

# Lecture: Terrestrial Planet Formation

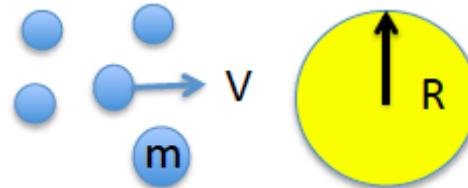
Summer School 2015

Professor Schlichting (MIT)

6<sup>th</sup> August 2015

# Growth Rates

Collision Rate  $= nV\pi R^2$



Growth Rate  $\frac{dM}{dt} = nmV\pi R^2$

Mass Doubling Rate  $\frac{1}{M} \frac{dM}{dt} \sim \frac{\sigma\Omega}{\rho R}$

@ 1 AU  $t_{\text{coll}} = M dt/dM \sim 10^7$  years

@ 20 AU  $t_{\text{coll}} = M dt/dM \sim 10^{11}$  years



Need to enhance collisional Cross-sections

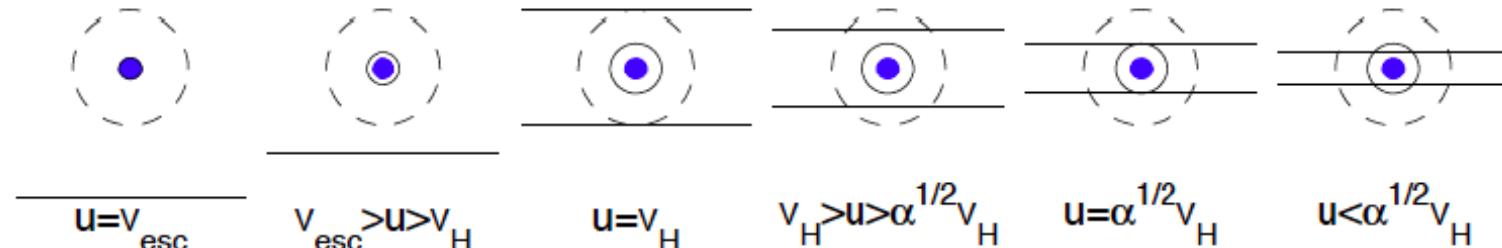


Figure 1: Geometry of disk scale height (*solid horizontal line*), body size (*filled circle*), its Hill sphere (*dashed circle*), and its effective size for accretion (*solid circle*)

# Gravitational Focusing



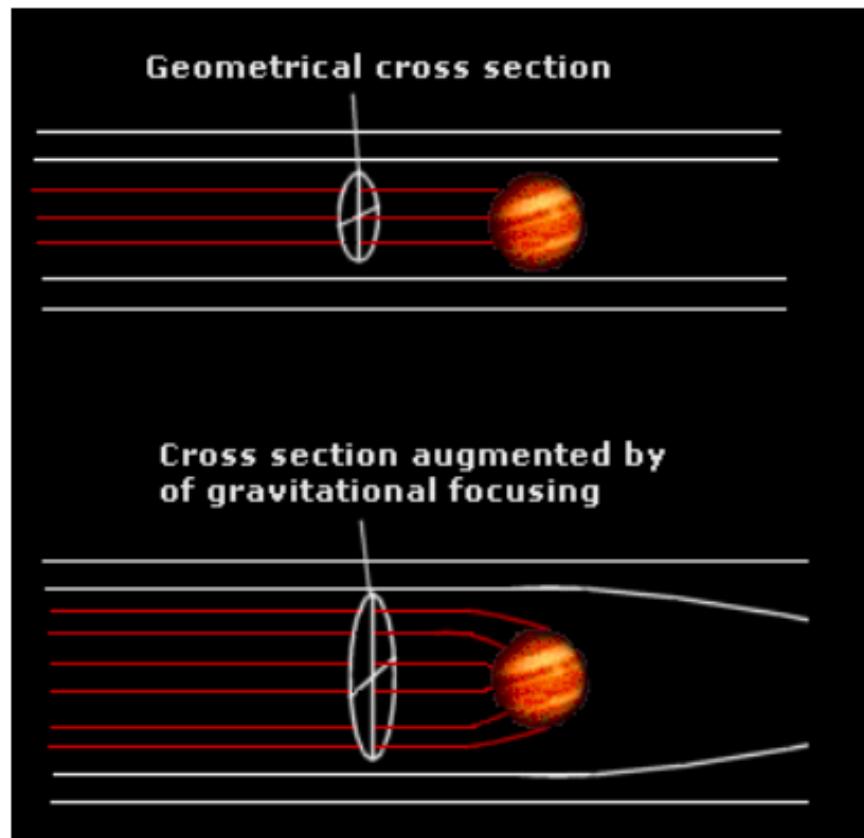
$$\frac{dm_p}{dt} = \rho_s V \pi b^2$$

Without gravitational focusing:

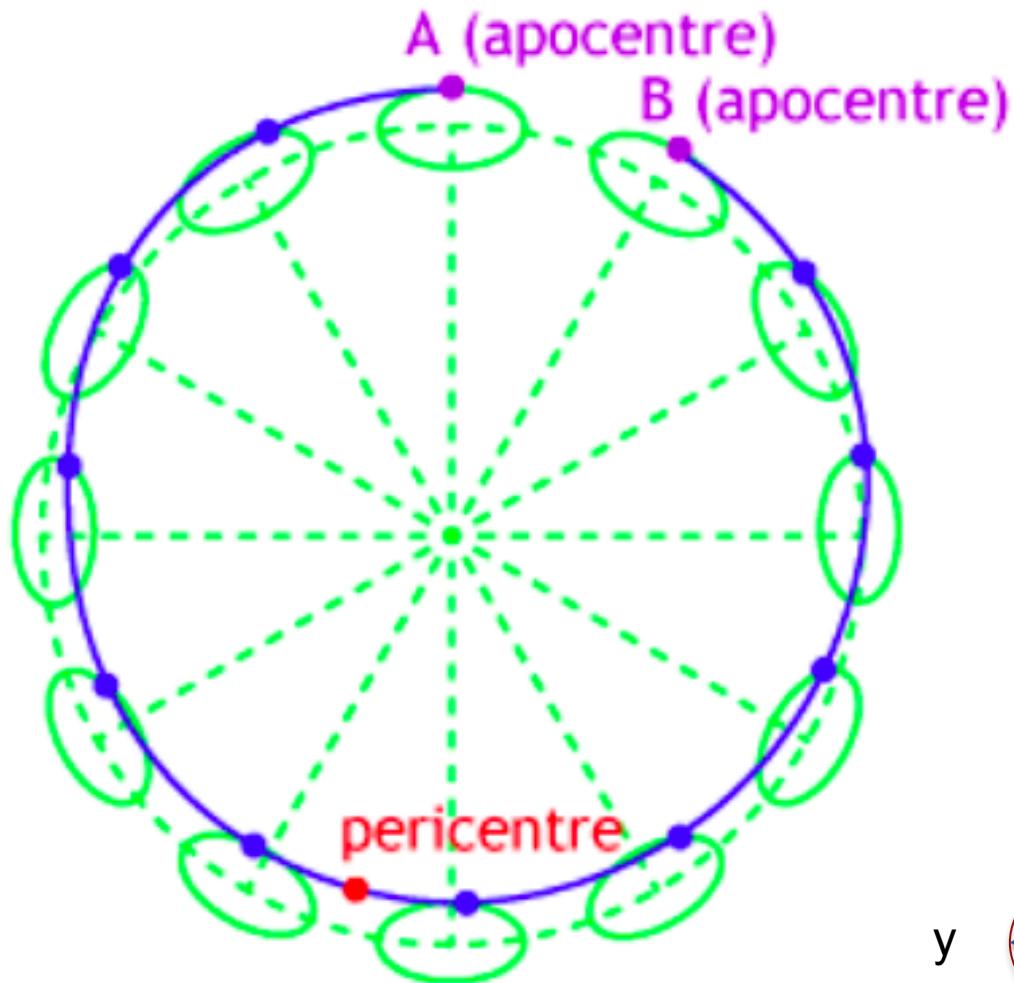
$$\frac{dm_p}{dt} = \rho_s \cdot V \cdot \pi a^2$$

With gravitational focusing:

$$\frac{dm_p}{dt} = \rho_s \cdot V \cdot \pi a^2 \left[ 1 + \frac{v_e^2}{V^2} \right]$$



# Motion in the Rotating Frame



Positions

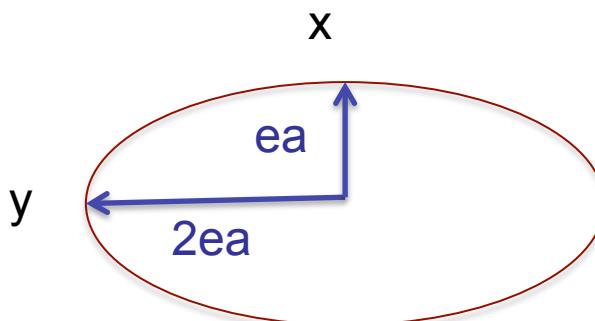
$$x = ae \sin \Omega t$$

$$y = 2ae \cos \Omega t$$

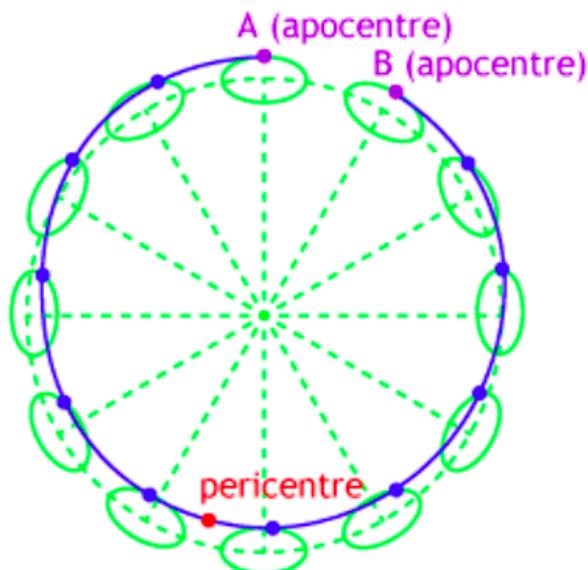
Velocities

$$\dot{x} = ae\Omega \cos \Omega t$$

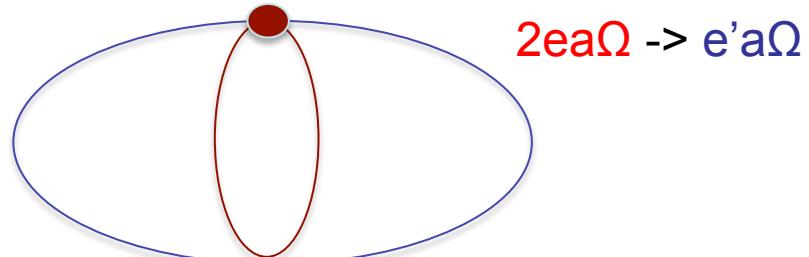
$$\dot{y} = -2ae\Omega \sin \Omega t$$



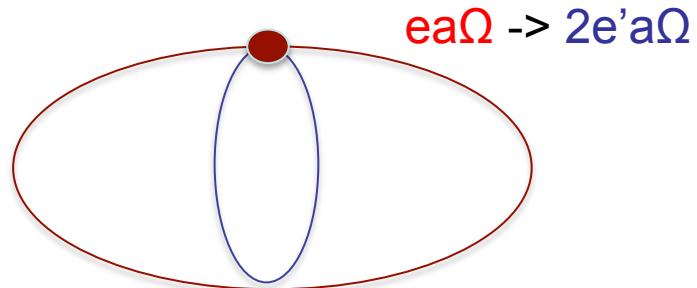
# Viscous Stirring



- Velocity doubles at quadrature



- Velocity halves at periape and apoaope



Safronov (1972) assumed an equal number of collisions in each direction:

→ Average Increase in k.e.  
 $= (4+1/4)/2 = 17/8$

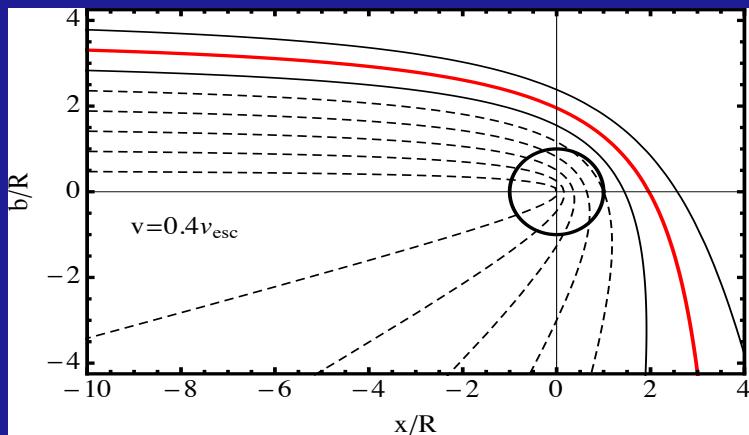
$$\text{cross section} = \pi R^2 \left( \frac{v_{esc}}{v_\infty} \right)^4$$

# Viscous Stirring

Viscous stirring tends to increase the random kinetic energy all all bodies in the disk

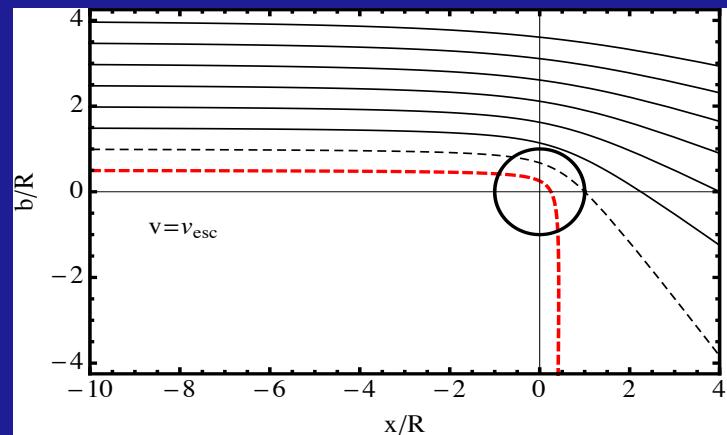
For  $v \ll v_{esc}$

$$\sigma_{stirr} \gg \sigma_{coll}$$



For  $v \gg v_{esc}$

$$\sigma_{stirr} \ll \sigma_{coll}$$

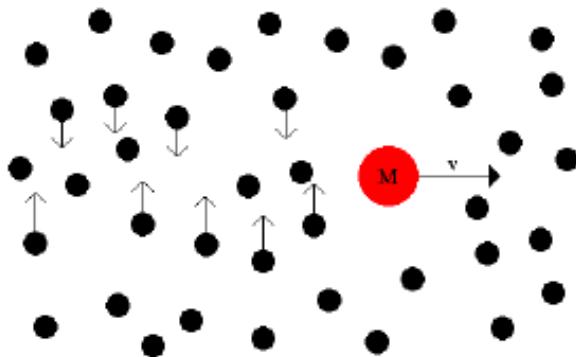


$$\sigma_{stirr} \sim \pi R^2 \left( \frac{v}{v_{esc}} \right)^{-4}$$

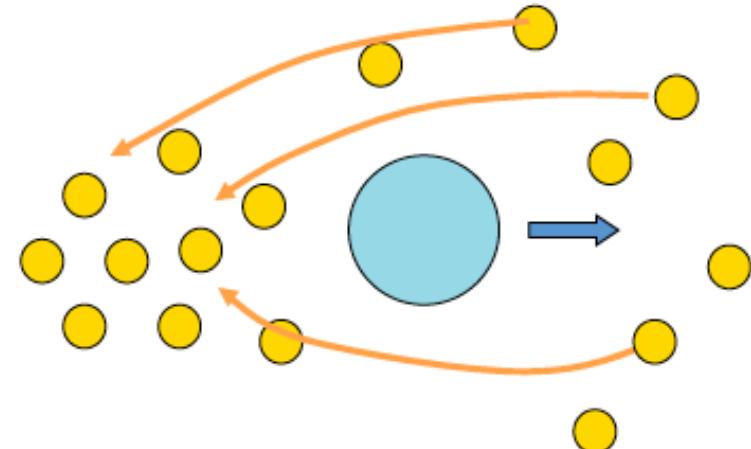
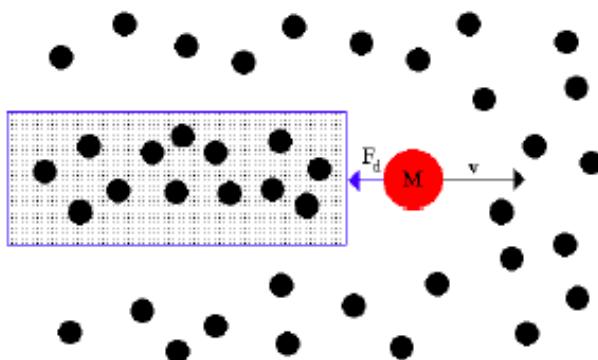
$$\sigma_{coll} \sim \pi R^2 \left( 1 + \left( \frac{v}{v_{esc}} \right)^{-2} \right)$$

# Dynamical Friction

consider a mass,  $M$ , moving through a uniform sea of stars. Stars in the wake are displaced inward.



this results in an enhanced region of density behind the mass, with a drag force,  $F_d$ , known as dynamical friction.



$$F_{df} \sim \left( \frac{GM}{v} \right)^2 \frac{\Sigma}{b} \sim GM\Sigma \quad \text{2D}$$

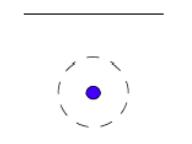
# Growth Rates & Stirring Rates

$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R}$$

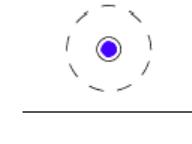
$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R} \left( \frac{u}{v_{esc}} \right)^{-2}$$

$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R} \alpha^{-1/2} \frac{v_{esc}}{u}$$

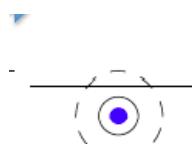
$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R} \alpha^{-3/2}$$



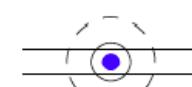
$$u = v_{esc}$$



$$v_{esc} > u > v_H$$



$$v_H > u > \alpha^{1/2} v_H$$



$$u < \alpha^{1/2} v_H$$

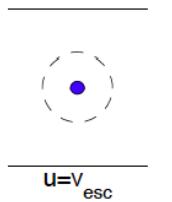
$$\frac{1}{u} \frac{du}{dt} \sim \frac{\Sigma\Omega}{\rho R} \left( \frac{u}{v_{esc}} \right)^{-4} \quad u > v_H$$

$$\frac{1}{u} \frac{du}{dt} \sim \frac{\Sigma\Omega}{\rho R} \alpha^{-3/2} \left( \frac{u}{v_{esc}} \right)^{-1} \quad u < v_H$$

# Growth Regimes

$$\frac{1}{R} \frac{dR}{dt} \propto R^{-1}$$

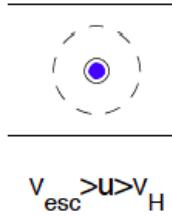
Orderly Growth



$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R}$$

$$\frac{1}{R} \frac{dR}{dt} \propto R$$

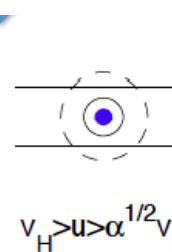
Runaway Growth



$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R} \left( \frac{u}{v_{esc}} \right)^{-2}$$

$$\frac{1}{R} \frac{dR}{dt} \propto 1$$

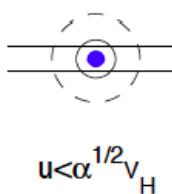
Neutral Growth



$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R} \alpha^{-1/2} \frac{v_{esc}}{u}$$

$$\frac{1}{R} \frac{dR}{dt} \propto R^{-1}$$

Orderly Growth



$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma\Omega}{\rho R} \alpha^{-3/2}$$

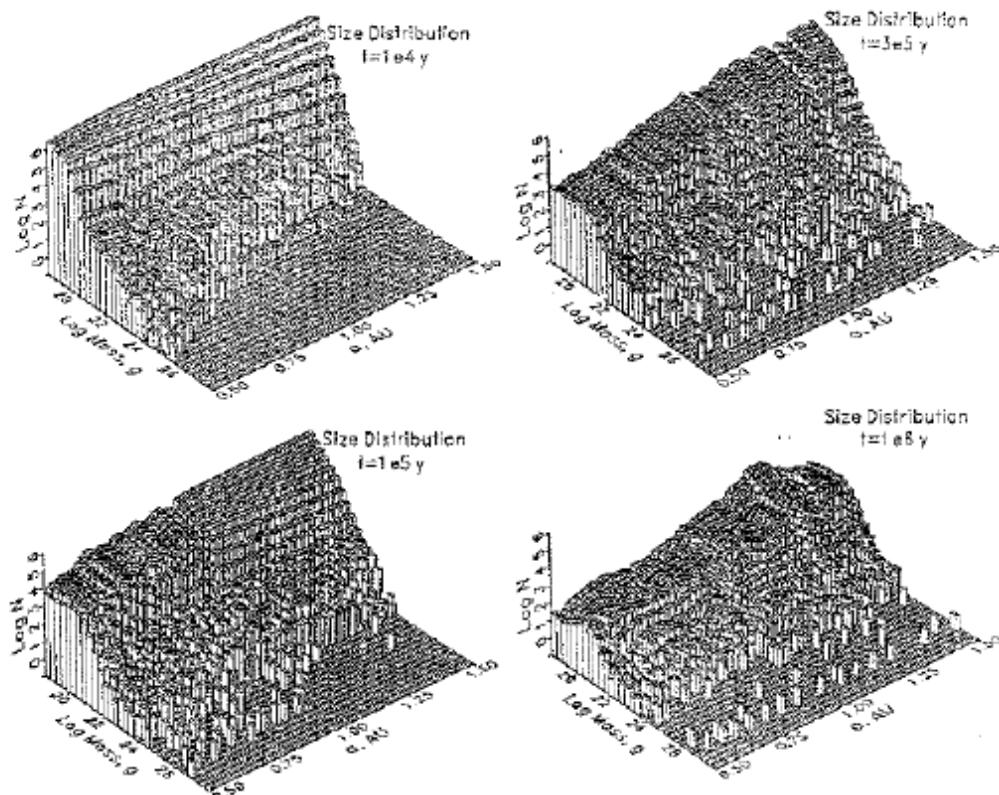
# Runaway Growth

$$\frac{1}{R_1} \frac{dR_1}{dt} \propto R_1$$

For  $v_{\text{esc}} > u$  gravitational focusing enhances the accretion rate

$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma \Omega}{\rho R} \left( \frac{v_{\text{esc}}}{u} \right)^2 \propto R$$

Small number of large bodies run-away from the rest of the planetesimal population



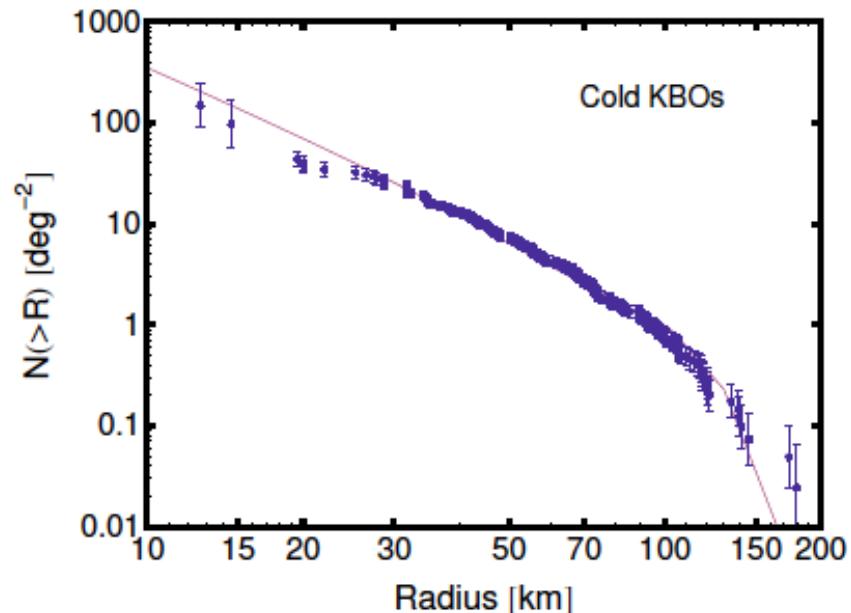
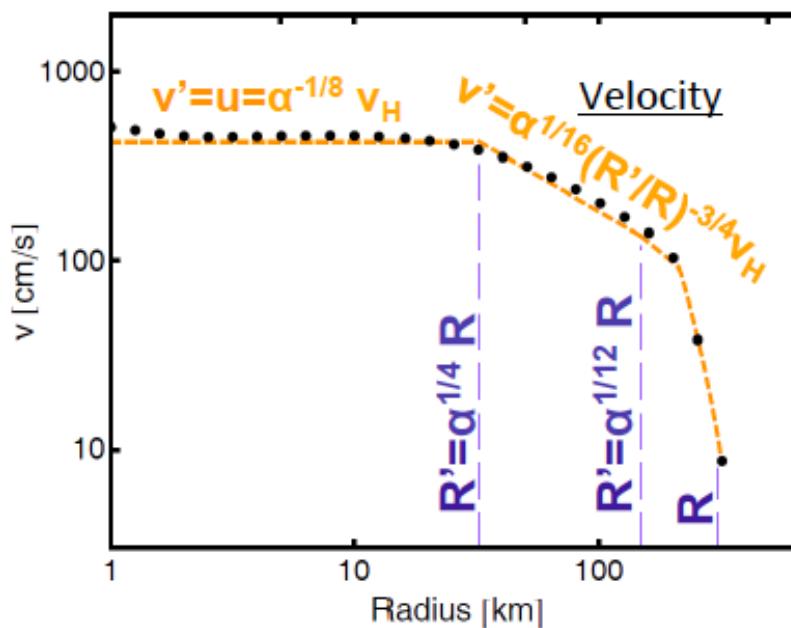
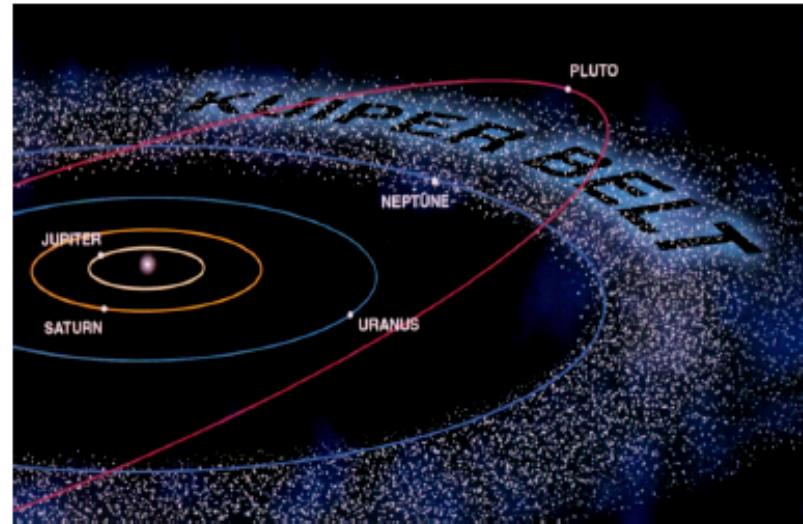
# Runaway Growth in Solar System

$$\alpha \sim R_{Sun} / a$$

$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma \Omega}{\rho R} \left( \frac{u}{v_{esc}} \right)^{-2} + \frac{\Sigma \Omega}{\rho R} \alpha^{-3/2}$$

Accretion of small bodies

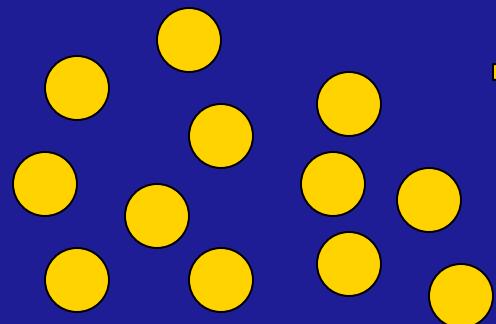
Accretion of big bodies



(Schlichting & Sari, 2011)

# Two-Group-Approximation

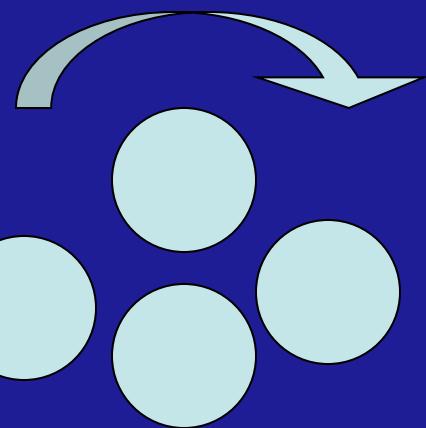
Small bodies



damp  $v$ , accrete onto big bodies

Big bodies

excite  $v$ , accrete themselves



Mass surface density =  $\sigma$

mass =  $m$ , radius =  $r$

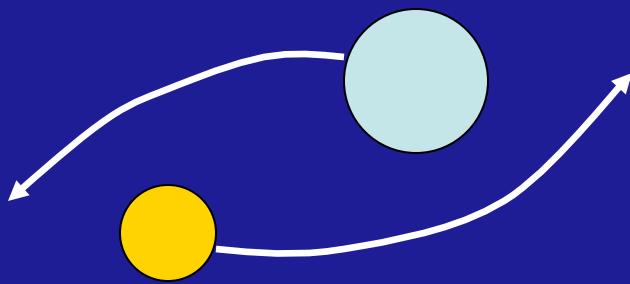
velocity dispersion =  $u$

Mass surface density =  $\Sigma$

Mass =  $M$ , Radius =  $R$

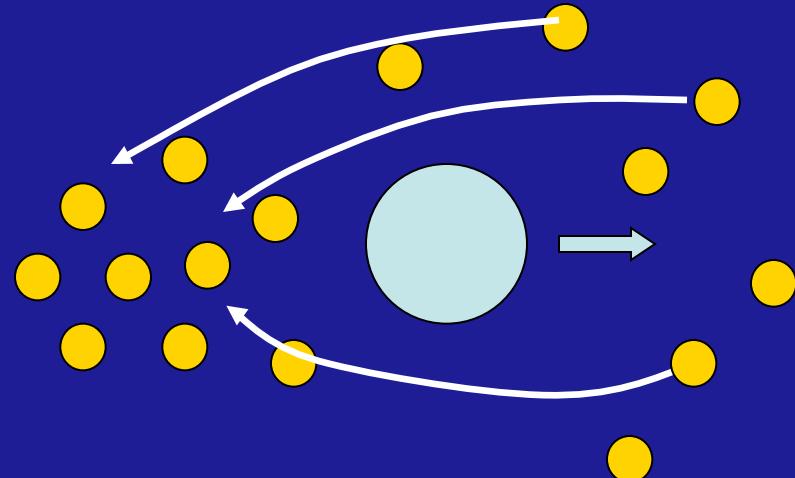
velocity dispersion =  $v$

# Velocity Evolution

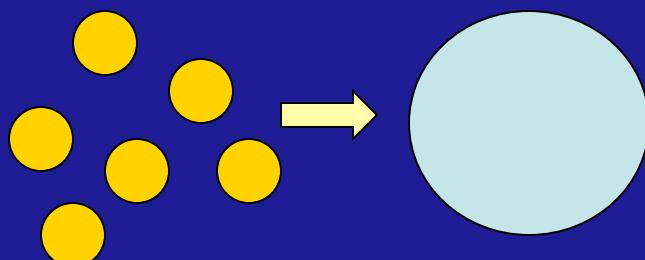
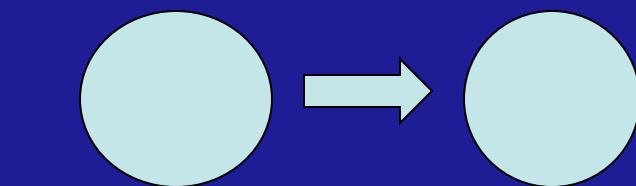


## Growth

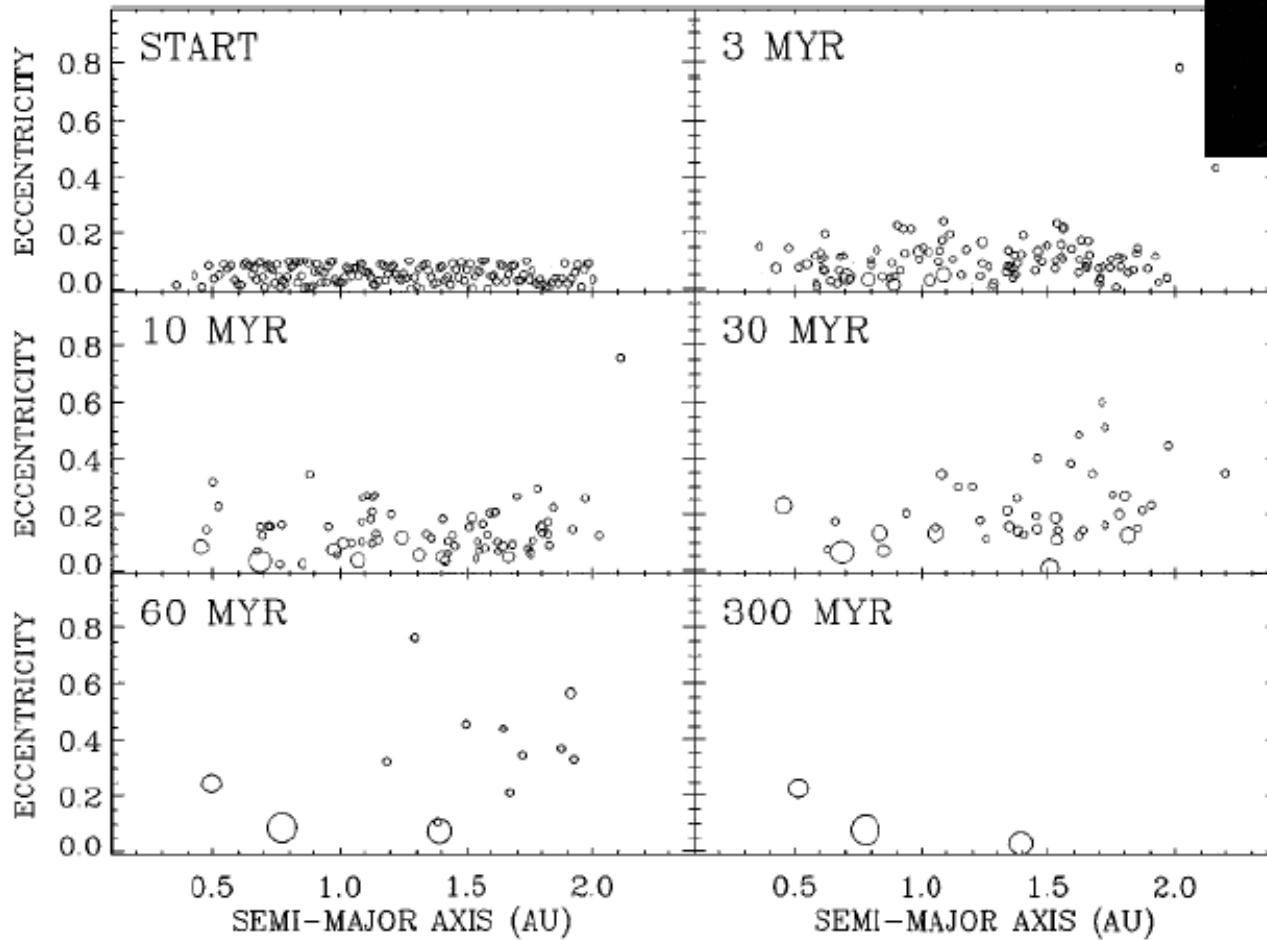
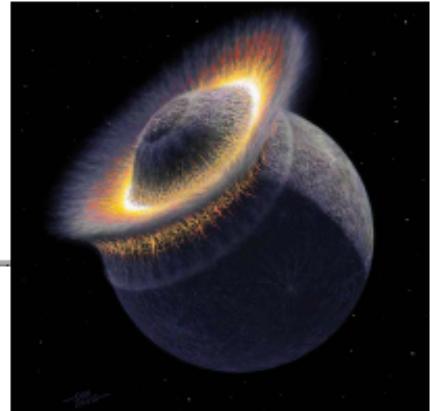
Viscous Stirring



Dynamical Friction



# Giant Impacts



Chambers 2002

# Oligarchic Growth

- Runaway growth ends when
  - a) velocities of small bodies become sub-Hill
  - b) velocities of small bodies become super-escape
  - c) when Oligarchic growth begins

When Oligarchic growth ends

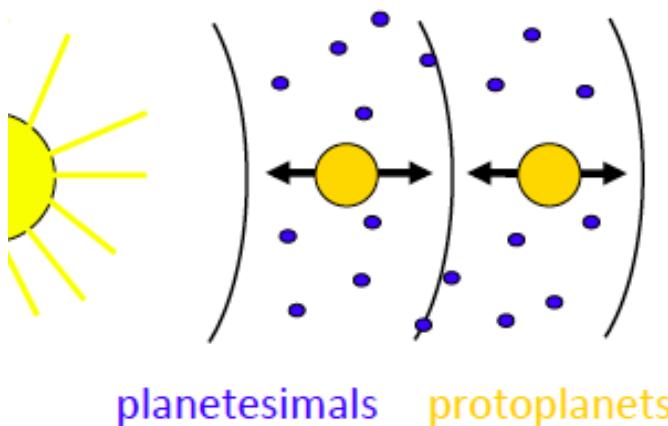
$$\Sigma \sim \sigma$$

Oligarchic growth maybe the final stage of planet formation  
in the outer Solar system

# Isolation Masses

$$r_{\text{H}} = r \left( \frac{m_p}{3M_*} \right)^{1/3}$$

$$\begin{aligned} m_p &= 2\pi r \cdot 2Br_{\text{H}} \cdot \Sigma_s = 4\pi Br^2 \Sigma_s \left( \frac{m_p}{3M_*} \right)^{1/3} \\ &= \frac{(4\pi Br^2 \Sigma_s)^{3/2}}{(3M_*)^{1/2}}. \end{aligned}$$

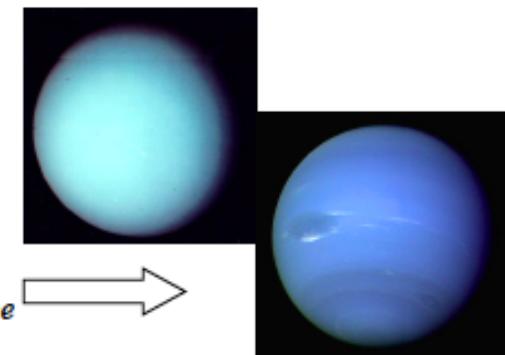


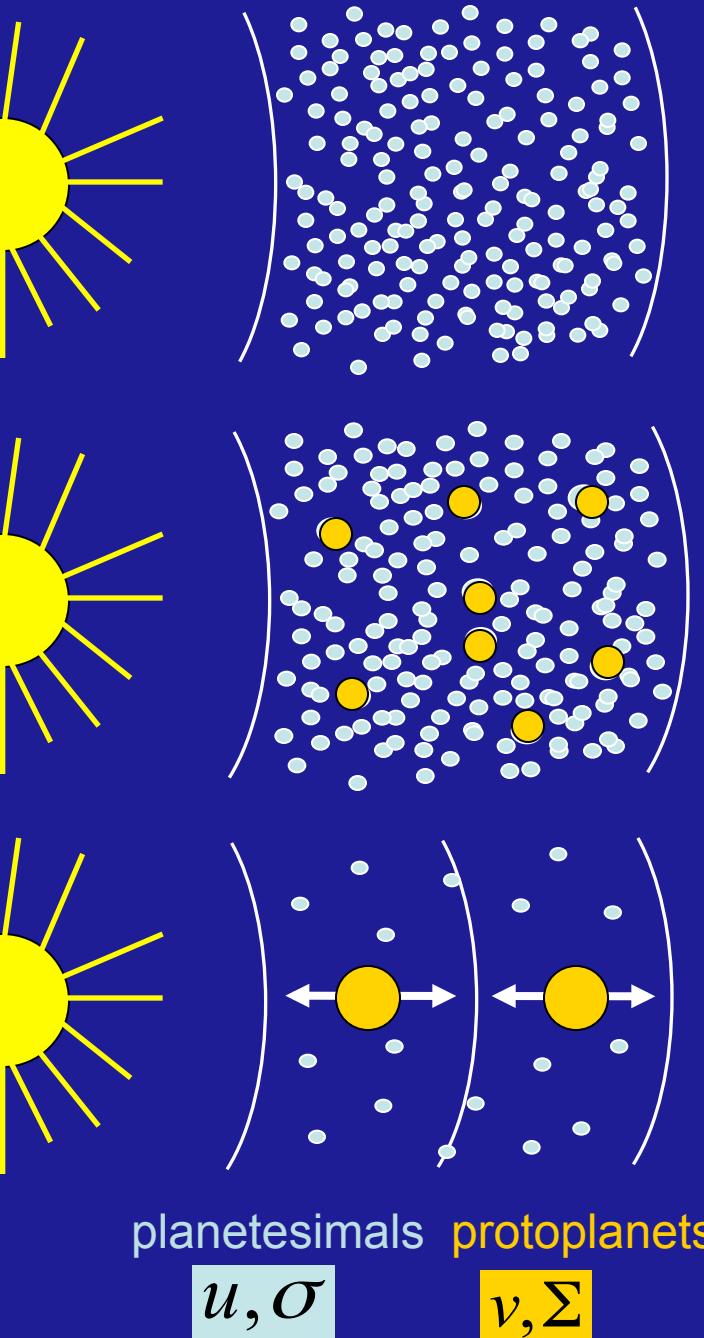
@ 20 AU

$$M_{\text{iso}} \approx 2\pi a (\Delta a_{\text{zone}}) \Sigma \sim M_{\text{Neptune}} \longrightarrow$$

@ 1 AU

$$M_{\text{iso}} \approx 2\pi a (\Delta a_{\text{zone}}) \Sigma \sim 0.01 M_{\text{Earth}}$$





### 1. Planetesimal formation:

### 2. Runaway growth:

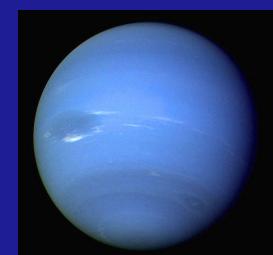
$$\frac{1}{R_1} \frac{dR_1}{dt} \propto R_1$$

For  $v_{\text{esc}} > u$  gravitational focusing enhances the accretion rate

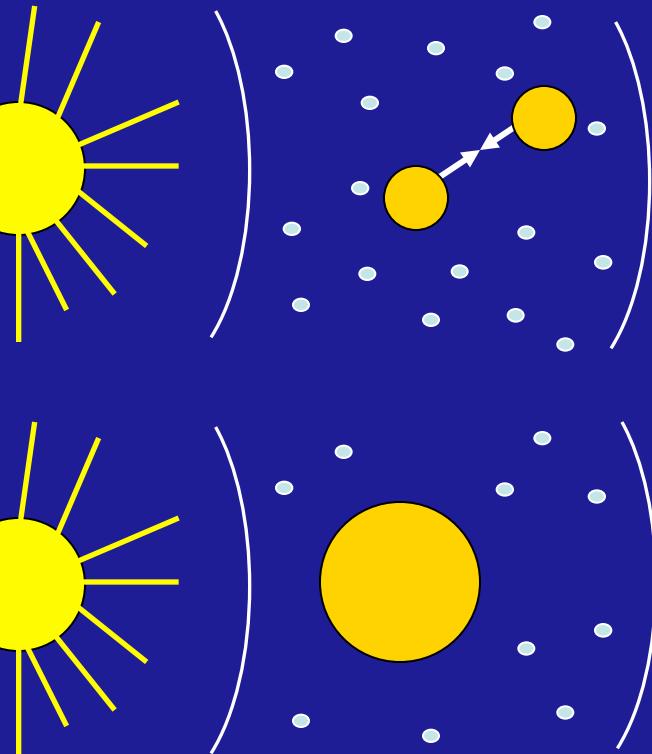
$$\frac{1}{R} \frac{dR}{dt} \sim \frac{\sigma \Omega}{\rho R} \left( \frac{v_{\text{esc}}}{u} \right)^2 \longrightarrow t_{\text{grow}} \sim 10^7 \text{ years}$$

### 3. Oligarchic growth & Isolation:

$$M_{\text{iso}} \approx 2\pi a (\Delta a_{\text{zone}}) \Sigma \sim M_{\text{Neptune}}$$



# Last Stages of Terrestrial Planet Formation



## Giant Impacts:

Protoplanets' velocity dispersion increases

→ Giant Impacts

$$t_{\text{Giant-Impacts}} \sim 10^8 \text{ years (1AU)}$$

## Clean up:

- Orbits planar & circular
- Accretion & ejection of remaining planetesimals