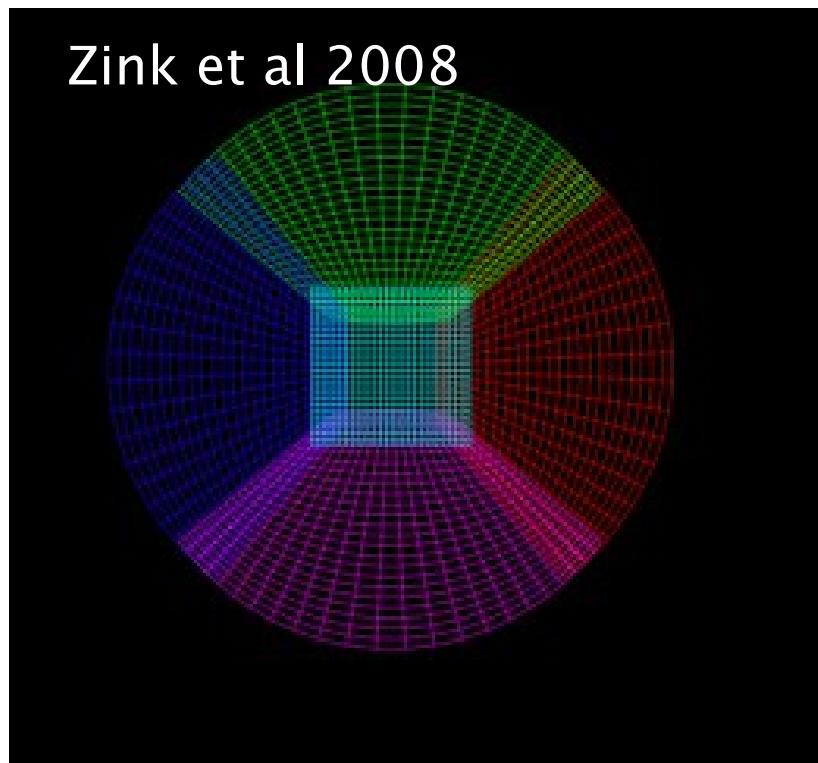
Examples:

- Carpet (Schnetter et al 2006) used in Cactus/Whisky
- HAD used for GRMHD simulations(Anderson et al 2006)



Multi-patched grids

- + Very well adapted to the problem
- + No time-step problem
- Still complicated
- Patch boundaries



Examples:

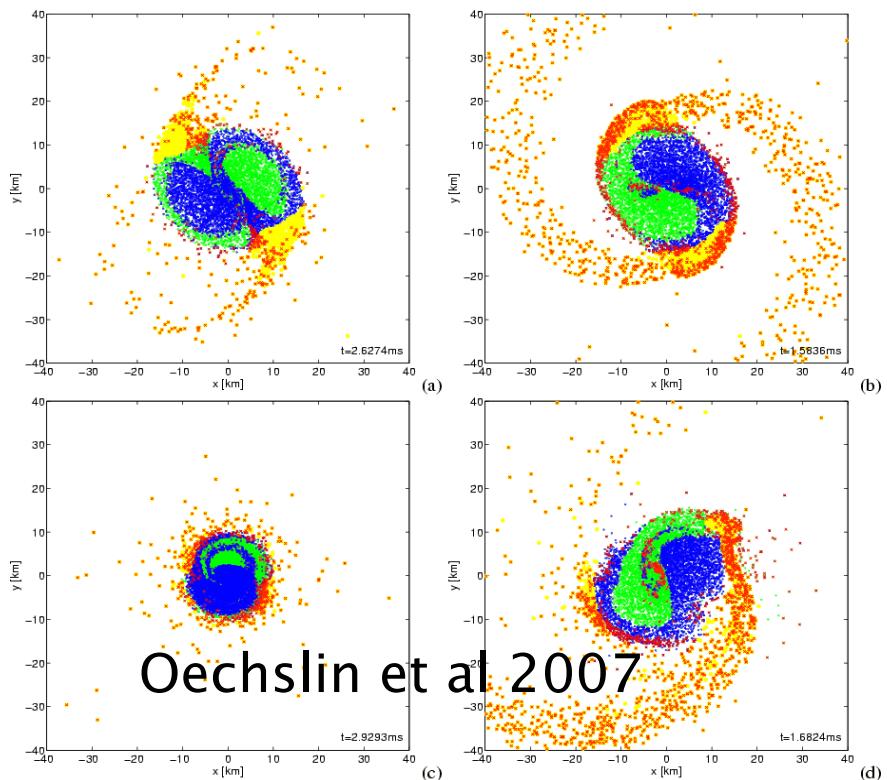
- Zink et al 2008
(cube-sphere)
- Scheidegger et al 2008
(3D Cartesian embeded in 1D spherical)

Smoothed particle hydrodynamics (SPH)

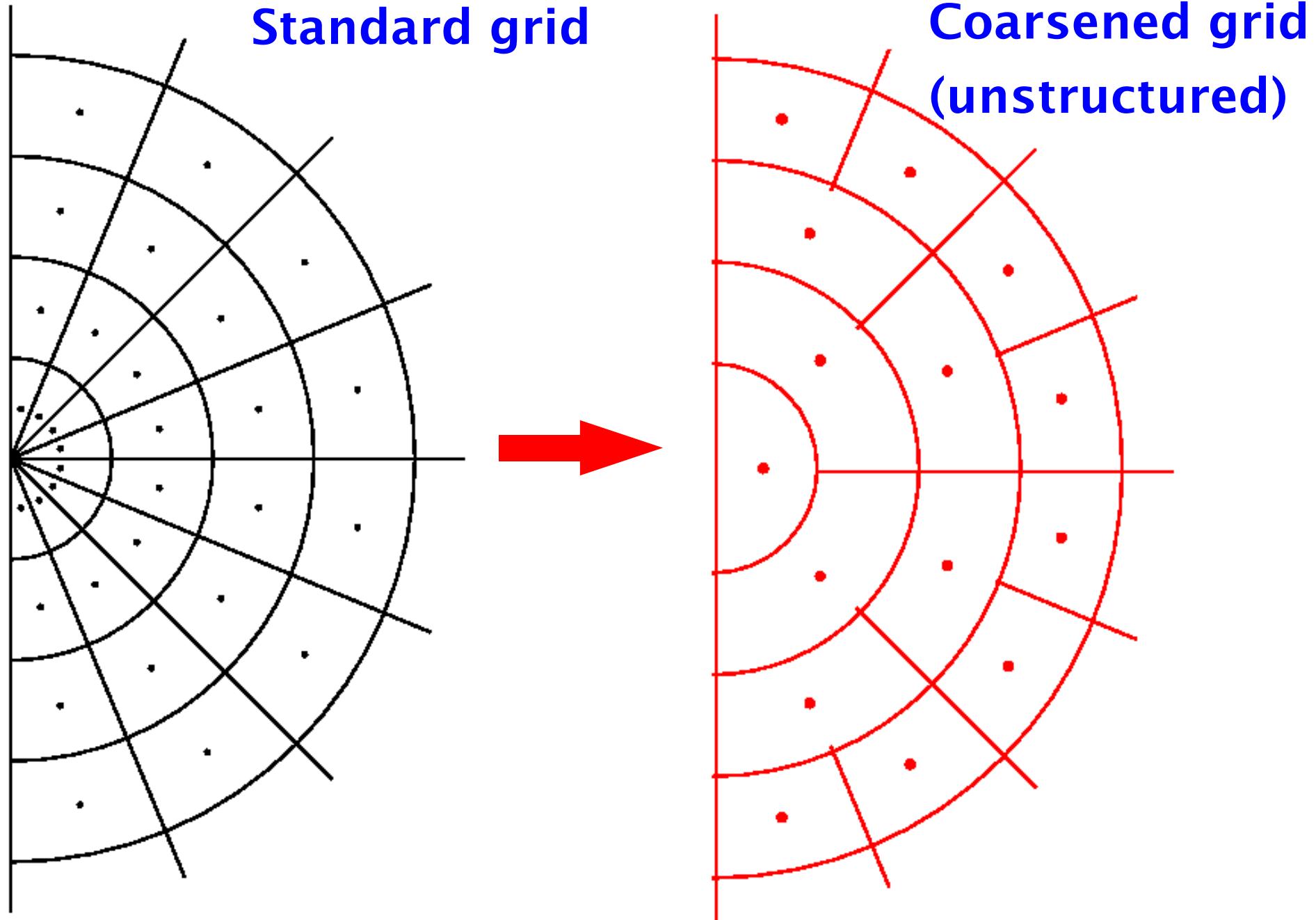
- + Well adapted to the problem
- + No time-step problem
- GR gravity is not a force
(Auxiliary grid for the metric)
- Difficult to handle discontinuities

Examples:

- Oechslin et al 2001



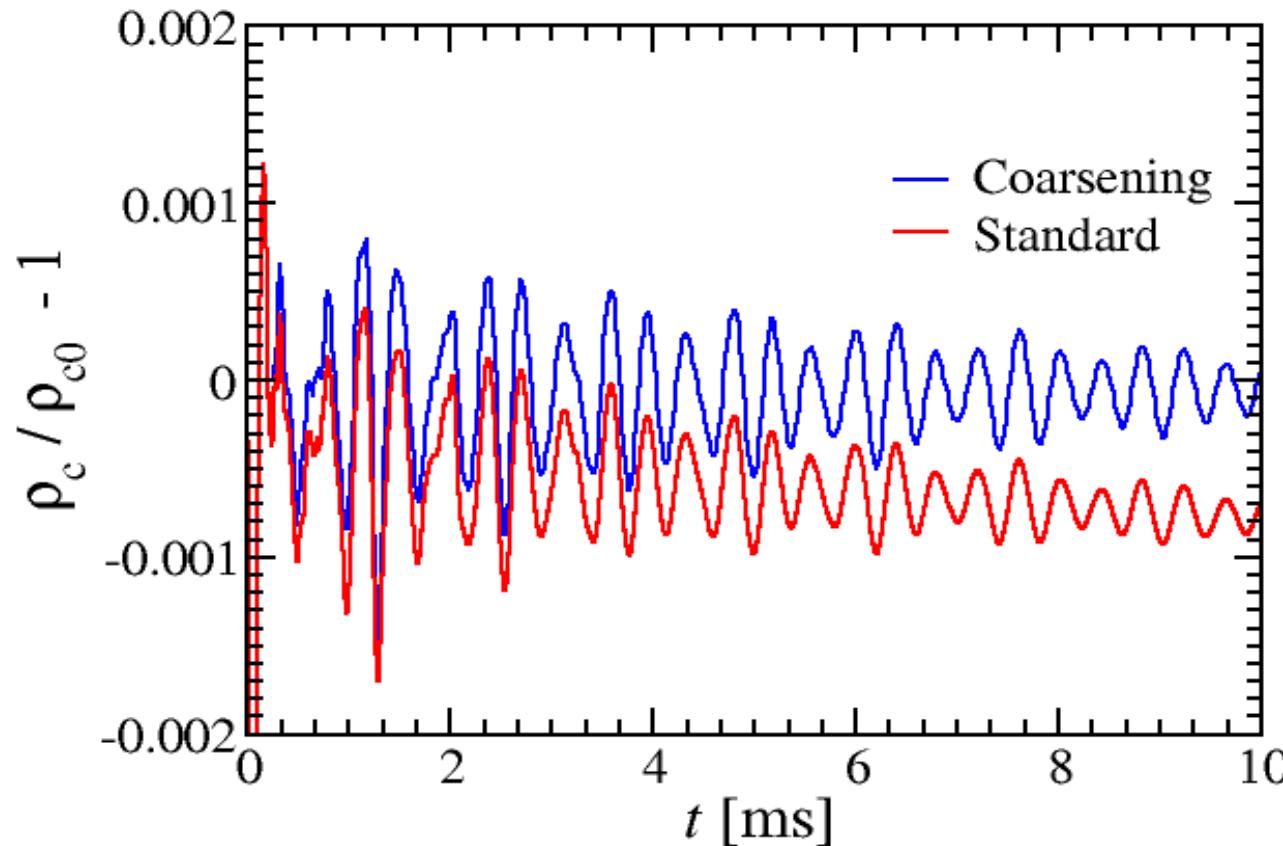
Mesh coarsening (preliminary)



Mesh coarsening (preliminary)

Central density evolution: coarsened vs standard

Rotating neutron star: $M=1.63M_{\odot}$, $R_e=17.3$, $T/|W|=0.074$, rigid rot., 80×32

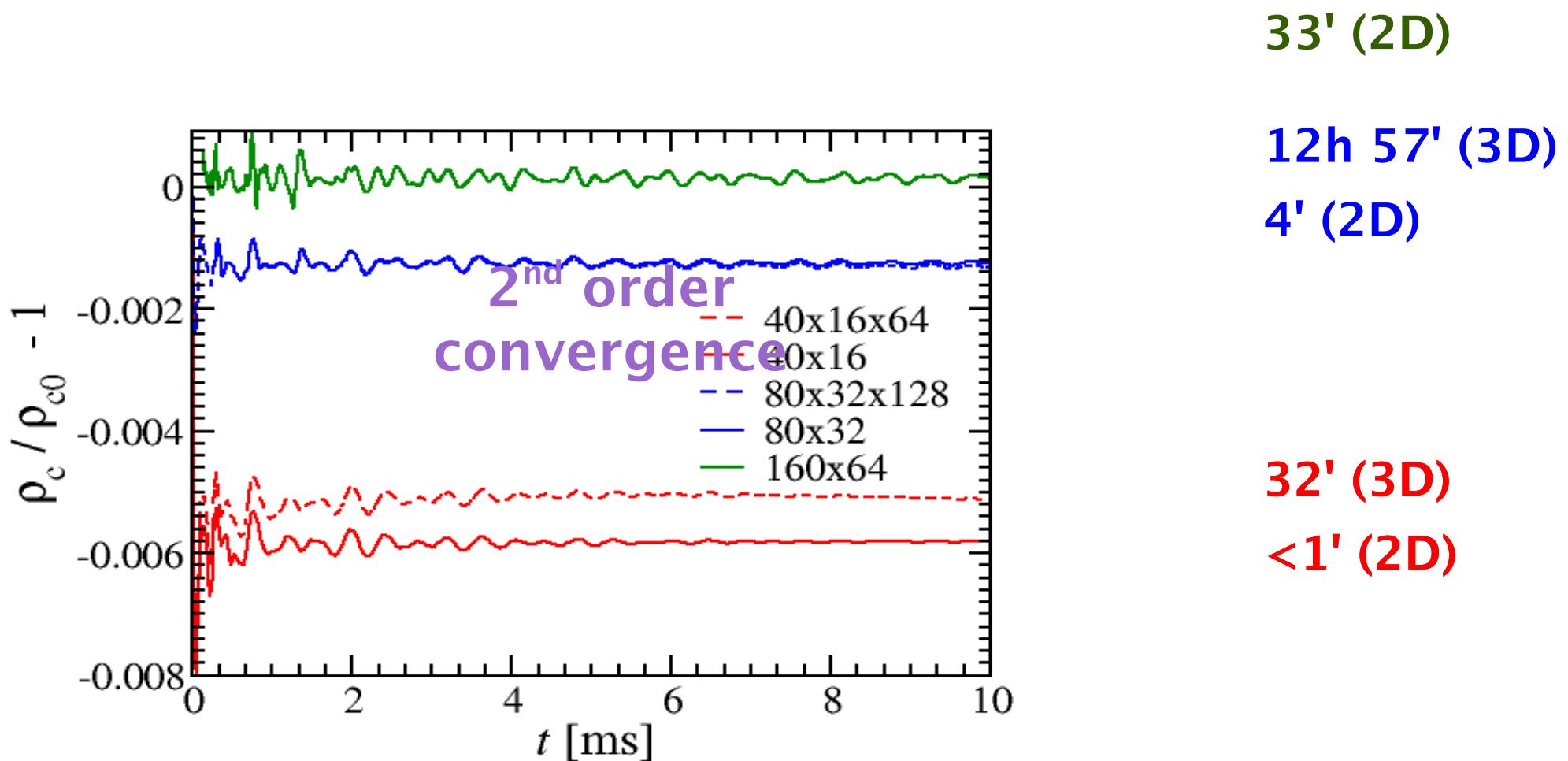


$T_{CPU} < 4'$

$T_{CPU} = 1h\ 44'$

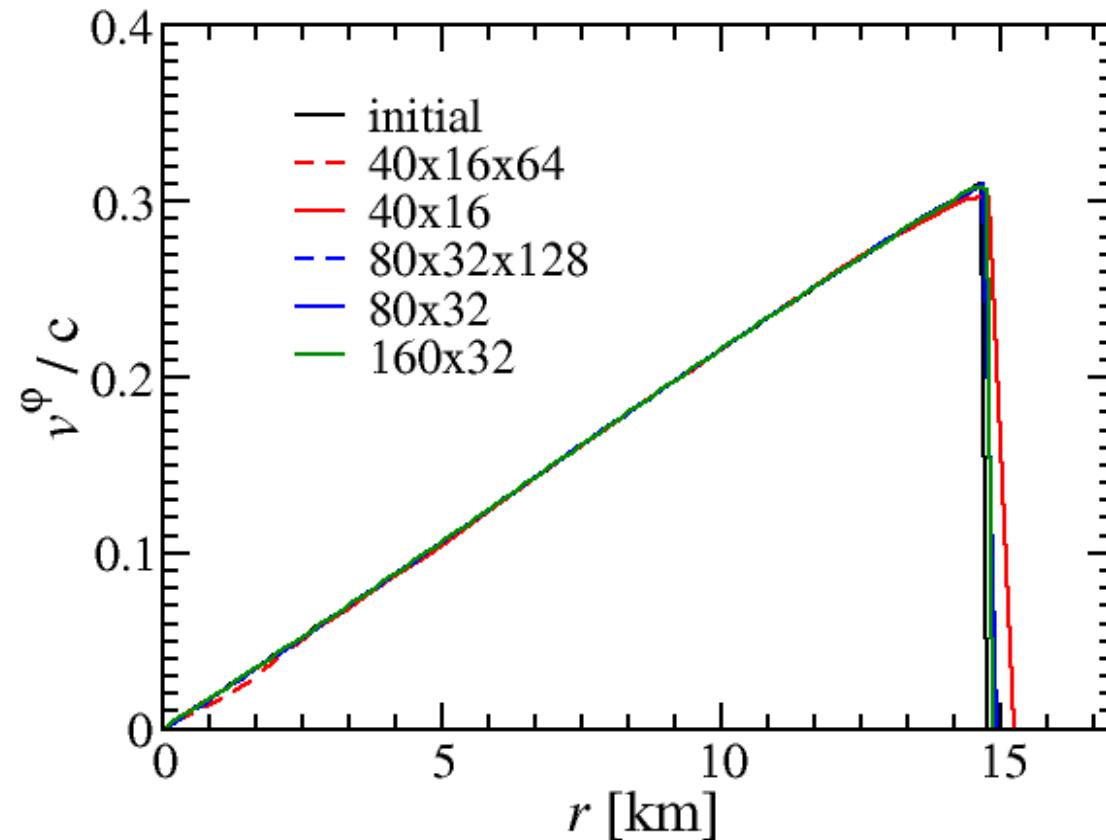
Mesh coarsening (preliminary)

Central density evolution: convergence (PPM+MC)



Mesh coarsening (preliminary)

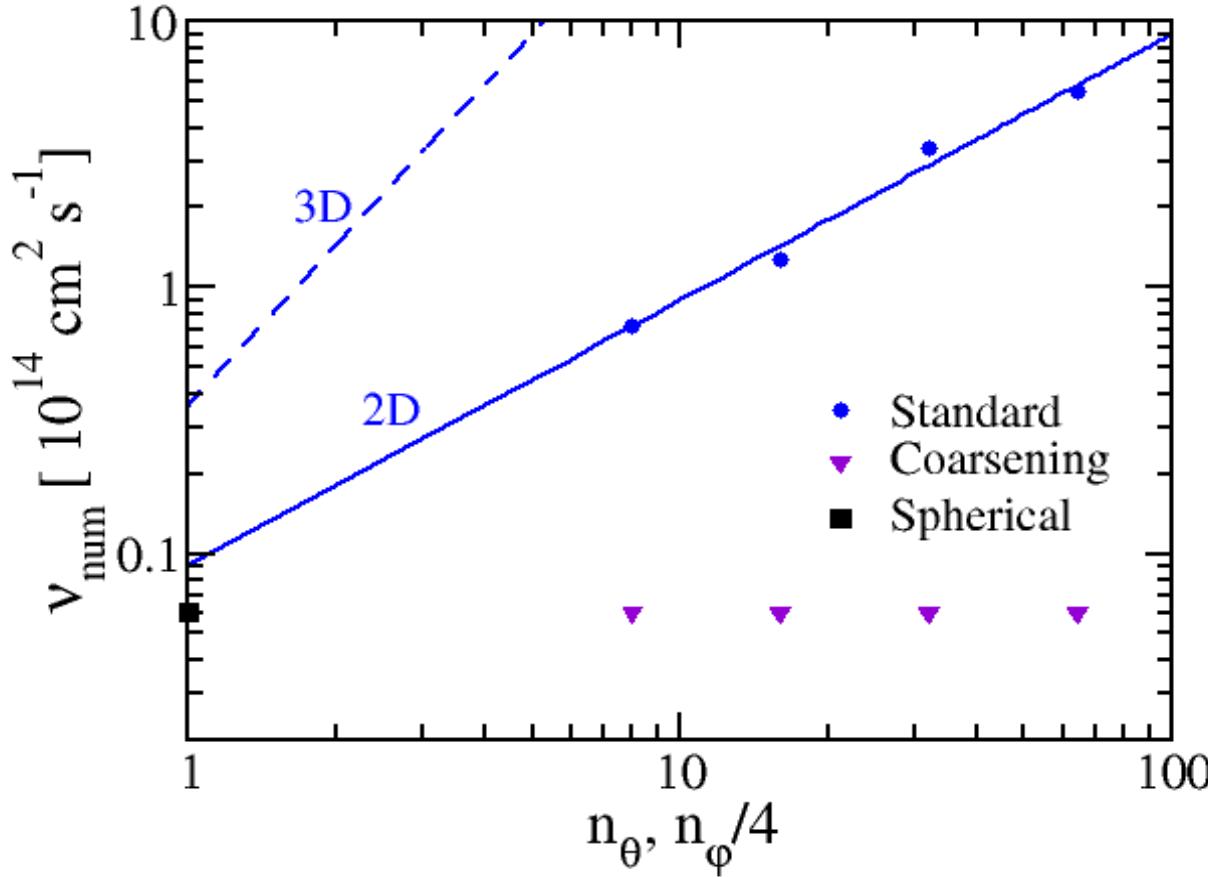
Rotation profile (after 10 ms)



Mesh coarsening (preliminary)

Numerical viscosity (damping of radial modes)

$$\nu \sim \frac{R^2}{\tau_{damping}}$$



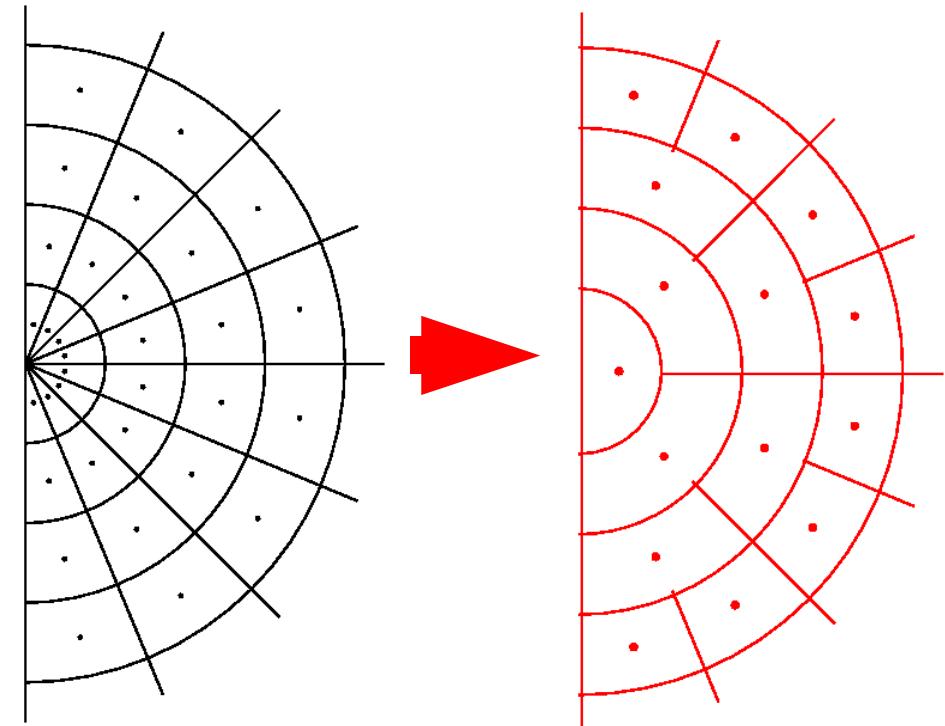
**Non-rotating
NS (2D)**

M=1.4 M

R=14.2 km

Mesh coarsening (preliminary)

- + Well adapted to the problem
(spherical coordinates)
- + No time-step problem
- + Small modification in spherical
coordinates code
- Still experimental



Conclusions

Rotating compact objects (CC, NS) can be much more difficult than standard Supernovae

General relativity

MHD and small scale instabilities

Recovery (maybe no)

Mesh coarsening: poor man's approach to 3D