

# Modelling our Galaxy

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# Outline

- Surveys
- Cosmological simulations
- Our alternative

# Surveys

- Photometry
  - 2MASS, SDSS, Vista, PanStarrs, Sky Mapper
- Spectroscopy
  - Geneva-Copenhagen S, RAVE, SEGUE, LAMOST, Gaia-ESO, Galah, LSST
- Astrometry
  - Gaia (2017, 2019,..)
- A massive job to synthesise into
  - a working model of Galaxy
  - a picture of how the Galaxy was assembled
  - Chemodynamical models are central to this effort

# Cosmological simulations

- Outside Oxford efforts focus on cosmological simulations
  - Huge N-body + hydrodynamic computations
  - One simulation requires many months of hundreds of CPUs
- Simulations of Dark-Matter clustering form bedrock of cosmology
  - Physics incredibly simple but took 15 yrs for simulations to get basic picture right (over-merging problem sorted ~2000)
- For 15 years the community has struggled with simulations that include gas & star formation
  - Without powerful feedback from star formation galaxies become too luminous and too spheroidal
  - There is direct observational evidence for powerful outflow from star-forming galaxies but the mechanisms of gas exchange between inter-galactic and interstellar media not really understood
  - Consensus is that core problem is ~ fractal structure of ISM, which extends to scales that won't be resolved numerically any time soon
- So cosmological simulations lack predictive power
  - They are far too costly to fit to data
  - It's hard even to characterise them

# Orbits

- Most orbits in plausible  $\Phi(R,z)$  are quasiperiodic
  - $\Rightarrow \exists$  angle-action coords  $(\theta, J)$
  - The momenta  $J$  are consts of motion
    - $J_r$  quantifies excursions in  $r$
    - $J_\phi = L_z$  angular momentum
    - $J_z$  quantifies excursions  $\perp$  Galactic plane
- We have developed techniques for computing  $J(x,v)$ ,  $\theta(x,v)$  and  $x(\theta, J)$ ,  $v(\theta, J)$

# (Extended) Distribution Functions

- Any non-negative  $f(\mathbf{J})$  defines an equilibrium dynamical model
- We think of galaxies as built up out of components/populations
- DFs  $f_i(\mathbf{J})$  can be straightforwardly added
- $f(\mathbf{J}) = \sum_i f_i(\mathbf{J})$
- Our basic DF for disc is “quasi isothermal”

$$f(J_r, J_z, L_z) = f_{\sigma_r}(J_r, L_z) f_{\sigma_z}(J_z, L_z),$$

$$f_{\sigma_r}(J_r, L_z) \equiv \frac{\Omega \Sigma}{\pi \sigma_r^2 \kappa} [1 + \tanh(L_z/L_0)] e^{-\kappa J_r / \sigma_r^2}$$

$$f_{\sigma_z}(J_z, L_z) \equiv \frac{\nu}{2\pi \sigma_z^2} e^{-\nu J_z / \sigma_z^2} \quad \Sigma(L_z) = \Sigma_0 e^{-R_c/R_d}$$

# Metallicity-blind models

- 2010 – 2012 we demonstrated good fits to GCS stars ( $s < 120$  pc) with superposition of a quasi-isothermal for stars of each age
- 2014 we showed that these models could predict to good (but imperfect) accuracy kinematics of stars with  $s \sim 2000$  pc (RAVE data)
- 2015 we fitted DF to kinematics of RAVE stars and from vertical density profile determined by SDSS determined parameters of dark halo under assumption that it is a (flattened) NFW model

# The thin disc DF

- A sum over ages of stars with velocity dispersion increasing with age

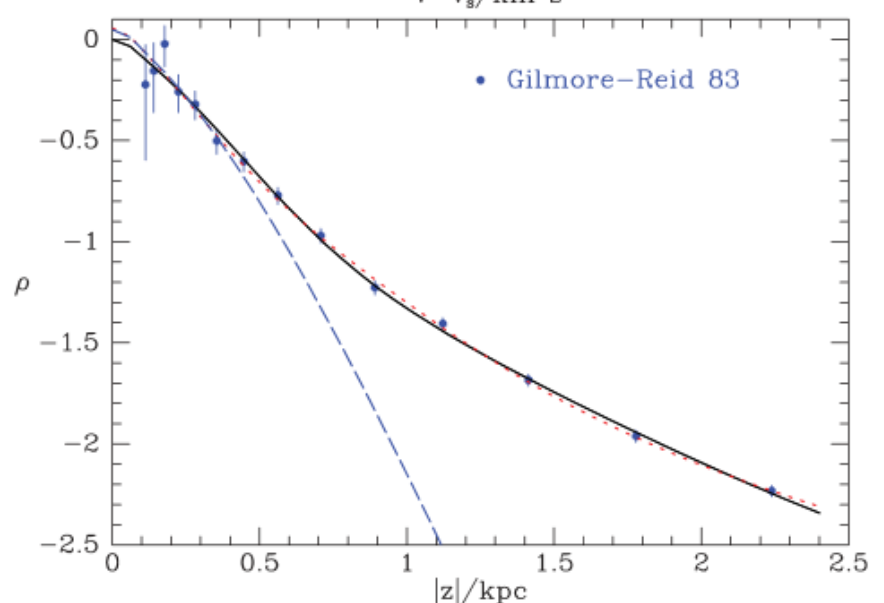
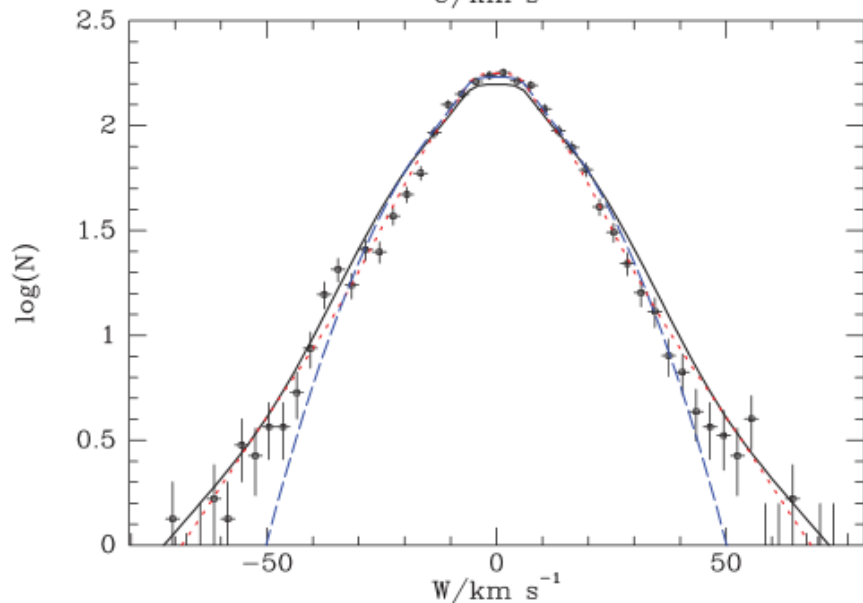
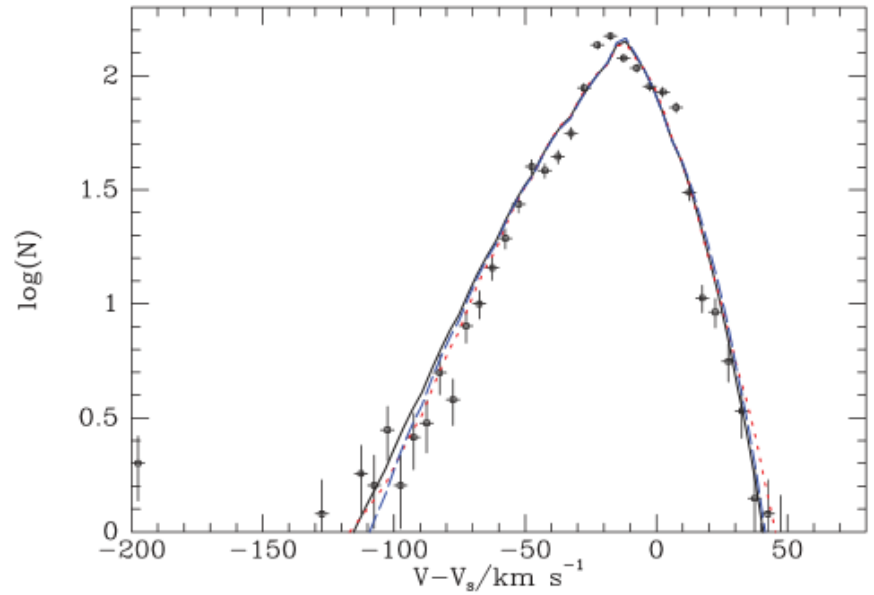
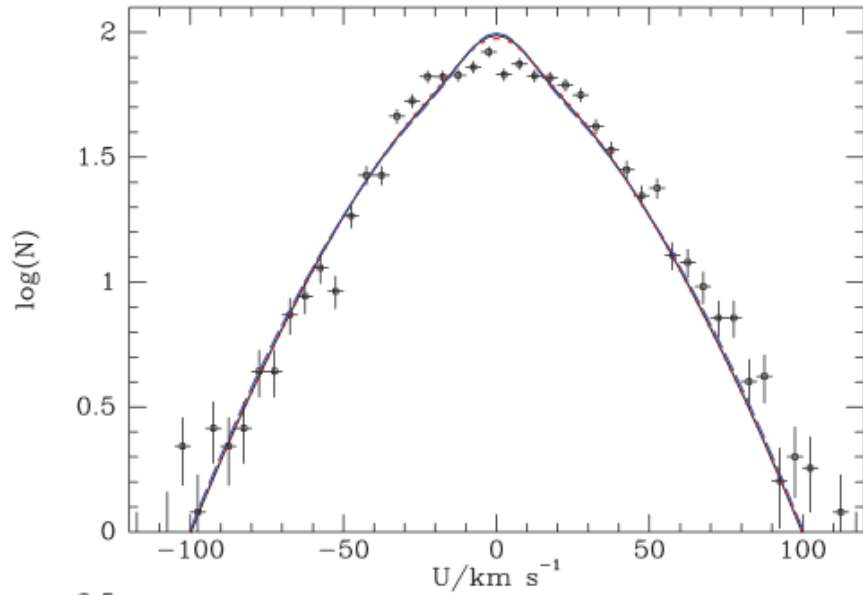
$$f_{\text{thin}}(J_r, J_z, L_z) = \frac{\int_0^{\tau_m} d\tau e^{\tau/t_0} f_{\sigma_r}(J_r, L_z) f_{\sigma_z}(J_z, L_z)}{t_0(e^{\tau_m/t_0} - 1)}$$

$$\sigma_r(L_z, \tau) = \sigma_{r0} \left( \frac{\tau + \tau_1}{\tau_m + \tau_1} \right)^\beta e^{q(R_0 - R_c)/R_d}$$

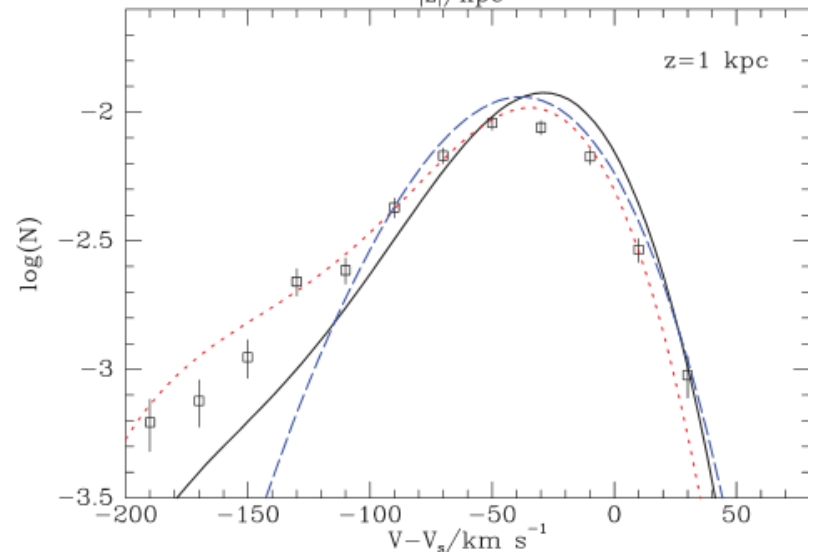
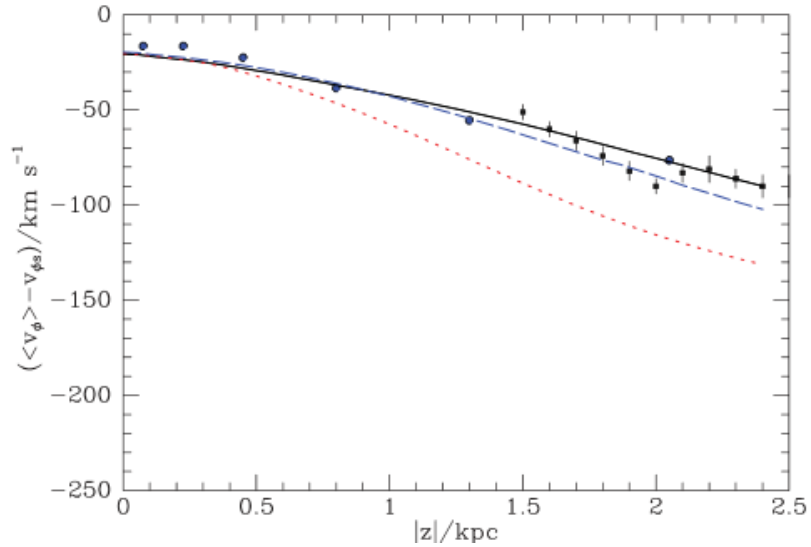
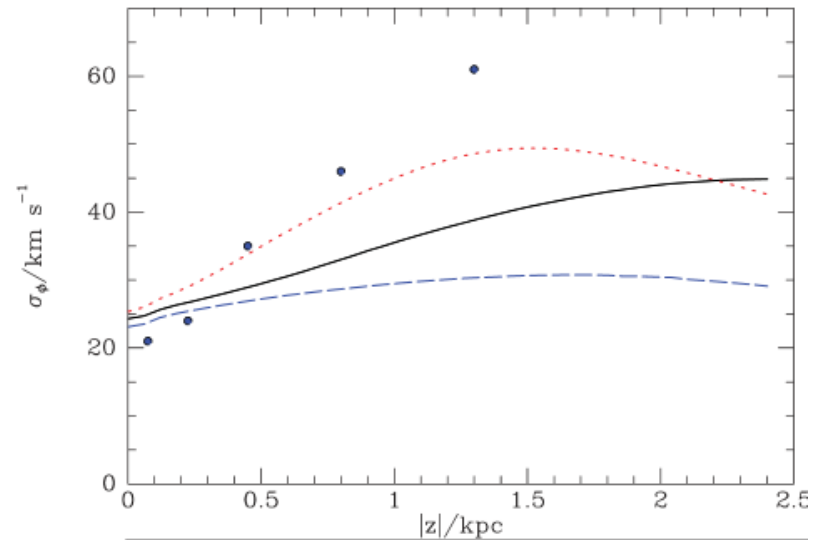
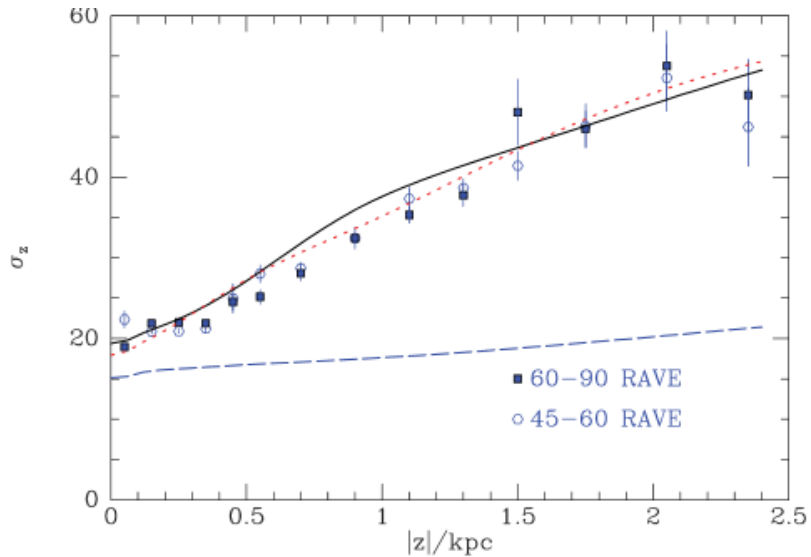
$$\sigma_z(L_z, \tau) = \sigma_{z0} \left( \frac{\tau + \tau_1}{\tau_m + \tau_1} \right)^\beta e^{q(R_0 - R_c)/R_d}.$$



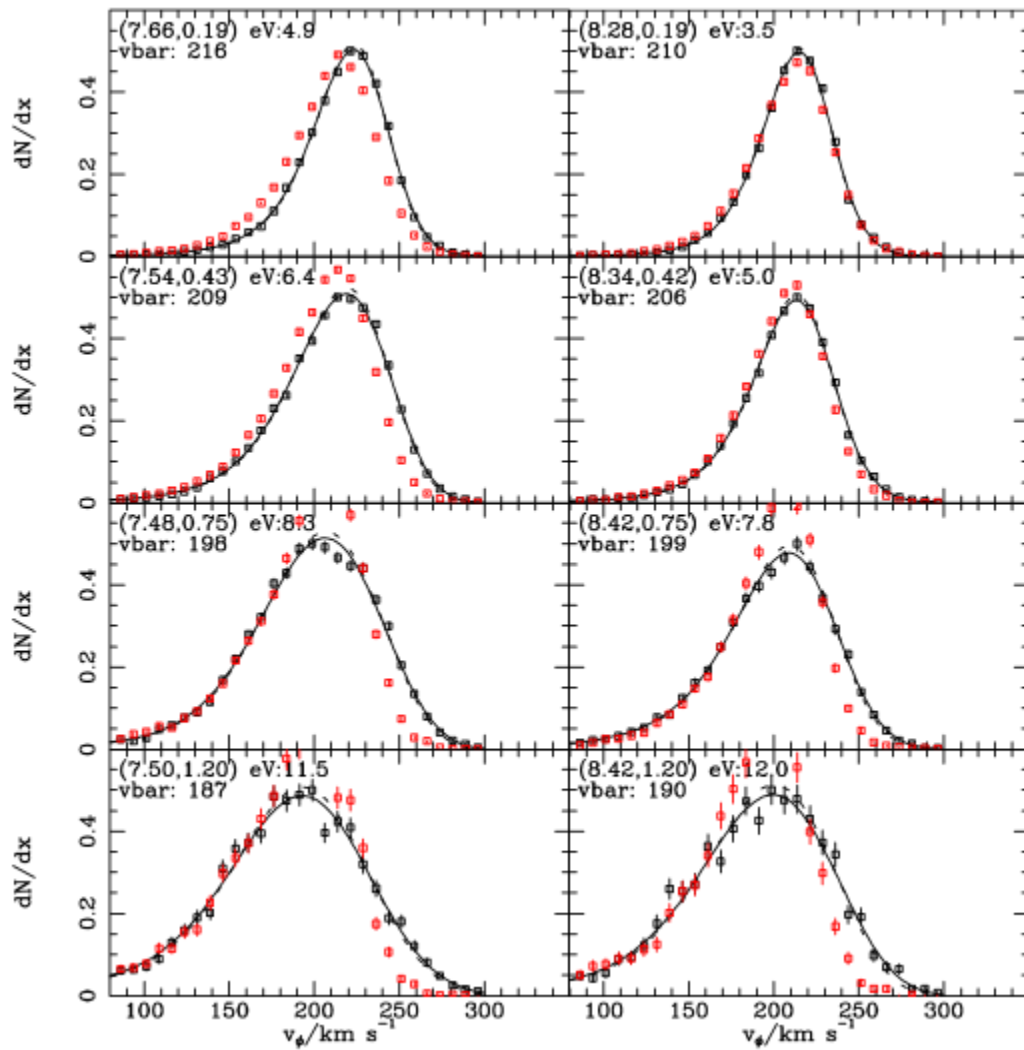
# GCS kinematics (+ Gilmore Reid density)



# Predictions

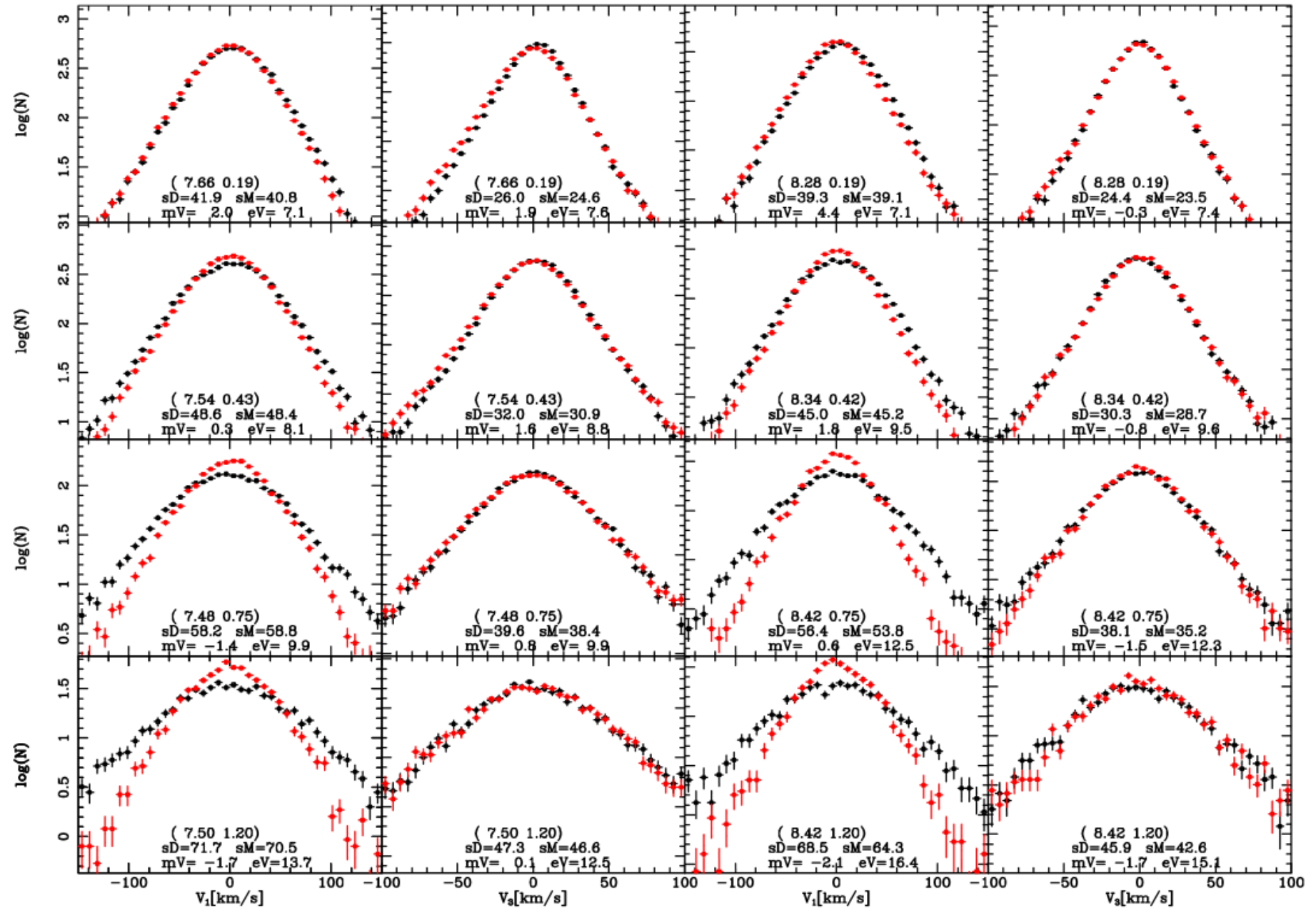


# More predictions

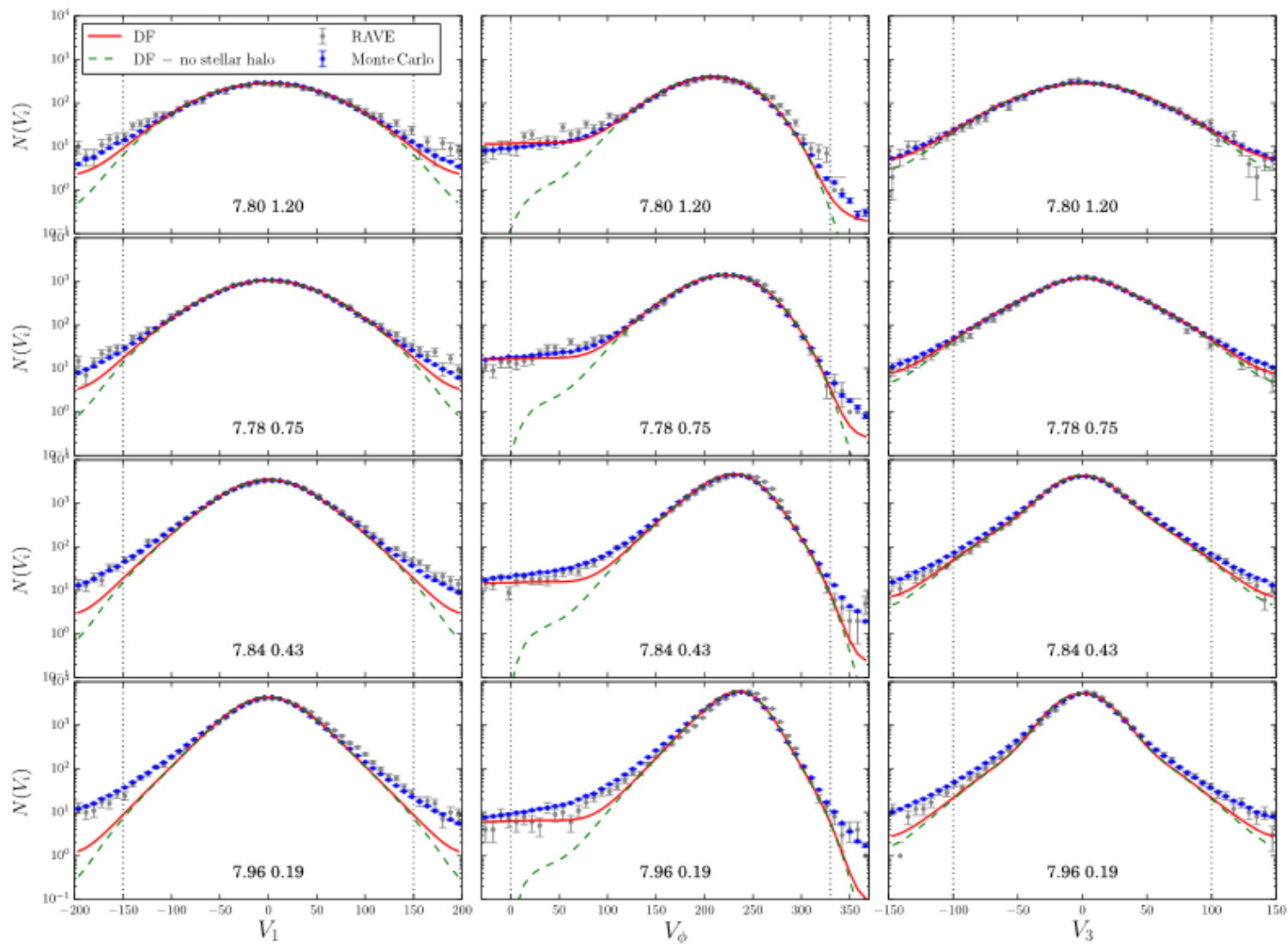


$V_\phi$

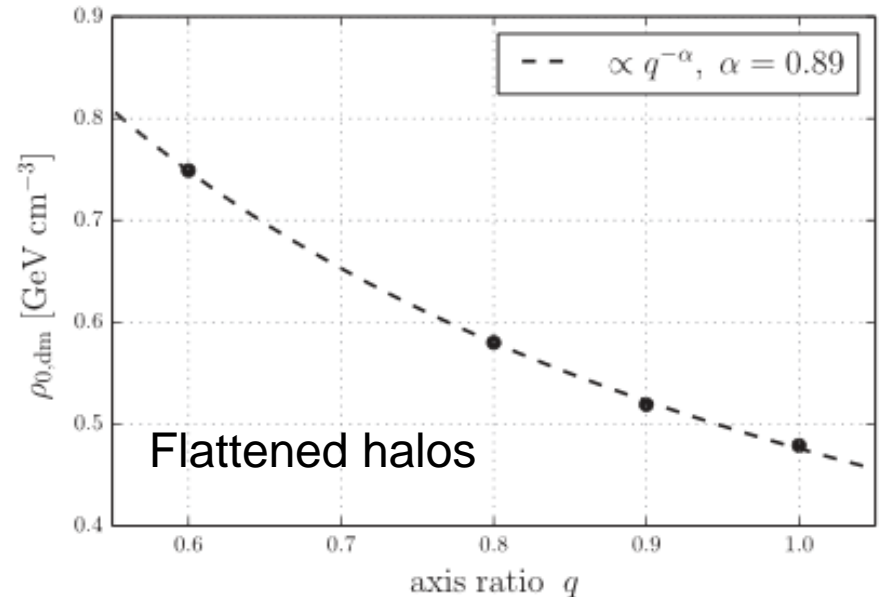
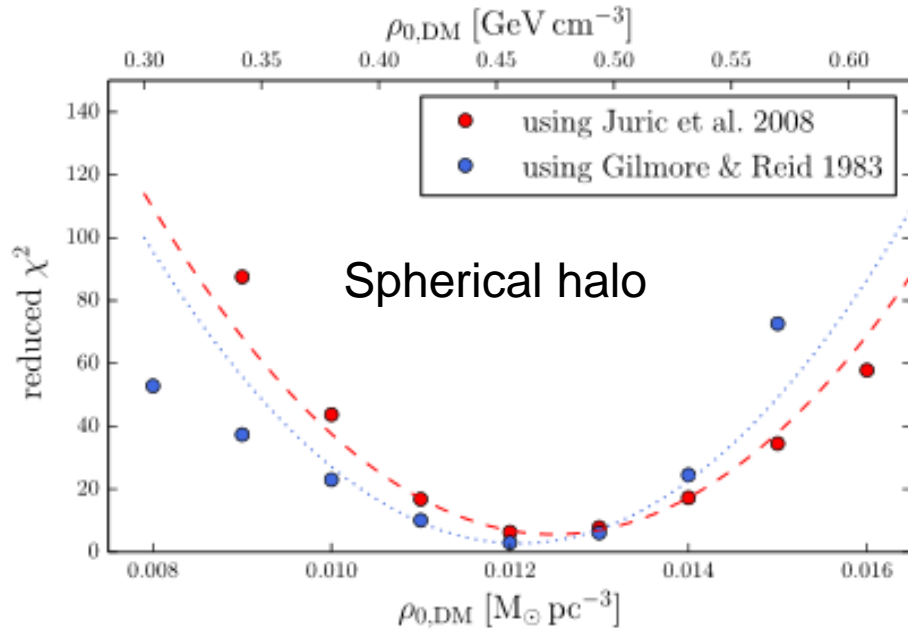
# More predictions ( $V_R$ , $V_z$ )



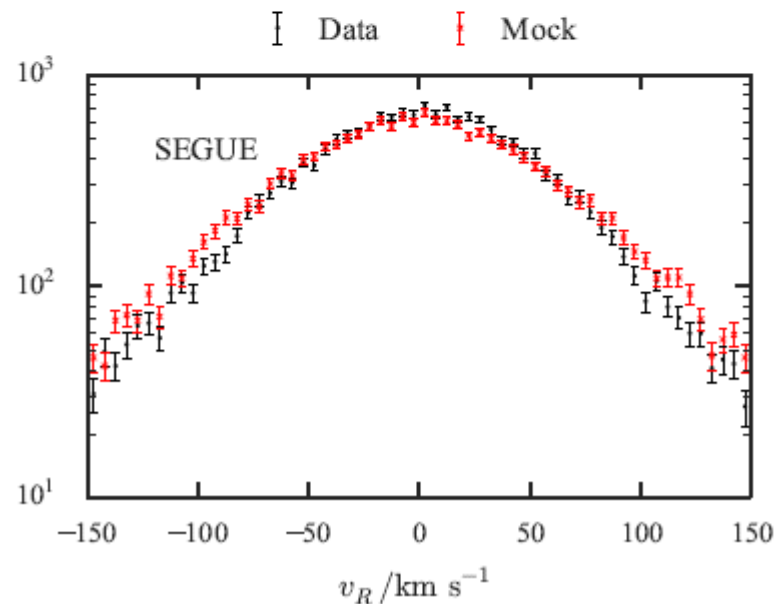
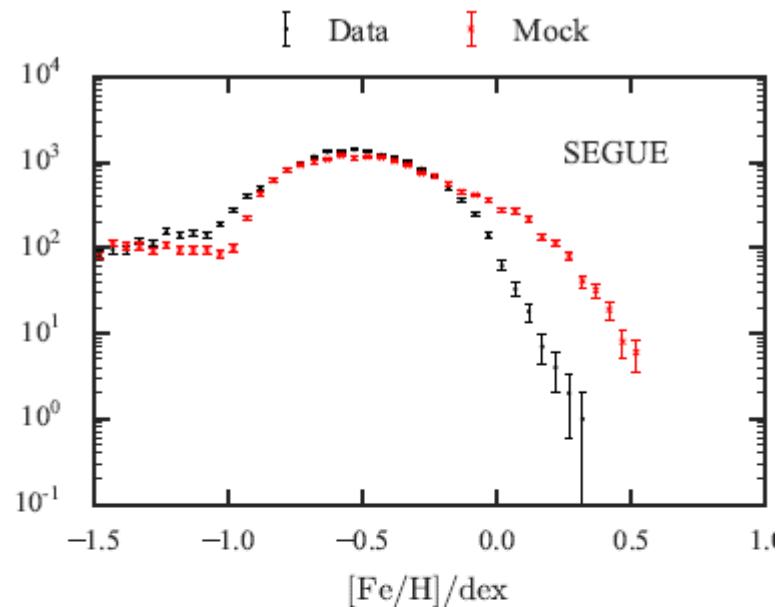
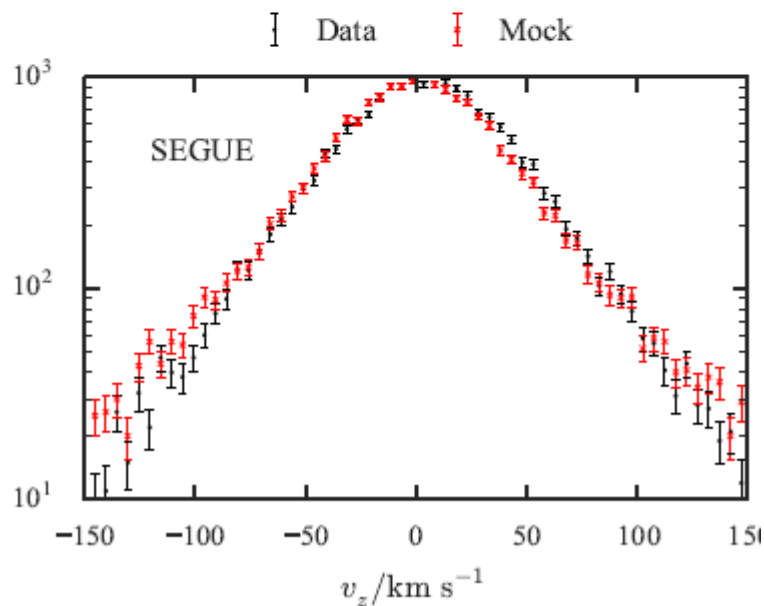
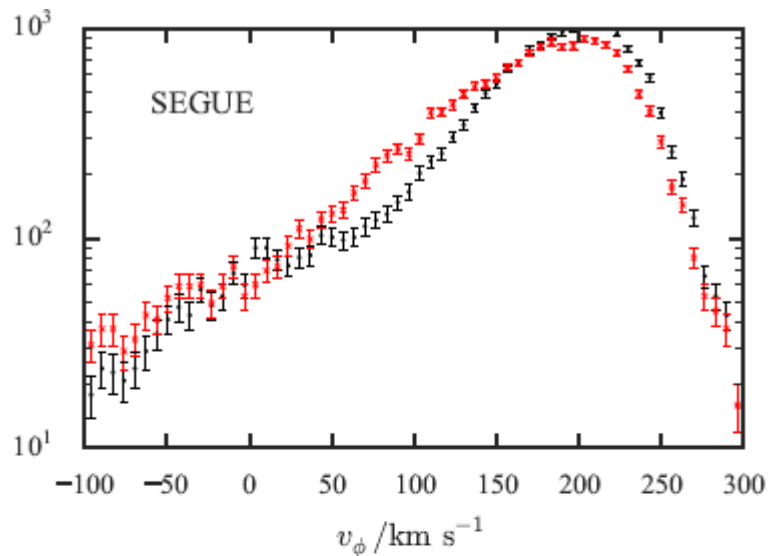
# Model fitted to RAVE



# Strongest constraints yet on local DM density (Piffl + 2014)



# Predictions



# Self-consistency

- This work all done with  $\Phi(R,z)$  computed from specified  $\rho(R,z)$
- $\rho_{\text{disc}}$  assumed to be exponential in  $R$  and double exponential in  $z$
- The disc self-consistent with its assumed contribution to  $\Phi(R,z)$  only to extent that DF produces roughly exponential + double exponential  $\rho(R,z)$
- Current models specify dark halo by  $f(J)$  rather than  $\rho(R,z)$  and we compute the potential self-consistently generated by  $f_{\text{total}}(J)$

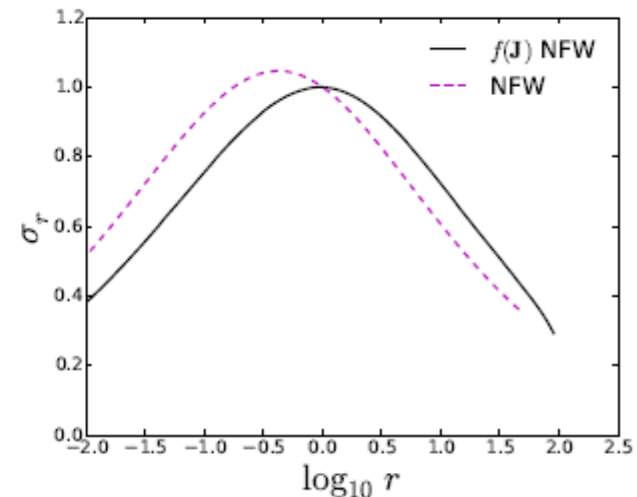
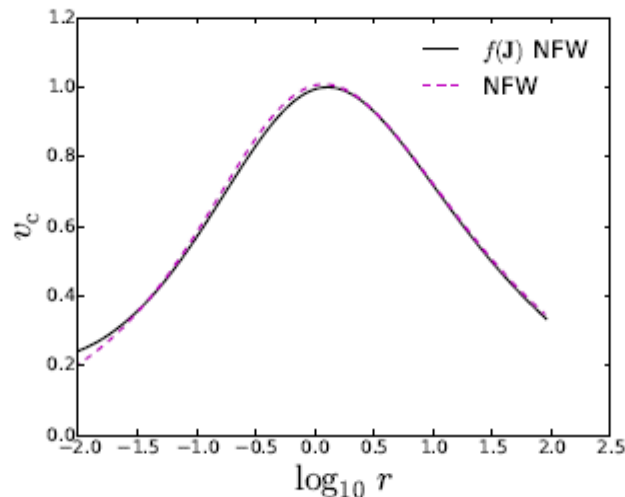
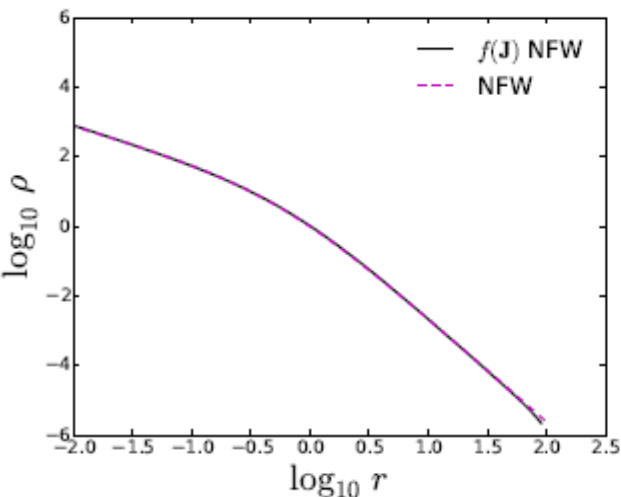


# A live NFW halo (Posti + 2015)

- A simple  $f_{\text{DM}}(\mathbf{J})$  self-consistently produces a near NFW density profile

$$f(\mathbf{J}) = \frac{M_0 [1 + J_0/h(\mathbf{J})]^{5/3}}{J_0^3 [1 + g(\mathbf{J})/J_0]^{2.9}}$$

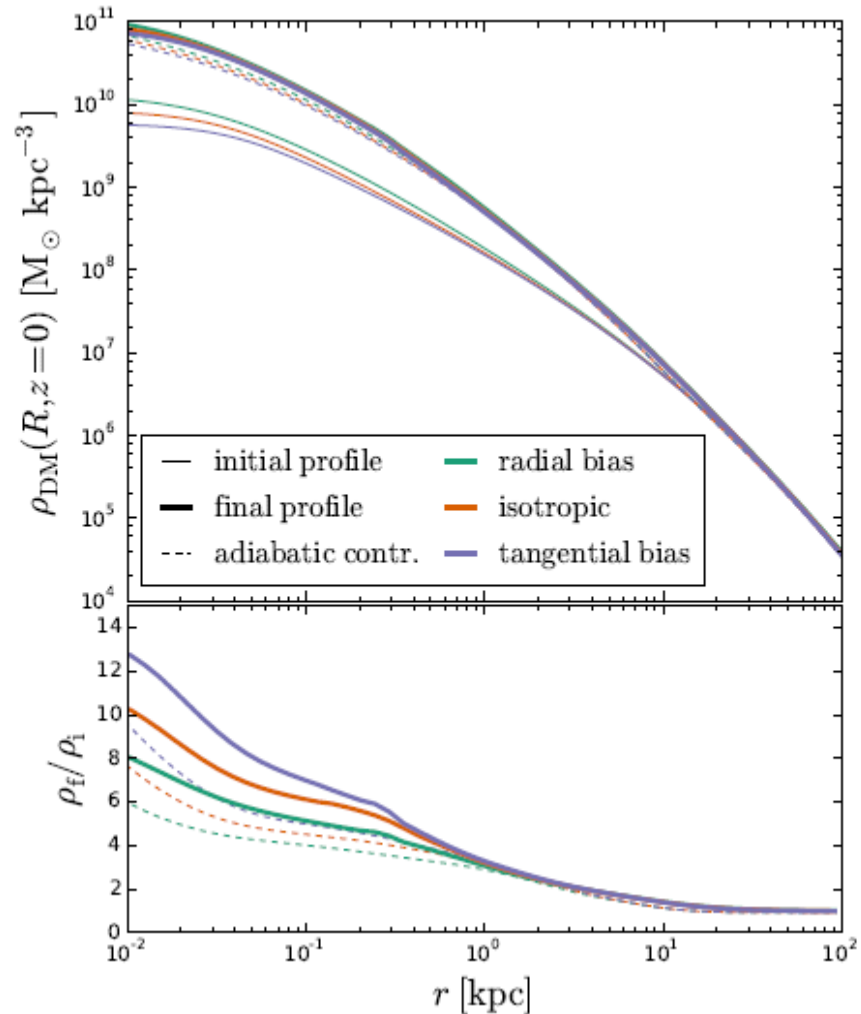
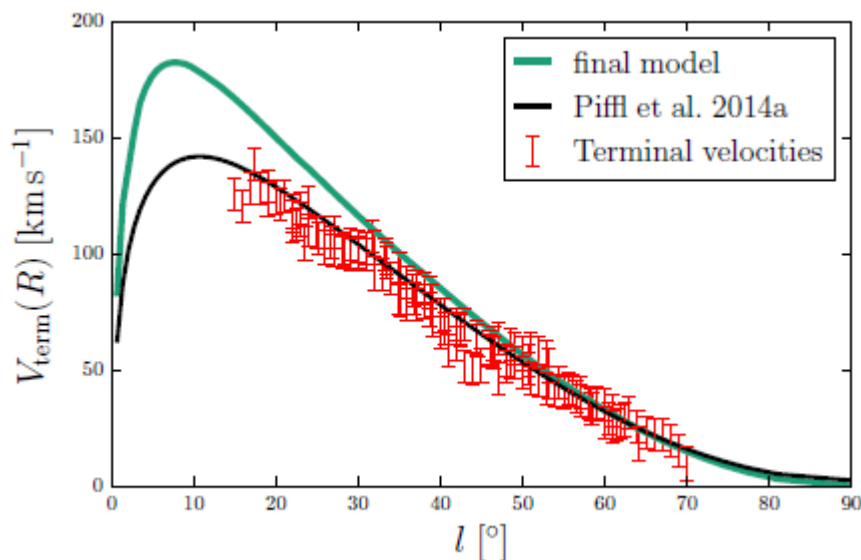
- $g$  &  $h$  homogeneous fns of order unity – in spherical case the same linear function
- $J_0$  sets the linear scale



# Adiabatic addition of the disc

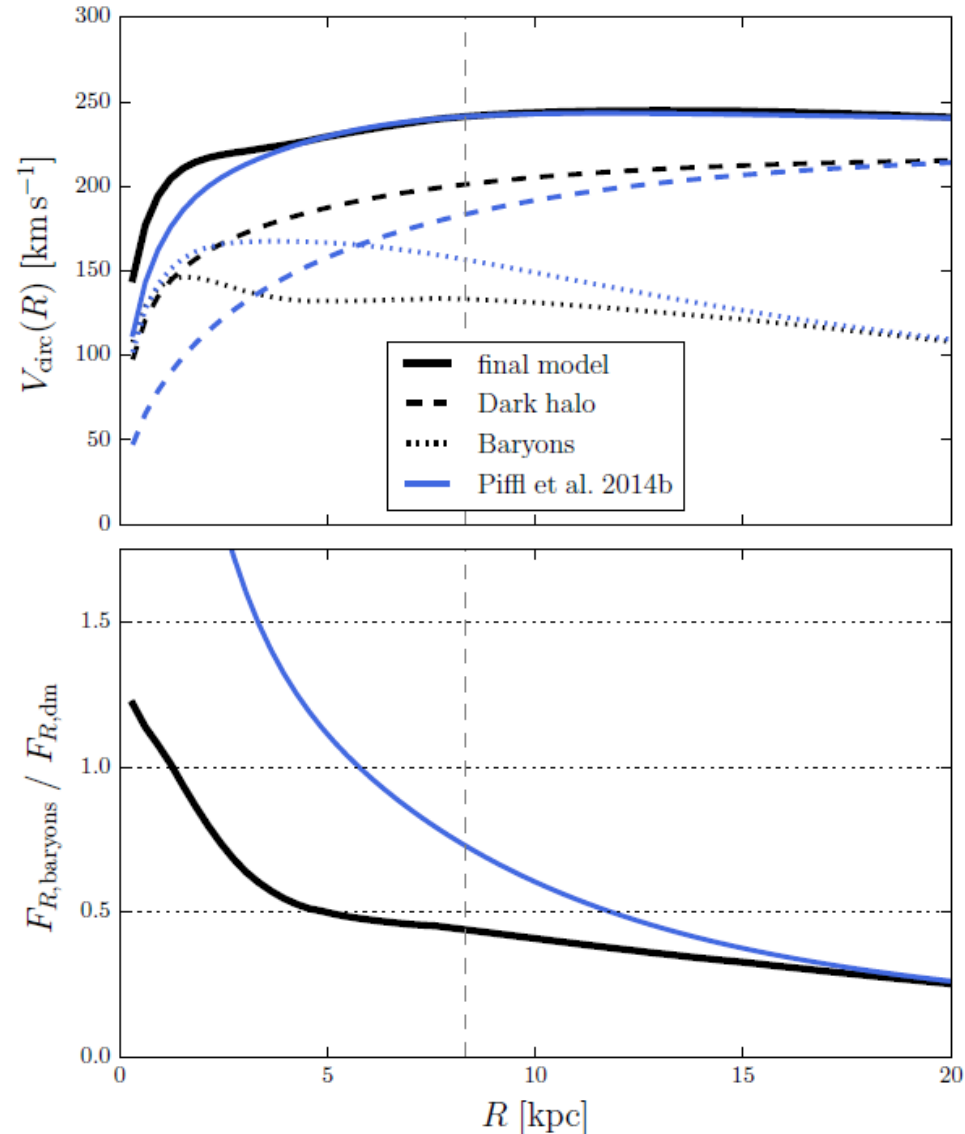
(Piffl + 2015)

- The disc has grown over many dynamical times, and  $J_i$  adiabatic invariants, so perhaps  $f_{\text{DM}}(\mathbf{J})$  same after disc growth as in DM-only simulation
- When we adopt this  $f_{\text{DM}}$  and  $f_{\text{disc}}$  fitted to RAVE data with NFW  $\rho_{\text{DM}}$ , we no longer match  $v_c(R)$  curve because DM has been strongly compressed by gravity of disc



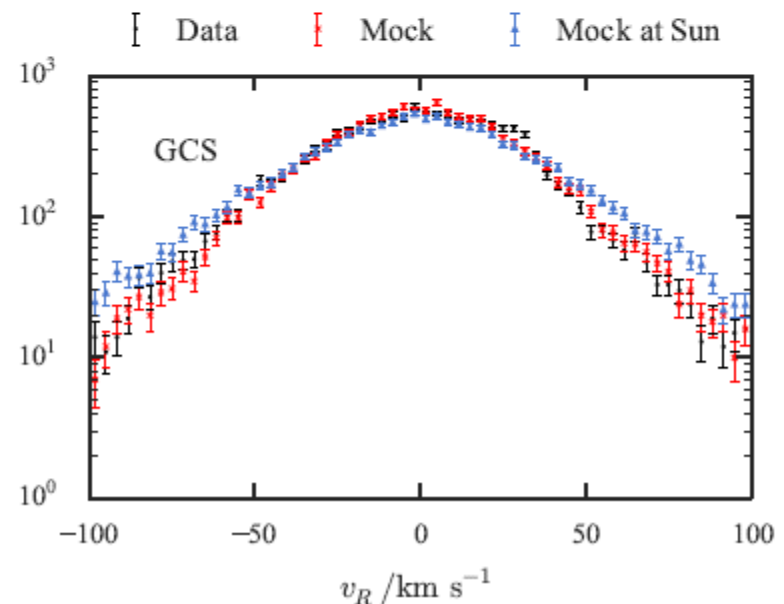
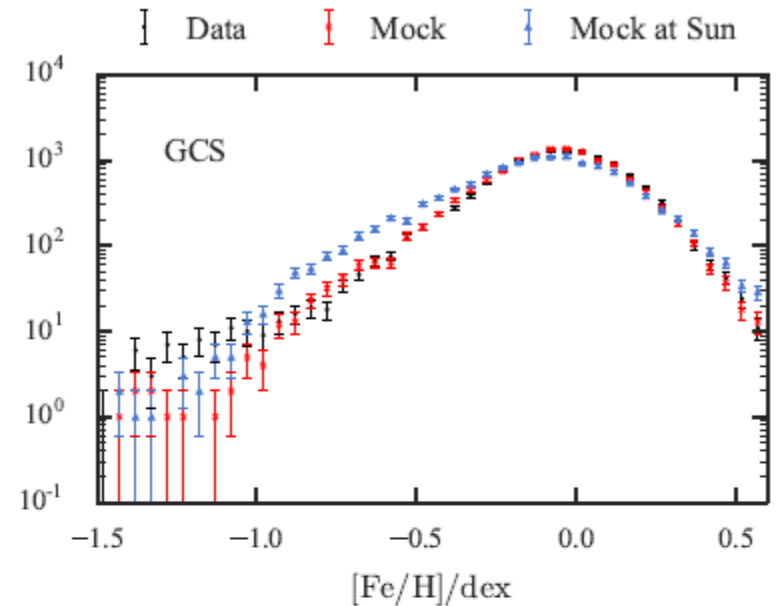
# Currently (Piffi & B 2015)

- Investigating whether data can be fitted using expected NFW  $f(J)$  – scale length of disc has to be increased and mass reduced. Then problematic to get  $\rho_{\text{disc}}$  sufficiently peaked to plane
- Likely that we have evidence for central heating of DM by bar/spiral structure...



# EDFs (Sanders & B 2015)

- Luminosity and colour essential for predicted observables
  - They depend of  $[\text{Fe}/\text{H}]$  as well as age
  - So one has to extend to  $f(J, \tau, [\text{Fe}/\text{H}])$
- Our choice of  $f(J, \tau, [\text{Fe}/\text{H}])$  is motivated by a model of how spiral structure heats disc and drives “radial migration” (shifts in  $L_z$ )
  - But  $f(J, \tau, [\text{Fe}/\text{H}])$  ultimately independent of model



# Conclusions

- With analytic EDFs for disc and a DF for DM one can fit large bodies of data and successfully predict other data
- These models are cheaper than cosmological simulation by several orders of magnitude and no less/more rigorous
- Their physical content is much easier to appreciate & I believe the way forward is to fit our EDFs to simulations
- We are mapping the Galaxy's DM distribution in some detail
- We may soon have evidence of DM heating by baryons