

Formation and Evolution of Compact Binaries

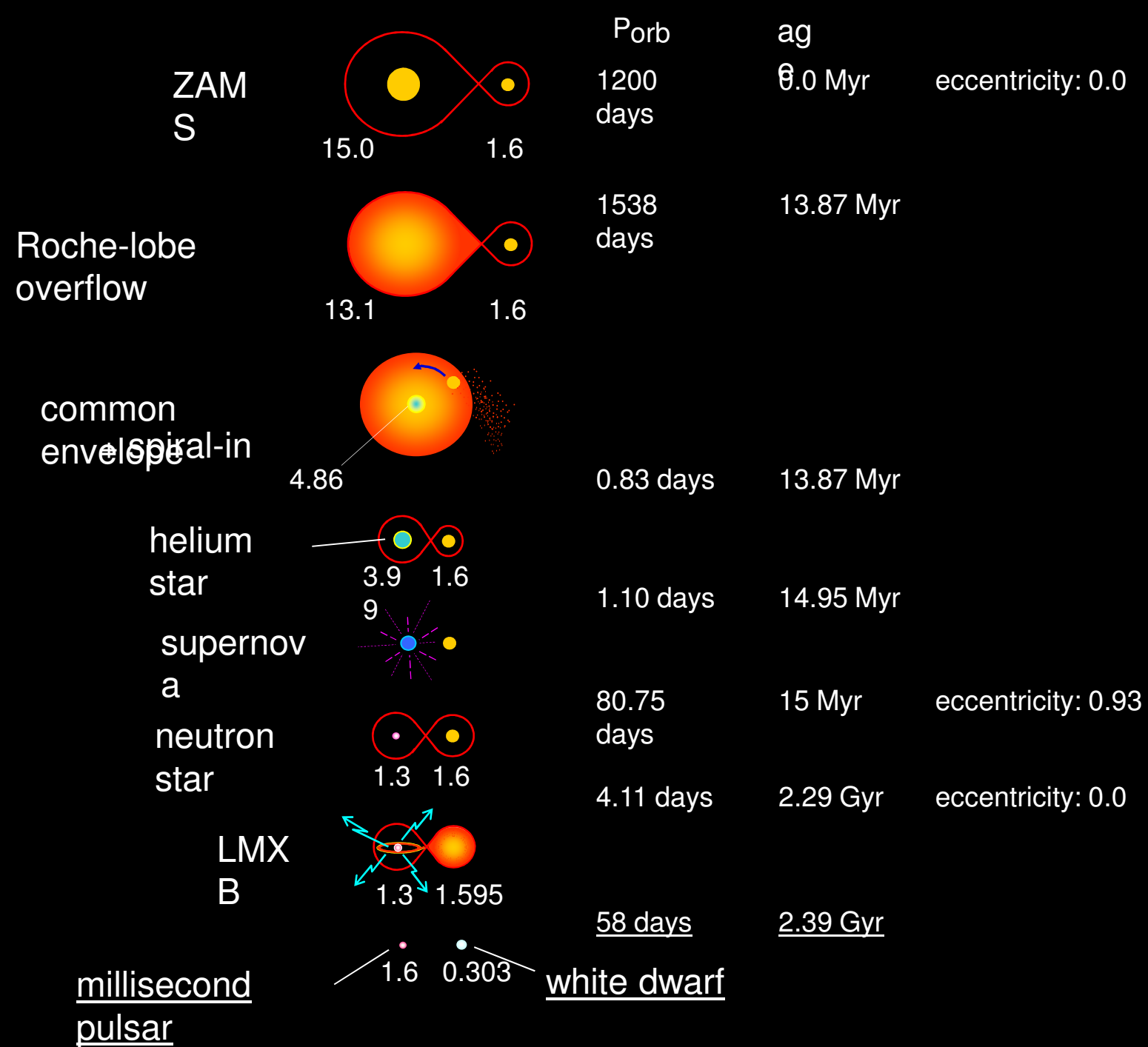


Niels Bohr International Academy Summer School 2009

Thomas Tauris

Lecture topics:

- ❖ HMXBs and LMXBs – a cartoon overview
- ❖ Roche-lobe overflow (RLO): case A, case B, case C
- ❖ Stability criteria for mass transfer / stellar evolution
- ❖ The orbital angular momentum balance equation
- ❖ Common envelope + spiral-in
- ❖ Formation of millisecond pulsars
 - Detailed evolution of LMXBs



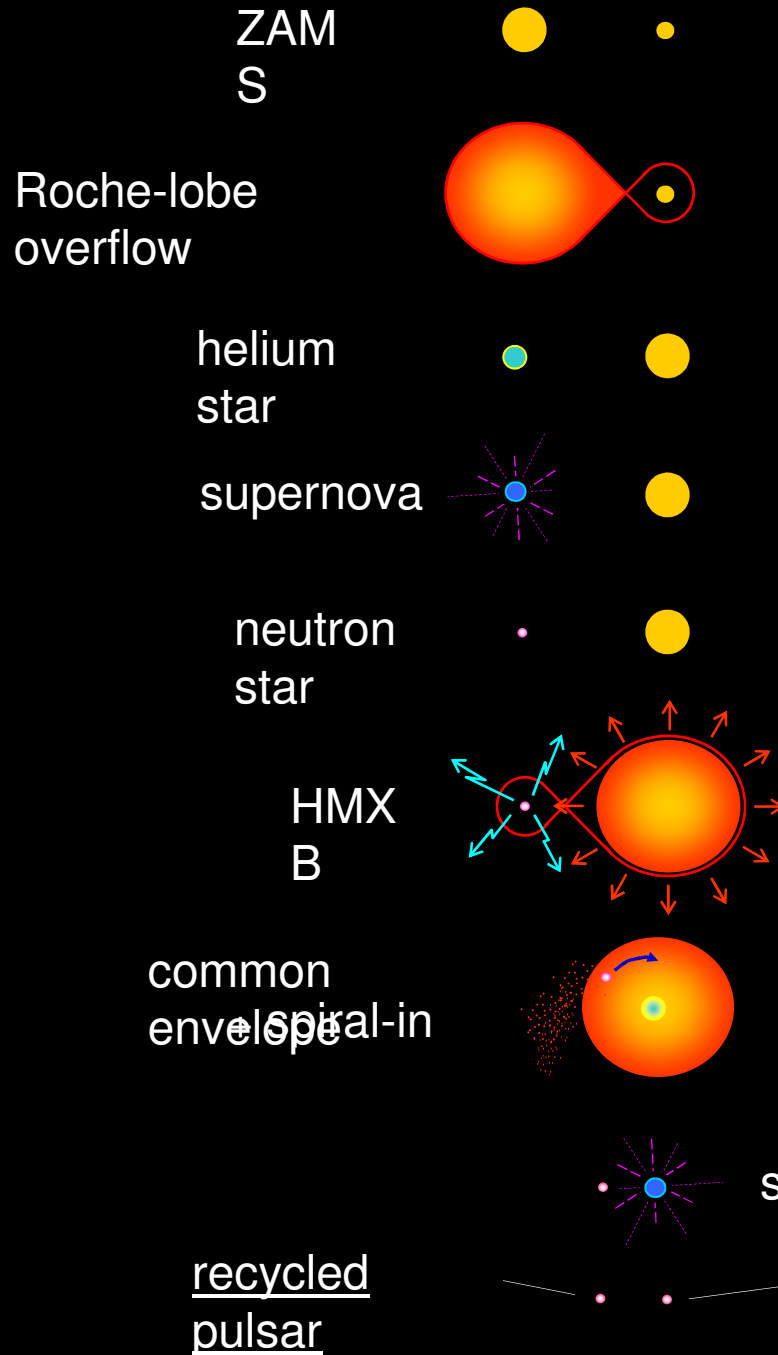
How to make a double pulsar

Perfect playground for:

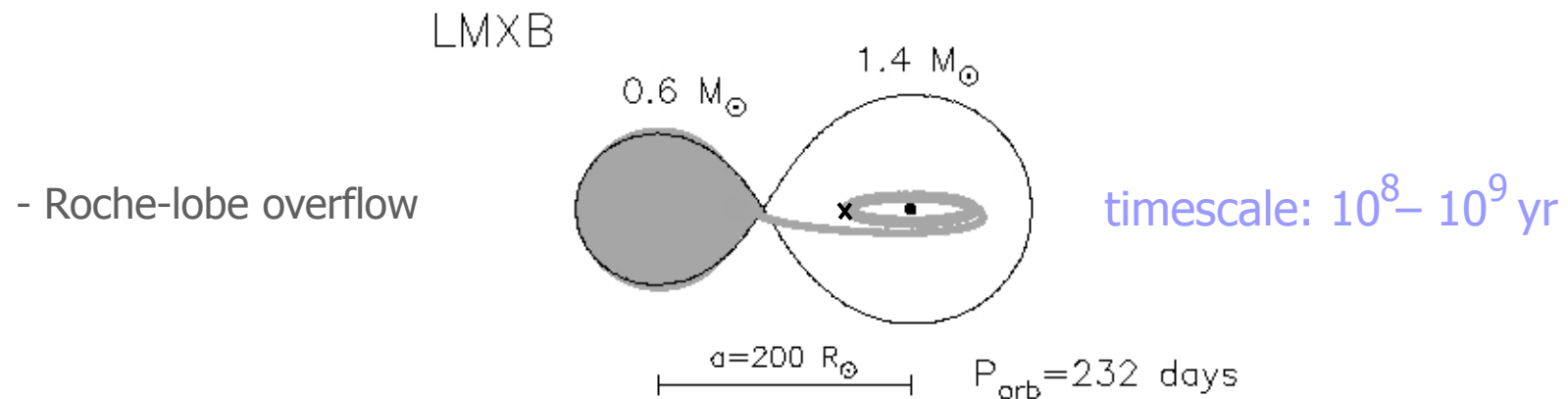
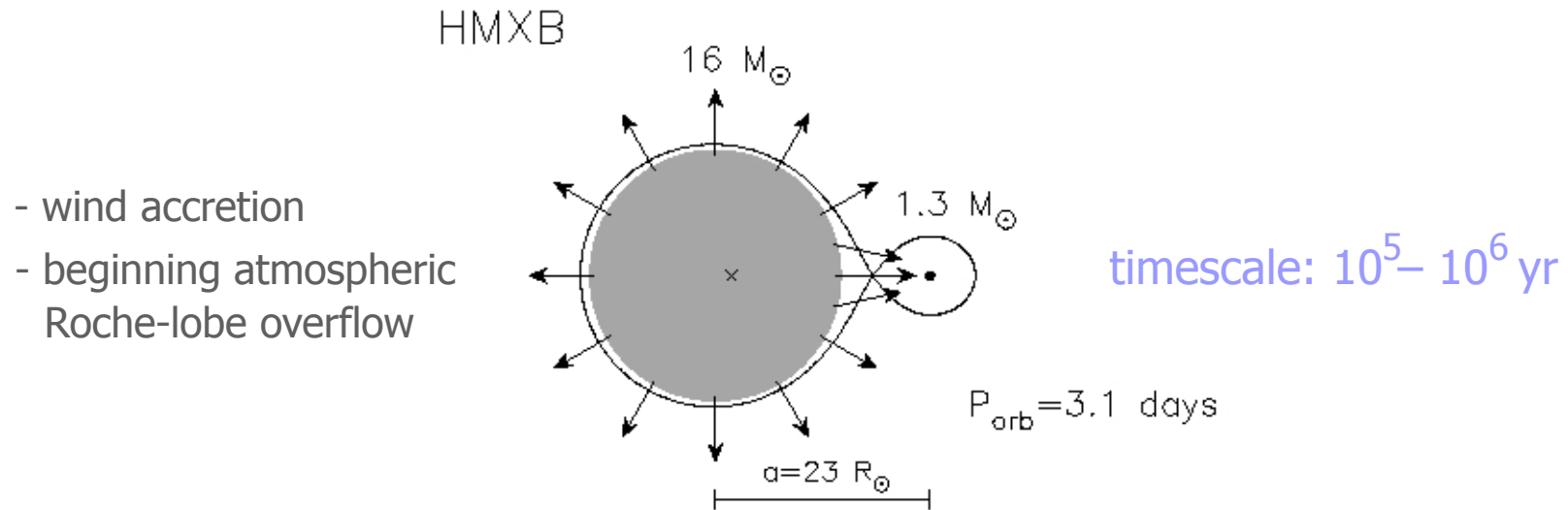
- population synthesis
- Monte Carlo simulations

Things you need to consider:

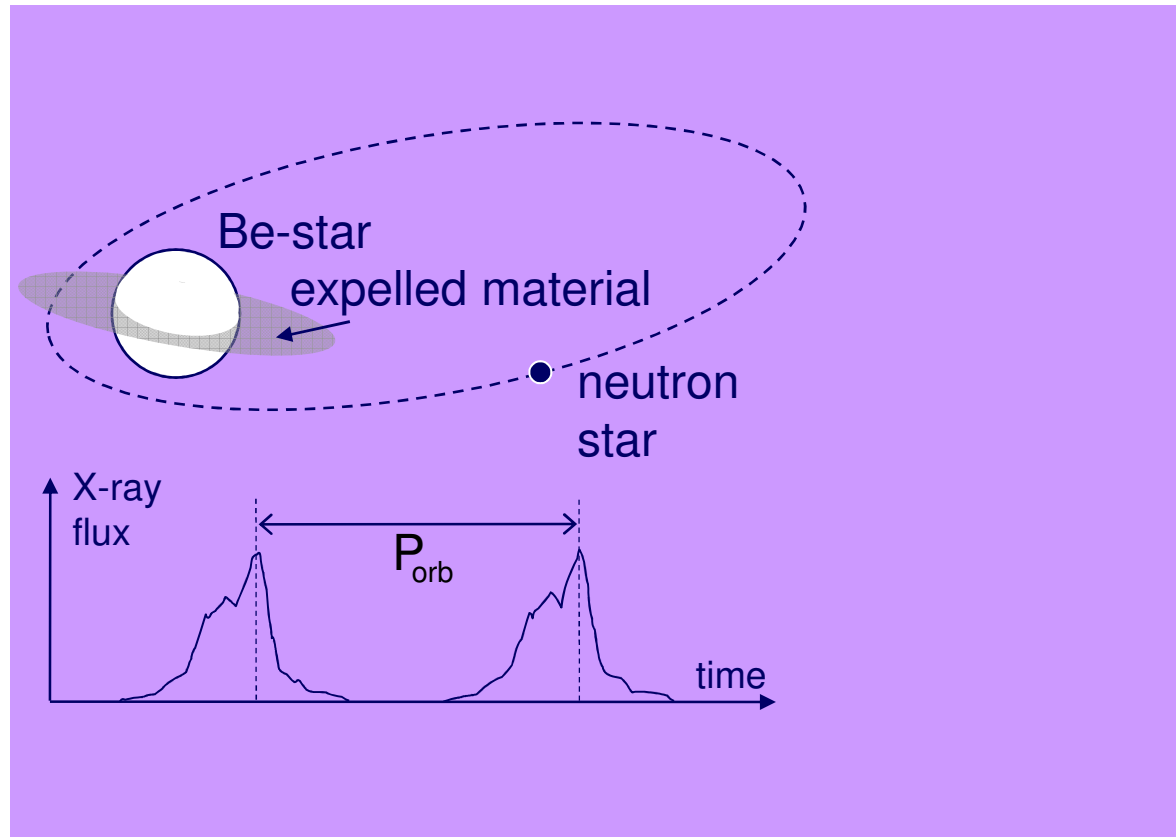
- stellar evolution (incl. He-star evolution)
- various binary interactions (mass transfer, accretion, CE, tidal interactions, GWR)
- asymmetric supernovae



standard X-ray binaries



Be-star X-ray binary



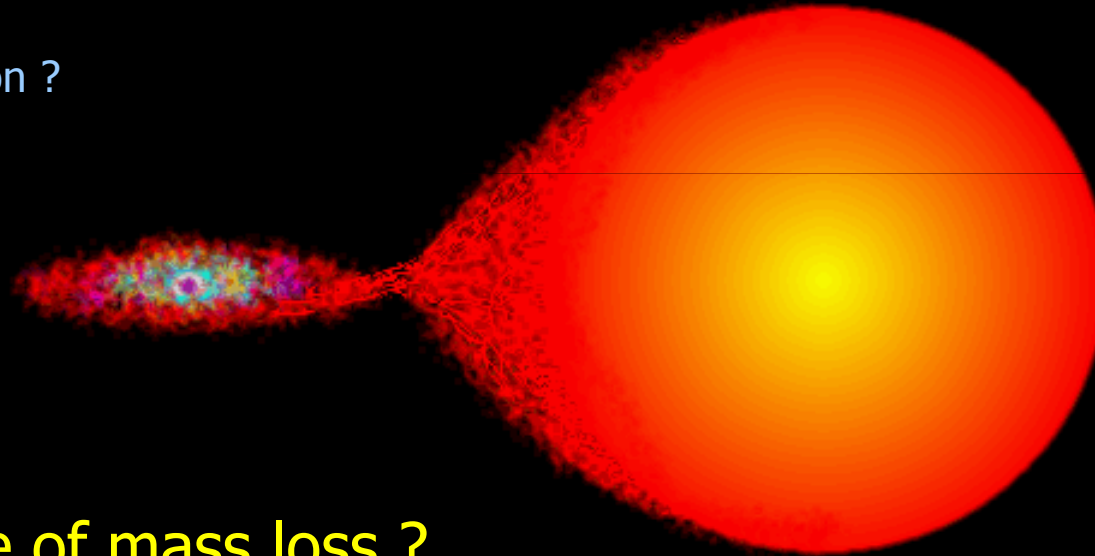
Aspects of Roche-lobe overflow

Accretion ?

super-Eddington ?
jet ?
B-field, spin ?

Stability ?

response of donor star ?
dynamically stable ?



Mode of mass loss ?

specific orbital angular momentum ?

Equipotential surfaces

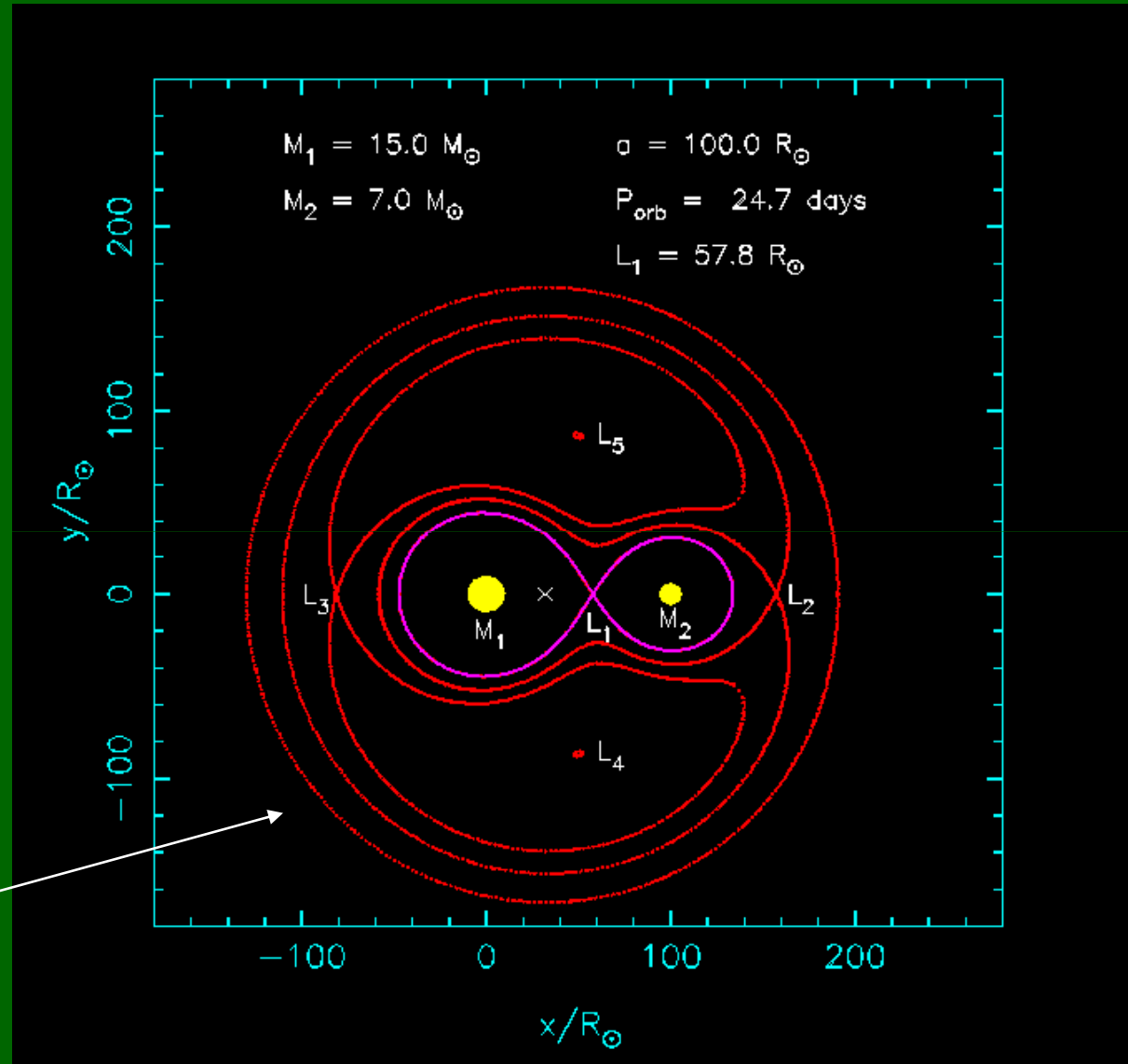
effective gravitational potential:

$$\Phi = -\frac{GM_1}{r_1} - \frac{GM_2}{r_2} - \frac{\Omega^2 r_3^2}{2}$$

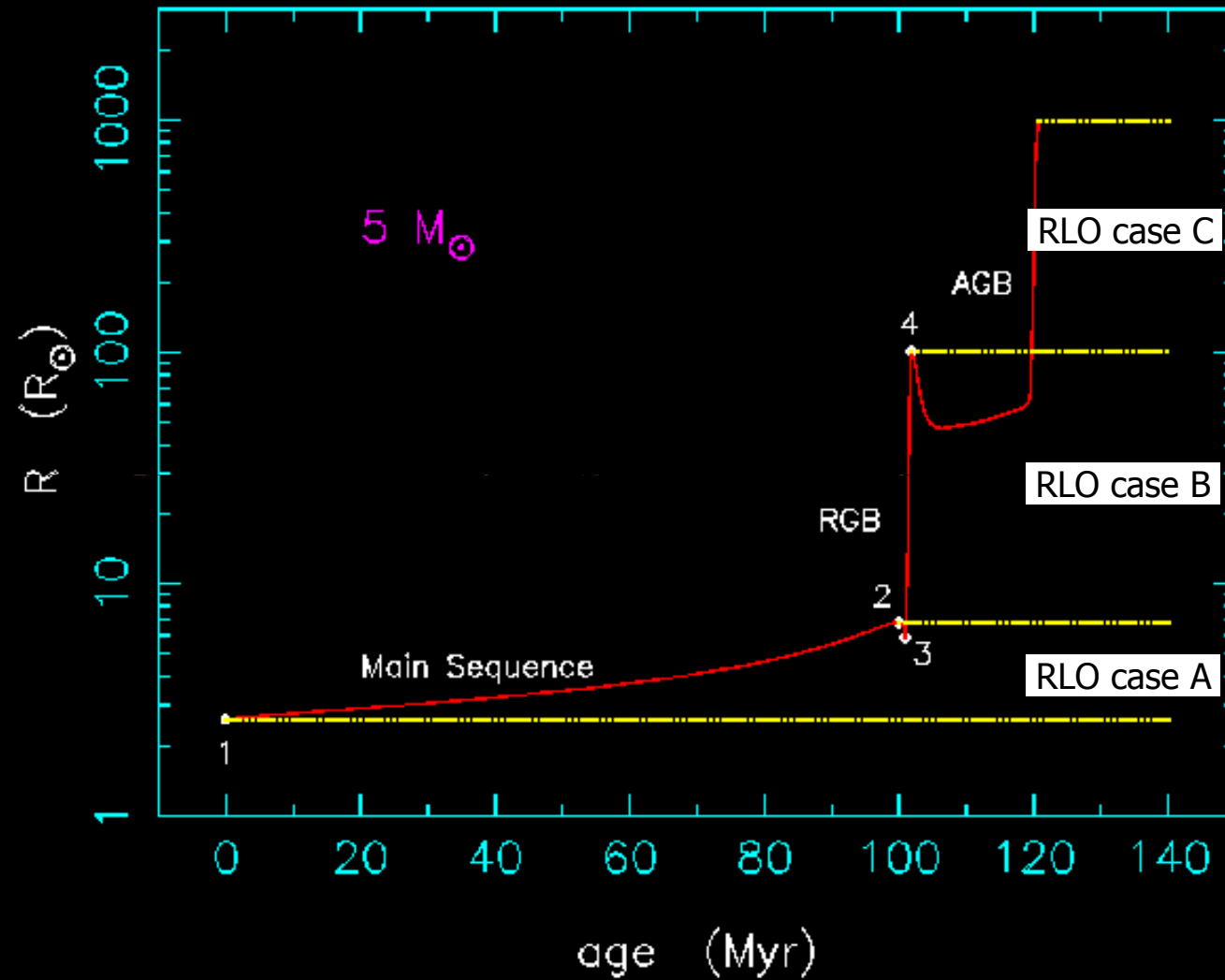
Roche-lobe radius:

$$\frac{R_L}{a} = \frac{0.49 q^{2/3}}{0.6 q^{2/3} + \ln(1 + q^{1/3})}$$

co-moving frame



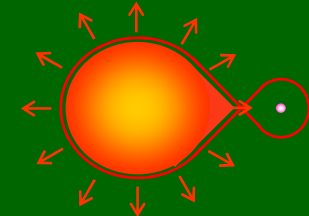
Stellar evolution



Stability criteria for mass transfer

exponents of radius to mass: $R \propto M^\zeta$

$$\zeta_{donor} \equiv \frac{\partial \ln R_2}{\partial \ln M_2} \quad \wedge \quad \zeta_L \equiv \frac{\partial \ln R_L}{\partial \ln M_2}$$



adiabatic or thermal response of the donor star to mass loss

initial stability criteria: $\zeta_L \leq \zeta_{donor}$

$$\dot{R}_2 = \left. \frac{\partial R_2}{\partial t} \right|_{M_2} + R_2 \zeta_{donor} \frac{\dot{M}_2}{M_2}$$

nuclear burning

$$\dot{R}_L = \left. \frac{\partial R_L}{\partial t} \right|_{M_2} + R_L \zeta_L \frac{\dot{M}_2}{M_2}$$

tidal spin-orbit couplings
gravitational wave radiation

$\dot{R}_2 = \dot{R}_L$ yields mass loss rate!

The Orbital Angular Momentum Balance Equation

$$J_{orb} = \frac{M_1 M_2}{M} \Omega a^2 \sqrt{1-e^2}$$

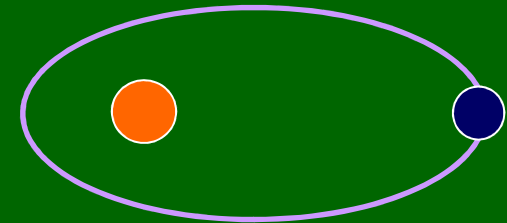
orbital angular momentum

logarithmic differentiation

($e=0$, tidal circularization)



$$\frac{\dot{a}}{a} = 2 \frac{\dot{J}_{orb}}{J_{orb}} - 2 \frac{\dot{M}_1}{M_1} - 2 \frac{\dot{M}_2}{M_2} + \frac{\dot{M}_1 + \dot{M}_2}{M}$$

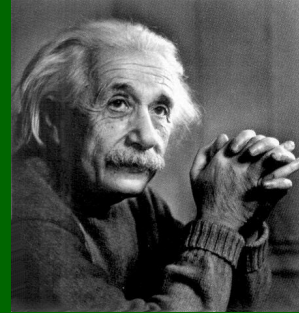


$$J_{orb} = |\vec{r} \times \vec{p}|$$

$$\frac{\dot{J}_{orb}}{J_{orb}} = \frac{\dot{J}_{gwr}}{J_{orb}} + \frac{\dot{J}_{mb}}{J_{orb}} + \frac{\dot{J}_{ls}}{J_{orb}} + \frac{\dot{J}_{ml}}{J_{orb}}$$

Gravitational wave radiation:

$$\frac{\dot{J}_{gwr}}{J_{orb}} = - \frac{32}{5} \frac{G^3}{c^5} \frac{M_1 M_2 M}{a^4} \quad s^{-1}$$



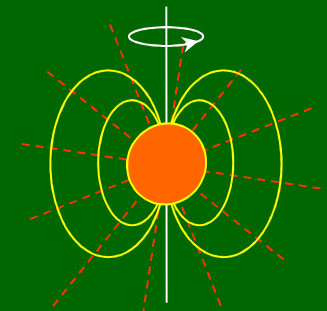
Magnetic braking:

$$\frac{\dot{J}_{mb}}{J_{orb}} \propto - \frac{k^2 R^4}{a^5} \frac{GM^3}{M_1 M_2}$$

(uncertain)

Low-mass stars: magnetic wind!
 \Rightarrow loss of spin angular momentum

In tight binaries the system
 is tidally locked (synchronized)
 and spin-orbit couplings operate
 \Rightarrow loss of orbital angular momentum



$$M_2 < 1.5 M_{\odot}$$

$$P_{orb} < 2 \text{ days}$$

Spin-orbit couplings:

$$\frac{\dot{J}_{ls}}{J_{orb}}$$

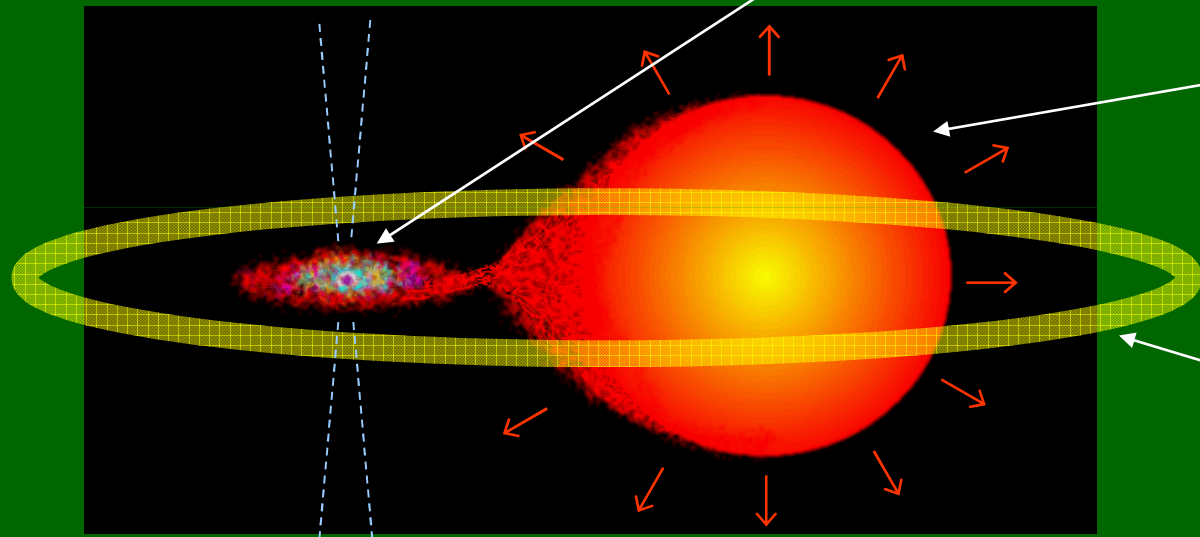
- fx . change in stellar moment of inertia
 (as a result of nuclear burning or mass loss)

Mass loss:

$$\frac{\dot{J}_{ml}}{J_{orb}} = \frac{\alpha + \beta q^2 + \delta \gamma (1+q)^2}{1+q} \frac{\dot{M}_2}{M_2}$$

$$\beta = \max\left(\frac{|\dot{M}_2| - \dot{M}_{Edd}}{|\dot{M}_2|}, 0\right)$$

β : mass ejected from accretor
- fx. in a relativistic jet
(isotropic re-emission)



α : direct fast wind

δ : mass loss via circumbinary
coplanar toroid with
radius: $\gamma^2 a$

accretion efficiency: $\epsilon = 1 - \alpha - \beta - \delta$ $(\partial M_{NS} = -\epsilon \partial M_2)$

Common envelope + spiral-in evolution

dynamically unstable mass transfer:

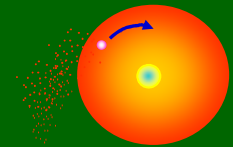
- deep convective envelope of donor star
(rapid expansion in response to mass loss)
- $M_{\text{donor}} > M_{\text{accretor}}$
(orbit shrinks in response to mass loss)



Run-away process!!



common envelope



drag force → dissipation of orb. ang. mom. + deposition of E_{orb} in the envelope

outcome:

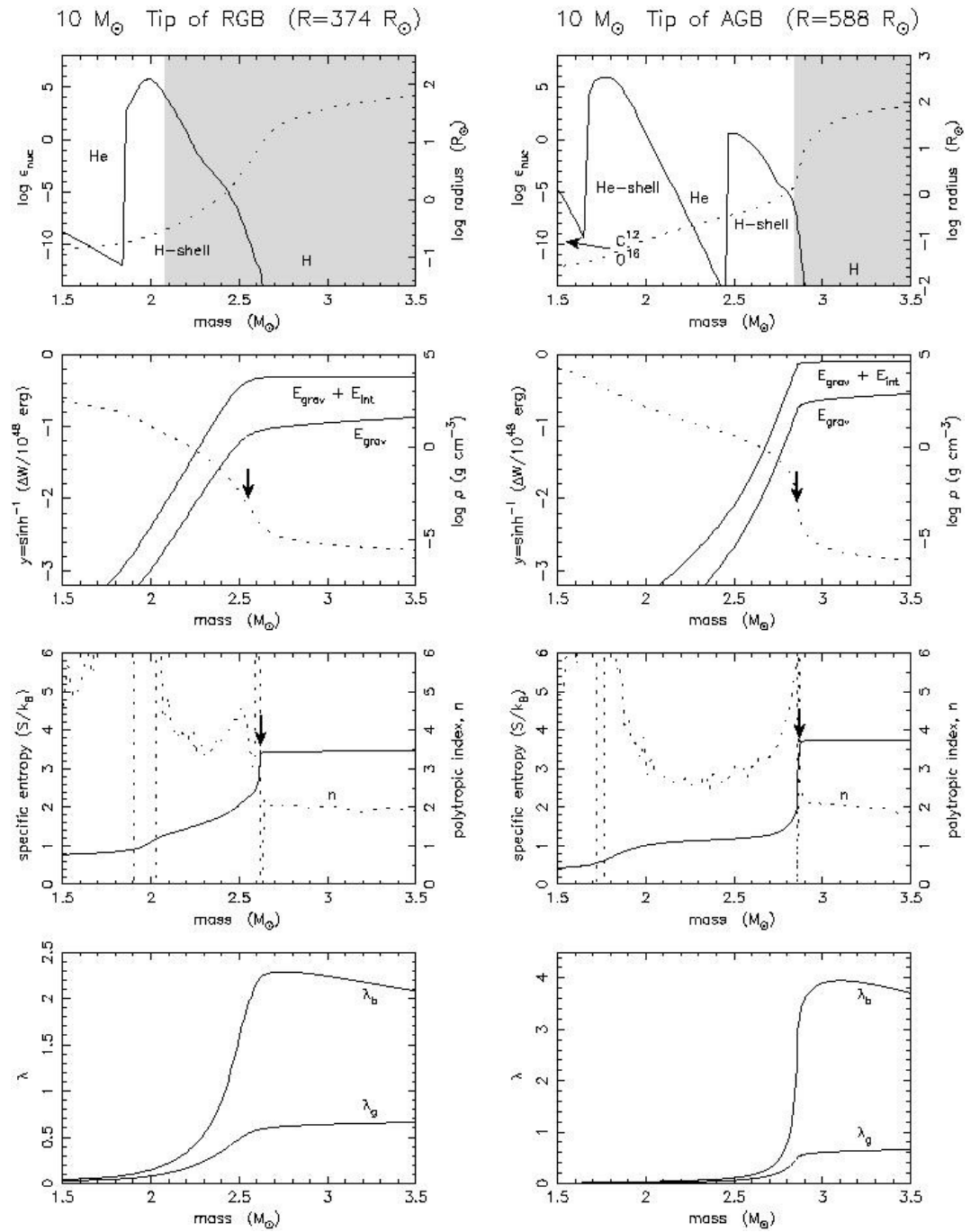
- huge reduction of orbital separation

rejection of stellar envelope
(NS orbiting a naked helium star)

merging of NS + core
(Thorne-Zytkow object / black hole)

Needed to explain observed tight systems!!

How to define M_{core} ?
 - quite important!



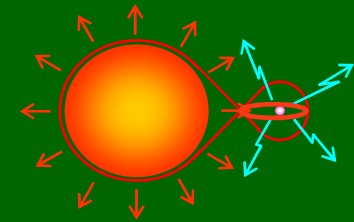
Tauris & Dewi (2001)

Intermediate Mass X-ray Binaries

Why are so few IMXBs observed ?

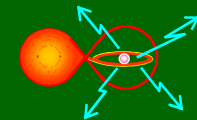
HMXB: wind accretion (beginning atmospheric RLO)
RLO is dynamically unstable and a CE forms

$$M_2 > 10 M_{\odot}$$



LMXB: stable RLO

$$M_2 \leq 1.5 M_{\odot}$$



IMXB: wind accretion is too weak, and
RLO is often unstable (or very short)

Formation of millisecond pulsars (MSPs)

First pulsar discovered in 1967 (today ~1800 pulsars are known)

First binary pulsar discovered in 1974: PSR 1913+16 (today ~50 systems are known)

First millisecond pulsar discovered in 1982: PSR 1937+21 ($P=1.5$ msec)

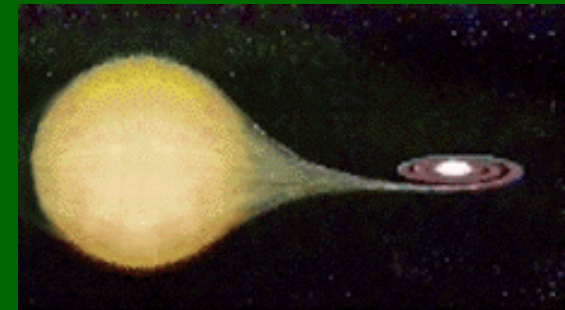
Millisecond pulsar characteristics:

P-Pdot diagram

- high incidence of binaries among millisecond pulsars
- rapid spin ($P < 20$ msec)
- relatively weak magnetic fields ($B=10^8 - 10^9$ G)



old neutron stars which have been "recycled"
via accretion of mass and angular momentum



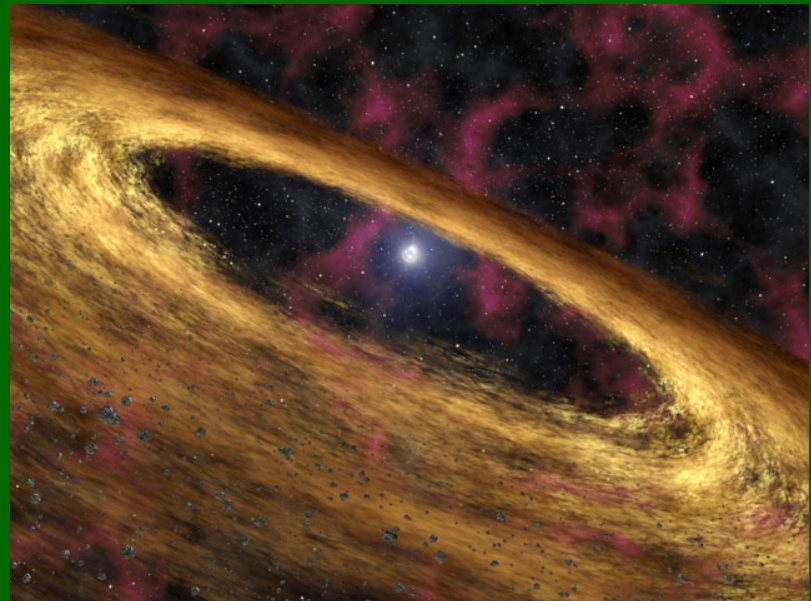
Single millisecond pulsars

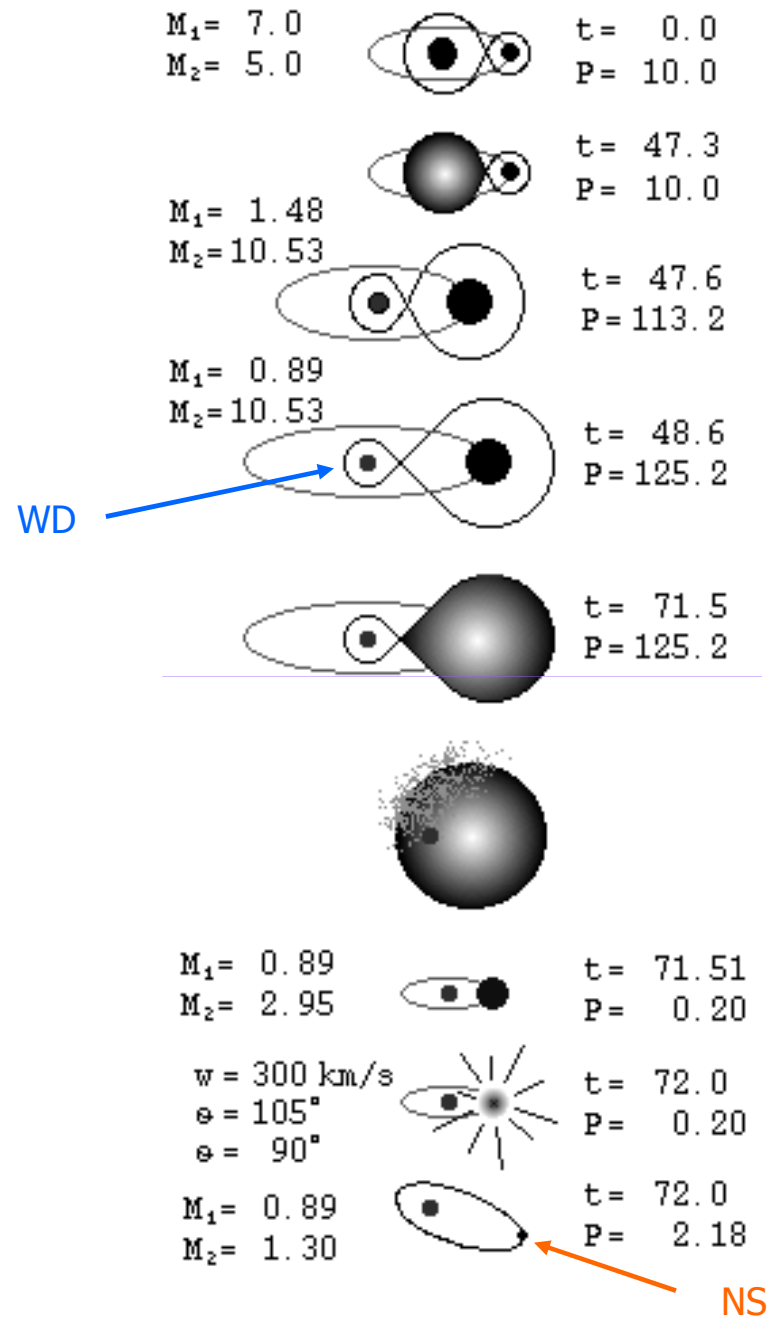
All single millisecond pulsars are born in a binary system.

Once a recycled millisecond pulsar turns on its emission of ultra-relativistic particles, it is often able to completely evaporate its companion and thus end up as an isolated millisecond pulsar.

Observational evidence:

- ✓ eclipsing MSPs with $0.02 M_{\odot}$ companions
- ✓ the “planetary pulsar”, PSR 1257+12





Summary:

- ❖ To study the formation and evolution of compact binaries one must have a detailed knowledge stellar evolution and binary interactions
- ❖ Millisecond pulsars are old “rejuvenated” pulsars
- ❖ Common envelope + spiral-in explain very narrow orbits
- ❖ HMXBs and LMXBs are wellknown b/c they are detected
 - IMXBs are not detected
- ❖ Perfect scenario for Monte Carlo simulations

