General relativistic equilibrium configurations with purely toroidal magnetic fields and its stability

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Talk Plan

- Introduction
- (Part 1) Guide line to construct equilibrium configuration of magnetized star
- (Part 2) Stability of magnetic field in neutron star
- Summary

Introduction



Pulsars are believed to be neutron stars which have strong magnetic field ($B \sim 10^{12}$ G)



There are some speical classes called magnetars which have a very strong magnetic field $B \sim 10^{14-15}$ G.

What is an origin of such stong magnetic fields ?

Magnetically driven SN (many references) ? or dynamo process in proto NS ?(Thompson & Duncun 93 etc.).

Introduction

We don't have information of interior magnetic fields because

$$\frac{dE}{dt} = \frac{d}{dt} \left(\frac{1}{2} I \Omega^2 \right) = -4\pi^2 I \frac{\dot{P}}{P^3}$$
assuming dipole radiation
$$\Rightarrow \quad B_p \propto (IP\dot{P})^{1/2} R^{-3}$$

Note that 10^{14-15} G is poloidal component.

Toroidal magnetic field (Takiwaki 04')



Also, toroidal field is easily generated by magnetic winding.

So, there is a possibility that toroidal field dominates the poloidal component.

How to construct magnetized star in equilibrium

Basic assumptions

- 0. Stationary and axisymmetry
- 1. Perfect fluid
- 2. No meridional flow
- 3. Ideal MHD
- 4. Pure toroidal magnetic field
- 5. Barotropic EOS

How to construct magnetized star in equilibrium

Basic equations



Simple analogy for intergrability conditions

Newtonian rotating equilibrium configurations without magnetic field

hydrostatic eq.
$$\frac{\nabla P}{\rho} + \nabla \phi_g + \varpi \Omega^2 \nabla \varpi = 0$$

if $\Omega = \Omega(\varpi)$
Bernoulli eq.
$$\int \frac{dP}{\rho} + \phi_g + \int \Omega^2 \varpi d\varpi = C$$

Purely toroidal magnetic field in general relativistic equilibrium is given by

Faraday tensor
$$F_{r\theta} = (\text{metric}) \times K(u)$$

where $u = \rho_0 h(r \sin \theta)^2 \times (\text{metric})$. We simplify choose K(u) as

 $K(u) = bu^k$ with $k \ge 1$

How to construct magnetized star in equilibrium

Basic equations Einstein equation Maxwell equation Hydrostatic equilibrium (Gravitational field) (Magnetic field) equation (Matter field) For pure toroidal magnetic field case, Integrability condition 4 elliptic equations Bernoulli eq. (1st integral of EOM) for metric potentials

Numerical scheme : KEH (Komatsu et al. 1987)

Magnetized Neutron Stars

Stellar structure on the meridional plane



- 1. Concentration of toroidal field deep inside the star
- 2. Prolate stellar configuration due to the strong magnetic fields
- 3. It is possible to construct hyper strong magnetized star, e.g., H/|W|=O(0.1)

It is known that pure toroidal magnetic fields are unstable (Tayler instability).

Tayler instability = interchange type instability



Set up

- 1. Prepare the equilibrium configurations as initial conditions (with varying the value of k)
- 2. Start the axisymmetric GRMHD simulation

k=1 configuration (Stable)



k=2 configuration (Unstable)



Unstable configuration (k=2) comes to the stable one (k=1)

Magnetic energy -> Kinetic energy



Tayler instability induces the circular motion on the meridional plane and settles down to a quasi equilibrium state.

Energy / M.

<u>Tayler instability in 3D GRMHD simulation</u> (preliminary)

Recall that m = 1 perturbation inevitabley induces the instability because the instability condition is $k > m^2/4$.



<u>Summary</u>

- 1. Part 1: We formulate GR equilibrium configuration with purely toroidal magnetic field.
- 2. Part 2: With 2D GRMHD simulation, we explore the Tayler instability and confirm the linear analysis prediction.
- 3. Part 2: Toroidal magnetic fields settle down to "new" quasi equilibrium state with the circular motion.

Future issues

- 1. Effect of magnetic field on EOS for hyper strong magnetic field case
- 2. Tayler instability for non-axisymmetric perturbation (on going)
- 3. GR equilibrium configuration with poloidal-toroidal magnetic field.