Holographic Conductivity

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Based on work with Gary Horowitz, Jorge Santos and Mike Blake







Basics of Conductivity

$$\vec{j} = \vec{j}(\omega)e^{-i\omega t}$$

Obm'de a Wodel:

Drude Model: AC Conductivity $= \frac{1}{\rho} \frac{1}{1-i\omega\tau}$

Fourier Transform:
$$\vec{j}(\omega) = \sigma(\omega) \vec{E}(\omega) \quad \vec{j}(\omega) = \vec{j}(\omega) \vec{e}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{e}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) = \vec{j}($$

$$\sigma(\omega) = \frac{1}{\rho} \frac{1}{1 - i\omega\tau} \qquad \sigma(\omega) = \frac{1}{\rho} \frac{1}{1 - i\omega\tau}$$

$$ec{E} = N ec{p}(\omega) e^{-i\omega t} + i/\omega \quad \text{as} \quad \omega o \infty$$

$$ec{j} = ec{j}(\omega) e^{-i\omega t} \qquad \qquad \sigma(\omega) o i/\omega \quad \text{as} \quad \omega o \infty$$

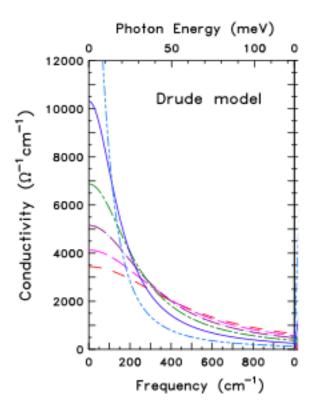
Drude Model

$$m\frac{d\vec{v}}{dt} + \frac{m}{\tau}\vec{v} = q\vec{E}$$

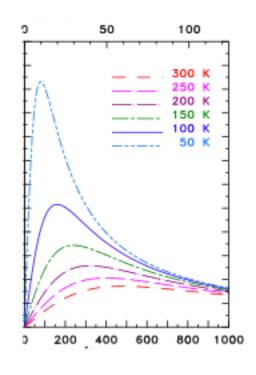
$$\vec{j} = nq\vec{v}$$

$$\vec{j} = nq\vec{v}$$
 \Longrightarrow $\sigma(\omega) = \left(\frac{nq^2\tau}{m}\right)\frac{1}{1 - i\omega\tau}$

Drude Model

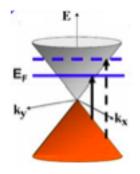


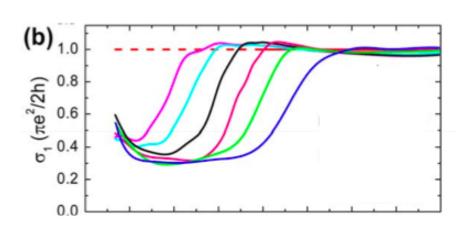
$$Re(\sigma) = \frac{\sigma_0}{1 + \omega^2 \tau^2}$$

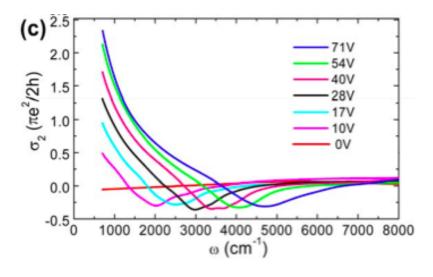


$$Im(\sigma) = \frac{\sigma_0 \omega \tau}{1 + \omega^2 \tau^2}$$

Graphene

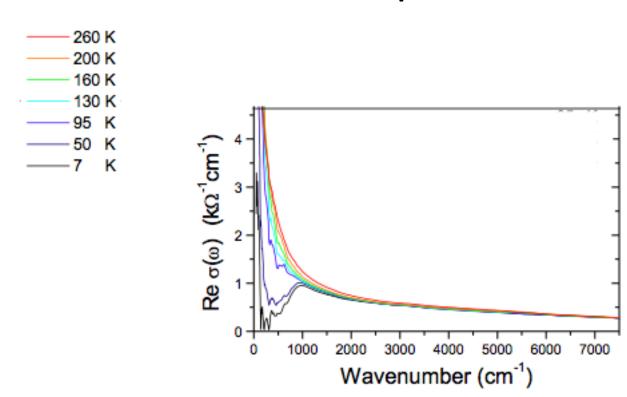






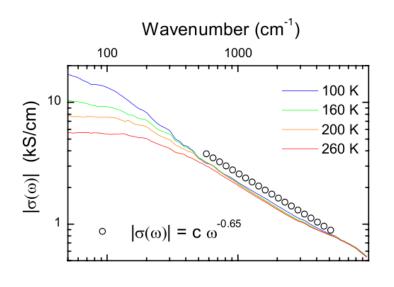
Cuprates

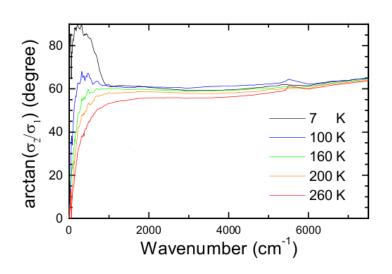
van der Marel et al. (2004)



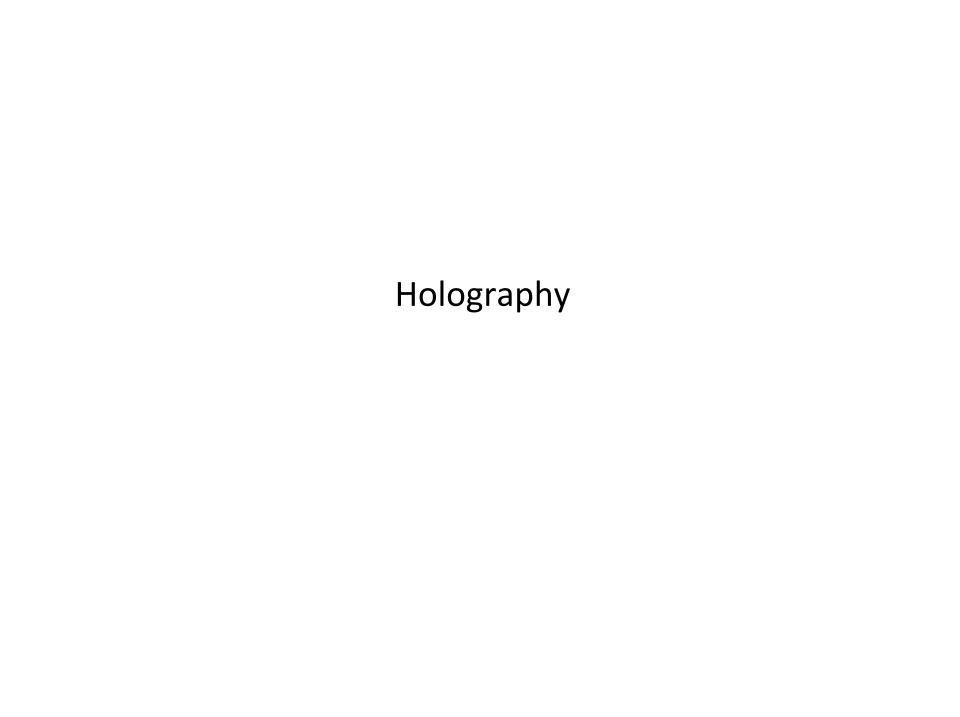
 $Bi_2Sr_2Ca_{0.92}Y_{0.08}Cu_2O_{8+\delta}$

Cuprates

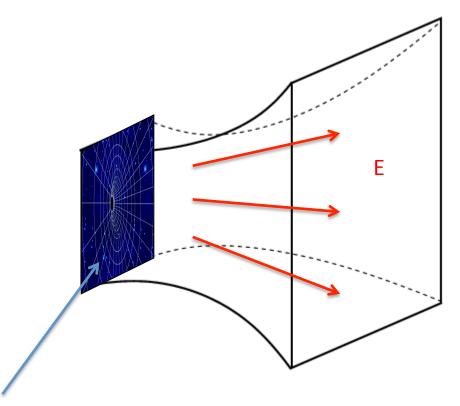




$$|\sigma(\omega)| \sim \frac{1}{\omega^{2/3}}$$

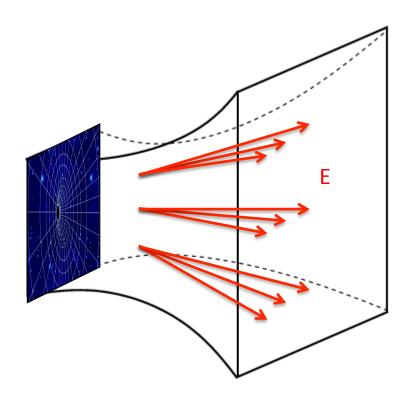


Finite Density Matter



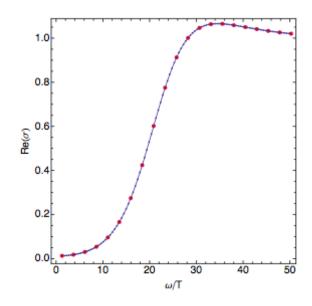
Boundary field theory d=2+1

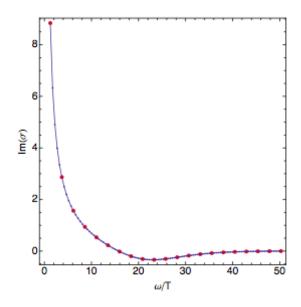
- d=3+1 bulk
- Hawking radiation = finite temperature
- Electric field = chemical potential



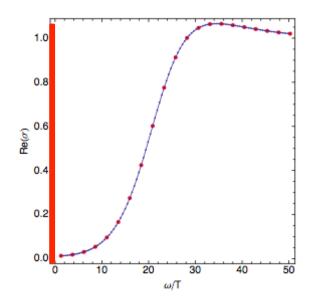
$$A_x = \frac{E}{i\omega} + \langle J_x \rangle r + \dots$$

$$j(\omega) = \sigma(\omega)E(\omega)$$

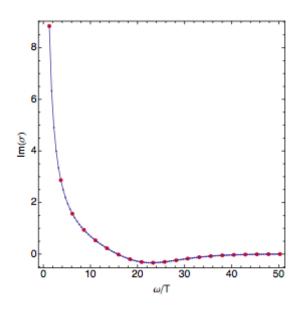




$$j(\omega) = \sigma(\omega)E(\omega)$$



$$\operatorname{Re} \sigma(\omega) \sim K \delta(\omega)$$



$$\operatorname{Im}(\sigma) \to \frac{K}{\pi\omega}$$

The Delta-Function

$$\operatorname{Re} \sigma(\omega) \sim K \delta(\omega)$$

- Due to: Finite density of charge carriers
 - Translational invariance

Story 1:

Breaking Translational Invariance

How to Build a Lattice

Horowitz, Santos and Tong (2012)

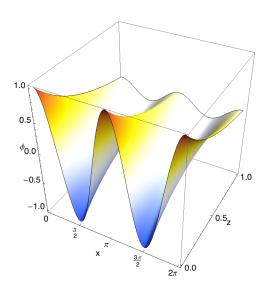
$$\mathcal{L} = \mathcal{L}_{QFT} + \mu(x, y)\mathcal{Q}$$

Spatially dependent chemical potential

Aside: We have also built other lattices which have smooth charge density. The following results remain unchanged

The lattice black hole

Stripes:
$$\mu = \mu[1 + Acos(k_L x)]$$



Parameters

 T, μ, k_L, A

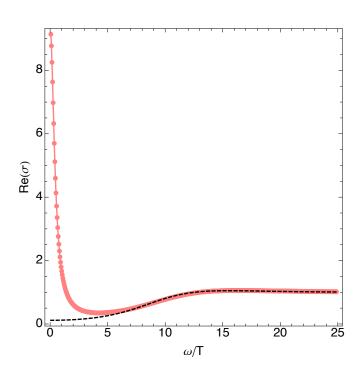
$$ds^{2} = \frac{L^{2}}{z^{2}} \left[-g_{tt}(z,x)dt^{2} + g_{zz}(z,x)dz^{2} + g_{xx}(z,x)(dx + a(z,x)dz)^{2} + g_{yy}(z,x)dy^{2} \right]$$

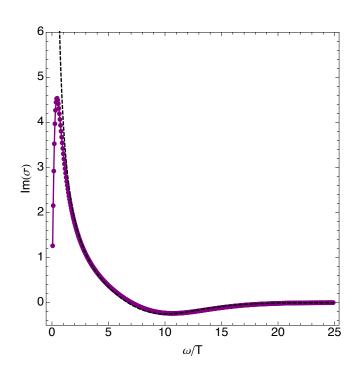
Perturbing the Lattice

$$\delta g_{tt}, \ \delta g_{tz}, \ \delta g_{tx}, \ \delta g_{zz}, \ \delta g_{zx}, \ \delta g_{xx}, \ \delta g_{yy}$$

$$\delta A_t, \ \delta A_z, \ \delta A_x, \ \delta \Phi$$

$$j(\omega) = \sigma(\omega)E(\omega)$$

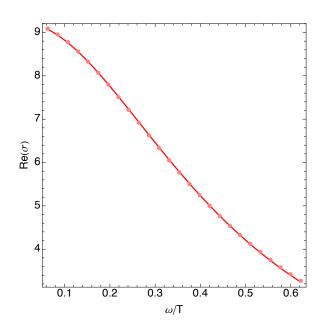


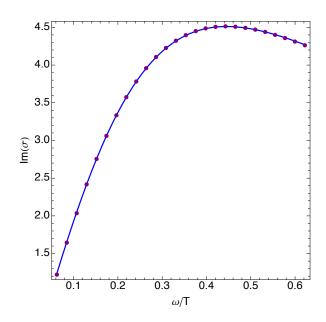


$$\mu = 1.4$$
 $T = 0.115\mu$ $k_L = 2$ $A = 1.5$

Low Frequency Drude Behaviour $\omega \lesssim T$







$$\sigma(\omega) = \frac{K\tau}{1 - i\omega\tau}$$

Note: No quasiparticles

DC Resistivity

$$\rho \sim T^{2\nu - 1}$$

• Exponent depends on lattice spacing:
$$\nu=\frac{1}{2}\sqrt{5+2(k/\mu)^2-4\sqrt{1+(k/\mu)^2}}$$

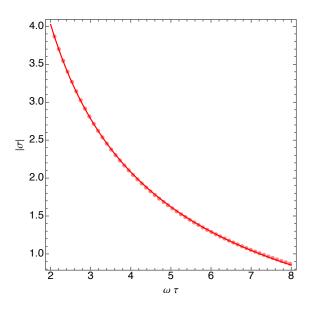
• This is characteristic of a *locally critical* theory i.e. geometry: $AdS_2 \times {f R}^2$ field theory $z \to \infty$

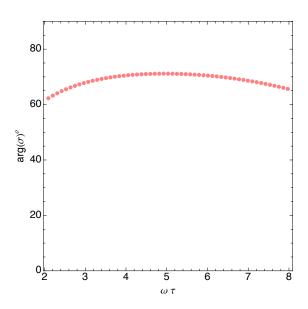
Hartnoll, Hofman (2012)

Model building: a mechanism suggested to drive this to linear resistivity.

Mid-Frequency Behaviour





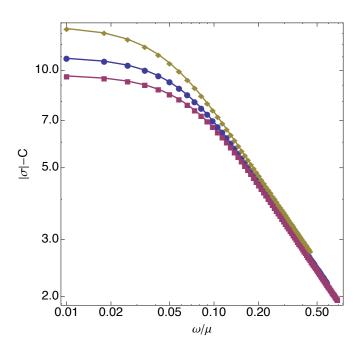


$$|\sigma(\omega)| = \frac{B}{\omega^{2/3}} + C$$

Robust Power-Law

$$T = 0.98\mu, T = 0.115\mu, T = 0.13\mu$$

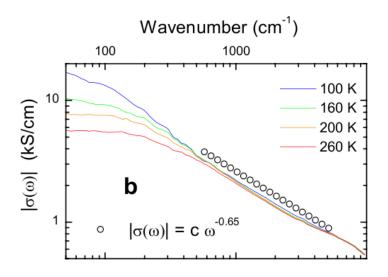
Log-log plots



$$|\sigma(\omega)| = \frac{B}{\omega^{2/3}} + C$$

B temperature independent

Comparison to Cuprates

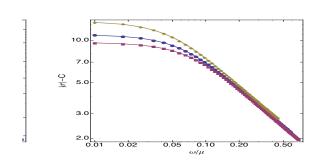


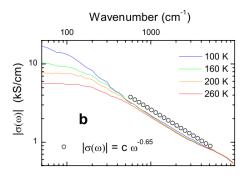
$$|\sigma(\omega)| \sim \frac{1}{\omega^{2/3}}$$

More Scaling

- Also seen in thermoelectric conductivity
- Also seen in AdS₅

Why?





Story 2:

A Holographic Model for Momentum Dissipation

Momentum Dissipation in Holography

Vegh (2013)

Bulk diffeomorphisms
$$\Longrightarrow$$
 $\partial_{\mu}T^{\mu\nu}=0$

No diffeomorphisms
$$\implies \partial_{\mu}T^{\mu\nu} \neq 0$$

This means massive gravity. (Gulp!)

Is massive gravity an effective holographic theory for models with disorder or a lattice?

Massive Gravity

Expanding around flat space: $g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}$

$$S_{FP} = \int d^4x \sqrt{-g} \, \frac{m^2}{4} \left(h^{\mu\nu} h_{\mu\nu} - h^{\mu}_{\ \mu} h^{\nu}_{\ \nu} \right)$$

But expanding around a general background



Add further non-linear terms to try to remedy this...

Massive Gravity...ugly, but it works

de Rahm, Gabadadze, Tolley (2010) Hassan, Rosen (2011)

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left(R + \Lambda + \sum_{i=1}^4 \alpha_i \, e_i[\mathcal{K}] \right)$$

$$\mathcal{K}^2 = g^{-1}f$$

$$e_1 = \operatorname{Tr} \mathcal{K}$$

$$e_2 = (\operatorname{Tr} \mathcal{K})^2 - \operatorname{Tr} (\mathcal{K}^2)$$

Fixed background metric

$$e_3 = (\operatorname{Tr} \mathcal{K})^3 - 3(\operatorname{Tr} \mathcal{K})\operatorname{Tr} (\mathcal{K}^2) + 2\operatorname{Tr} (\mathcal{K}^3)$$

$$e_4 = (\operatorname{Tr} \mathcal{K})^4 - 6(\operatorname{Tr} \mathcal{K})^2 \operatorname{Tr} (\mathcal{K}^2) + 3\operatorname{Tr} (\mathcal{K}^2)^2 + 8(\operatorname{Tr} \mathcal{K}) \operatorname{Tr} (\mathcal{K}^3) - 6\operatorname{Tr} (\mathcal{K}^4)$$

Holographic Massive Gravity

Does it give sensible answers?

Black Hole Thermodynamics

Blake and Tong (to appear)

Yes

$$S = S_{BH} \sim r_h^2$$

But horizon radius is $r_h = r_h(T, \mu, m^2)$

graviton mass parameters (typically depend on position in bulk)

Hydrodynamics

Davison (2013)

$$\partial_{\mu} T^{\mu i} = -\tau^{-1} T^{0i}$$

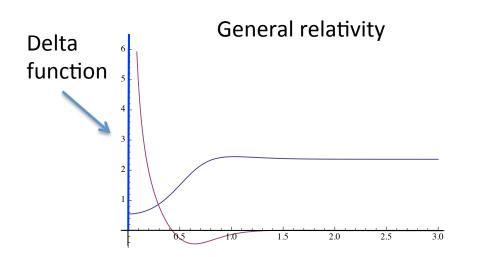
with relaxation time $au^{-1} \sim r_h m^2(r_h)$

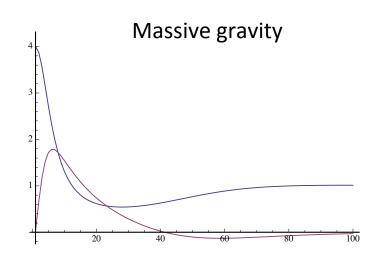
DC Conductivity

Blake and Tong (to appear)

$$\sigma_{DC} = \frac{\mu^2}{m^2(r_h)} + 1$$

Universal behaviour, governed only by horizon

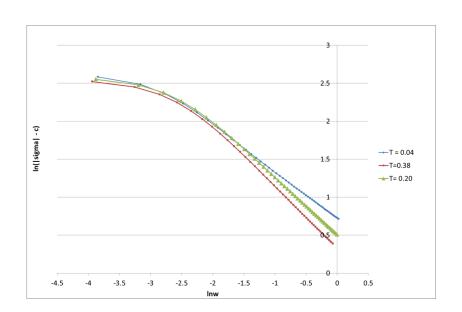




For small frequencies:
$$\sigma(\omega) \sim \frac{\sigma_{
m DC}}{1-i\omega au}$$

Vegh (2013) Davison (2013)

Mid-Frequency Behaviour



Vegh (2013) Blake and Tong

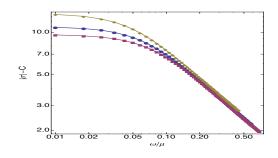
$$|\sigma(\omega)| = \frac{B}{\omega^{\gamma}} + C$$

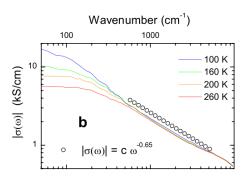
Some evidence for scaling, but all parameters now temperature dependent.

Summary

Story 1: General Relativity with a Lattice

Why Scaling?





Story 2: Massive Gravity

- Limited scaling, but sensible results.
- What is this a theory of? Disorder?