The Last Word on Large Scale CMB Anomalies

Assaf Ben-David

Niels Bohr International Academy and Discovery Center, The Niels Bohr Institute, Copenhagen

NBIA–APCTP Workshop August, 2014





Outline

Introduction: Large scales – potential and problems

- Review of some large scale CMB anomalies:
 - The "Axis of Evil"
 - Low power on large scales
 - Lack of large-angular correlations
 - Hemispherical power asymmetry
 - Point parity
 - Mirror parity

[ABD & Kovetz, 2014, arXiv:1403.2104]

- The Integrated Sachs–Wolfe Effect
 - Estimation
 - New insight on large scale anomalies?
- Future prospects

The Standard Model

Assumption:

On large enough scales the universe is homogeneous and isotropic

- Quantum fluctuations of the inflaton field seed Gaussian fluctuations in the metric perturbation field Φ

$$\langle \Phi_{\mathbf{k}} \rangle = 0$$
 $\langle \Phi_{\mathbf{k}} \Phi_{\mathbf{k}'}^* \rangle = (2\pi)^3 \delta(\mathbf{k} - \mathbf{k}') P_{\Phi}(k)$

 Projected onto the sphere, we get the CMB temperature anisotropy

$$T(\hat{\mathbf{n}}) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\hat{\mathbf{n}})$$

 The coefficients are uncorrelated Gaussian random variables

$$\langle a_{\ell m} \rangle = 0 \qquad \langle a_{\ell m} a_{\ell' m'}^* \rangle = \delta_{\ell \ell'} \delta_{m m'} C_{\ell}$$



[Planck Collaboration XV, 2013, arXiv:1303.5075]

Why are Large Scales Interesting?



- Short inflation is theoretically preferred.
- If some pre-inflationary signal exists, it will affect the largest scales.

- Largest cosmic variance Anomalies are limited to $\sim 3\sigma$.
- Use large scale anomalies as clues.
- A-posteriori choices affect the statistical interpretation of the results!

Deformations on Large Scales



- Pre-inflationary relics
- Example the PIP model: [Fialkov et al., JCAP 2010]
 - A relic pre-inflationary particle (PIP) creates a one-point contribution on large scales, $\Phi \propto \log r$.
 - Does not affect the two-point function.
 - Search the CMB for a set of giant concentric rings. [Kovetz, ABD & Itzhaki, ApJ 2010]



- A non-trivial topology:
 - Change the global structure of the Universe by identifying points in space.
 - Compact dimensions limit the power on large scales.
 - Example A compact 3-torus with dimensions $L_x \times L_y \times L_z$:

$$\mathbf{k} = 2\pi \left(\frac{n_x}{L_x}, \frac{n_y}{L_y}, \frac{n_z}{L_z}\right) \qquad n_i \in \mathbb{Z}$$

 Example – "Stringy" topologies: orbifold point, orbifold line. [ABD, Rathaus & Itzhaki, JCAP 2012; Rathaus, ABD & Itzhaki, JCAP 2013]

Large Scale CMB Data



Contributions to Large Scale Signal



Galactic Foregrounds

- Component separation methods provide high-quality CMB maps.
- Maps are still contaminated in the Galactic plane area.
- In pixel-space, masking is straightforward. However, masking breaks isotropy!
- In harmonic-space, masking is problematic, since the spherical harmonics are not orthogonal on the masked sky.



[Bennett et al., ApJS 2013]



Large Scale CMB Anomalies

The "Axis of Evil"

Low Power on Large Scales

Lack of Large-Angular Correlations

Hemispherical Power Asymmetry

Point Parity

Mirror Parity

The "Axis of Evil"





The "Axis of Evil"

- Quadrupole and octupole are both "planar" – dominated by $|m| = \ell$ modes.
- [de Oliveira-Costa et al., PRD 2004 Copi et al., PRD 2004 Copi et al., MNRAS 2006 Abramo et al., PRD 2006 Copi et al., Advances in Astronomy 2010 Planck Collaboration XXIII, 2013, arXiv:1303.5083 Rassat and Starck, A&A, 2013 Rassat et al., A&A, 2013 Copi et al., 2013, arXiv:1311.4562 Rassat et al., 2014, arXiv:1405.1844

- Their two planes are closely aligned.
- Statistic "Maximum Angular Momentum Dispersion":

Find
$$\hat{\mathbf{n}}_{\ell} = \operatorname*{arg\,max}_{\hat{\mathbf{n}}} \sum_{m=-\ell}^{\ell} m^2 \left|a_{\ell m}(\hat{\mathbf{n}})\right|^2$$
 and test $\hat{\mathbf{n}}_2 \cdot \hat{\mathbf{n}}_3$.

- Significance: As high as 3.8σ on WMAP, 2.7–3.2σ on Planck data.
 [Copi et al., 2013, arXiv:1311.4562]
- If allowing dominance by <u>any</u> m, all *e* = 2–5 are found to be correlated. [Land & Magueijo, PRL 2005; PRD 2005; MNRAS 2007]

The "Axis of Evil" – Multipole Vectors

- Irreducible representation of the rotation group independent of coordinate system. [Maxwell, 1873]
- Instead of 2ℓ +1 coefficients, each multipole ℓ is decomposed to ℓ axes and an amplitude:

$$T_{\ell}(\hat{\mathbf{n}}) = \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\hat{\mathbf{n}}) = A_{\ell} \prod_{i=1}^{\ell} \left(\hat{\mathbf{v}}_{i}^{(\ell)} \cdot \nabla \right) \left. \frac{1}{r} \right|_{r=1}$$

• Each pair defines a plane $\mathbf{w}_{i,j}^{(\ell)} = \hat{\mathbf{v}}_i^{(\ell)} \times \hat{\mathbf{v}}_j^{(\ell)}$.

- In this way each multipole defines $\ell(\ell-1)/2$ planes.
 - The normal to the quadrupole plane is the MAMD axis.
 - Octupole planarity \rightarrow 3 octupole planes are aligned.
 - Quadrupole plane is aligned with the 3 octupole planes.
- Statistic:

$$D_i = \left| \mathbf{w}_{1,2}^{(2)} \cdot \mathbf{w}_{j,k}^{(3)} \right|$$

 $S = \frac{1}{3} \sum_{i=1}^{3} D_i$

Significance of 2.9σ on WMAP and 2.1–2.6σ on Planck data. [Copi et al., 2013, arXiv:1311.4562]

The "Axis of Evil"



[Copi et al., 2013, arXiv:1311.4562]

Low Power on Large Scales

- Cosmological model is determined almost exclusively by small scales.
- It appears that the quadrupole amplitude is too low.
- WMAP: Quadrupole still within cosmic variance (<2σ). [Bennett et al., ApJS 2013]
- Planck: Using the statistic $s_1 = \max_{r \in [0,1]} \frac{1}{\sqrt{\ell_{\max}}} \sum_{\ell=2}^{\lfloor \ell_{\max}r \rfloor} \frac{\hat{C}_{\ell} C_{\ell}}{\Delta \hat{C}_{\ell}}$ report 2.5 σ .

[Planck Collaboration XV, 2013, arXiv:1303.5075]

- Only $\ell = 2: C_2$ is χ^2 -distributed in ΛCDM .
 - On WMAP and Planck data:
 1.7–2.3σ, depending on masking.
 [Rassat et al., A&A 2013; Rassat et al., 2014, arXiv:1405.1844]



[Planck Collaboration XV, 2013, arXiv:1303.5075]

Lack of Large-Angular Correlations

- Angular correlation function $C(\theta) = \langle T(\hat{\mathbf{n}}_1)T(\hat{\mathbf{n}}_2) \rangle_{\hat{\mathbf{n}}_1 \cdot \hat{\mathbf{n}}_2 = \cos \theta}$ is near zero for large scales.
- Related to low-power anomaly: $C(\theta) = \sum_{\ell} \frac{2\ell+1}{4\pi} C_{\ell} P_{\ell}(\cos \theta)$

• Statistic: [Spergel et al., ApJS 2003] $S_{1/2} = \int_{-1}^{1/2} C(\theta)^2 d\cos\theta$

- Significance of up to 3.3σ on masked
 WMAP and Planck maps.
 [Copi et al., 2013, arXiv:1310.3831]
- Limit of 1/2 chosen a-posteriori.
- Compares C(θ) to zero, not to ΛCDM expectation.



[Copi et al., 2013, arXiv:1310.3831]

• A more robust statistic: [Planck Collaboration XXIII, 2013, arXiv:1303.5083]

$$\chi^2 = \sum_{i,j=1}^{N_{\text{bin}}} (C(\theta_i) - \langle C(\theta_i) \rangle) M_{ij}^{-1} (C(\theta_j) - \langle C(\theta_j) \rangle)$$

No significant deviation from ΛCDM expectation.

Hemispherical Power Asymmetry

- When the large-scale power spectrum is estimated locally, it is not isotropic.
- Maximal isotropy at (l, b) = (57°, 10°), near the ecliptic pole.



About 2.7σ on WMAP data. [Eriksen et al., ApJ 2004; Hansen et al., MNRAS 2004]

- Can also be seen in two-, three- and four-point angular correlation functions. [Planck Collaboration XXIII, 2013, arXiv:1303.5083]
 - Using a χ^2 statistic, significance > 3.1 σ .
- A phenomenological model dipole modulation: [Gordon, ApJ 2007; Eriksen et al., ApJL 2007; Akrami et al., ApJL 2014]

$$T(\hat{\mathbf{n}}) = T^{\text{iso}}(\hat{\mathbf{n}})(1 + \mathbf{d} \cdot \hat{\mathbf{n}})$$

Provides a better fit on large scales than the isotropic model.

Hemispherical Power Asymmetry



[Planck Collaboration XXIII, 2013, arXiv:1303.5083]

Point Parity

- Power spectrum shows preference for odd multipoles on large scales. [Land & Magueijo, PRD 2005; Kim & Naselsky, ApJL 2010; Kim & Naselsky, PRD 2010]
- Equivalent to tendency for odd point-parity: under $\hat{\mathbf{n}} \to -\hat{\mathbf{n}}$, the harmonic coefficients transform as $a_{\ell m} \to (-1)^{\ell} a_{\ell m}$.



[Kim & Naselsky, PRD 2010]

• Statistic:

$$P^{\pm}(\ell_{\max}) = \sum_{\ell=2}^{\ell_{\max}} \frac{1}{2} \left[1 \pm (-1)^{\ell} \right] \frac{\ell(\ell+1)}{2\pi} C_{\ell} \qquad g(\ell_{\max}) = \frac{P^{+}(\ell_{\max})}{P^{-}(\ell_{\max})}$$

- Maximal significance is for ℓ_{max} = 22. Both WMAP and Planck data show a ~ 2.8σ deviation. [Kim & Naselsky, PRD 2010; Planck Collaboration XXIII, 2013, arXiv:1303.5083]
- Is the low quadrupole part of a broader odd-parity preference?

Mirror Parity

[ABD & Kovetz, 2014, arXiv:1403.2104]

- Parity with respect to reflections through a plane, $\hat{\mathbf{r}} \rightarrow \hat{\mathbf{r}}_{\hat{\mathbf{n}}} = \hat{\mathbf{r}} 2 \left(\hat{\mathbf{r}} \cdot \hat{\mathbf{n}} \right) \hat{\mathbf{n}}$.
- Pixel-space estimator: [de Oliveira-Costa et al., ApJ 1996; PRD 2004; Finelli et al., JCAP 2012; Planck Collaboration XXIII, 2013, arXiv:1303.5083]
 - Average difference between hemispheres:
 - Seems easy to apply a Galactic mask.
- Harmonic-space estimator: [ABD, Kovetz & Itzhaki, ApJ 2012; Rassat & Starck, A&A, 2013]
 - Under reflection through *z*-axis, $Y_{\ell m}(\hat{\mathbf{r}}) \rightarrow (-1)^{\ell+m} Y_{\ell m}(\hat{\mathbf{r}})$.
 - For each direction, compare for each ℓ the distribution of power between even and odd $\ell+m$ multipoles.
 - Normalized

$$S_{\rm h}(\hat{\mathbf{n}}) = \sum_{\ell=2}^{\ell_{\rm max}} \sum_{m=-\ell}^{\ell} (-1)^{\ell+m} \frac{|a_{\ell m}(\hat{\mathbf{n}})|^2}{\hat{C}_{\ell}} - (\ell_{\rm max} - 1)$$

 $S_{\rm p}^{\pm}(\hat{\mathbf{n}}) = \left[\frac{T(\hat{\mathbf{r}}) \pm T(\hat{\mathbf{r}}_{\hat{\mathbf{n}}})}{2}\right]^2$

- Easy to test scale dependence.
- Masking the Galactic plane is not straightforward.

$$\hat{C}_{\ell} = \frac{1}{2\ell+1} \sum_{m} |a_{\ell m}|^2$$

Anomaly Report Card

The "Axis of Evil"	3.2σ
Low Power on Large Scales	2.5 σ
Lack of Large-Angular Correlations	3.3 σ
Hemispherical Power Asymmetry	3.Iσ
Point Parity	2.8 σ
Mirror Parity	2σ

Contributions to Large Scale Signal



The Integrated Sachs-Wolfe Effect

[Sachs & Wolfe, ApJ 1967]

$$\delta^{\rm ISW} = -2 \int_{\eta_{\rm LS}}^{\eta_0} \frac{\partial \Phi}{\partial \eta} \,\mathrm{d}\eta$$

- On large scales, this is the main secondary source of anisotropy.
- The energy of the photons is changed when crossing a potential well which evolves in time.



- When the Universe is:
 - Matter dominated \rightarrow linear growth = expansion rate, $\Phi = \text{const.} \rightarrow \text{no ISW.}$
 - Dark-energy dominated → linear growth < expansion rate,
 Φ ≠ const. → late-time ISW effect.

Estimation of the ISW Contribution

- Use galaxy surveys to estimate the local density perturbations.
- Gravitational potential is given by the Poisson equation $\nabla^2 \Phi(\mathbf{x}, t) = 4\pi G \bar{\rho}(t) a^2 \delta(\mathbf{x}, t)$
- Use cross-correlation between CMB and LSS maps.
- Alternatively, build ISW map directly from LSS map. [Francis & Peacock, MNRAS 2010]



ISW Estimation Methods

- Directly from LSS: [Francis & Peacock, MNRAS 2010]
 - Split 2MASS data to 3 redshift shells.
 - Need to estimate the galaxy bias $g_{\ell m} = b_g \delta^m_{\ell m}$ in each redshift shell.
 - Need to assume cosmological parameters.
- Using only cross-correlation: [Rassat et al., A&A 2013; Rassat & Starck, A&A 2013; Rassat et al., 2014, arXiv:1405.1844]
 - ISW map is independent of galaxy bias.
 - Spectra can be calculated assuming cosmological parameters.
 - Spectra can also be estimated from the data. This assumes the form of the ISW signal, not its presence.

$$\delta_{\ell m}^{\rm ISW} = \frac{C_{\ell}^{\rm g, ISW}}{C_{\ell}^{\rm gg}} g_{\ell m}$$

ISW Estimation Methods (cont.)

• Using also CMB data:

$$\delta^{\rm obs} = \delta^{\rm p} + \delta^{\rm ISW}$$

[Barreiro et al., 2008, 2013; Planck Collaboration XIX, 2013, arXiv:1303.5079]

$$\delta_{\ell m}^{\rm ISW} = \frac{C_{\ell}^{\rm g,ISW}}{C_{\ell}^{\rm gg}} g_{\ell m} + \frac{\widetilde{C}_{\ell}^{\rm ISW,ISW}}{\widetilde{C}_{\ell}^{\rm ISW,ISW} + C_{\ell}^{\rm pp}} \left[\delta_{\ell m}^{\rm obs} - \frac{C_{\ell}^{\rm g,ISW}}{C_{\ell}^{\rm gg}} g_{\ell m} \right]$$

$$\widetilde{C}_{\ell}^{\text{ISW,ISW}} = C_{\ell}^{\text{ISW,ISW}} - \frac{\left(C_{\ell}^{\text{g,ISW}}\right)^2}{C_{\ell}^{\text{gg}}}$$

- Maximum likelihood solution, assuming a Gaussian prior.
- Can be extended to include more correlated datasets, such as the gravitational lensing potential map.

[Manzotti & Dodelson, 2014, arXiv:1407.5623]

Existing Works

Work	Method	Data	Tested Anomalies	
Francis & Peacock, 2010	Direct integration	2MASS (3 slices)	✓ WMAP	
Barreiro et al., 2013	Cross + CMB	WMAP + NVSS	×	
Planck Collaboration XIX	Cross + CMB	(i) Planck + NVSS (ii) Planck + Lensing	×	
Rassat et al., 2013; 2014	Only Cross	NVSS + 2MASS	✔ WMAP, Planck	
Manzotti & Dodelson, 2014	Cross (2 tracers) + CMB	Planck + NVSS - Lensing	×	

Currently only $\ell \geq 10$

Anomalies After ISW Subtraction

		Francis & Peacock, 2010	Rassat et <i>al.</i> , 2014
The "Axis of Evil"	3.2 <i>σ</i>	×	×
Low Power on Large Scales	2.5 <i>σ</i>	×	~
Lack of Large-Angular Correlations	3.3 <i>σ</i>		
Hemispherical Power Asymmetry	3.I <i>σ</i>		
Point Parity	2.8σ		
Mirror Parity	2σ		×

Summary & Future Prospects

- Large scale anomalies can provide clues to the very early Universe.
 - "Fight against cosmic variance."
 - Significance is debated and sometimes strongly depends on dataset, estimator choice and masking technique.
 - A-posteriori choices can hinder analyses.
- Planck CMB maps are soon becoming extremely clean of Galactic foregrounds.
 - Could eliminate the need for masking!
- ISW is the main secondary source of anisotropy on large scales.
 - Will the CMB anomalies be alleviated once the ISW signal is removed?

- Results may depend on assumptions used in ISW estimation.
- Recovery of the primordial signal should be done even if one doesn't consider the anomalies significant!
- Polarization maps from Planck will allow improvement of the lensing reconstruction on large scales.
- Future LSS surveys are expected to improve ISW reconstruction accuracy.
 - Only 40% correlation signal up to z = 0.3.
 However, 90% signal up to z = 1.3.
 - Need extensive sky coverage to get large scale signal.
 - DES, LSST, Euclid