

SOLVABLE MATTER ON 2D CAUSAL DYNAMICAL TRIANGULATION?

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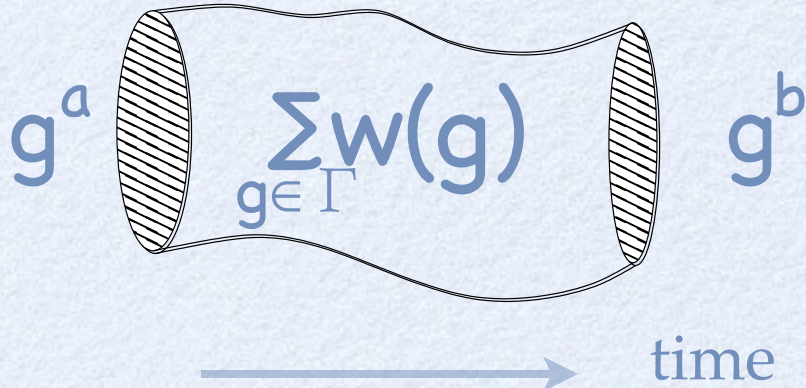
1. QUANTUM GRAVITY

Gravity's dynamical degree of freedom is the metric $g_{\mu\nu}(x,t)$

Classically $g_{\mu\nu}(x,t)$ obeys Einstein's equations:

$$g_{\mu\nu}(x,0) \longrightarrow g_{\mu\nu}(x,t)$$

Quantum mechanics is different:

$$\langle g^b(x), t=T \mid g^a(x), t=0 \rangle \sim \int_{g \in \Gamma} \Sigma w(g)$$


The diagram illustrates the evolution of the metric tensor g over time. It features a cylinder with a wavy surface, representing the space of metrics. The left circular face is labeled g^a and the right circular face is labeled g^b . A horizontal arrow at the bottom points to the right and is labeled "time". A vertical arrow on the left points upwards and is labeled "space". The cylinder's surface is labeled with the integral $\int_{g \in \Gamma} \Sigma w(g)$.

Probability amplitude for evolution from g^a to g^b

How are Γ and w defined ?

Several approaches; some non-stringy ones

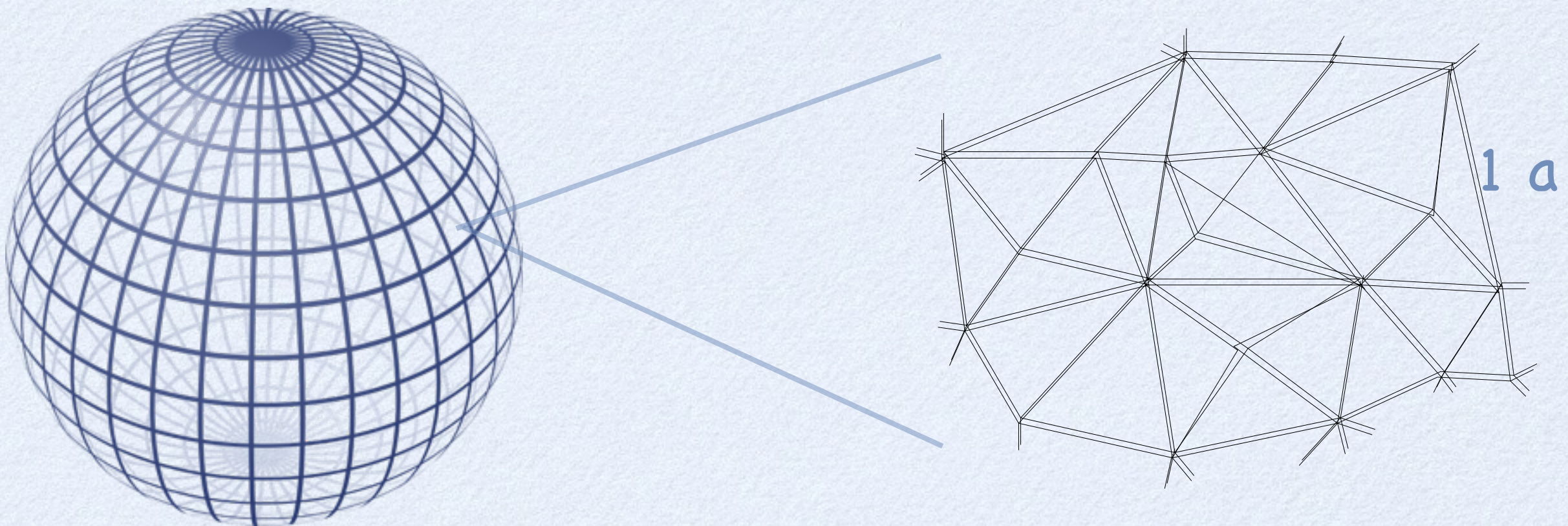
1. Continuum field theoretical -- exact RG looking for non-Gaussian fixed points where QG non-perturbatively renormalizable (asymptotic safety)

2. Discretized -- Causal Dynamical Triangulations Γ is a set of graphs, $w(g)$ defined in terms of graph quantities, we look for a critical point (or line) where a continuum limit can be taken.

Recently much discussed also in context of Horava/Lifshitz gravity

2. TRIANGULATIONS & METRICS

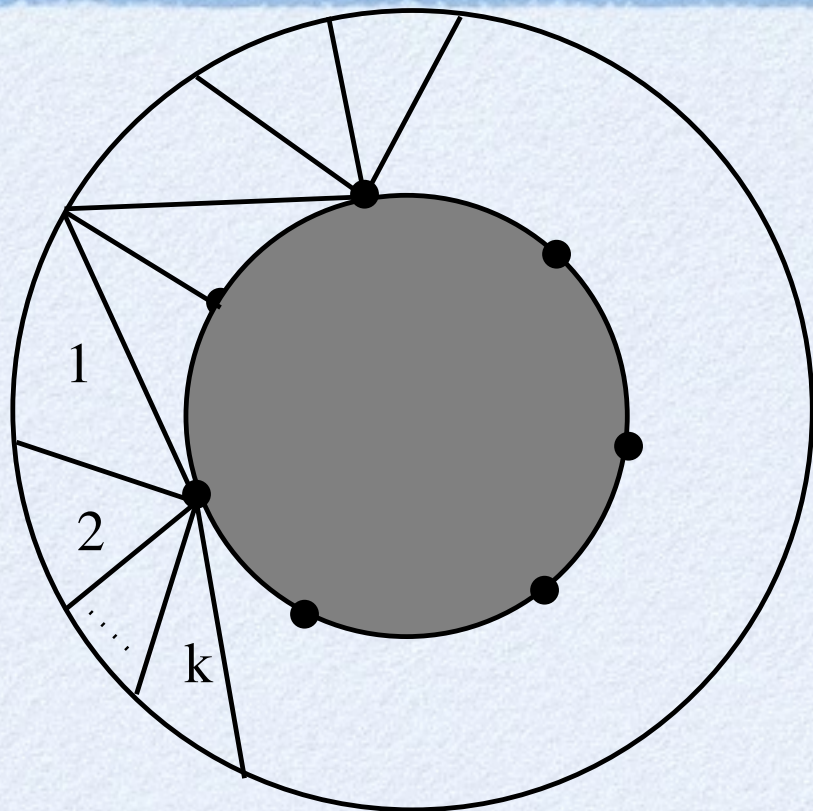
- So how do we do Σ_g ?
- Note $g(x,t) \leftrightarrow \{\text{geodesic distances}\} \leftrightarrow \{\text{graph distances}\}$
- Think about 2-dim space with spherical topology



$$\Sigma_g \leftrightarrow \Sigma_{\text{triangulations}} , a \rightarrow 0$$

3. CAUSAL TRIANGULATION & TREE BIJECTION

arXiv:0908.3643



$$w_G = \prod_{v \in G} g^{k_v + 1}$$

$$Z(g) = \sum_G w_G$$

Critical at $g = g_c = 1/2$

at g_c offspring probability

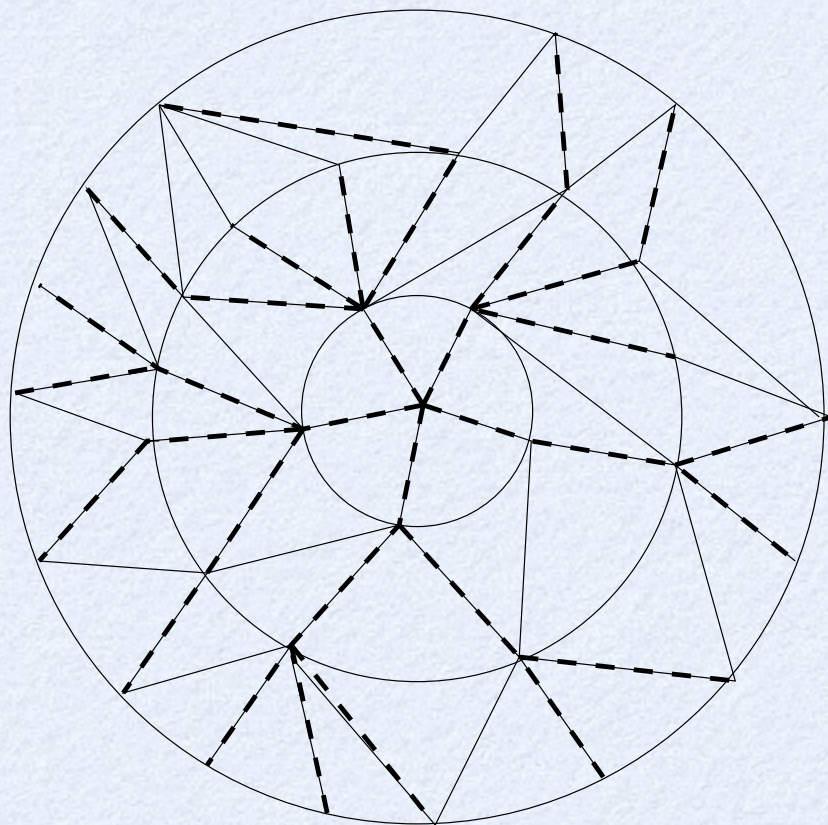
$$p_n = (1/2)^{n+1}$$

so critical Galton Watson

$$\langle n \rangle = 1$$

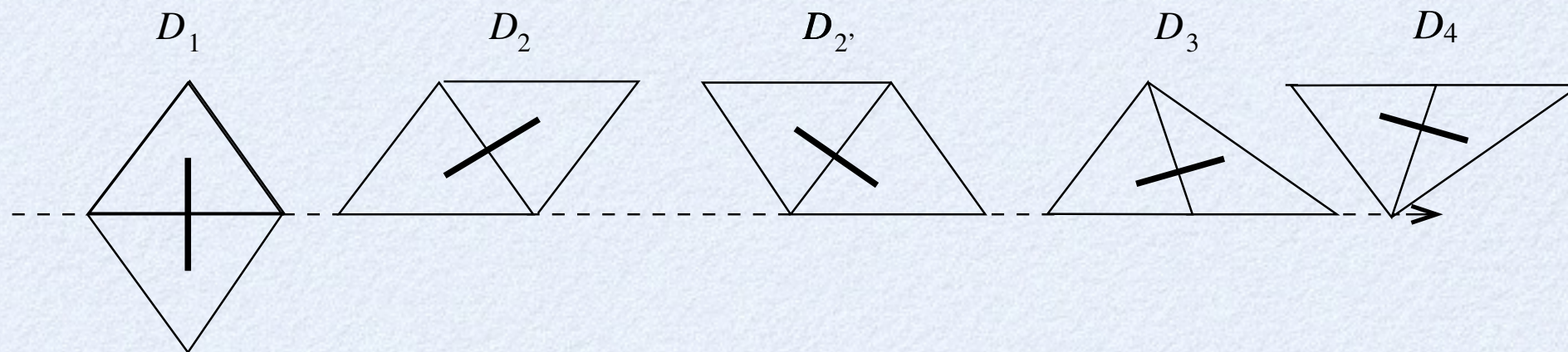
$$\mu(\infty \text{ CDT}) \Leftrightarrow \mu(\text{URT})$$

Uniform RT is a particular GRT



3. DIMERS ON CDT arXiv:1201.4322 Atkin & Zohren

Introduce dimers - possible configurations



Two dimers may not share a triangle

$$w_G = \prod_{v \in G} g^{k_v + 1}$$

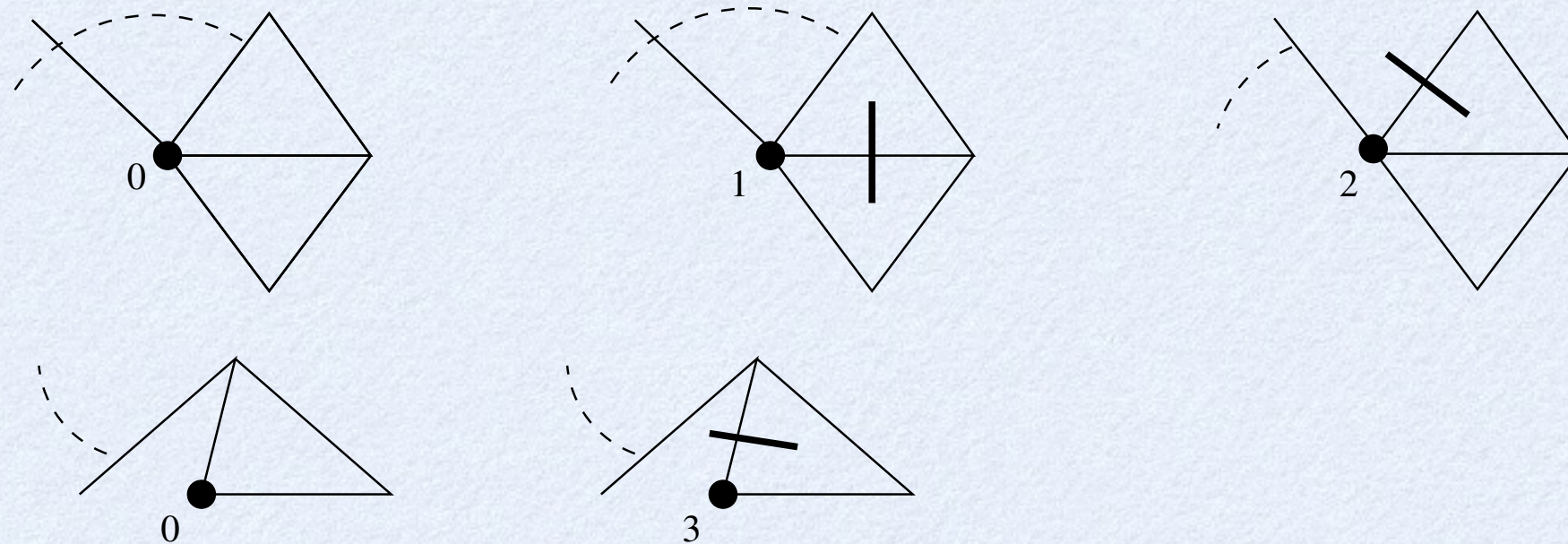
$$Z(\{\xi\}; g) = \sum_{G, D} w_G \xi_1^{D_1} \xi_2^{D_2} \xi_{2'}^{D_{2'}} \xi_3^{D_3} \xi_4^{D_4}$$

Z fn of $\xi_2 + \xi_{2'}$ so set $\xi_{2'} = 0$

A & Z: $\xi_3 = \xi_4 = 0$ allows bijection to labelled trees

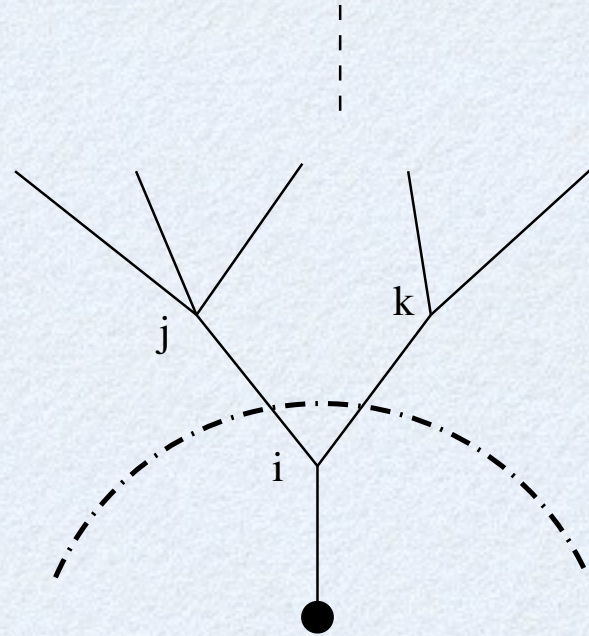
4. LABELLED TREES

A, D & JW: $\xi_4 = 0$ allows bijection to labelled trees



Label allocation \mathbf{L} survives bijection —
trees $\boldsymbol{\tau}$ inherit a set of labelling rules from dimer rule

τ can be decomposed



$$W_i(\{\xi\}; g) = \sum_{\tau, L; i} g^\tau \xi_1^{L_1} \xi_2^{L_2} \xi_3^{L_3}$$

$$= F_i(\{W\}; \{\xi\}; g)$$

these are geometric series so F are rational fns of W
and

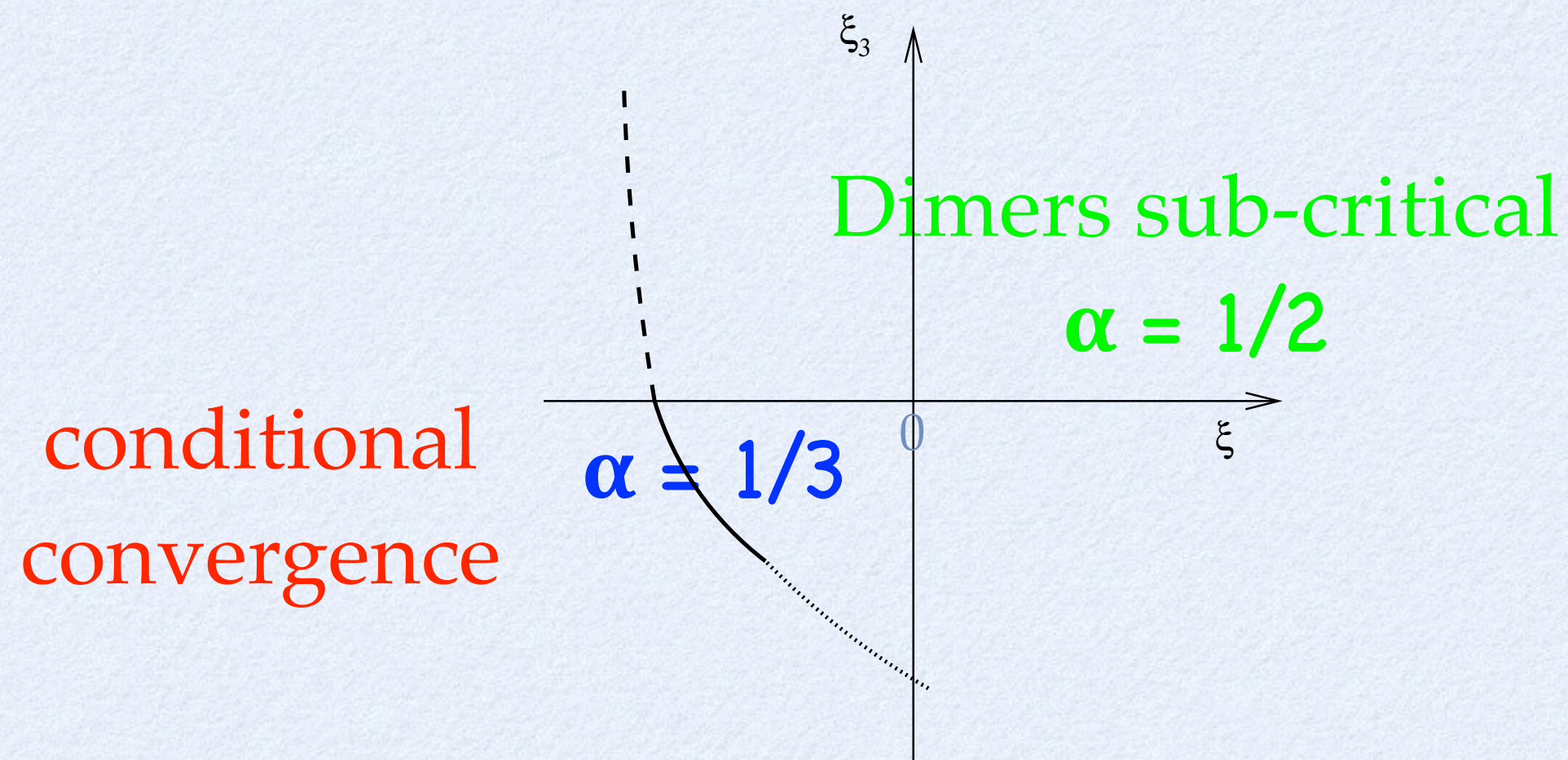
$$Z = g^{-1} W_0 - 1$$

and W_0 satisfies a cubic — we are in business

5. PHASE DIAGRAM

find $W_0 \sim W_{0c}(\{\xi\}) - A(\{\xi\}) (g_c(\{\xi\}) - g)^\alpha$

$g_c(\{\xi\})$ is free energy; set $\xi = \xi_1 = \xi_2$

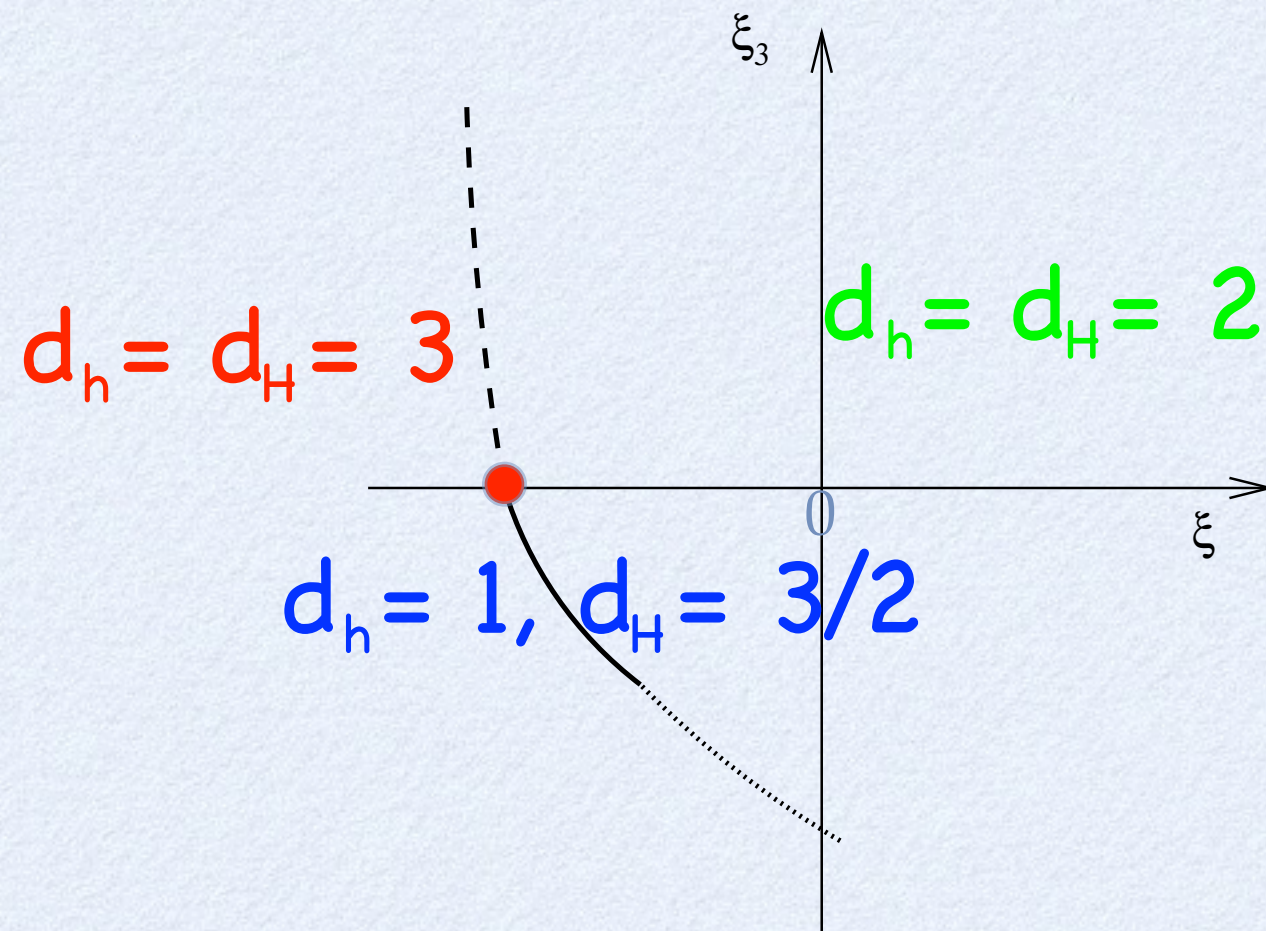


6. GEOMETRIC CHARACTERIZATION

Hausdorff dimension:

finite tree $N \sim R^{d_H}$

infinite tree $B_R \sim R^{d_h}$



9. CONCLUDING REMARKS

1. Away from dimer criticality the continuum theory is just (Horava-Lifshitz) 2d gravity
2. The segment of critical line $d_h = 1$, $d_H = 3/2$ is probably critical dimer CFT ($c = -22/5$) coupled to H-L 2d gravity
3. The AZ point $d_h = d_H = 3$ is special, but not that special — it is still a one-parameter family. Presumably it describes a CFT coupled to H-L gravity but we do not know which CFT
4. Systems with negative weights are intricate and can easily confound out intuition