

Core-Collapse Supernovae

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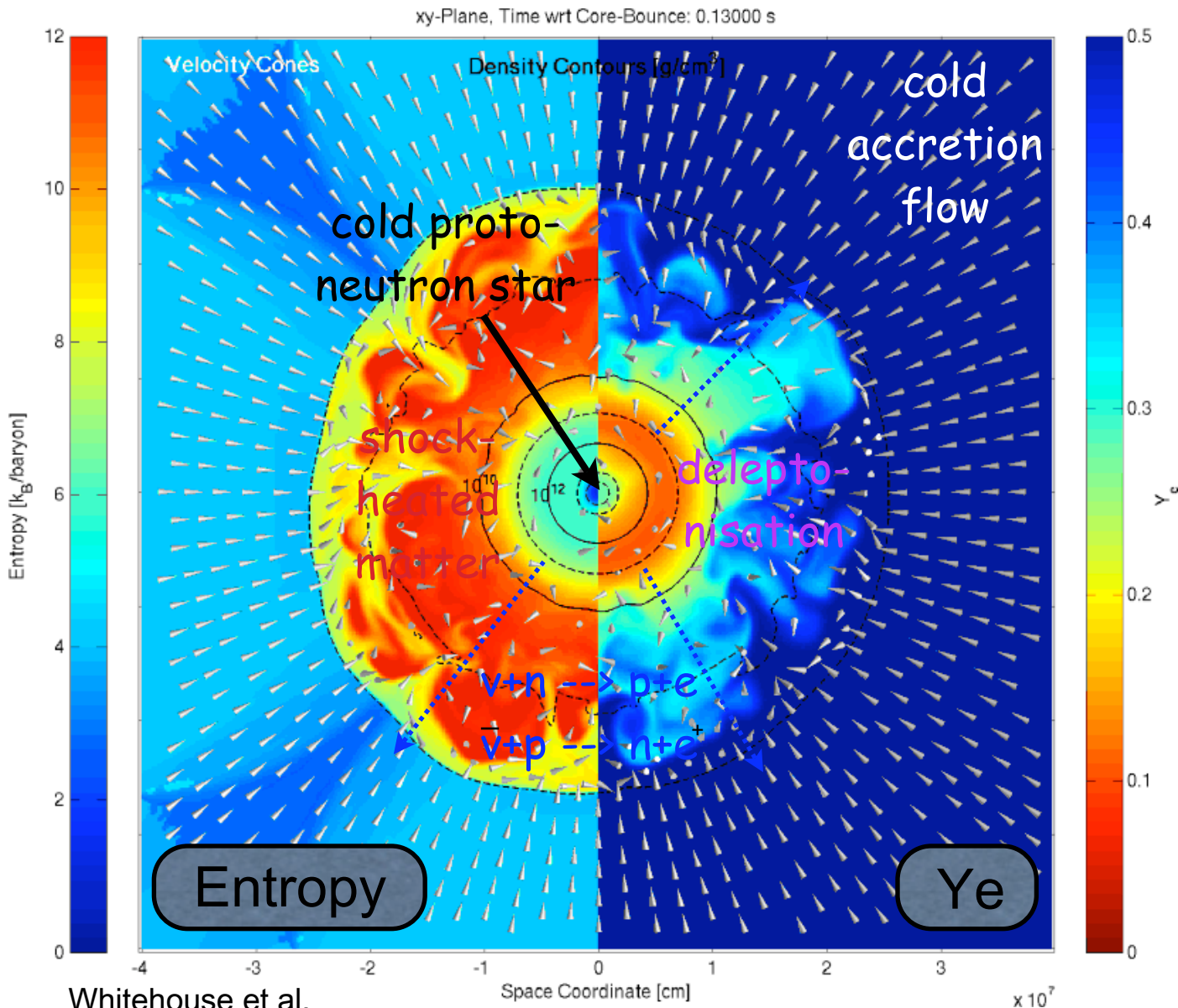
- Collapse phase: Dynamics & ν -interactions
- Postbounce phase: ν -transport & explosion mechanisms
- Models: Approximations & prediction of observables

Large cancellation effects in the total energy budget:

- Huge energy in!
- Huge energy out!
- The rest makes the supernova!

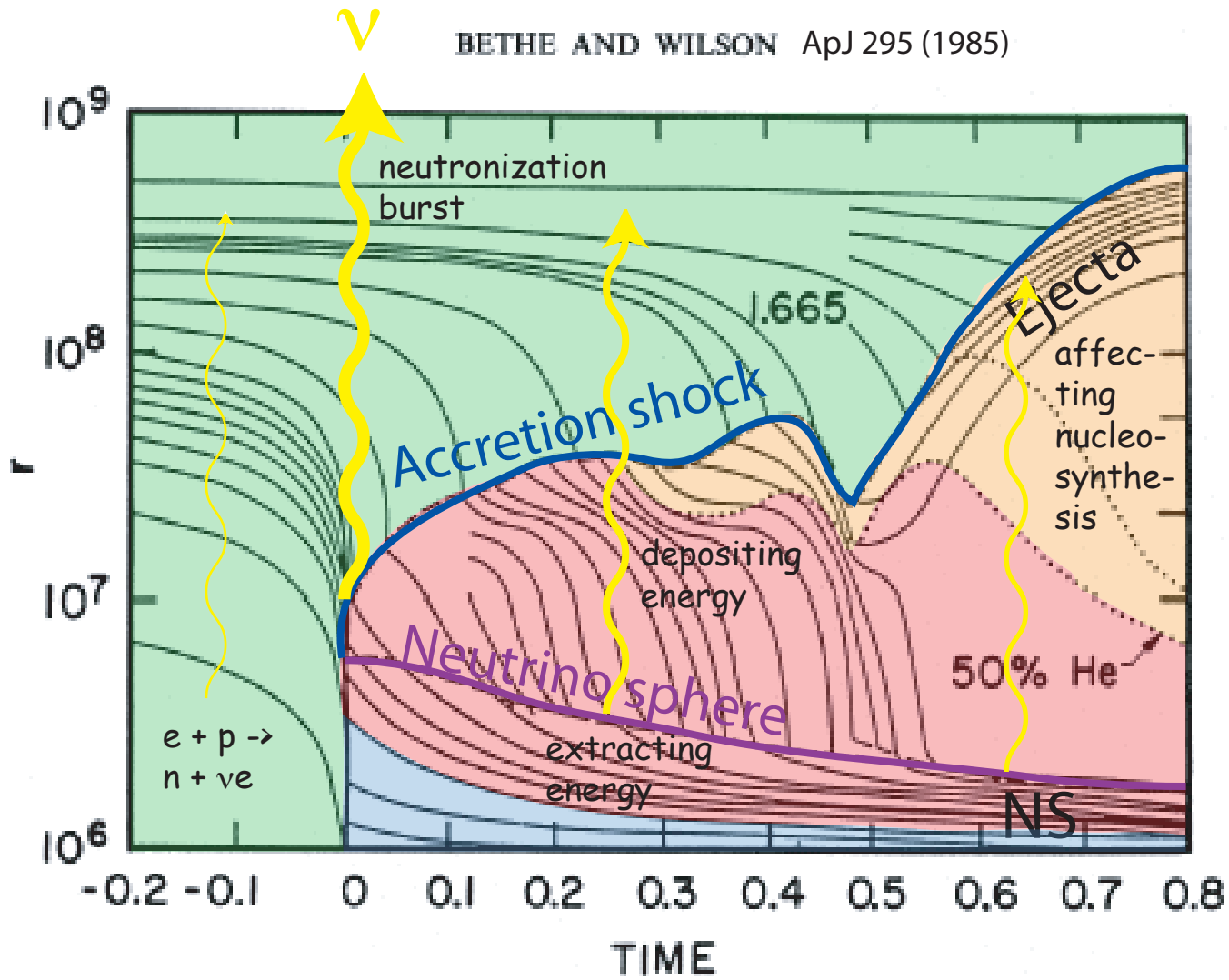
- Leading order contributions from many fields of physics possible...

The Supernova Problem



- How does the collapse of single stars lead to explosions that outshine a galaxy?
- Which new physics is observable in the extreme conditions of matter during the explosion?
- Does the nucleosynthesis of heavy elements explain the abundances on Earth, the Sun and distant stars?

Delayed explosion: 4 phases



collapse phase || postbounce accretion phase | explosion phase
 bounce

Ensemble of nuclei

Cool bulk nuclear matter

Hot dissociated matter

Freeze-out of nuclei

Discussed Explosion Mechanisms



Energy scales:

- Gravitational
~ $3E+53$ erg
- Explosion
~ $1E+51$ erg

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1) Prompt explosion mechanism, E(bounce)

(e.g. Baron et al. 1985)

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2) Neutrino-driven explosion mechanism, E(therm.)

(Colgate 1966, Arnett, Bruenn, Burrows, Mezzacappa ... Marek & Janka 2009)

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3) Magneto-rotational explosion mechanism, E(rot.)

(Bisnovatyi-Kogan 1976, Leblanc & Wilson 1979, ...)

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(Burrows et al. 2006)

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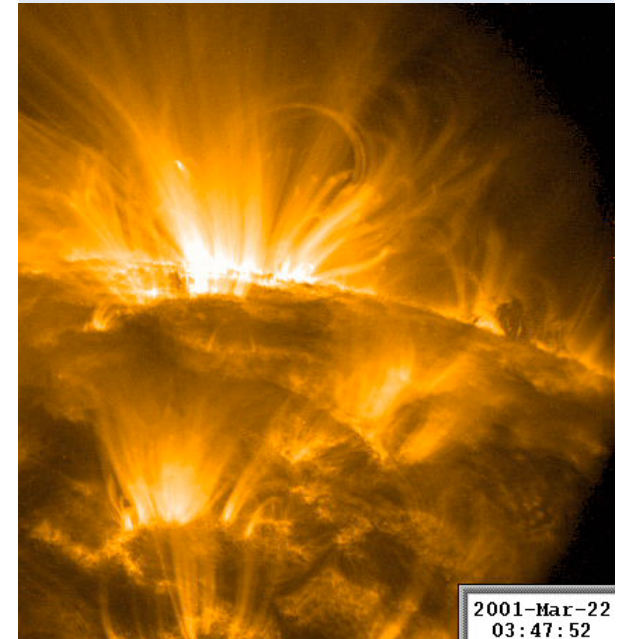
(Burrows et al. 2006)

5) Magneto-sonic/viscous expl. mech., E(buoyancy)

(Akiyama et al. 2003, Thompson et al. 2003, Socrates et al. 2005)

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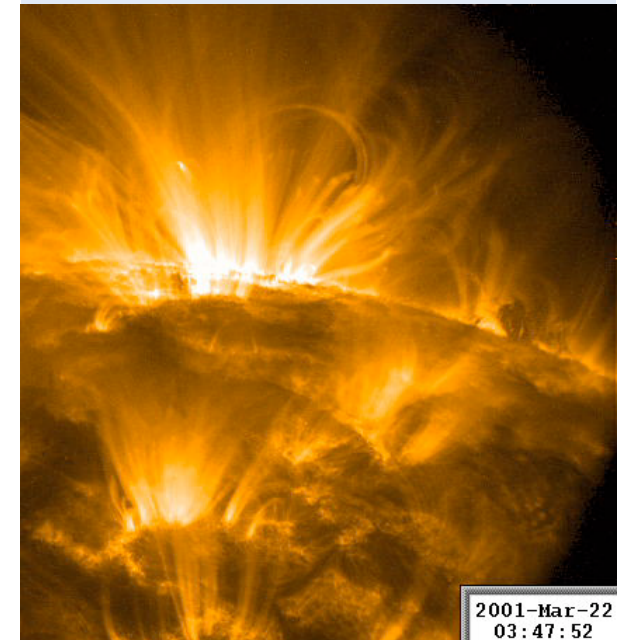
(Akiyama et al. 2003, Thompson et al. 2003, Socrates et al. 2005)

6) Phase transition induced expl. mech., E(compact)

(Migdal et al. 1971, ... Sagert et al. 2009)

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Input <--> Model <--> Observation

Physics

Modelling challenges

Observations

Reaction network

- uncertainties
- stiff partial diff'eqs.

Magneto-hydrodynamics

- resolution
- time scales

Gravity

- NR: elliptic equations
- GR: metric/horizons

Radiative transfer

- dimensionality
- non-locality

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Lightcurve
Polarisation
Spectra/abundances
Event rates
Galactic Evolution
NS kick & spin
neutrino signal
Grav. wave signal
GRB's

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Transient searches

Input <--> Model <--> Observation

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Transient searches

→ Quality must match!

Solving the Boltzmann equation



$$\frac{\partial F}{\alpha c \partial t} + \frac{\partial (4\pi r^2 \alpha \rho \mu F)}{\alpha \partial m}$$

$$= \frac{\dot{j}}{\rho} - \tilde{\chi} F$$

$$\frac{\partial Y_e}{\partial t} = -\frac{2\pi m_B}{h^3 c^2} \int E^2 dE d\mu \left(\frac{\dot{j}}{\rho} - \tilde{\chi} F \right) \quad \frac{\partial e}{\partial t} = \dots \quad \frac{\partial u}{\partial t} = \dots$$

(Mezzacappa & Bruenn 1993, Liebendörfer 2000, Liebendörfer et al. 2004)

Evolution of specific neutrino distr. function:

$$F(t, m, \mu, E) = f(t, r, \mu, E) / \rho$$

=> 3D implicit problem

Comoving metric:

$$ds^2 = -\alpha^2 dt^2 + \left(\frac{1}{\Gamma} \frac{\partial r}{\partial a} \right)^2 + r^2 (d\vartheta^2 + \sin^2 \vartheta d\varphi^2)$$

Stress-energy tensor:

$$\begin{aligned} T^{tt} &= \rho(1 + e + J) \\ T^{ta} = T^{at} &= \rho H \\ T^{aa} &= p + \rho K \\ T^{\vartheta\vartheta} = T^{\varphi\varphi} &= p + \frac{1}{2}\rho(J - K) \end{aligned}$$

Solving the Boltzmann equation



$$\frac{\partial F}{\alpha c \partial t} + \frac{\partial (4\pi r^2 \alpha \rho \mu F)}{\alpha \partial m}$$

$$\begin{aligned} &= \frac{j}{\rho} - \tilde{\chi} F + \frac{1}{h^3 c^4} E^2 \int d\mu' R_{is}(\mu, \mu', E) F(\mu', E) \\ &- \frac{1}{h^3 c^4} E^2 F \int d\mu' R_{is}(\mu, \mu', E) \\ &+ \frac{1}{h^3 c^4} \left[\frac{1}{\rho} - F(\mu, E) \right] \int E'^2 dE' d\mu' \tilde{R}_{nes}^{in}(\mu, \mu', E, E') F(\mu', E) \\ &- \frac{1}{h^3 c^4} F(\mu, E) \int E'^2 dE' d\mu' \tilde{R}_{nes}^{out}(\mu, \mu', E, E') \left[\frac{1}{\rho} - F(\mu', E') \right] \end{aligned}$$

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$$\begin{aligned} & \frac{\partial F}{\alpha c \partial t} + \frac{\partial (4\pi r^2 \alpha \rho \mu F)}{\alpha \partial m} + \Gamma \left(\frac{1}{r} - \frac{\partial \alpha}{\alpha \partial r} \right) \frac{\partial [(1 - \mu^2) F]}{\partial \mu} \\ & + \left[-\mu \Gamma \frac{\partial \alpha}{\alpha \partial r} \right] \frac{1}{E^2} \frac{\partial (E^3 F)}{\partial E} \\ & = \frac{j}{\rho} - \tilde{\chi} F + \frac{1}{h^3 c^4} E^2 \int d\mu' R_{is}(\mu, \mu', E) F(\mu', E) \\ & - \frac{1}{h^3 c^4} E^2 F \int d\mu' R_{is}(\mu, \mu', E) \\ & + \frac{1}{h^3 c^4} \left[\frac{1}{\rho} - F(\mu, E) \right] \int E'^2 dE' d\mu' \tilde{R}_{nes}^{in}(\mu, \mu', E, E') F(\mu', E) \\ & - \frac{1}{h^3 c^4} F(\mu, E) \int E'^2 dE' d\mu' \tilde{R}_{nes}^{out}(\mu, \mu', E, E') \left[\frac{1}{\rho} - F(\mu', E') \right] \\ & \frac{\partial Y_e}{\partial t} = -\frac{2\pi m_B}{h^3 c^2} \int E^2 dE d\mu \left(\frac{j}{\rho} - \tilde{\chi} F \right) \quad \frac{\partial e}{\partial t} = \dots \quad \frac{\partial u}{\partial t} = \dots \end{aligned}$$

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$$\begin{aligned}
 & \frac{\partial F}{\alpha c \partial t} + \frac{\partial (4\pi r^2 \alpha \rho \mu F)}{\alpha \partial m} + \Gamma \left(\frac{1}{r} - \frac{\partial \alpha}{\alpha \partial r} \right) \frac{\partial [(1 - \mu^2) F]}{\partial \mu} \\
 & + \left(\frac{\partial \ln \rho}{\alpha c \partial t} + \frac{3u}{r c} \right) \frac{\partial [\mu (1 - \mu^2) F]}{\partial \mu} \\
 & + \left[\mu^2 \left(\frac{\partial \ln \rho}{\alpha c \partial t} + \frac{3u}{r c} \right) - \frac{1u}{r c} - \mu \Gamma \frac{\partial \alpha}{\alpha \partial r} \right] \frac{1}{E^2} \frac{\partial (E^3 F)}{\partial E} \\
 & = \frac{j}{\rho} - \tilde{\chi} F + \frac{1}{h^3 c^4} E^2 \int d\mu' R_{is}(\mu, \mu', E) F(\mu', E) \\
 & - \frac{1}{h^3 c^4} E^2 F \int d\mu' R_{is}(\mu, \mu', E) \\
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Finite differencing of time evolution



Forward differencing:
evaluate slope with
current state vector

Backw. differencing:
evaluate slope with
future state

Let's see...

Explicit finite differencing



Forward differencing:
evaluate slope with
current state vector

- simple
- accurate for small time steps
- limited by characteristic time scale

Go faster...

Explicit finite differencing



Forward differencing:
evaluate slope with
current state vector

- simple
- inaccurate for large time steps
- even catastrophic!

Think...

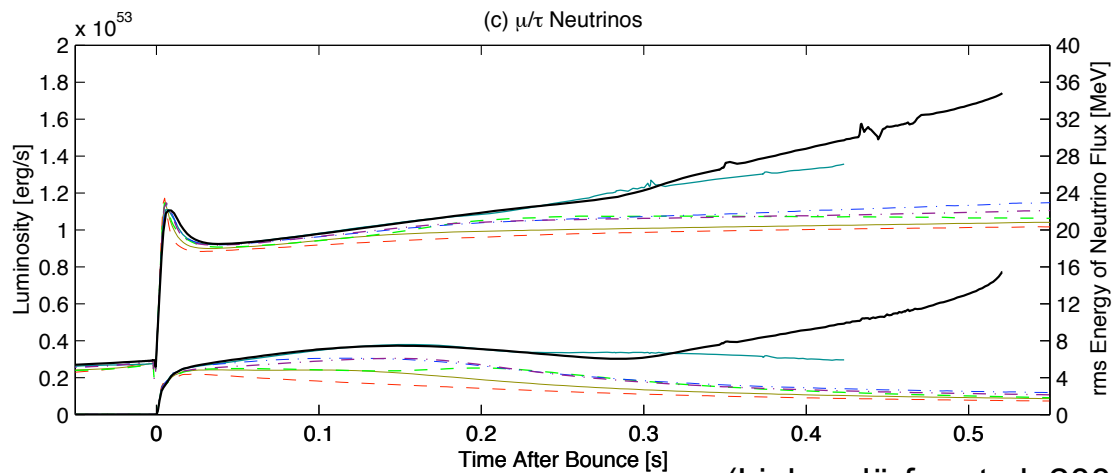
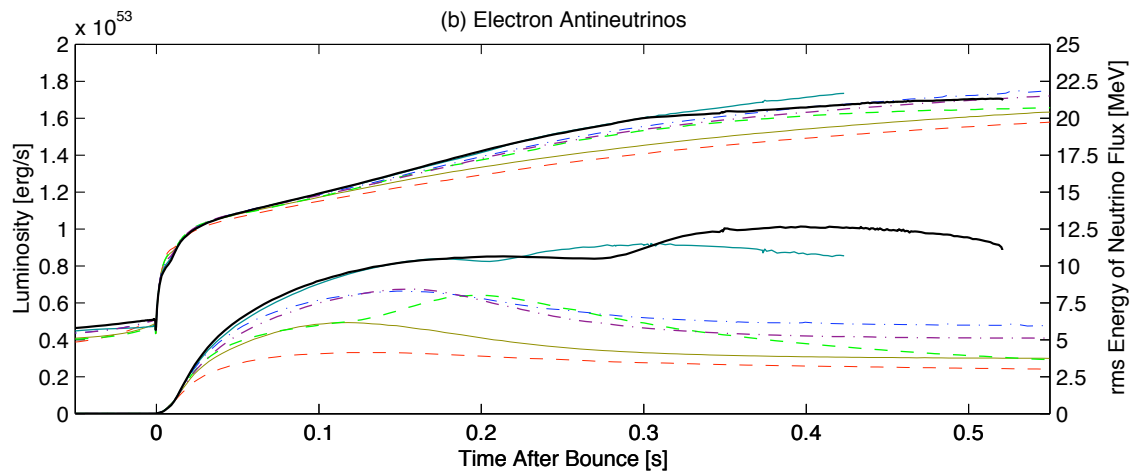
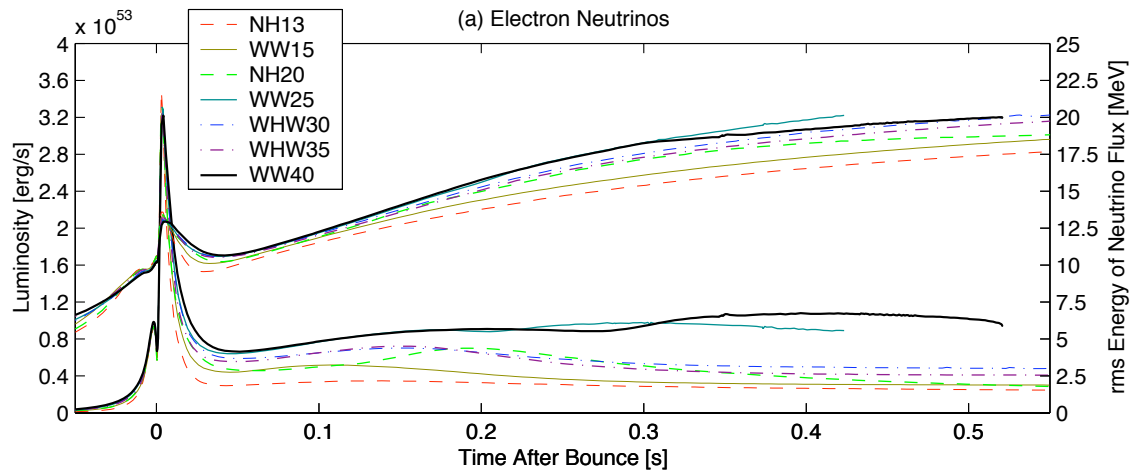
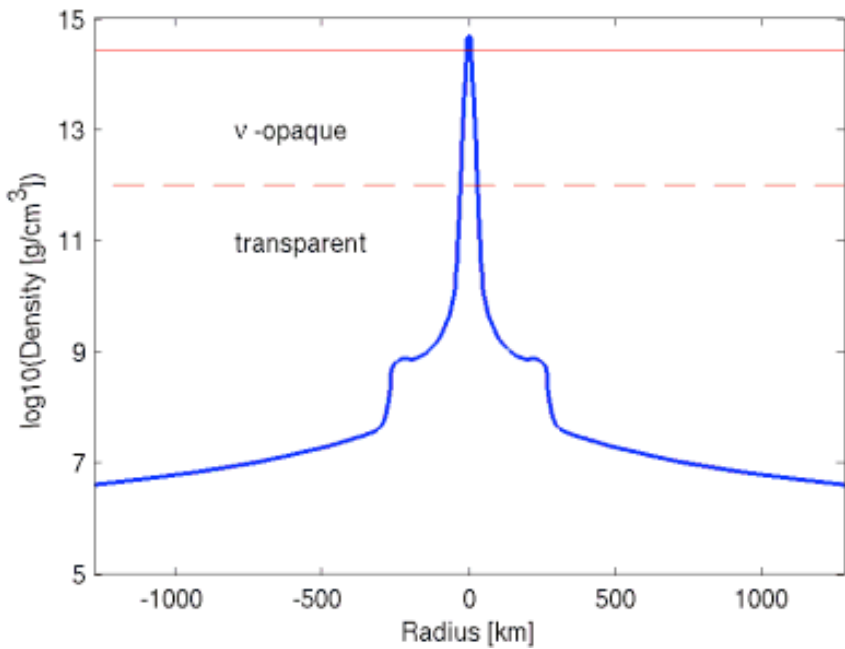
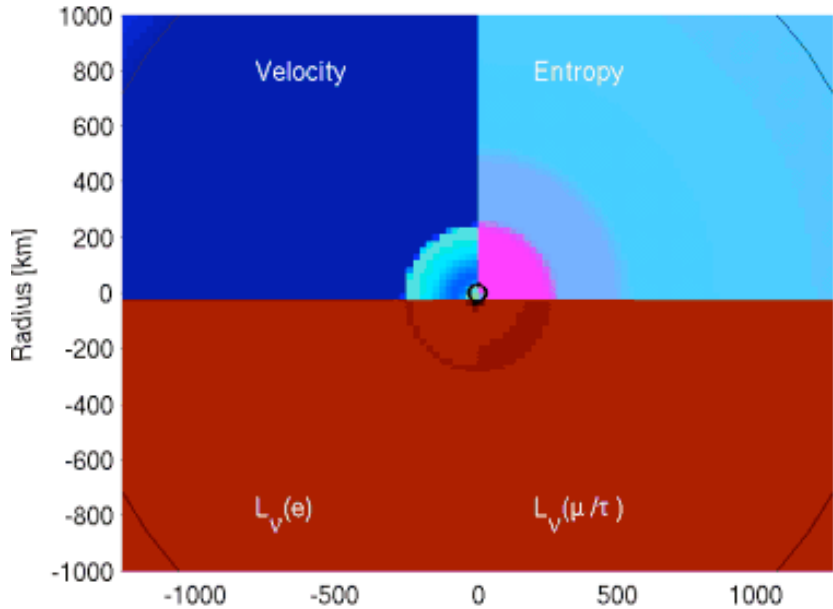
Implicit finite differencing



Backw. differencing:
evaluate slope with
future state vector

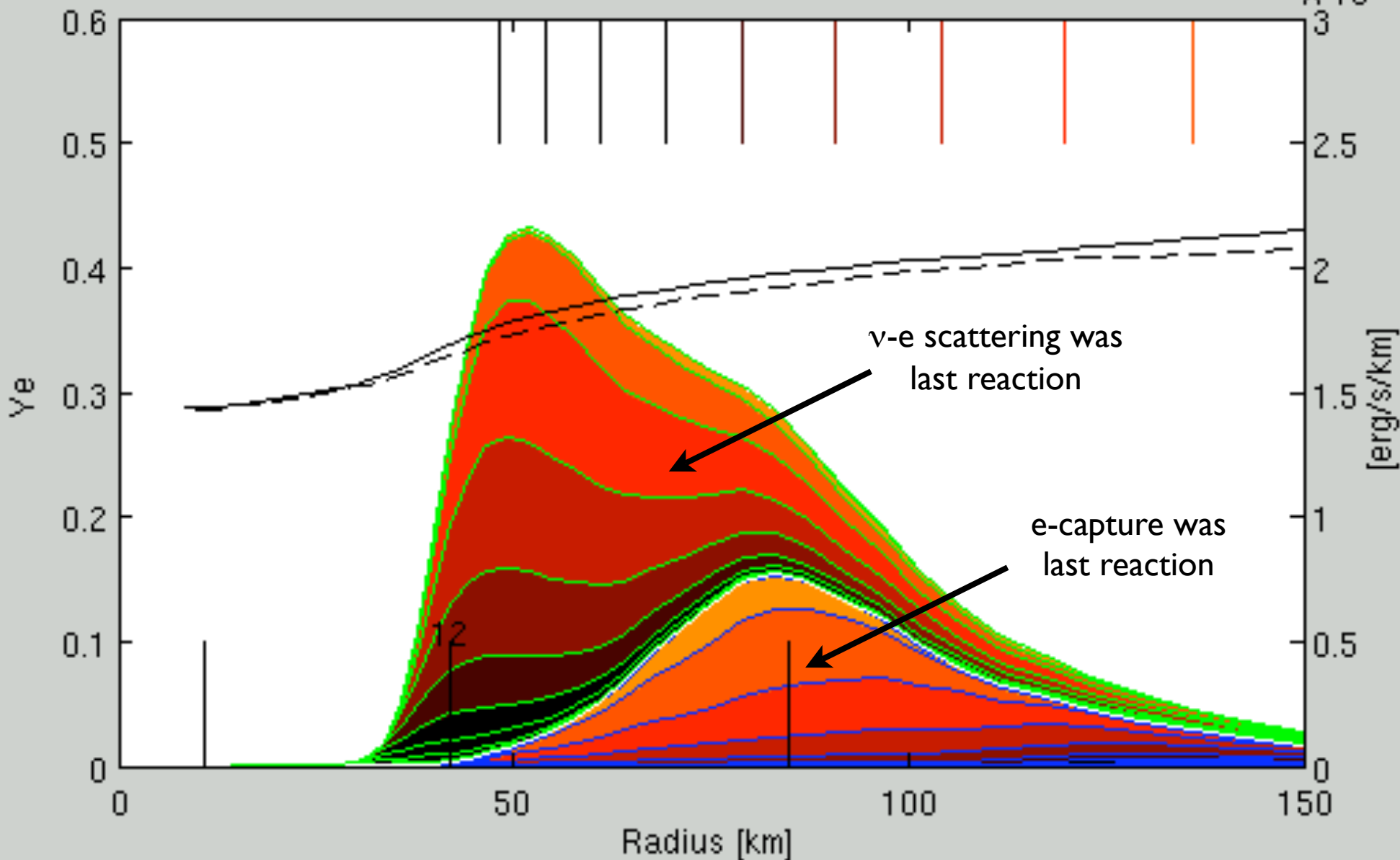
- Long time steps possible
- Follows 'average' evolution
- nonlinear system
- computationally expensive!

Time: 0.255 s

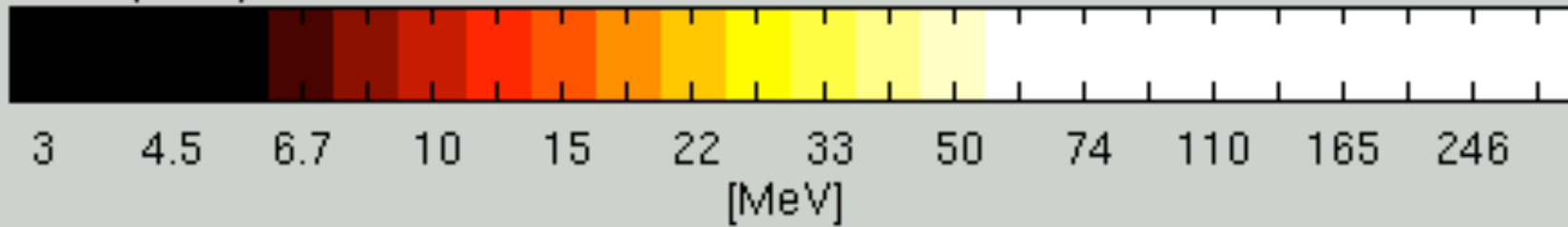


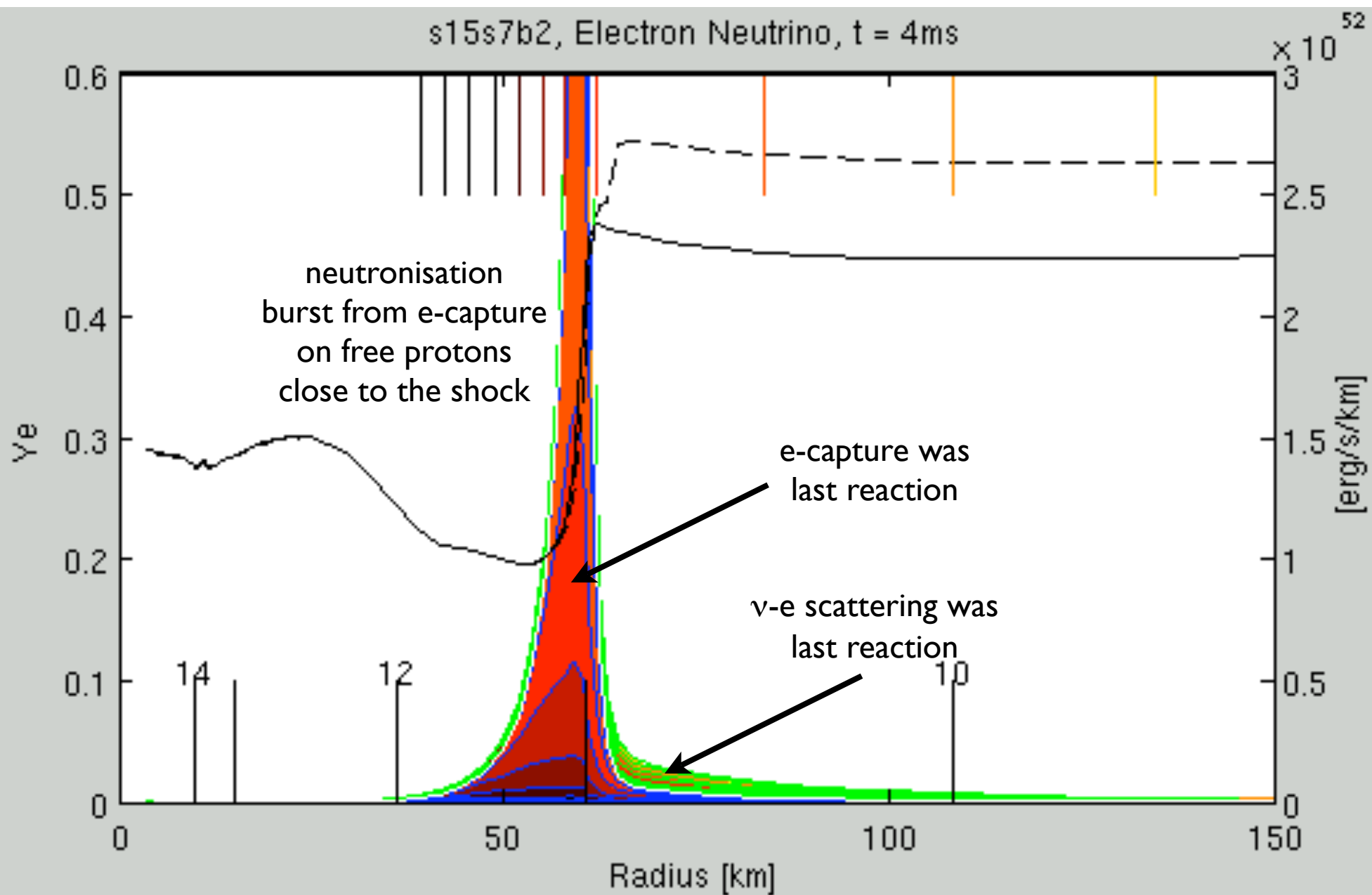
(Liebendörfer et al. 2003)

s15s7b2, Electron Neutrino, t = -2ms

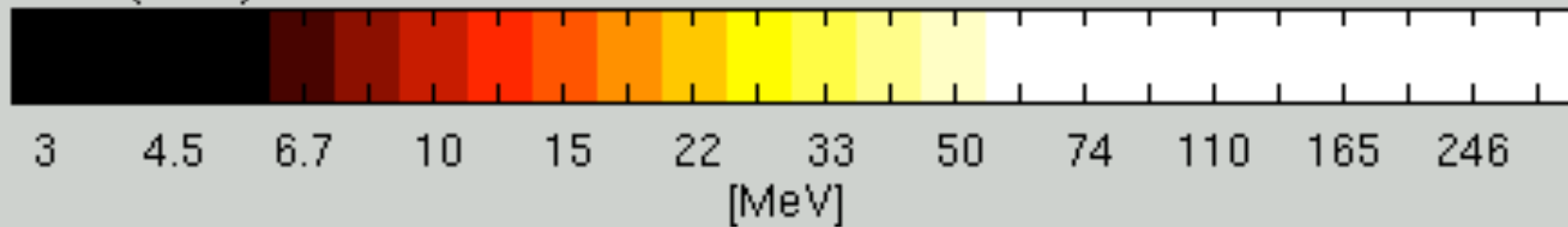


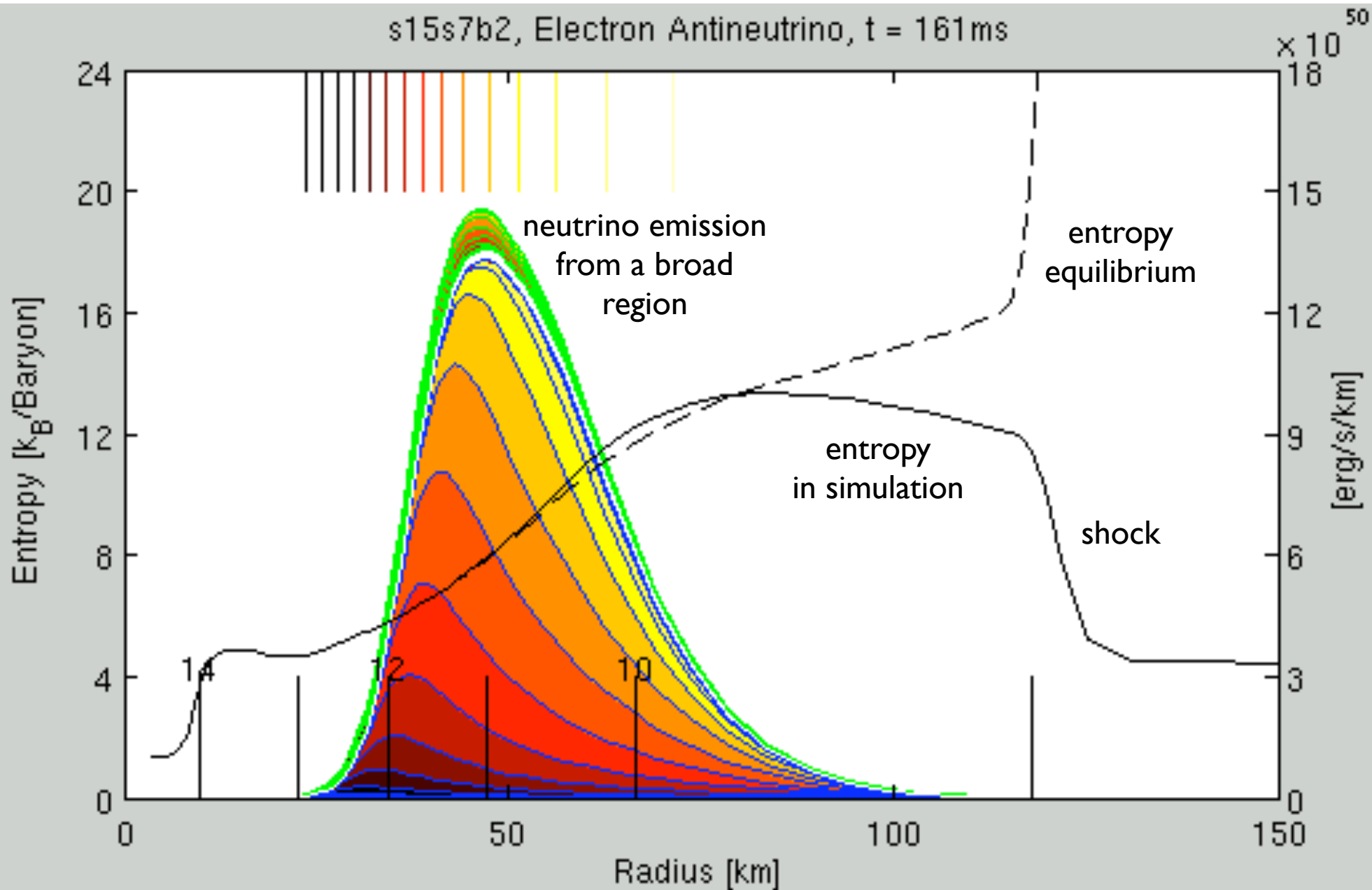
CITA (2004)



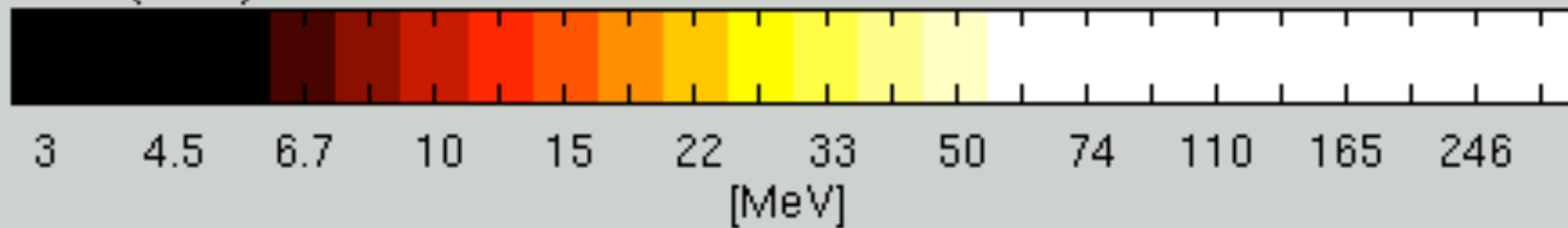


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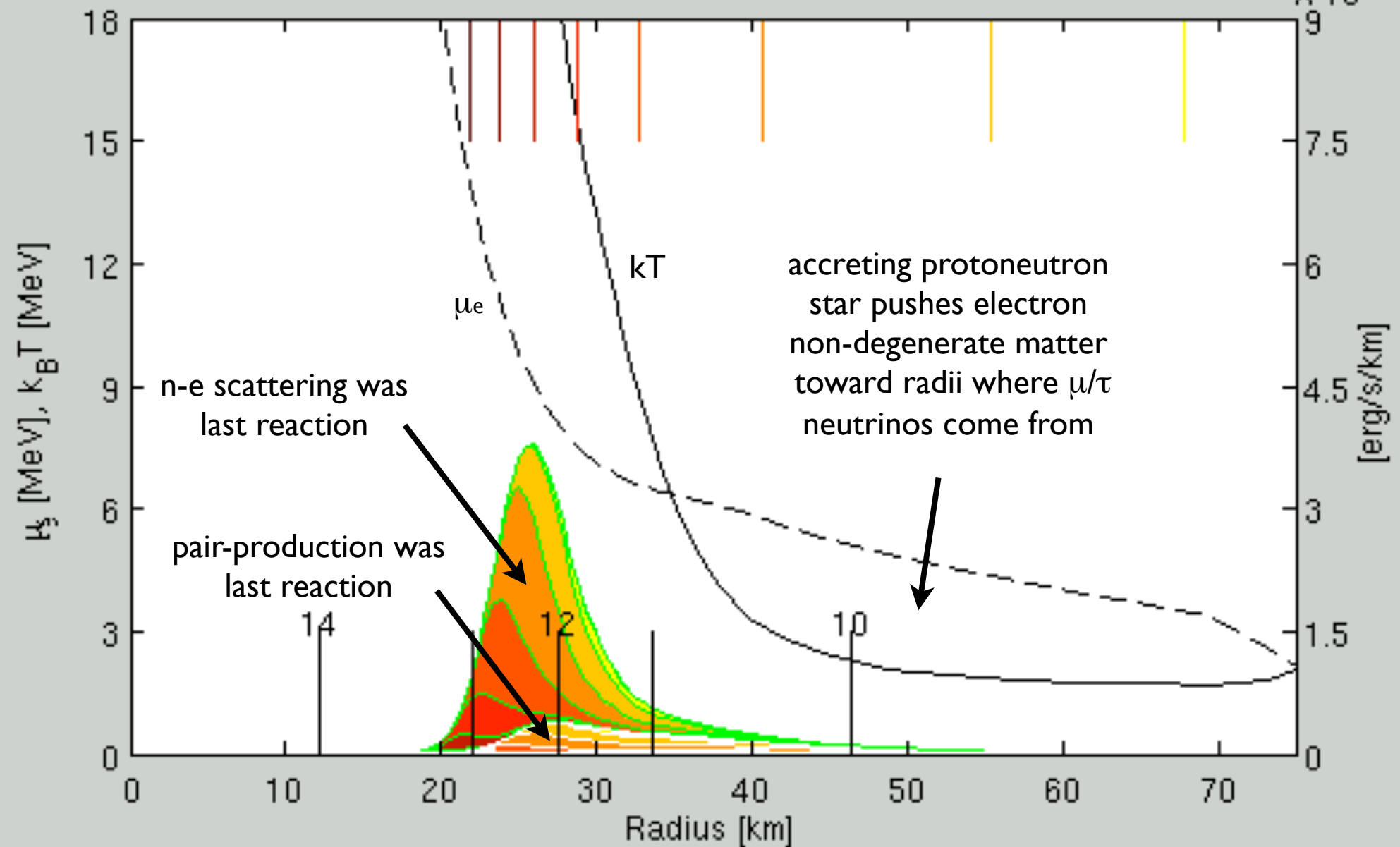




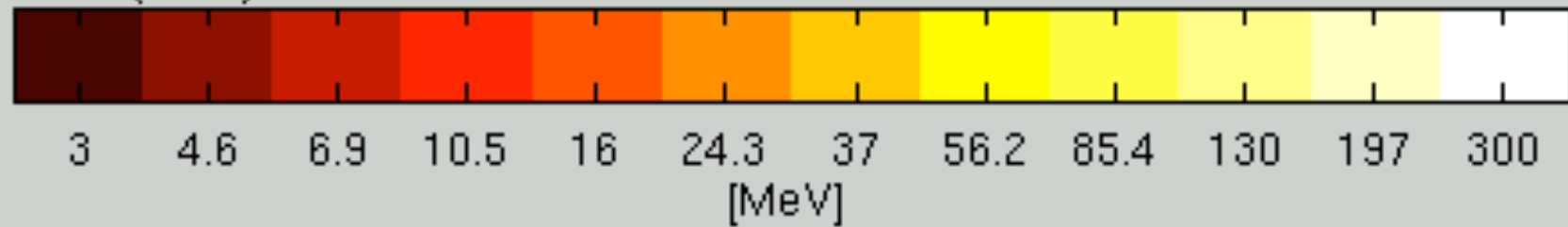
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μ/τ Neutrino, $t = 296\text{ms}$



CITA (2003)



Comparison among independent groups



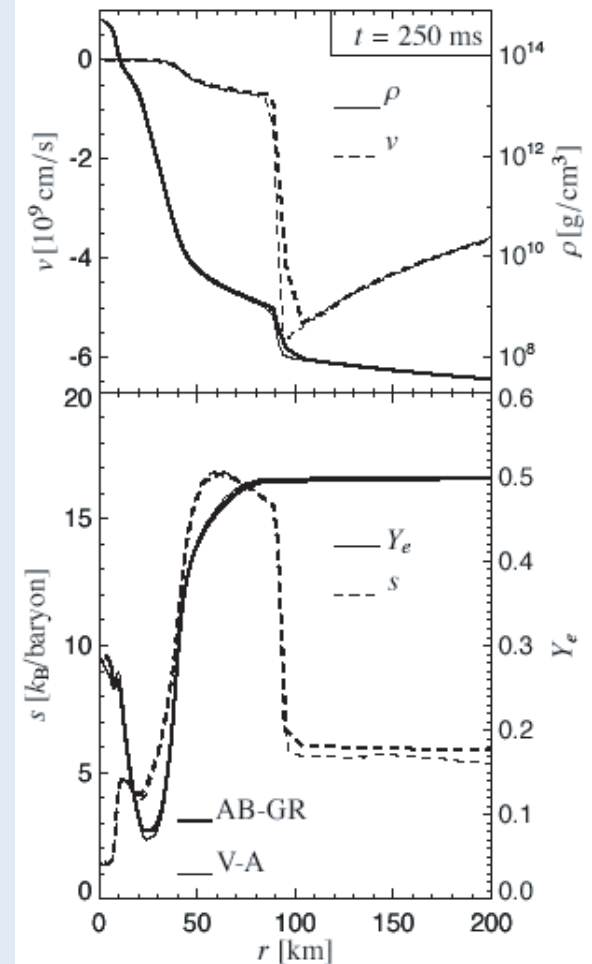
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Comparison of spherically symmetric simulations:
Oak Ridge/Basel group and Garching group

Liebendörfer, Rampp, Janka, Mezzacappa, ApJ 620 (2005)

Explosions only in exceptional cases!

excellent agreement:



(Marek et al., A&A 2006)

Comparison among independent groups

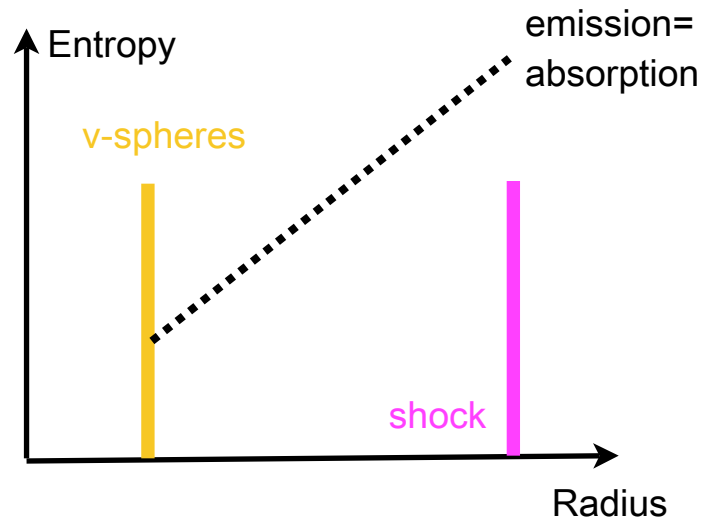
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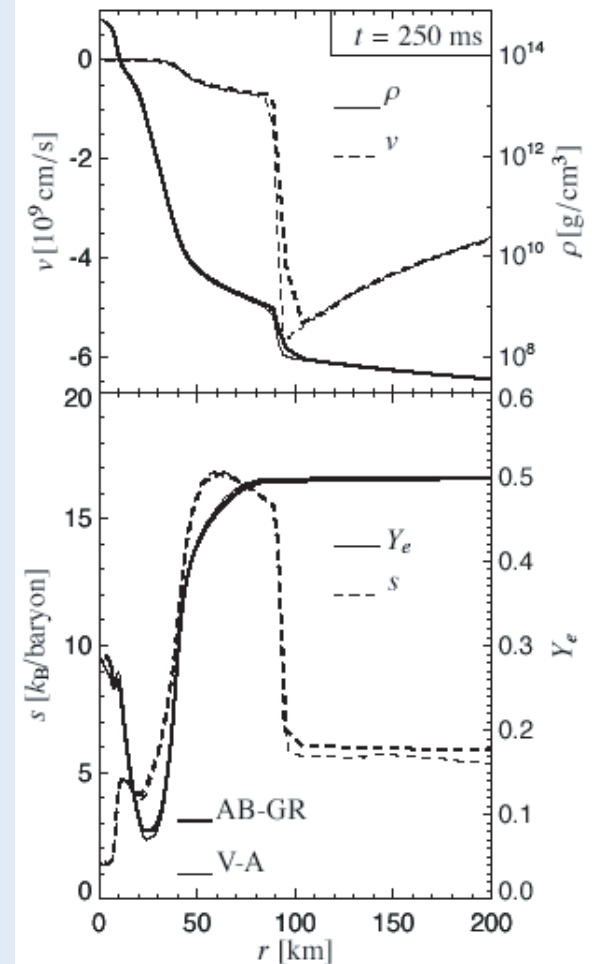
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Problem 1: Exposure
to heating limited by
high infall velocity



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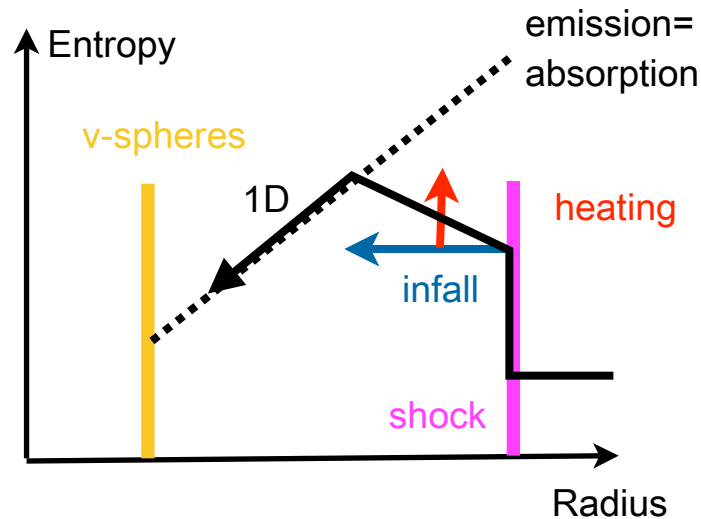
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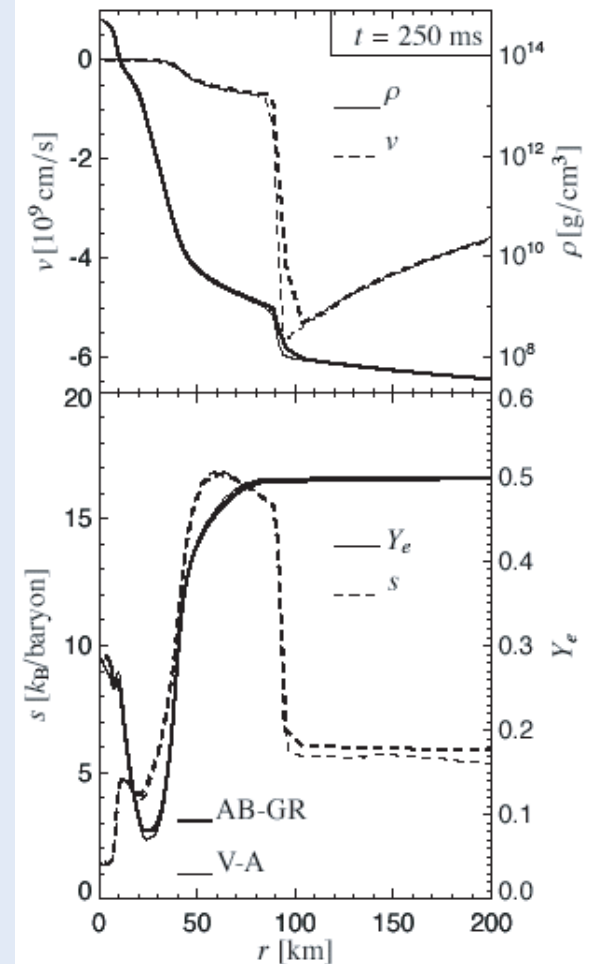
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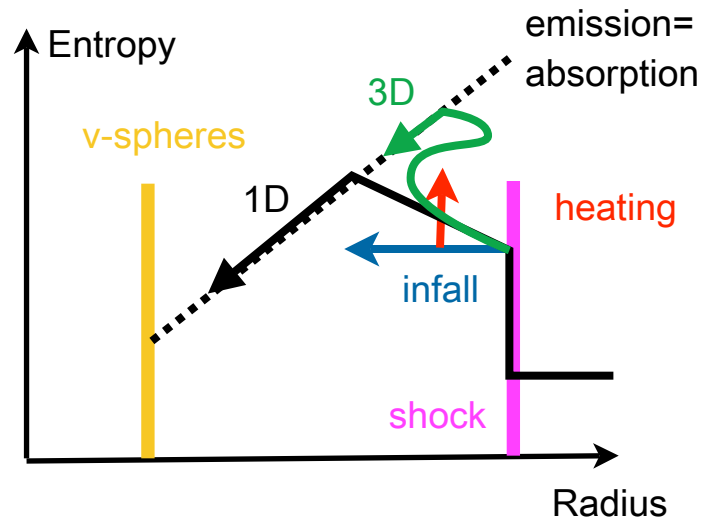
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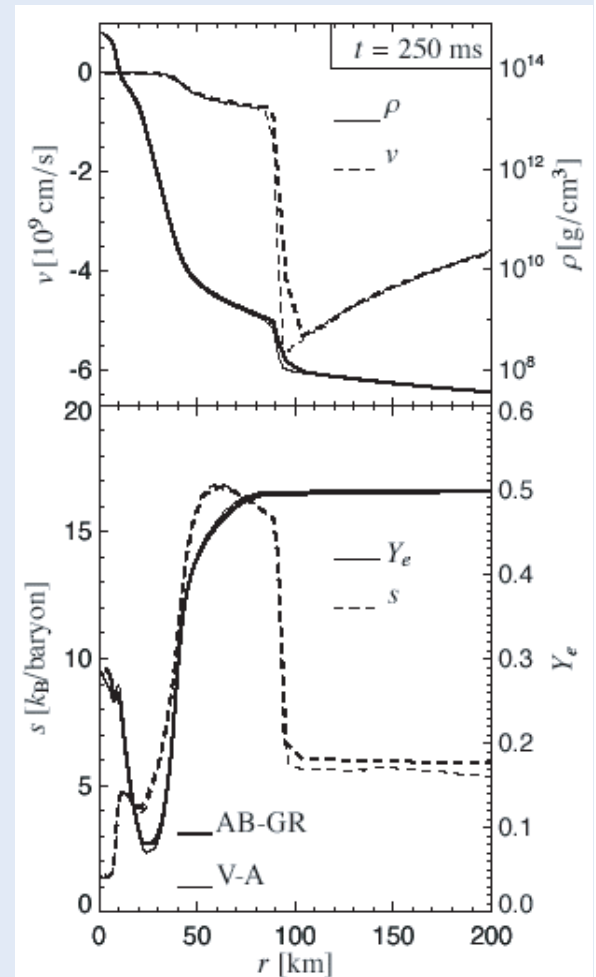
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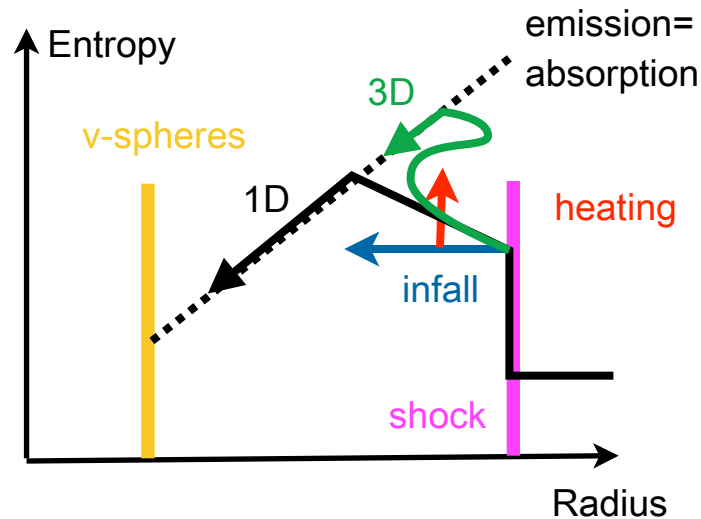
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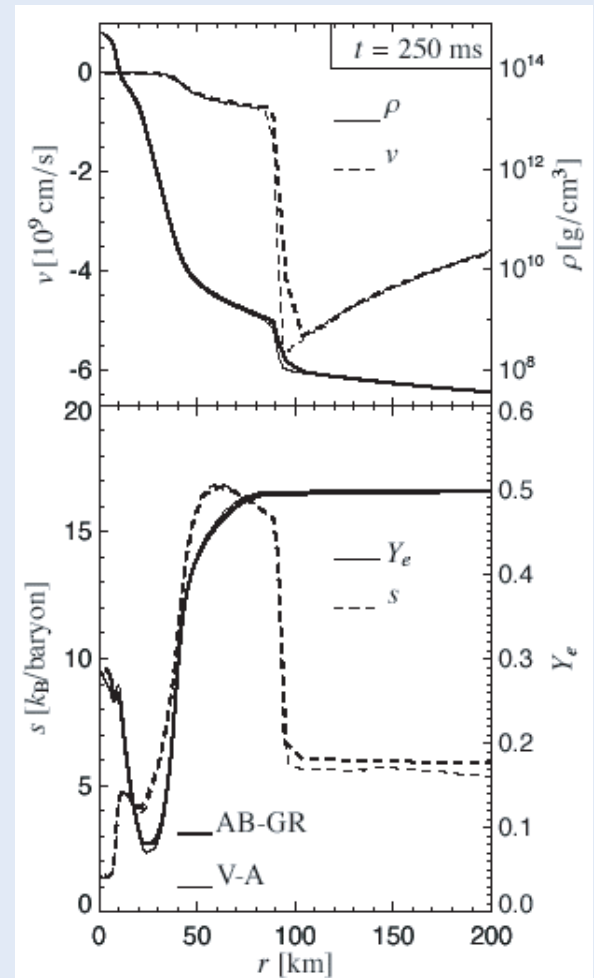
Explosions only in exceptional cases!

Problem 1: Exposure to heating limited by high infall velocity



Problem 2: Accretion (-luminosity) shuts off after onset of explosion

excellent agreement:



(Marek et al., A&A 2006)