

# *The Vernacular of the $S$ -Matrix*

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NBIA-Oxford Colloquium



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International Academy



UNIVERSITY OF  
OXFORD

# Organization and Outline

- 1 *Spiritus Movens*: a Parable from Collider Physics
  - A Simple, Practical Problem in Quantum Chromodynamics
- 2 The *Physical Vernacular* of the  $S$ -Matrix: On-Shell Observables
  - Physically Observable Data Describing Asymptotic States
  - Beyond (Mere) Scattering Amplitudes: On-Shell Functions
  - Basic Building Blocks:  $S$ -Matrices for Three Massless Particles
- 3 The On-Shell Analytic  $S$ -Matrix: All-Loop Recursion Relations
  - Building-up Diagrams with “BCFW” Bridges
  - On-Shell (Recursive) Representations of Scattering Amplitudes
  - *Exempli Gratia*: On-Shell Manifestations of Tree Amplitudes
- 4 The Combinatorics of Scattering (and Grassmannian Geometry)
  - Combinatorial Classification of On-Shell Functions in Planar SYM
  - Canonical Coordinates, Computation, & Auxiliary Grassmannian
- 5 The Ongoing Revolution: Toward a Complete Reformulation of QFT

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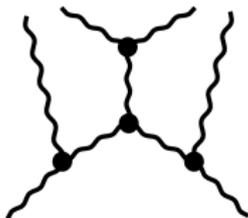
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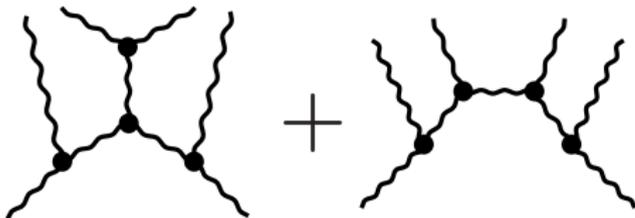
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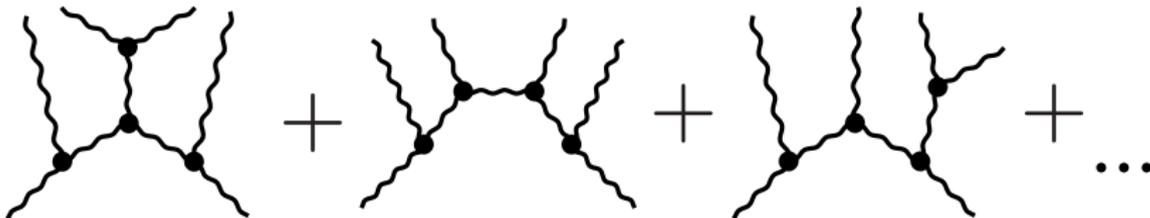
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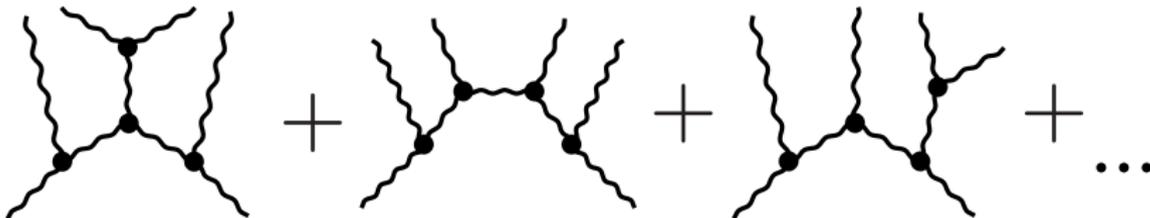
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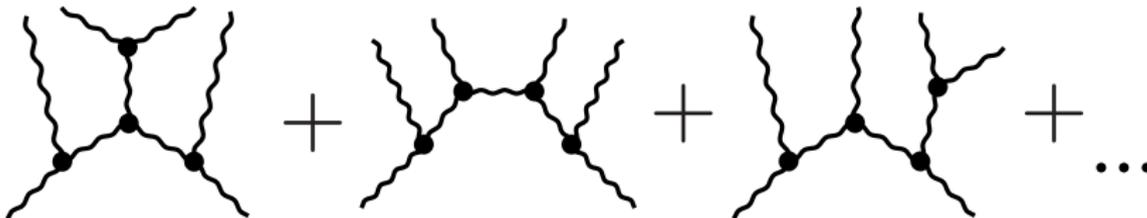
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## Supercollider physics

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Eichten *et al.* summarize the motivation for exploring the 1-TeV ( $\sim 10^{12}$  eV) energy scale in elementary particle interactions and explore the capabilities of proton-antiproton colliders with beam energies between 1 and 50 TeV. The authors calculate the production rates and characteristics for a number of conventional processes, and discuss their intrinsic physics interest as well as their role as backgrounds to more exotic phenomena. The authors review the theoretical motivation and expected signatures for several new phenomena which may occur on the 1-TeV scale. Their results provide a reference point for the choice of machine parameters and for experiment design.

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Eichten *et al.*: Supercollider physics

617

TeV. From Fig. 76 we find the corresponding two-jet cross section (at  $p_T = 0.5 \text{ TeV}/c$ ) to be about  $7 \times 10^{-2}$  nb/CeV, which is larger by an order of magnitude. Let us next consider the cross section in the neighborhood of the peak in Fig. 80. The integrated cross section in the bin  $0.3 \text{ GeV} < E_T < 0.4$  is approximately 0.1 nb/CeV with transverse energy given (roughly by  $|E_T| \approx |T_{xy}| \times \cos(\theta) \approx 350 \text{ GeV}$ . The corresponding two-jet cross section, again from Fig. 76, is approximately 10 nb/CeV, which is larger by 2 orders of magnitude. In fact, we have certainly underestimated  $\langle E_T \rangle$  and thus somewhat overestimated the two-jet/one-jet ratio in this second case.

We draw two conclusions from this very casual analysis:

At least at small-to-moderate values of  $E_T$ , two-jet events should account for most of the cross section. The observed cross section is large enough that a detailed study of this topology should be possible.

$$\sigma(E_T) \approx \int_{\theta_{min}}^{\theta_{max}} \int_{\phi_{min}}^{\phi_{max}} d\theta_{12} \int_{\phi_{min}}^{\phi_{max}} d\theta_{13} \frac{\sigma(E_T, \theta_{12}, \theta_{13}, \theta_{23})}{\sigma(E_T)} \quad (8.47)$$

where  $\sigma(E_T, \theta_{ij})$  is the two-jet cross section and  $\theta_{ij}$  is the minimum  $E_T$  required for a discernible two-jet event. For a more careful study of double-parton scattering at QPC and Tevatron energies, see Flemer and Teplitz (1983).

In view of the progress that multiple-scattering helps, improving our understanding of the QCD background is an urgent priority for further study.

### D. Summary

We conclude this section with a brief summary of the range of jet energy which are accessible for various beam energies and luminosities. We find essentially no differences between  $pp$  and  $p\bar{p}$  collisions, so only  $pp$  results will be given except at  $\sqrt{s} = 2 \text{ TeV}$  where  $p\bar{p}$  runs are contemplated. Figure 30A shows the  $E_T$  range which can be explored at the level of at least one event per (GeV) of  $E_T$  per unit rapidity at 90° in the c.m. (compare Figs. 77–79 and 83). The results are presented in terms of the transverse energy per event  $E_T$  which corresponds to twice the transverse momentum  $p_T$  of a jet. In Fig. 100 we plot the value of  $E_T$  that distinguishes the region to which the two-gluon, quark-gluon, and quark-quark final states are dominant. Comparing with Fig. 30A, we find that while the accessible range of  $E_T$  are large, it is more extremely difficult to obtain a clean sample of quark jets. Useful for understanding trigger rates is the total cross section for two jets (integrated over  $E_T = 0.2 - 1.2 \text{ GeV}$  for both jets) as a function of invariant  $\sqrt{s} = 2.5$  to  $2.5 \times 2.5$ . This is shown for  $pp$  collisions in Fig. 30B.

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It is apparent that these questions are amenable to detailed investigation with the aid of realistic Monte Carlo simulations. Given the elementary two-color cross sections and reasonable parametrizations of the fragmentation functions, this exercise can be carried out with some degree of confidence.

For double events containing more than three jets, the theoretical situation is considerably more primitive. A specific question of interest concerns the QCD four-jet background to the detection of  $H^0 \rightarrow \mu\mu$  pairs in their isotropic decay. The cross sections for the elementary two-color process have not been calculated, and their complexity is such that they may not be evaluated in the foreseeable future. It is worthwhile to seek estimates of the four-jet cross section, even if these are only reliable in restricted regions of phase space.

Another background source of four-jet events is double parton scattering, as shown in Fig. 103. If all the parton momenta fractions are small, the two interactions may be treated as uncorrelated. The resulting four-jet cross section with transverse energy  $E_T$  may then be approximated by

### IV. ELECTROWEAK PHENOMENA

In this section we discuss the supercollider processes associated with the standard model of the weak and electroweak interactions (Glashow, 1961; Weinberg, 1967; Salam, 1960). By "standard model" we understand the SU(2) $\times$ U(1) theory applied to three quark and lepton doublets, and with the gauge symmetry broken by a single complex Higgs doublet. The particles associated with the electroweak interaction are identified with the self-handed charged intermediate bosons  $W^\pm$ , the neutral intermediate

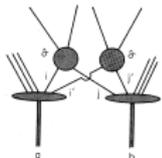


FIG. 303. Four-jet topology arising from two independent parton interactions.

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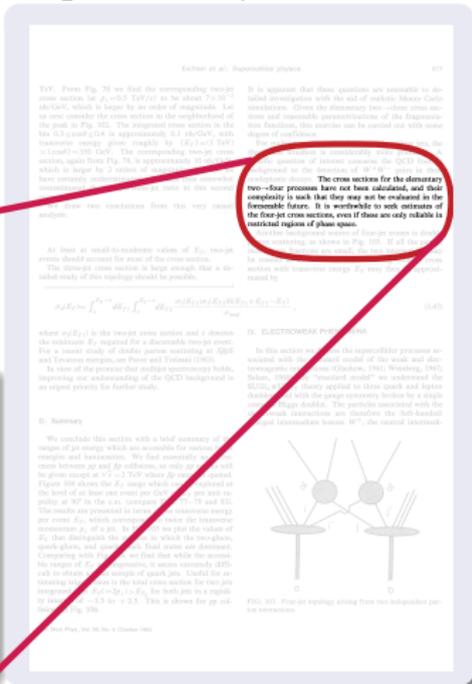
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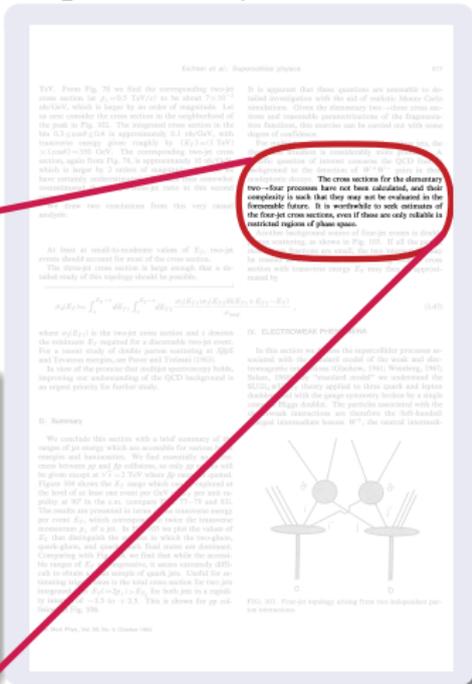


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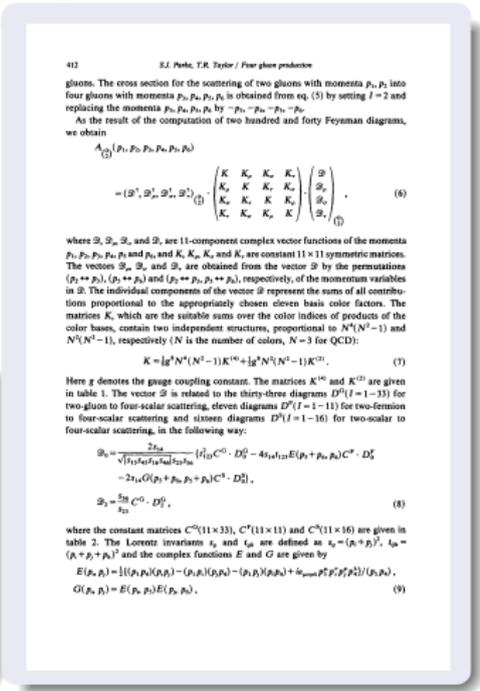
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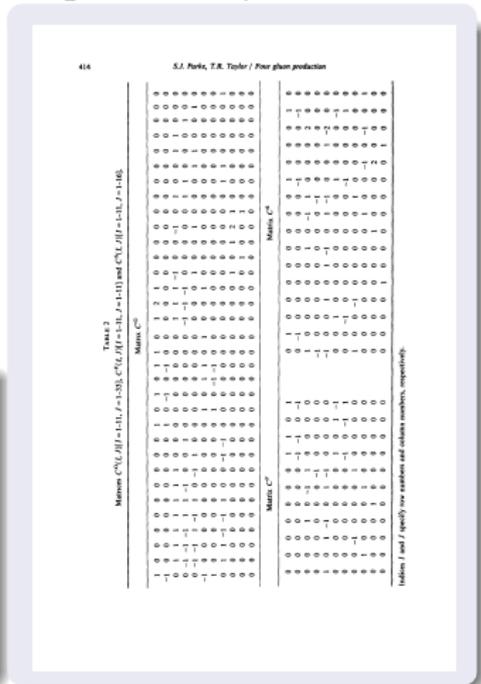
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416 S.J. Parke, T.R. Taylor / Four gluon production

$$\begin{aligned}
 D_1^2(9) &= \frac{4}{s_{12}s_{34}s_{13}} [(s_1 - p_1 + p_2)(p_1 + p_2 - p_3)] E(p_3, p_4) \\
 &\quad - [(p_1 - p_2 + p_3)(p_1 - p_1 + p_3)] E(p_2, p_3) + [(p_1 + p_2 - p_3)] E(p_2, p_3 - p_4), \\
 D_1^2(10) &= \frac{4}{s_{12}s_{34}s_{13}} [(s_1 + p_1 - p_2)(p_1 - p_1 + p_3)] E(p_3, p_4) \\
 &\quad - [(p_1 - p_2 + p_3)(p_1 - p_1 + p_3)] E(p_2, p_3) + [(p_1 + p_2 - p_3)] E(p_2 - p_3, p_4), \\
 D_1^2(11) &= \frac{s_1}{s_{12}s_{13}} [s_{12} - s_{34} + s_{34}], \\
 D_1^2(12) &= \frac{-s_1}{s_{12}s_{13}} [s_{12} - s_{34} - s_{34}], \\
 D_1^2(13) &= \frac{s_1}{s_{12}s_{13}} [s_{12} + s_{34}] [s_{12} - s_{34} + s_{34}], \\
 D_1^2(14) &= \frac{-s_1}{s_{12}s_{13}} [s_{12} + s_{34}] [s_{12} - s_{34} - s_{34}], \\
 D_1^2(15) &= \frac{-s_1}{s_{12}s_{34}} (p_1 - p_3)(p_2 - p_4), \\
 D_1^2(16) &= \frac{-4}{s_{12}s_{34}s_{13}} [s_{12} - s_{34} + s_{34}] E(p_1, p_2), \\
 D_1^2(17) &= \frac{4}{s_{12}s_{34}s_{13}} [s_{12} - s_{34} - s_{34}] E(p_1, p_2), \\
 D_1^2(18) &= \frac{-4}{s_{12}s_{34}s_{13}} [2(p_1 + p_2)(p_3 - p_4) + s_{12}] E(p_3, p_4), \\
 D_1^2(19) &= \frac{-2}{s_{12}s_{34}} E(p_2, p_1 - p_3), \\
 D_1^2(20) &= \frac{2}{s_{12}s_{34}} E(p_1 - p_3, p_2), \\
 D_1^2(21) &= \frac{-4}{s_{12}s_{34}s_{13}} [s_{12} - s_{34} + s_{34}] E(p_1, p_2), \\
 D_1^2(22) &= \frac{4}{s_{12}s_{34}s_{13}} [s_{12} - s_{34} - s_{34}] E(p_1, p_2), \\
 D_1^2(23) &= \frac{4}{s_{12}s_{34}s_{13}} [2(p_1 + p_2)(p_3 - p_4) + s_{12}] E(p_3, p_4),
 \end{aligned}$$

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S.J. Parke, T.R. Taylor / Four gluon production 417

$$\begin{aligned}
 D_1^0(24) &= \frac{-2}{t_{12}t_{34}} E(p_2 - p_1, p_1), \\
 D_1^0(25) &= \frac{2}{t_{14}t_{32}} E(p_2, p_1 - p_1), \\
 D_1^0(26) &= \frac{-2}{t_{12}t_{32}} E(p_2, p_1 - p_1), \\
 D_1^0(27) &= \frac{2}{t_{34}t_{12}} E(p_1 - p_2, p_1), \\
 D_1^0(28) &= \frac{2}{t_{14}t_{32}} E(p_2, p_1 - p_1), \\
 D_1^0(29) &= \frac{-2}{t_{34}t_{12}} E(p_1 - p_2, p_1), \\
 D_1^0(30) &= \frac{4}{t_{12}t_3 t_{12}^2} [(p_1 + p_2 - p_1)(p_2 + p_1 - p_1) - t_{32}] E(p_2, p_1), \\
 D_1^0(31) &= \frac{4}{t_{12}t_3 t_{12}^2} [(p_1 + p_2 - p_1)(p_2 - p_1 + p_1) + t_{32}] E(p_2, p_1), \\
 D_1^0(32) &= \frac{4}{t_{12}t_3 t_{12}^2} [(p_1 - p_2 + p_1)(p_2 + p_1 - p_1) + t_{32}] E(p_2, p_1), \\
 D_1^0(33) &= \frac{4}{t_{12}t_3 t_{12}^2} [(p_1 - p_2 + p_1)(p_2 - p_1 + p_1) - t_{32}] E(p_2, p_1),
 \end{aligned} \tag{11}$$

where  $\delta_2 = 1$ .  
 The diagrams  $D_1^0$  are obtained from  $D_1^0$  by replacing  $\delta_2$  by  $\delta_2 = 0$  and the functions  $E(p_i, p_j)$  by  $G(p_i, p_j)$ .  
 The diagrams  $D_1^0$  are listed below:

$$\begin{aligned}
 D_1^0(1) &= \frac{4}{t_{12}t_3 t_{12}^2} [F(p_1, p_2)E(p_2, p_1) - F(p_2, p_1)E(p_2, p_1) \\
 &\quad + \{F(p_2, p_1) + t_{32}\}E(p_1, p_1)], \\
 D_1^0(2) &= \frac{-4}{t_{34}t_2 t_{34}^2} [F(p_1, p_2) + \{t_{34}\}E(p_2, p_1) \\
 &\quad + \{F(p_2, p_1) + \{t_{34}\}E(p_2, p_1) - F(p_2, p_1)E(p_2, p_1)], \\
 D_1^0(3) &= \frac{4}{t_{12}t_3 t_{12}^2} [F(p_1, p_2)E(p_2, p_1) - F(p_2, p_1)E(p_2, p_1) \\
 &\quad - \{F(p_2, p_1) - \{t_{32} - \{t_{32}\} + \{t_{32}\}E(p_2, p_1)].
 \end{aligned}$$

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$$\begin{aligned}
 D_1^2(7) &= \frac{1}{232541121} [s_{12} - s_{14} + s_{24}] [s_{12} - s_{14} - s_{24}], \\
 D_1^2(8) &= \frac{1}{248202144} [s_{23} + s_{24} - s_{32}] [s_{12} - s_{14} + s_{24}], \\
 D_1^2(9) &= \frac{1}{232541121} [s_{34} + s_{42} - s_{23}] [s_{23} - s_{24} + s_{32}], \\
 D_1^2(10) &= \frac{1}{248202144} (p_1 - p_2)(p_2 - p_3), \\
 D_1^2(11) &= \frac{1}{248202144} (p_1 - p_2)(p_2 - p_3), \\
 D_1^2(12) &= \frac{1}{248202144} (p_1 - p_2)(p_2 - p_3), \\
 D_1^2(13) &= \frac{1}{248202144} (p_1 - p_2)(p_2 - p_3), \\
 D_1^2(14) &= \frac{1}{248202144} (p_1 - p_2)(p_2 - p_3), \\
 D_1^2(15) &= -\frac{1}{248202144} [(p_2 + p_3)(p_2 - p_3)] [(p_1 - p_2)(p_2 - p_3)] \\
 &\quad + [(p_2 - p_3)(p_2 - p_3)] [(p_1 - p_2)(p_2 + p_3)] \\
 &\quad + [(p_2 + p_3)(p_2 - p_3)] [(p_1 - p_2)(p_2 - p_3)], \\
 D_1^2(16) &= \frac{2}{248202144} [(p_2 - p_3)(p_2 + p_3)] [(p_1 - p_2)(p_2 - p_3)] \\
 &\quad + [(p_2 + p_3)(p_2 - p_3)] [(p_1 - p_2)(p_2 - p_3)] \\
 &\quad + [(p_1 - p_2)(p_2 + p_3)] [(p_1 - p_2)(p_2 - p_3)]. \tag{13}
 \end{aligned}$$

The preceding list completes the result. Let us recapitulate now the numerical procedure of calculating the full cross section. First the diagrams  $D$  are calculated by using eqs. (11)–(13). The result is substituted into eq. (8) to obtain the vectors  $\mathcal{D}_i$  and  $\mathcal{D}_j$ . After generating the vectors  $\mathcal{D}_k, \mathcal{D}_l, \mathcal{D}_m, \mathcal{D}_n, \mathcal{D}_o, \mathcal{D}_p, \mathcal{D}_q, \mathcal{D}_r, \mathcal{D}_s$  and  $\mathcal{D}_t$  by the appropriate permutations of momenta, eq. (6) is used to obtain the functions  $A_i$  and  $A_j$ . Finally, the total cross section is calculated by using eq. (5). The FORTRAN 5 program based on such a scheme generates ten Monte Carlo points in less than a second on the heterotic CDC CYBER 175/875.

Given the complexity of the final result, it is very important to have some reliable testing procedures available for numerical calculations. Usually in QCD, the multi-gluon amplitudes are tested by checking the gauge invariance. Due to the specific

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## THE CROSS SECTION FOR FOUR-GLUON PRODUCTION BY GLUON-GLUON FUSION

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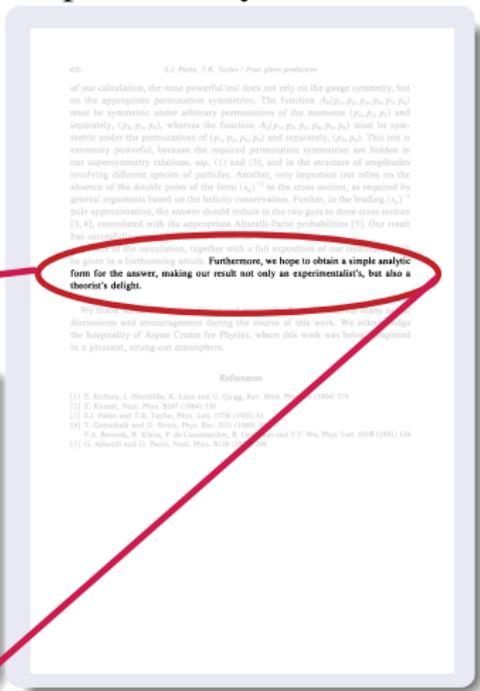
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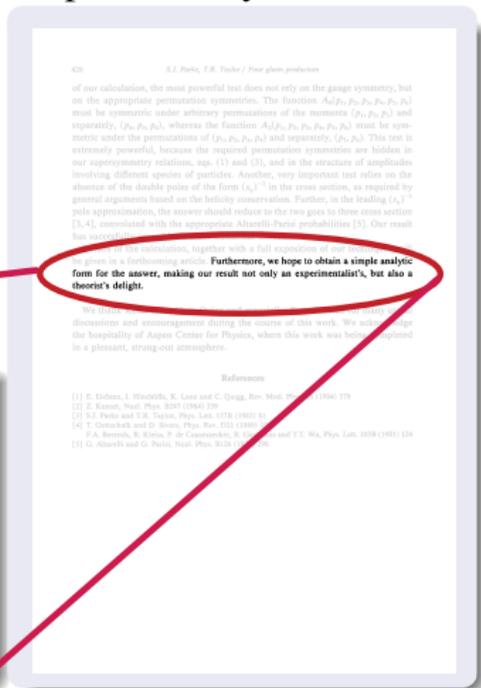
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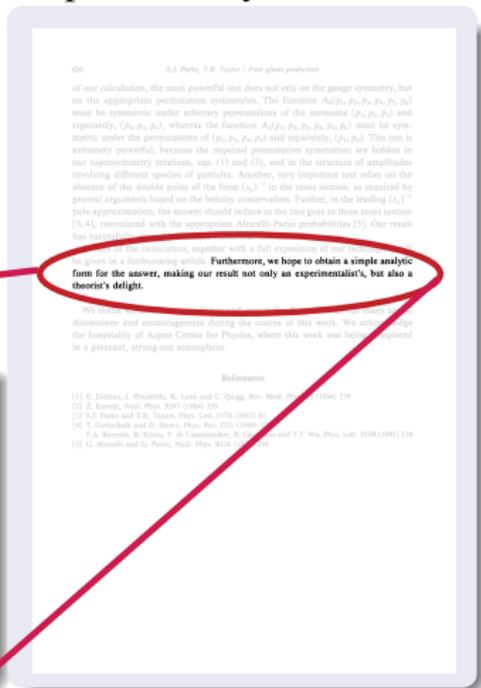
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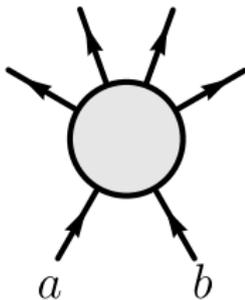
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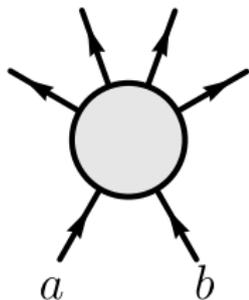
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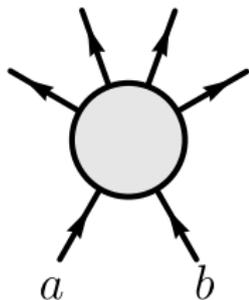
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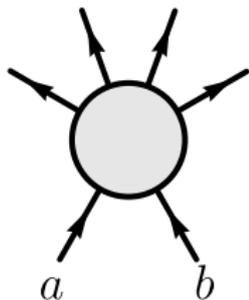
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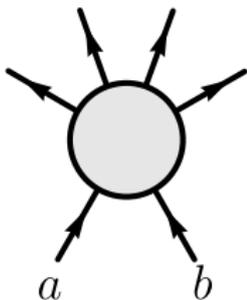
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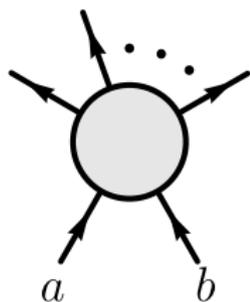
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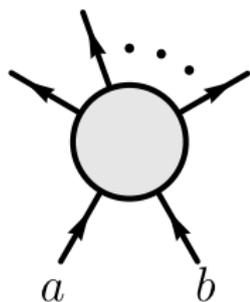
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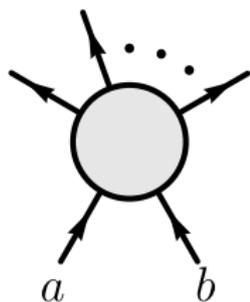
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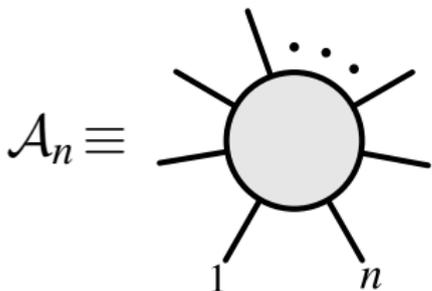
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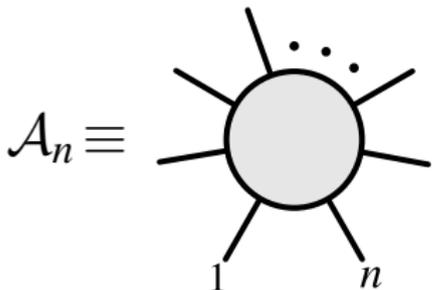
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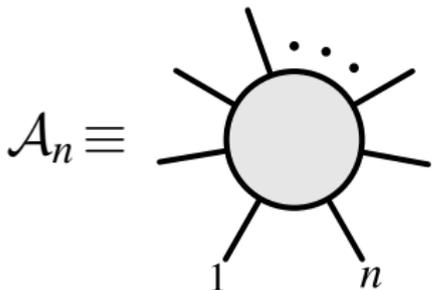
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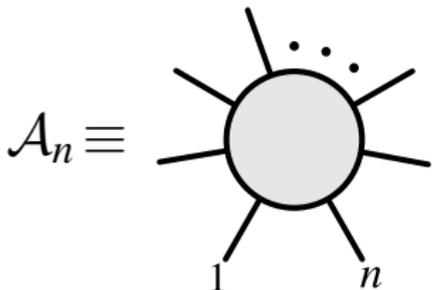
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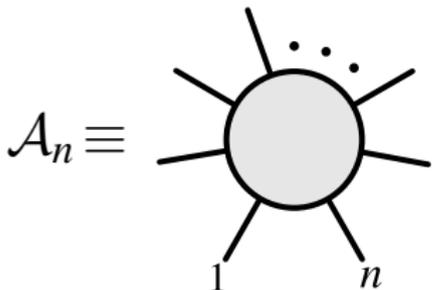


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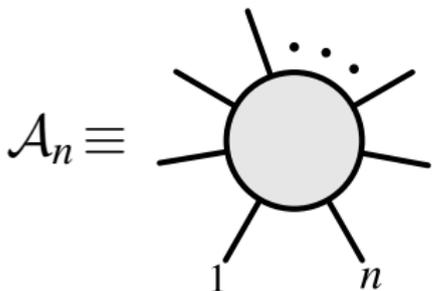


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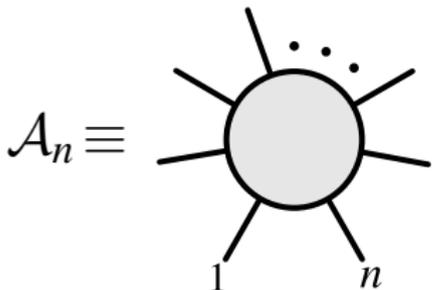


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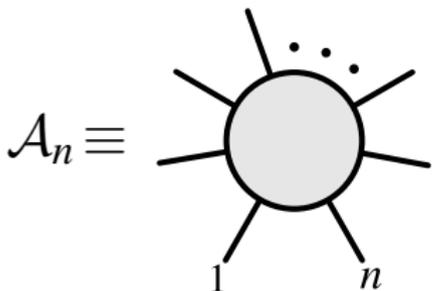


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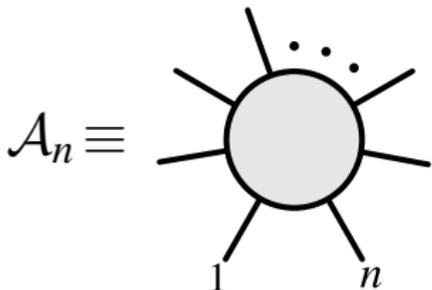


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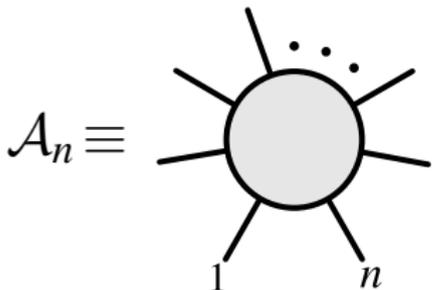


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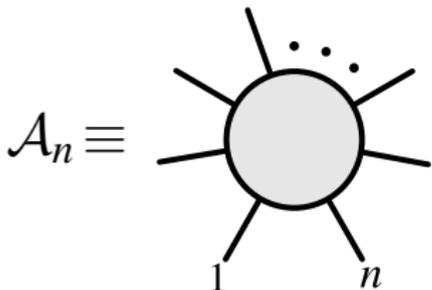


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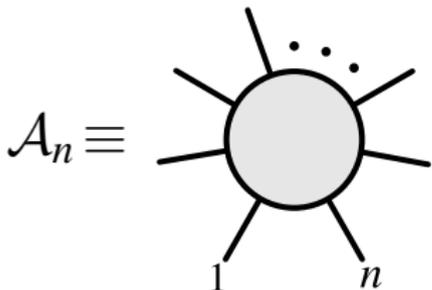


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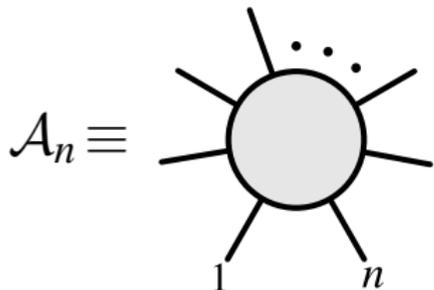


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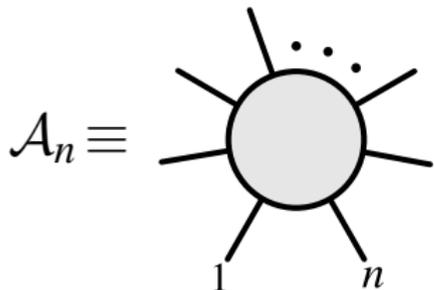


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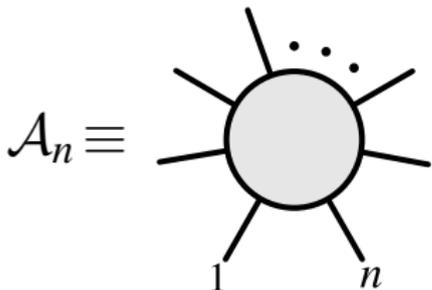


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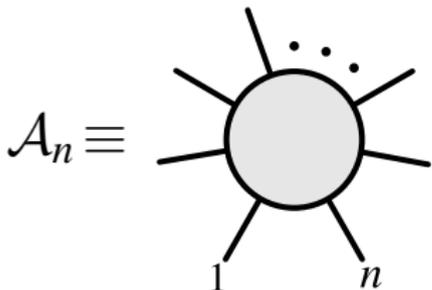
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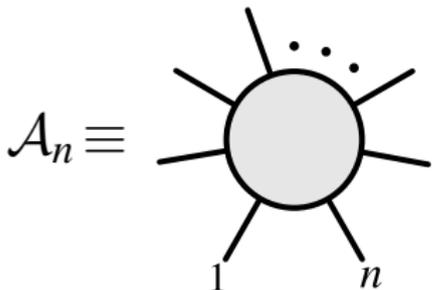
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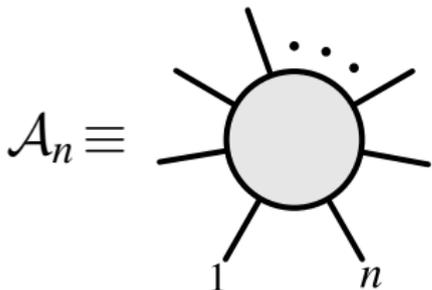
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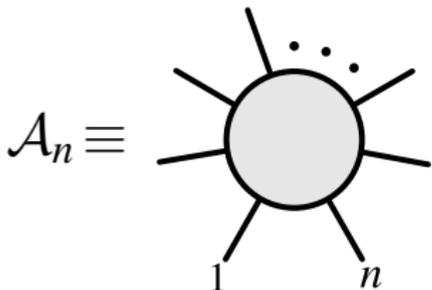
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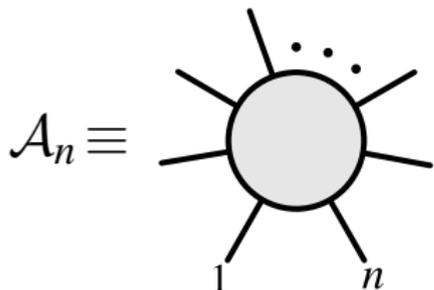
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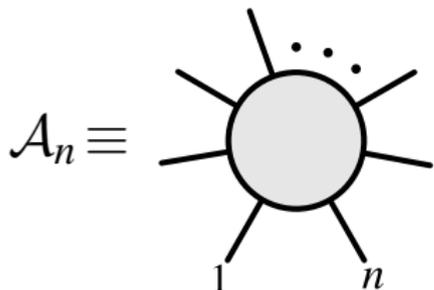
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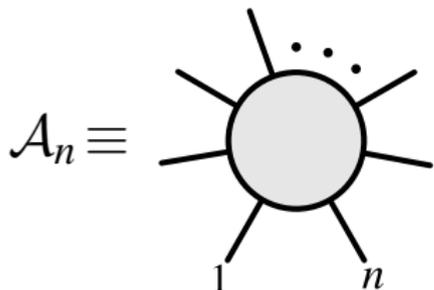
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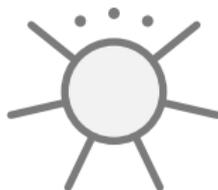
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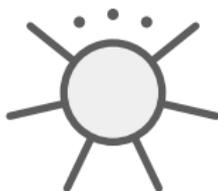
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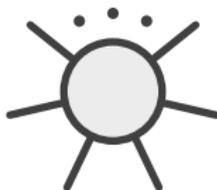
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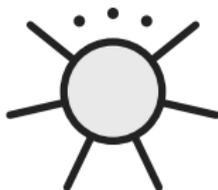
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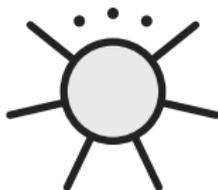
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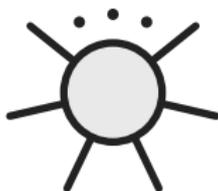
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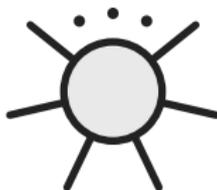
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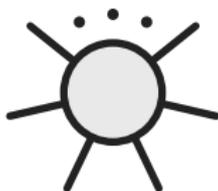
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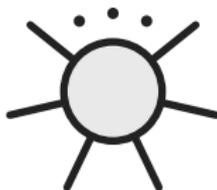
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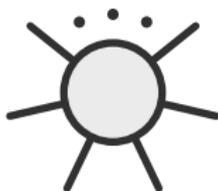
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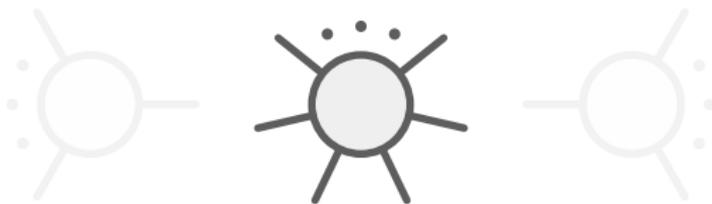
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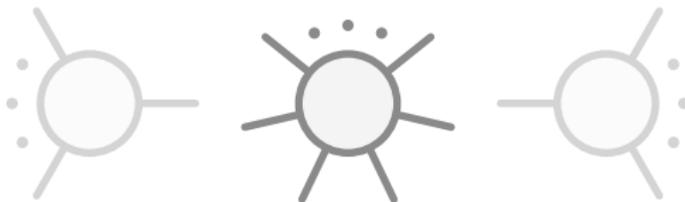
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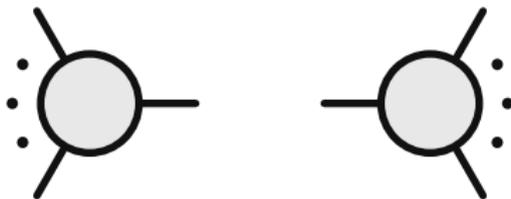
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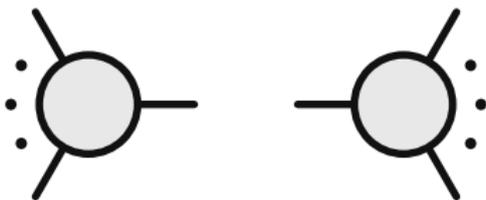
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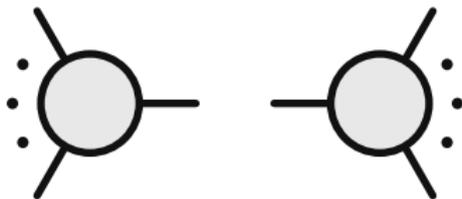
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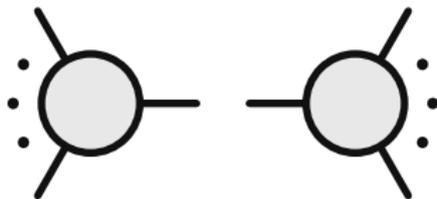
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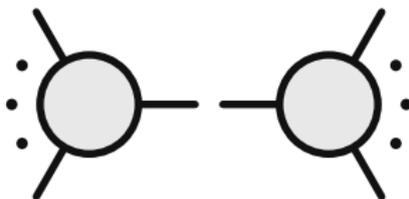
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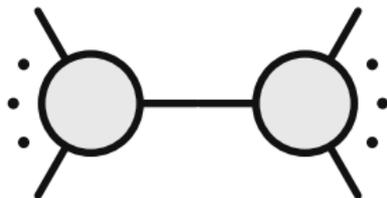
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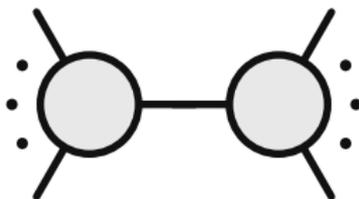
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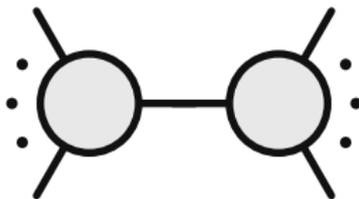
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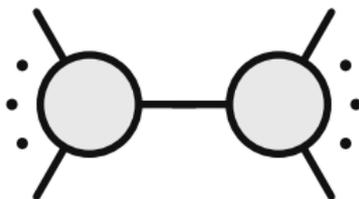
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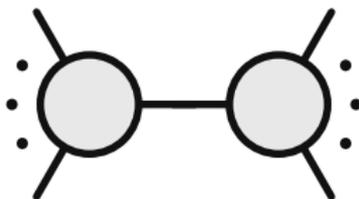
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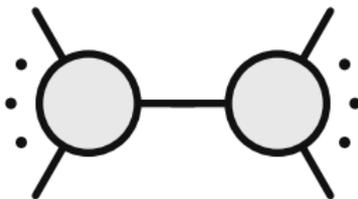
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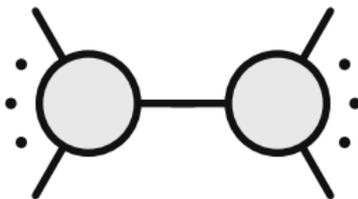
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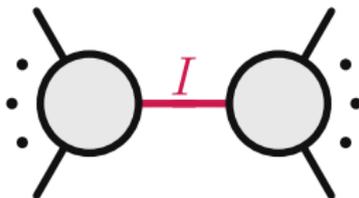
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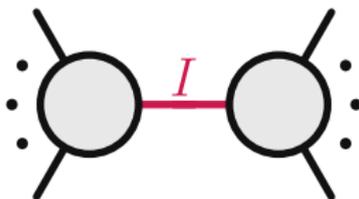
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**Internal Particles:**

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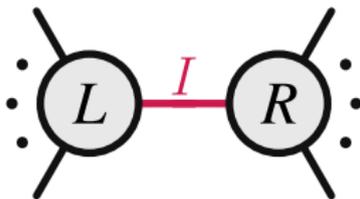
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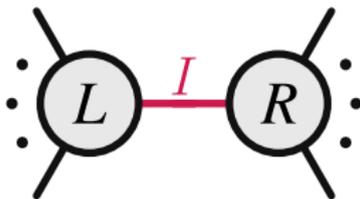


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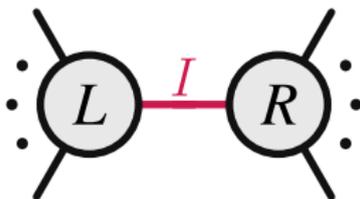


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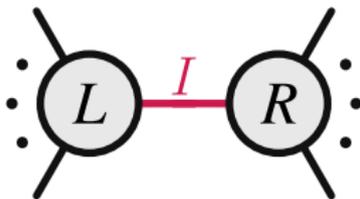


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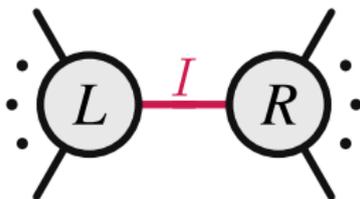


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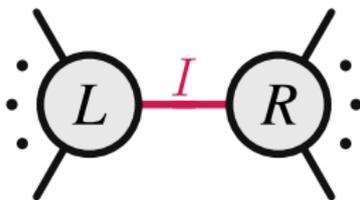


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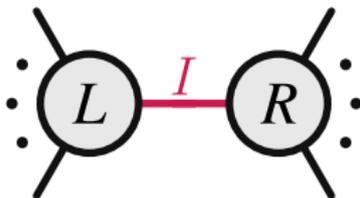


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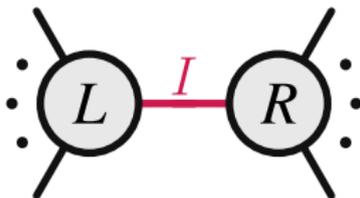


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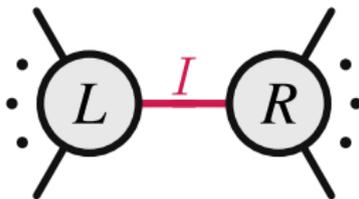


**Internal Particles:** **locality** dictates that we multiply each amplitude, and **unitarity** dictates that we marginalize over unobserved states—integrating over the Lorentz-invariant phase space (“LIPS”) for each particle *I*, and summing over the possible states (helicities, masses, colours, etc.).

$$\sum_{\text{states } I} \int d^3\text{LIPS}_I \mathcal{A}_L(\dots, I) \times \mathcal{A}_R(I, \dots)$$

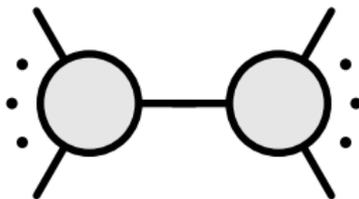
## Broadening the Class of Physically Meaningful Functions

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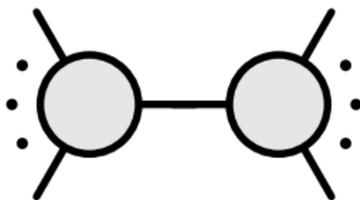
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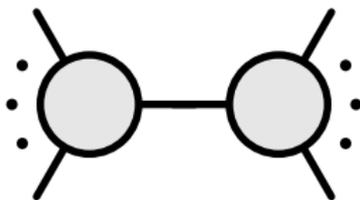
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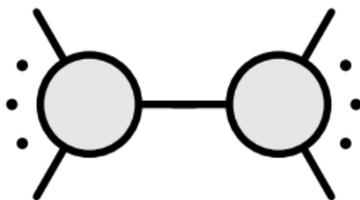


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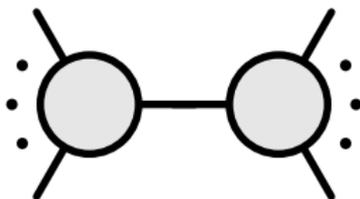


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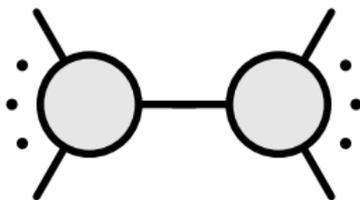
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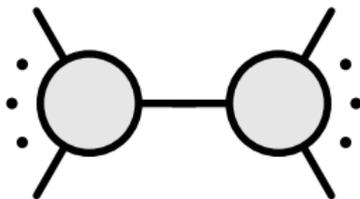
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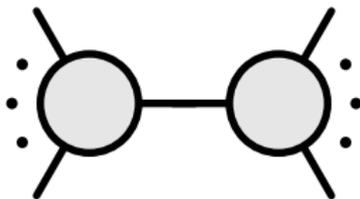
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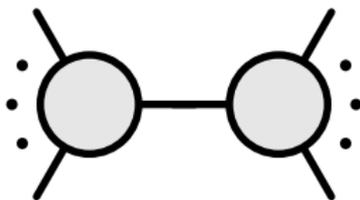
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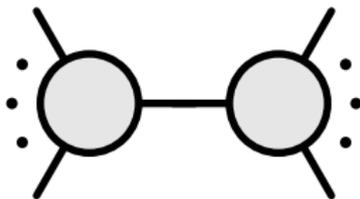
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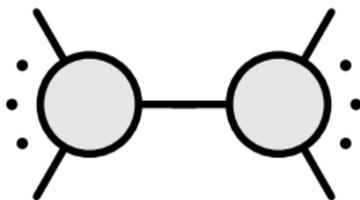
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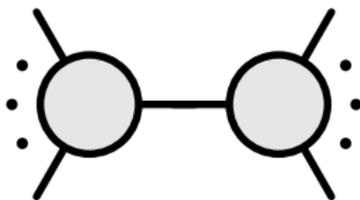
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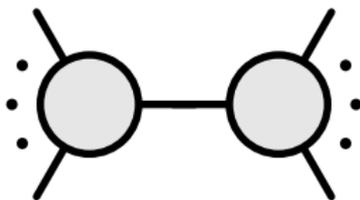
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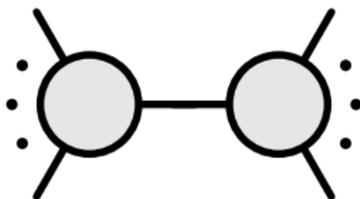
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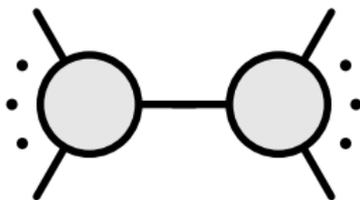
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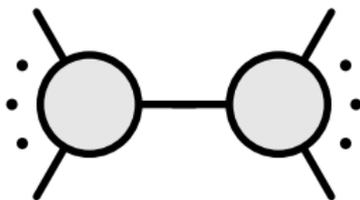
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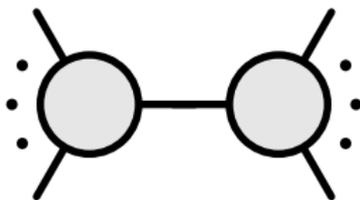
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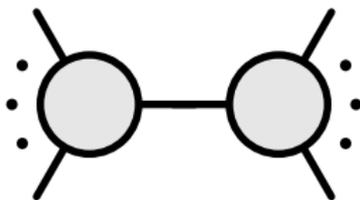
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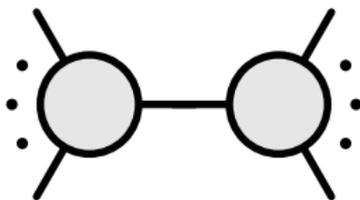
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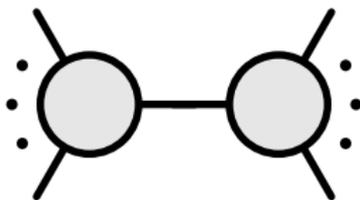
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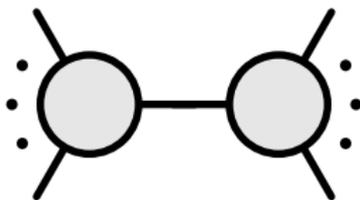
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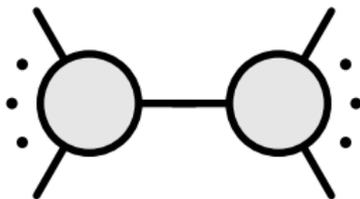
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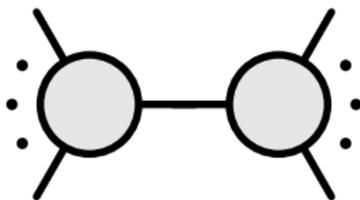
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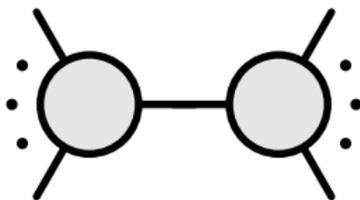
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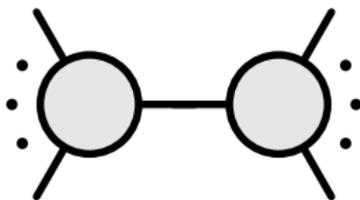


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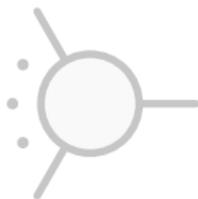
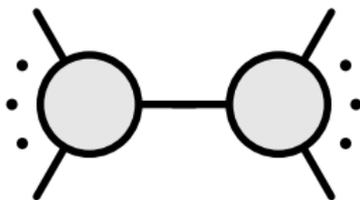


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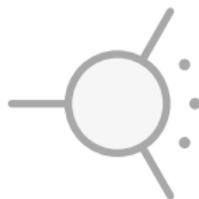
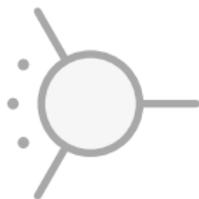
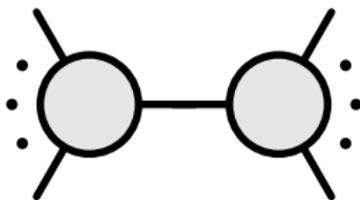
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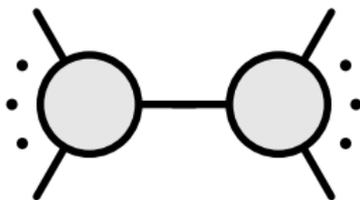


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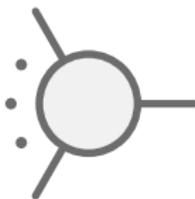
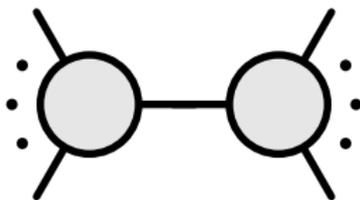


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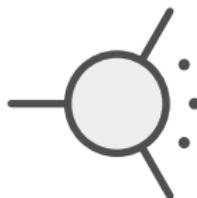
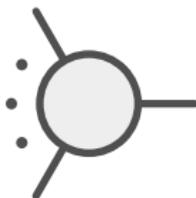
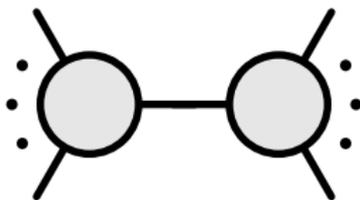


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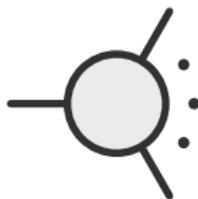
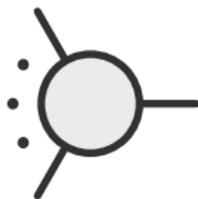
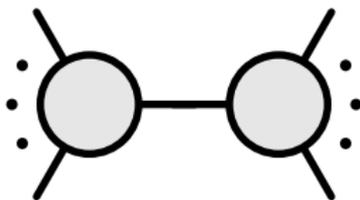


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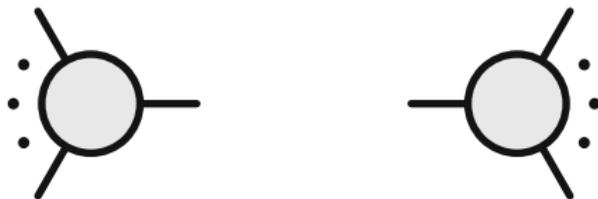
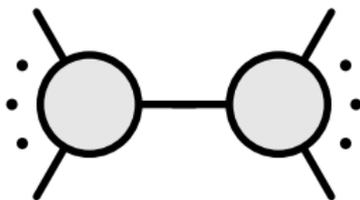


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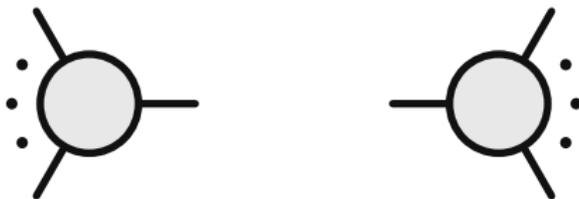
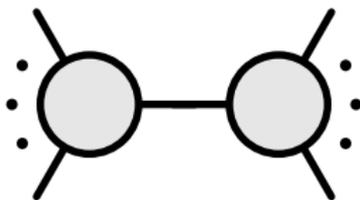


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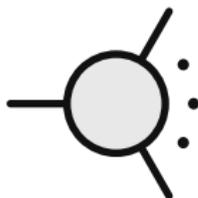
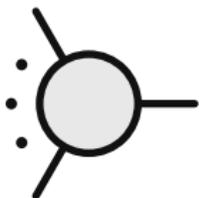
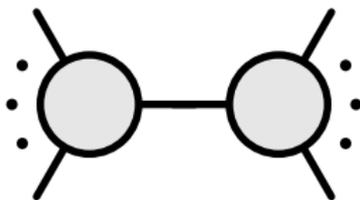
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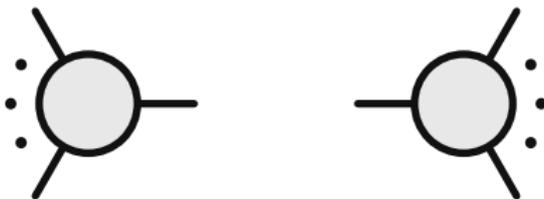
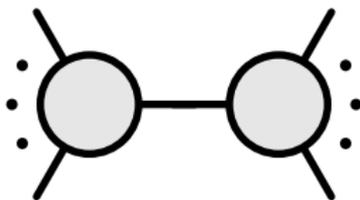
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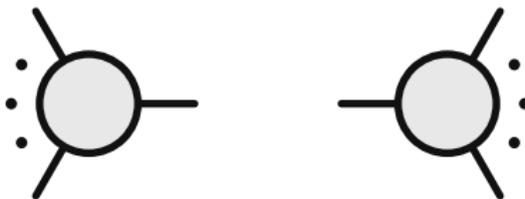
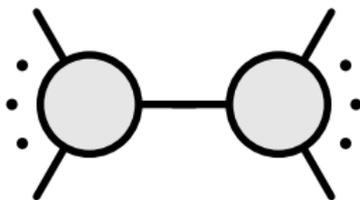


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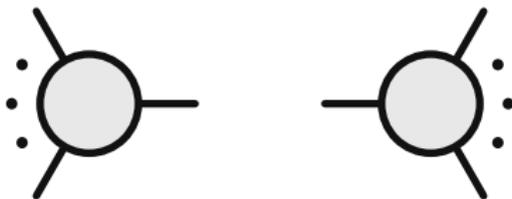
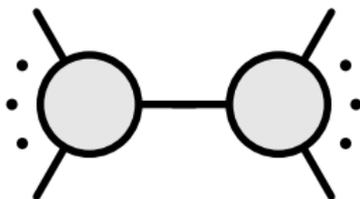
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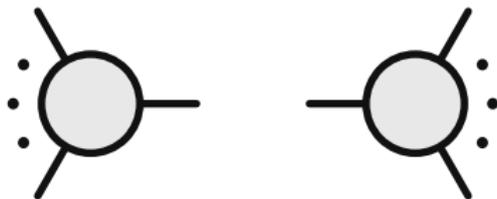
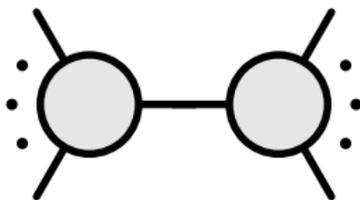


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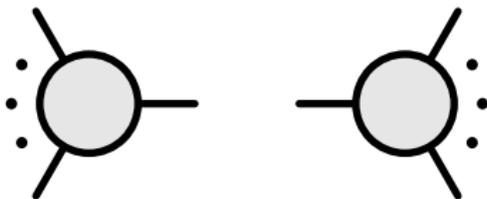
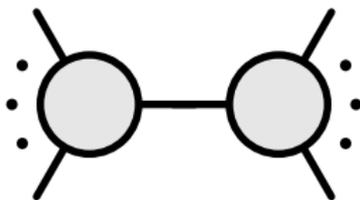
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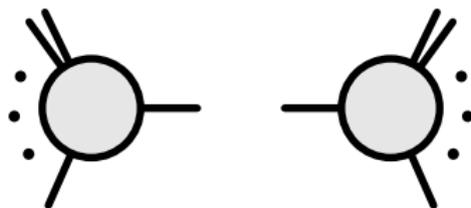
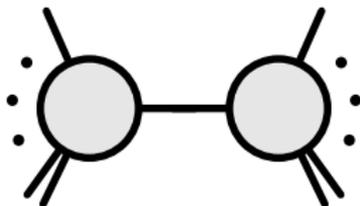


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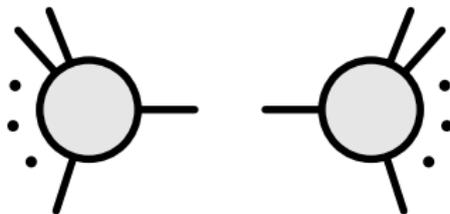
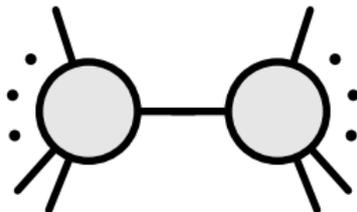


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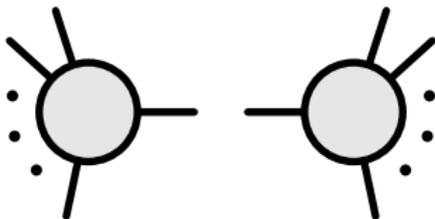
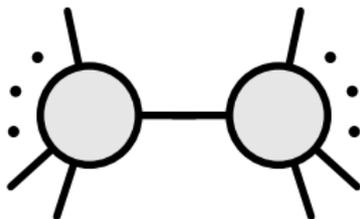
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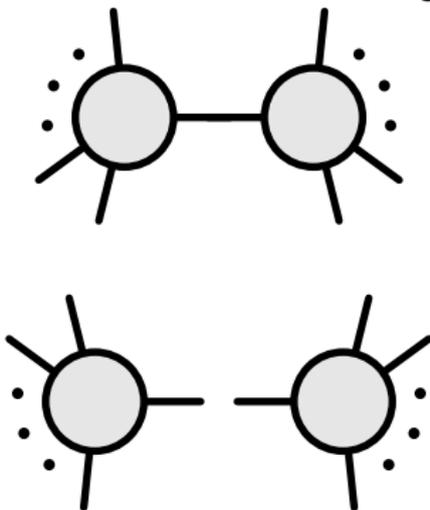
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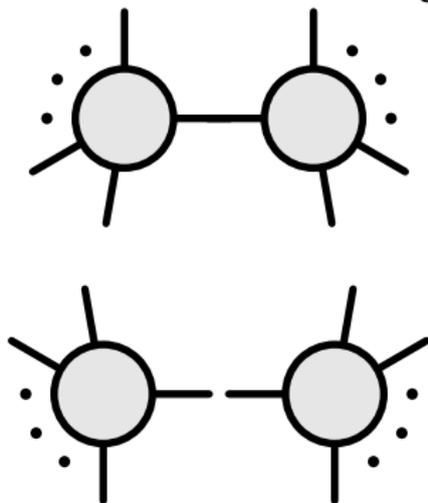


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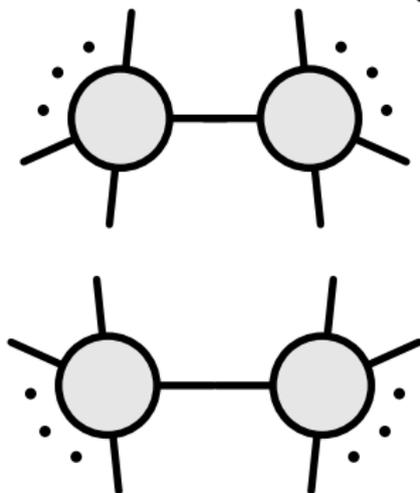


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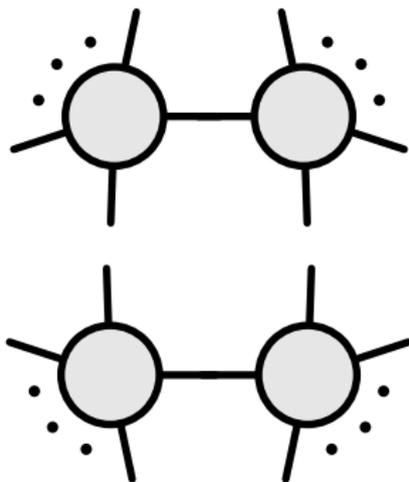


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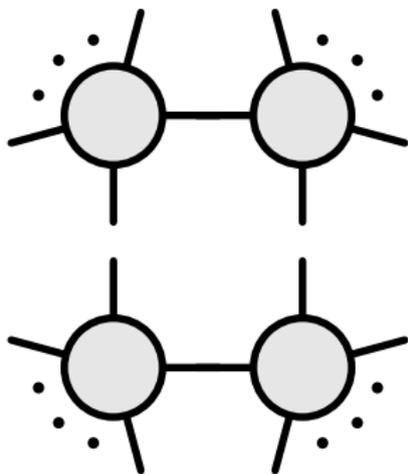
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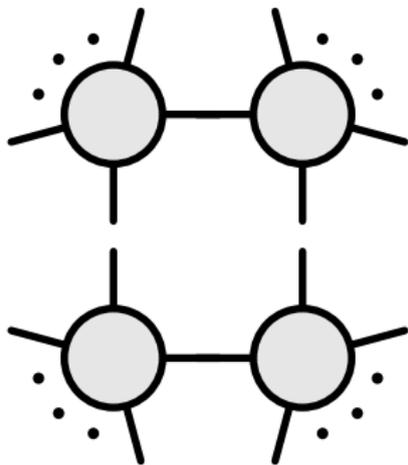


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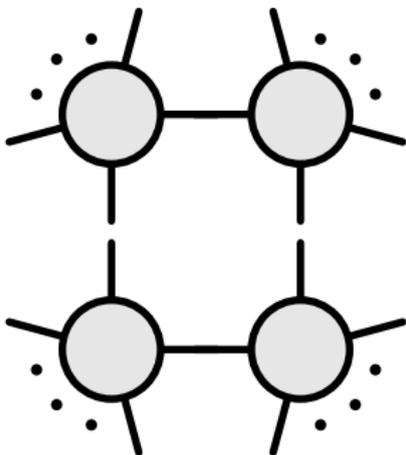


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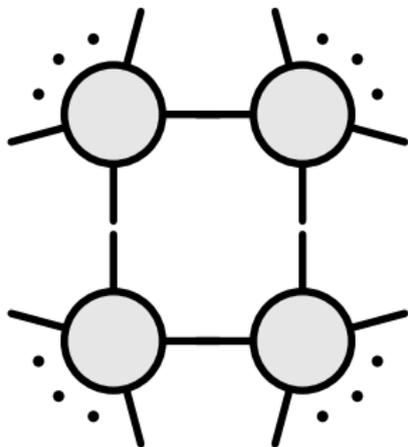


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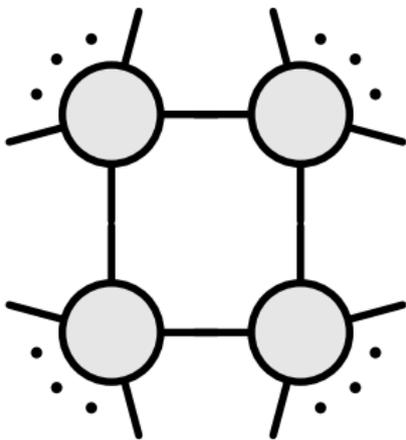


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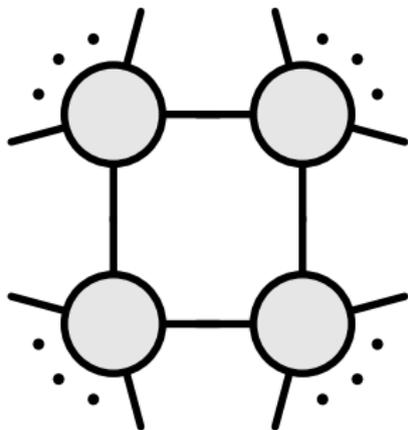


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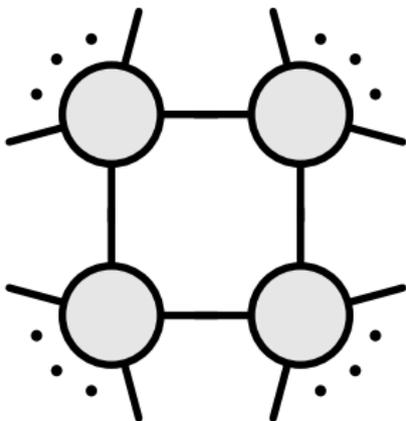


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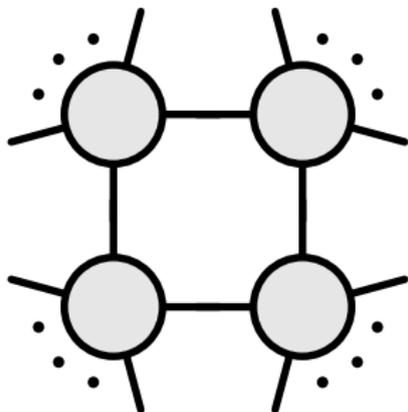


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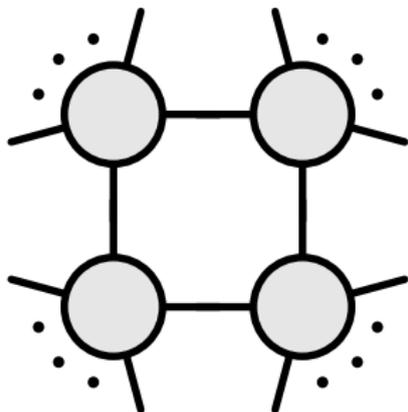


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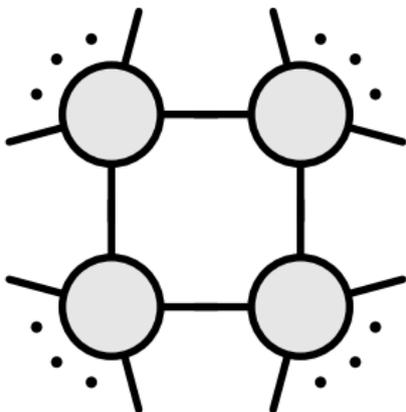


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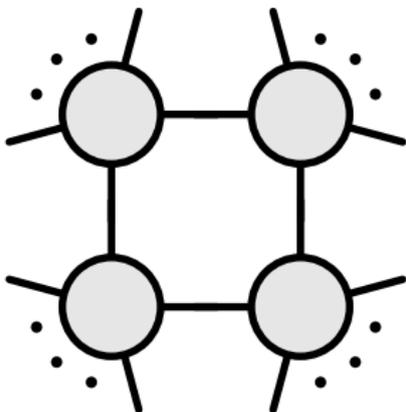


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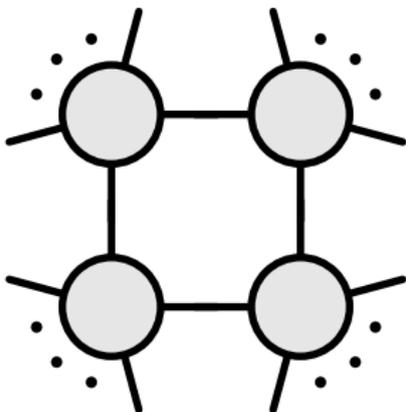


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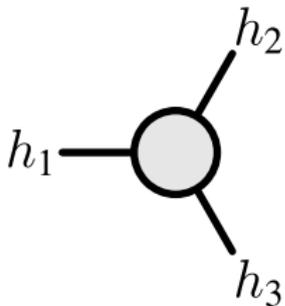
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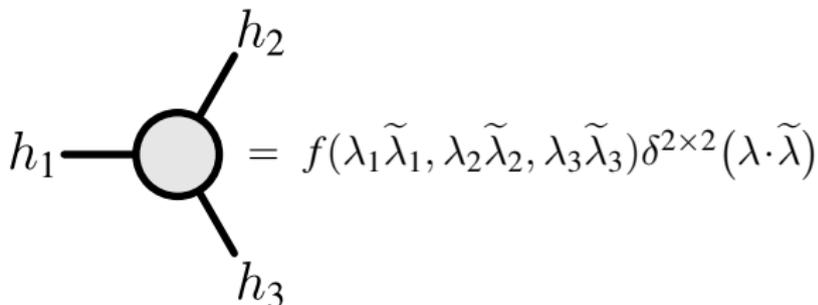
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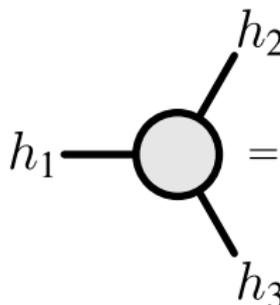
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$$h_1 \text{---} \bigcirc \begin{matrix} \nearrow h_2 \\ \searrow h_3 \end{matrix} = f(\lambda_1 \tilde{\lambda}_1, \lambda_2 \tilde{\lambda}_2, \lambda_3 \tilde{\lambda}_3) \delta^{2 \times 2}(\lambda \cdot \tilde{\lambda})$$

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A diagram showing a central grey circle representing a vertex. Three lines extend from the circle: one to the left labeled  $h_1$ , one to the top-right labeled  $h_2$ , and one to the bottom-right labeled  $h_3$ .

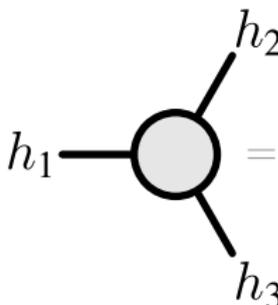
$$= f(\lambda_1 \tilde{\lambda}_1, \lambda_2 \tilde{\lambda}_2, \lambda_3 \tilde{\lambda}_3) \delta^{2 \times 2}(\lambda \cdot \tilde{\lambda})$$

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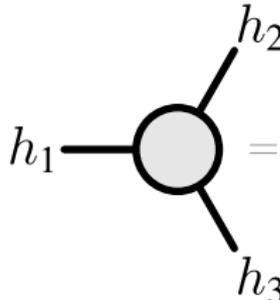
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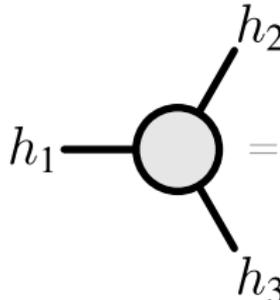
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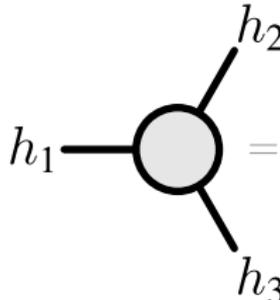
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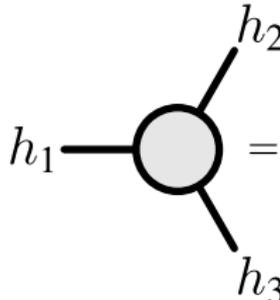
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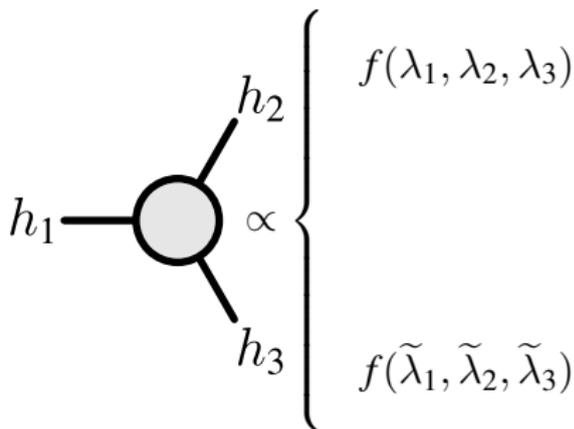
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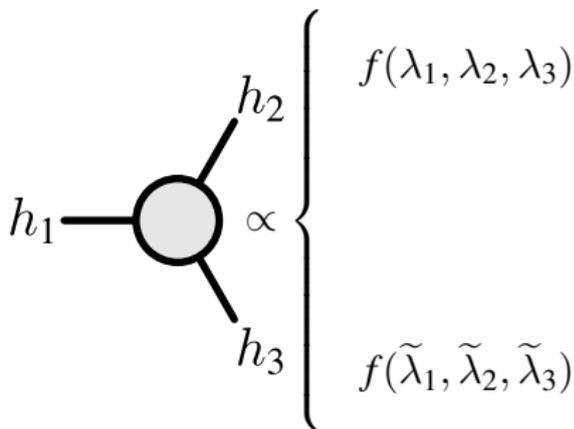
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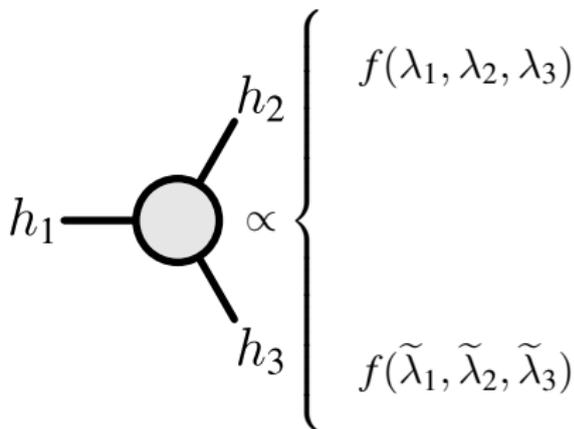
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$$h_1 + h_2 + h_3 \leq 0$$

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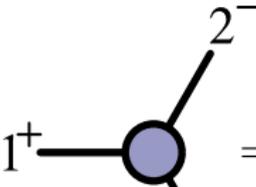
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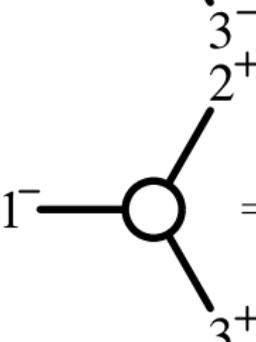
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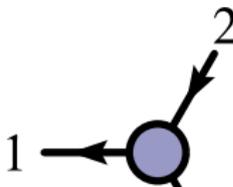
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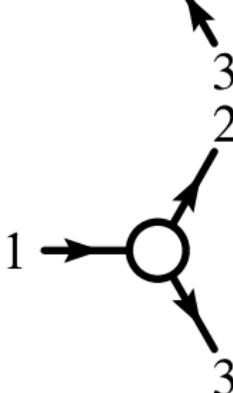
$$\begin{aligned}
 & \text{Diagram 1: } 1^+ \text{ (in), } 2^- \text{ (out), } 3^- \text{ (out), } 2^+ \text{ (out)} \\
 & \qquad = \frac{\langle 23 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 31 \rangle} \delta^{2 \times 2}(\lambda \cdot \tilde{\lambda}) \\
 & \text{Diagram 2: } 1^- \text{ (in), } 3^+ \text{ (out), } 2^+ \text{ (out)} \\
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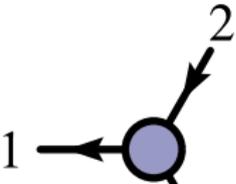
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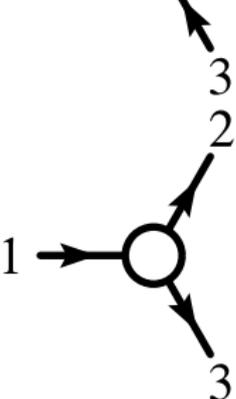
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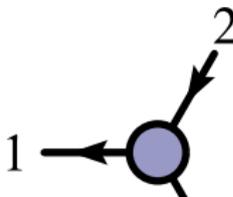
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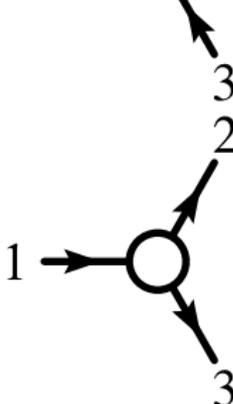
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 & \text{Diagram 2 (White Vertex)} = \frac{[3 1] [2 3]^3}{[1 2] [2 3] [3 1]} \delta^{2 \times 2}(\lambda \cdot \tilde{\lambda}) \equiv \mathcal{A}_3\left(-\frac{1}{2}, +\frac{1}{2}, +\right)
 \end{aligned}$$

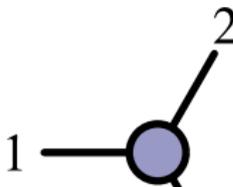
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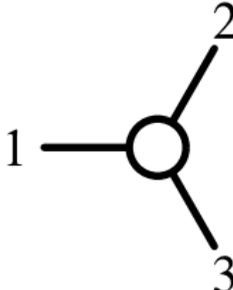
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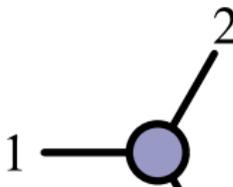
$$= \frac{\delta^{2 \times 4}(\lambda \cdot \tilde{\eta})}{\langle 12 \rangle \langle 23 \rangle \langle 31 \rangle} \delta^{2 \times 2}(\lambda \cdot \tilde{\lambda}) \equiv \mathcal{A}_3^{(2)}$$



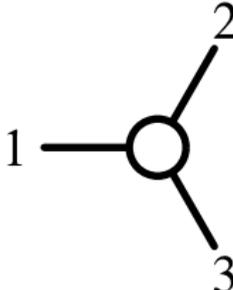
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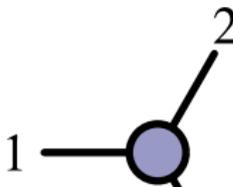
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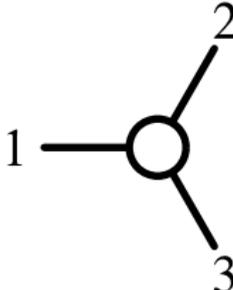
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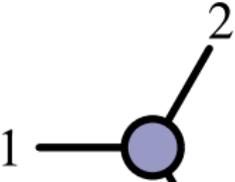
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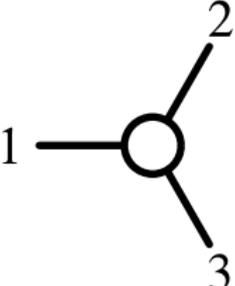
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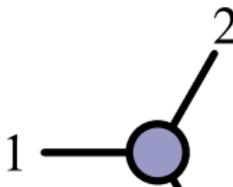
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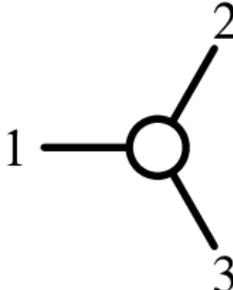
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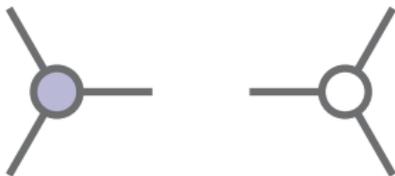
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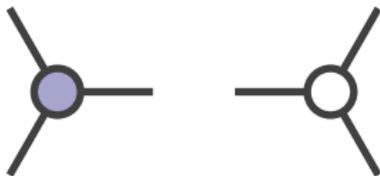
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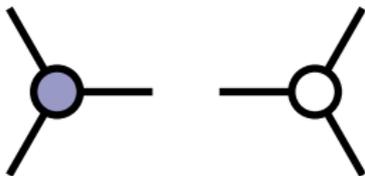
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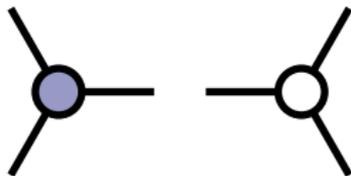
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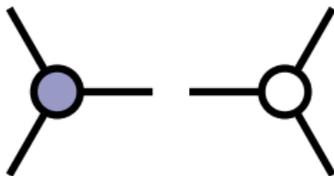
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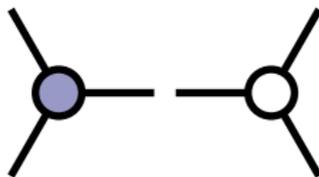
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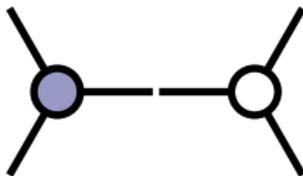
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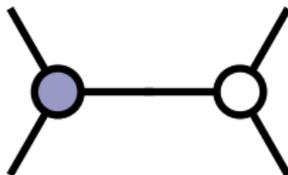
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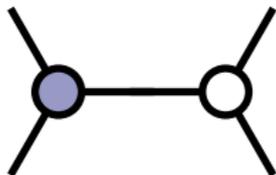
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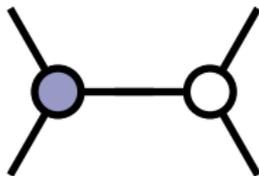
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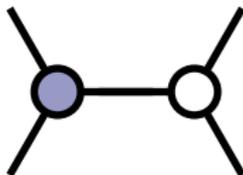
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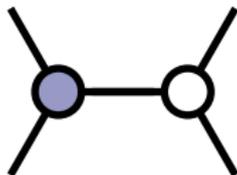
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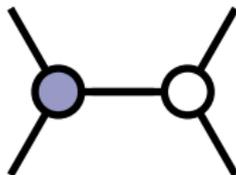
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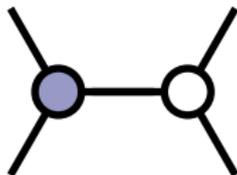
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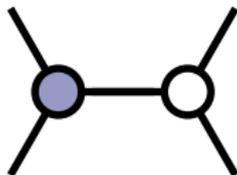
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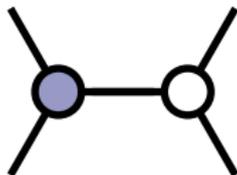
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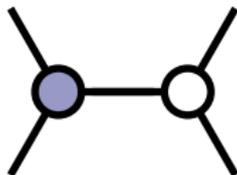
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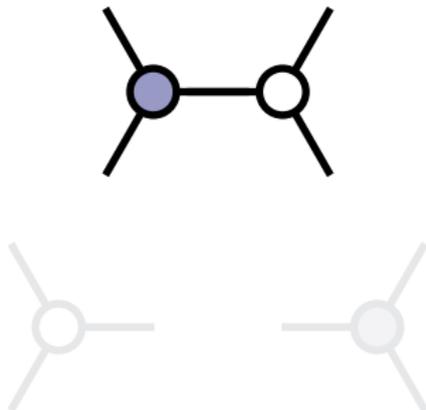
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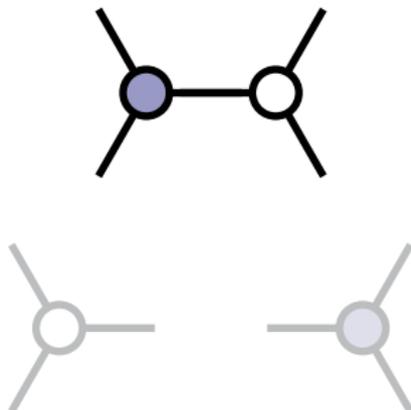
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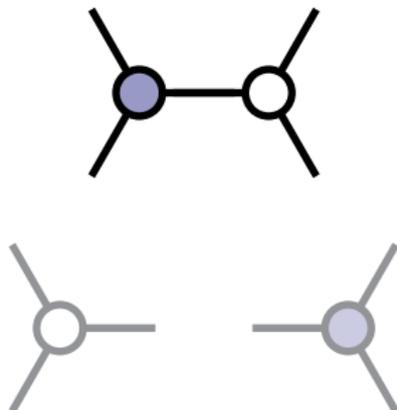
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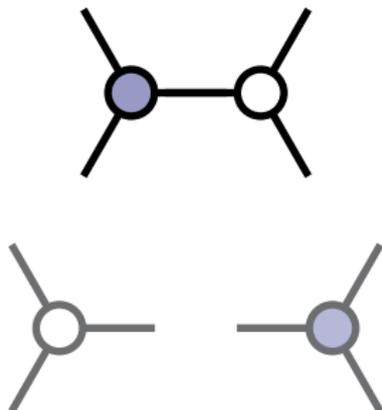
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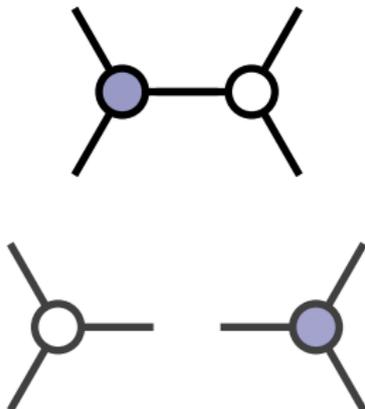
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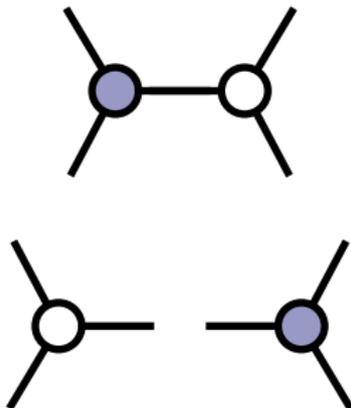
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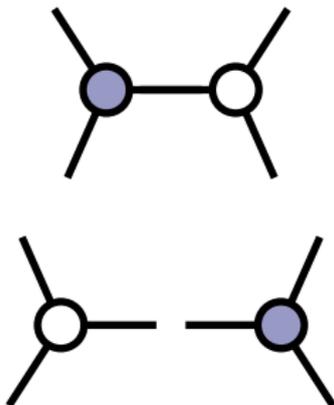
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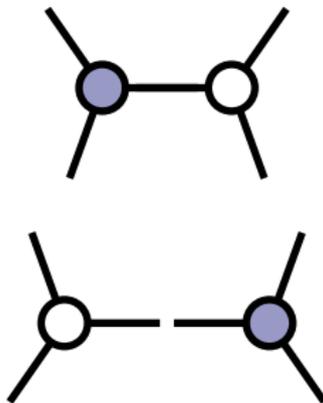
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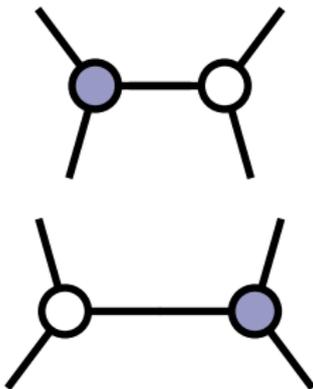
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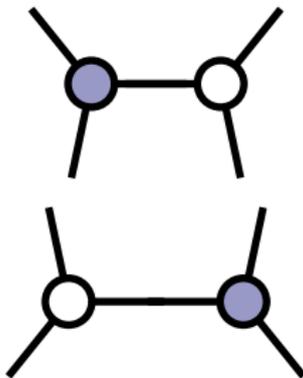
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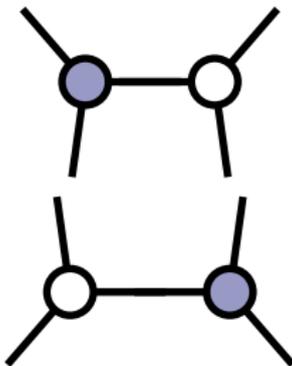
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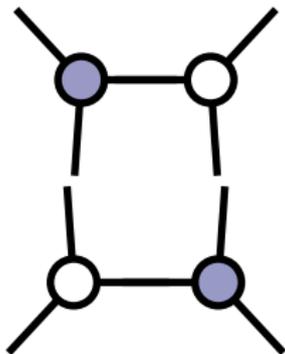
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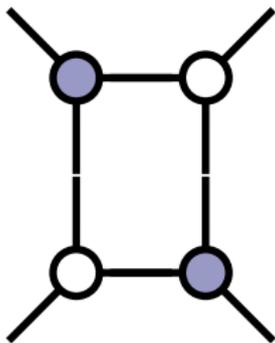
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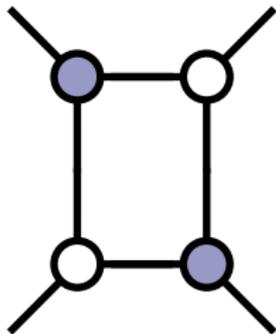
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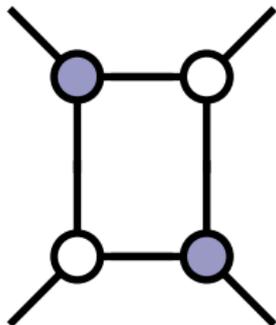
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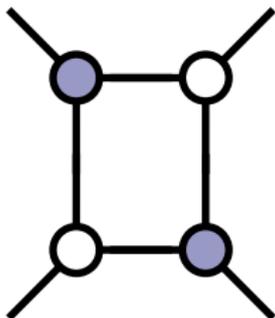
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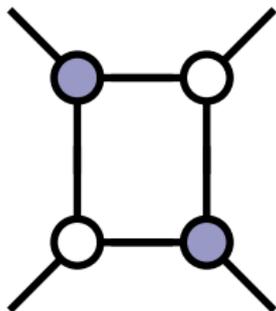
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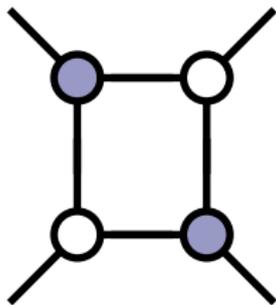
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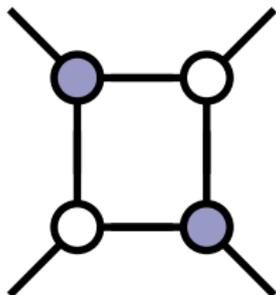
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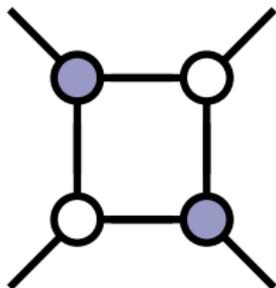
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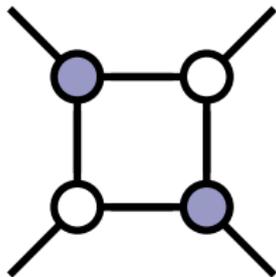
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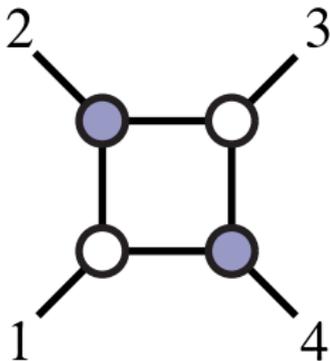
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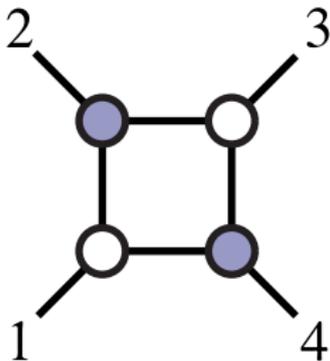
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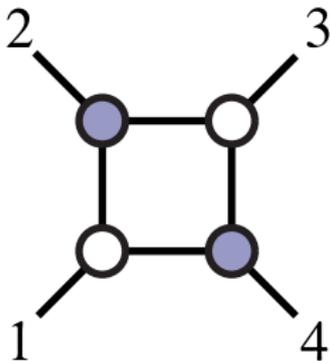
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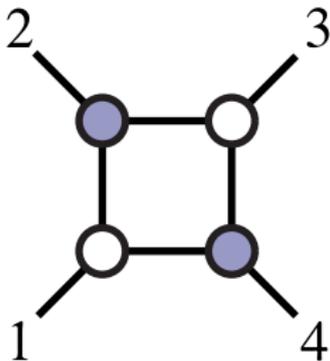
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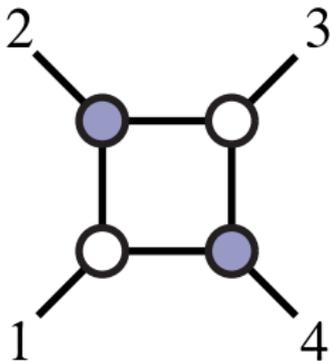
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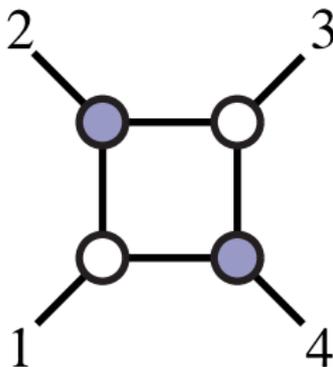
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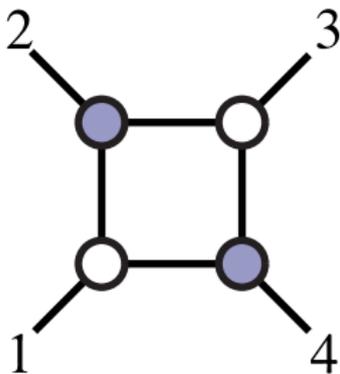
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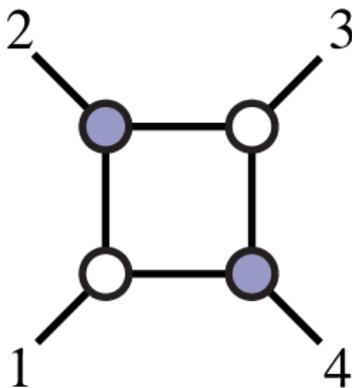
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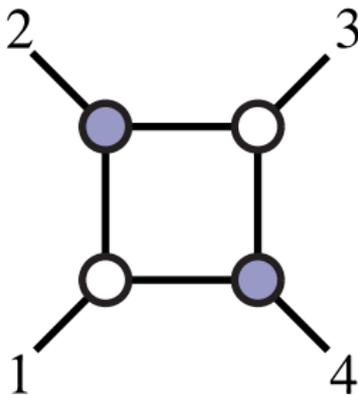
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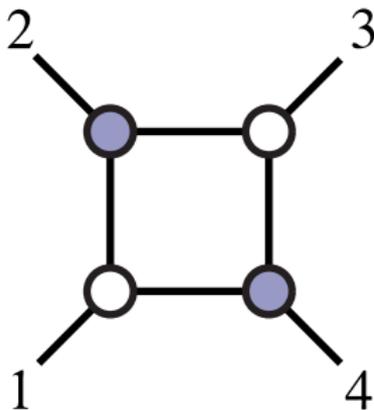
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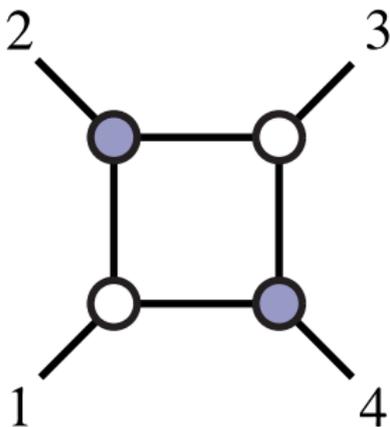
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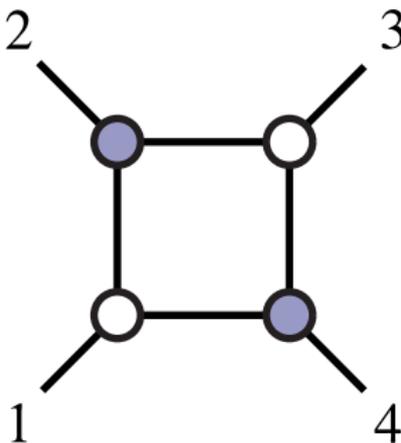
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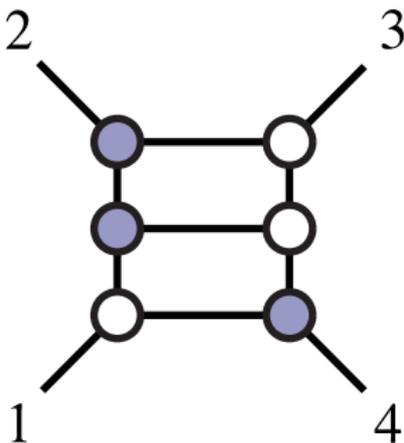
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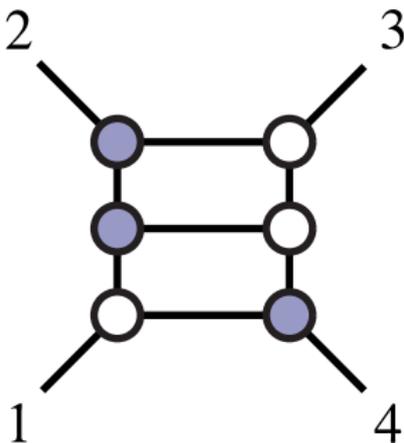
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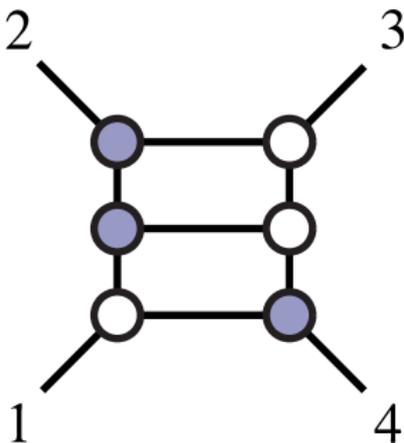
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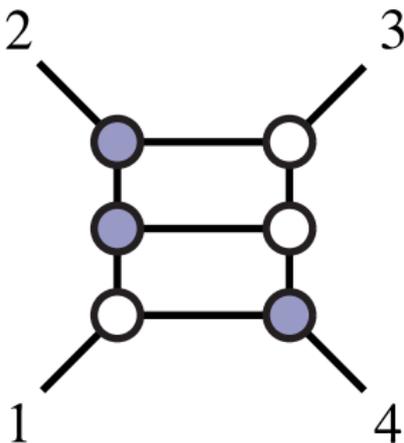
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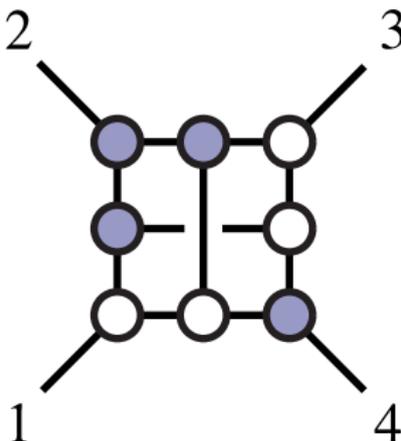
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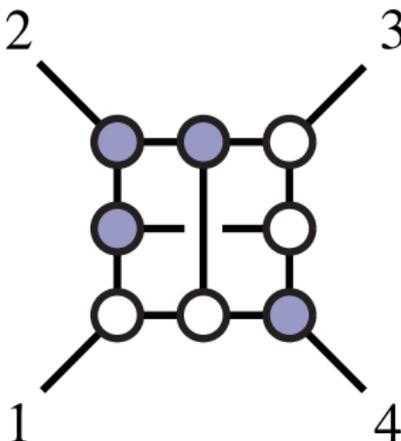
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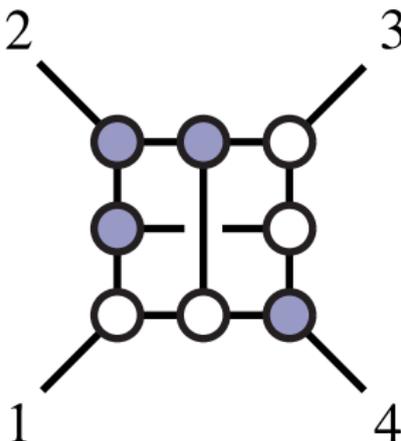
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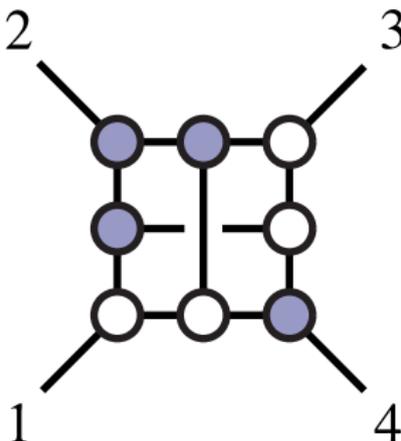
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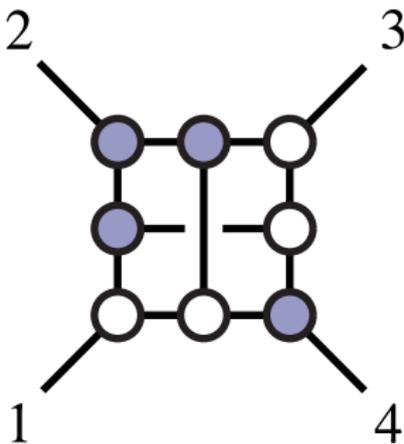
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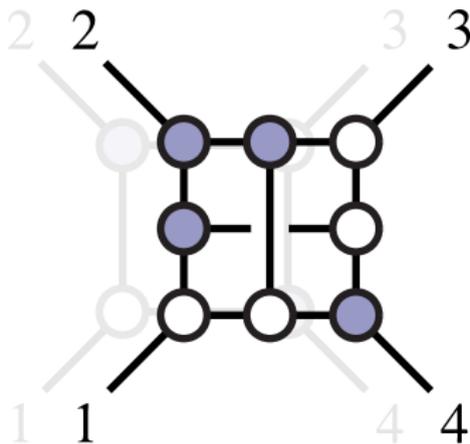
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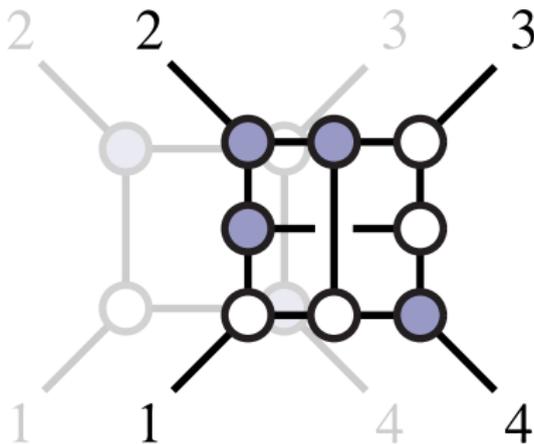
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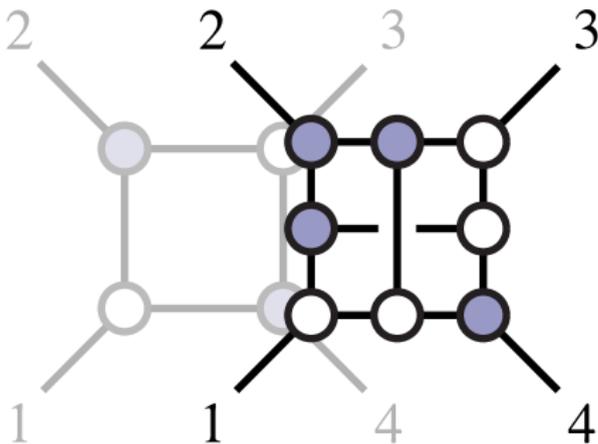
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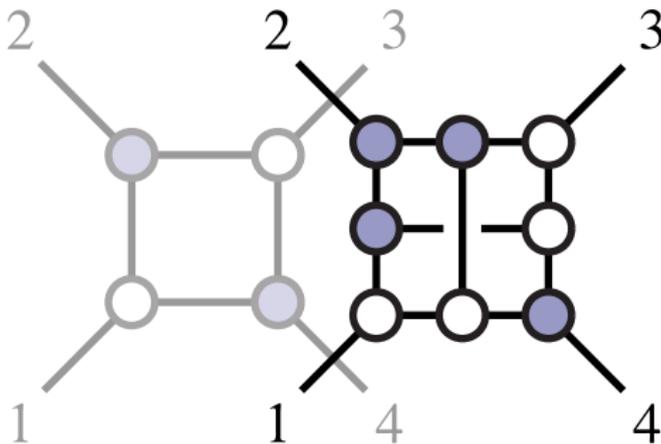
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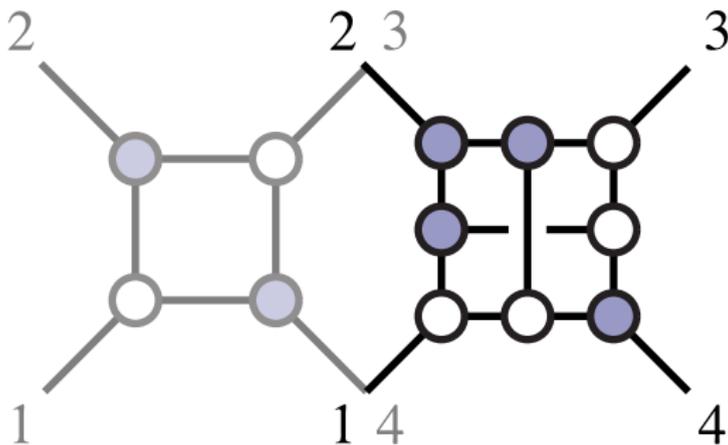
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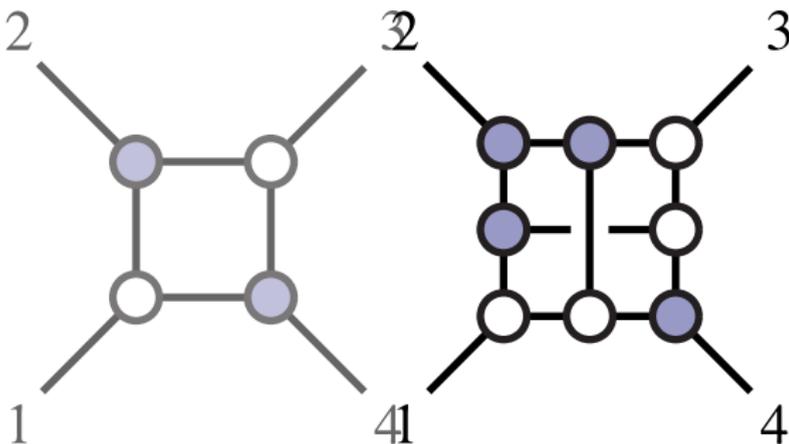
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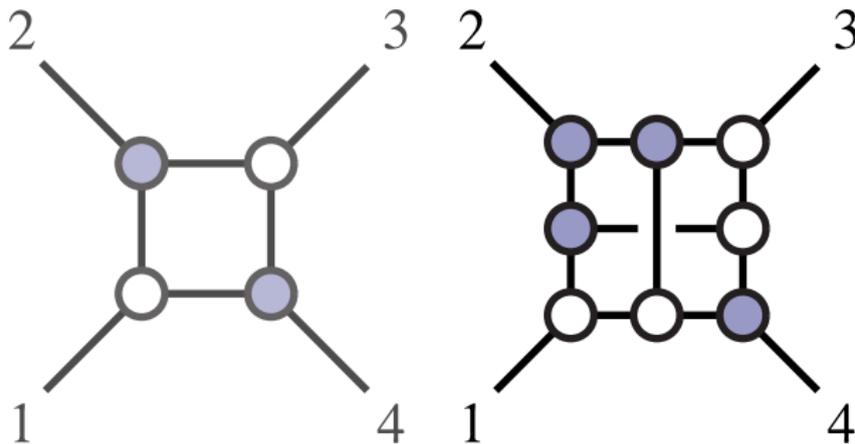
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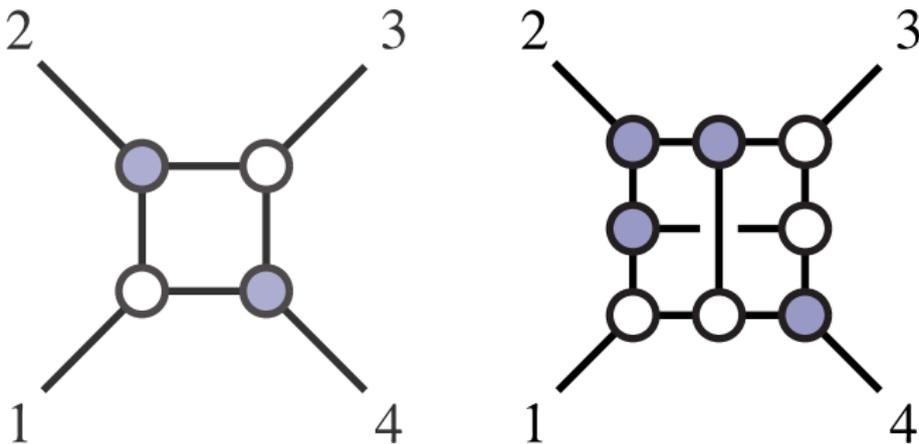
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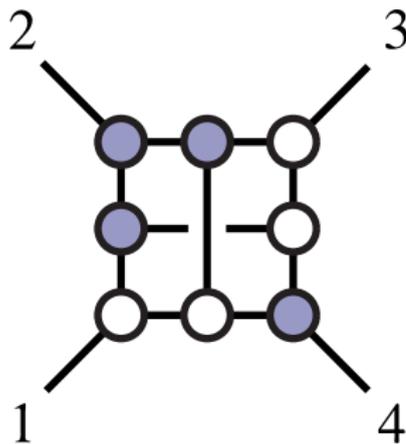
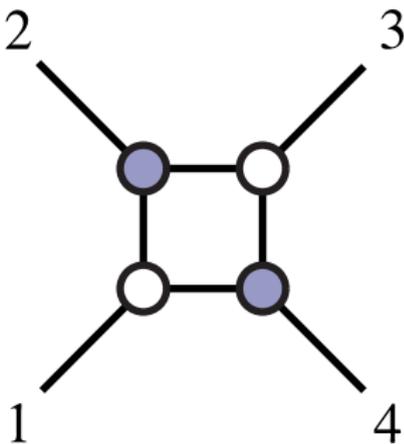
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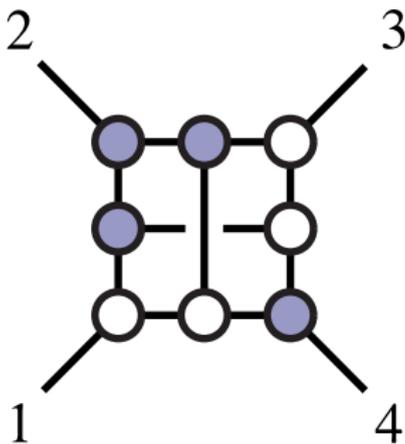
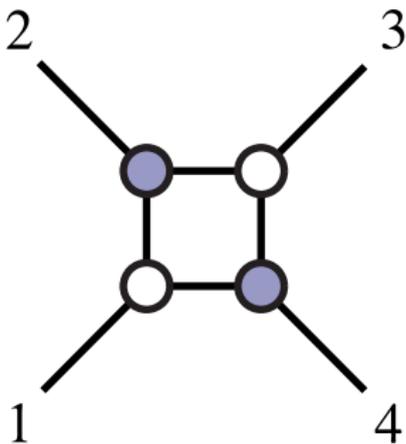
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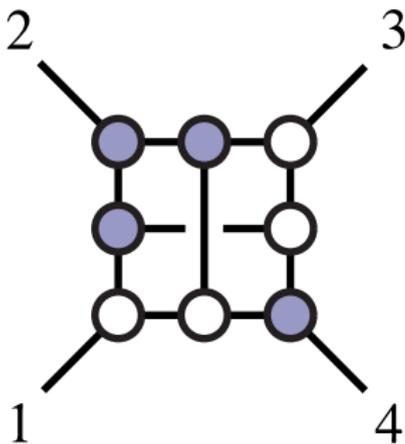
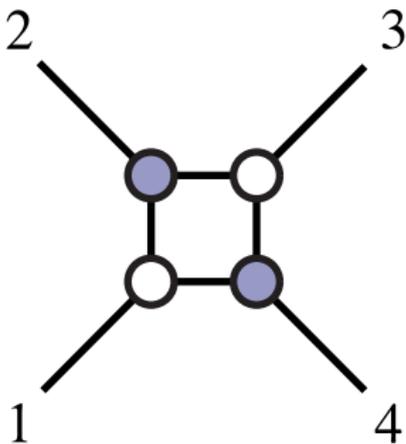
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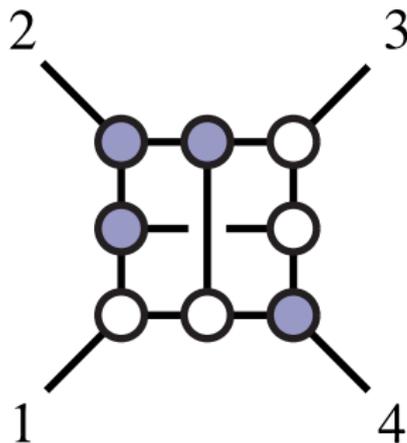
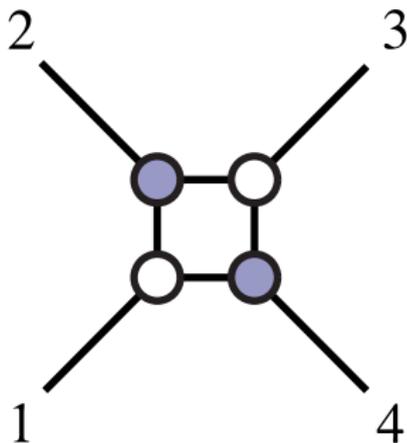
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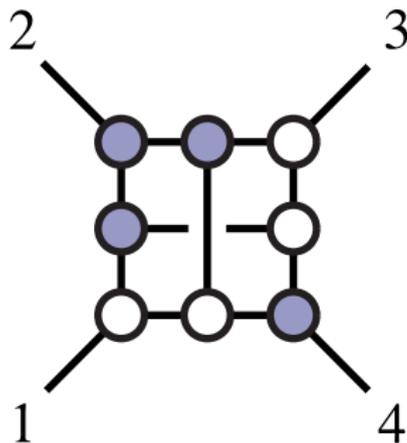
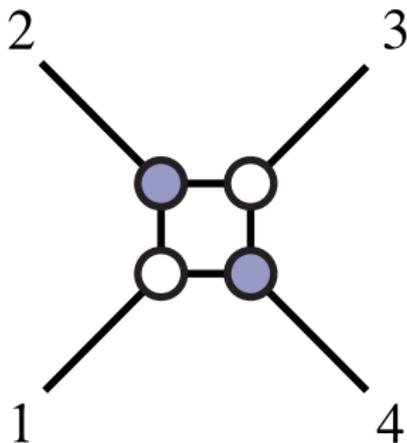
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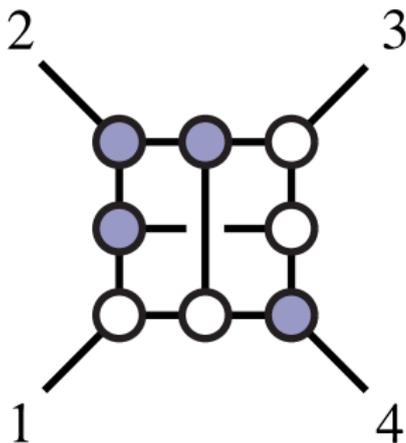
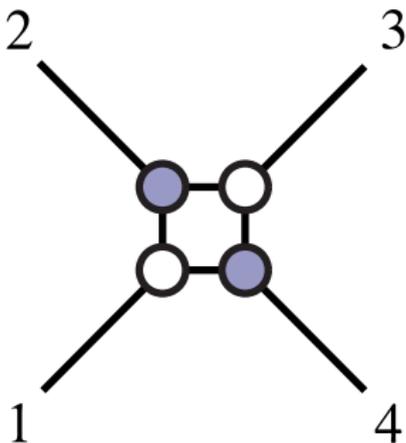
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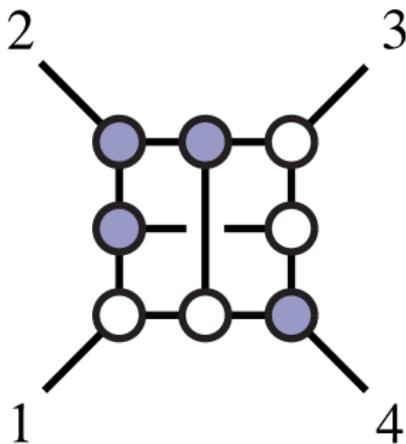
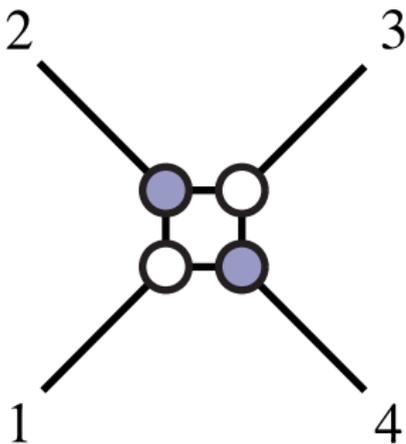
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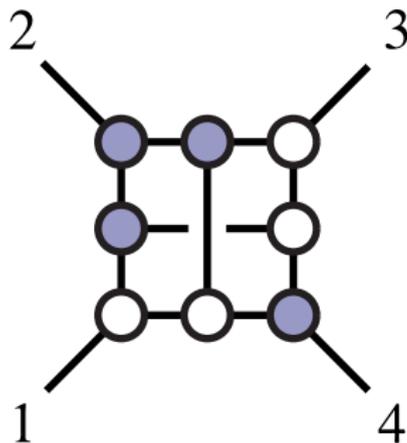
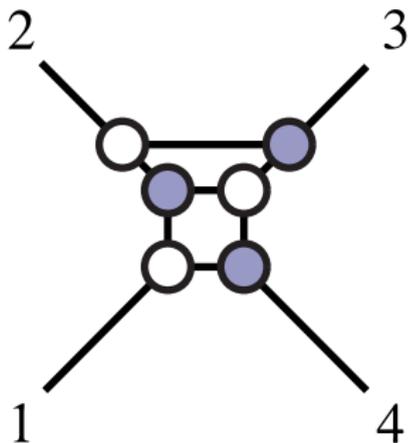
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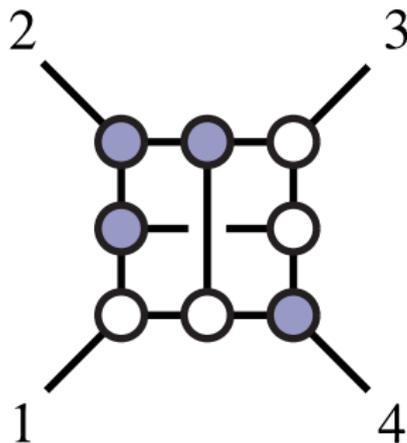
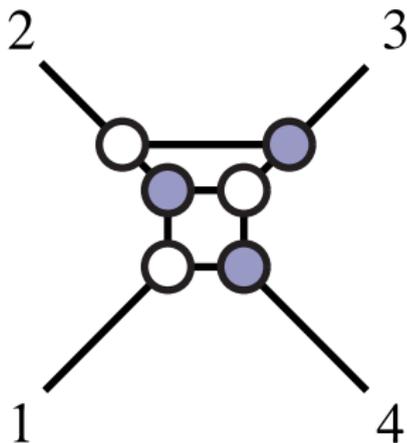
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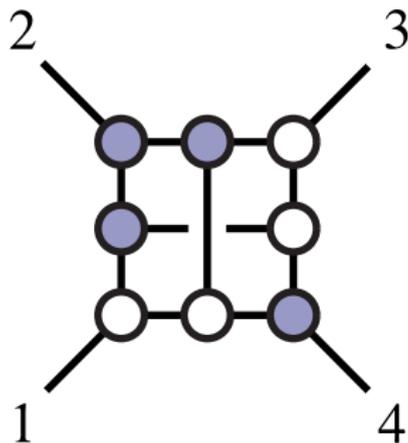
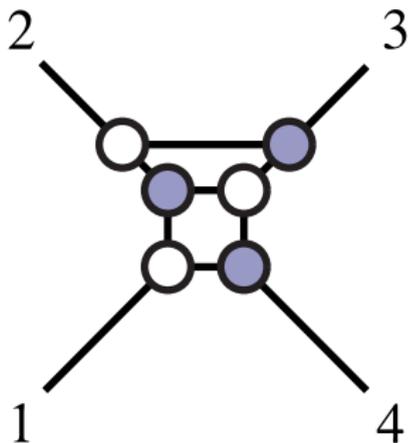
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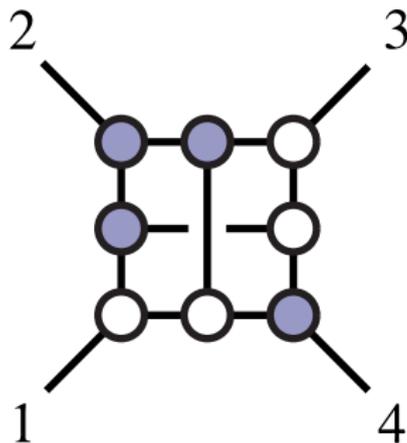
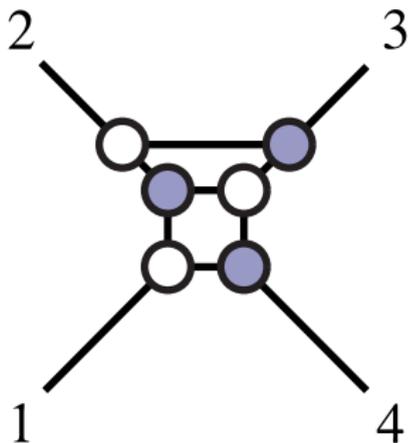
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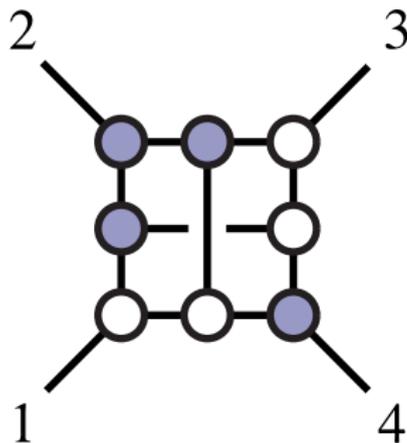
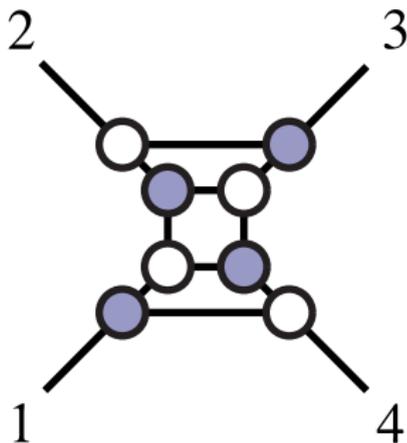
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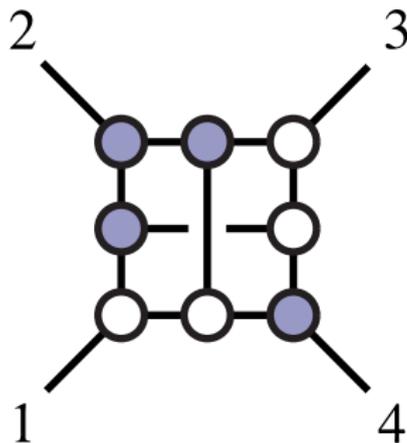
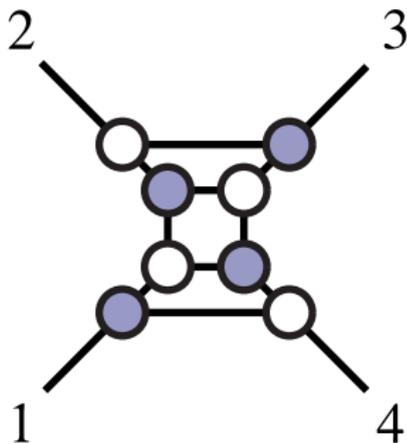
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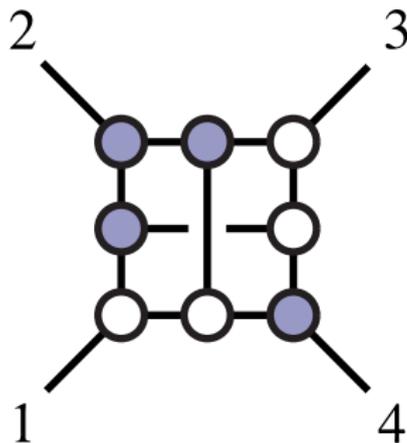
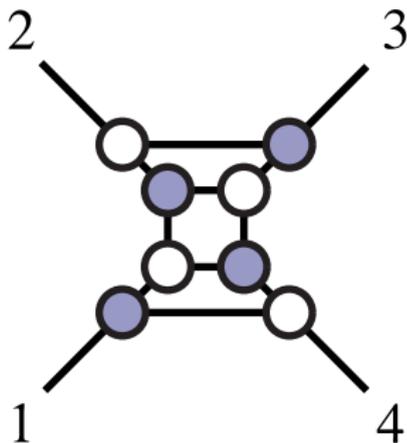
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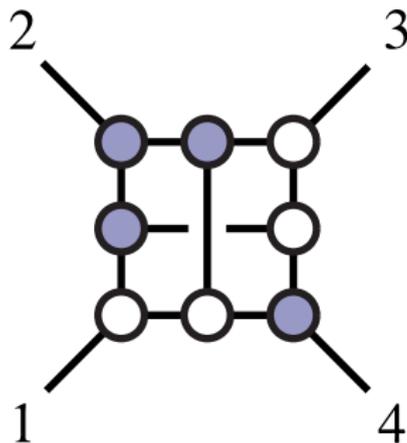
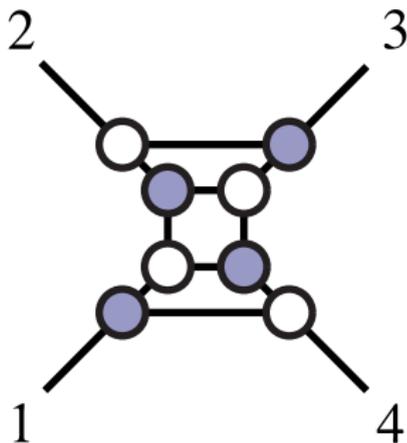
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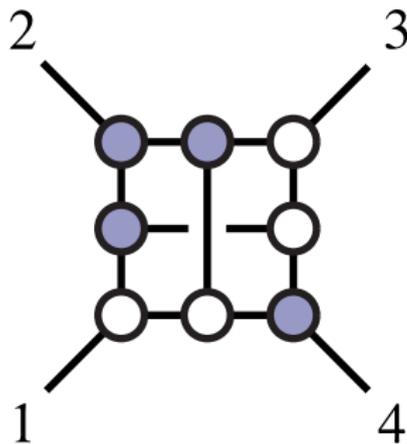
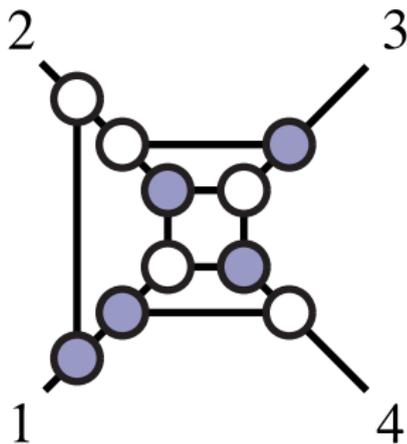
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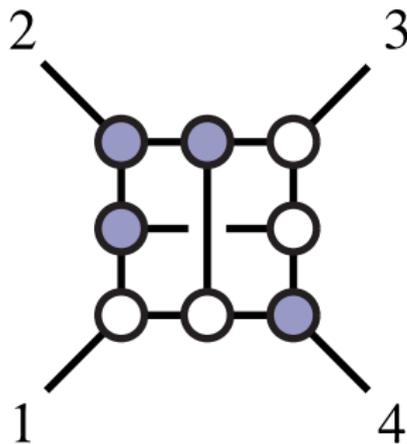
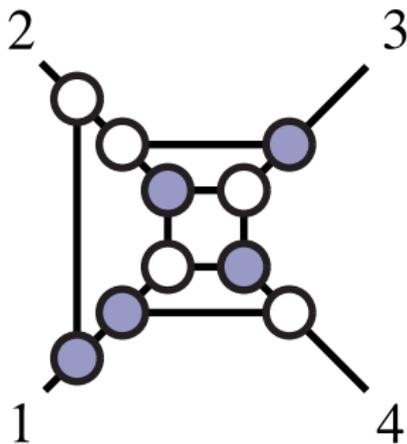
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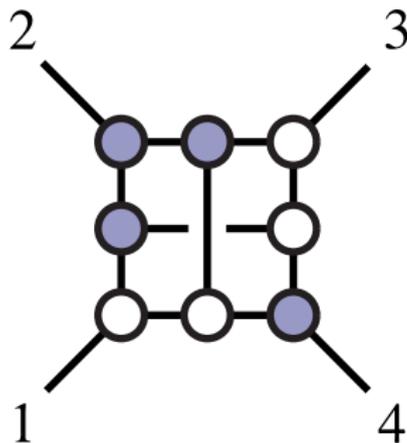
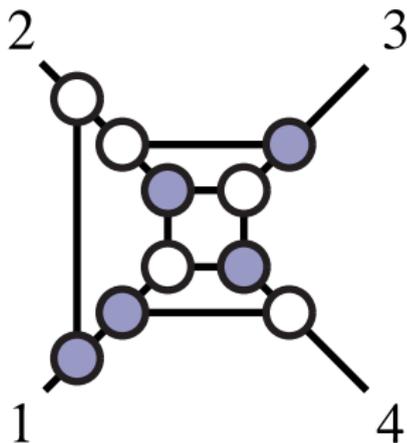
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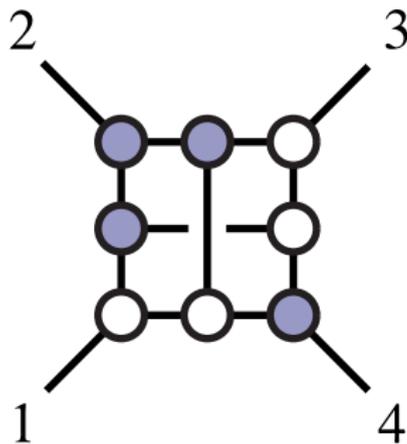
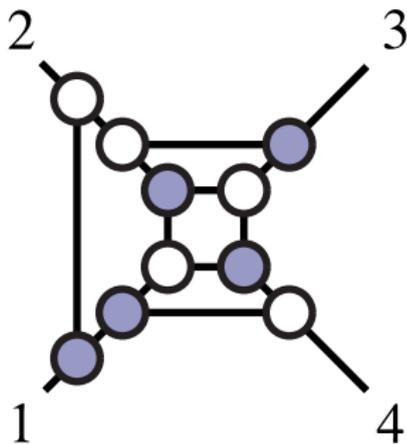
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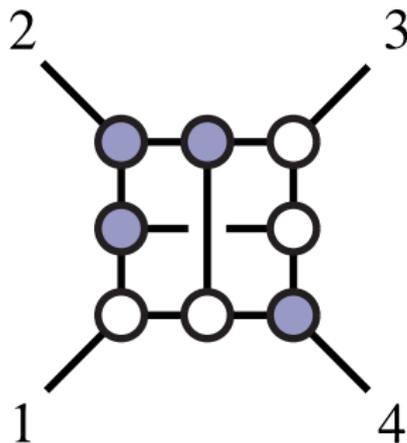
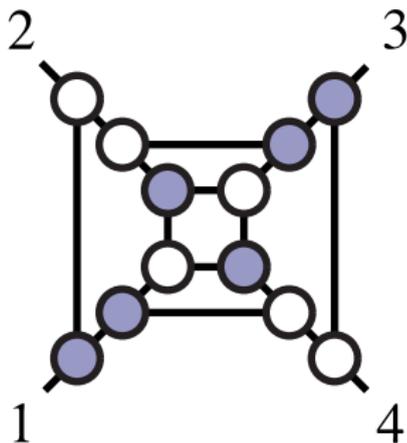
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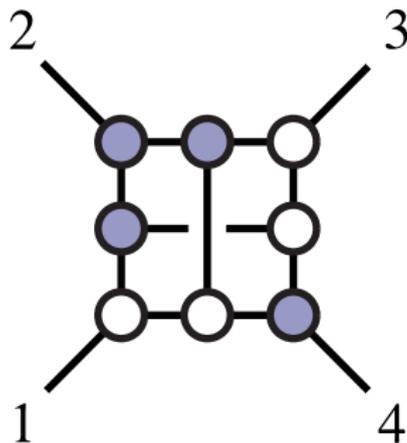
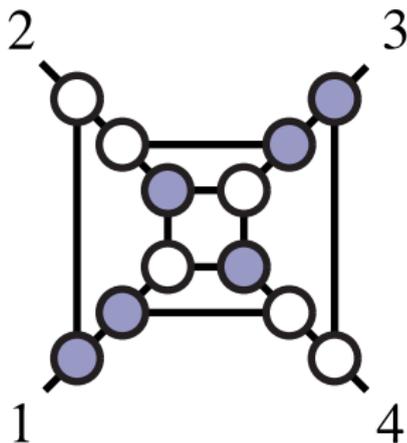
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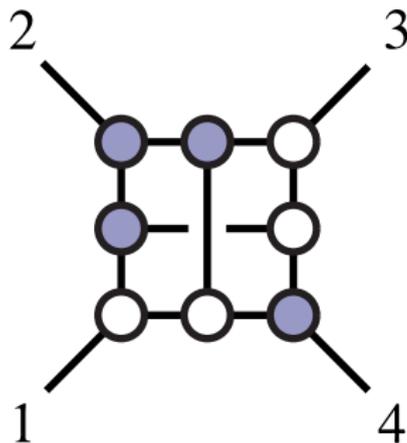
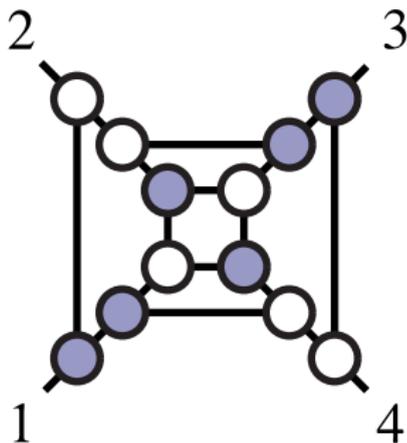
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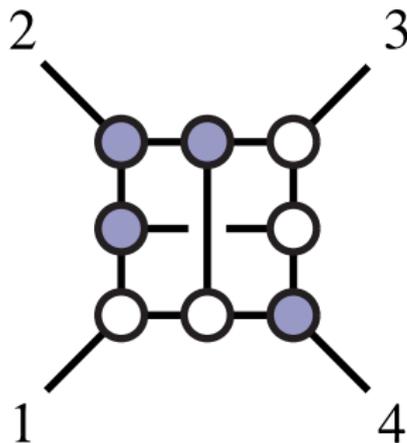
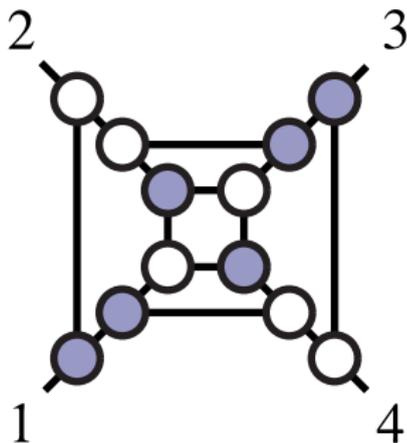
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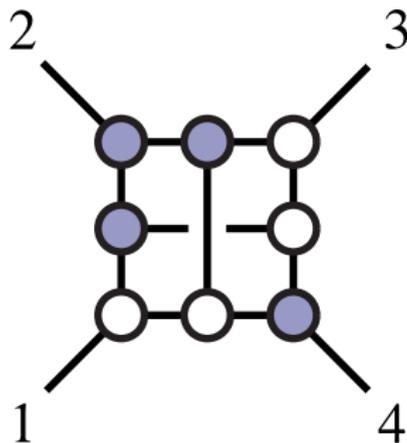
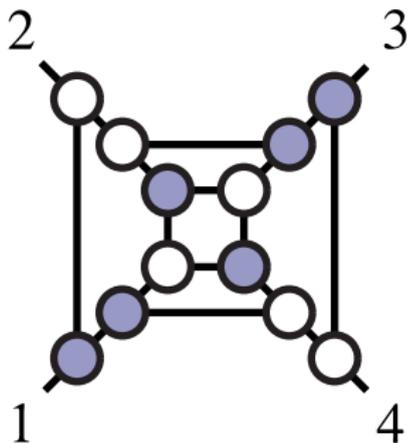
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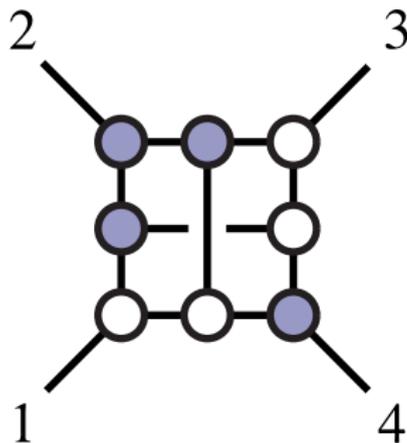
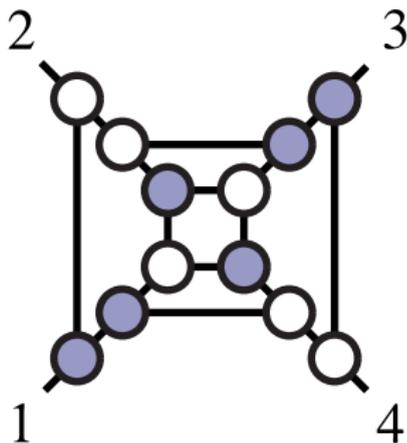
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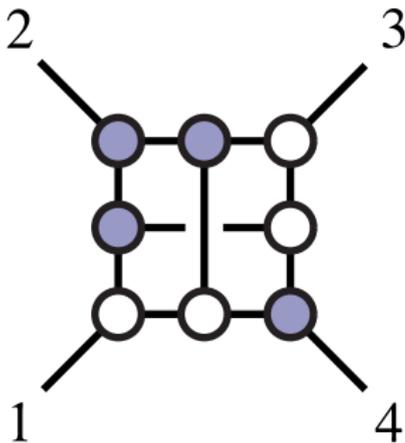
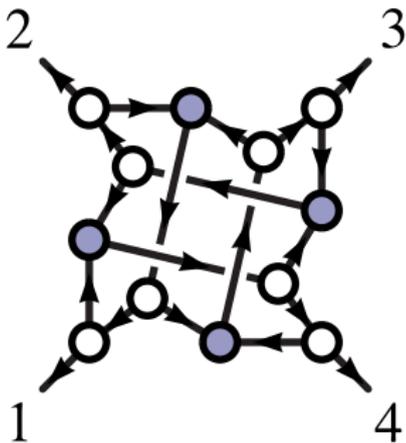
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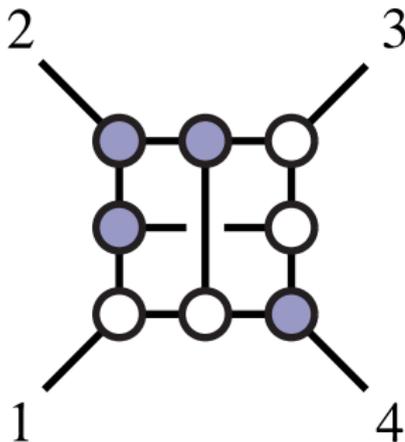
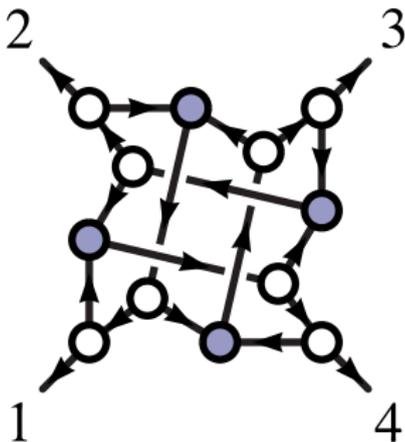
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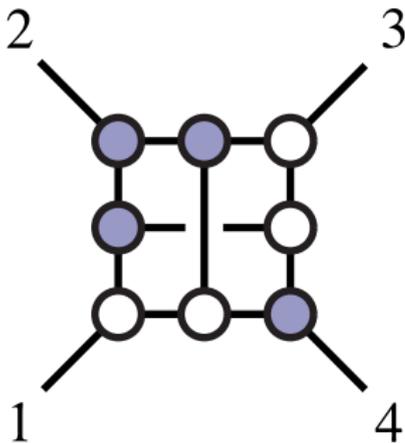
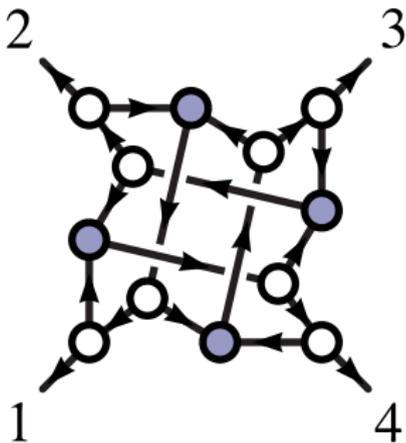
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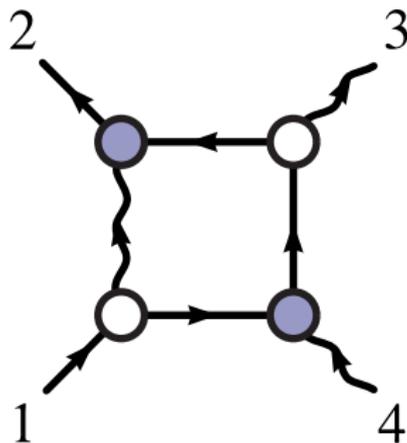
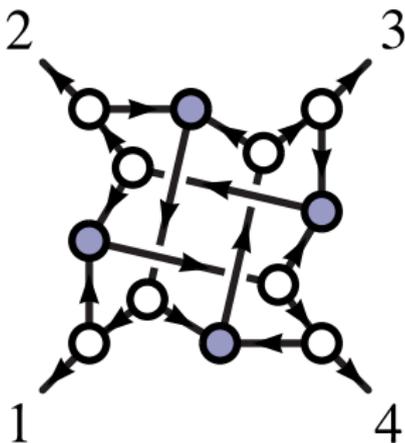
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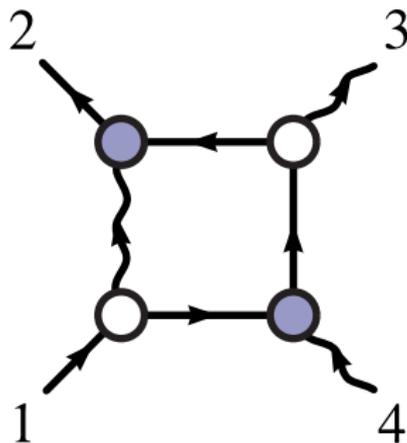
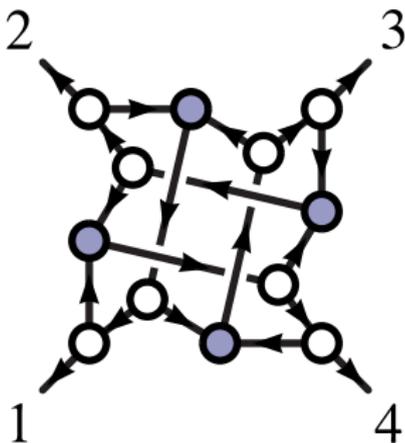
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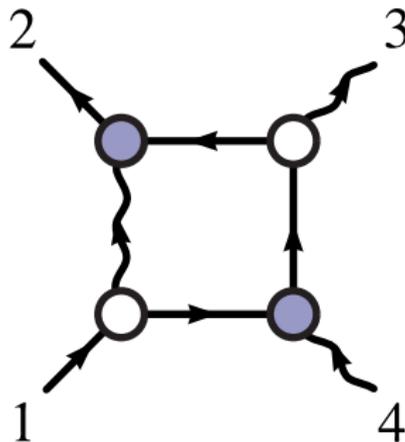
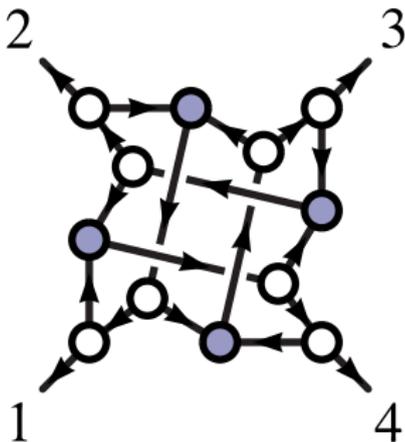
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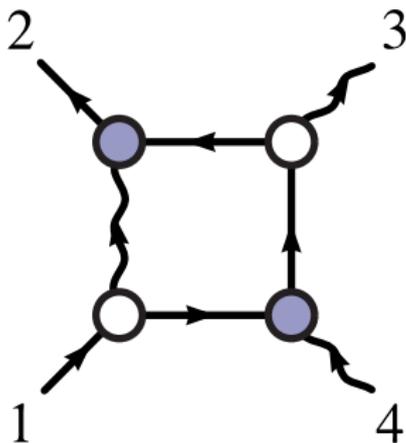
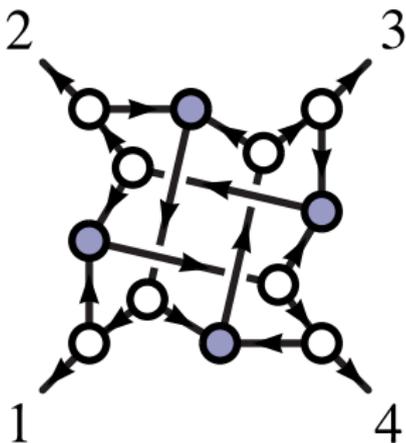
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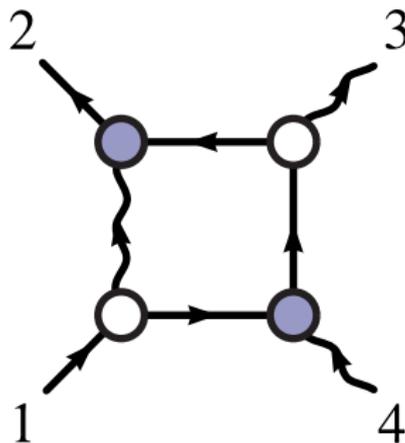
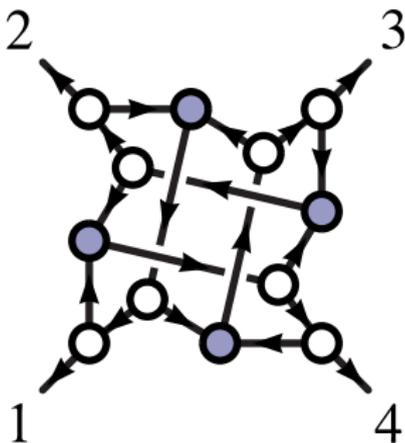
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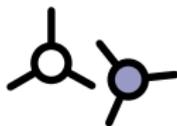
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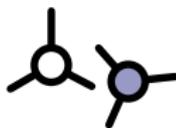
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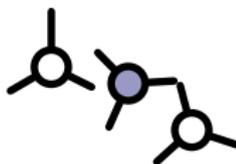
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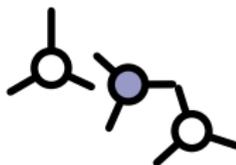
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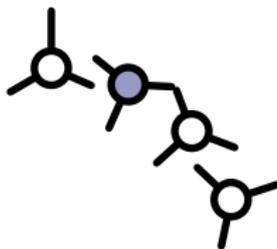
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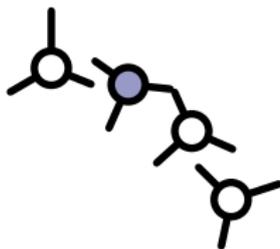
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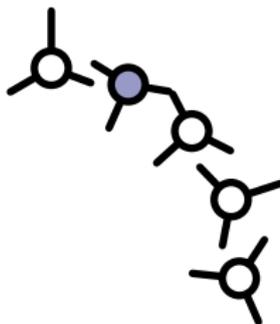
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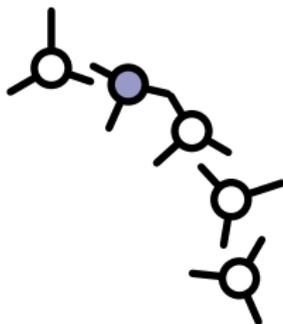
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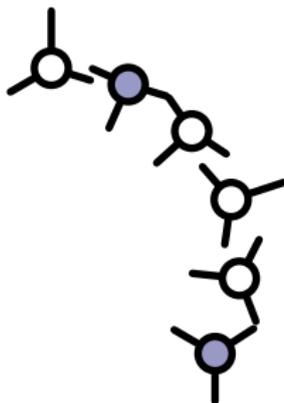
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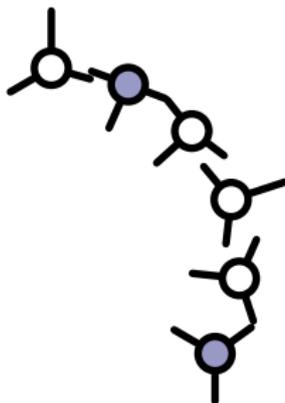
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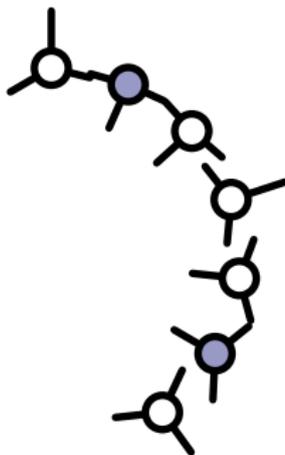
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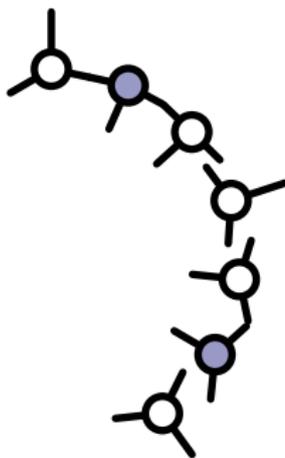
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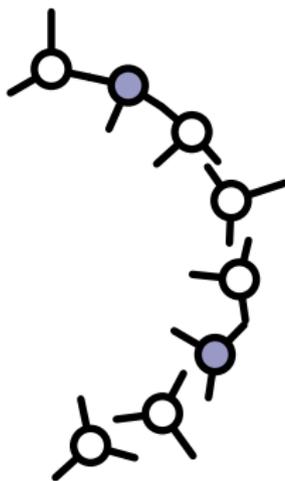
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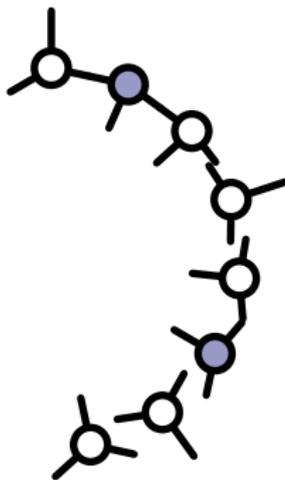
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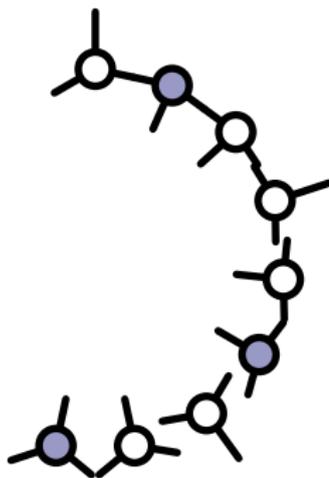
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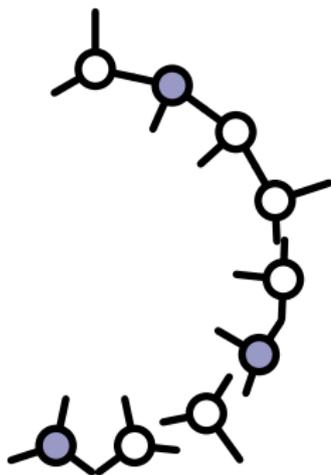
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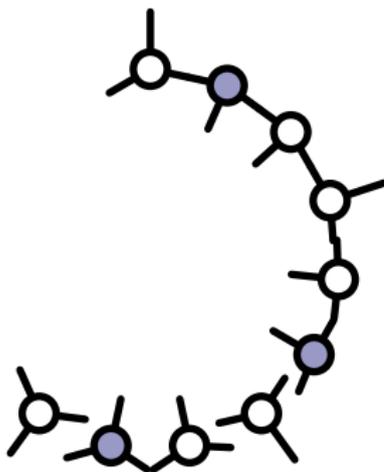
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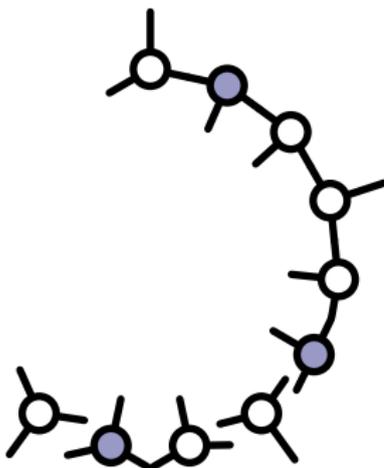
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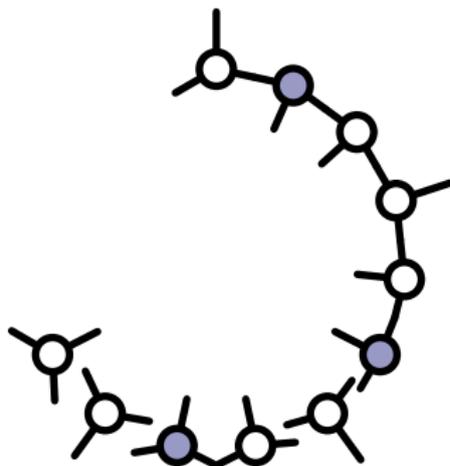
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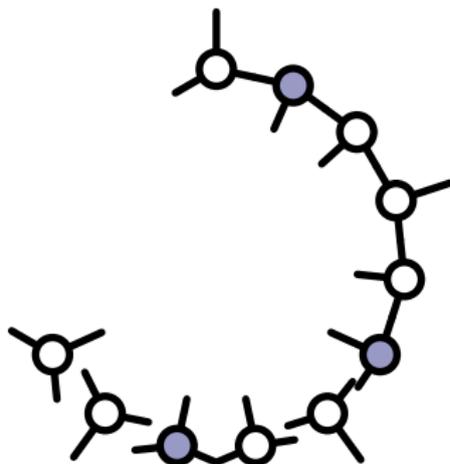
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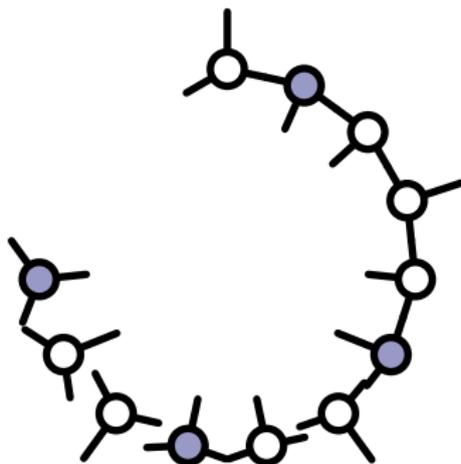
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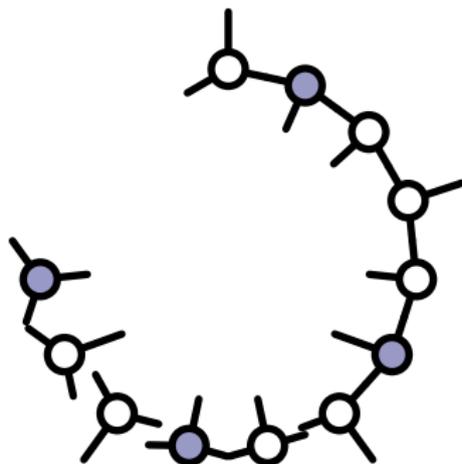
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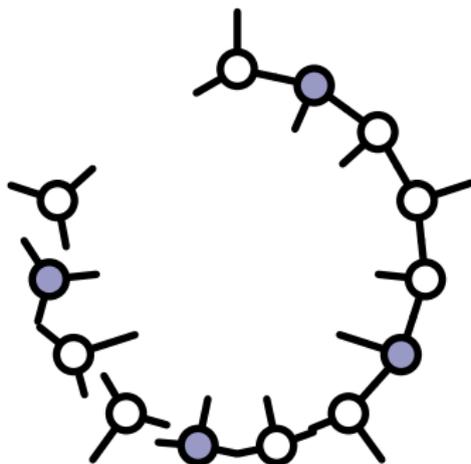
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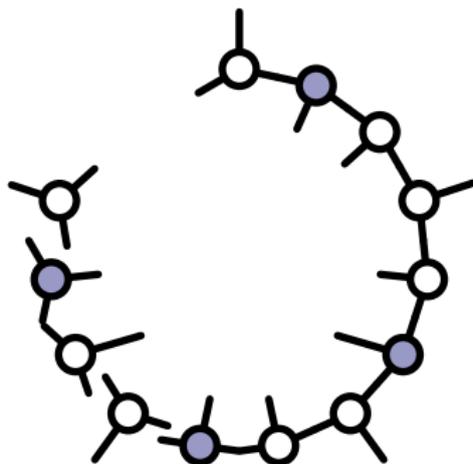
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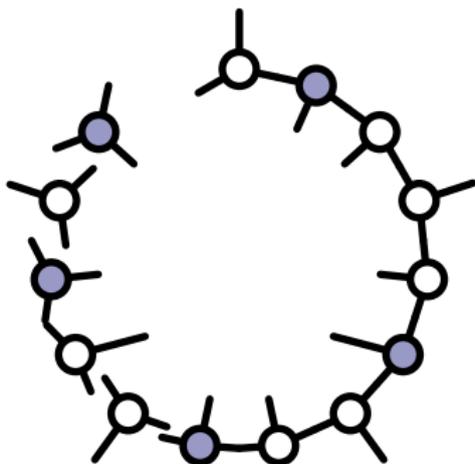
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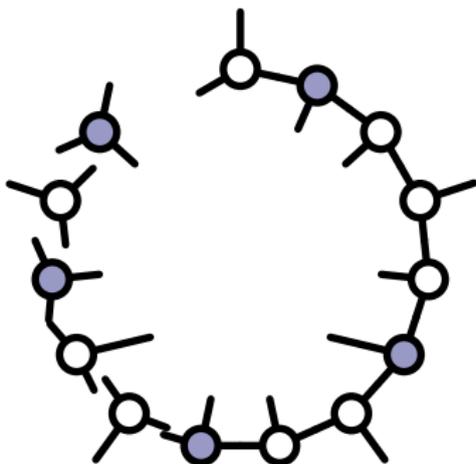
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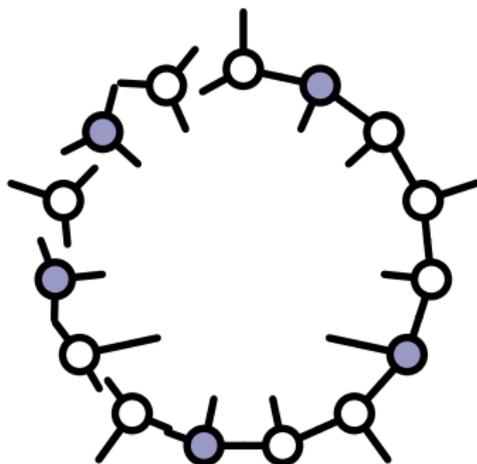
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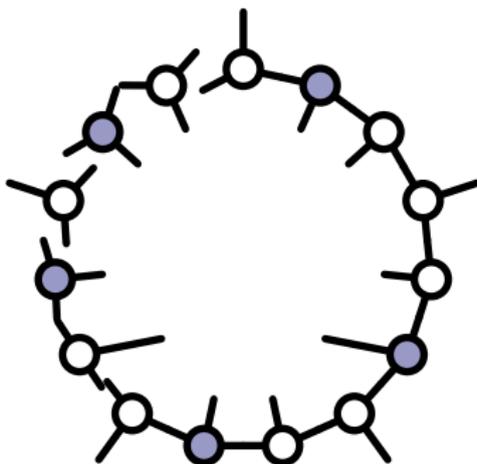
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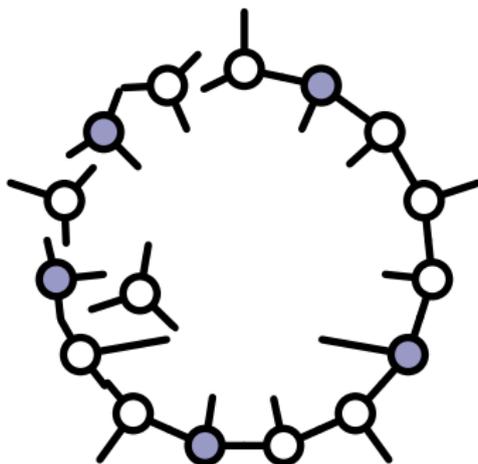
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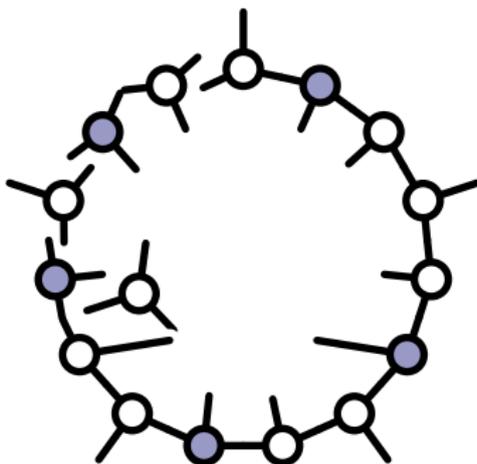
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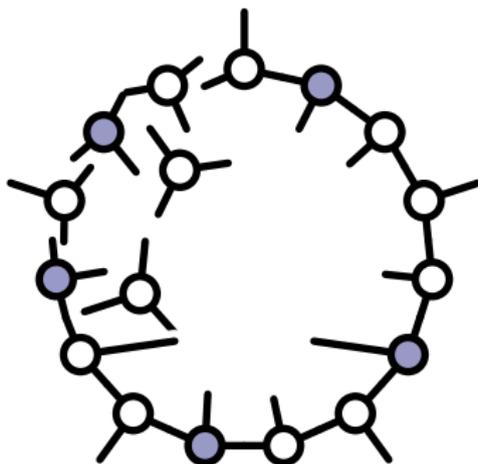
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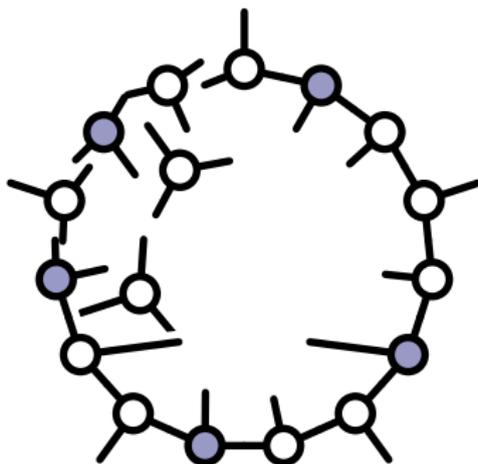
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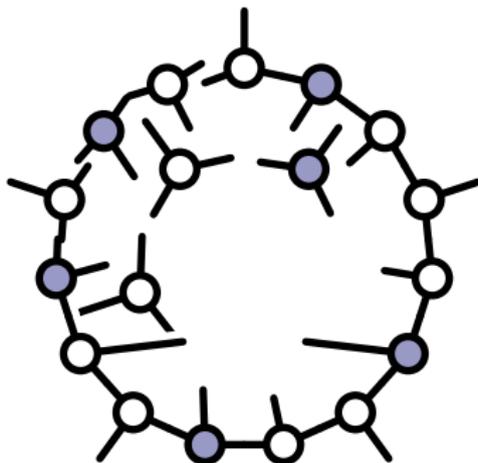
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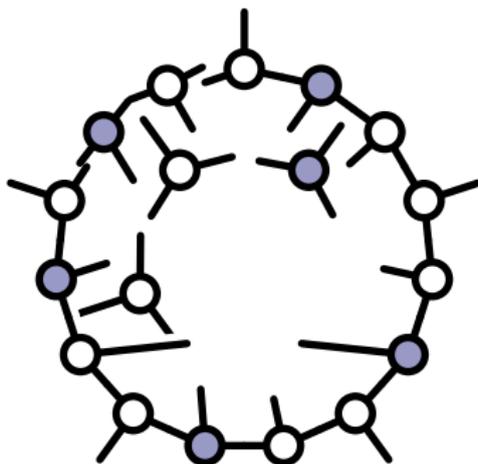
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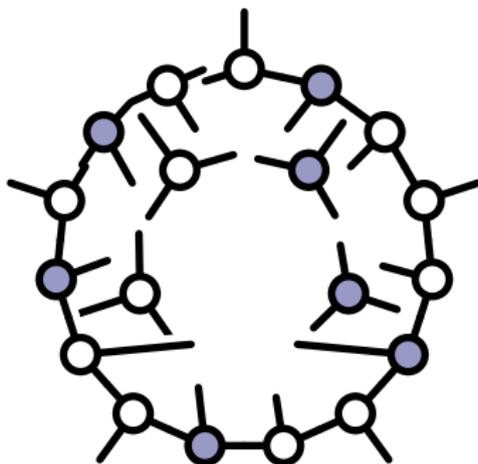
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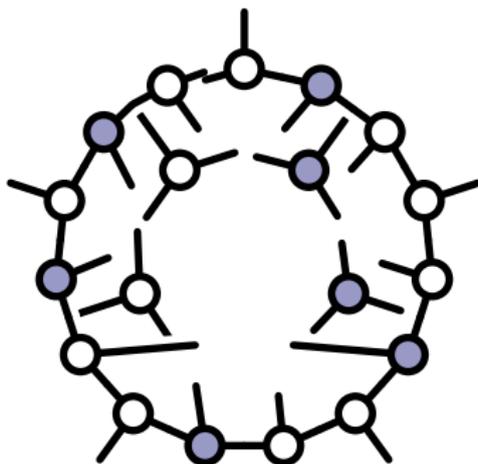
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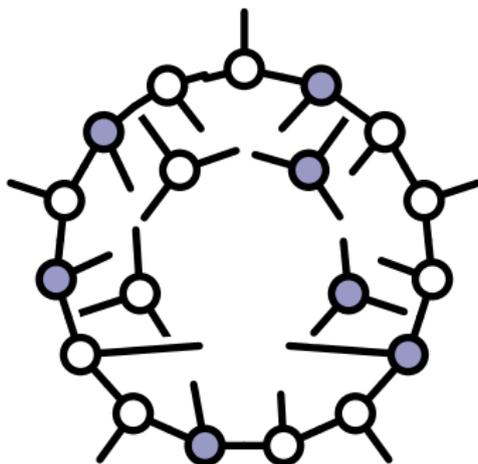
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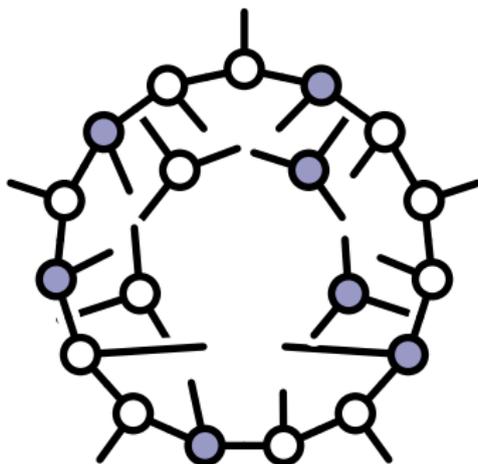
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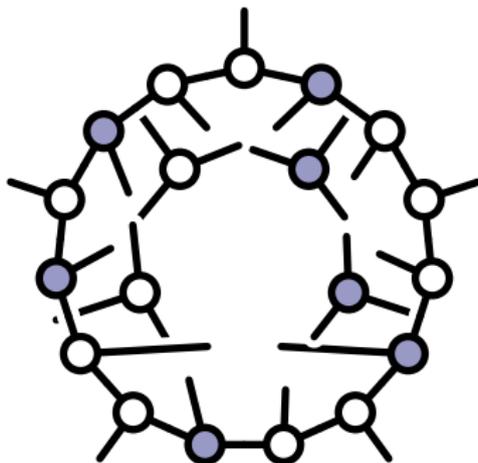
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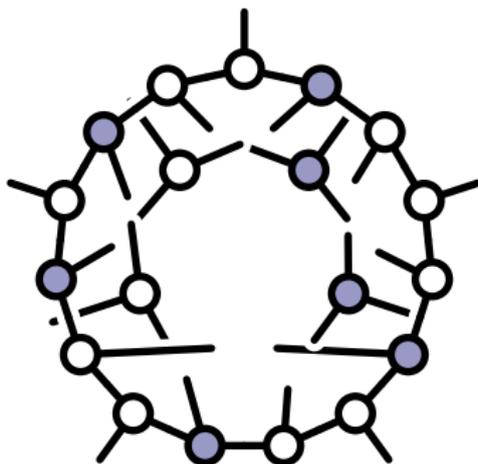
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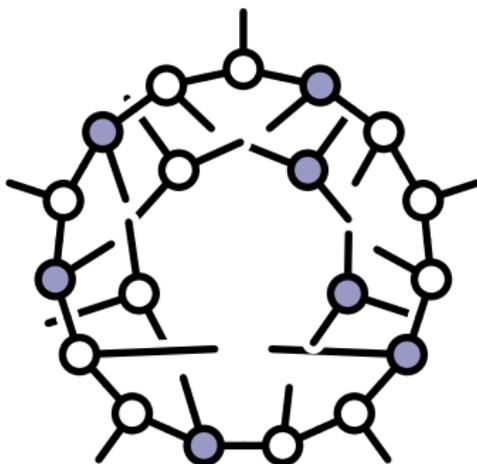
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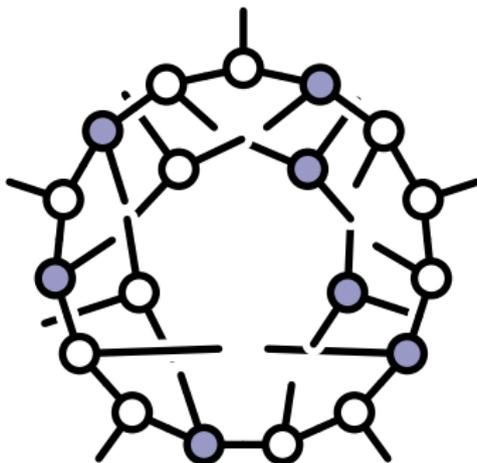
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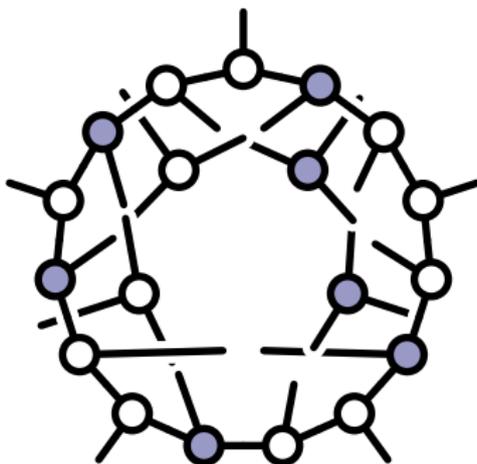
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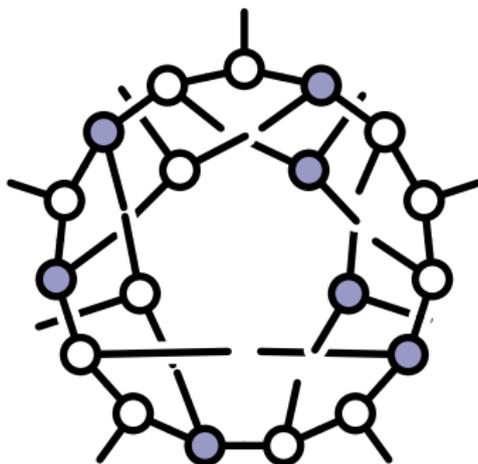
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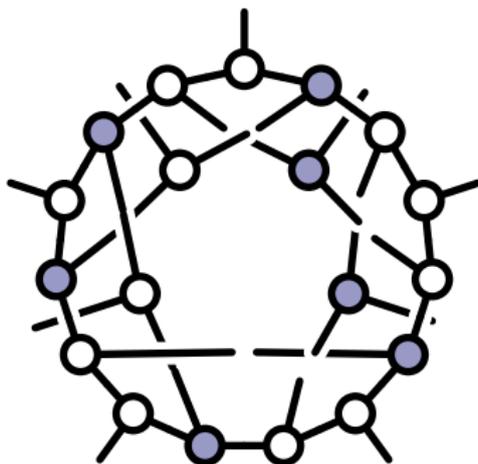
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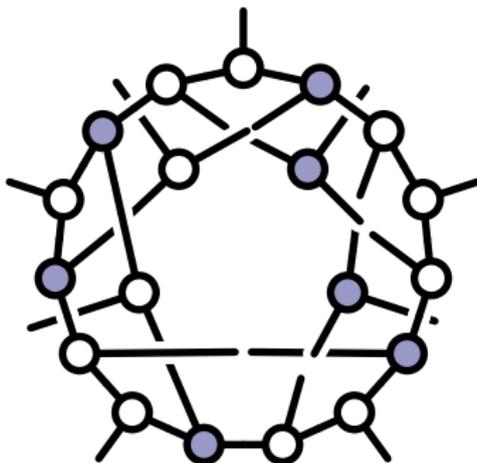
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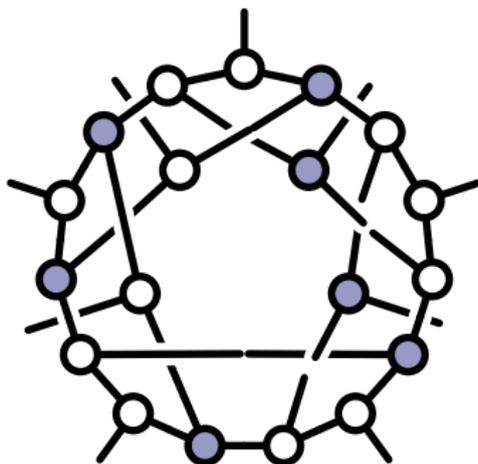
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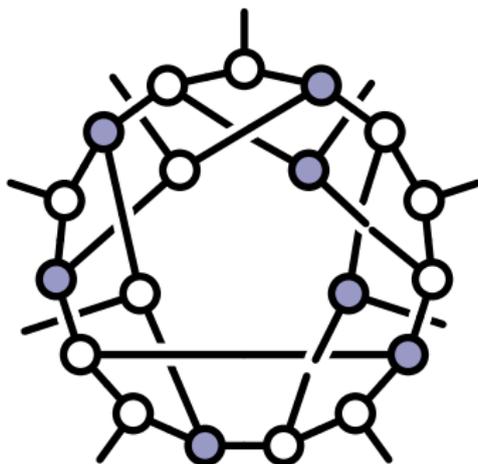
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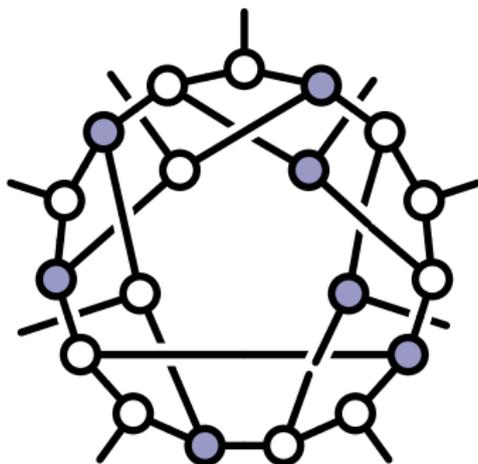
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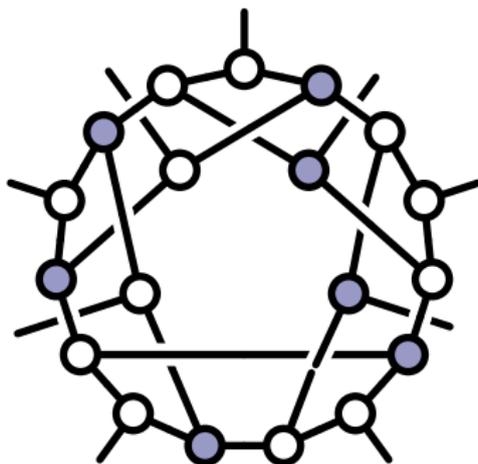
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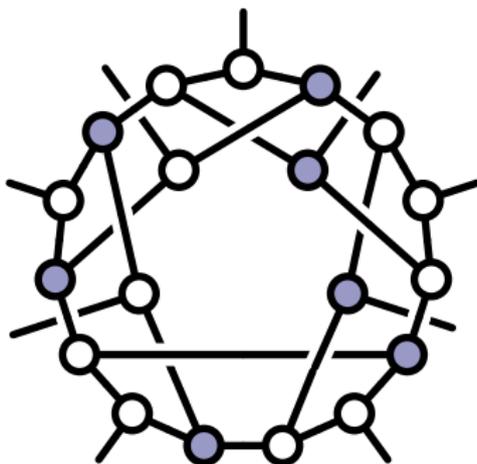
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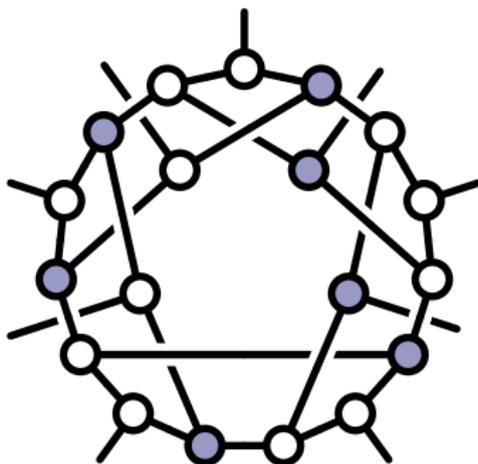
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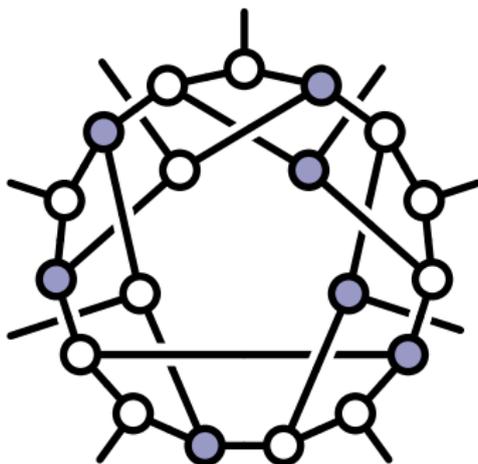
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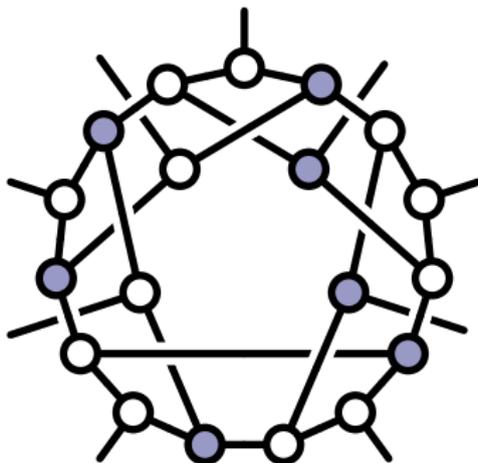
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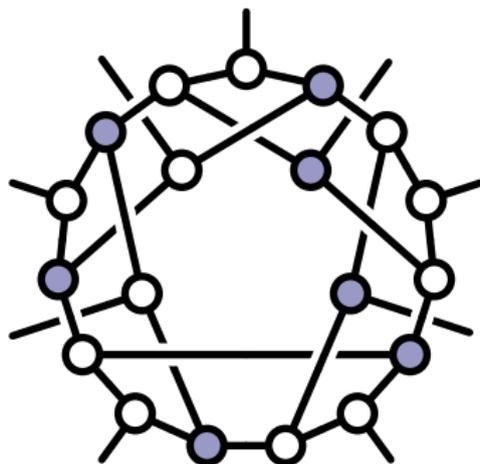
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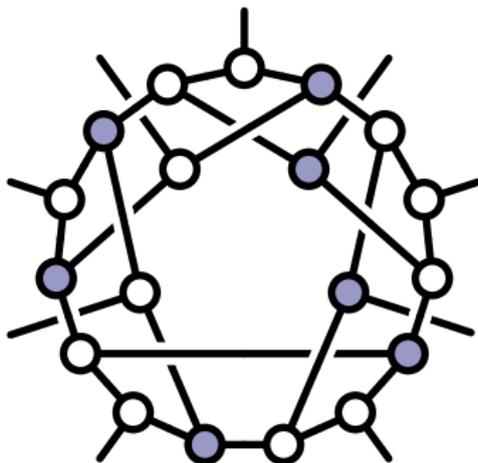
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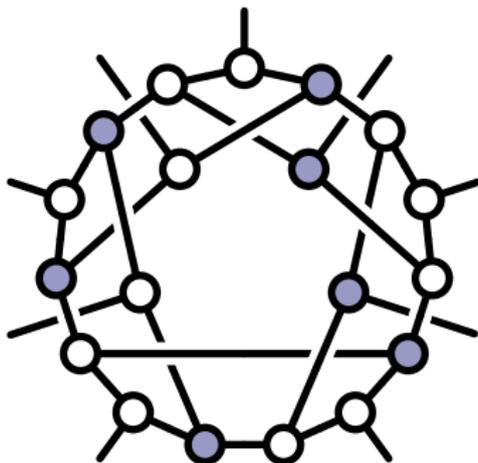
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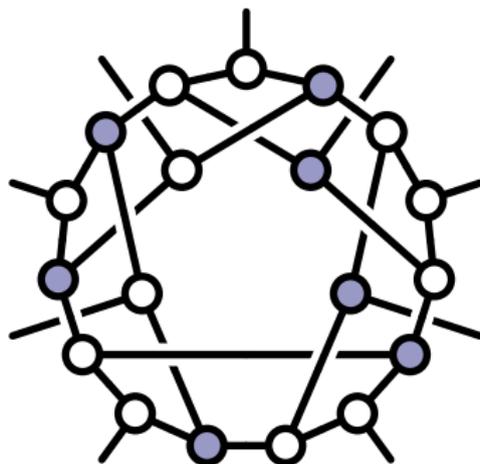
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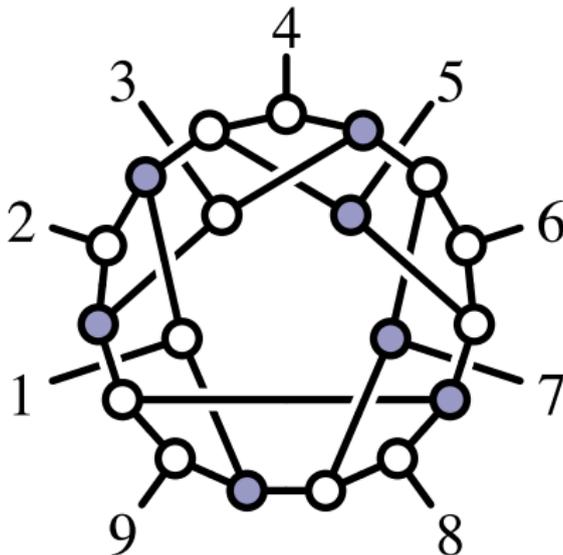
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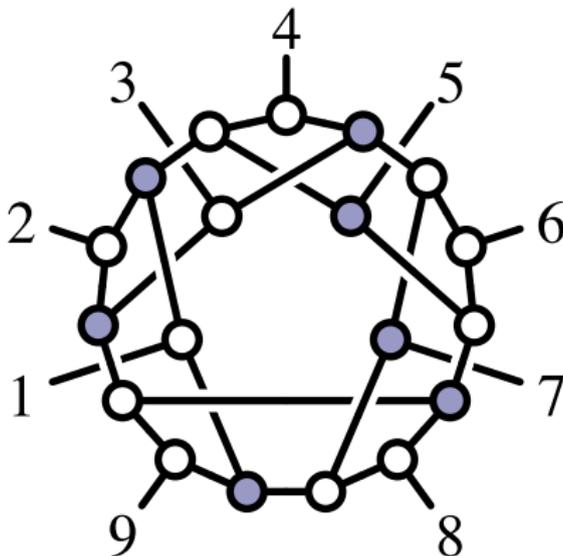
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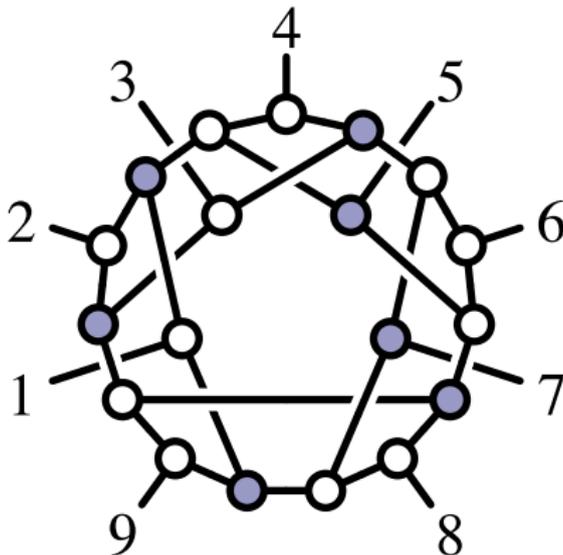
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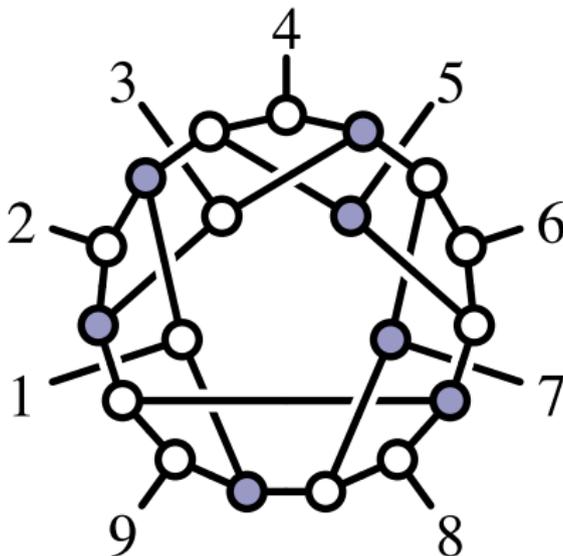
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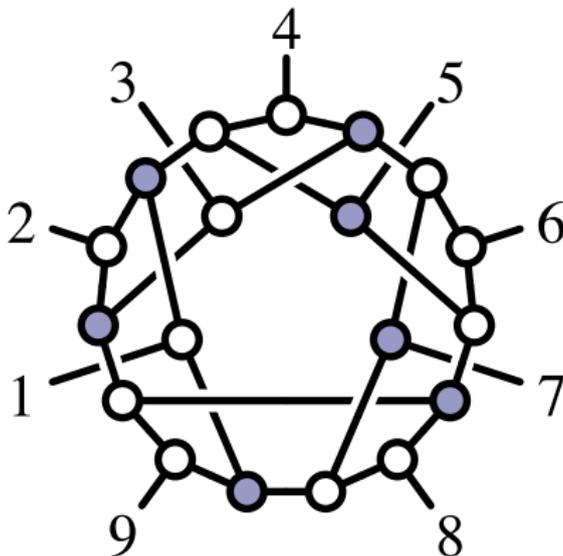
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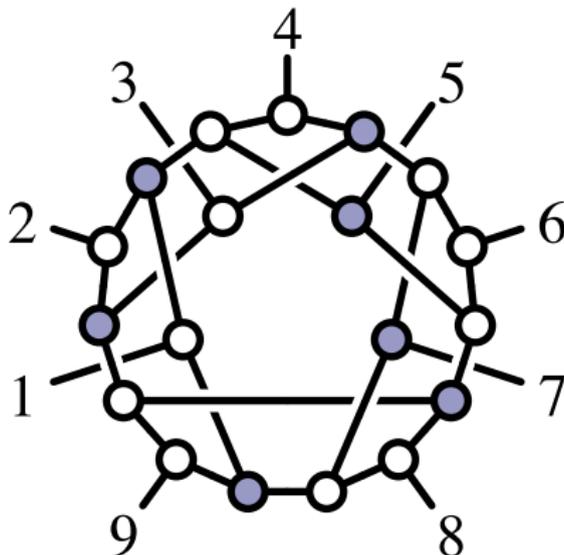
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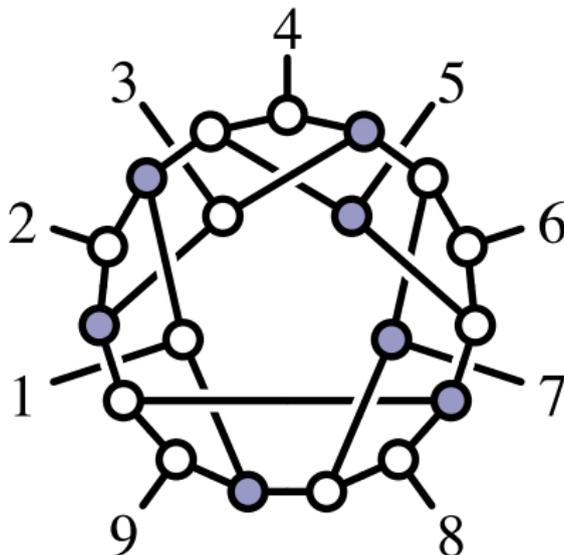
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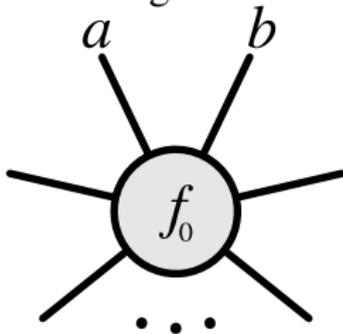
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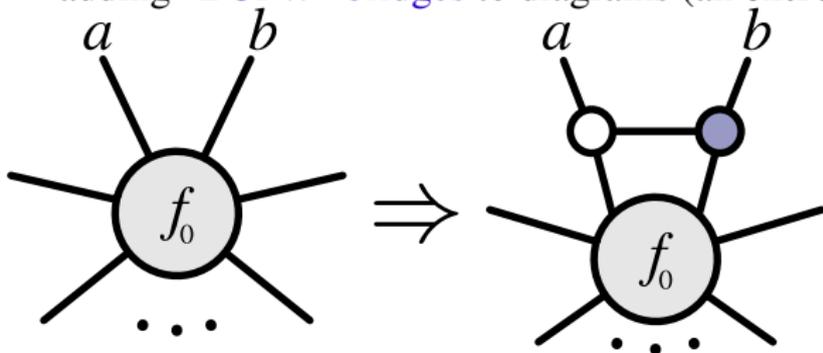
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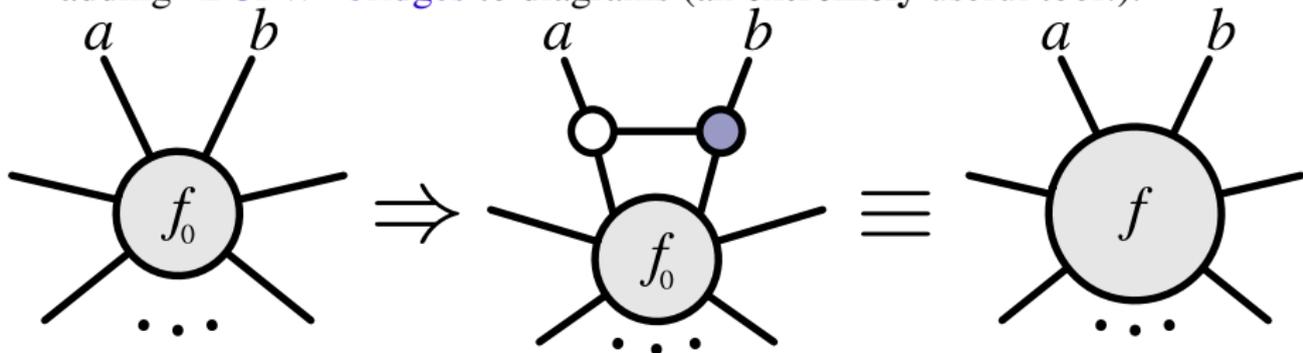
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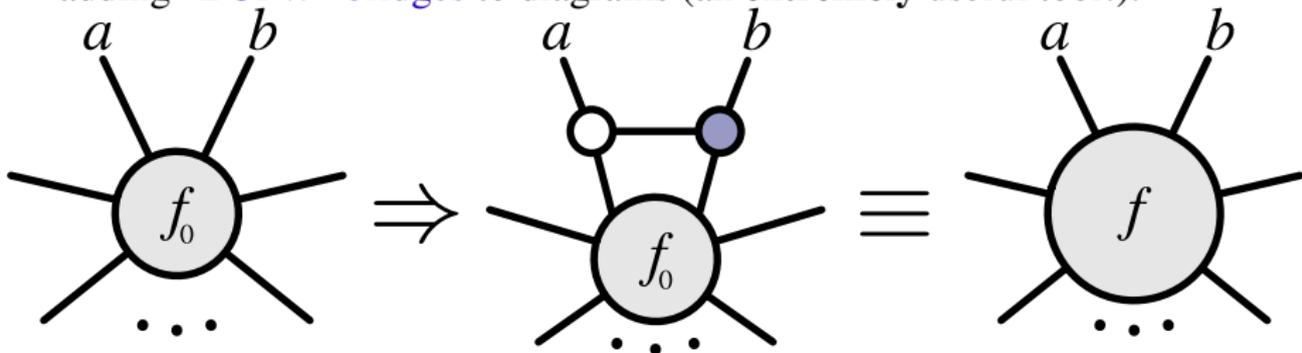
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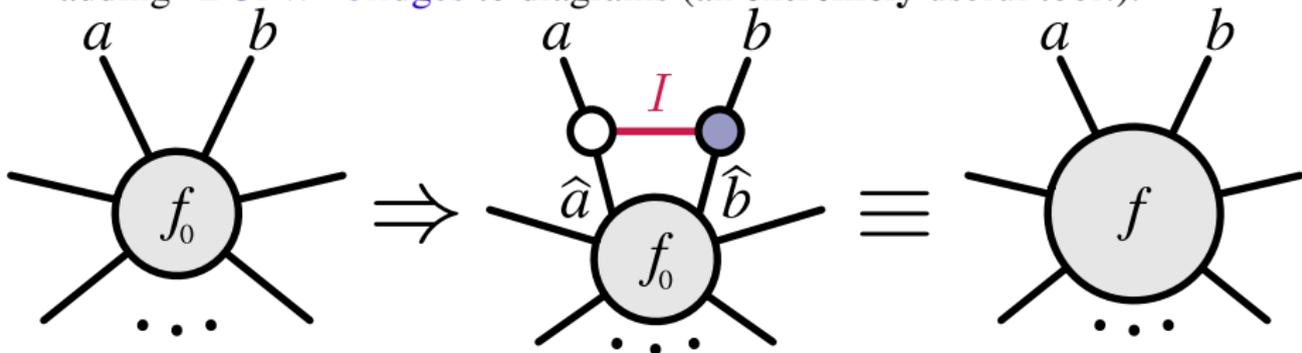
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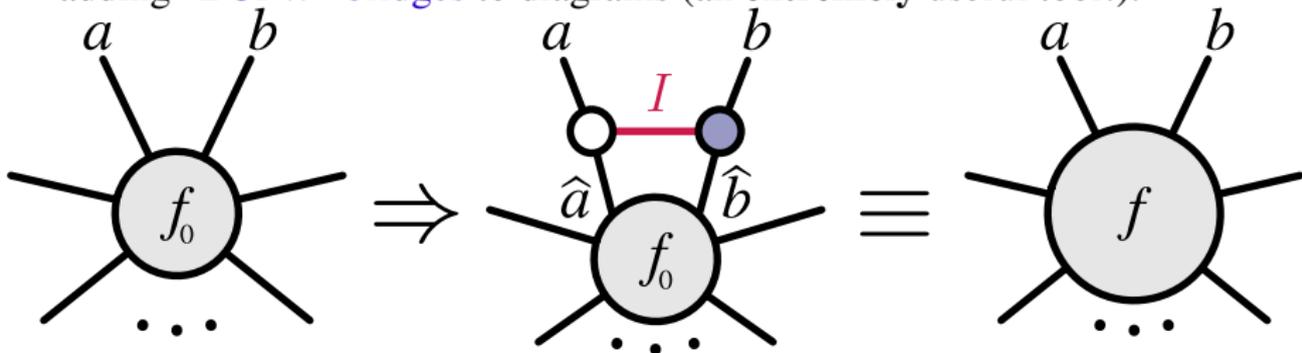
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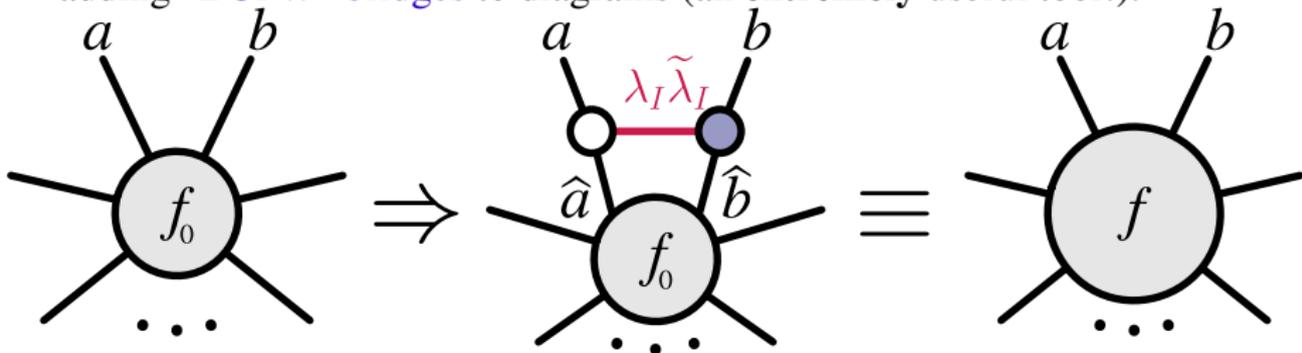


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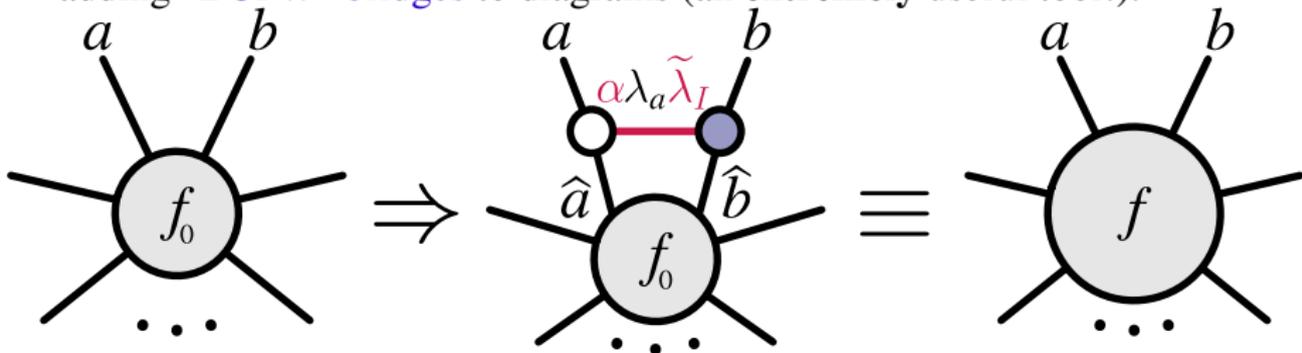


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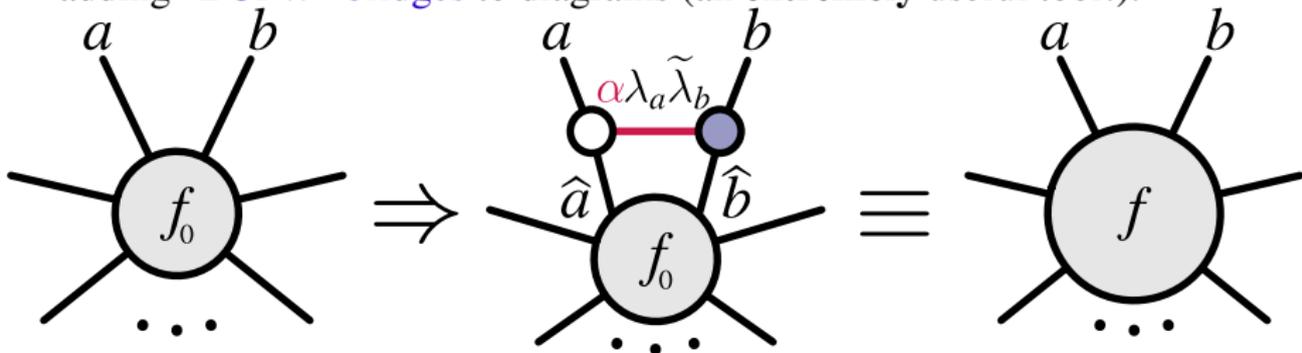


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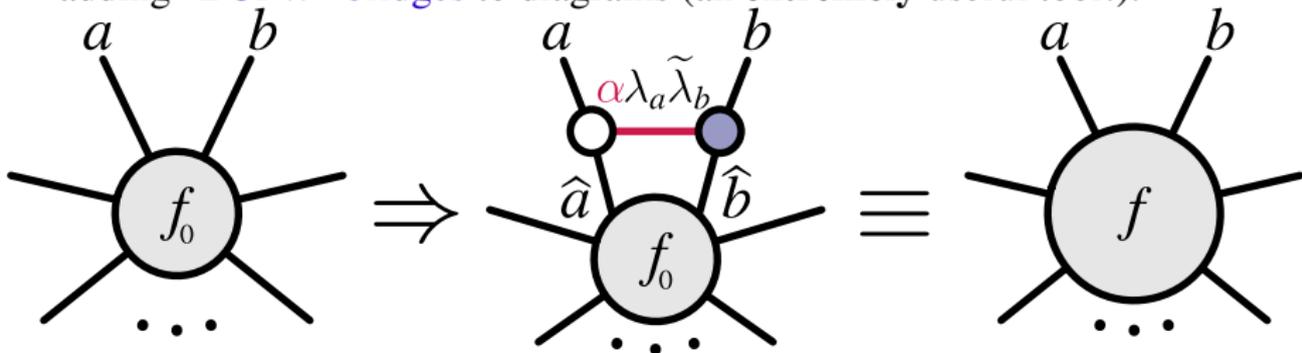


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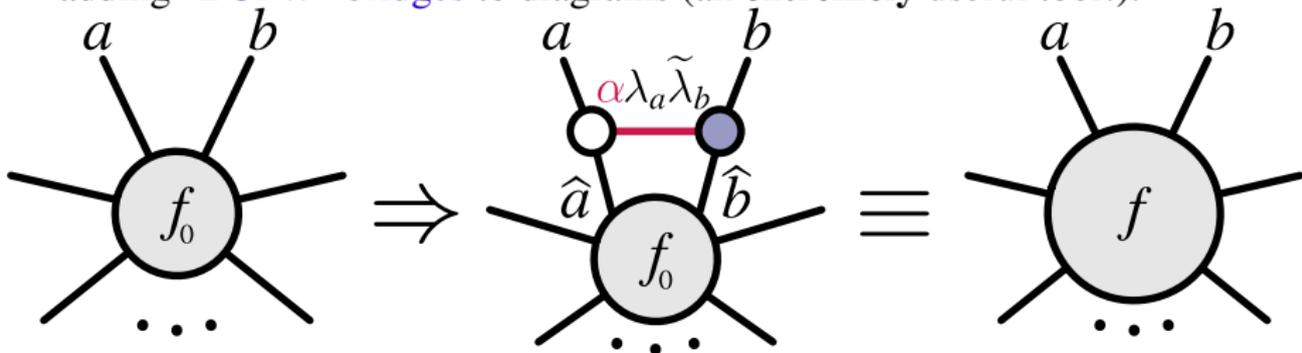


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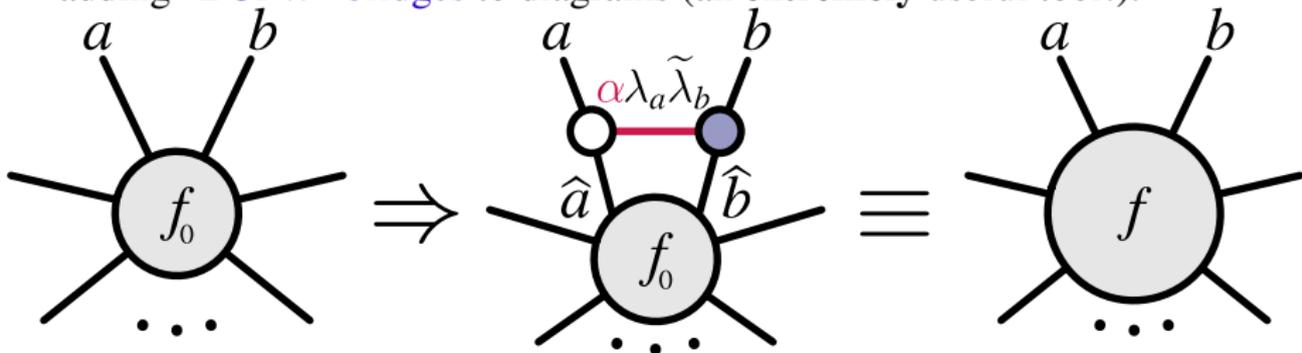


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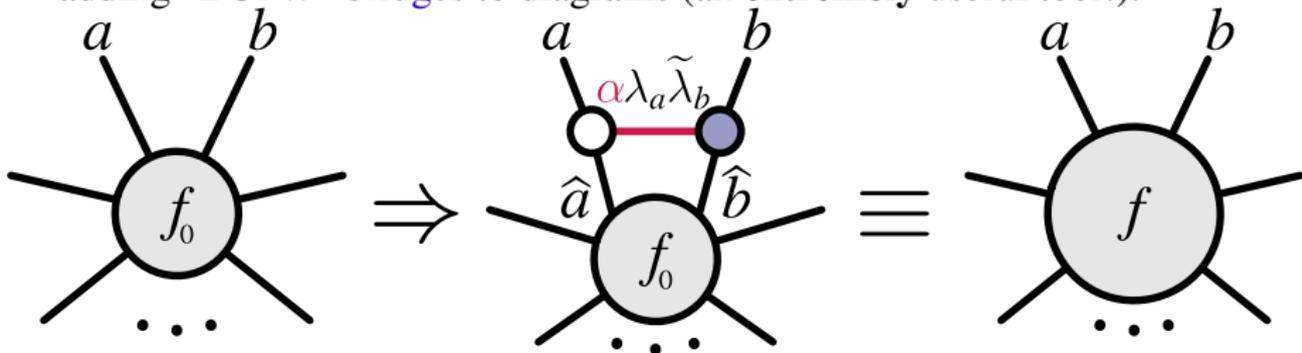
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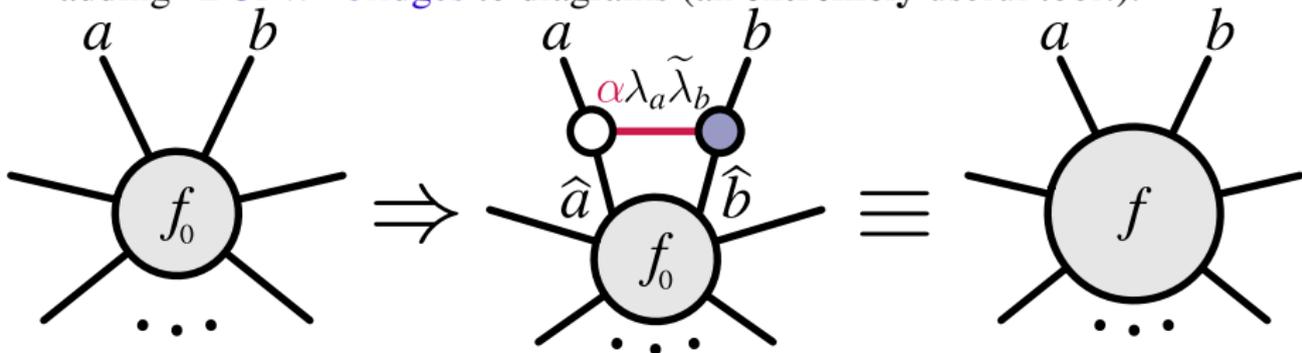
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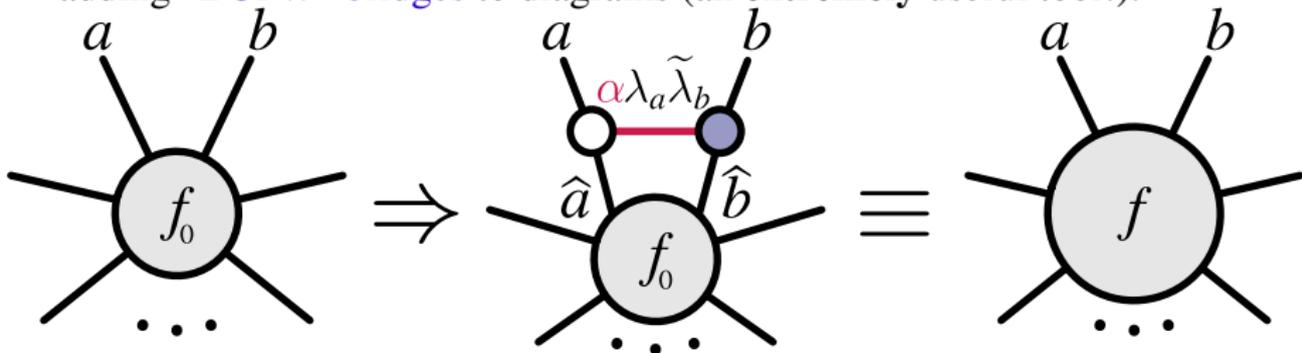
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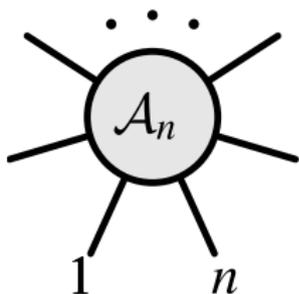
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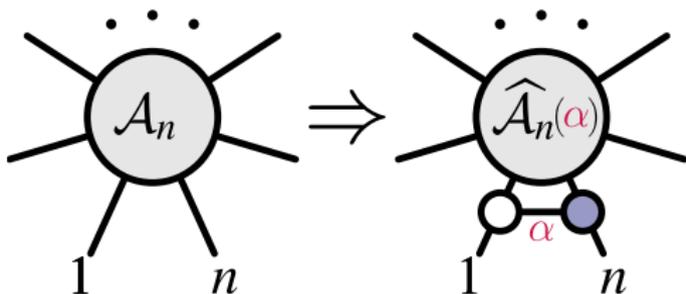
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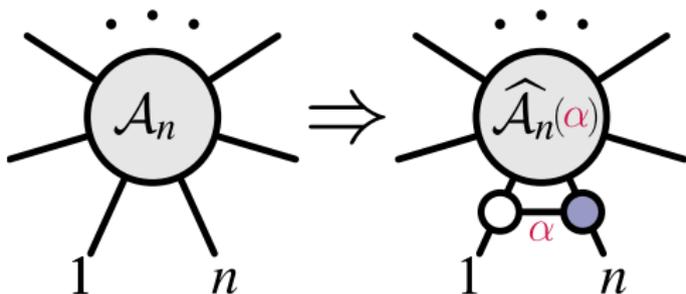
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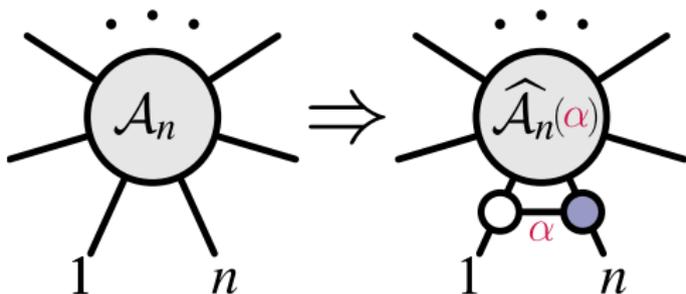
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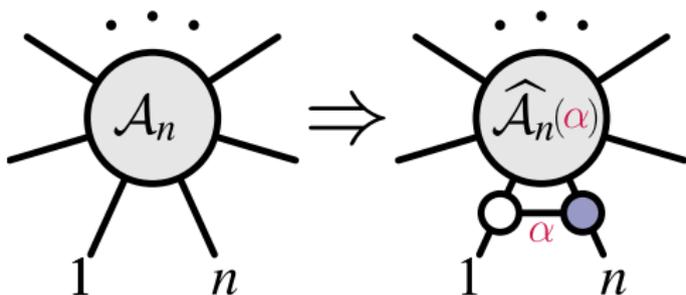


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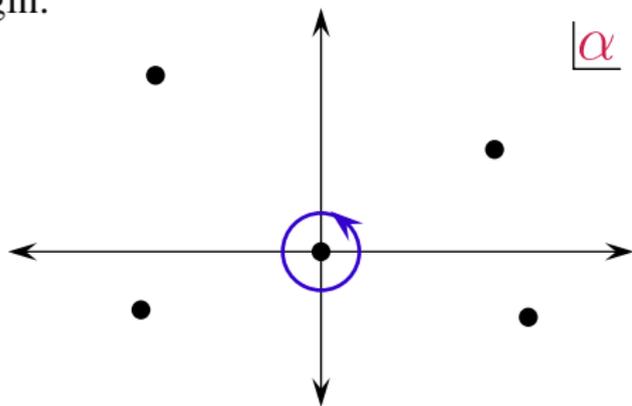
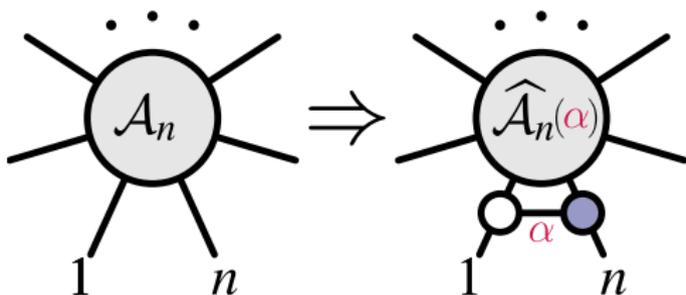


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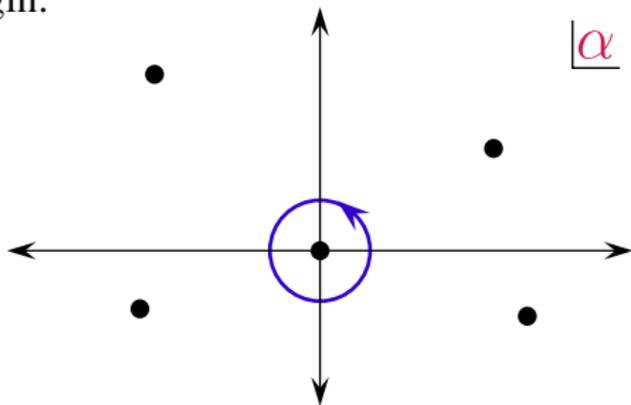
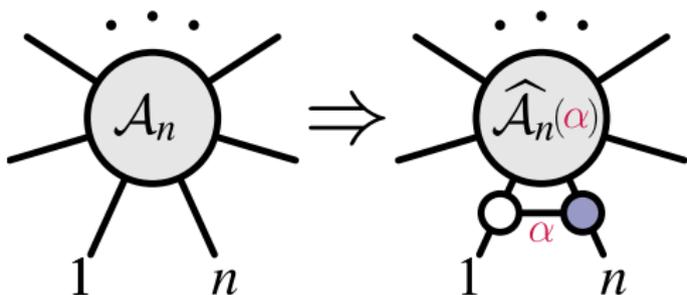


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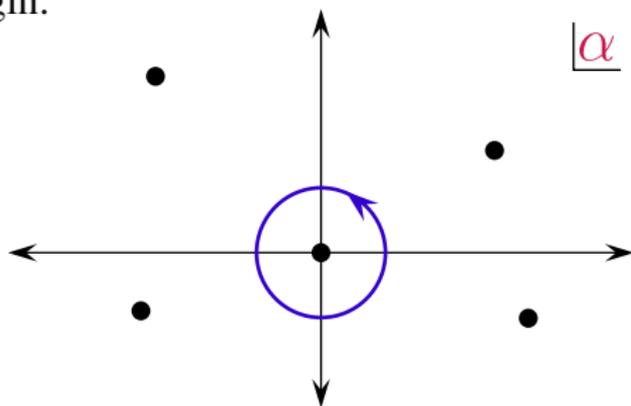
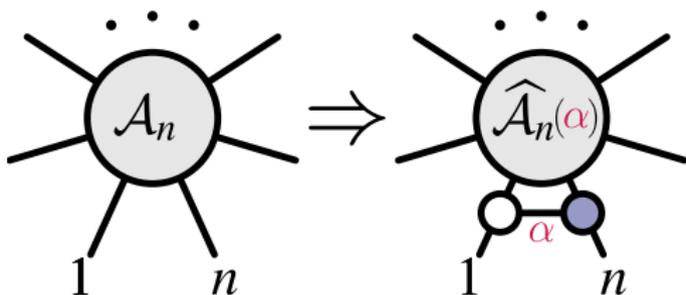


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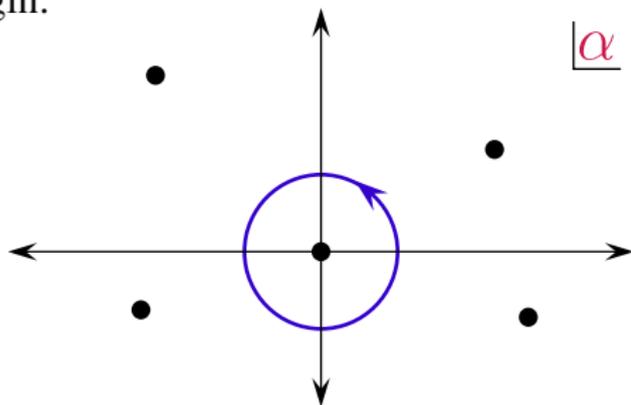
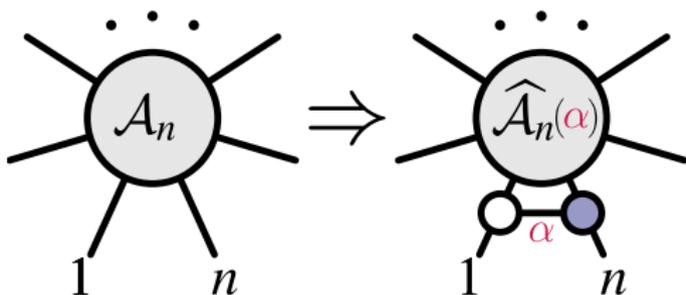


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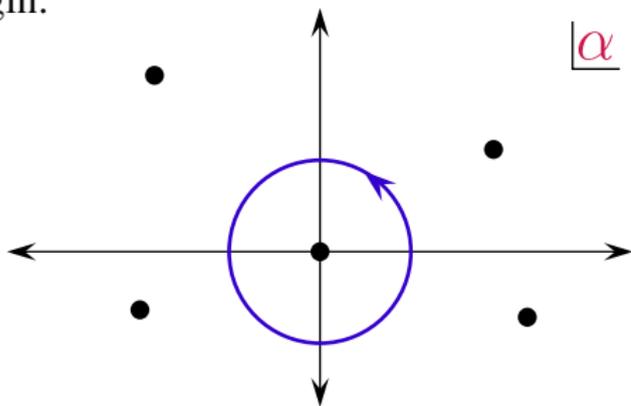
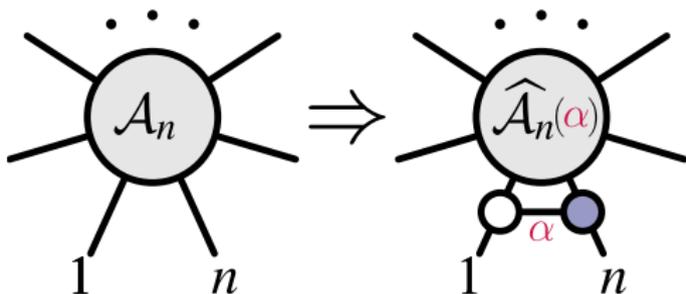


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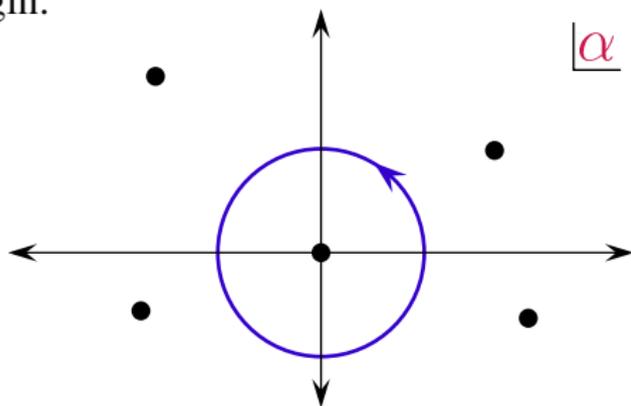
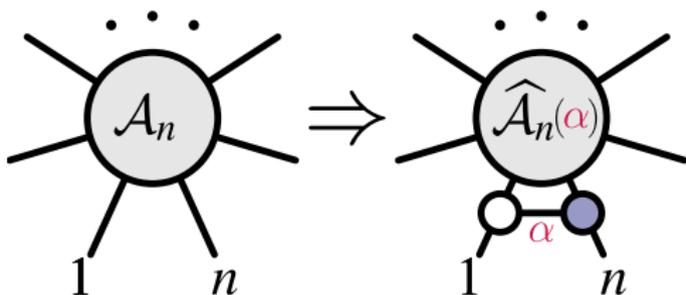


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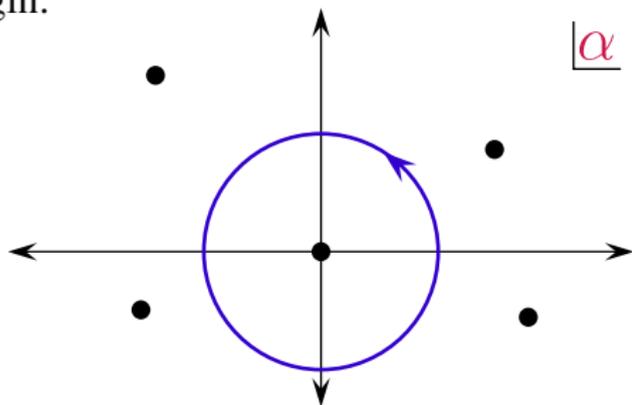
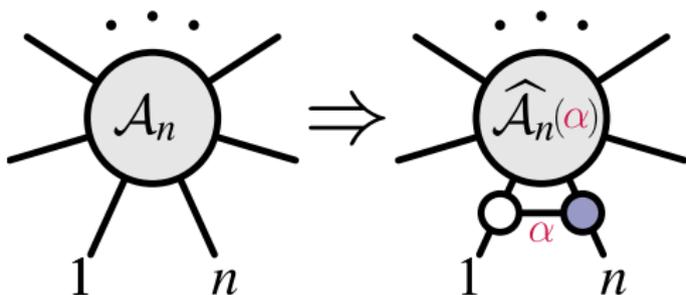


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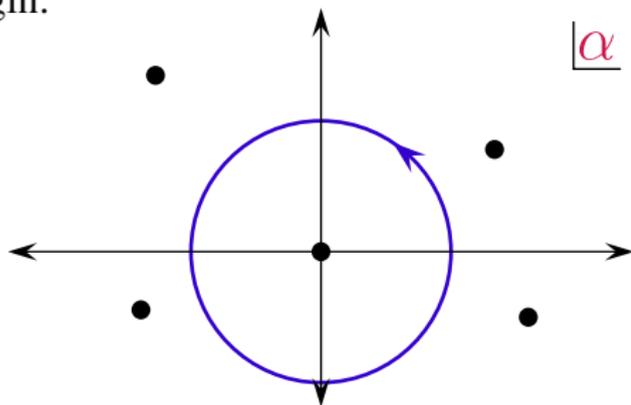
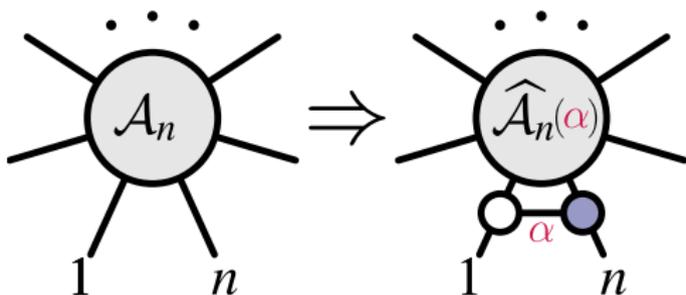


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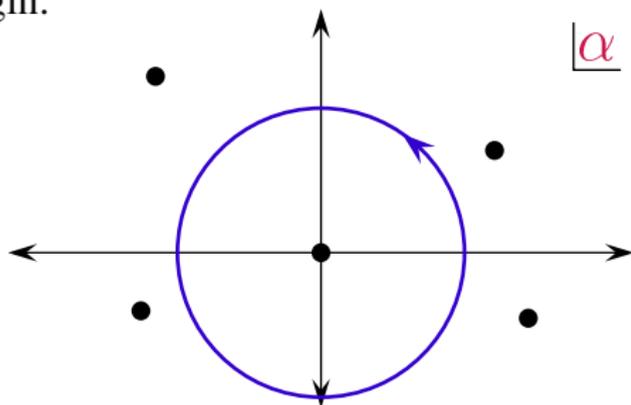
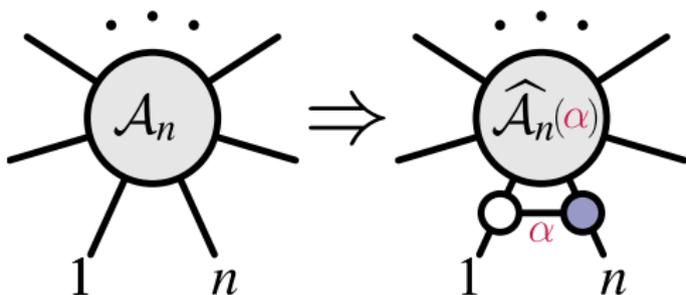


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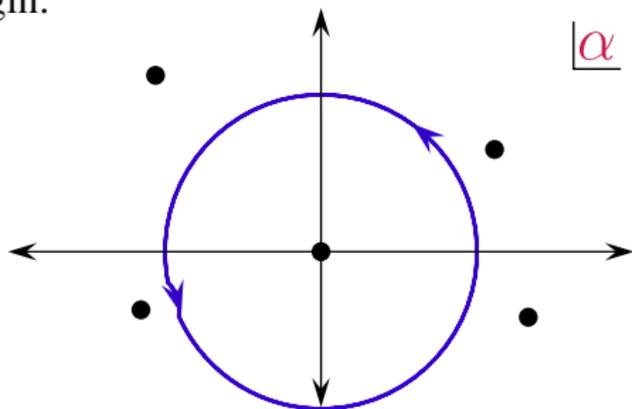
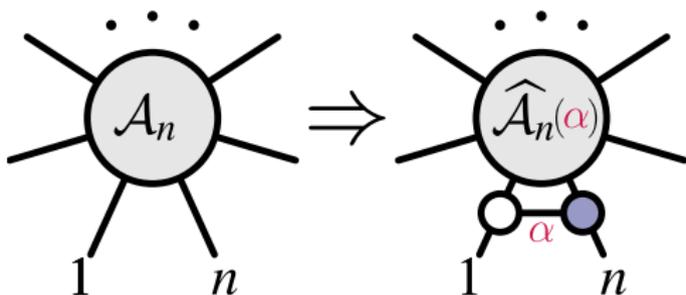


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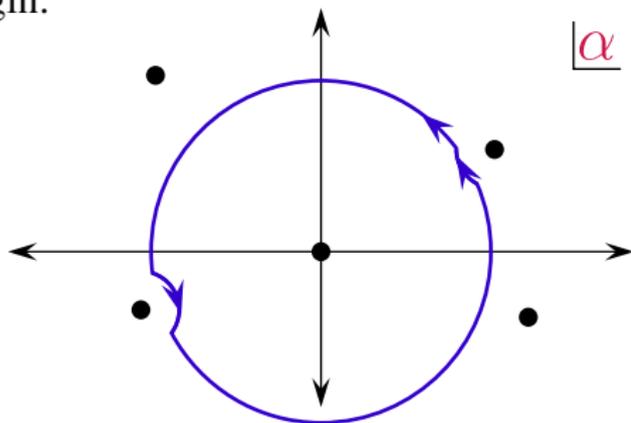
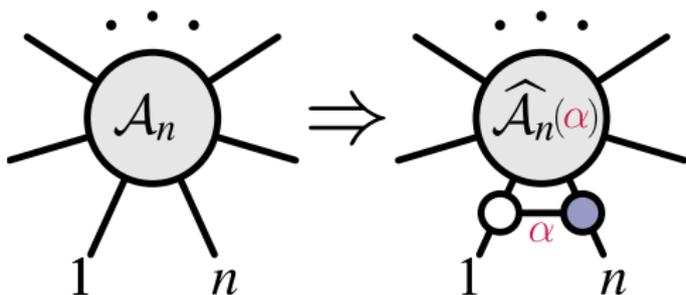


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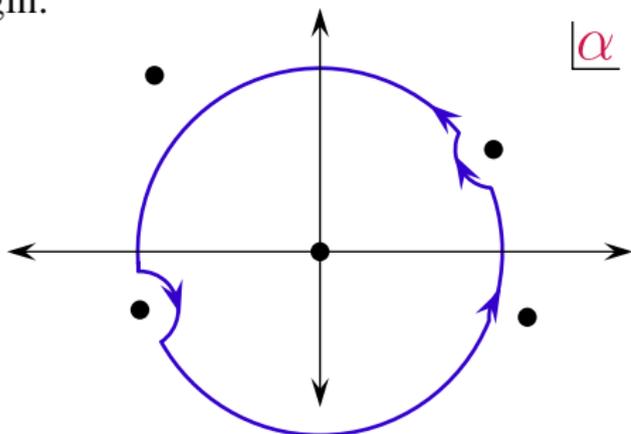
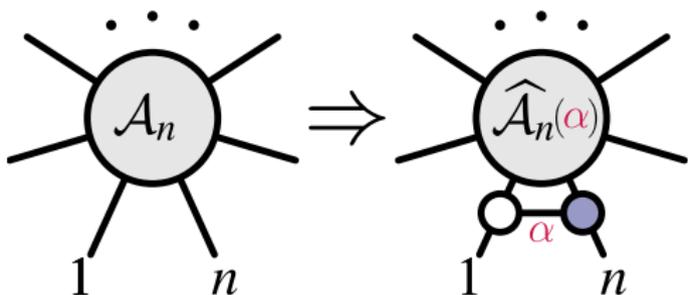


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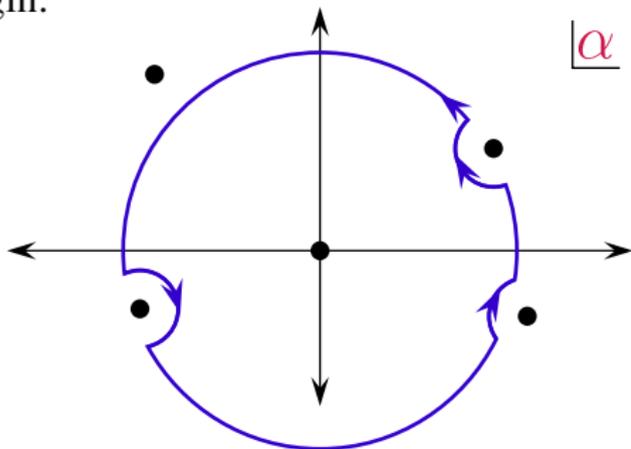
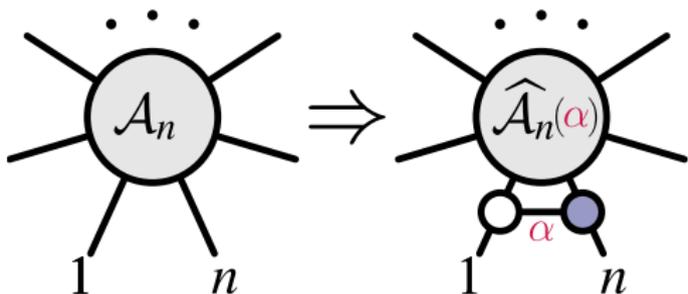


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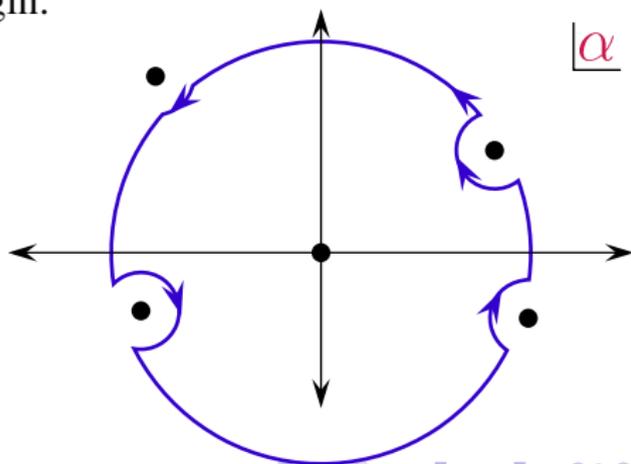
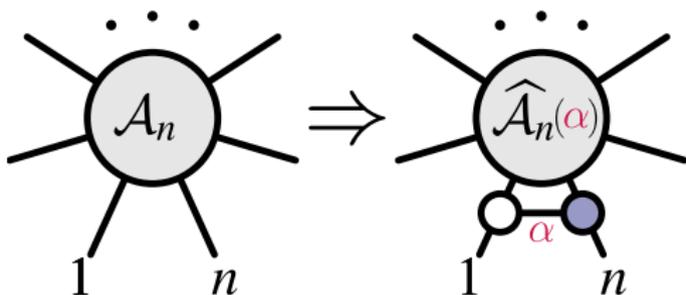


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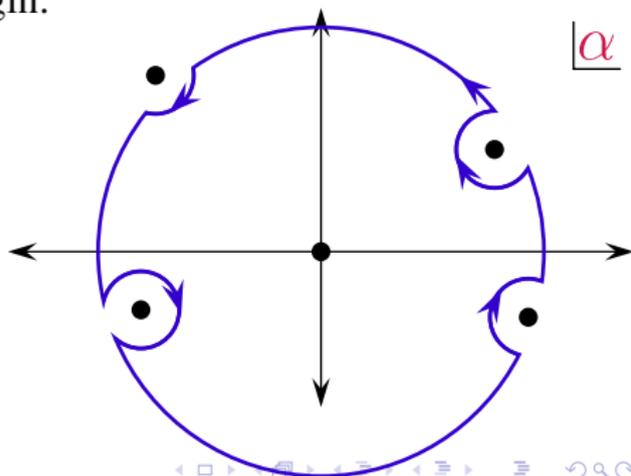
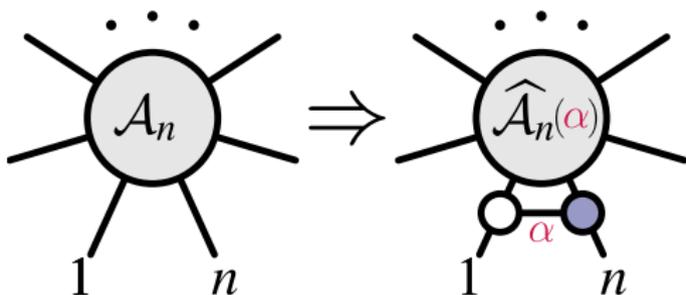


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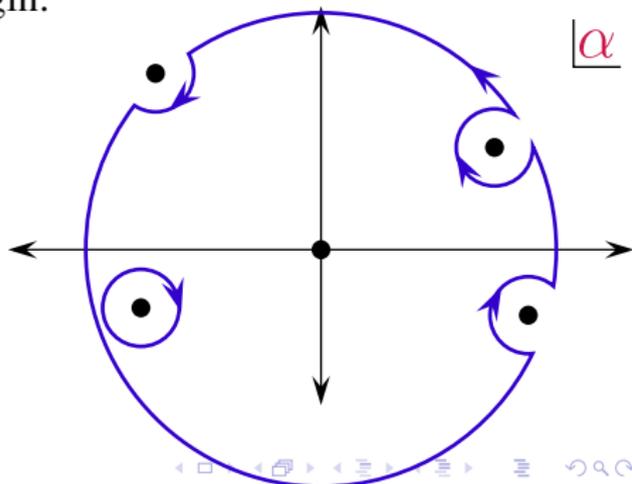
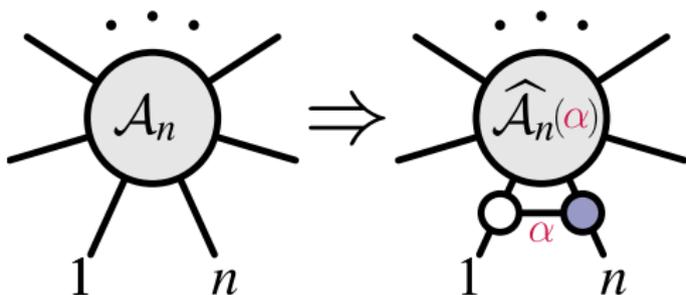


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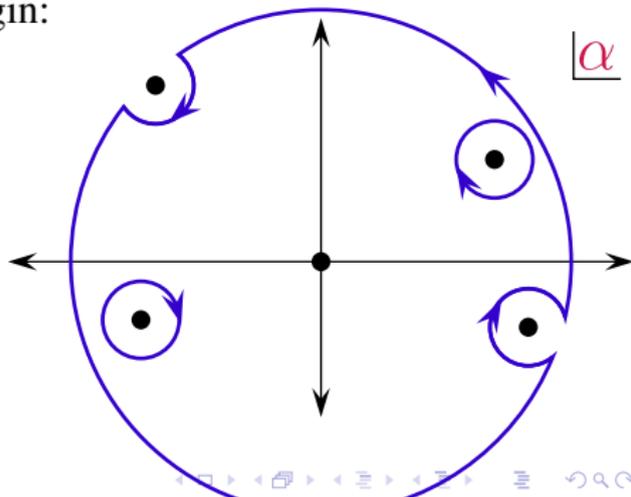
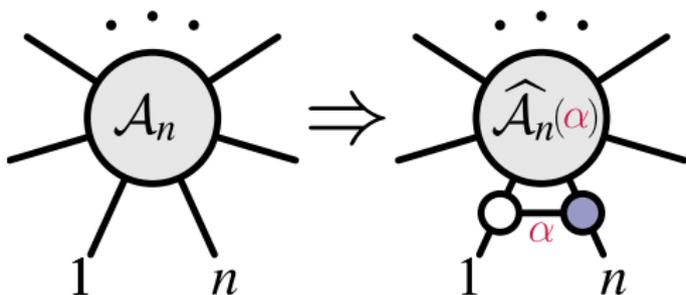


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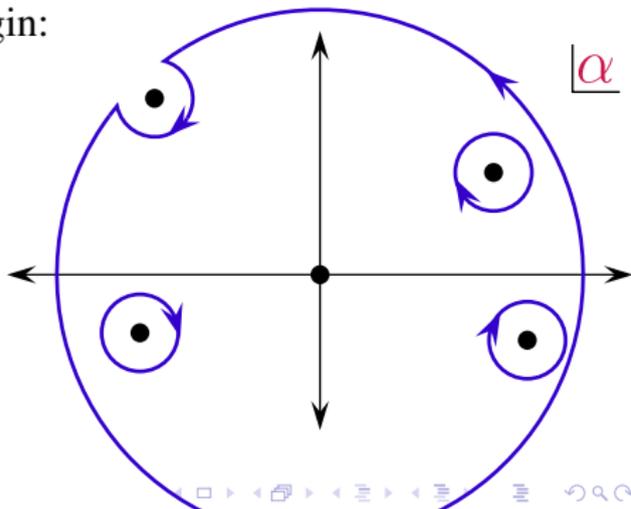
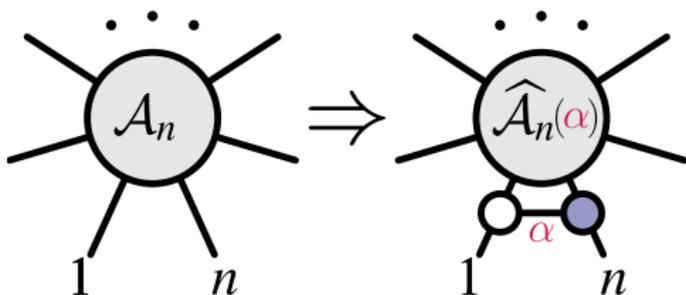


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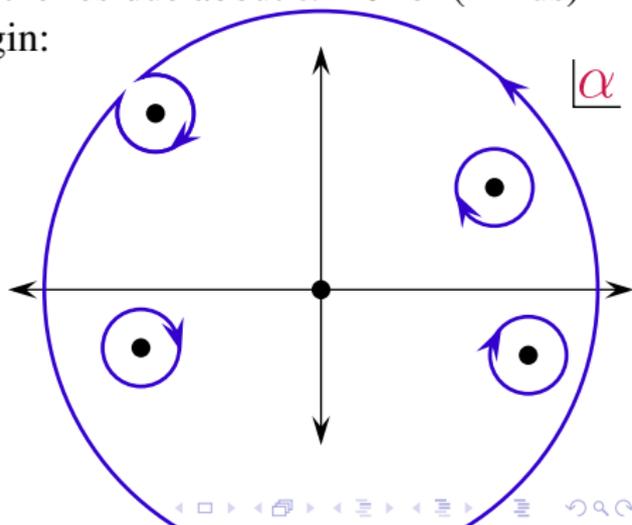
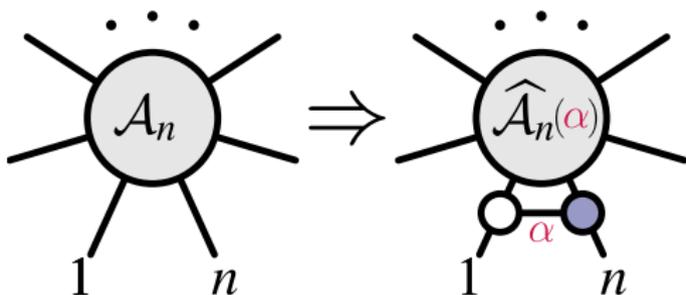


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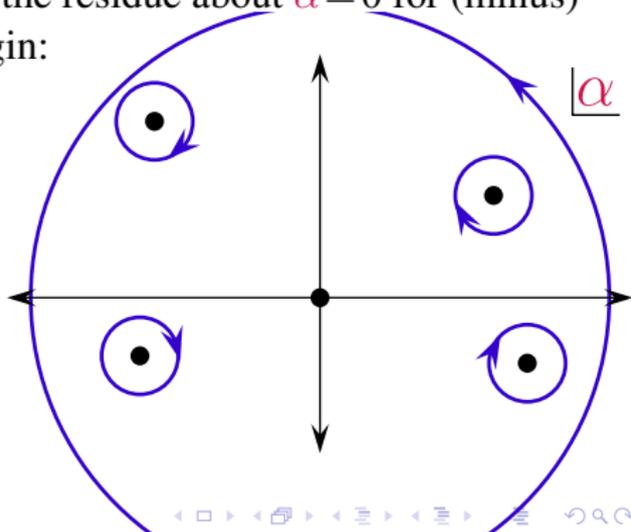
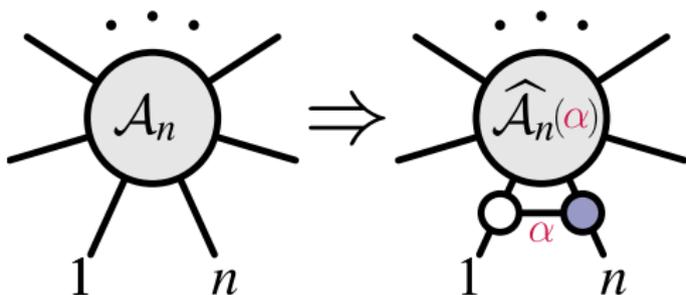


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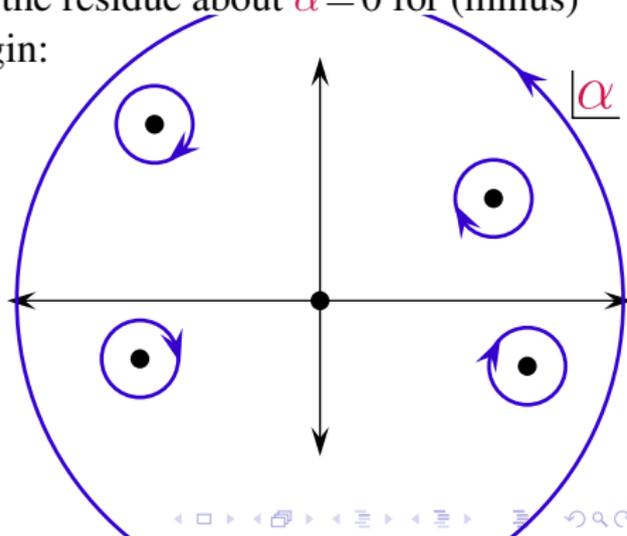
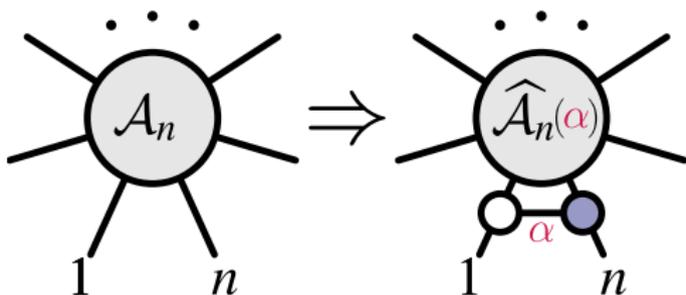


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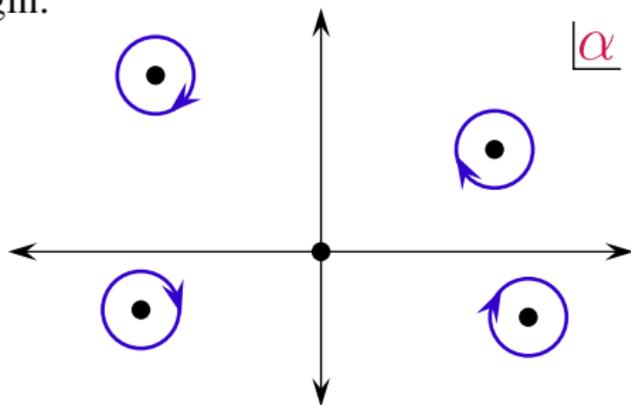
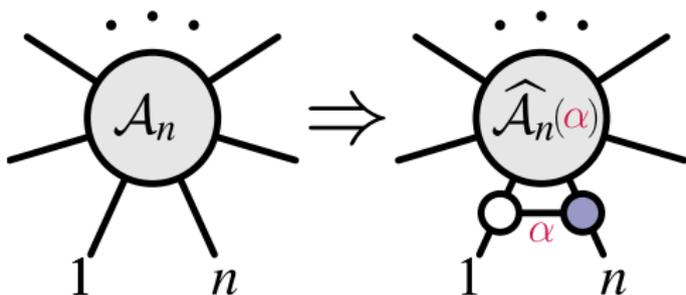


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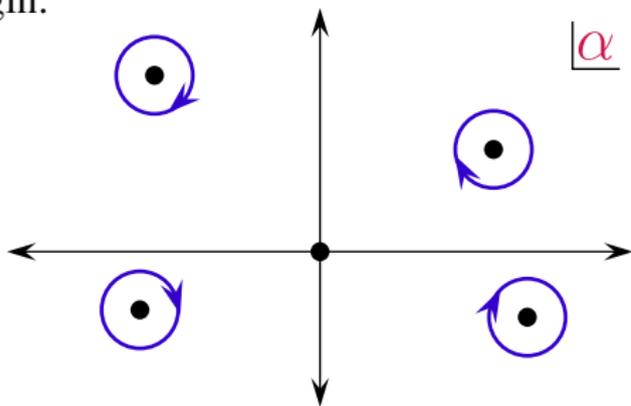
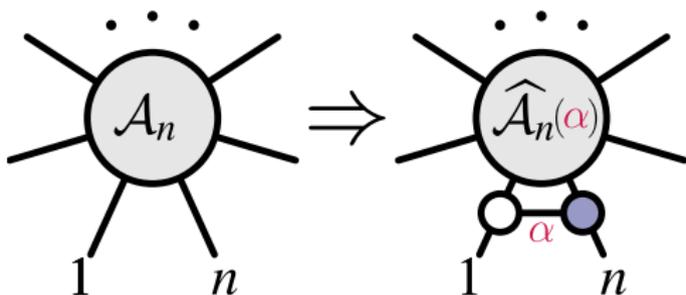


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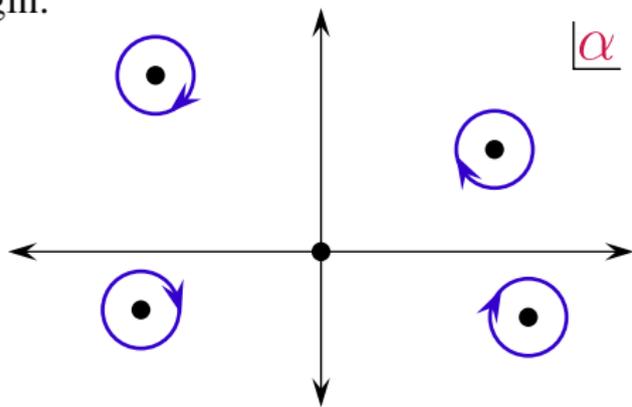
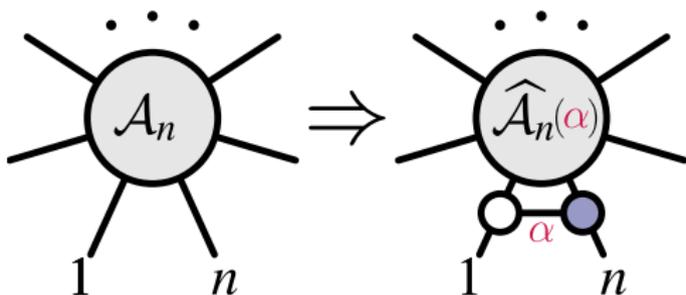


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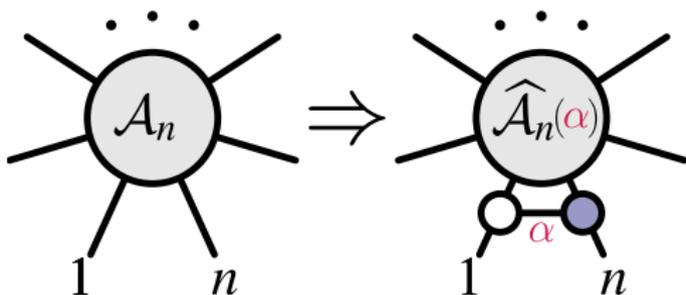


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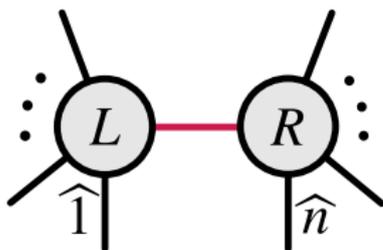


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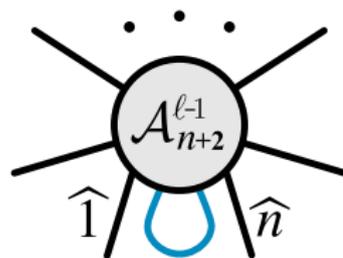
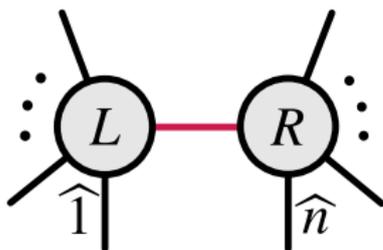


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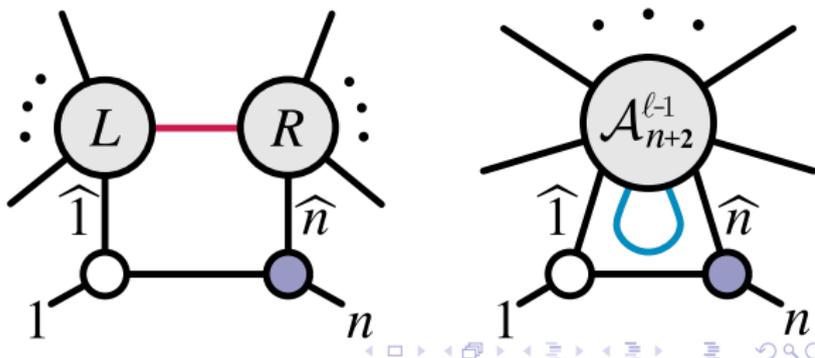


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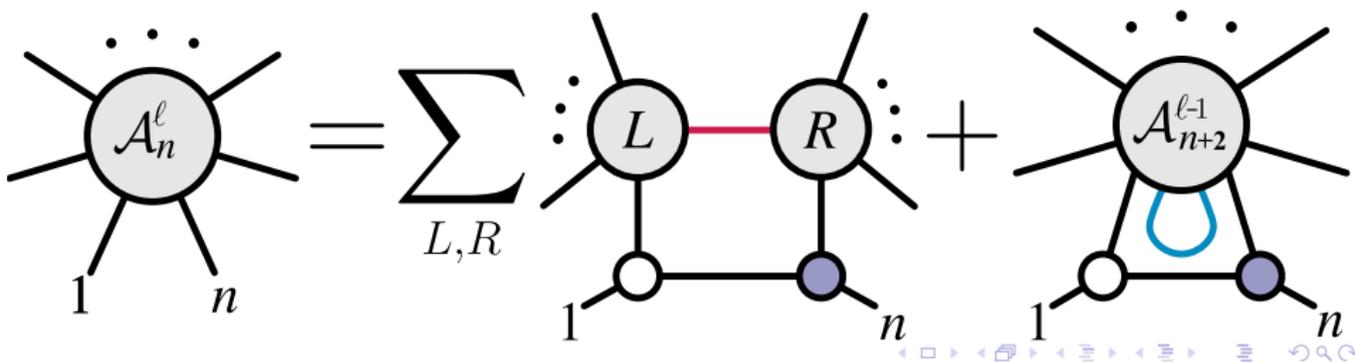


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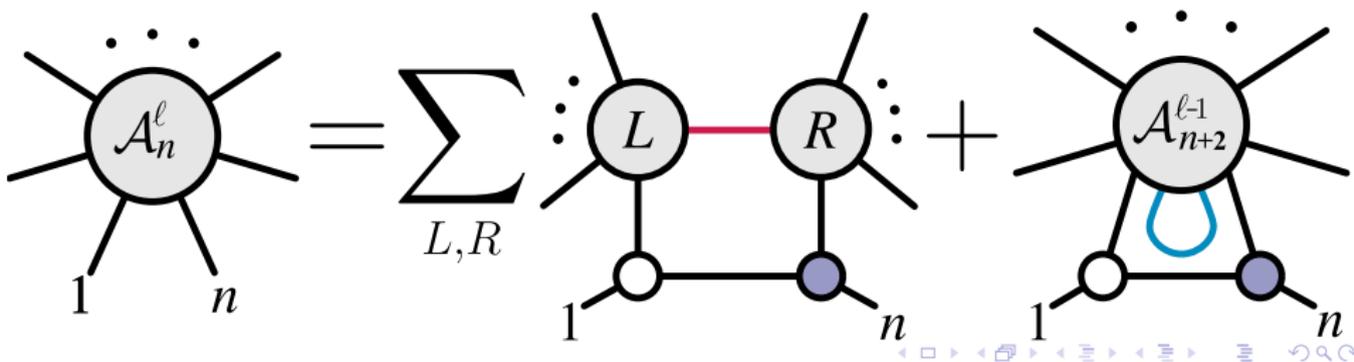
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$$\mathcal{A}_n = \widehat{\mathcal{A}}_n(\alpha \rightarrow 0) \propto \oint_{\alpha=0} \frac{d\alpha}{\alpha} \widehat{\mathcal{A}}_n(\alpha)$$

We can use **Cauchy’s theorem** to trade the residue about  $\alpha = 0$  for (minus)  
 the sum of residues away from the origin—these come in two types:  
**factorization-channels** and **forward-limits**

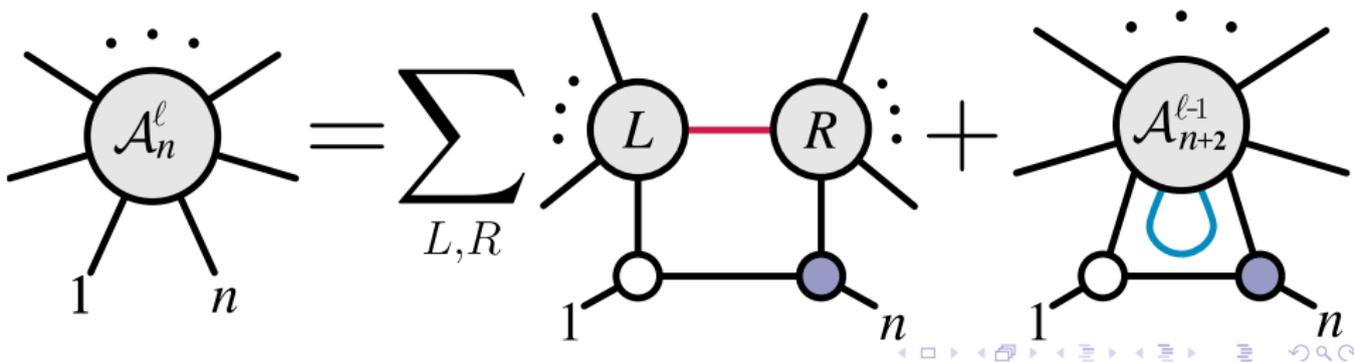


# The Analytic Boot-Strap: All-Loop Recursion Relations



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**Forward-limits** and loop-momenta:

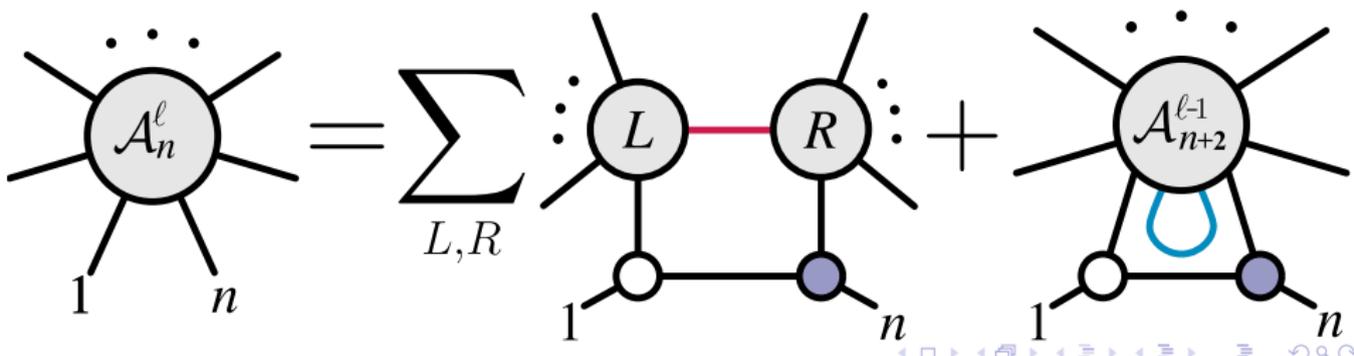


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the familiar “off-shell” loop-momentum is represented by on-shell data as:

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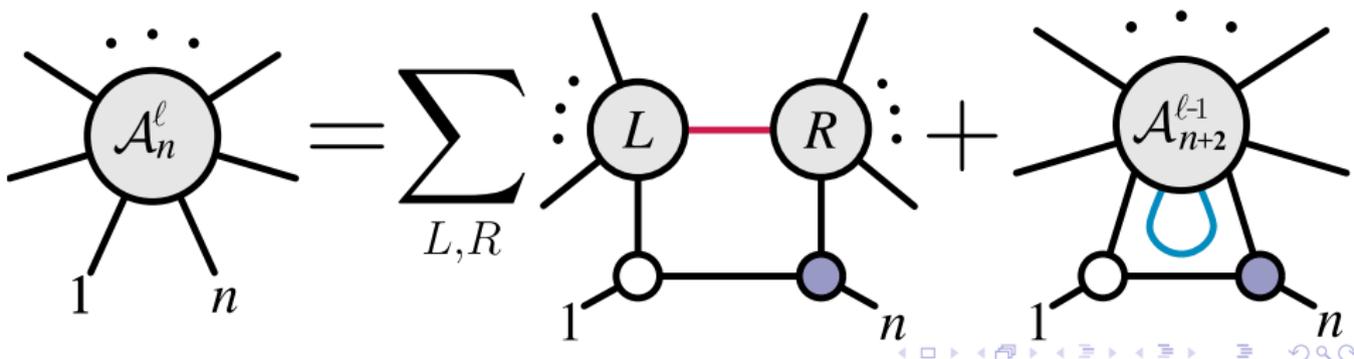


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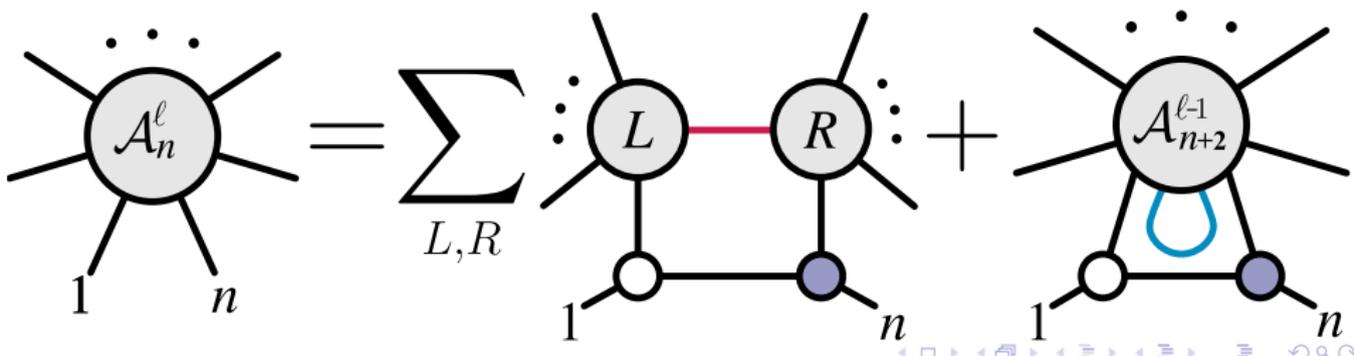


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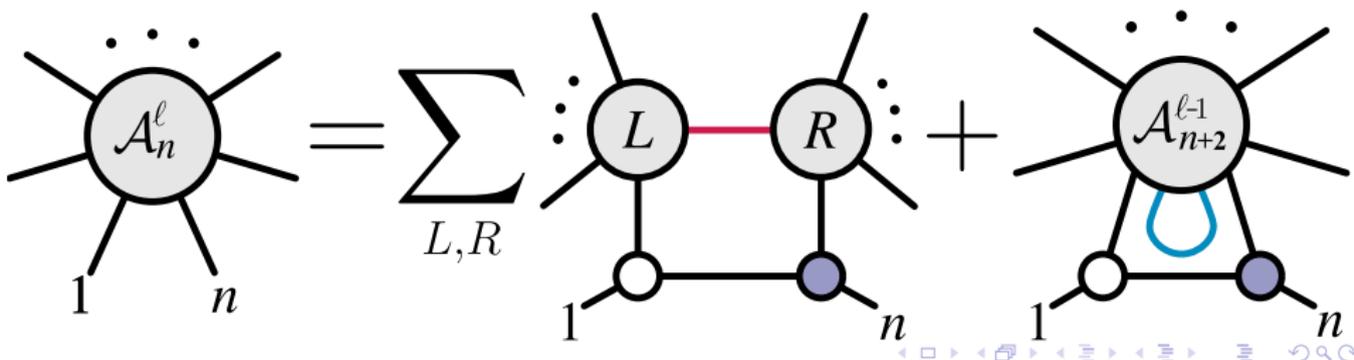


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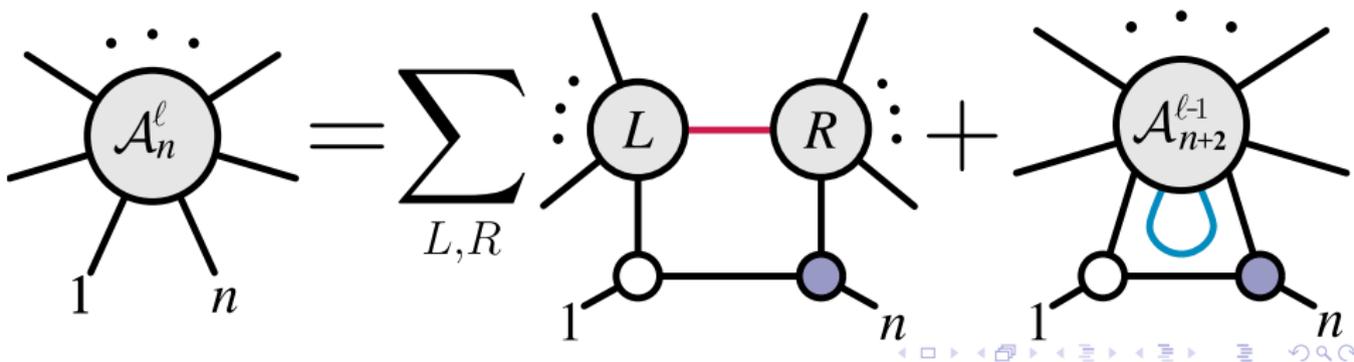


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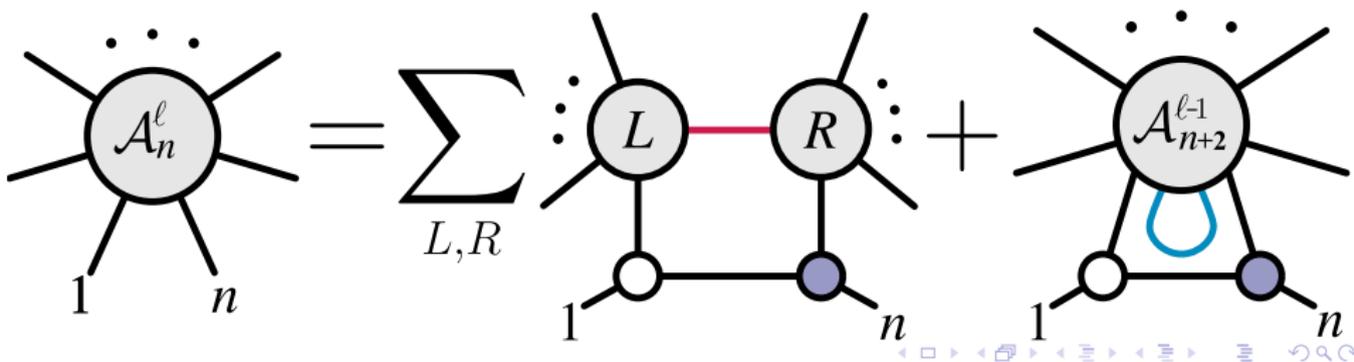


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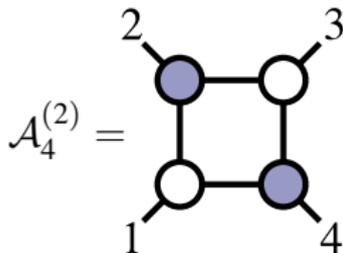
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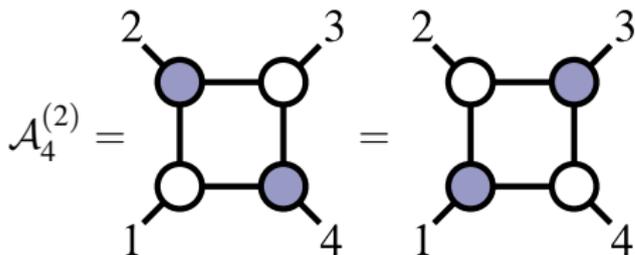
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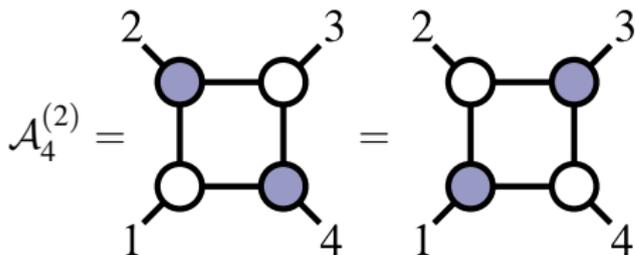
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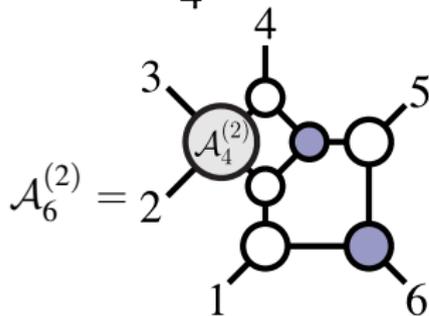
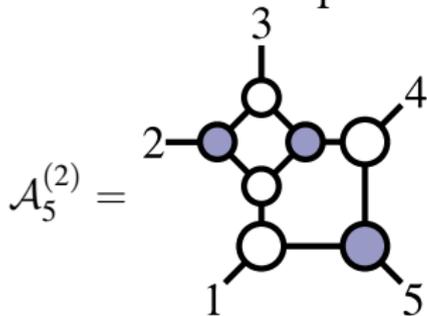
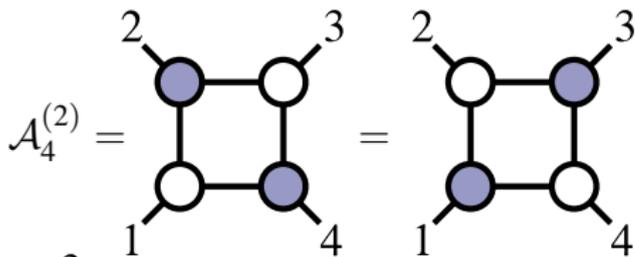
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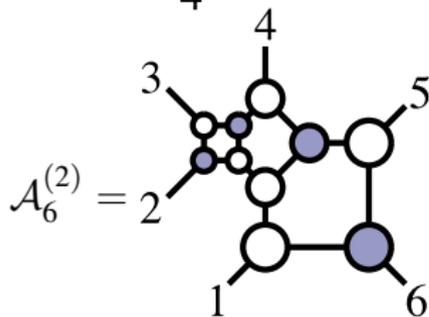
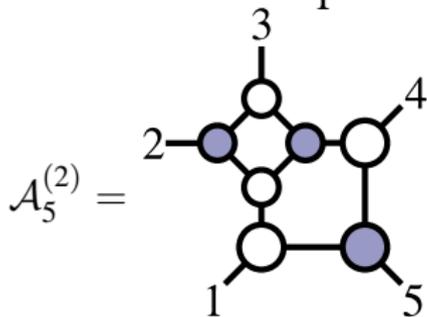
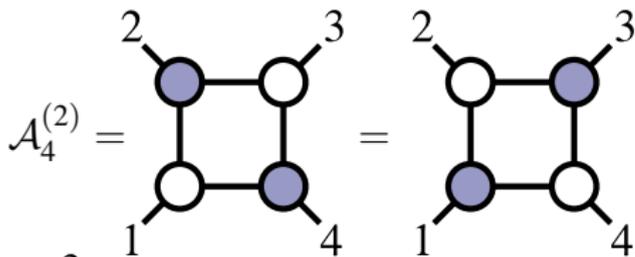
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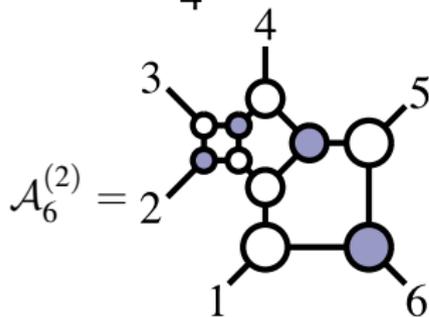
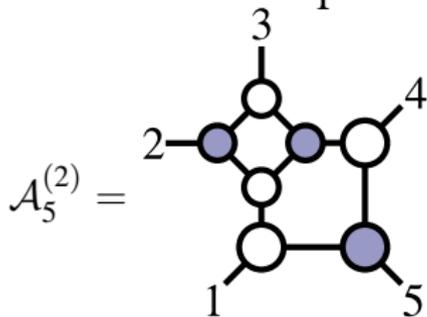
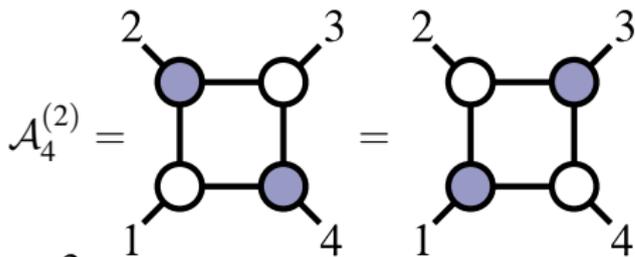
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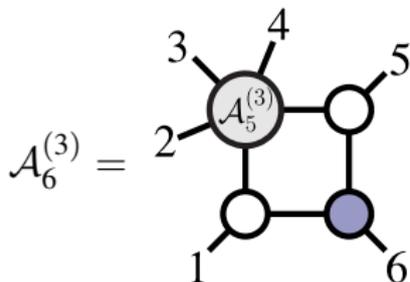
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$$\mathcal{A}_6^{(3)} = \text{Diagram 1} + \text{Diagram 2}$$

The diagram shows the equation  $\mathcal{A}_6^{(3)} = \text{Diagram 1} + \text{Diagram 2}$ .  
 Diagram 1: A square loop with four vertices. The top-left vertex is a grey circle labeled  $\mathcal{A}_5^{(3)}$  with external legs 2, 3, and 4. The top-right vertex is a white circle with external leg 5. The bottom-left vertex is a white circle with external leg 1. The bottom-right vertex is a blue circle with external leg 6. The legs are connected in a square loop.  
 Diagram 2: A square loop with four vertices. The top-left vertex is a grey circle labeled  $\mathcal{A}_4^{(2)}$  with external legs 2 and 3. The top-right vertex is a grey circle labeled  $\mathcal{A}_4^{(2)}$  with external legs 4 and 5. The bottom-left vertex is a white circle with external leg 1. The bottom-right vertex is a blue circle with external leg 6. The legs are connected in a square loop.

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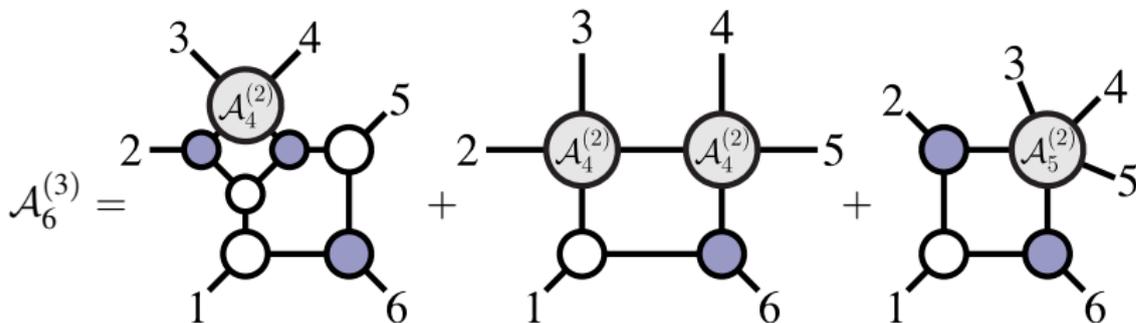
$$\mathcal{A}_6^{(3)} = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3}$$

The diagrammatic equation shows the 6-point tree amplitude  $\mathcal{A}_6^{(3)}$  as a sum of three terms. The first term is a square loop with a shaded blue vertex at the bottom-right, labeled  $\mathcal{A}_5^{(3)}$  at the top-left vertex. The second term is a chain of two shaded blue vertices, each labeled  $\mathcal{A}_4^{(2)}$ , with external legs 1, 2, 3, 4, 5, 6. The third term is a chain of two shaded blue vertices, each labeled  $\mathcal{A}_5^{(2)}$ , with external legs 1, 2, 3, 4, 5, 6.

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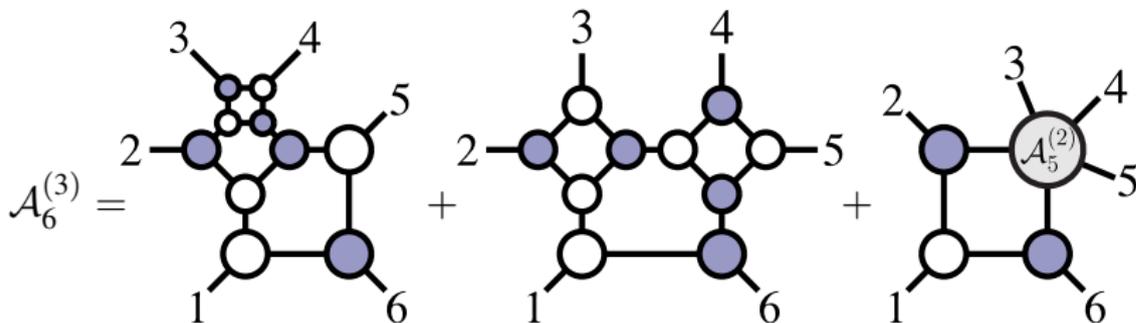
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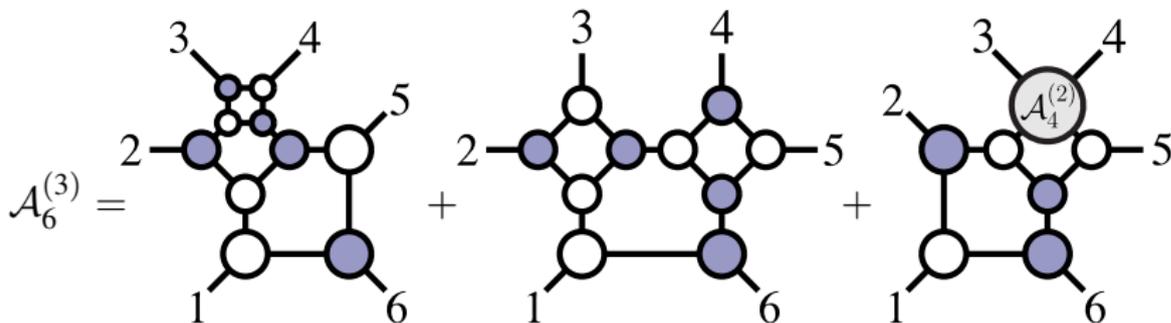
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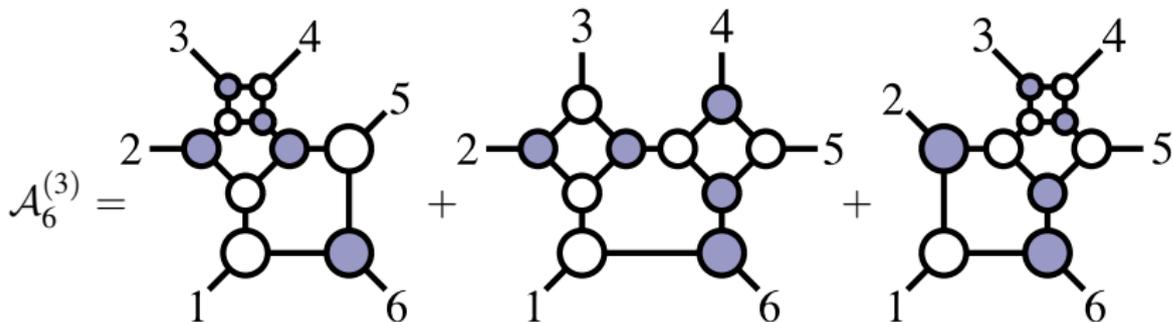
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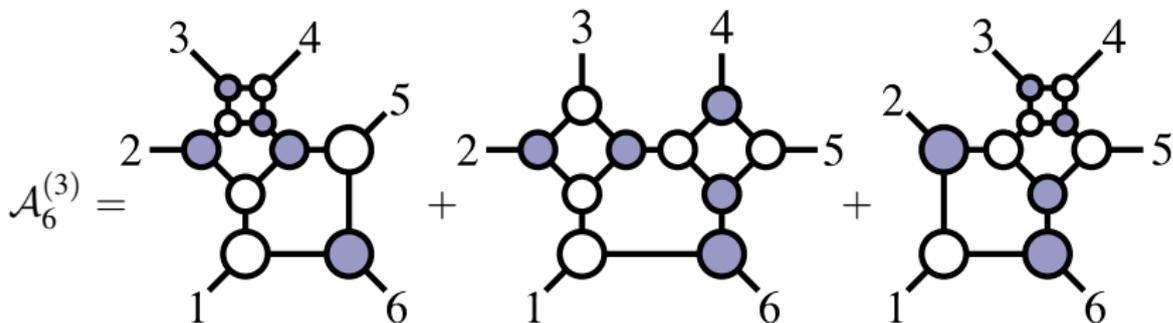
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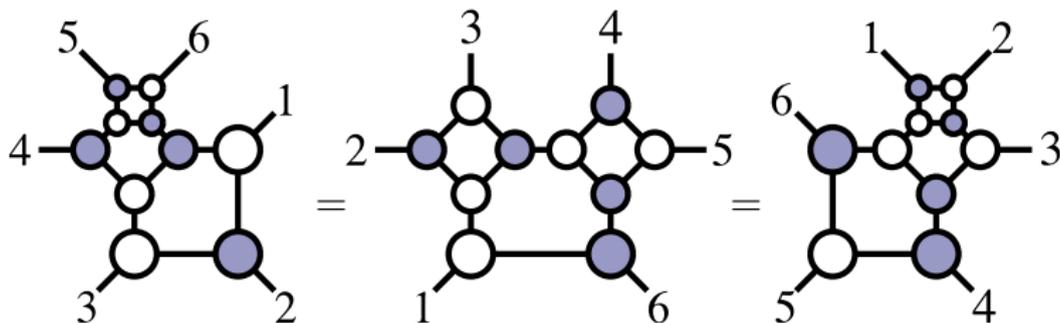
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The image shows three Feynman diagrams for the six-point tree-level amplitude  $\mathcal{A}_6^{(3)}$ . Each diagram has external legs labeled 1 through 6. The diagrams are summed together to give the total amplitude. The diagrams are:

- Diagram 1: A central loop with a bridge between two vertices. The bridge connects the two vertices on the right side of the loop.
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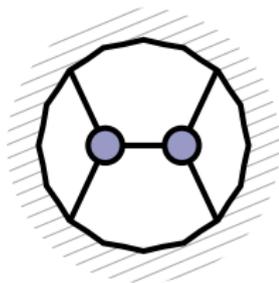
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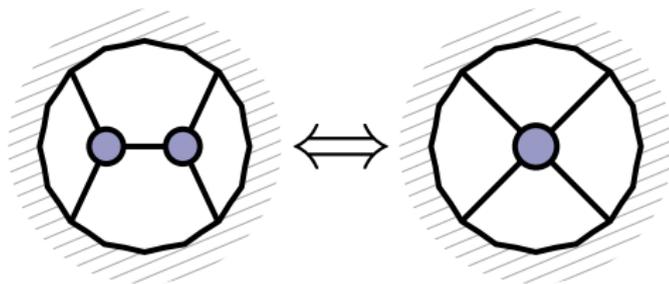
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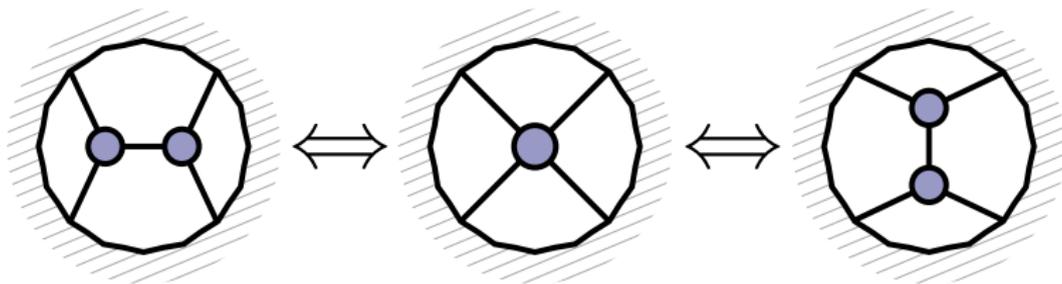
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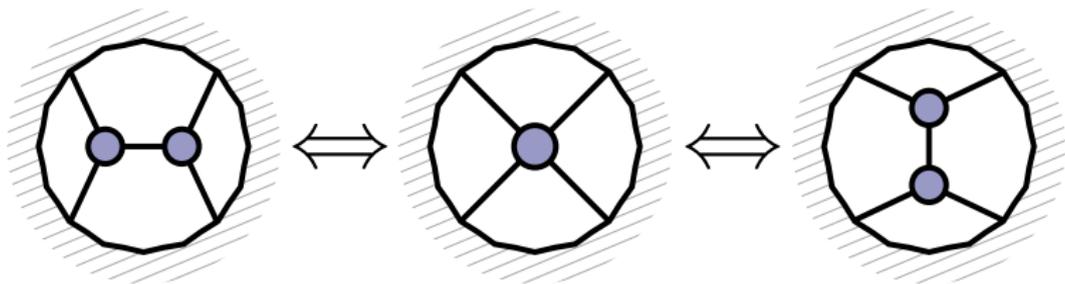
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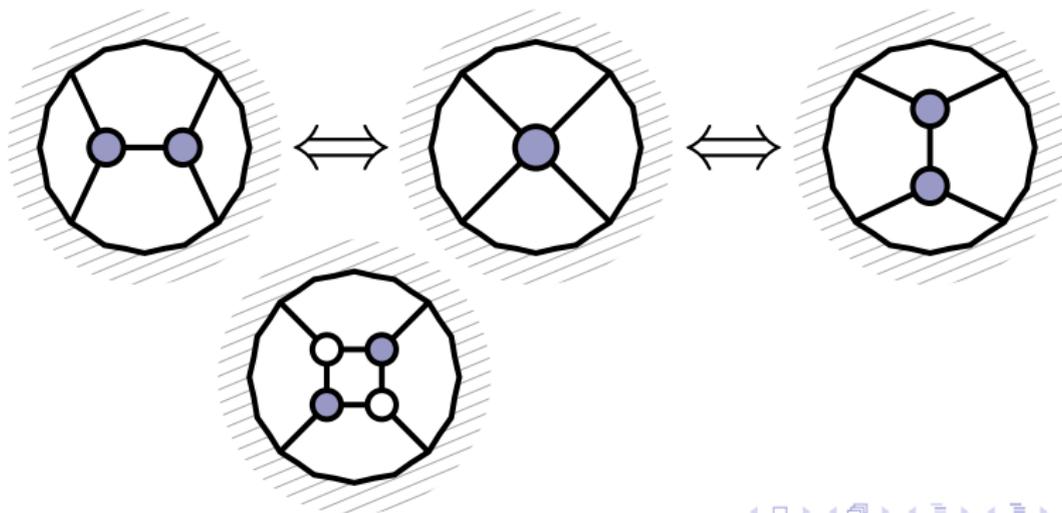
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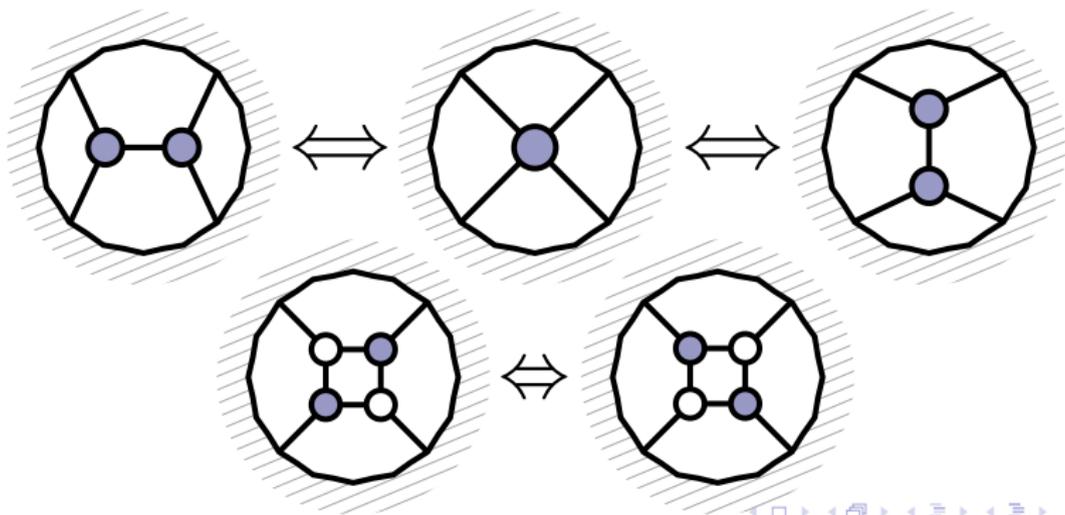
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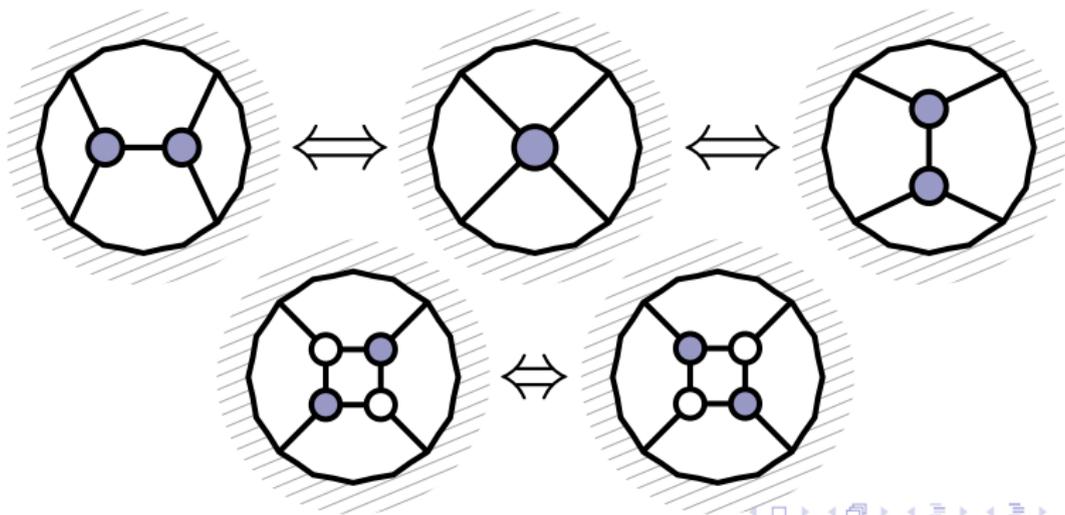
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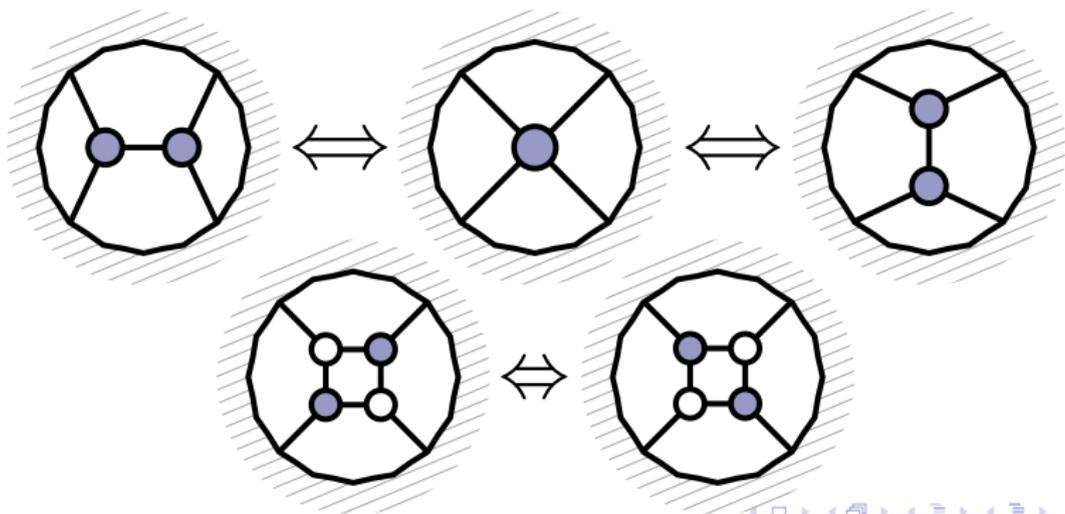
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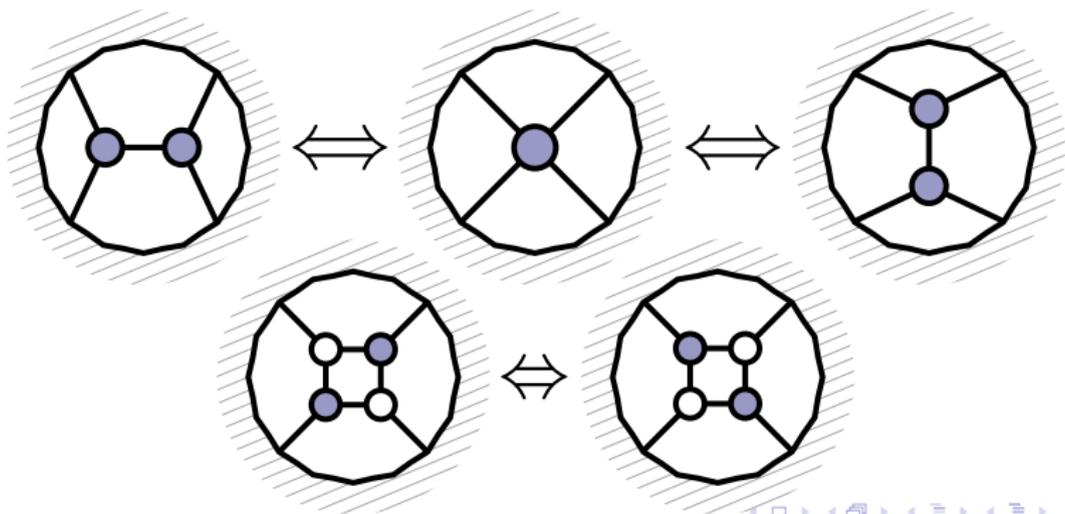


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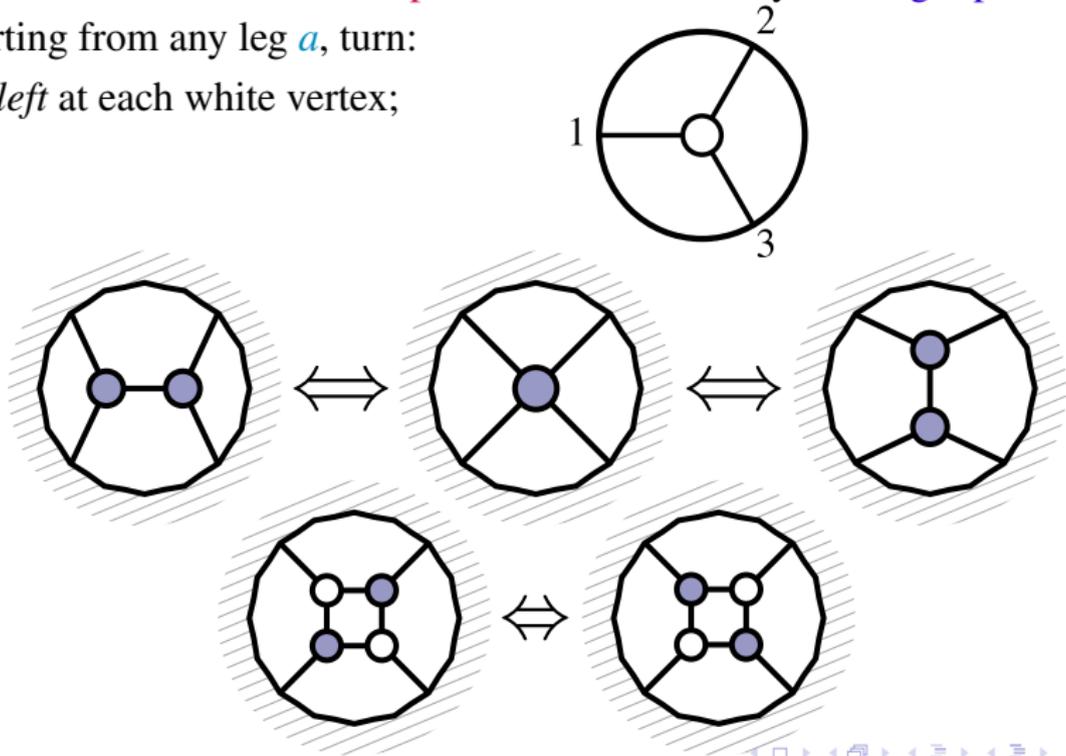


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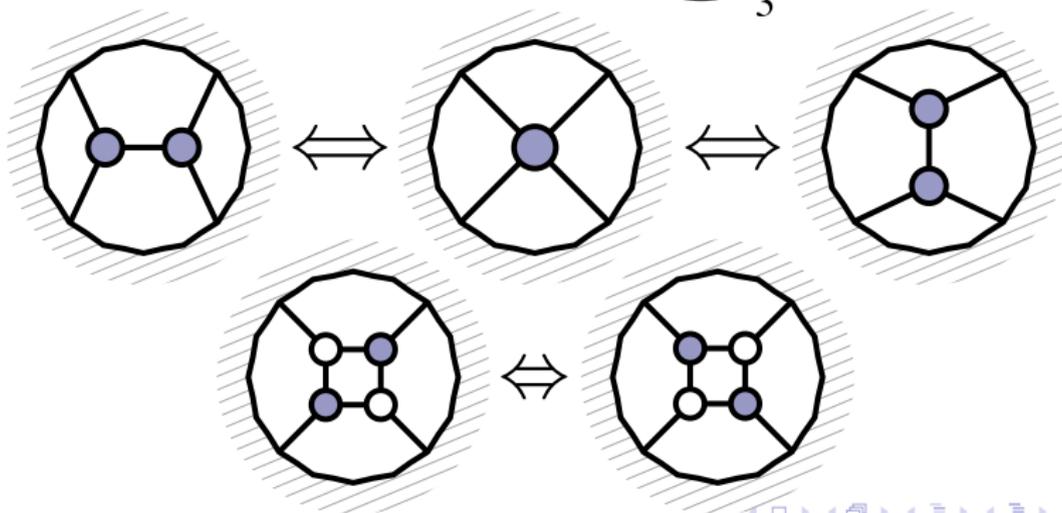
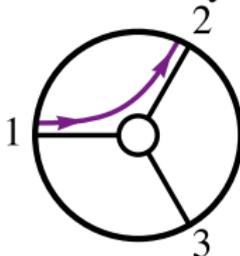


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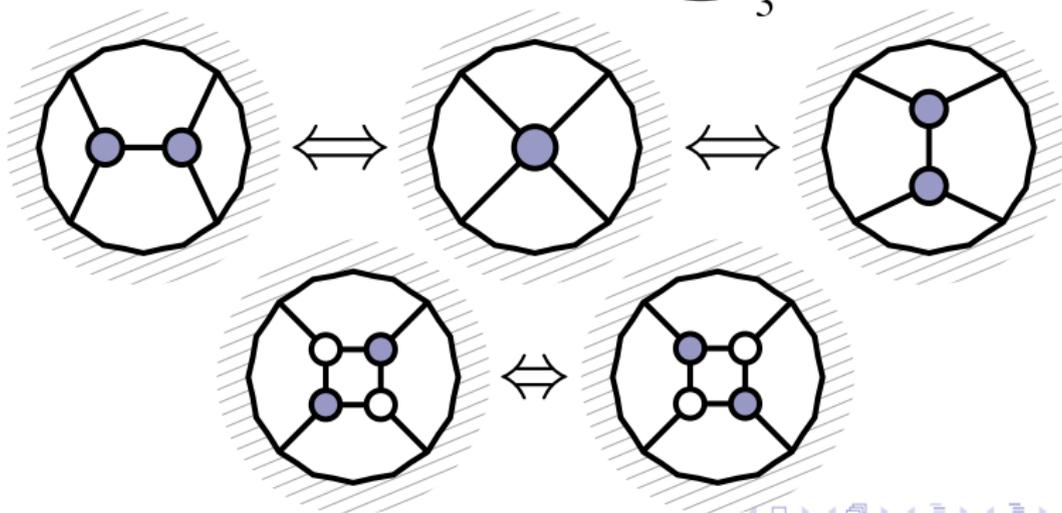
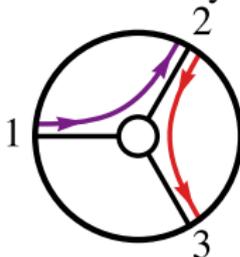


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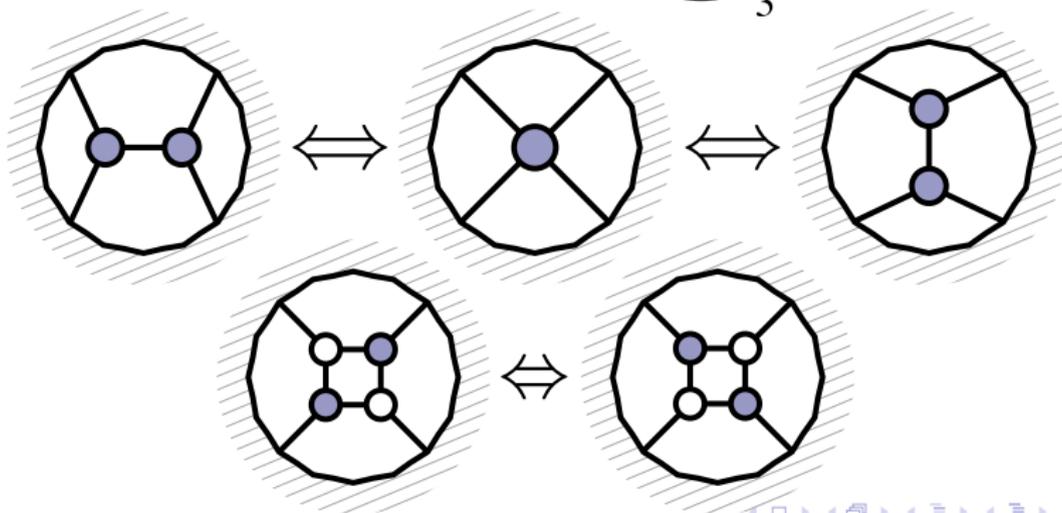
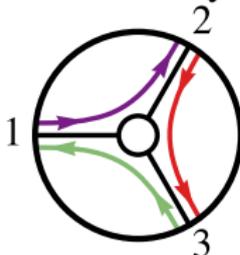


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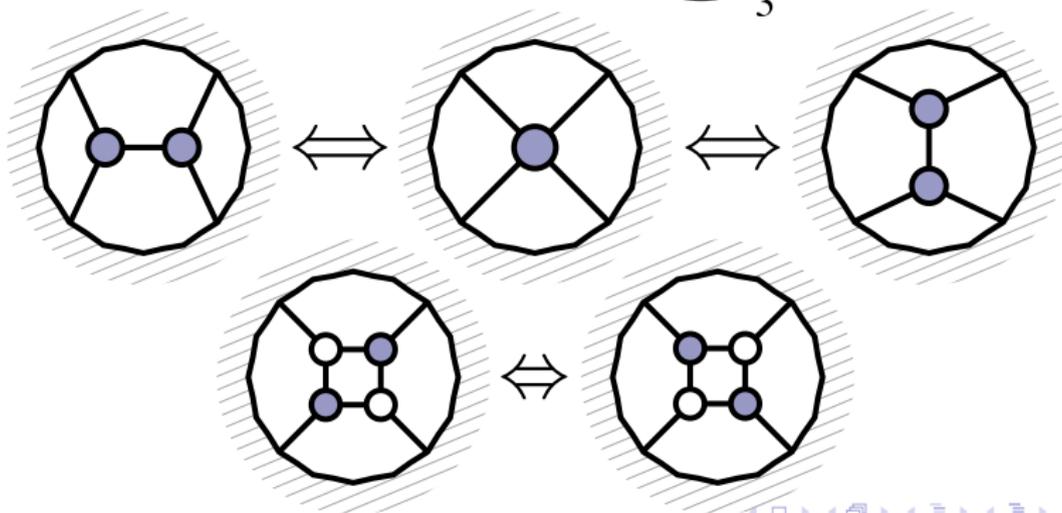
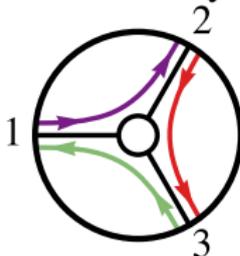


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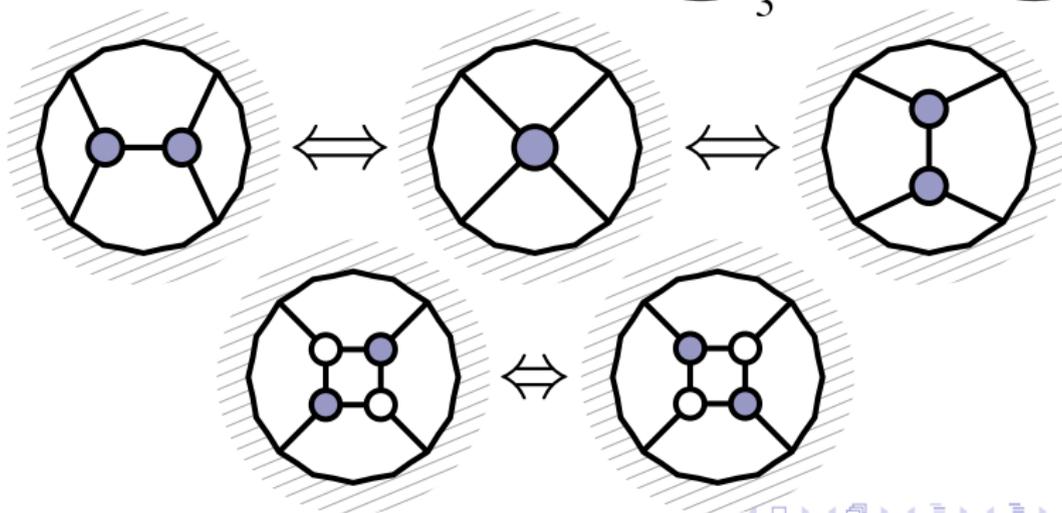
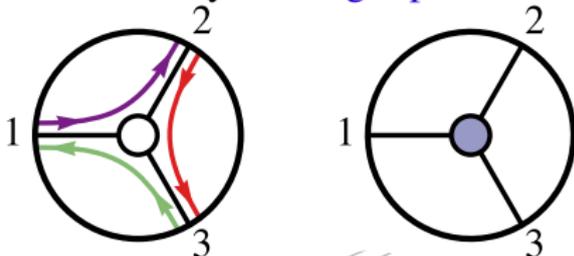


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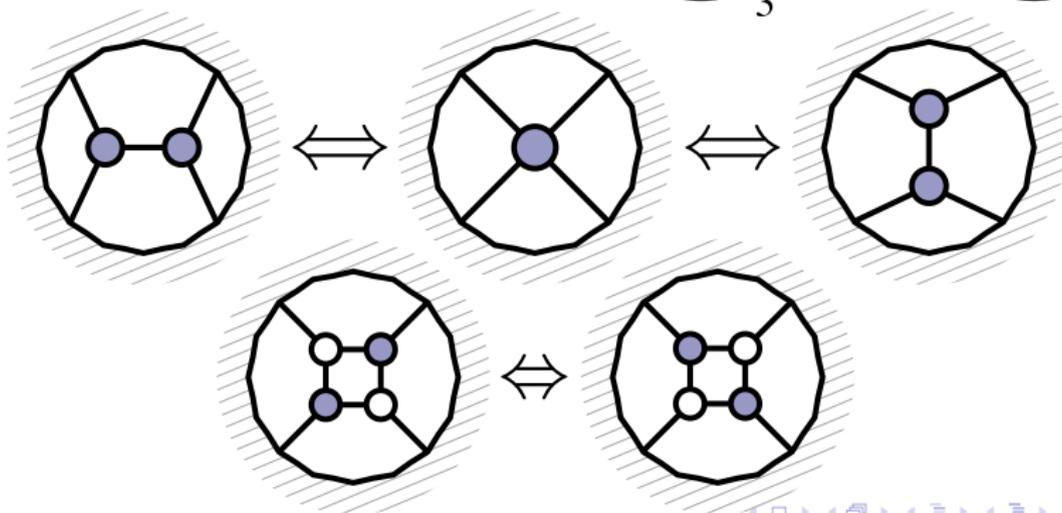
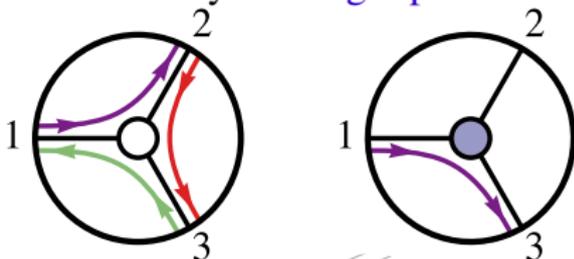


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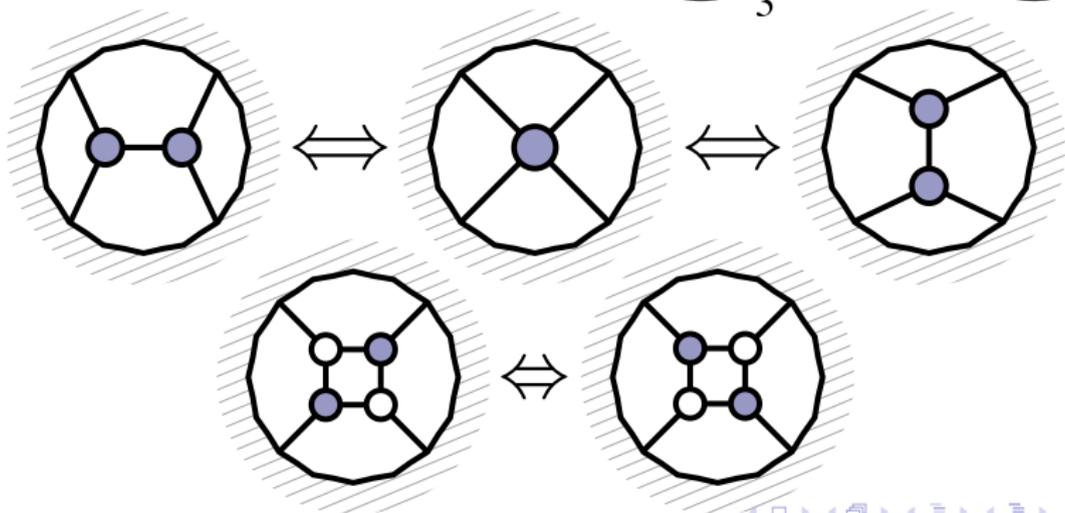
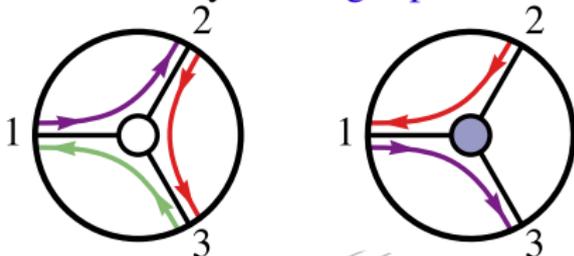


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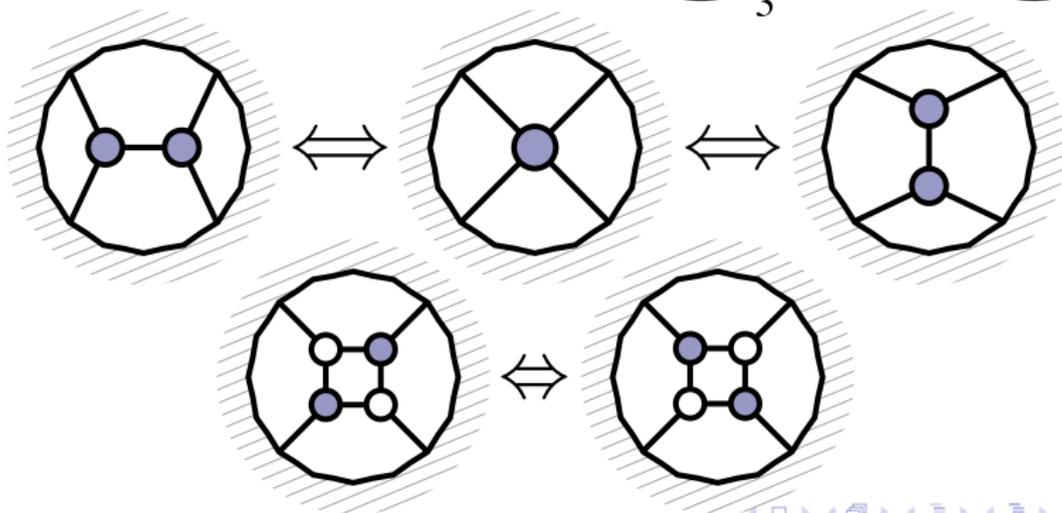
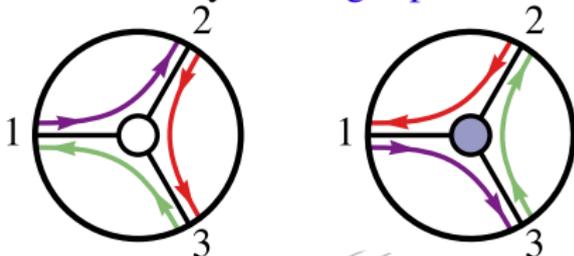


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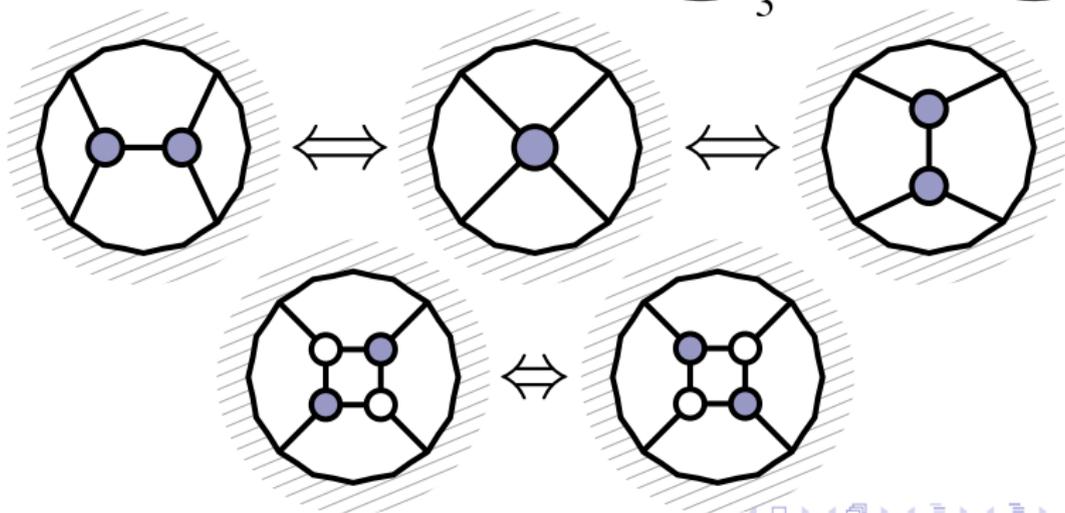
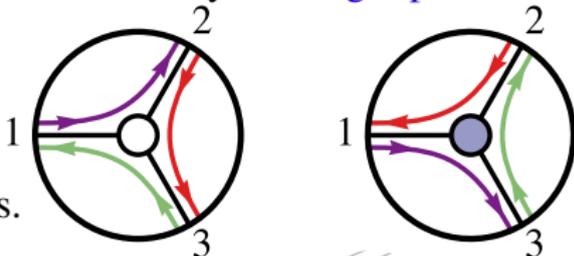
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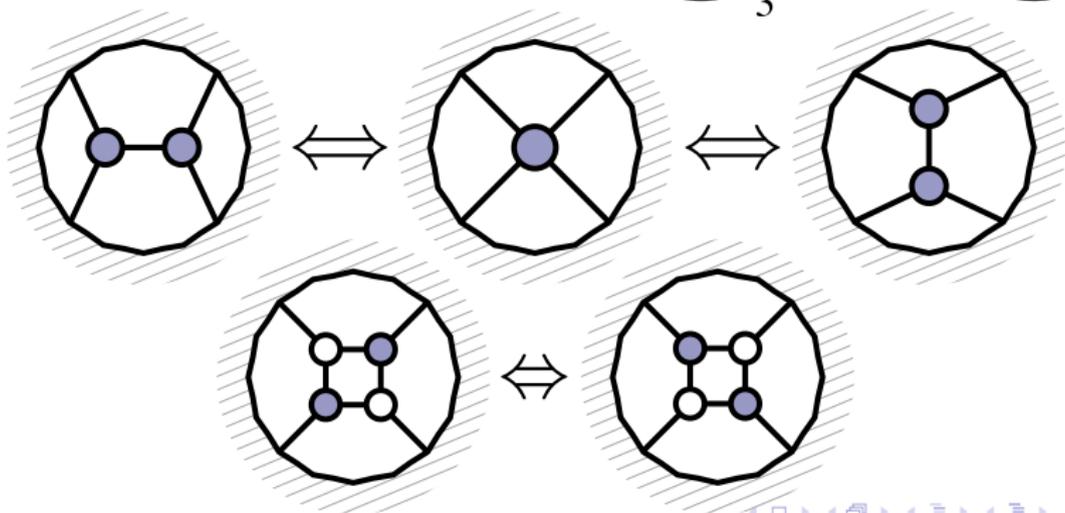
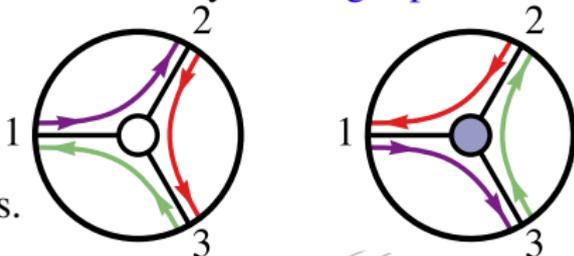
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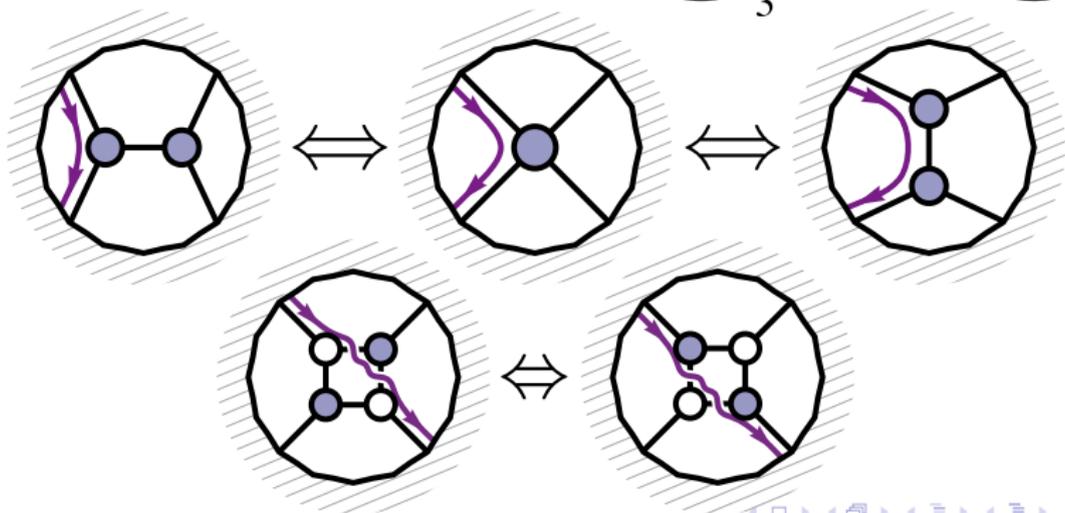
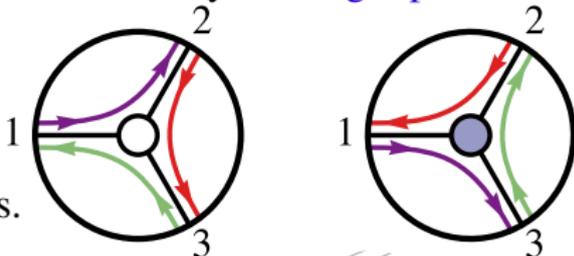
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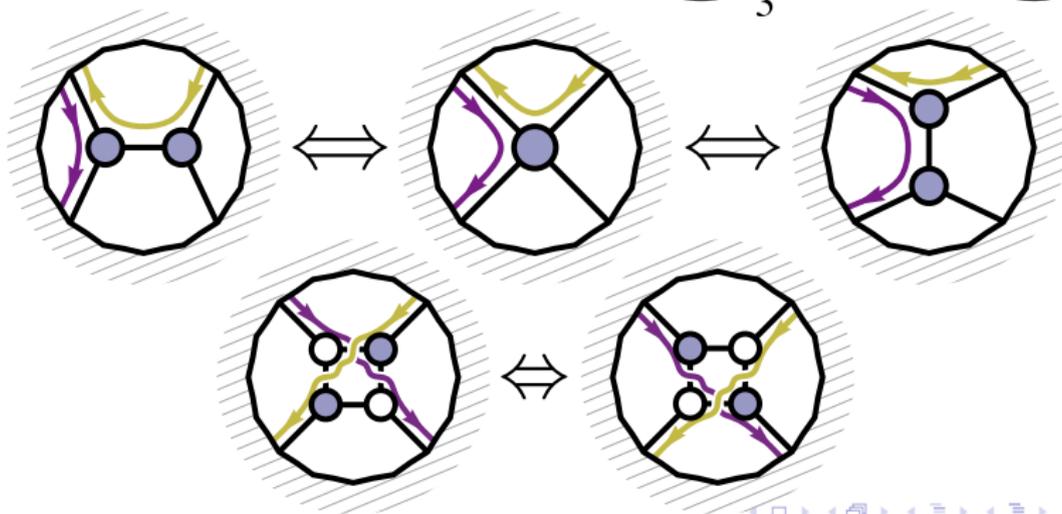
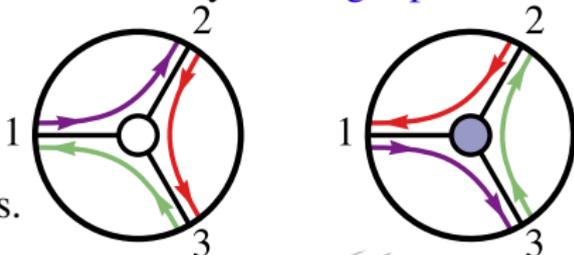
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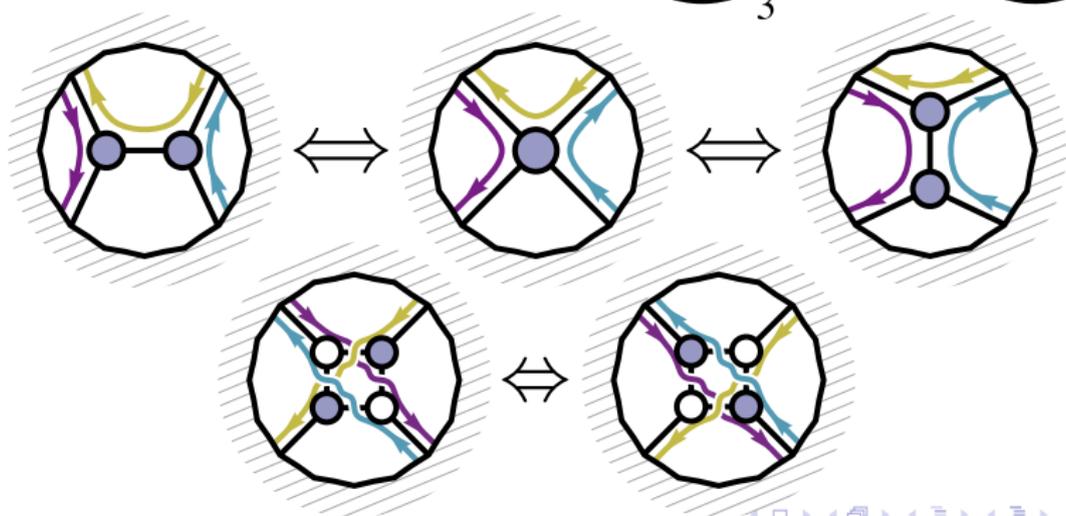
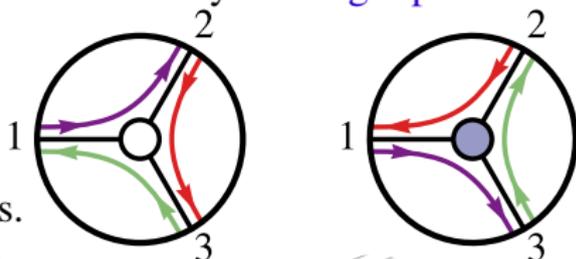
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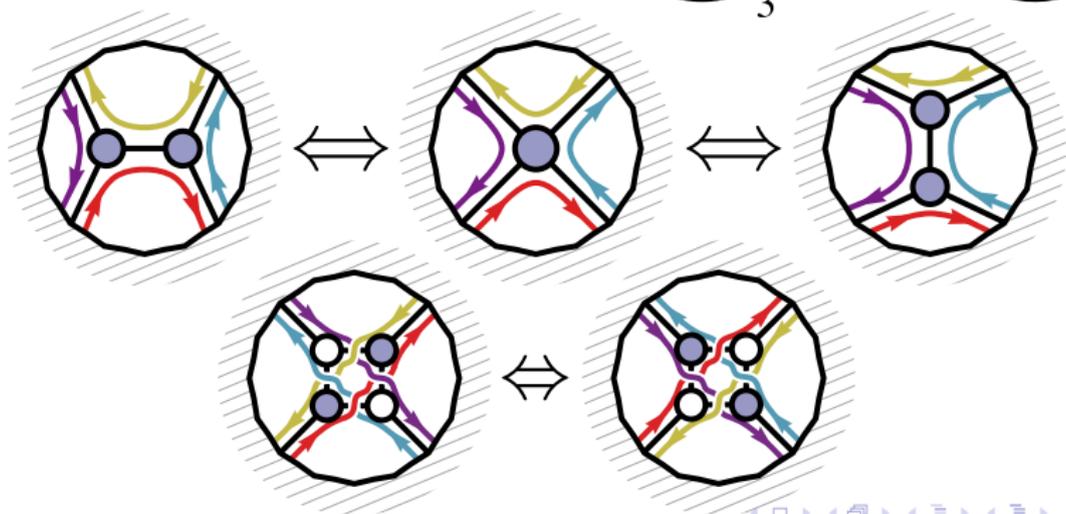
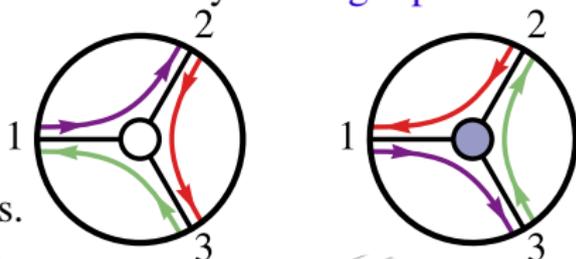
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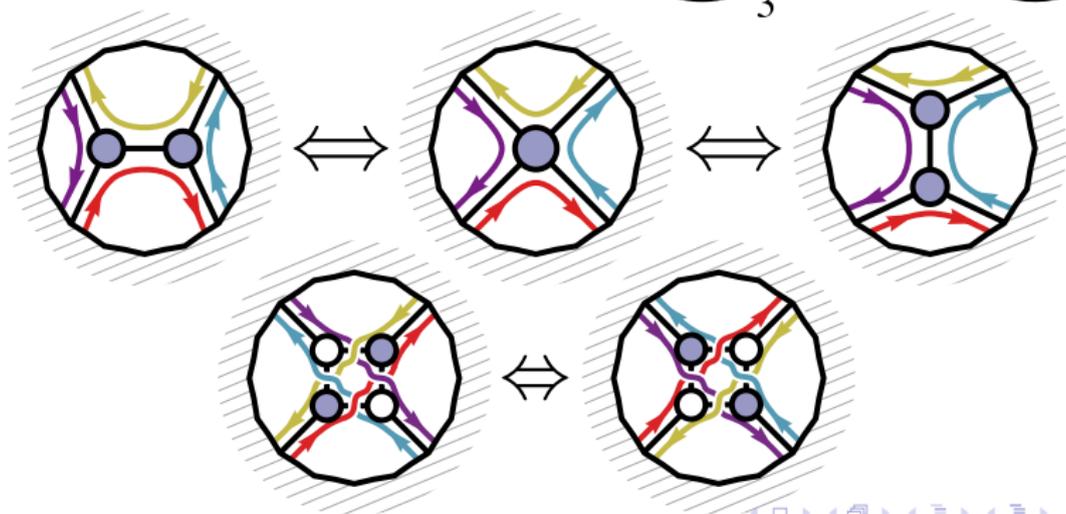
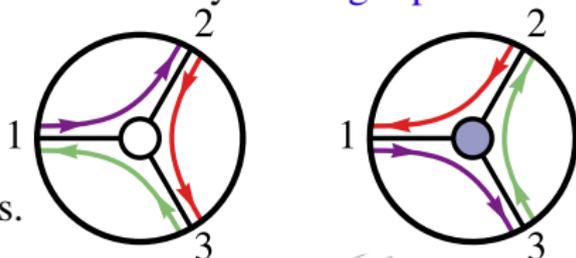
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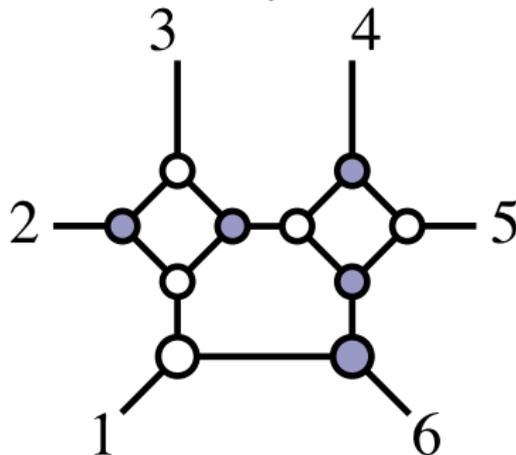
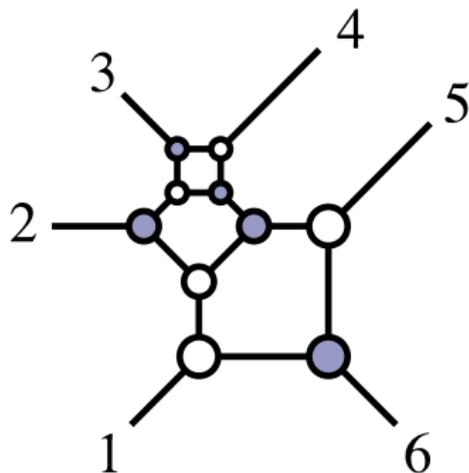
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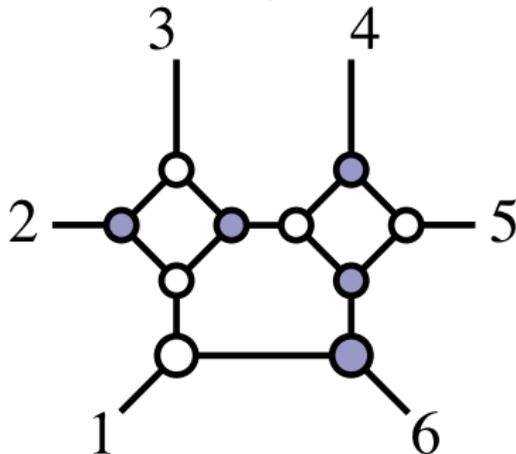
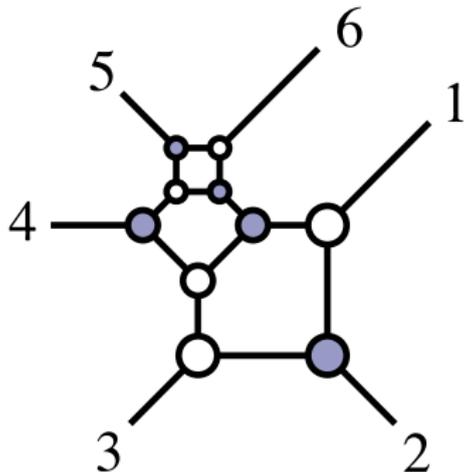
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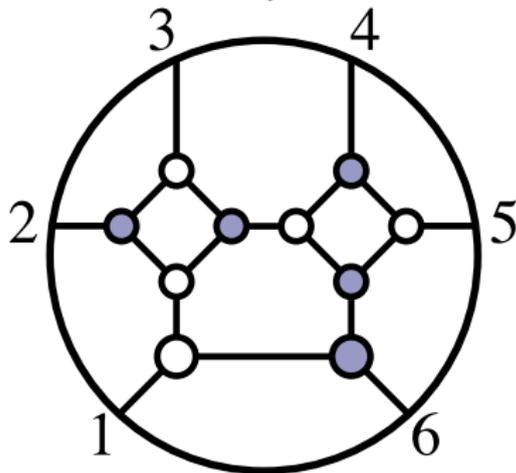
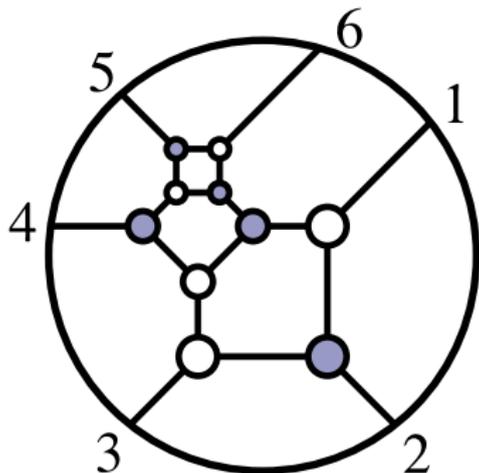
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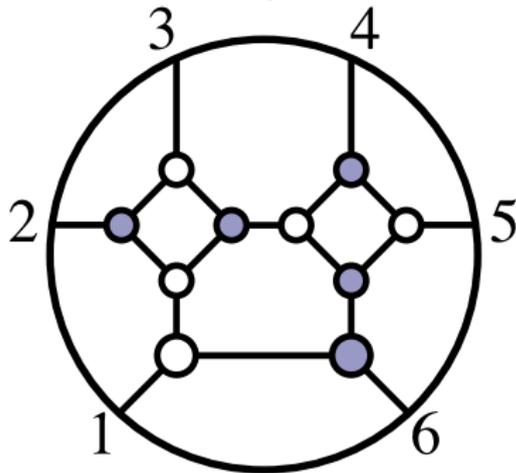
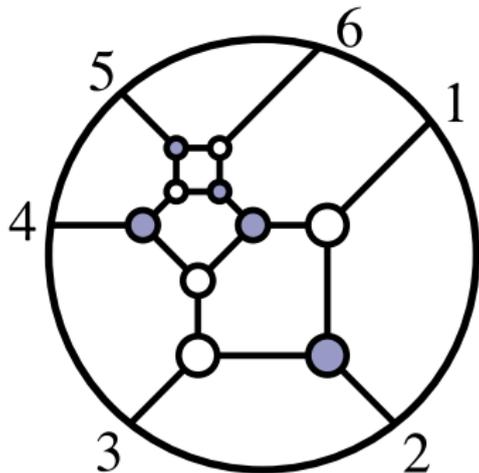
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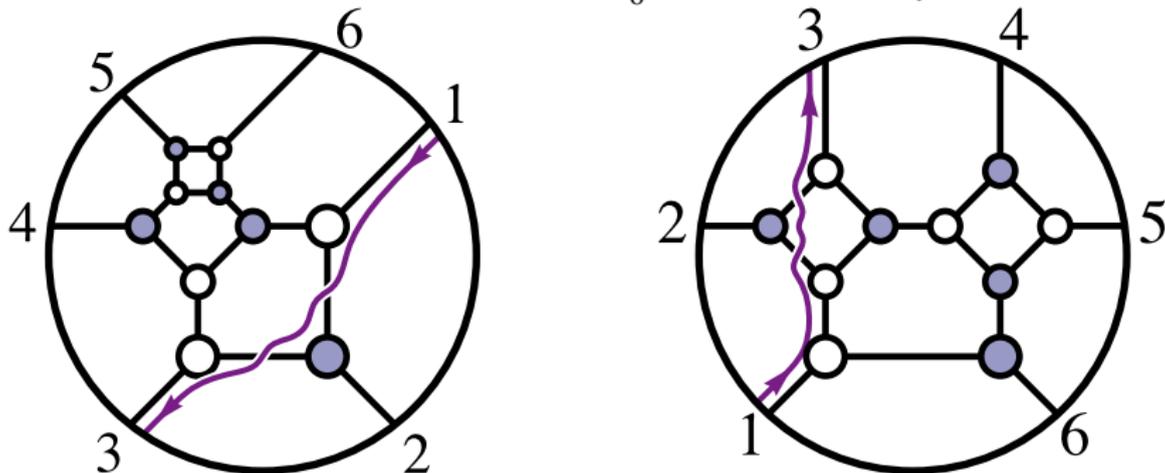
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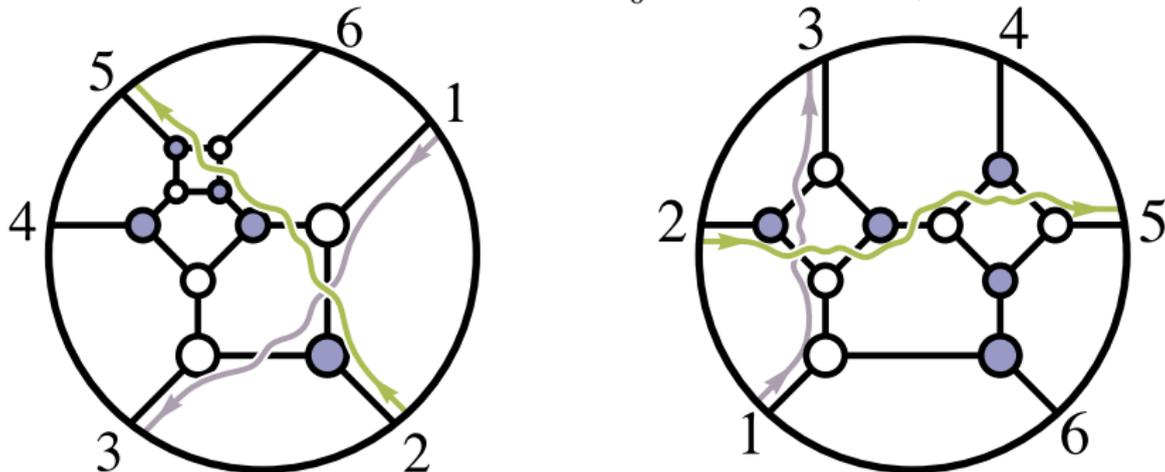
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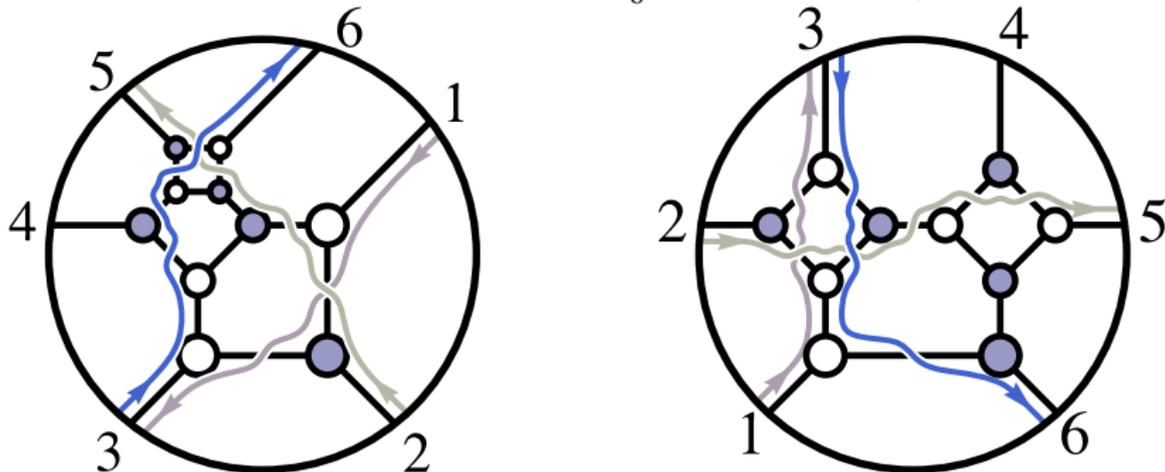
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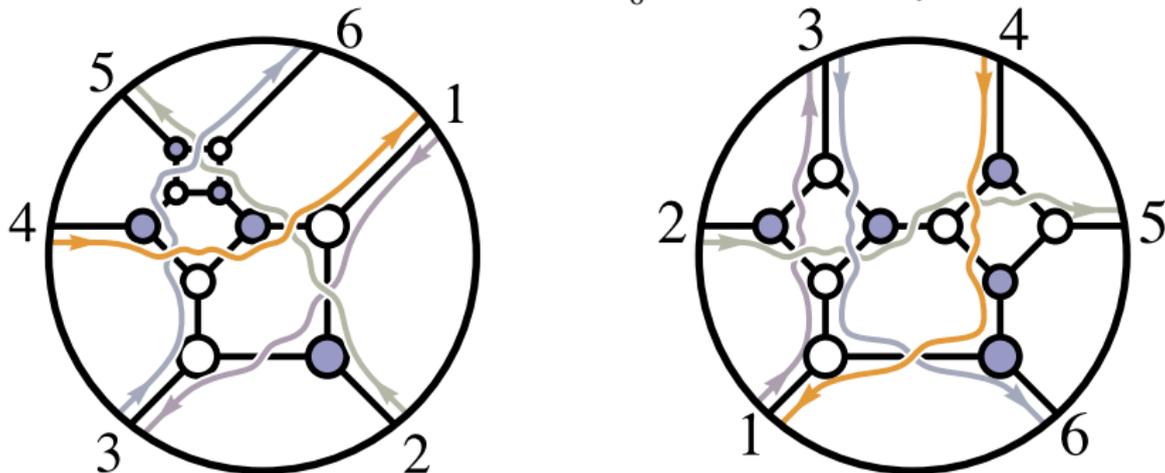
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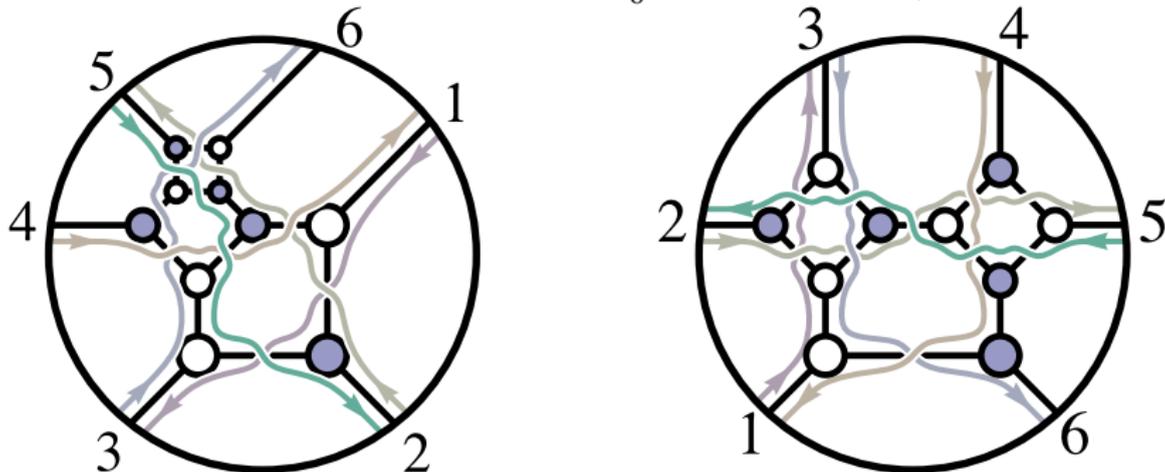
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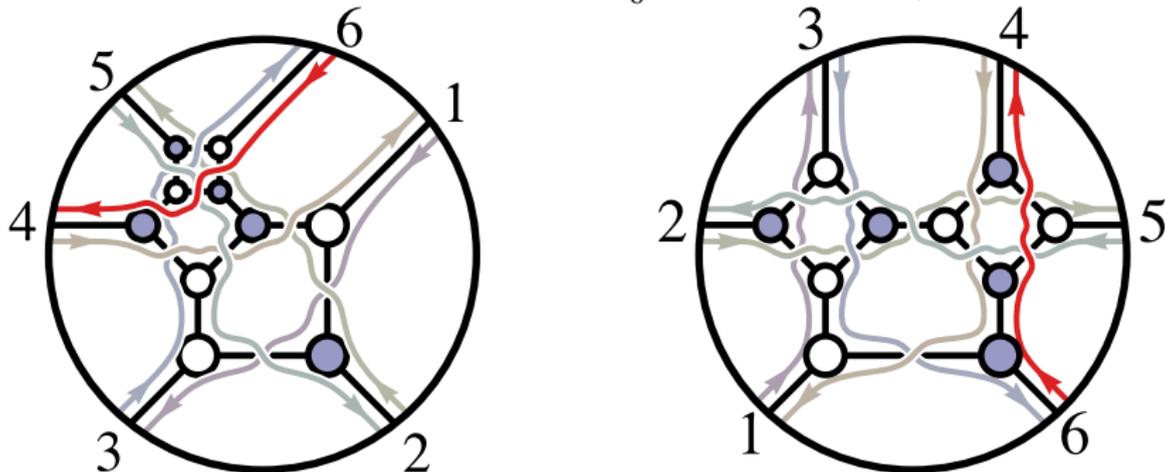
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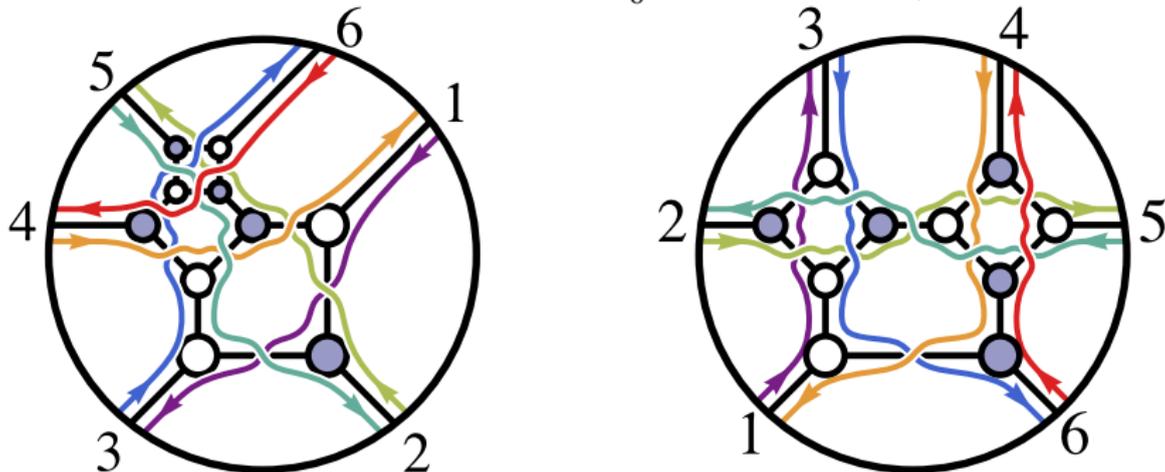
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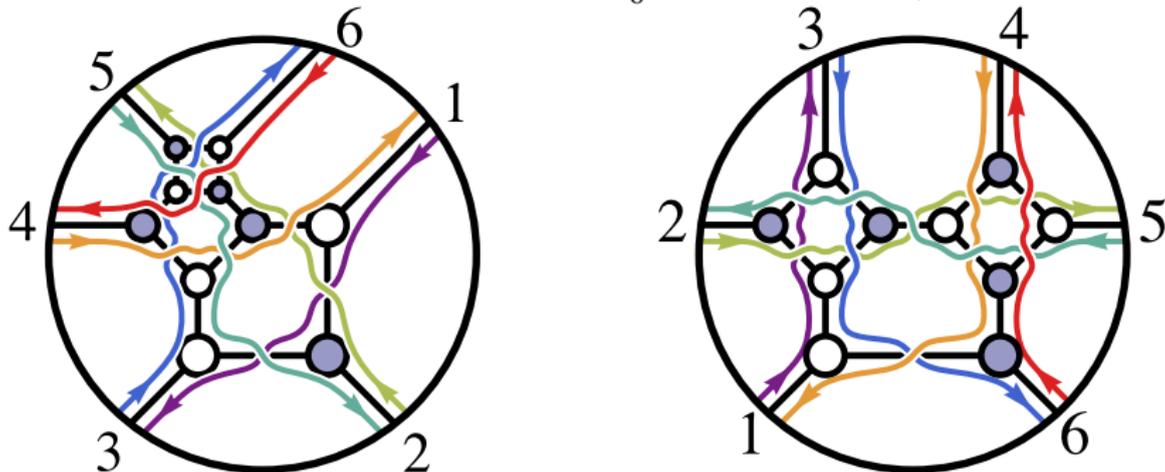
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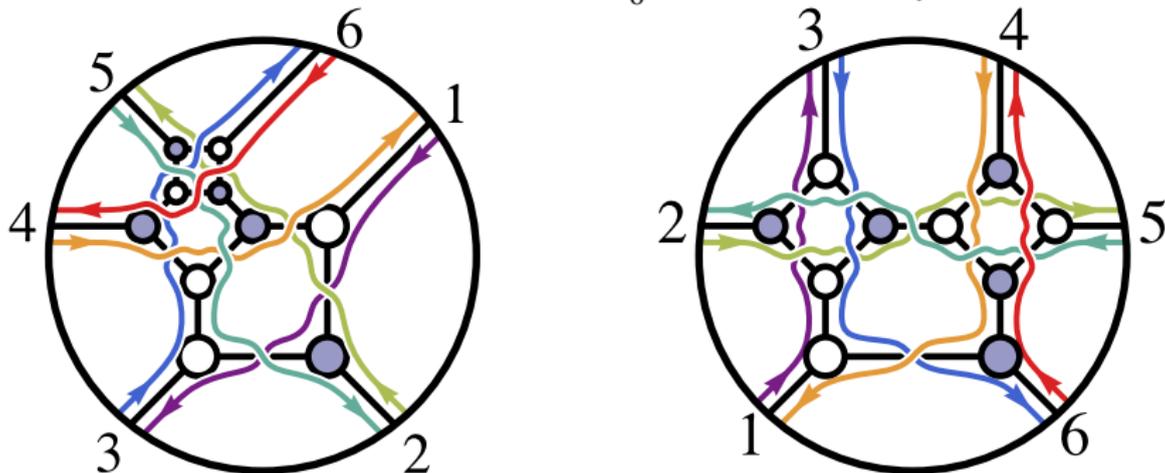
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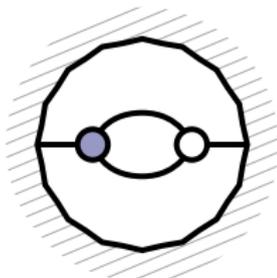
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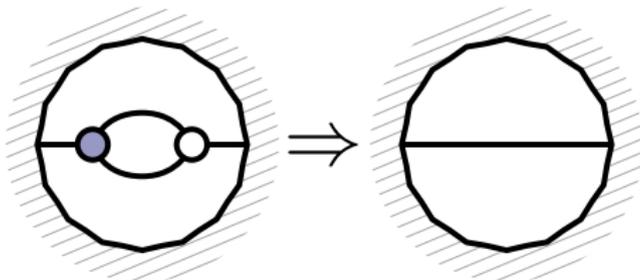
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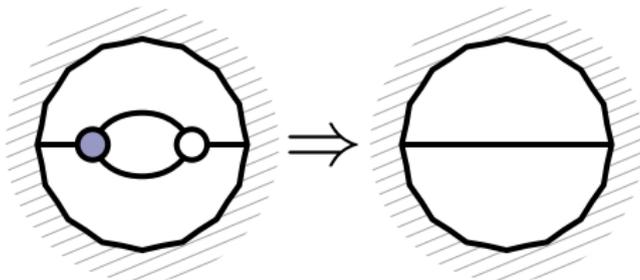
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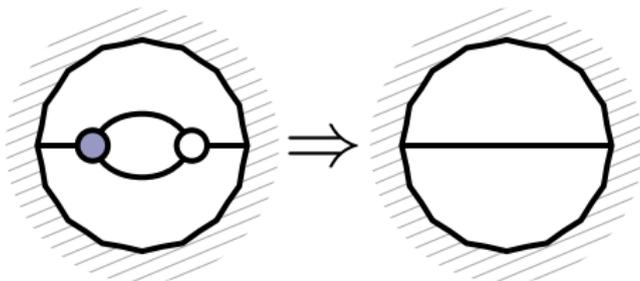


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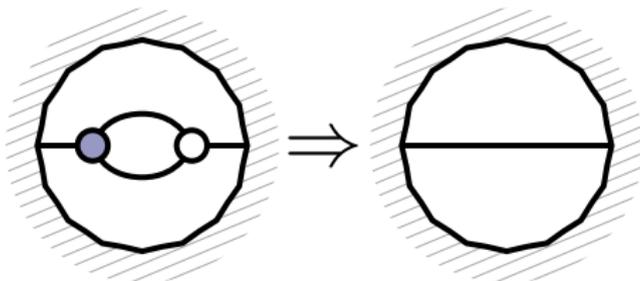


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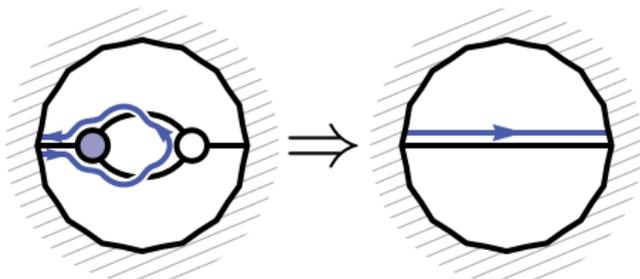


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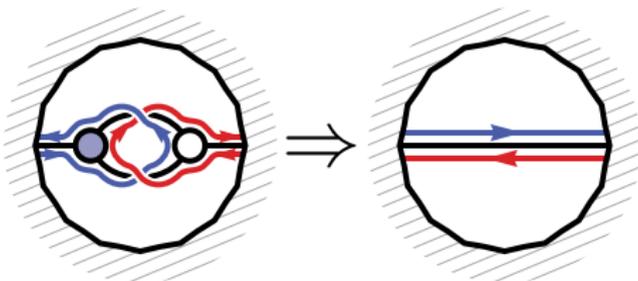


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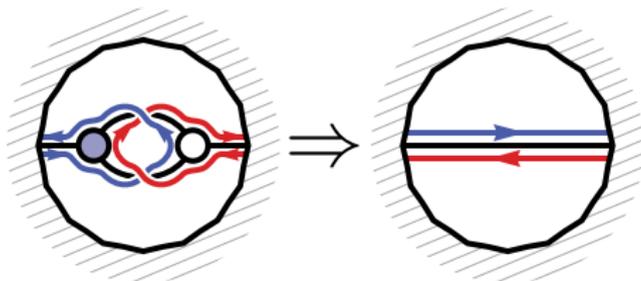
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Such factors of  $d\alpha/\alpha$  arising from bubble deletion encode **loop integrands!**



# Canonical Coordinates for Computing On-Shell Functions

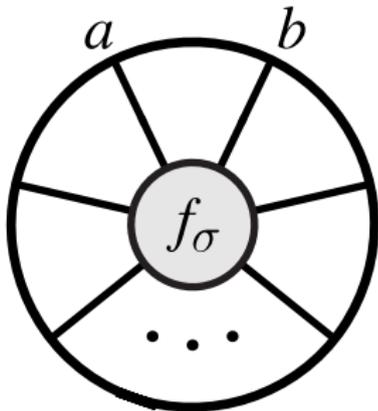
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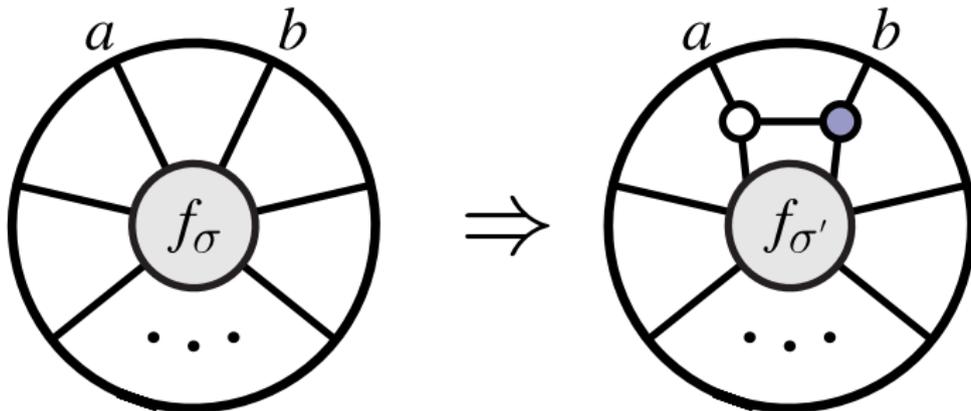
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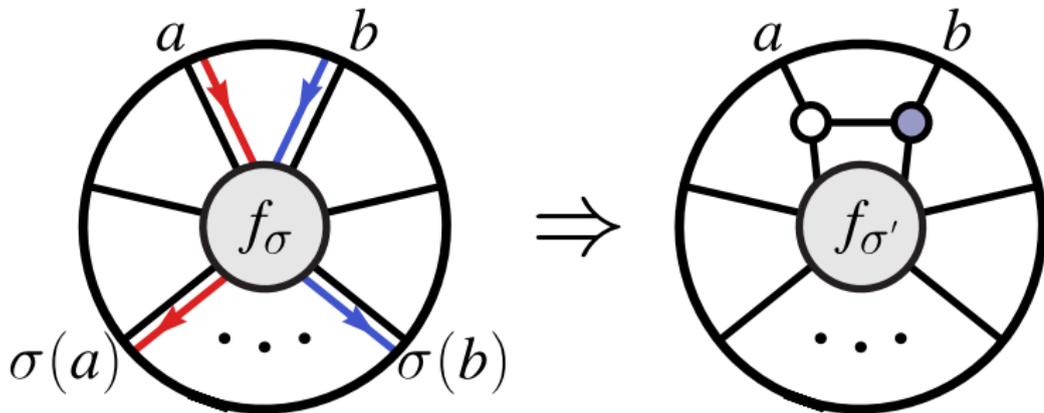
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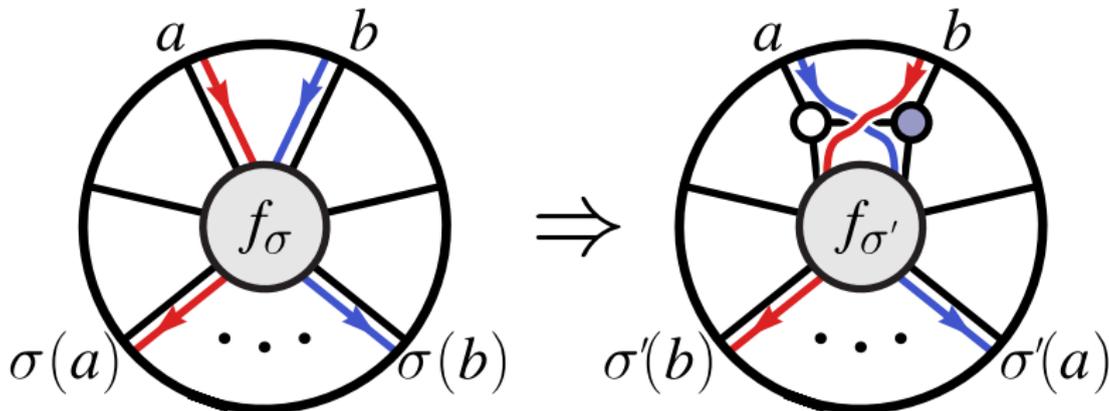
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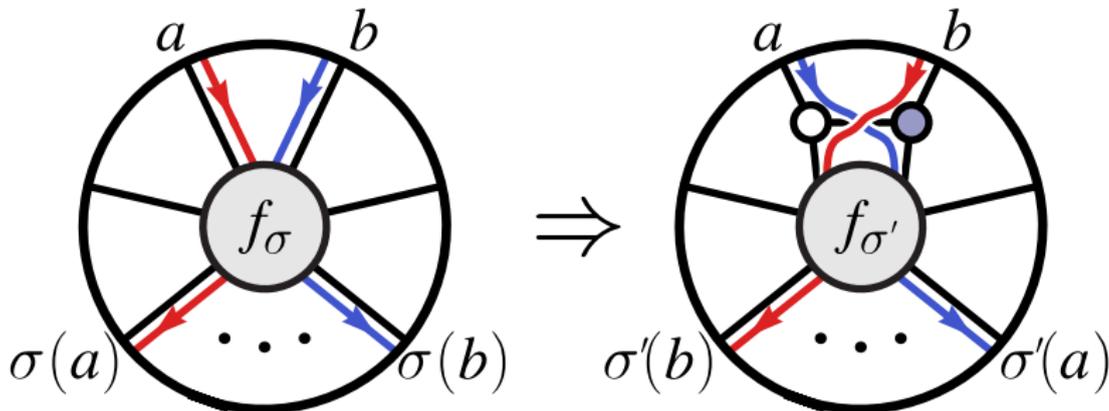
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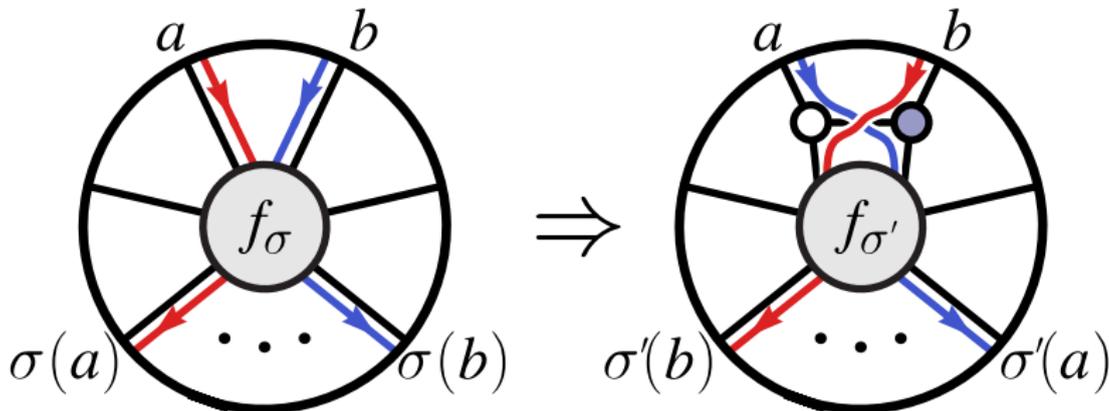
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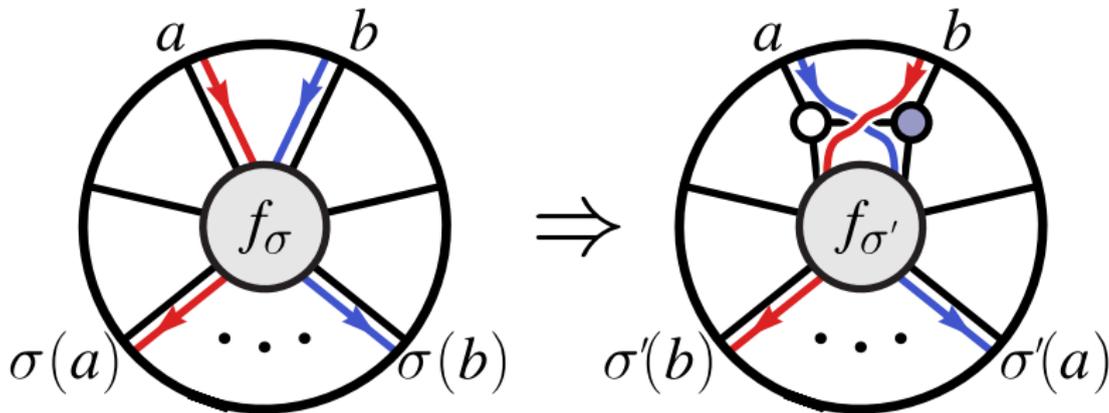
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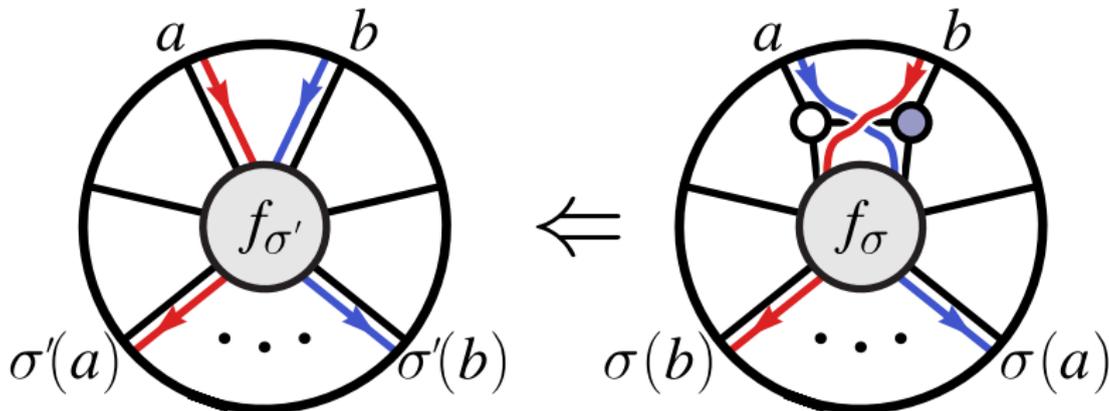
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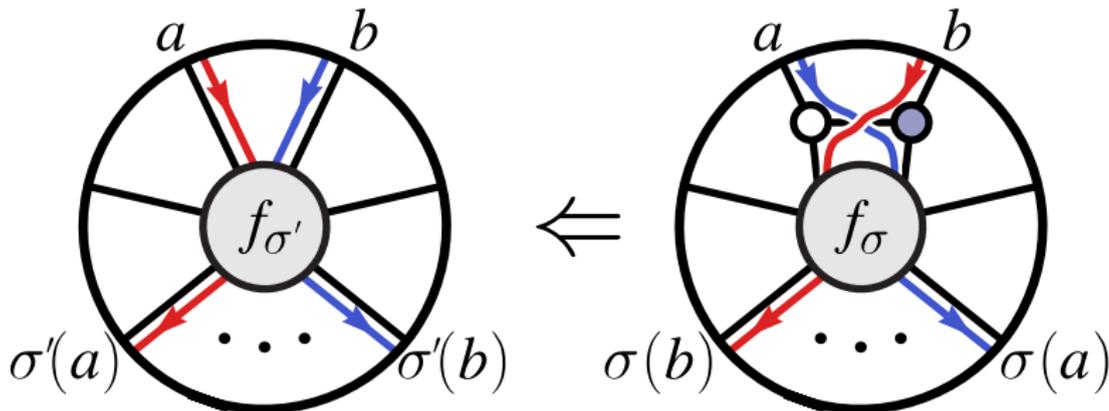
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Read the other way, we can ‘peel-off’ bridges and thereby **decompose** a permutation into transpositions according to  $\sigma = (ab) \circ \sigma'$



# Canonical Coordinates for Computing On-Shell Functions

There are many ways to decompose a permutation into transpositions

‘Bridge’ Decomposition

$$\sigma: \begin{pmatrix} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ 3 & 5 & 6 & 7 & 8 & 10 \end{pmatrix}$$

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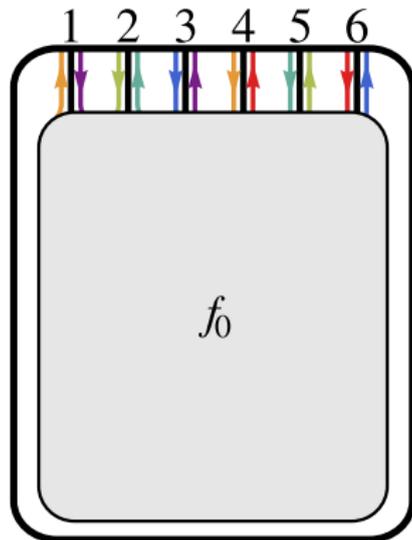
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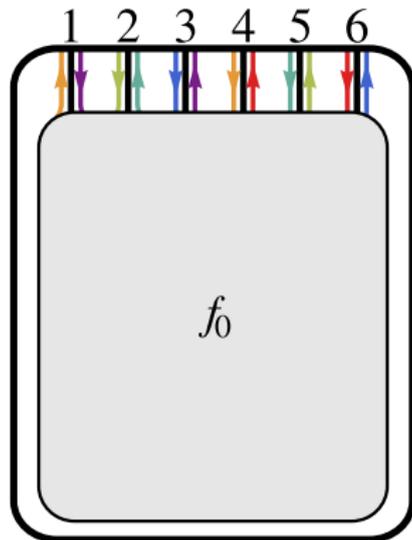


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$$f_0 \begin{matrix} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \\ \{3 & 5 & 6 & 7 & 8 & 10\} \end{matrix} \tau$$

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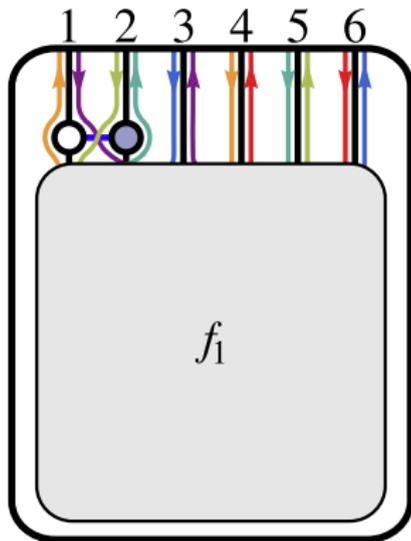
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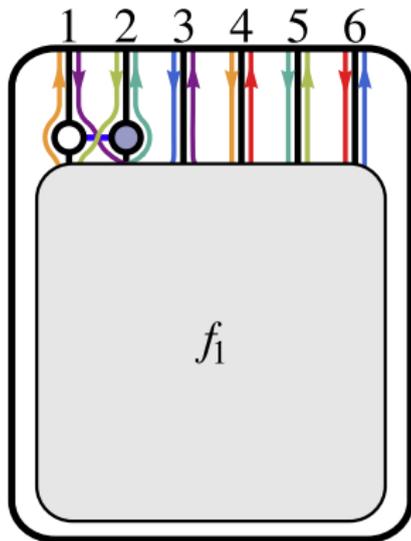


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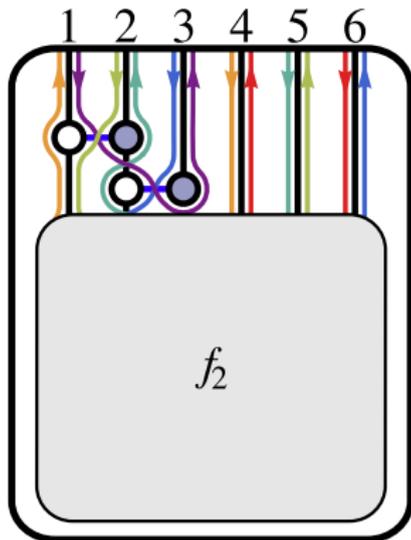


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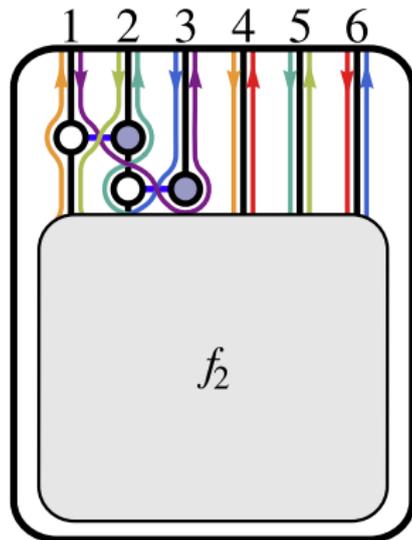


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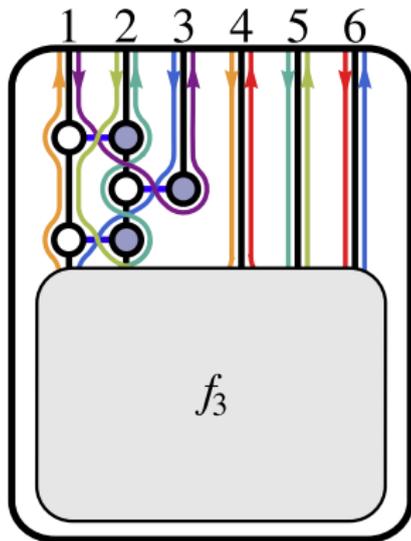


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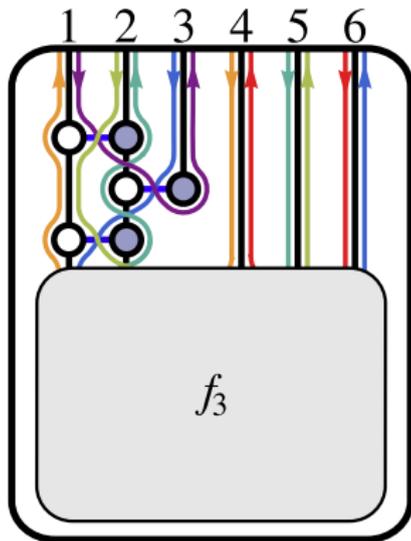


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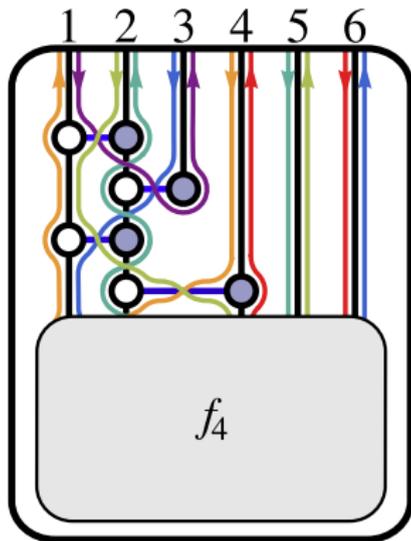
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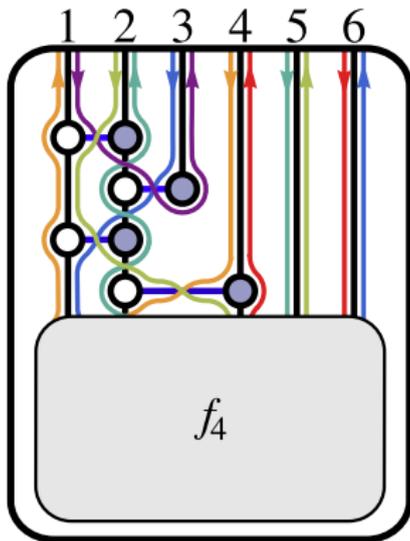


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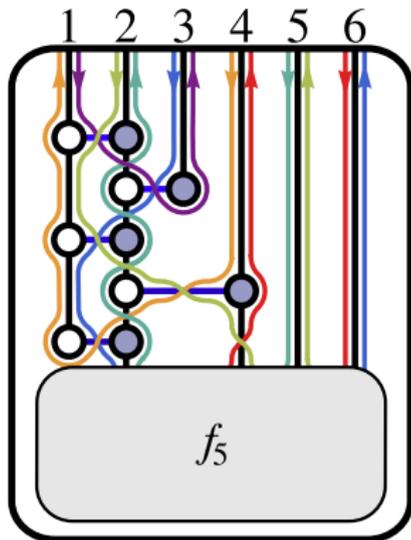


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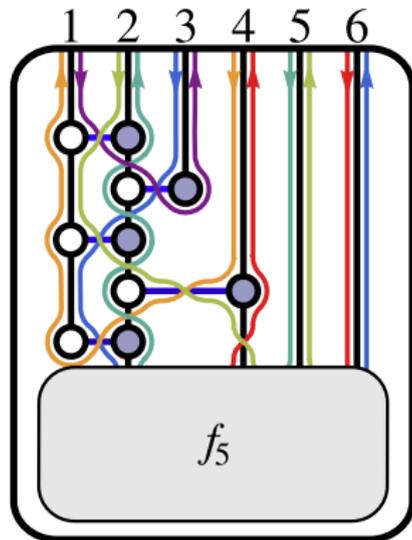


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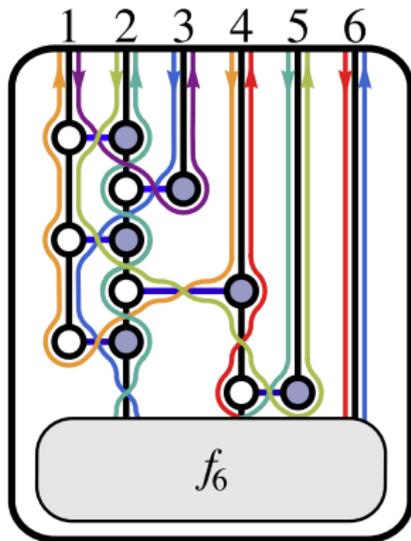
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	{7	6	3	5	8	10}	(4 5)

# Canonical Coordinates for Computing On-Shell Functions

There are many ways to decompose a permutation into transpositions—*e.g.*, always choose the **first** transposition  $\tau \equiv (ab)$  such that  $\sigma(a) < \sigma(b)$ :

$$f_0 = \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \frac{d\alpha_3}{\alpha_3} \frac{d\alpha_4}{\alpha_4} \frac{d\alpha_5}{\alpha_5} \frac{d\alpha_6}{\alpha_6} f_6$$

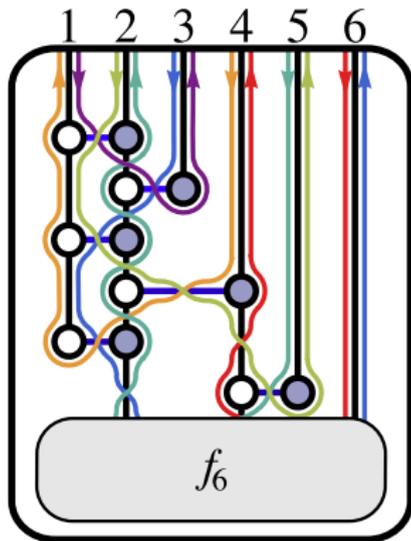


‘Bridge’ Decomposition							
	1	2	3	4	5	6	$\tau$
$f_0$	↓	↓	↓	↓	↓	↓	
$f_1$	{3	5	6	7	8	10}	(1 2)
$f_2$	{5	3	6	7	8	10}	(2 3)
$f_3$	{5	6	3	7	8	10}	(1 2)
$f_4$	{6	5	3	7	8	10}	(2 4)
$f_5$	{6	7	3	5	8	10}	(1 2)
$f_6$	{7	6	3	5	8	10}	(4 5)

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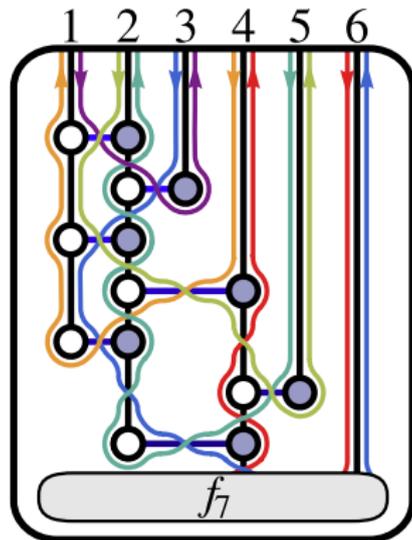


'Bridge' Decomposition							
	1	2	3	4	5	6	$\tau$
$f_0$	↓	↓	↓	↓	↓	↓	
$f_1$	{3	5	6	7	8	10}	(1 2)
$f_2$	{5	3	6	7	8	10}	(2 3)
$f_3$	{5	6	3	7	8	10}	(1 2)
$f_4$	{6	5	3	7	8	10}	(2 4)
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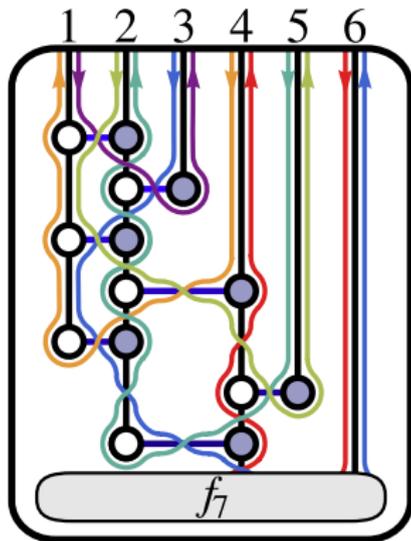
## 'Bridge' Decomposition

	1	2	3	4	5	6	$\tau$
$f_0$	↓	↓	↓	↓	↓	↓	
$f_1$	{3	5	6	7	8	10}	(1 2)
$f_2$	{5	3	6	7	8	10}	(2 3)
$f_3$	{5	6	3	7	8	10}	(1 2)
$f_4$	{6	5	3	7	8	10}	(2 4)
$f_5$	{6	7	3	5	8	10}	(1 2)
$f_6$	{7	6	3	5	8	10}	(4 5)
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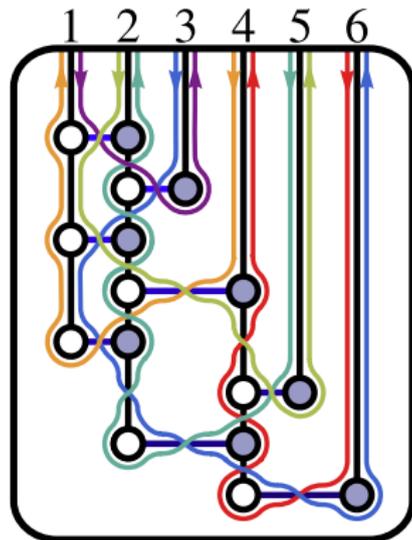
## 'Bridge' Decomposition

	1	2	3	4	5	6	$\tau$
$f_0$	↓	↓	↓	↓	↓	↓	
$f_1$	{3	5	6	7	8	10}	(1 2)
$f_2$	{5	3	6	7	8	10}	(2 3)
$f_3$	{5	6	3	7	8	10}	(1 2)
$f_4$	{6	5	3	7	8	10}	(2 4)
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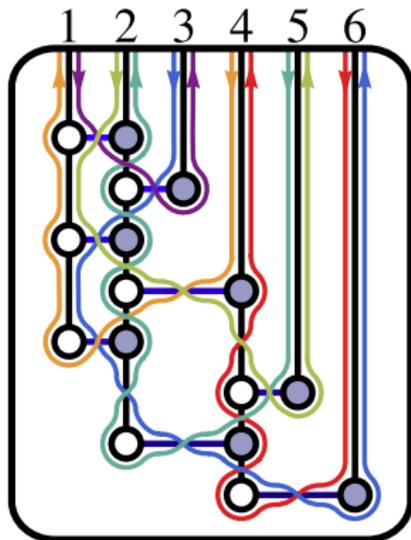
## 'Bridge' Decomposition

	1	2	3	4	5	6	$\tau$
$f_0$	↓	↓	↓	↓	↓	↓	
$f_1$	{3	5	6	7	8	10}	(1 2)
$f_2$	{5	3	6	7	8	10}	(2 3)
$f_3$	{5	6	3	7	8	10}	(1 2)
$f_4$	{6	5	3	7	8	10}	(2 4)
$f_5$	{6	7	3	5	8	10}	(1 2)
$f_6$	{7	6	3	5	8	10}	(4 5)
$f_7$	{7	6	3	8	5	10}	(2 4)
$f_8$	{7	8	3	6	5	10}	(4 6)
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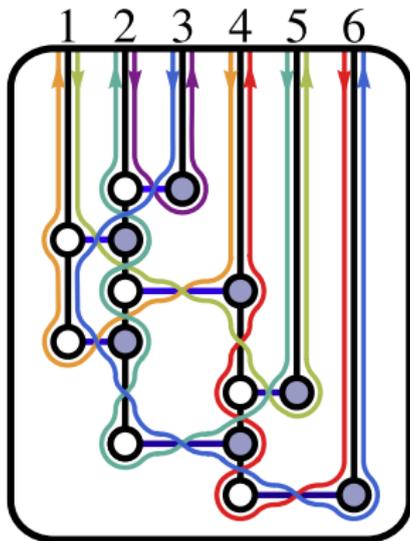
## 'Bridge' Decomposition

	1	2	3	4	5	6	$\tau$
$f_0$	{3	5	6	7	8	10}	(1 2)
$f_1$	{5	3	6	7	8	10}	(2 3)
$f_2$	{5	6	3	7	8	10}	(1 2)
$f_3$	{6	5	3	7	8	10}	(2 4)
$f_4$	{6	7	3	5	8	10}	(1 2)
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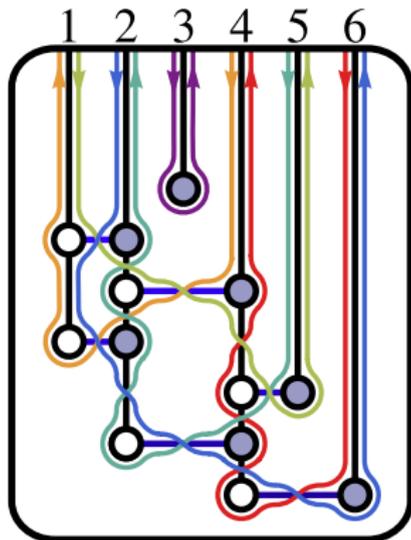
## 'Bridge' Decomposition

	1	2	3	4	5	6	$\tau$
	↓	↓	↓	↓	↓	↓	
$f_1$	{5	3	6	7	8	10}	(23)
$f_2$	{5	6	3	7	8	10}	(12)
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## 'Bridge' Decomposition

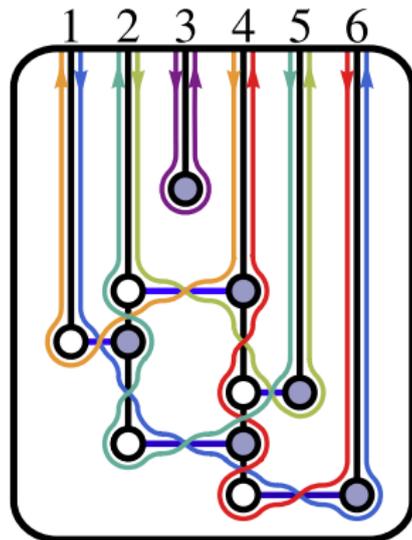
1	2	3	4	5	6	
↓	↓	↓	↓	↓	↓	$\tau$

$f_2$	{	5	6	3	7	8	10	}	(1 2)
$f_3$	{	6	5	3	7	8	10	}	(2 4)
$f_4$	{	6	7	3	5	8	10	}	(1 2)
$f_5$	{	7	6	3	5	8	10	}	(4 5)
$f_6$	{	7	6	3	8	5	10	}	(2 4)
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## 'Bridge' Decomposition

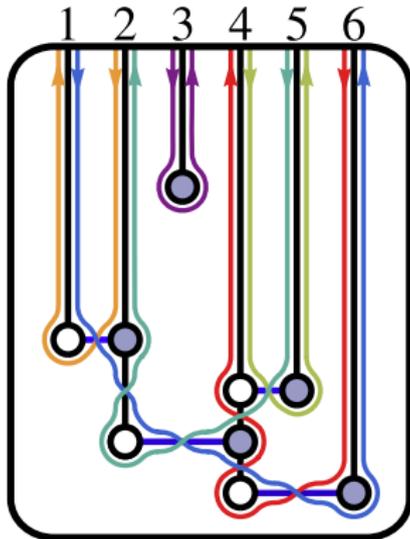
1	2	3	4	5	6	
↓	↓	↓	↓	↓	↓	$\tau$

$f_3$	{6	5	3	7	8	10}	(24)
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## 'Bridge' Decomposition

1	2	3	4	5	6	
↓	↓	↓	↓	↓	↓	$\tau$

$f_4$	{	6	7	3	5	8	10	}	(1 2)
$f_5$	{	7	6	3	5	8	10	}	(4 5)
$f_6$	{	7	6	3	8	5	10	}	(2 4)
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$$f_8 = \prod_{a=\sigma(a)+n} \left( \delta^4(\tilde{\eta}_a) \delta^2(\tilde{\lambda}_a) \right) \prod_{b=\sigma(b)} \left( \delta^2(\lambda_b) \right)$$

$$C \equiv \begin{pmatrix} \mathbf{1} & \mathbf{2} & \mathbf{3} & \mathbf{4} & \mathbf{5} & \mathbf{6} \\ \hline 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \end{pmatrix}$$

## 'Bridge' Decomposition

$$\begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{array} \quad \tau$$

$$f_8 \{ \mathbf{7} \ \mathbf{8} \ \mathbf{3} \ \mathbf{10} \ \mathbf{5} \ \mathbf{6} \}$$

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$$f_8 = \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^\perp)$$

$$C \equiv \begin{pmatrix} \mathbf{1} & \mathbf{2} & \mathbf{3} & \mathbf{4} & \mathbf{5} & \mathbf{6} \\ \hline 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \end{pmatrix}$$

‘Bridge’ Decomposition

$$\begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{array} \quad \tau$$

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‘Bridge’ Decomposition

$$\begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{array} \quad \tau$$

$$f_8 \{ \mathbf{7} \ \mathbf{8} \ \mathbf{3} \ \mathbf{10} \ \mathbf{5} \ \mathbf{6} \}$$

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$$f_7 = \frac{d\alpha_8}{\alpha_8} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^\perp)$$

$$C \equiv \begin{pmatrix} 1 & 2 & 3 & 4 & 5 & 6 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & \alpha_8 \end{pmatrix}$$

(46):  $c_6 \mapsto c_6 + \alpha_8 c_4$

## 'Bridge' Decomposition

$$\begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{array} \quad \tau$$

$$f_7 \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} (46)$$

$$f_8 \{7 \ 8 \ 3 \ 10 \ 5 \ 6\}$$

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There are many ways to decompose a permutation into transpositions—*e.g.*, always choose the **first** transposition  $\tau \equiv (ab)$  such that  $\sigma(a) < \sigma(b)$ :

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$$f_6 = \frac{d\alpha_7}{\alpha_7} \frac{d\alpha_8}{\alpha_8} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^\perp)$$

$$C \equiv \begin{pmatrix} 1 & 2 & 3 & 4 & 5 & 6 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & \alpha_7 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & \alpha_8 \end{pmatrix}$$

(24):  $c_4 \mapsto c_4 + \alpha_7 c_2$

## 'Bridge' Decomposition

$$\begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{array} \quad \tau$$

$$\begin{array}{l} f_6 \{7 \ 6 \ 3 \ 8 \ 5 \ 10\} (24) \\ f_7 \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} (46) \\ f_8 \{7 \ 8 \ 3 \ 10 \ 5 \ 6\} \end{array}$$

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$$f_5 = \frac{d\alpha_6}{\alpha_6} \frac{d\alpha_7}{\alpha_7} \frac{d\alpha_8}{\alpha_8} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^\perp)$$

$$C \equiv \begin{pmatrix} 1 & 2 & 3 & 4 & 5 & 6 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & \alpha_7 & \alpha_6 \alpha_7 & 0 \\ 0 & 0 & 0 & 1 & \alpha_6 & \alpha_8 \end{pmatrix}$$

(45):  $c_5 \mapsto c_5 + \alpha_6 c_4$

## 'Bridge' Decomposition

$$\begin{array}{cccccc} 1 & 2 & 3 & 4 & 5 & 6 \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{array} \quad \tau$$

$$\begin{array}{l} f_5 \{7 \ 6 \ 3 \ 5 \ 8 \ 10\} \\ f_6 \{7 \ 6 \ 3 \ 8 \ 5 \ 10\} \\ f_7 \{7 \ 8 \ 3 \ 6 \ 5 \ 10\} \\ f_8 \{7 \ 8 \ 3 \ 10 \ 5 \ 6\} \end{array} \begin{array}{l} (45) \\ (24) \\ (46) \end{array}$$







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There are many ways to decompose a permutation into transpositions—*e.g.*, always choose the **first** transposition  $\tau \equiv (ab)$  such that  $\sigma(a) < \sigma(b)$ :

$$f_0 = \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \frac{d\alpha_3}{\alpha_3} \frac{d\alpha_4}{\alpha_4} \frac{d\alpha_5}{\alpha_5} \frac{d\alpha_6}{\alpha_6} \frac{d\alpha_7}{\alpha_7} \frac{d\alpha_8}{\alpha_8} f_8$$

$$f_1 = \frac{d\alpha_2}{\alpha_2} \dots \frac{d\alpha_8}{\alpha_8} \delta^{3 \times 4} (C \cdot \tilde{\eta}) \delta^{3 \times 2} (C \cdot \tilde{\lambda}) \delta^{2 \times 3} (\lambda \cdot C^\perp)$$

$$C \equiv \begin{pmatrix} 1 & 2 & 3 & 4 & 5 & 6 \\ 1 & (\alpha_3 + \alpha_5) & \alpha_2(\alpha_3 + \alpha_5) & \alpha_4\alpha_5 & 0 & 0 \\ 0 & 1 & \alpha_2 & (\alpha_4 + \alpha_7) & \alpha_6\alpha_7 & 0 \\ 0 & 0 & 0 & 1 & \alpha_6 & \alpha_8 \end{pmatrix}$$

(23):  $c_3 \mapsto c_3 + \alpha_2 c_2$

‘Bridge’ Decomposition							
	1	2	3	4	5	6	$\tau$
	↓	↓	↓	↓	↓	↓	
$f_1$	{5	3	6	7	8	10}	(23)
$f_2$	{5	6	3	7	8	10}	(12)
$f_3$	{6	5	3	7	8	10}	(24)
$f_4$	{6	7	3	5	8	10}	(12)
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# Canonical Coordinates for Computing On-Shell Functions

There are many ways to decompose a permutation into transpositions—*e.g.*, always choose the **first** transposition  $\tau \equiv (ab)$  such that  $\sigma(a) < \sigma(b)$ :

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$$C \equiv \begin{pmatrix} 1 & 2 & 3 & 4 & 5 & 6 \\ 1 & (\alpha_1 + \alpha_3 + \alpha_5) & \alpha_2(\alpha_3 + \alpha_5) & \alpha_4\alpha_5 & 0 & 0 \\ 0 & 1 & \alpha_2 & (\alpha_4 + \alpha_7) & \alpha_6\alpha_7 & 0 \\ 0 & 0 & 0 & 1 & \alpha_6 & \alpha_8 \end{pmatrix}$$

(12):  $c_2 \mapsto c_2 + \alpha_1 c_1$

‘Bridge’ Decomposition							
	1	2	3	4	5	6	$\tau$
$f_0$	↓	↓	↓	↓	↓	↓	
$f_0$	{3	5	6	7	8	10}	(12)
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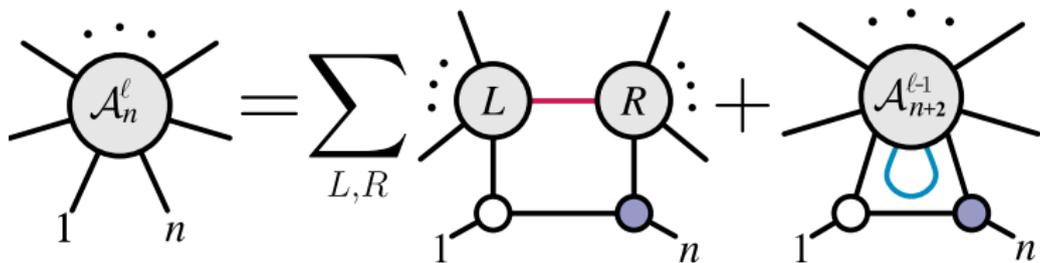
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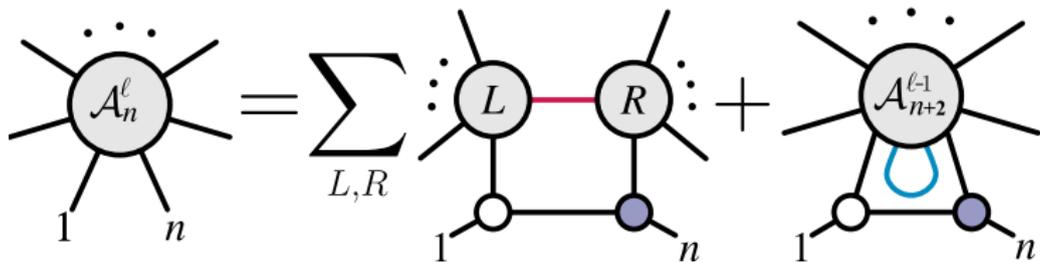
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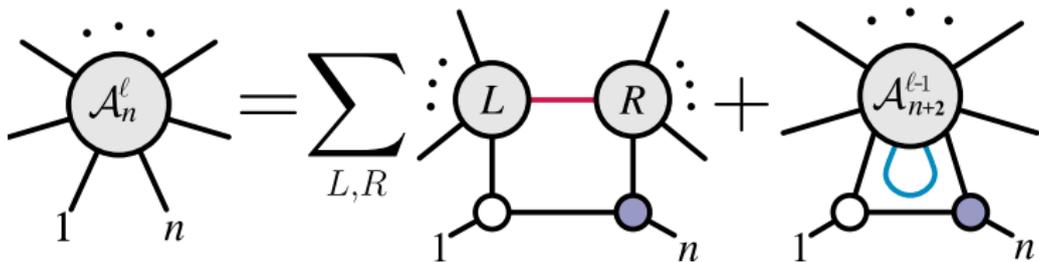
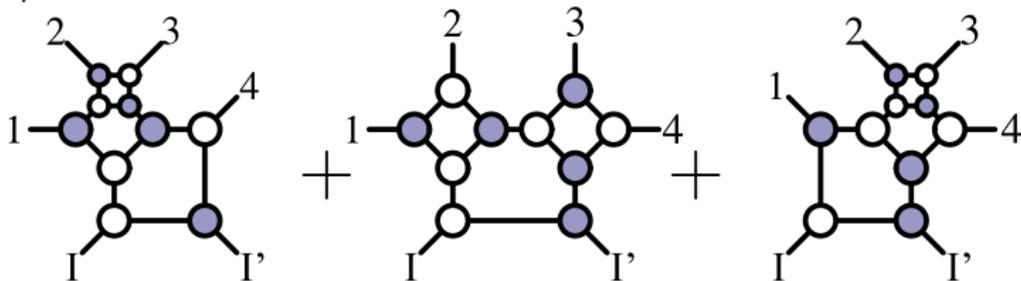
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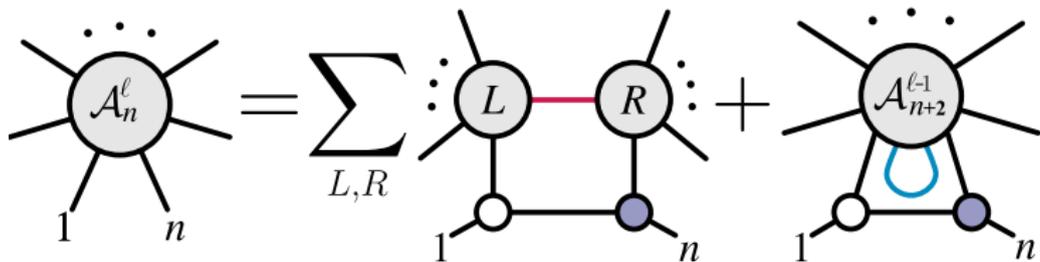
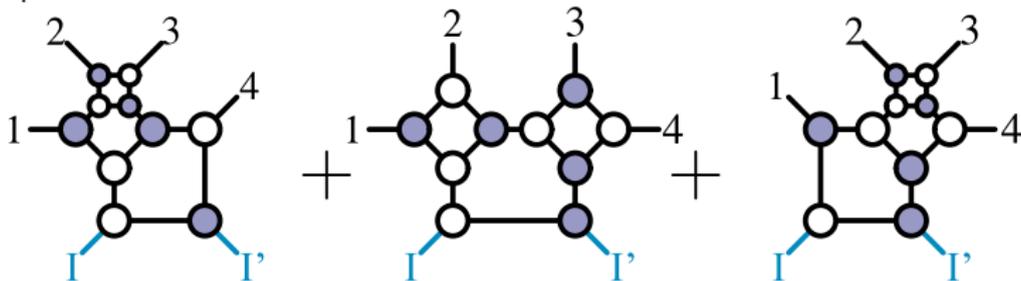
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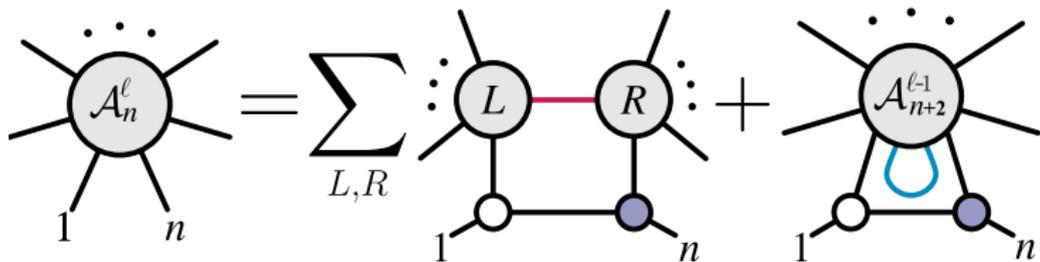
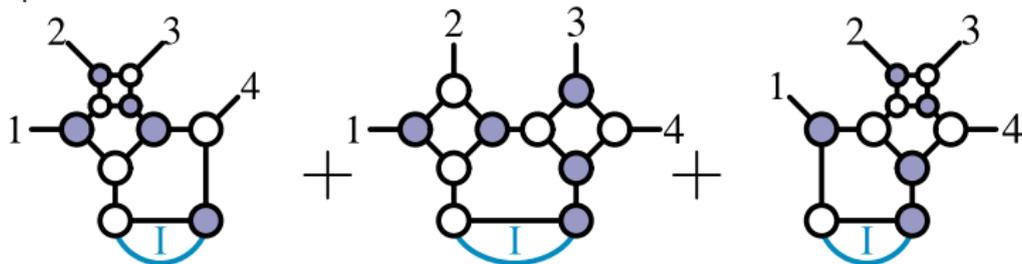
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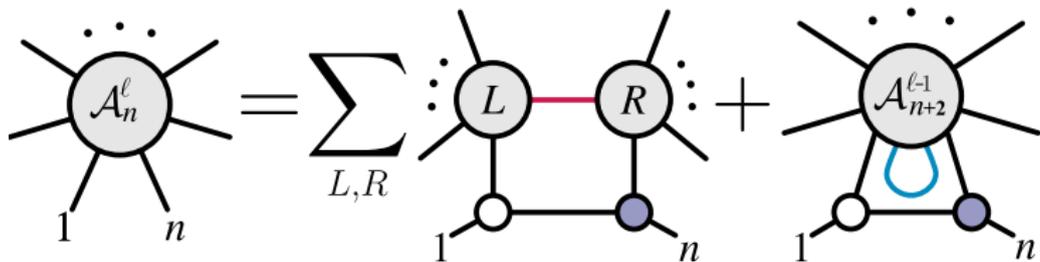
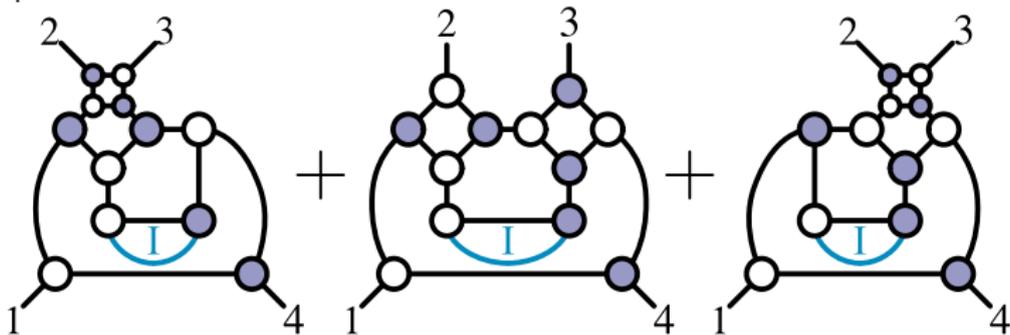
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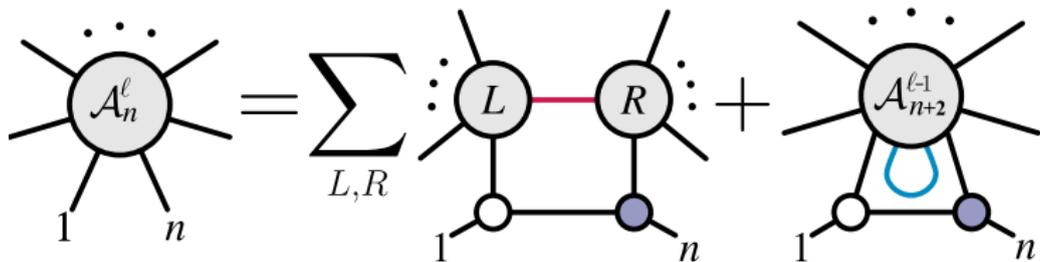
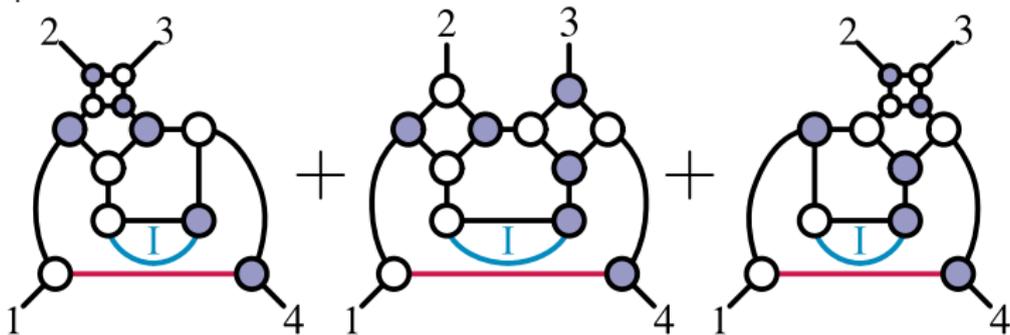
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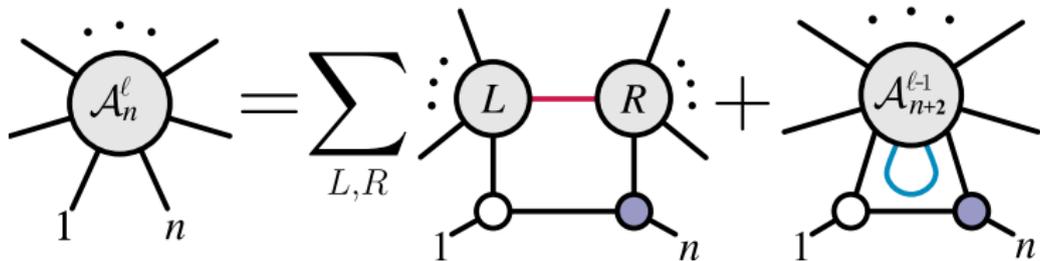
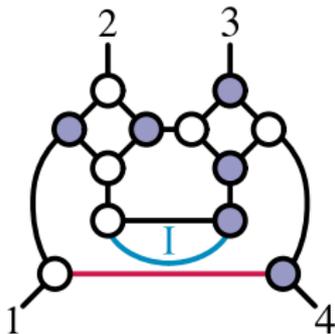
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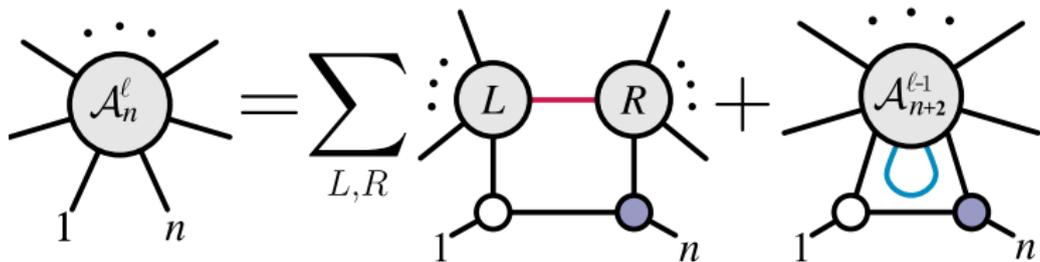
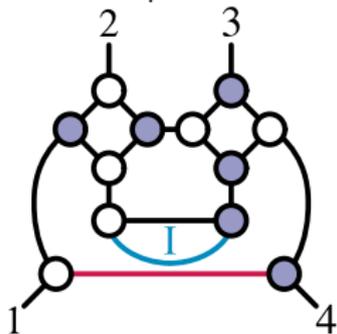
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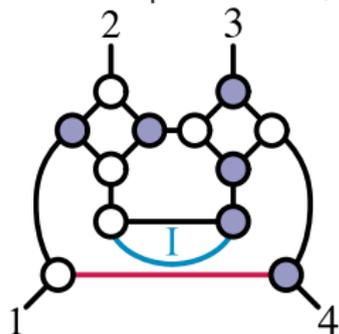
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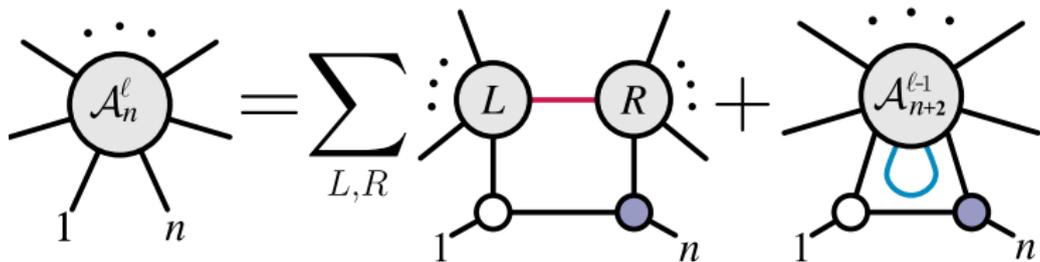
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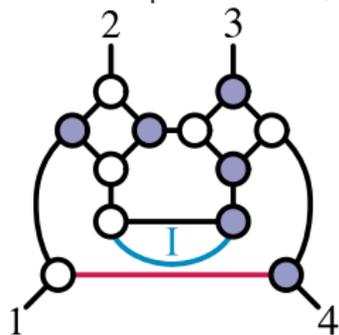
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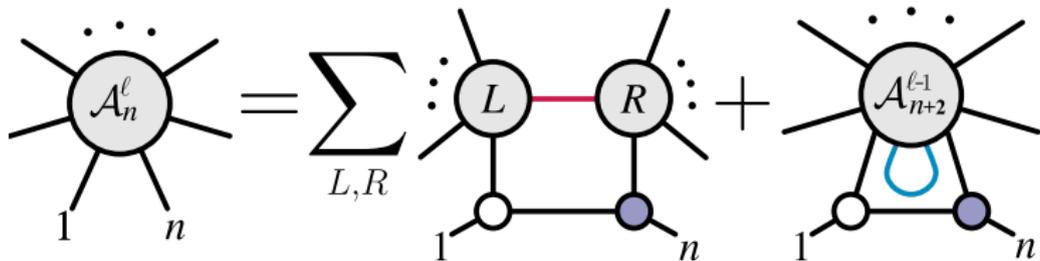
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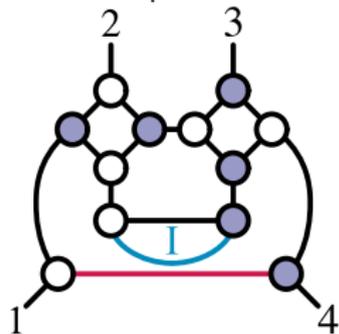
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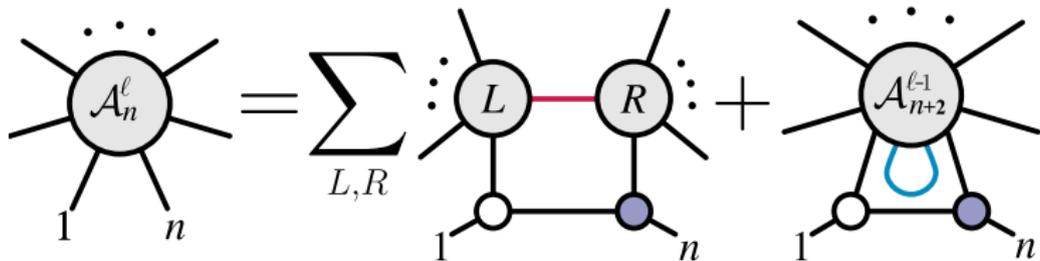
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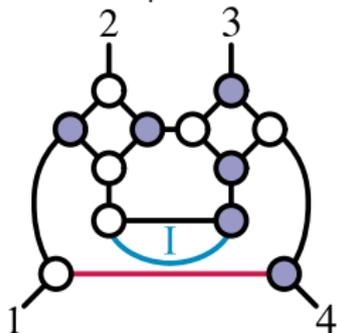
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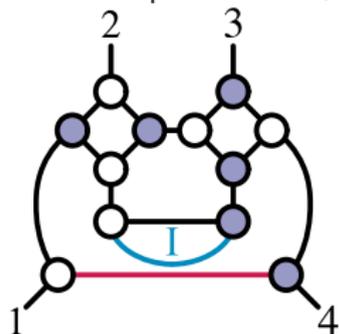
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$$= \mathcal{A}_4^{(2),0} \times \int_{\ell \in \mathbb{R}^{3,1}} d^4 \ell \frac{(p_1 + p_2)^2 (p_3 + p_4)^2}{\ell^2 (\ell + p_1)^2 (\ell + p_1 + p_2)^2 (\ell - p_4)^2}$$

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complexity of the computations. It has also been useful to use the results for the cuts already computed when computing the coefficients of integrals detected by new cuts. In this way, one can insure the consistency of results from different cuts and reduce the number of unknowns at the same time.

Let us make a further comment about our computation procedure. The conformal integrals with pentagon loops have numerators containing the loop momenta in combinations like  $(k+l)^2$ , where  $l$  is the loop momentum and  $k$  is an external on-shell momentum. If the propagator with momentum  $l$  is cut then, on that cut, one cannot distinguish between  $(k+l)^2$  and  $2kl-l^2$ . However, it is easy to see that one can choose to cut another propagator and in that case this ambiguity does not arise and the numerator factor is uniquely defined.

#### IV. RESULTS

We use dual variable notation (see Ref. [18]) for the integrals. The external dual variables are listed in clockwise direction. To the left loop we associate the dual variable  $x_p$  and to the right loop we associate the dual variable  $x_q$ . We use the notation  $x_{ij} = x_i - x_j$ .

We introduce the following notation which will be useful in the following

$$\begin{bmatrix} a & b & c & \cdots \\ a' & b' & c' & \cdots \end{bmatrix} = x_{a'}^2 x_{b'}^2 x_{c'}^2 \cdots \pm (\text{permutations of } \{a', b', c', \dots\}). \quad (6)$$

The sign  $\pm$  above takes into account the signature of the permutation of  $\{a', b', c', \dots\}$ . It is easy to show that

$$\begin{bmatrix} a & b & c & \cdots \\ a' & b' & c' & \cdots \end{bmatrix} = \det_{\substack{a, b, c, \dots \\ a', b', c', \dots}} x_{a'}^2. \quad (7)$$

For some topologies, the expansion of the  $\begin{bmatrix} \phantom{a} & \phantom{b} & \phantom{c} & \cdots \\ \phantom{a'} & \phantom{b'} & \phantom{c'} & \cdots \end{bmatrix}$  symbol yields terms that would cancel propagators. For those cases we make the convention that all the terms that would cancel propagators are absent. In fact, as we will see, terms that would cancel propagators of the double pentagon topologies naturally yield coefficients for some of the topologies with a smaller number of propagators.

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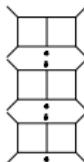
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### A. Double box topologies

In the case of the double box topologies the massive legs attached to the vertices incident with the common edge have to be a sum of at least three massless momenta. The cases where these massive legs are the sum of two massless momenta are treated separately in the subsection IV A 7. This distinction only arises for the double box topologies.

#### 1. No legs attached

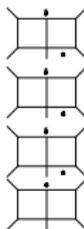


$$\frac{1}{2} (s_{2,3}^2)^2 s_{2,1,3}^2 \quad (8)$$

$$\frac{1}{4} (s_{23}^2)^2 s_{2,1,3}^2 \quad (9)$$

$$-\frac{1}{4} s_{23}^2 (s_{23}^2 - s_{2,1,3}^2 - s_{2,1,3}^2)^2 \quad (10)$$

#### 2. One massless leg attached



$$\frac{1}{4} (s_{2,1,3}^2 s_{2,1,3}^2 - s_{23}^2 s_{2,1,3}^2) s_{2,2,3}^2 \quad (11)$$

$$\frac{1}{4} (-s_{2,1,3}^2 s_{2,1,3}^2 s_{2,1,3}^2 + s_{2,1,3}^2 s_{2,1,3}^2 s_{2,1,3}^2 - s_{2,1,3}^2 s_{2,1,3}^2 s_{2,1,3}^2) \quad (12)$$

$$-\frac{1}{4} s_{23}^2 s_{2,1,3}^2 s_{2,1,3}^2 \quad (13)$$

$$-\frac{1}{4} s_{2,1,3}^2 s_{2,1,3}^2 s_{2,1,3}^2 \quad (14)$$

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- using techniques based on 'generalized unitarity', and
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### Abstract

We compute the even part of the planar two-loop MHV amplitude in  $\mathcal{N} = 4$  supersymmetric Yang-Mills theory, for an arbitrary number of external particles. The answer is expressed as a sum of conformal integrals.



$$\frac{1}{4}(x_{12}^2 - x_{34}^2 - x_{56}^2 - 2x_{13}^2 x_{24}^2 - x_{25}^2 x_{36}^2 - x_{35}^2 x_{46}^2) x_{12}^2 x_{34}^2 \quad (15)$$

3. Two massless legs attached



$$\frac{1}{4}(x_{12}^2 x_{34}^2 x_{56}^2 + x_{13}^2 x_{24}^2 x_{56}^2 - x_{23}^2 x_{14}^2 x_{56}^2 - x_{14}^2 x_{23}^2 x_{56}^2 + x_{15}^2 x_{26}^2 x_{34}^2 + x_{25}^2 x_{16}^2 x_{34}^2) \quad (16)$$



$$\frac{1}{4}(-x_{12}^2 x_{34}^2 x_{56}^2 + x_{13}^2 x_{24}^2 x_{56}^2 + x_{23}^2 x_{14}^2 x_{56}^2 - x_{14}^2 x_{23}^2 x_{56}^2 - x_{15}^2 x_{26}^2 x_{34}^2 - x_{25}^2 x_{16}^2 x_{34}^2) \quad (17)$$



$$\frac{1}{4}(x_{12}^2 x_{34}^2 x_{56}^2 - x_{13}^2 x_{24}^2 x_{56}^2 - x_{23}^2 x_{14}^2 x_{56}^2 + x_{14}^2 x_{23}^2 x_{56}^2 + x_{15}^2 x_{26}^2 x_{34}^2 - x_{25}^2 x_{16}^2 x_{34}^2) \quad (18)$$

4. One massive leg attached



$$\frac{1}{4}x_{12}^2 x_{34}^2 x_{56}^2 - x_{13}^2 x_{24}^2 x_{56}^2 \quad (19)$$



$$\frac{1}{4}(x_{12}^2 x_{34}^2 x_{56}^2 - x_{13}^2 x_{24}^2 x_{56}^2 - x_{23}^2 x_{14}^2 x_{56}^2 + x_{14}^2 x_{23}^2 x_{56}^2) \quad (20)$$



$$0 \quad (21)$$



$$\frac{1}{4}(x_{12}^2 x_{34}^2 x_{56}^2 - x_{13}^2 x_{24}^2 x_{56}^2 - x_{23}^2 x_{14}^2 x_{56}^2 + x_{14}^2 x_{23}^2 x_{56}^2) \quad (22)$$

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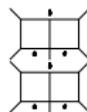
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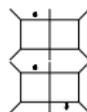
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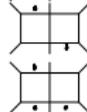


$$0 \quad (23)$$

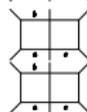
5. One massless leg and one massive leg attached



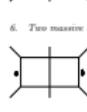
$$0 \quad (24)$$



$$-\frac{1}{4}x_{2-2a}^2 x_{1,a}^2 x_{2+1,a}^2 \quad (25)$$



$$0 \quad (26)$$



$$\frac{1}{4}x_{2-2a}^2 (x_{2-1,b}^2 x_{+1,b-1}^2 - x_{2-1,b-1}^2 x_{2+1,b}^2) \quad (27)$$



$$\frac{1}{4}(-x_{2a+1,b}^2 x_{+1,b-1}^2 + x_{2a+1,b-1}^2 x_{2+1,b}^2 + x_{2a-1,b}^2 x_{2+1,b}^2 - x_{2a-1,b-1}^2 x_{2+1,b}^2) \quad (28)$$



$$0 \quad (29)$$



$$0 \quad (30)$$



$$0 \quad (31)$$

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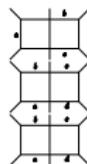
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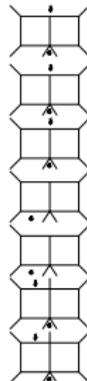


$$0 \quad (32)$$

$$0 \quad (33)$$

$$0 \quad (34)$$

2. Extra double boxes



$$\frac{1}{4} \left( -x_{2-2a}^2 x_{2-1,a}^2 x_{2+1,b}^2 + x_{2-2,a}^2 x_{2-1,b}^2 x_{2+1,a}^2 + x_{2-2,b}^2 x_{2-1,a}^2 x_{2+1,b}^2 - x_{2-2,a}^2 x_{2-1,b}^2 x_{2+1,a}^2 \right) \quad (35)$$

$$-\frac{1}{4} \begin{bmatrix} a+1 & b-1 & b \\ b & b+1 & a-1 \end{bmatrix} \quad (36)$$

$$0 \quad (37)$$

$$-\frac{1}{4} \begin{bmatrix} a & a+1 & a+2 \\ a+2 & a+3 & a-2 \end{bmatrix} \quad (38)$$

$$\frac{1}{4} \left( x_{2-2,a+1}^2 x_{2-2,a+2}^2 - x_{2-2,a}^2 x_{2-2,a+1}^2 \right) x_{2,a+2}^2 \quad (39)$$

$$-\frac{1}{4} \begin{bmatrix} a+1 & b-1 & b \\ b+1 & b+2 & a-1 \end{bmatrix} \quad (40)$$

$$0 \quad (41)$$

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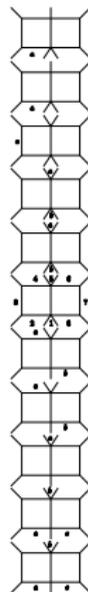
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$$-\frac{1}{4} \begin{bmatrix} a & a+1 & a+2 \\ a+3 & a+4 & a-2 \end{bmatrix} \quad (42)$$

$$\frac{1}{2} (x_{1-2}^2 x_{2-3}^2 x_{3-4}^2 - x_{2-3}^2 x_{3-4}^2 x_{4-1}^2) x_{4-1}^2 \quad (43)$$

$$-\frac{1}{4} \begin{bmatrix} a-1 & a & a+1 \\ a+3 & a-4 & a-3 \end{bmatrix} \quad (44)$$

$$0 \quad (45)$$

$$0 \quad (46)$$

$$-\frac{1}{2} \begin{bmatrix} 2 & 3 & 4 \\ 6 & 7 & 8 \end{bmatrix} \quad (47)$$

$$0 \quad (48)$$

$$-\frac{1}{4} \begin{bmatrix} a-2 & a-1 & a \\ a+2 & b-1 & b \end{bmatrix} \quad (49)$$

$$-\frac{1}{4} \begin{bmatrix} a-3 & a-2 & a-1 \\ a+1 & a+2 & a+3 \end{bmatrix} \quad (50)$$

$$0 \quad (51)$$

$$0 \quad (52)$$

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$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 \\ b+1 & c-1 & c \end{bmatrix} \quad (53)$$

B. Kissing double-box topologies



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b \\ b & b+1 & a-1 & a \end{bmatrix} + \frac{1}{4} \begin{bmatrix} a & a+1 \\ b-1 & b \end{bmatrix} \begin{bmatrix} b & b+1 \\ a-1 & a \end{bmatrix} - \frac{1}{4} (x_{12}^2 - x_{34}^2)^2 - x_{13}^2 x_{24}^2 + x_{14}^2 x_{23}^2 + x_{12}^2 x_{34}^2 - x_{13}^2 x_{24}^2 - x_{14}^2 x_{23}^2 + x_{12}^2 x_{34}^2 + x_{13}^2 x_{24}^2 + x_{14}^2 x_{23}^2 \quad (54)$$



$$-\frac{1}{4} \begin{bmatrix} a+1 & a+2 & b-1 & b \\ b & b+1 & a-1 & a \end{bmatrix} + \frac{1}{4} \begin{bmatrix} a-1 & a \\ b & b+1 \end{bmatrix} \begin{bmatrix} a+1 & a+2 \\ b-1 & b \end{bmatrix} \quad (55)$$



$$-\frac{1}{4} \begin{bmatrix} a+1 & a+2 & b-1 & b \\ b+1 & b+2 & a-1 & a \end{bmatrix} + \frac{1}{4} \begin{bmatrix} a+1 & a+2 \\ b-1 & b \end{bmatrix} \begin{bmatrix} b+1 & b+2 \\ a-1 & a \end{bmatrix} \quad (56)$$



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b \\ b & b+1 & c-1 & c \end{bmatrix} + \frac{1}{4} \begin{bmatrix} a & a+1 \\ b-1 & b \end{bmatrix} \begin{bmatrix} b & b+1 \\ c-1 & c \end{bmatrix} \quad (57)$$



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b \\ b+1 & b+2 & c-1 & c \end{bmatrix} + \frac{1}{4} \begin{bmatrix} a & a+1 \\ b-1 & b \end{bmatrix} \begin{bmatrix} b+1 & b+2 \\ c-1 & c \end{bmatrix} \quad (58)$$



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b \\ c & c+1 & d-1 & d \end{bmatrix} + \frac{1}{4} \begin{bmatrix} a & a+1 \\ b-1 & b \end{bmatrix} \begin{bmatrix} c & c+1 \\ d-1 & d \end{bmatrix} \quad (59)$$

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### C. Box-Pentagon topologies

#### 1. No legs attached



$$\frac{1}{4} s_{23}^2 s_{2+1,4}^2 (s_{2+1,4}^2 s_{2-1,3}^2 - s_{2+1,4}^2 s_{2-1,3}^2) \quad (90)$$

$$\frac{1}{2} s_{2+1,4}^2 s_{2+1,4}^2 (s_{2-1,3}^2 s_{2+1,4}^2 - s_{2-1,3}^2 s_{2+1,4}^2) \quad (91)$$

#### 2. One massless leg attached



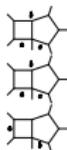
$$\frac{1}{4} (s_{2-1,3}^2 s_{2+1,4}^2 - s_{2-1,3}^2 s_{2+1,4}^2) (s_{2+1,4}^2 s_{2+1,4}^2 - s_{2+1,4}^2 s_{2+1,4}^2) \quad (92)$$

$$\frac{1}{4} s_{2-1,3}^2 (s_{2+1,4}^2 s_{2+1,4}^2 + s_{2+1,4}^2 s_{2+1,4}^2 - s_{2+1,4}^2 s_{2+1,4}^2) \quad (93)$$



$$\frac{1}{4} (s_{2-1,3}^2 s_{2+1,4}^2 - s_{2-1,3}^2 s_{2+1,4}^2 - s_{2-1,3}^2 s_{2+1,4}^2 + 2s_{2-1,3}^2 s_{2+1,4}^2 - s_{2-1,3}^2 s_{2+1,4}^2) \quad (94)$$

#### 3. One massive leg attached



$$0 \quad (95)$$

$$\frac{1}{4} (s_{2+1,4}^2 s_{2+1,4}^2 - s_{2+1,4}^2 s_{2+1,4}^2) (s_{2+1,4}^2 s_{2+1,4}^2 - s_{2+1,4}^2 s_{2+1,4}^2) \quad (96)$$

$$\frac{1}{4} s_{2-1,3}^2 s_{2+1,4}^2 (s_{2+1,4}^2 s_{2+1,4}^2 - s_{2+1,4}^2 s_{2+1,4}^2) \quad (97)$$

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4. One massless, one massive leg attached



$$0 \quad (68)$$

$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b & b+1 \\ b+2 & c-1 & c & q \end{bmatrix} \quad (69)$$

Note that in the previous formula we suppress the terms containing  $x_{2+1a}^2$  which would otherwise cancel a propagator of the underlying topology. When expanded out, the expression above has 12 terms.



$$-\frac{1}{4} \begin{bmatrix} a-2 & a-1 & a & a+1 \\ a+2 & b-1 & b & q \end{bmatrix} \quad (70)$$

In the previous formula we suppress the terms containing  $x_{2+1a}^2$  which would otherwise cancel a propagator of the underlying topology.

5. Two massless legs attached



$$\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b \\ b+1 & b+2 & a-1 & q \end{bmatrix} \quad (71)$$

In the previous formula we suppress the terms containing  $x_{2+1a}^2$  which would otherwise cancel a propagator of the underlying topology.

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$$\frac{1}{4} \left[ \begin{matrix} a-2 & a-1 & a & a+1 \\ a+2 & a+3 & a-3 & q \end{matrix} \right] - \frac{1}{4} \left[ \begin{matrix} a-1 & a & a+1 & a+2 \\ a+3 & a-3 & a-2 & q \end{matrix} \right] -$$

$$\frac{1}{4} \left( -x_{-3,a+1}^2 x_{-2,a+2}^2 x_{-1,a+3}^2 + x_{-2,a-1}^2 x_{-3,a+2}^2 x_{-1,a+3}^2 - \right.$$

$$- x_{-2,-3,a+2}^2 x_{-2,a+1}^2 x_{-1,a+2}^2 + 2x_{-2,a+1}^2 x_{-2,a+2}^2 x_{-1,a+3}^2 +$$

$$+ x_{-2,-3,a+1}^2 x_{-2,a+2}^2 x_{-1,a+2}^2 + x_{-2,-3,a+2}^2 x_{-2,a+1}^2 x_{-1,a+3}^2 -$$

$$- 2x_{-2,a+1}^2 x_{-2,a+2}^2 x_{-1,a+3}^2 + x_{-2,-3,a+2}^2 x_{-2,a+1}^2 x_{-1,a+3}^2 -$$

$$- 2x_{-2,-3,a+2}^2 x_{-2,a+1}^2 x_{-1,a+2}^2 + 2x_{-2,-3,a-1}^2 x_{-2,a+2}^2 x_{-1,a+3}^2 -$$

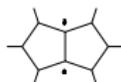
$$- x_{-2,-3,a+2}^2 x_{-2,a+1}^2 x_{-1,a+2}^2 + 1x_{-2,-3,a+1}^2 x_{-2,a+2}^2 x_{-1,a+3}^2 +$$

$$+ 2x_{-2,-3,a+2}^2 x_{-2,a+1}^2 x_{-1,a+2}^2 + 1x_{-2,-3,a-1}^2 x_{-2,a+2}^2 x_{-1,a+3}^2 \Big). \quad (72)$$

We have written down this formula to emphasize how nontrivial it is. We suppress the terms containing  $x_{-2,a+1}^2$  and  $x_{-2,a+2}^2$ , respectively. These terms would otherwise cancel a propagator of the underlying topology. We will see below that the box-pentagon topologies with massless legs attached to the vertices of the edge common to both loops can in fact be seen to originate in double-pentagon topologies, by cancelling some propagators.

## D. Double pentagon topologies

### 1. No legs attached



$$-\frac{1}{4} \left[ \begin{matrix} a & a+1 & b-1 & b & p \\ b & b+1 & a-1 & a & q \end{matrix} \right] \quad (73)$$

In the expansion of the above formula we drop terms that would cancel propagators (in this case, the terms containing  $x_{-2,a}^2$ ,  $x_{-2,b}^2$ ,  $x_{-2,p}^2$ , or  $x_{-2,q}^2$ ). This expression has 6 terms when expanded.

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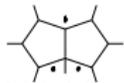
2. One massless leg attached



$$-\frac{1}{4} \begin{bmatrix} a+1 & a+2 & b-1 & b & p \\ b & b+1 & a-1 & a & q \end{bmatrix} \quad (74)$$

In the formula above we drop terms that would cancel propagators (in this case, the terms are  $x_{ab}^2$ ,  $x_{bc}^2$  and  $x_{ca}^2$ ). This expression has 15 terms when expanded.

3. One massive leg attached



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b & p \\ b & b+1 & c-1 & c & q \end{bmatrix} \quad (75)$$

In the formula above we drop terms that would cancel propagators (in this case, the terms containing  $x_{ab}^2$ ,  $x_{bc}^2$  or  $x_{ca}^2$ ). This expression has 16 terms when expanded.

4. Two massless legs attached



$$-\frac{1}{4} \begin{bmatrix} a+1 & a+2 & b-1 & b & p \\ b+1 & b+2 & a-1 & a & q \end{bmatrix} \quad (76)$$

In the formula we drop terms that would cancel propagators (in this case, the terms containing  $x_{ab}^2$ ). This expression has 64 terms when expanded.

5. One massless, one massive leg attached



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b & p \\ b+1 & b+2 & c-1 & c & q \end{bmatrix} \quad (77)$$

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In the formula above we drop terms that would cancel propagators (in this case, the terms containing  $x_{ab}^2$ ). This expression has 78 terms when expanded.

6. Two massive legs attached



$$-\frac{1}{4} \begin{bmatrix} a & a+1 & b-1 & b & p \\ c & c+1 & d-1 & d & q \end{bmatrix} \quad (78)$$

In the formula above we drop terms that would cancel propagators (in this case, the terms containing  $x_{ab}^2$ ). When expanded, the above expression contains 96 terms. The number of conformal drawings is 160 (the number of coefficients unrelaxed by symmetries is lower).

### E. Assembly of the result

As explained in Sec. II, for the MHV amplitudes the ratio between the  $l$ -loop amplitude and the tree-level amplitude can be written as a sum between parity even and parity odd contributions

$$M_l^{(0)} = M_l^{(0),\text{even}} + M_l^{(0),\text{odd}}. \quad (79)$$

Then, the even part can be written

$$M_l^{(0),\text{even}} = -g^{-2D} e^{2\gamma_E} \int d^D x_i d^D x_j \sum_{\sigma \in \mathbb{S}_{2n}} \sum_{i \in \mathbb{S}_{2n}} s_i c_i I_i, \quad (80)$$

where the first sum runs over cyclic and anti-cyclic permutations of the external legs, the second sum runs over all the topologies,  $s_i$  is a symmetry factor associated to topology  $i$ ,  $c_i$  is the numerator of the topology  $i$ , as listed in Sec. IV and  $I_i$  is the denominator or the product of propagators in the topology  $i$ .

Apart from the parity odd part which we have not computed, there is also a contribution which is not detectable from four-dimensional cuts, denoted by  $M_l^{(0),e}$ . This part of the result is such that its integral vanishes in four dimensions, but the integral itself can give contributions to the divergent and finite parts. In Ref. [32], for  $n=6$  case, this part of the result was found to be closely related to  $\mathcal{O}(\epsilon)$  contributions at one loop,  $M_l^{(1),e}$ .

Based on previous computations we expect that the odd part and the  $\mu$  integrals will not be needed in order to compare with the Wilson loop results. The odd parts could be

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## The two-loop MHV amplitudes in $\mathcal{N} = 4$ supersymmetric Yang-Mills theory

C. Vergu\*

*Physics Department, Brown University, Providence, RI 02912, USA*

### Abstract

We compute the even part of the planar two-loop MHV amplitude in  $\mathcal{N} = 4$  supersymmetric Yang-Mills theory, for an arbitrary number of external particles. The answer is expressed as a sum of conformal integrals.

computed by using the leading singularity method (see Ref. [33] and also [34, 35]) or the technique of maximal cuts of Ref. [31]. In order to compute  $M^{2D}$ , one would have to compute  $D$ -dimensional cuts. In practice this is done by computing the cuts of  $\mathcal{N} = 1$  super-Yang-Mills in ten dimensions, dimensionally reduced to  $D$  dimensions.

### V. DISCUSSION

In this paper we computed the even part of the two-loop planar MHV scattering amplitudes in  $\mathcal{N} = 4$  super Yang-Mills. The answer can be expressed in terms of a finite (and relatively small) number of two-loop pseudo-conformal integrals.

A computation of these integrals in dimensional regularization through the finite parts (of order  $\mathcal{O}(\epsilon^0)$ ) would be very interesting and would allow a comparison with the results of Ref. [28], where the corresponding Wilson loop computation was performed.

However, a computation of these integrals seems to be rather difficult. In Ref. [28] the Wilson loop result was expressed in terms of some master integrals called: "hard", "curtain", "cross", "Y" and "factorized cross." These master integrals depend on whether some momenta are zero, massless or massive (this is similar to the situation for scattering amplitudes; in that case also, the value of the integral depends on whether the external legs are massive or massless).

It is interesting to note that for the Wilson loop computation, there are no new master integrals beyond nine sides (this number arises by considering the "hard" integral where all the momenta  $Q_1, Q_2$  and  $Q_3$  are massive). For the scattering amplitude, however, new integrals appear until twelve points, as shown in this paper. It would be interesting to get a deeper understanding of this "mismatch."

The results presented in this paper hint that a different organization of the result may be possible. For example, the coefficients written down using the square brackets symbols can be assembled over a common denominator whose topology is that of a double pentagon. Sometimes, the coefficient of a given topology needs to be split into two contributions which get assembled into different double pentagon topologies (see Eq. (72) for an example).

It is also noteworthy that part of the kissing double boxes coefficient neatly combines with a double pentagon topology after multiplying the numerator and denominator by  $x_{pq}^2$ , while the remaining part has a factorized form. This factorized form is a product of "one-mass"

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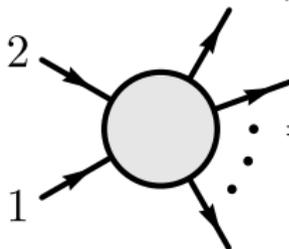
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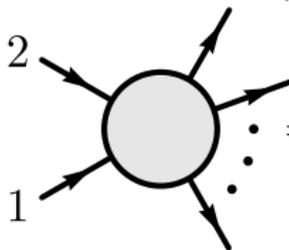
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$$\times \left\{ 1 + \dots \right.$$

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 & \times \left\{ 1 + \sum_{a < b} \text{Diagram} + \dots \right\}
 \end{aligned}$$

The diagram on the left is a circle with  $n$  external legs. Two legs are labeled 1 and 2. The diagram in the sum is a square with a wavy line connecting two vertices, labeled  $a$  and  $b$ .

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## The Two-Loop Hexagon Wilson Loop in $\mathcal{N} = 4$ SYM

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$$G(-\alpha, \alpha^2, \alpha^2, 1; 1) - H(-1, -1, -1, -1, -\frac{1}{\alpha}) + H(-1, -1, 0, -1, -\frac{1}{\alpha}) \quad (G.242)$$

$$\begin{aligned} & - H(-1, -1, 0, 1, \frac{1}{\alpha}) + H(-1, 0, -1, -1, \frac{1}{\alpha}) - H(-1, 0, -1, 1, \frac{1}{\alpha}) \\ & - H(-1, 0, 1, -1, \frac{1}{\alpha}) + H(-1, 0, 1, 1, \frac{1}{\alpha}) - H(0, -1, -1, -1, \frac{1}{\alpha}) \\ & - H(0, -1, -1, 1, \frac{1}{\alpha}) - 2H(0, -1, 0, -1, \frac{1}{\alpha}) + 2H(0, -1, 0, 1, \frac{1}{\alpha}) \\ & - H(0, -1, 1, -1, \frac{1}{\alpha}) + H(0, -1, 1, 1, \frac{1}{\alpha}) - 4H(0, 0, -1, -1, \frac{1}{\alpha}) \\ & + 4H(0, 0, -1, 1, \frac{1}{\alpha}) + 4H(0, 0, 1, -1, \frac{1}{\alpha}) - 4H(0, 0, 1, 1, \frac{1}{\alpha}) \\ & - H(0, 1, -1, -1, \frac{1}{\alpha}) + H(0, 1, -1, 1, \frac{1}{\alpha}) + H(0, 1, 1, -1, \frac{1}{\alpha}) \\ & - H(0, 1, 1, 1, \frac{1}{\alpha}) - 2H(1, 0, -1, -1, \frac{1}{\alpha}) - 2H(1, 0, 0, -1, \frac{1}{\alpha}) \\ & + 4H(1, 0, 0, 1, \frac{1}{\alpha}) - 2H(1, 1, 0, -1, \frac{1}{\alpha}) + H(1, 1, 0, 1, \frac{1}{\alpha}) \end{aligned}$$

$$G(\alpha, \alpha^2, \alpha^2, 1; 1) - H(-1, -1, 0, -1, \frac{1}{\alpha}) + 2H(-1, -1, 0, 1, \frac{1}{\alpha}) \quad (G.243)$$

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### H. The analytic expression of the remainder function

In this appendix we present the full analytic expression of the remainder function. The result is also available in electronic form from [www.arXiv.org](http://www.arXiv.org). Using the notation introduced in Eqs. (3.23) and (5.7), the full expression reads,



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$$\begin{aligned} & \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, 0, 1, 1\right) - \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, 0, \frac{1}{1-s_1}, 1\right) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, 1, 0, 1\right) - \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_1}, 0, 1\right) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_1}, 1, 1\right) - \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_1}, \frac{1}{1-s_1}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, \frac{s_2-1}{s_1+s_2-1}, 1, 1\right) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_1}, \frac{s_2-1}{s_1+s_2-1}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_1}, 1\right) - G\left(\frac{1}{s_1}, 0, 0, \frac{1}{s_2}, 1\right) + \\ & \frac{1}{2}G\left(\frac{1}{s_1}, 0, 0, \frac{1}{s_1+s_2}, 1\right) - G\left(\frac{1}{s_1}, 0, 0, \frac{1}{s_2}, 1\right) + \frac{1}{2}G\left(\frac{1}{s_1}, 0, 0, \frac{1}{s_1+s_2}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{s_1}, 0, \frac{1}{s_1}, \frac{1}{s_1+s_2}, 1\right) - \frac{1}{4}G\left(\frac{1}{s_1}, 0, \frac{1}{s_1}, \frac{1}{s_2}, 1\right) - \frac{1}{4}G\left(\frac{1}{s_1}, 0, \frac{1}{s_2}, \frac{1}{s_1+s_2}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{s_1}, 0, \frac{1}{s_2}, \frac{1}{s_1+s_2}, 1\right) - \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, 0, 1, 1\right) + \\ & \frac{1}{2}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, 1, 1\right) + \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, 0, 1, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, 1, 1\right) + \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, 0, 1, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, 0, 1\right) + \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, 1, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, \frac{1}{1-s_2}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, \frac{1}{1-s_2}, 1\right) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_3-1}{s_2+s_3-1}, \frac{1}{1-s_2}, \frac{1}{1-s_2}, 1\right) - G\left(\frac{1}{s_2}, 0, 0, \frac{1}{s_1}, 1\right) + \\ & \frac{1}{2}G\left(\frac{1}{s_2}, 0, 0, \frac{1}{s_1+s_2}, 1\right) - G\left(\frac{1}{s_2}, 0, 0, \frac{1}{s_3}, 1\right) + \frac{1}{2}G\left(\frac{1}{s_2}, 0, 0, \frac{1}{s_1+s_2}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{s_2}, 0, \frac{1}{s_2}, \frac{1}{s_2+s_3}, 1\right) - \frac{1}{4}G\left(\frac{1}{s_2}, 0, \frac{1}{s_2}, \frac{1}{s_3}, 1\right) - \frac{1}{4}G\left(\frac{1}{s_2}, 0, \frac{1}{s_3}, \frac{1}{s_2+s_3}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{s_2}, 0, \frac{1}{s_3}, \frac{1}{s_2+s_3}, 1\right) - \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 0, 1\right) + \\ & \frac{1}{2}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 1, 1\right) + \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, 0, 1, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 1, 1\right) + \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 0, 1, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 0, 1\right) + \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 1, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, \frac{1}{1-s_3}, 1\right) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{s_4-1}{s_3+s_4-1}, 1, 1\right) - \frac{799^4}{360} + \\ & \frac{1}{4}G\left(\frac{1}{1-s_3}, \frac{s_4-1}{s_3+s_4-1}, \frac{s_4-1}{s_3+s_4-1}, \frac{1}{1-s_3}, 1\right) - G\left(\frac{1}{s_3}, 0, 0, \frac{1}{s_1}, 1\right) - \end{aligned}$$

## The Two-Loop Hexagon Wilson Loop in $\mathcal{N} = 4$ SYM

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# Spiritus Movens: Even More Shocking Simplicity Exists...

Quite recently, del Duca, Duhr, and Smirnov determined the 2-loop, 6-particle amplitude  $\mathcal{A}_6^{(2),2}$  **analytically**—a truly heroic computation on par with Parke and Taylor's

- dimensionally regulating thousands of separately divergent integrals
- final formula: 18 pages of so-called “Goncharov” polylogarithms

$$\begin{aligned}
 & G\left(\frac{1}{s_3}, 0, 0, \frac{1}{s_1}\right) + \frac{1}{2}G\left(\frac{1}{s_3}, 0, 0, \frac{1}{s_1 + s_2}\right) + \frac{1}{2}G\left(\frac{1}{s_3}, 0, 0, \frac{1}{s_2 + s_3}\right) - \\
 & \frac{1}{4}G\left(\frac{1}{s_3}, \frac{1}{s_1}, \frac{1}{s_2}, \frac{1}{s_1 + s_2}\right) - \frac{1}{4}G\left(\frac{1}{s_3}, \frac{1}{s_2}, \frac{1}{s_1 + s_2}, \frac{1}{s_1}\right) + \frac{1}{4}G\left(\frac{1}{s_3}, 0, \frac{1}{s_2}, \frac{1}{s_1 + s_2}\right) - \\
 & \frac{1}{4}G\left(\frac{1}{s_3}, \frac{1}{s_1}, \frac{1}{s_2}, \frac{1}{s_2 + s_3}\right) - \frac{1}{24}{}^2\mathcal{G}\left(\frac{1}{1-s_1}, s_{221}; 1\right) + \frac{1}{2}{}^2\mathcal{G}\left(\frac{1}{1-s_2}, s_{121}; 1\right) + \\
 & \frac{1}{8}{}^2\mathcal{G}\left(\frac{1}{1-s_1}, s_{121}; 1\right) - \frac{1}{24}{}^2\mathcal{G}\left(\frac{1}{1-s_1}, s_{221}; 1\right) + \frac{1}{8}{}^2\mathcal{G}\left(\frac{1}{1-s_2}, s_{121}; 1\right) + \\
 & \frac{1}{8}{}^2\mathcal{G}\left(\frac{1}{1-s_2}, s_{221}; 1\right) - \frac{1}{24}{}^2\mathcal{G}\left(\frac{1}{1-s_3}, s_{121}; 1\right) + \frac{1}{8}{}^2\mathcal{G}\left(\frac{1}{1-s_3}, s_{221}; 1\right) + \\
 & \frac{1}{2}{}^2\mathcal{G}\left(\frac{1}{1-s_3}, s_{221}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, 0, \frac{1}{1-s_1}, s_{121}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, 0, \frac{1}{1-s_2}, s_{121}; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, 0, \frac{1}{1-s_2}, s_{221}; 1\right) - \frac{1}{4}\mathcal{G}\left(0, 0, \frac{1}{1-s_2}, s_{221}; 1\right) - \frac{1}{4}\mathcal{G}\left(0, 0, \frac{1}{1-s_3}, s_{221}; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, 0, \frac{1}{1-s_3}, s_{221}; 1\right) - \frac{1}{4}\mathcal{G}\left(0, 0, s_{221}, \frac{1}{1-s_1}\right) + \mathcal{G}\left(0, 0, s_{121}, 0; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, 0, s_{121}, \frac{1}{1-s_1}\right) + \mathcal{G}\left(0, 0, s_{121}, 0; 1\right) - \frac{1}{4}\mathcal{G}\left(0, 0, s_{221}, \frac{1}{1-s_2}\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, 0, s_{221}, \frac{1}{1-s_2}\right) + \frac{1}{2}\mathcal{G}\left(0, 0, s_{121}, \frac{1}{1-s_3}\right) + \mathcal{G}\left(0, 0, s_{221}, 0; 1\right) - \\
 & \frac{1}{2}\mathcal{G}\left(0, 0, s_{221}, \frac{1}{1-s_3}\right) - \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_1}, 0, s_{121}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_2}, 0, s_{121}; 1\right) - \\
 & \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_3}, 0, s_{121}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_1}, \frac{1}{1-s_2}, s_{121}; 1\right) - \\
 & \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_2}, \frac{1}{1-s_3}, s_{121}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_1}, \frac{1}{1-s_3}, s_{121}; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_1}, s_{121}, 1; 1\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_2}, s_{121}, \frac{1}{1-s_1}\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{121}, 1; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_1}, s_{121}, \frac{1}{1-s_1}\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_2}, 0, s_{121}; 1\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, 0, s_{121}; 1\right) - \\
 & \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_2}, \frac{1}{1-s_3}, s_{121}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_2}, \frac{1}{1-s_3}, s_{221}; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, 1; 1\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, \frac{1}{1-s_2}\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, 1; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_2}, s_{221}, \frac{1}{1-s_2}\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_2}, 0, s_{221}; 1\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, 0, s_{221}; 1\right) - \\
 & \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_3}, \frac{1}{1-s_3}, s_{221}; 1\right) - \frac{1}{2}\mathcal{G}\left(0, \frac{1}{1-s_3}, \frac{1}{1-s_3}, s_{221}; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, 1; 1\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, \frac{1}{1-s_3}\right) - \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, 1; 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, \frac{1}{1-s_3}, s_{221}, \frac{1}{1-s_3}\right) - \frac{1}{4}\mathcal{G}\left(0, s_{121}, 0, \frac{1}{1-s_1}\right) - \frac{1}{4}\mathcal{G}\left(0, s_{121}, \frac{1}{1-s_1}, 0; 1\right) + \\
 & \frac{1}{4}\mathcal{G}\left(0, s_{221}, \frac{1}{1-s_1}, 1; 1\right) - \frac{1}{4}\mathcal{G}\left(0, s_{221}, \frac{1}{1-s_1}, \frac{1}{1-s_1}\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, s_{221}, \frac{s_2-1}{s_1+s_2-1}, 1; 1\right) + \frac{1}{4}\mathcal{G}\left(0, s_{221}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_1}, 1\right) - \\
 & \frac{1}{4}\mathcal{G}\left(0, s_{221}, \frac{1}{s_3}, 0; 1\right) - \frac{1}{4}\mathcal{G}\left(0, s_{221}, 0, \frac{1}{1-s_2}\right) - \frac{1}{4}\mathcal{G}\left(0, s_{221}, \frac{1}{s_1}, 0; 1\right) -
 \end{aligned}$$

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$$\begin{aligned} & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{223}, \frac{1}{1-s_1}, \frac{1}{1-s_1} \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{s_2-1}{s_1+s_2-1}, 1; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{223}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_1} \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{s_2-1}{s_1+s_2-1}, 0 \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 0, 0 \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 0, 1; 1 \right) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 0, \frac{1}{1-s_1} \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 1, 0; 1 \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 1, \frac{1}{1-s_1} \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{1}{1-s_1}, 0; 1 \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{1}{1-s_1}, 1; 1 \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{1}{1-s_1}, \frac{1}{1-s_1} \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 0, 0; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 0, 1; 1 \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 0, \frac{1}{1-s_1} \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 1, 0; 1 \right) - \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, 1, \frac{1}{1-s_1} \right) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{1}{1-s_1}, 0; 1 \right) - \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{1}{1-s_1}, 1; 1 \right) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_1}, s_{123}, \frac{1}{1-s_1}, \frac{1}{1-s_1} \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, 0, 0, s_{223}; 1 \right) - \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, 0, 0, s_{223}; 1 \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, 0, \frac{1}{1-s_2}, s_{223}; 1 \right) - \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, 0, \frac{1}{1-s_2}, s_{223}; 1 \right) - \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, 0, s_{223}, 1; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, 0, s_{223}, \frac{1}{1-s_2} \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, 0, s_{223}, 1; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, 0, s_{223}, \frac{1}{1-s_2} \right) - \frac{3}{4} \mathcal{G} \left( \frac{1}{1-s_2}, \frac{1}{1-s_2}, \frac{1}{1-s_2}, s_{223}; 1 \right) - \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, \frac{1}{1-s_2}, 0, s_{223}; 1 \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, \frac{1}{1-s_2}, \frac{1}{1-s_2}, s_{223}; 1; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, \frac{1}{1-s_2}, s_{223}, \frac{1}{1-s_2} \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, \frac{1}{1-s_2}, s_{223}, 1; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, \frac{1}{1-s_2}, s_{223}, \frac{1}{1-s_2} \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 0, 1; 1 \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 0, \frac{1}{1-s_2} \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 1, 0; 1 \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 1, \frac{1}{1-s_2} \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, \frac{1}{1-s_2}, 0; 1 \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, \frac{1}{1-s_2}, 1; 1 \right) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, \frac{1}{1-s_2}, \frac{1}{1-s_2} \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, \frac{s_2-1}{s_1+s_2-1}, 1; 1 \right) - \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, \frac{s_2-1}{s_1+s_2-1}, \frac{1}{1-s_2} \right) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 0, 0; 1 \right) - \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 0, 1; 1 \right) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 0, \frac{1}{1-s_2} \right) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 1, 0; 1 \right) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-s_2}, s_{223}, 1, \frac{1}{1-s_2} \right) + \end{aligned}$$





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- dimensionally regulating thousands of separately divergent integrals
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$$\begin{aligned} & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_1-1}{s_2+s_3-1}, 1; 1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_1-1}{s_2+s_3-1}, \frac{1}{1-s_2}\right) H(0; u_1) + \frac{1}{2}G\left(\frac{1}{s_2}, 0, \frac{1}{s_1}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{s_2}, 0, \frac{1}{s_1+s_2}\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{s_2}, \frac{1}{s_1}, \frac{1}{s_1+s_2}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_1-1}{s_2+s_3-1}, 0; 1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_1-1}{s_2+s_3-1}, \frac{1}{1-s_2}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, \frac{s_1-1}{s_2+s_3-1}, \frac{s_1-1}{s_1+s_2}\right) H(0; u_1) + \frac{1}{2}G\left(\frac{1}{s_2}, 0, \frac{1}{s_1}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{s_2}, 0, \frac{1}{s_1+s_2}\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{s_2}, \frac{1}{s_1}, \frac{1}{s_1+s_2}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(0, \frac{1}{1-s_1}, v_{123}; 1\right) H(0; u_2) + \frac{1}{4}G\left(0, \frac{1}{1-s_1}, v_{123}; 1\right) H(0; u_1) + \\ & \frac{1}{4}G\left(0, \frac{1}{1-s_2}, v_{231}; 1\right) H(0; u_1) - \frac{1}{4}G\left(0, \frac{1}{1-s_2}, v_{231}; 1\right) H(0; u_1) + \\ & \frac{1}{4}G\left(0, \frac{1}{1-s_1}, v_{321}; 1\right) H(0; u_1) - \frac{1}{4}G\left(0, \frac{1}{1-s_1}, v_{321}; 1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(0, v_{231}, \frac{1}{s_1}\right) H(0; u_1) - \frac{1}{4}G\left(0, v_{231}, \frac{1}{1-s_2}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(0, v_{321}, \frac{1}{1-s_1}\right) H(0; u_1) - \frac{1}{4}G\left(0, v_{321}, \frac{1}{s_1+s_2-1}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(0, v_{123}, \frac{1}{1-s_1}\right) H(0; u_1) - \frac{1}{4}G\left(0, v_{123}, \frac{1}{1-s_1}\right) H(0; u_1) - \\ & \frac{1}{2}G\left(0, v_{231}, \frac{1}{1-s_2}\right) H(0; u_1) + \frac{1}{2}G\left(0, v_{321}, \frac{1}{1-s_2}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_1}, 0, v_{231}; 1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-s_1}, 0, v_{321}; 1\right) H(0; u_1) + \\ & \frac{1}{2}G\left(\frac{1}{1-s_1}, \frac{1}{1-s_1}, v_{123}; 1\right) H(0; u_1) + \frac{1}{2}G\left(\frac{1}{1-s_1}, \frac{1}{1-s_1}, v_{123}; 1\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_1}, v_{123}, 1; 1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-s_1}, v_{123}, \frac{1}{1-s_1}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, 0, v_{231}; 1\right) H(0; u_1) - \frac{1}{4}G\left(\frac{1}{1-s_2}, 0, v_{231}; 1\right) H(0; u_1) + \\ & \frac{1}{2}G\left(\frac{1}{1-s_2}, \frac{1}{1-s_2}, v_{231}; 1\right) H(0; u_1) - \frac{1}{2}G\left(\frac{1}{1-s_2}, \frac{1}{1-s_2}, v_{231}; 1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, v_{231}, 1; 1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-s_2}, v_{231}, \frac{1}{s_1}\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{1-s_2}, v_{321}, \frac{1}{1-s_1}\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-s_2}, v_{321}, \frac{1}{s_1}\right) H(0; u_1) + \\ & \frac{1}{2}G\left(\frac{1}{1-s_2}, v_{231}, \frac{1}{1-s_2}\right) H(0; u_1) - \frac{1}{2}G\left(\frac{1}{1-s_2}, v_{231}, 0; 1\right) H(0; u_1) - \\ & \frac{1}{2}G\left(\frac{1}{1-s_2}, v_{321}, \frac{1}{1-s_2}\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-s_2}, 0, v_{321}; 1\right) H(0; u_1) - \end{aligned}$$

## The Two-Loop Hexagon Wilson Loop in $\mathcal{N} = 4$ SYM

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$$\begin{aligned} & \frac{1}{4} \mathcal{G} \left( 0, \frac{1}{1-u_1} v_{123}; 1 \right) H(0; u_2) - \frac{1}{4} \mathcal{G} \left( 0, \frac{1}{1-u_2} v_{123}; 1 \right) H(0; u_1) + \\ & \frac{1}{4} \mathcal{G} \left( 0, \frac{1}{1-u_2} v_{231}; 1 \right) H(0; u_2) + \frac{1}{4} \mathcal{G} \left( 0, \frac{1}{1-u_1} v_{231}; 1 \right) H(0; u_1) - \\ & \frac{1}{4} \mathcal{G} \left( 0, \frac{1}{1-u_1} v_{321}; 1 \right) H(0; u_2) + \frac{1}{4} \mathcal{G} \left( 0, \frac{1}{1-u_2} v_{321}; 1 \right) H(0; u_1) + \\ & \frac{1}{4} \mathcal{G} \left( 0, u_{123} \frac{1}{1-u_1}; 1 \right) H(0; u_2) - \frac{1}{4} \mathcal{G} \left( 0, u_{123} \frac{1}{u_1+u_2-1}; 1 \right) H(0; u_2) - \\ & \frac{1}{4} \mathcal{G} \left( 0, u_{321} \frac{1}{u_2}; 1 \right) H(0; u_2) - \frac{1}{4} \mathcal{G} \left( 0, u_{112} \frac{1}{1-u_2}; 1 \right) H(0; u_2) + \\ & \frac{1}{2} \mathcal{G} \left( 0, v_{123} \frac{1}{1-u_1}; 1 \right) H(0; u_2) + \frac{1}{2} \mathcal{G} \left( 0, v_{231}; \frac{1}{1-u_2}; 1 \right) H(0; u_1) + \\ & \frac{1}{4} \mathcal{G} \left( 0, v_{231} \frac{1}{1-u_2}; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( 0, v_{321}; \frac{1}{1-u_1}; 1 \right) H(0; u_1) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, 0, v_{123}; 1 \right) H(0; u_2) - \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_2}, 0, v_{123}; 1 \right) H(0; u_1) + \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1} \frac{1}{1-u_1}, v_{123}; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1} \frac{1}{1-u_1}, v_{123}; 1 \right) H(0; u_2) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, u_{123}, \frac{u_2-1}{u_1+u_2-1}; 1 \right) H(0; u_2) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, u_{123}, 0; 1 \right) H(0; u_2) + \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{123}, \frac{1}{1-u_1}; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{123}, 0; 1 \right) H(0; u_2) - \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{123}, \frac{1}{1-u_1}; 1 \right) H(0; u_1) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_2}, 0, v_{123}; 1 \right) H(0; u_2) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_2}, v_{123}, \frac{1}{1-u_1}; 1 \right) H(0; u_2) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_2}, v_{123}, 1; 1 \right) H(0; u_2) + \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_2}, v_{231}; \frac{1}{1-u_1}; 1 \right) H(0; u_1) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_2}, v_{231}, 1; 1 \right) H(0; u_1) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_2}, v_{321}; \frac{1}{1-u_1}; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_2}, 0, v_{321}; 1 \right) H(0; u_2) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_2}, 0, v_{321}; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, u_{123}, v_{123}; 1 \right) H(0; u_2) + \\ & \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, 1-u_1, v_{123}; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{123}, 1; 1 \right) H(0; u_2) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, u_{123}, \frac{1}{u_2}; 1 \right) H(0; u_2) + \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, u_{123}, \frac{1}{1-u_2}; 1 \right) H(0; u_2) - \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, v_{123}, 0; 1 \right) H(0; u_2) - \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{123}, \frac{1}{1-u_1}; 1 \right) H(0; u_2) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, v_{231}, 0; 1 \right) H(0; u_2) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{231}, \frac{1}{1-u_1}; 1 \right) H(0; u_2) + \\ & \frac{1}{4} \mathcal{G} \left( \frac{1}{1-u_1}, v_{321}, 0; 1 \right) H(0; u_2) + \frac{1}{2} \mathcal{G} \left( \frac{1}{1-u_1}, v_{321}, \frac{1}{1-u_1}; 1 \right) H(0; u_2) - \\ & \frac{1}{4} \mathcal{G} \left( v_{123}, \frac{1}{1-u_1}; 1 \right) H(0; u_2) + \frac{1}{2} \mathcal{G} \left( v_{123}, \frac{1}{1-u_1}; 1; 1 \right) H(0; u_2) - \end{aligned}$$

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$$\begin{aligned} & \frac{1}{4}G\left(v_{122}, 1, \frac{1}{1-v_{122}}\right) H(0; u_2) - \frac{1}{4}G\left(v_{122}, \frac{1}{1-v_{122}}, 1; 1\right) H(0; u_2) + \\ & \frac{1}{4}G\left(v_{212}, 1, \frac{1}{1-v_{212}}\right) H(0; u_2) + \frac{1}{4}G\left(v_{212}, \frac{1}{1-v_{212}}, 1; 1\right) H(0; u_2) + \\ & \frac{1}{4}G\left(v_{221}, 1, \frac{1}{1-v_{221}}\right) H(0; u_2) + \frac{1}{4}G\left(v_{221}, \frac{1}{1-v_{221}}, 1; 1\right) H(0; u_2) - \\ & \frac{3}{4}G\left(v_{122}, 1, \frac{1}{1-v_{122}}\right) H(0; u_2) - \frac{3}{4}G\left(v_{212}, \frac{1}{1-v_{212}}, 1; 1\right) H(0; u_2) + \\ & \frac{1}{4}G\left(v_{221}, 1, \frac{1}{1-v_{221}}\right) H(0; u_2) + \frac{1}{4}G\left(v_{221}, \frac{1}{1-v_{221}}, 1; 1\right) H(0; u_2) + \\ & \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) H(0; u_2) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) H(0; u_2) + \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}\right) H(0; u_1) H(0; u_2) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, u_{222}; 1\right) H(0; u_1) H(0; u_2) - \frac{1}{4}G\left(\frac{1}{1-u_1}, v_{222}; 1\right) H(0; u_1) H(0; u_2) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, v_{221}; 1\right) H(0; u_1) H(0; u_2) + \frac{5}{24}G^2 H(0; u_1) H(0; u_2) - \\ & \frac{1}{4}G\left(0, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \frac{1}{4}G\left(0, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \\ & \frac{1}{4}G\left(0, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1, \frac{1}{1-v_{122}}\right) H(0; u_1) - \frac{1}{4}G\left(0, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) - \\ & \frac{3}{4}G\left(0, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) - \frac{1}{4}G\left(0, \frac{u_2-1}{u_1+u_2-1}, \frac{1}{1-v_{221}}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, 1, \frac{1}{u_1}\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}, 1; 1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}, \frac{1}{1-u_1}\right) H(0; u_1) + \frac{1}{2}G\left(\frac{1}{u_1}, 0, \frac{1}{u_1}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{u_1}, 0, \frac{1}{u_1}, u_1+u_2-1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_2-1}{u_1+u_2-1}, 0; 1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_2-1}{u_1+u_2-1}, \frac{1}{1-u_2}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_2-1}{u_1+u_2-1}, \frac{u_2-1}{u_1+u_2-1}\right) H(0; u_1) + \frac{1}{2}G\left(\frac{1}{u_2}, 0, \frac{1}{u_2}\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{u_2}, 0, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) - \\ & \frac{3}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \\ & \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1}, \frac{1}{u_2}, u_1+u_2-1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{1}{1-u_1}, v_{222}; 1\right) H(0; u_1) + \frac{1}{4}G\left(\frac{1}{1-u_1}, v_{222}; 1\right) H(0; u_1) - \\ & \frac{1}{4}G\left(\frac{1}{1-u_1}, v_{221}; 1\right) H(0; u_1) + \frac{1}{4}G\left(0, \frac{1}{1-u_2}, v_{221}; 1\right) H(0; u_1) + \end{aligned}$$

## The Two-Loop Hexagon Wilson Loop in $\mathcal{N} = 4$ SYM

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$$\begin{aligned} & \frac{1}{12} {}^2H(0, 1; u_1 + u_2) + \frac{1}{12} {}^2H(0, 1; u_1) + \frac{1}{4} {}^2H(0; u_1) H(0; u_2) H\left(0, 1; \frac{u_1 + u_2 - 1}{u_1 - 1}\right) - \\ & \frac{1}{24} {}^2H(0, 1; \frac{u_1 + u_2 - 1}{u_1 - 1}) + \frac{1}{12} {}^2H(0, 1; (u_1 + u_2)) - \frac{1}{24} {}^2H\left(0, 1; \frac{u_2 + u_3 - 1}{u_3 - 1}\right) + \\ & \frac{1}{12} {}^2H(0, 1; (u_2 + u_3)) - \frac{1}{2} G\left(0, \frac{1}{u_1 + u_2}\right) H(1, 0; u_1) - \\ & \frac{1}{2} G\left(0, \frac{1}{u_1 + u_2}\right) H(1, 0; u_1) + \frac{1}{4} G\left(\frac{1}{u_1}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_1) + \\ & \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_1) + \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_1) + \\ & \frac{1}{4} G\left(\frac{1}{1 - u_2}, \frac{1}{u_1 + u_2 - 1}\right) H(1, 0; u_1) + \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_2 + u_3}\right) H(1, 0; u_1) - \\ & \frac{1}{4} G\left(\frac{1}{1 - u_3}, \frac{1}{u_2 + u_3 - 1}\right) H(1, 0; u_1) - \frac{3}{4} H(0, 0; u_2) H(1, 0; u_1) - \frac{3}{4} H(0, 0; u_2) H(1, 0; u_1) + \\ & \frac{1}{4} H\left(0, 1; \frac{u_1 + u_2 - 1}{u_1 - 1}\right) H(1, 0; u_1) - \frac{1}{4} {}^2H(1, 0; u_1) - \frac{1}{2} G\left(0, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) - \\ & \frac{1}{2} G\left(0, \frac{1}{u_2 + u_3}\right) H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{1 - u_1}, \frac{1}{u_1 + u_2 - 1}\right) H(1, 0; u_2) + \\ & \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) + \\ & \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_2 + u_3}\right) H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_2 + u_3}\right) H(1, 0; u_2) - \\ & \frac{1}{4} G\left(\frac{1}{1 - u_1}, \frac{1}{u_2 + u_3 - 1}\right) H(1, 0; u_2) - \frac{3}{4} H(0, 0; u_1) H(1, 0; u_2) - \frac{3}{4} H(0, 0; u_1) H(1, 0; u_2) + \\ & \frac{1}{4} H\left(0, 1; \frac{u_2 + u_3 - 1}{u_3 - 1}\right) H(1, 0; u_2) - \frac{1}{4} H(1, 0; u_1) H(1, 0; u_2) - \frac{1}{3} {}^2H(1, 0; u_2) - \\ & \frac{1}{2} G\left(0, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) - \frac{1}{2} G\left(0, \frac{1}{u_2 + u_3}\right) H(1, 0; u_2) + \\ & \frac{1}{4} G\left(\frac{1}{u_1}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{1 - u_2}, \frac{1}{u_2 + u_3 - 1}\right) H(1, 0; u_2) + \\ & \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) - \frac{1}{4} {}^2H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_1 + u_2}\right) H(1, 0; u_2) + \\ & \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_2 + u_3}\right) H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{u_2}, \frac{1}{u_2 + u_3}\right) H(1, 0; u_2) + \\ & \frac{1}{4} G\left(\frac{1}{1 - u_2}, \frac{1}{u_2 + u_3 - 1}\right) H(1, 0; u_2) + \frac{1}{4} G\left(\frac{1}{1 - u_2}, \frac{1}{u_2 + u_3 - 1}\right) H(1, 0; u_2) + \\ & \frac{1}{4} H(0; u_1) H(0; u_2) H(1, 0; u_1) - \frac{1}{4} H(0; u_1) H(1, 0; u_2) - \frac{3}{4} H(0; u_2) H(1, 0; u_1) + \\ & \frac{1}{4} H\left(0, 1; \frac{u_1 + u_2 - 1}{u_1 - 1}\right) H(1, 0; u_1) - \frac{1}{4} H(1, 0; u_1) H(1, 0; u_2) - \frac{1}{4} H(1, 0; u_2) H(1, 0; u_1) + \\ & \frac{1}{12} {}^2H(1, 1; u_1) + \frac{1}{24} {}^2H(1, 1; u_2) + \frac{1}{24} {}^2H(1, 1; u_3) + \frac{1}{2} H(0; u_2) H(0, 0; u_1) + \\ & \frac{1}{2} H(0; u_2) H(0, 0; u_3) + \frac{1}{2} H(0; u_1) H(0, 0; u_2) - \frac{1}{2} H(0; u_2) H\left(0, 0, 1; \frac{u_1 + u_2 - 1}{u_2 - 1}\right) - \\ & \frac{1}{2} H(0; u_1) H\left(0, 0, 1; \frac{u_1 + u_2 - 1}{u_2 - 1}\right) - H(0; u_1) H(0, 0, 1; (u_1 + u_2)) - \\ & H(0; u_2) H(0, 0, 1; (u_1 + u_2)) - \frac{1}{2} H(0; u_1) H\left(0, 0, 1; \frac{u_1 + u_2 - 1}{u_1 - 1}\right) - \end{aligned}$$

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$$\begin{aligned} & \frac{1}{2} H(0; w_2) H\left(0, 0, 1; \frac{w_1 + w_2 - 1}{w_1 - 1}\right) - H(0; w_1) H(0, 0, 1; (w_1 + w_2)) - \\ & H(0; w_2) H(0, 0, 1; (w_1 + w_2)) - \frac{1}{2} H(0; w_1) H\left(0, 0, 1; \frac{w_2 + w_3 - 1}{w_2 - 1}\right) - \\ & \frac{1}{2} H(0; w_1) H\left(0, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}\right) - H(0; w_2) H(0, 0, 1; (w_2 + w_3)) - \\ & H(0; w_2) H(0, 0, 1; (w_2 + w_3)) - \frac{1}{2} H(0; w_2) H(0, 1, 0; w_2) - \frac{1}{2} H(0; w_2) H(0, 1, 0; w_2) - \\ & \frac{1}{2} H(0; w_1) H(0, 1, 0; w_1) + \frac{1}{4} H(0; w_2) H\left(0, 1, 1; \frac{w_1 + w_2 - 1}{w_2 - 1}\right) - \\ & \frac{1}{4} H(0; w_2) H\left(0, 1, 1; \frac{w_1 + w_2 - 1}{w_2 - 1}\right) + \frac{1}{4} H(0; w_1) H\left(0, 1, 1; \frac{w_2 + w_3 - 1}{w_1 - 1}\right) - \\ & \frac{1}{4} H(0; w_2) H\left(0, 1, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}\right) - \frac{1}{4} H(0; w_1) H\left(0, 1, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}\right) + \\ & \frac{1}{2} H(0; w_2) H\left(0, 1, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}\right) + \frac{1}{2} H(0; w_2) H(1, 0, 0; w_1) - \frac{1}{2} H(0; w_2) H(1, 0, 0; w_1) - \\ & \frac{1}{2} H(0; w_1) H(1, 0, 0; w_2) + \frac{1}{2} H(0; w_1) H(1, 0, 0; w_2) + \frac{1}{2} H(0; w_1) H(1, 0, 0; w_3) - \\ & \frac{1}{2} H(0; w_2) H(1, 0, 0; w_3) - \frac{1}{2} H(0; w_2) H\left(1, 0, 1; \frac{w_1 + w_2 - 1}{w_2 - 1}\right) - \\ & \frac{1}{4} H(0; w_2) H\left(1, 0, 1; \frac{w_1 + w_2 - 1}{w_1 - 1}\right) - \frac{1}{4} H(0; w_1) H\left(1, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}\right) - \\ & 7H(0, 0, 0; w_1) - 7H(0, 0, 0; w_2) - 7H(0, 0, 0; w_3) + \frac{3}{2} H(0, 0, 0, 1; \frac{w_1 + w_2 - 1}{w_1 - 1}) + \\ & 3H(0, 0, 0, 1; (w_1 + w_2)) + \frac{3}{2} H(0, 0, 0, 1; \frac{w_1 + w_2 - 1}{w_1 - 1}) + 3H(0, 0, 0, 1; (w_1 + w_2)) + \\ & \frac{3}{2} H(0, 0, 0, 1; \frac{w_2 + w_3 - 1}{w_2 - 1}) + 3H(0, 0, 0, 1; (w_2 + w_3)) + \frac{3}{2} H(0, 0, 0, 1; w_1) + \\ & \frac{3}{2} H(0, 0, 1, 0; w_2) + \frac{3}{2} H(0, 0, 1, 0; w_2) - \frac{1}{2} H(0, 1, 0, 0; w_1) - \frac{1}{2} H(0, 1, 0, 0; w_2) - \\ & \frac{1}{2} H(0, 1, 0, 0; w_3) + \frac{1}{2} H(0, 1, 0, 1; \frac{w_1 + w_2 - 1}{w_2 - 1}) + \frac{1}{2} H(0, 1, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + \\ & \frac{1}{2} H(0, 1, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + H(0, 1, 1, 0; w_1) + H(0, 1, 1, 0; w_2) + H(0, 1, 1, 0; w_3) - \\ & \frac{1}{4} H(0, 1, 1, 1; \frac{w_1 + w_2 - 1}{w_2 - 1}) - \frac{1}{4} H(0, 1, 1, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) - \\ & \frac{1}{4} H(0, 1, 1, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + H(1, 0, 0, 1; \frac{w_1 + w_2 - 1}{w_2 - 1}) + H(1, 0, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + \\ & H(1, 0, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + 2H(1, 0, 1, 0; w_1) + 2H(1, 0, 1, 0; w_2) + 2H(1, 0, 1, 0; w_3) + \\ & \frac{1}{4} H(1, 1, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + \frac{1}{4} H(1, 1, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + \\ & \frac{1}{4} H(1, 1, 0, 1; \frac{w_2 + w_3 - 1}{w_3 - 1}) + \frac{1}{2} H(1, 1, 1, 0; w_1) + \frac{1}{2} H(1, 1, 1, 0; w_2) - \\ & \frac{1}{24} {}^2H(0; w_2) H\left(1; \frac{w_1}{w_{123}}\right) - \frac{1}{24} {}^2H(0; w_1) H\left(1; \frac{w_2}{w_{231}}\right) + \frac{1}{24} {}^2H(0; w_2) H\left(1; \frac{w_3}{w_{312}}\right) + \\ & \frac{1}{8} {}^2H(0; w_2) H\left(1; \frac{w_1}{w_{123}}\right) - \frac{1}{8} {}^2H(0; w_1) H\left(1; \frac{w_2}{w_{231}}\right) + \frac{1}{24} {}^2H(0; w_2) H\left(1; \frac{w_3}{w_{123}}\right) - \end{aligned}$$



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$$\begin{aligned} & \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) - \\ & \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) - \\ & \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_1)H\left(0,1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_1)H\left(0,1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{1}{w_{123}}\right) - \\ & \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) - \\ & \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{1}{w_{123}}\right) - \\ & \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{1}{w_{123}}\right) - \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{1}{w_{123}}\right) + H(0;u_2)H\left(1,1,1;\frac{1}{w_{123}}\right) - H(0;u_1)H\left(1,1,1;\frac{1}{w_{123}}\right) - \\ & H(0;u_1)H\left(1,1,1;\frac{1}{w_{123}}\right) + H(0;u_1)H\left(1,1,1;\frac{1}{w_{123}}\right) + H(0;u_1)H\left(1,1,1;\frac{1}{w_{123}}\right) - \\ & H(0;u_2)H\left(1,1,1;\frac{1}{w_{123}}\right) - \frac{3}{2}H\left(0,0,0,1;\frac{1}{w_{123}}\right) - \frac{3}{2}H\left(0,0,0,1;\frac{1}{w_{123}}\right) - \\ & \frac{3}{2}H\left(0,0,0,1;\frac{1}{w_{123}}\right) - 3H\left(0,0,0,1;\frac{1}{w_{123}}\right) - 3H\left(0,0,0,1;\frac{1}{w_{123}}\right) - 3H\left(0,0,0,1;\frac{1}{w_{123}}\right) - \\ & \frac{1}{2}H\left(0,0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(0,0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(0,0,1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{2}H\left(0,0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(0,0,1,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(0,0,1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}H\left(0,1,1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}H\left(0,1,1,1;\frac{1}{w_{123}}\right) + G_5H(0;u_1) + G_5H(0;u_2) + G_5H(0;u_1) + \\ & \frac{5}{2}G_5H(1;u_1) + \frac{5}{2}G_5H(1;u_2) + \frac{5}{2}G_5H(1;u_1) + \frac{1}{2}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{2}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{2}G_5H\left(1,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(1,0,0,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(1,0,0,1;\frac{1}{w_{123}}\right) - \frac{1}{2}H\left(1,0,0,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \\ & \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \frac{1}{4}G_5H\left(1,1;\frac{1}{w_{123}}\right) + \end{aligned}$$

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$$\frac{1}{2}H\left(1,1,0,1;\frac{1}{x_{23}}\right) + \frac{3}{2}H\left(1,1,1,1;\frac{1}{x_{23}}\right) + \frac{3}{2}H\left(1,1,1,1;\frac{1}{x_{34}}\right) + \frac{3}{2}H\left(1,1,1,1;\frac{1}{x_{12}}\right)$$

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$$\frac{1}{2}H\left(1,1,0,1; \frac{1}{x_{23}}\right) + \frac{3}{2}H\left(1,1,1,1; \frac{1}{x_{23}}\right) + \frac{3}{2}H\left(1,1,1,1; \frac{1}{x_{34}}\right) + \frac{3}{2}H\left(1,1,1,1; \frac{1}{x_{12}}\right)$$

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