

Terrestrial gamma-ray flashes, antimatter and hadrons correlated to lightning events

Christoph Köhn
Olivier Chanrion, Torsten Neubert
11.11.2015



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Antimatter caught streaming from thunderstorms on Earth

By Jason Palmer

Science and technology reporter, BBC News, Seattle



NASA

Electrons racing up electric field lines give rise to light, then particles, then light

A space telescope has accidentally spotted thunderstorms on Earth producing beams of antimatter.

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Other species?

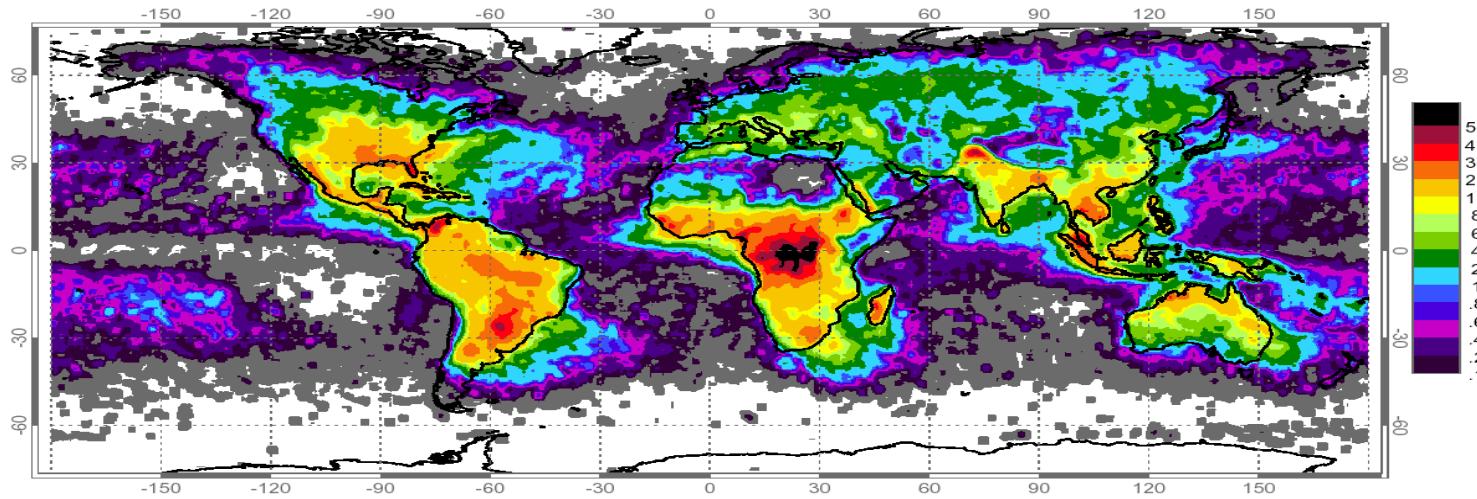


Electrons racing up electric field lines give rise to light, then particles, then light

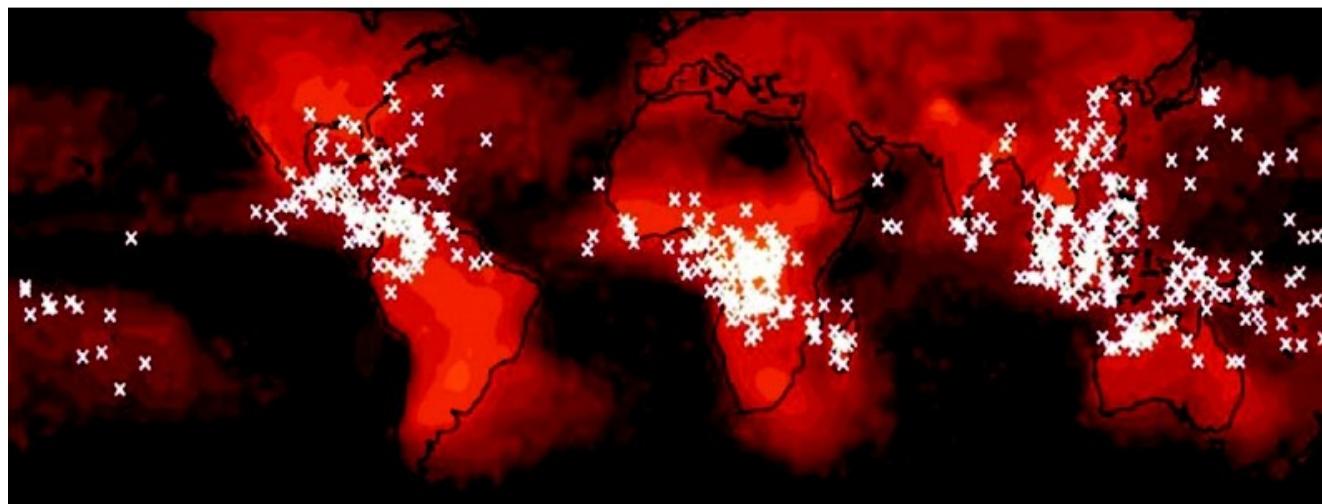
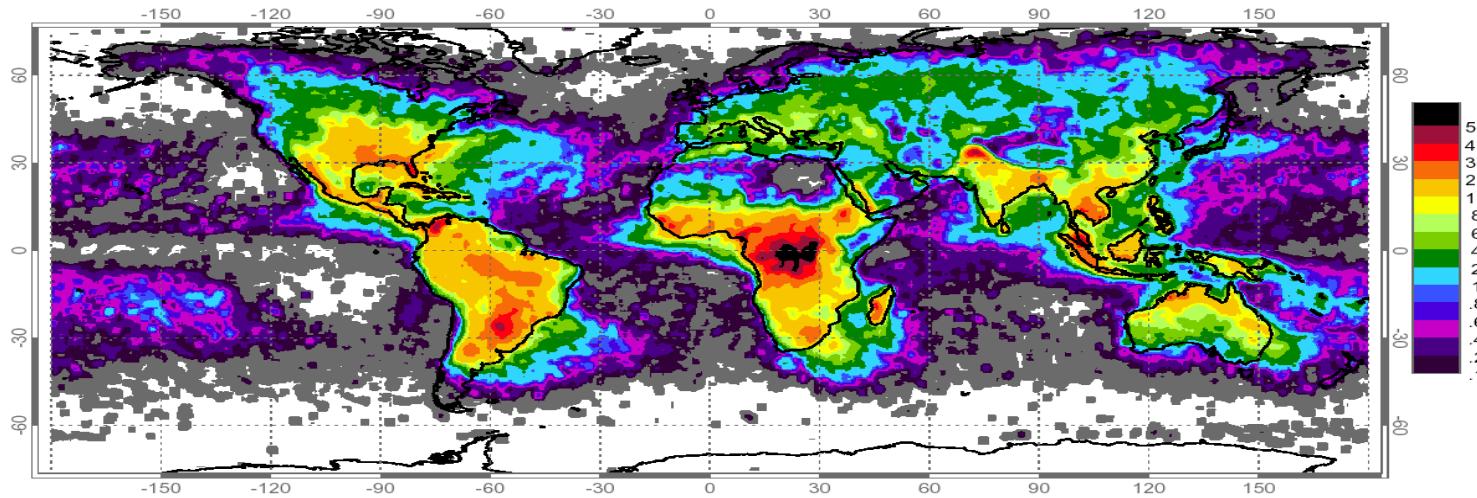
A space telescope has accidentally spotted thunderstorms on Earth producing beams of antimatter.

Related stories

Lightning and gamma-rays

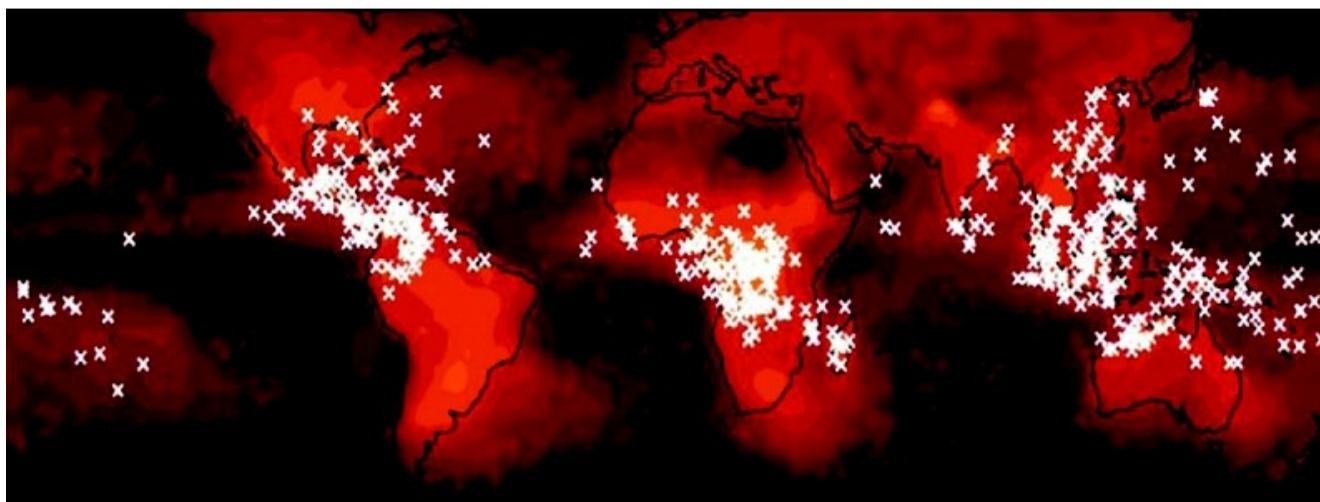
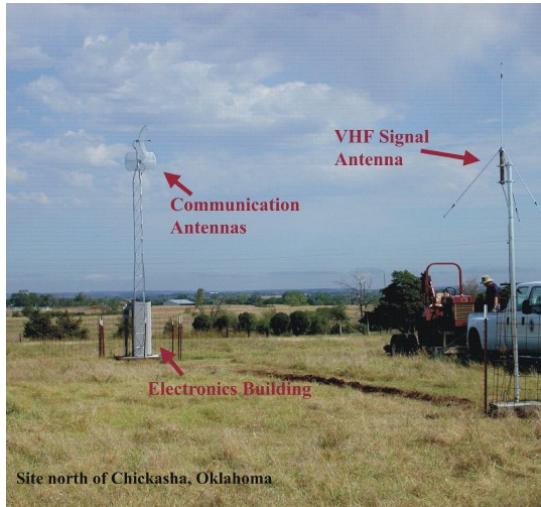


Lightning and gamma-rays

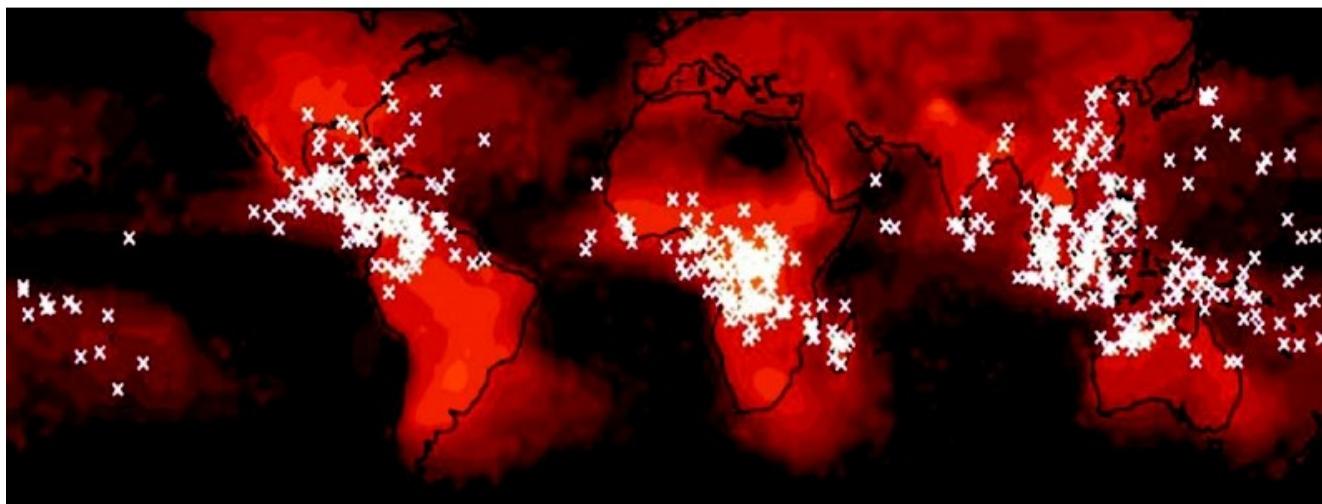
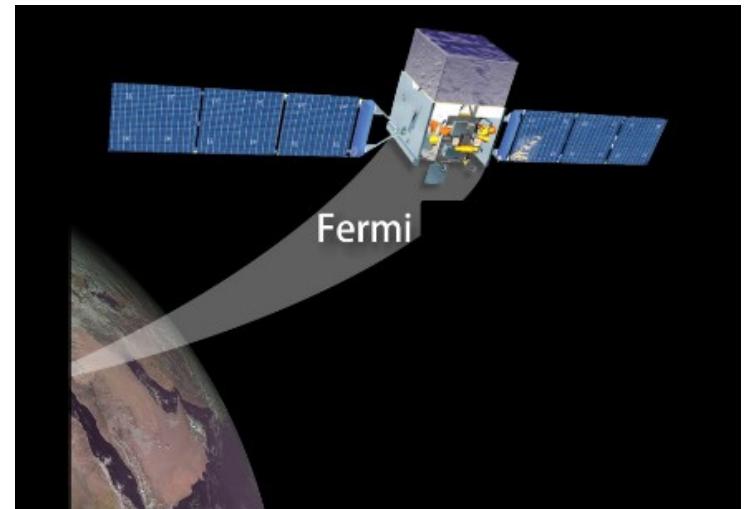
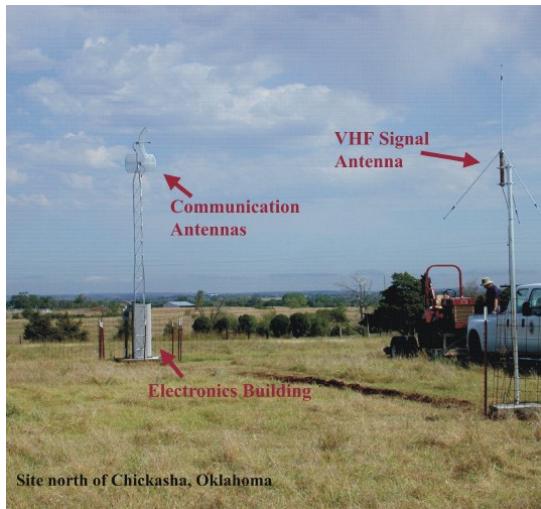


=> Terrestrial
gamma-ray flashes
(TGFs)

Lightning and gamma-rays



Lightning and gamma-rays

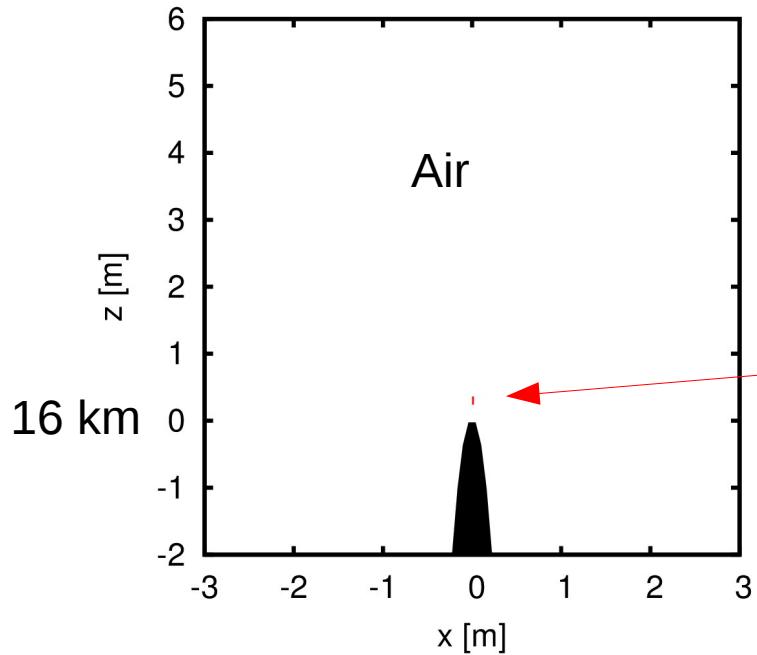


E_γ to 40 MeV (Fermi)
 E_γ to 100 MeV (Agile)

Stepping lightning leader



Test case: Negative stepped lightning leader



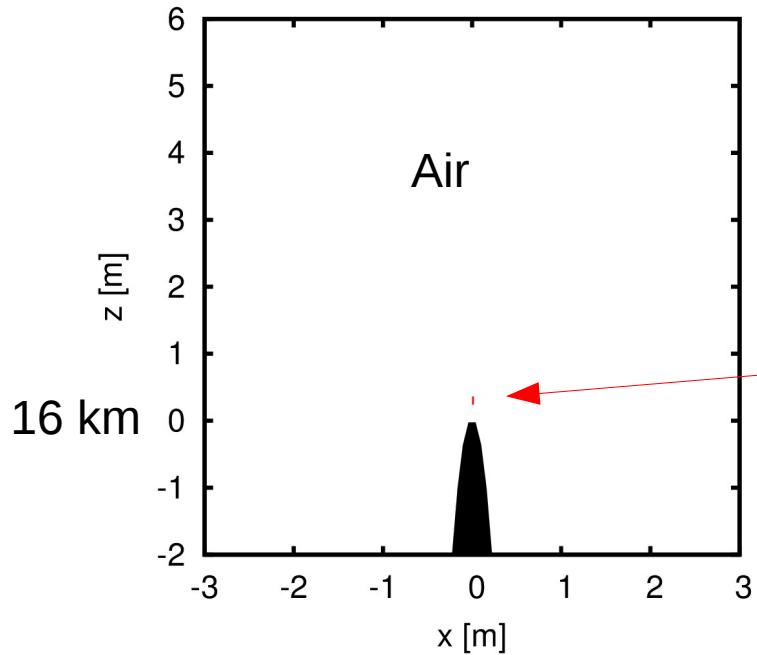
A hot, propagating and stepping plasma channel transporting charge through air

Here: $L = 4 \text{ km}$, $r = 1 \text{ cm}$

- Test electron(s) 30 cm ahead of leader tip

[Xu, Celestin and Pasko, Geophys. Res. Lett., vol. 39, L08801, 2012]

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Here: $L = 4 \text{ km}$, $r = 1 \text{ cm}$

- Test electron(s) 30 cm ahead of leader tip
- Calculate electric field of leader

[Xu, Celestin and Pasko, Geophys. Res. Lett., vol. 39, L08801, 2012]

Electric field of a leader

$$E_0 = 0.5 \text{ kV/cm}$$

$$L = 4 \text{ km}$$

$$r_c = 1 \text{ cm}$$

Simulate negative leader as an ellipsoid

Potential of a charged ellipsoid in an ambient field:

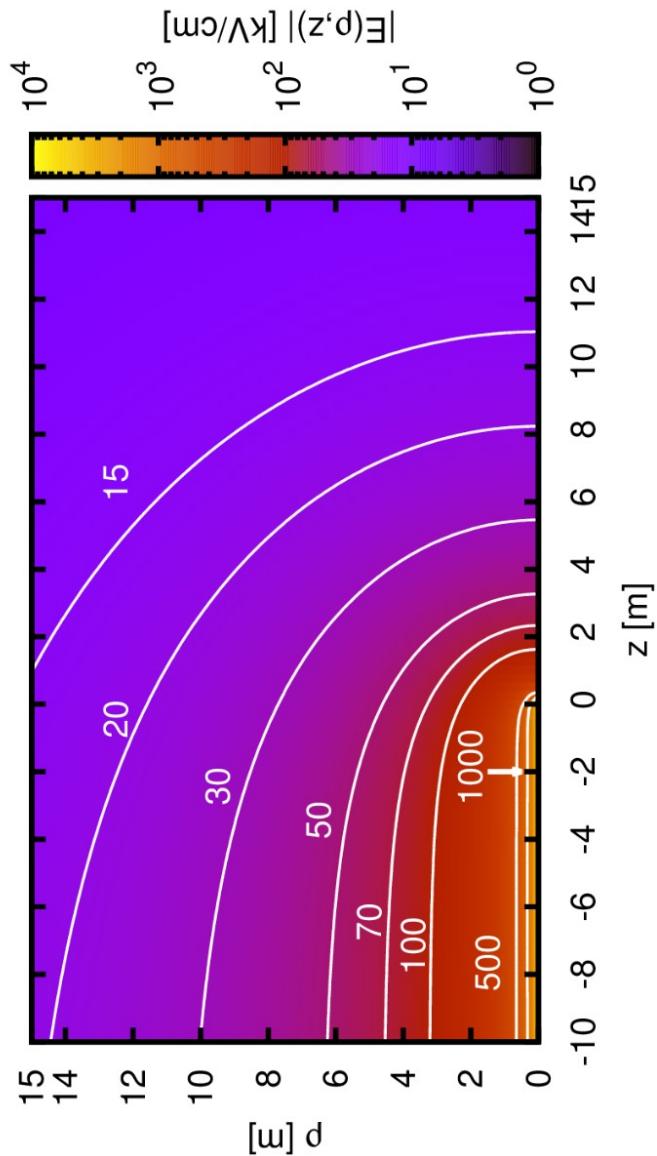
$$\Phi(\vec{r}) = -E_0 z \left[1 - \frac{\int_{\xi(\vec{r})}^{\infty} \frac{ds}{(s+L^2)^{3/2}(s+a^2)}}{\int_0^{\infty} \frac{ds}{(s+L^2)^{3/2}(s+a^2)}} \right]$$

where $\xi(\vec{r})$ is one of the ellipsoidal coordinates

[Landau and Lifschitz, Electrodynamics of Continuous Media, 1984]

=> Large field enhancement close to the leader tip

Electric field of a leader



$$\begin{aligned}E_0 &= 0.5 \text{ kV/cm} \\L &= 4 \text{ km} \\r_c &= 1 \text{ cm}\end{aligned}$$

Simulate negative leader as an ellipsoid

Potential of a charged ellipsoid in an ambient field:

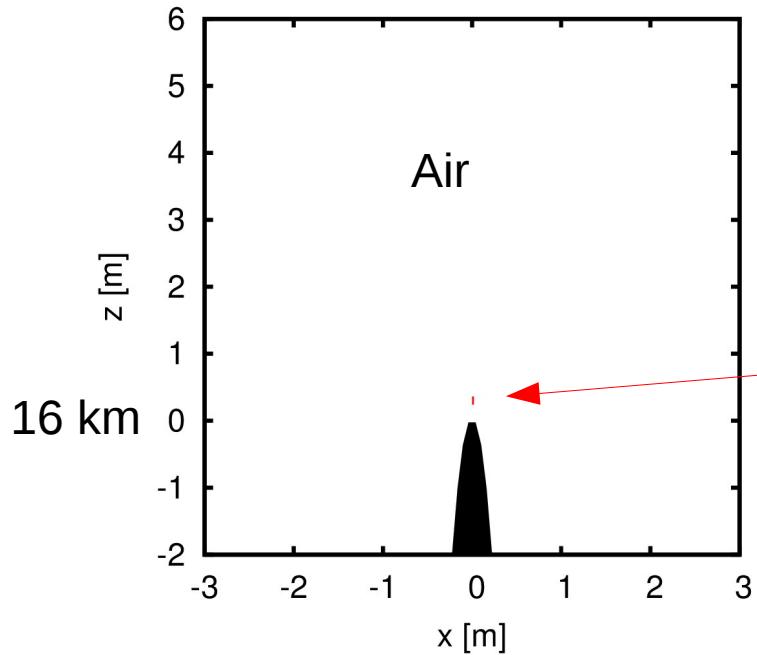
$$\Phi(\vec{r}) = -E_0 z \left[1 - \frac{\int_{\xi(\vec{r})}^{\infty} \frac{ds}{(s+L^2)^{3/2}(s+a^2)}}{\int_0^{\infty} \frac{ds}{(s+L^2)^{3/2}(s+a^2)}} \right]$$

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A hot, propagating and stepping plasma channel transporting charge through air

Here: $L = 4 \text{ km}$, $r = 1 \text{ cm}$

- Test electron(s) 30 cm ahead of leader tip
- Calculate electric field of leader
- Accelerate electrons in this field

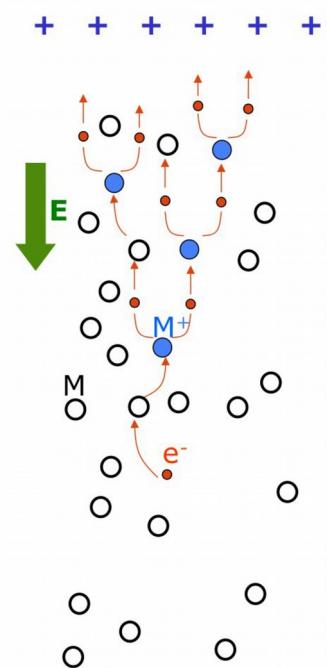
[Xu, Celestin and Pasko, Geophys. Res. Lett., vol. 39, L08801, 2012]

3D relativistic MC Code

- Relativistic leap-frog between two collisions:

$$\vec{r}_{n+1} = \vec{r}_n + \Delta t \vec{v}_{n+\frac{1}{2}} \quad \text{with} \quad \vec{p}_{n+\frac{1}{2}} = \vec{p}_{n-\frac{1}{2}} + q \Delta t \vec{E}(\vec{r}_n, t_n)$$

$$\vec{v}_{n+\frac{1}{2}} = \frac{c^2}{E_{n+\frac{1}{2}}} \cdot \vec{p}_{n+\frac{1}{2}} \quad E_{n+\frac{1}{2}} = \sqrt{c^2 |\vec{p}_{n+\frac{1}{2}}|^2 + m^2 c^4}$$



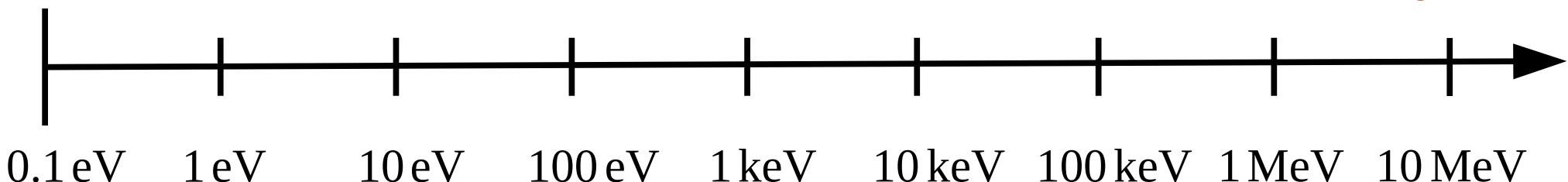
Collisions of electrons with air molecules:

Electron impact ionization

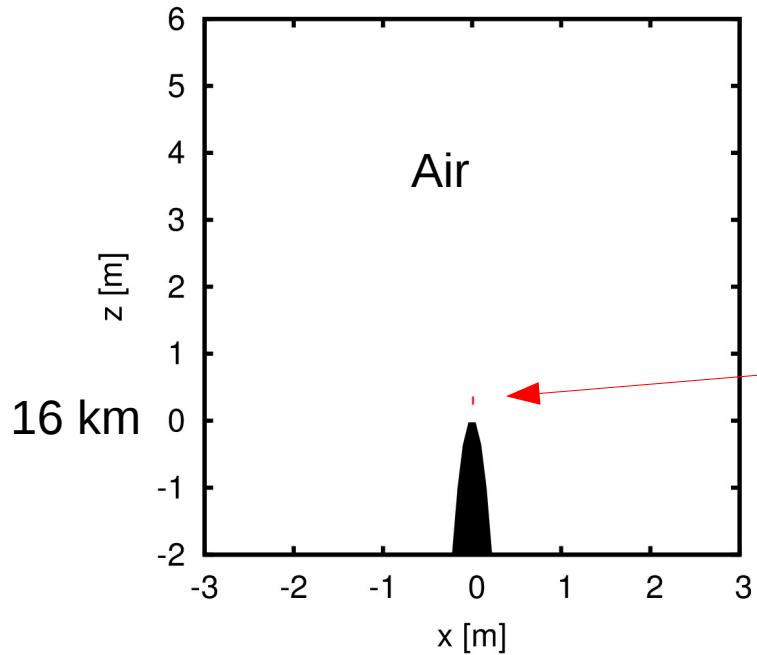
Attachment Excitation

Elastic scattering

Bremsstrahlung



Test case: Negative stepped lightning leader



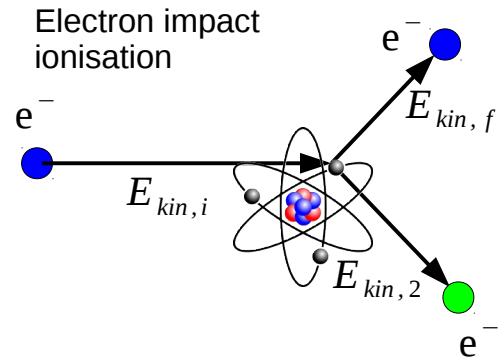
A hot, propagating and stepping plasma channel transporting charge through air

Here: $L = 4 \text{ km}$, $r = 1 \text{ cm}$

- Test electron(s) 30 cm ahead of leader tip
- Calculate electric field of leader
- Accelerate electrons in this field
- **Create new electrons and Bremsstrahlung photons**

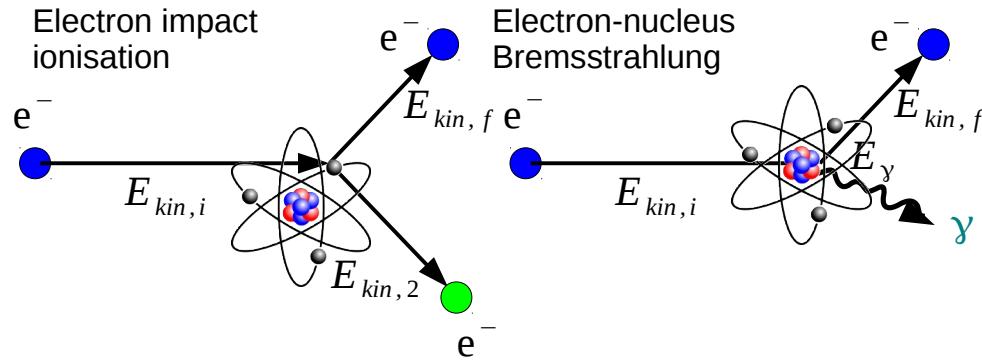
[Xu, Celestin and Pasko, Geophys. Res. Lett., vol. 39, L08801, 2012]

Source of electrons and photons



[Y. K. Kim, J. Paulo Santos, 2000.
Phys. Rev. A, vol. 62, 052710]

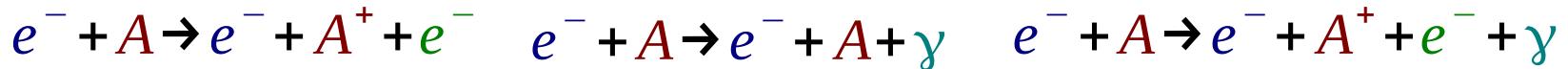
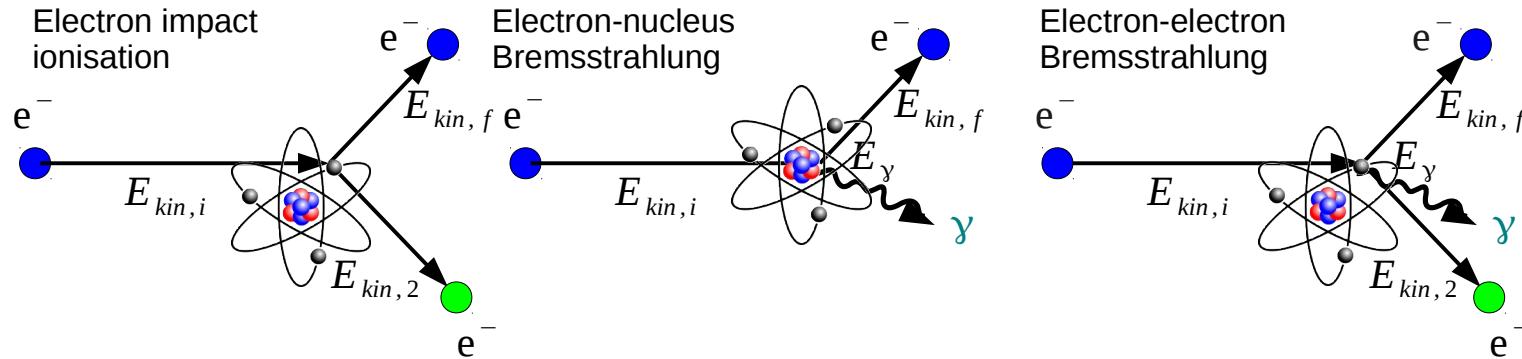
Source of electrons and photons



[Y. K. Kim, J. Paulo Santos, 2000.
Phys. Rev. A, vol. 62, 052710]

[C. Köhn and U. Ebert, 2014. Atmos.
Res., vol. 135-136, pp. 432-465]

Source of electrons and photons



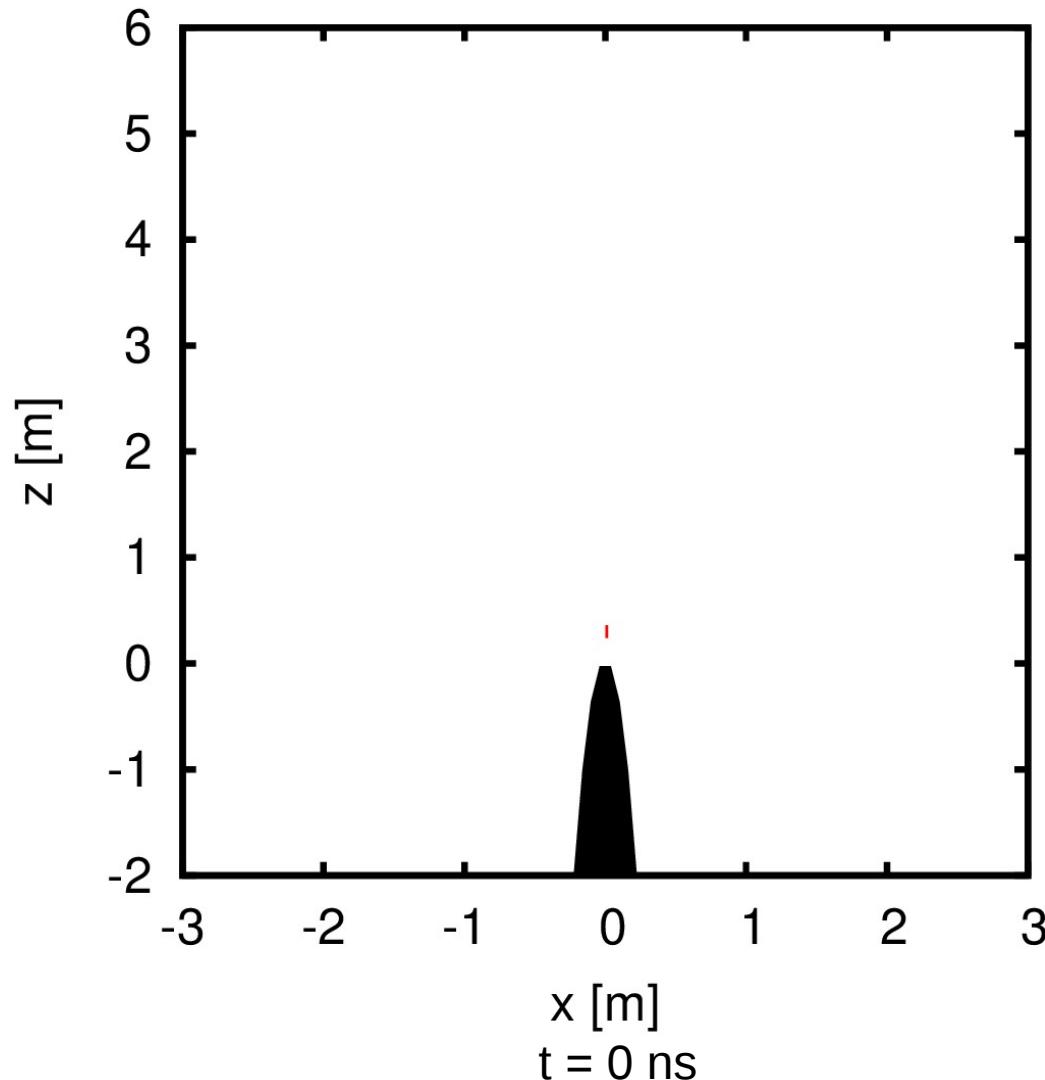
[Y. K. Kim, J. Paulo Santos, 2000. Phys. Rev. A, vol. 62, 052710]

[C. Köhn and U. Ebert, 2014. Atmos. Res., vol. 135-136, pp. 432-465]

[F. Tessier and I. Kawrakow, 2007. NIM Phys. Res. B, vol. 266, pp. 625-634]

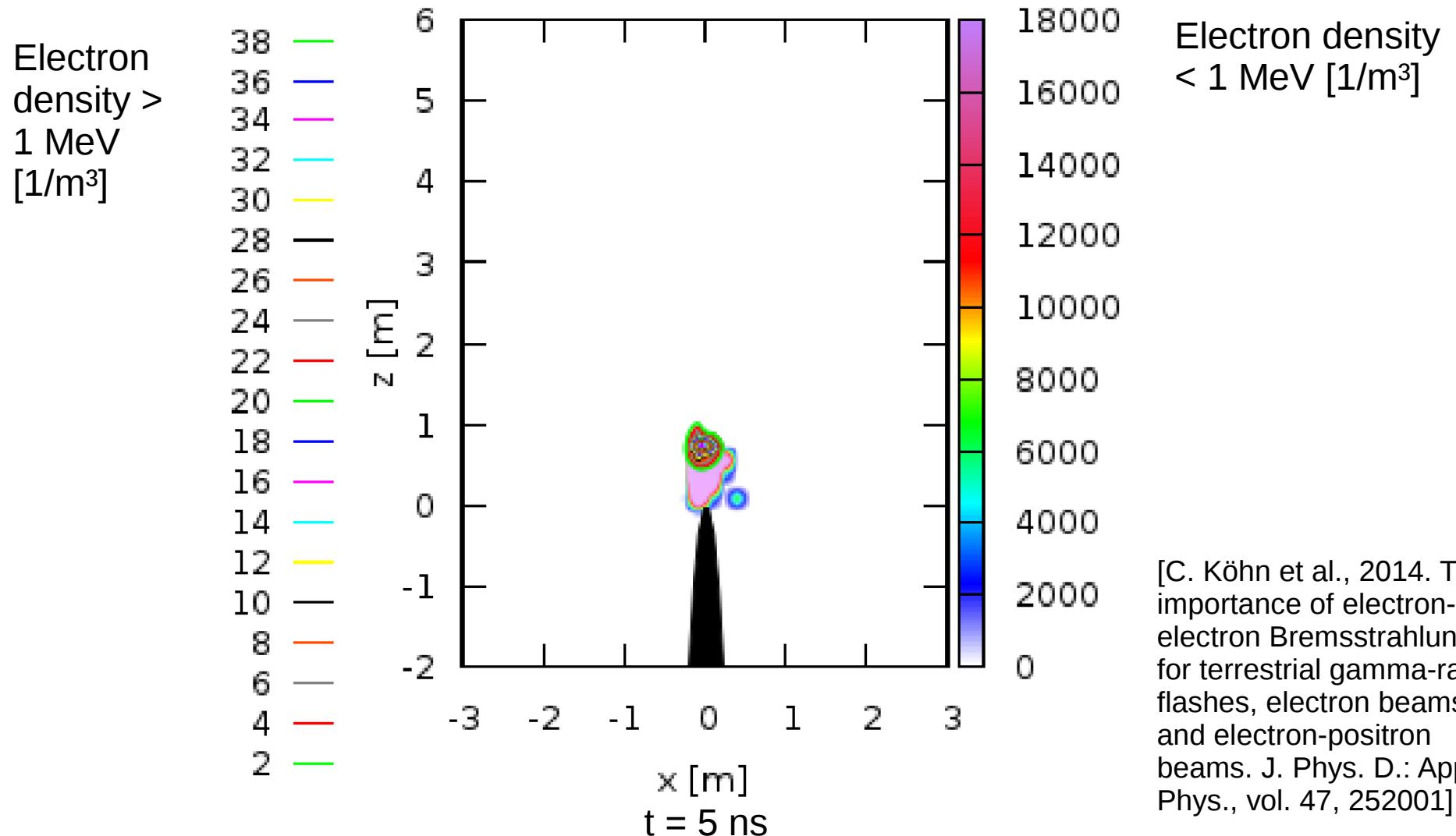
Dominant for energies above 1 MeV
 => Enrichment of high-energy electrons
 => More high-energy photons

Temporal evolution of electrons



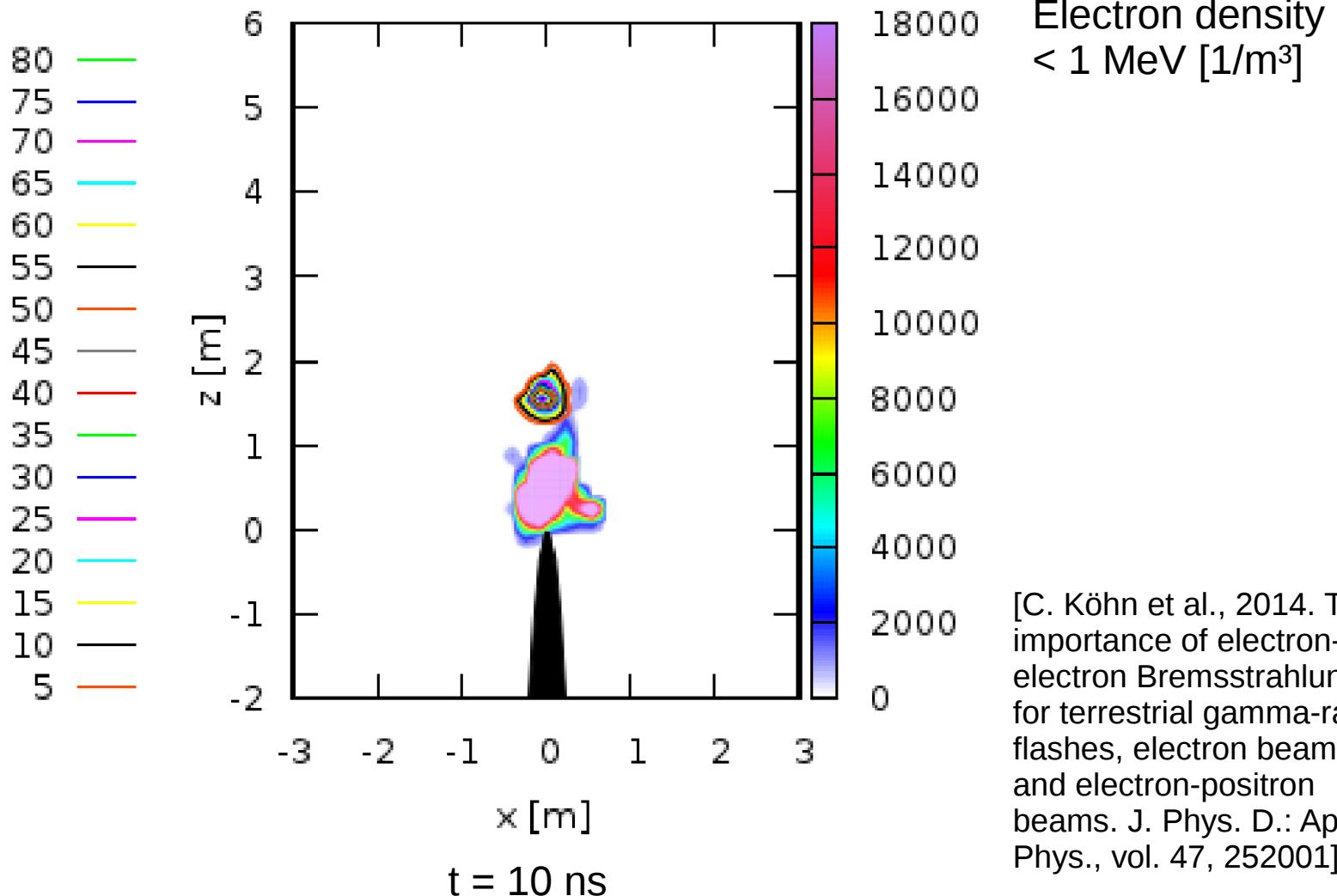
[C. Köhn et al., 2014. The importance of electron-electron Bremsstrahlung for terrestrial gamma-ray flashes, electron beams and electron-positron beams. *J. Phys. D.: Appl. Phys.*, vol. 47, 252001]

Temporal evolution of electrons



Temporal evolution of electrons

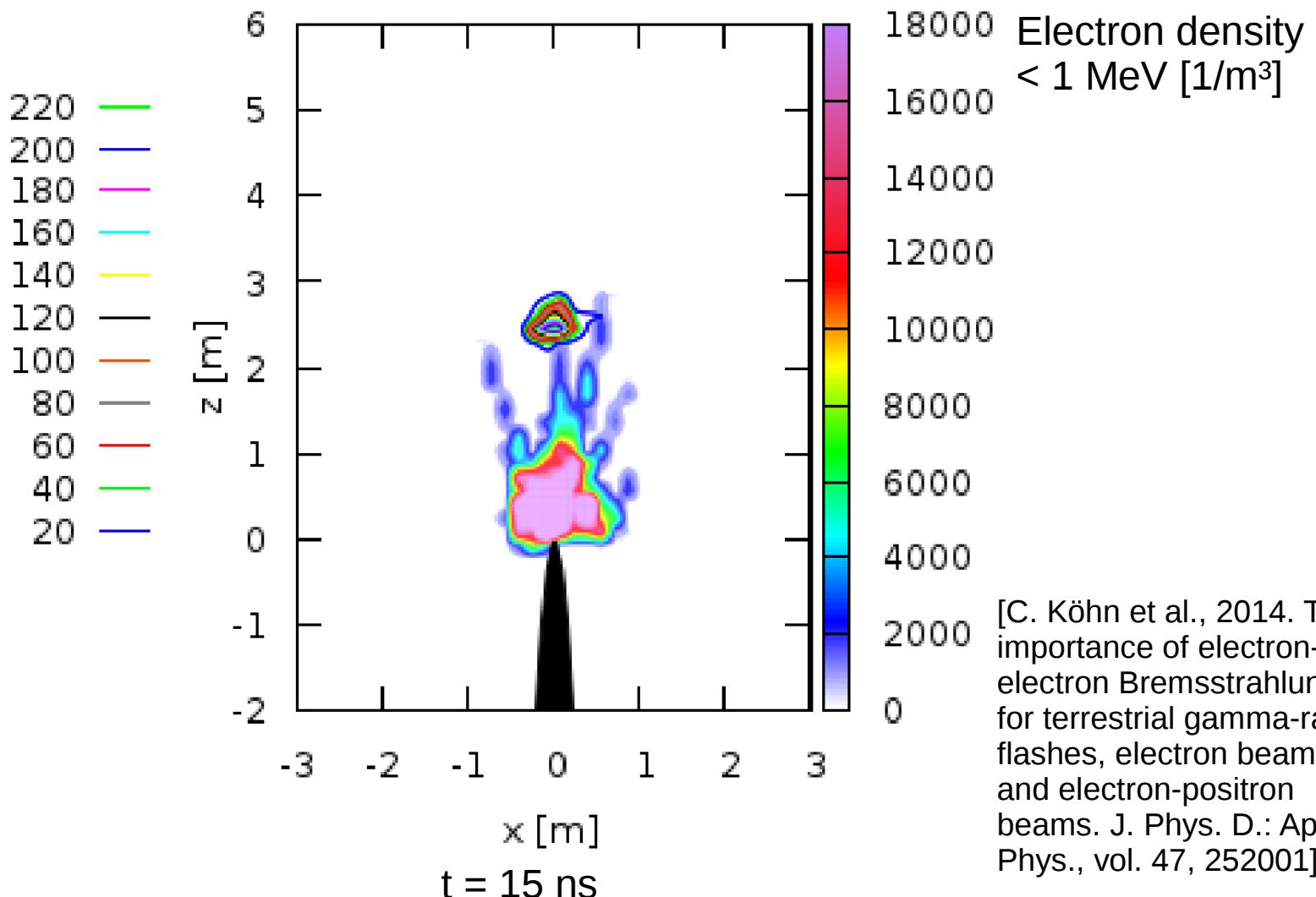
Electron
density >
1 MeV
[1/m³]



[C. Köhn et al., 2014. The importance of electron-electron Bremsstrahlung for terrestrial gamma-ray flashes, electron beams and electron-positron beams. J. Phys. D.: Appl. Phys., vol. 47, 252001]

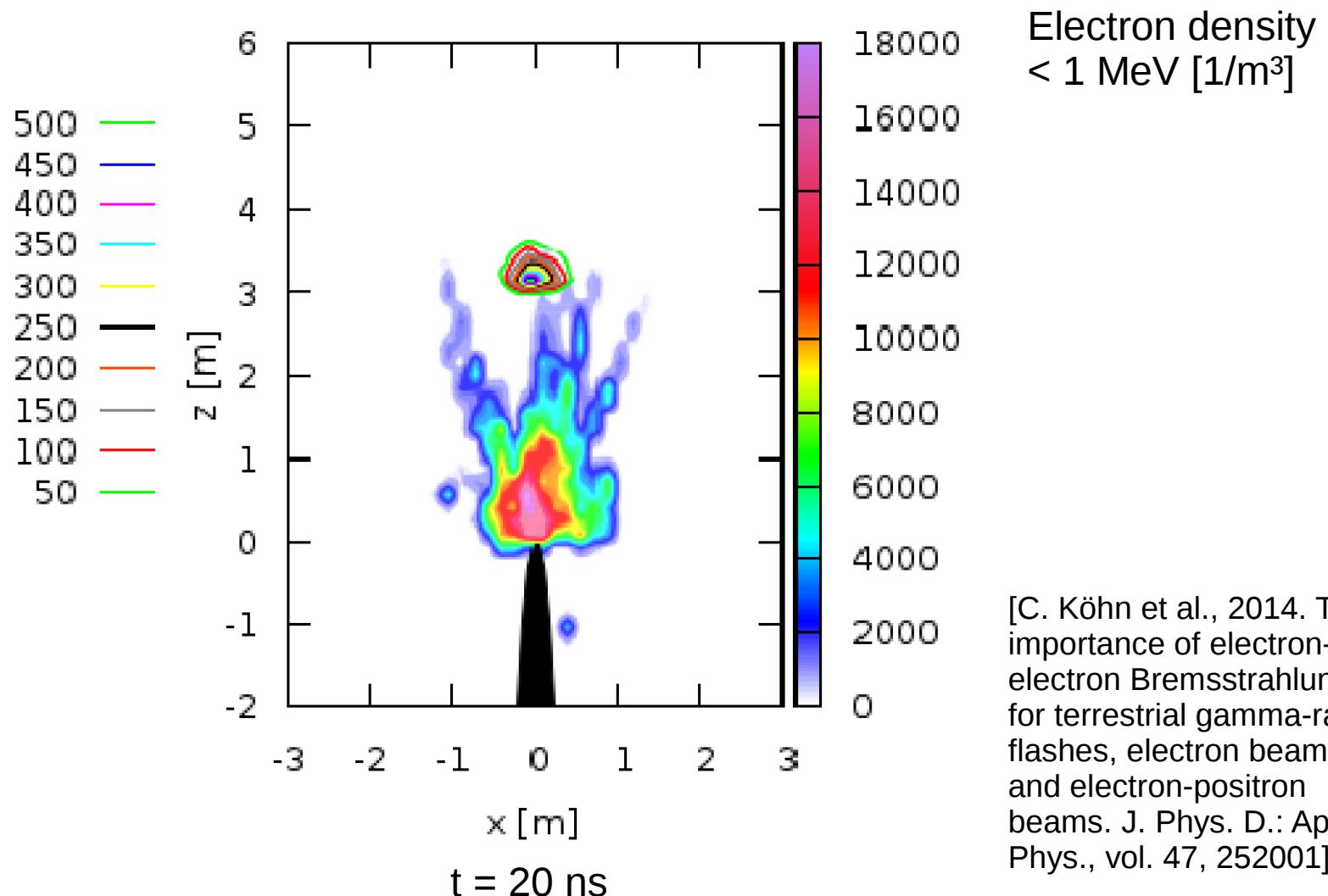
Temporal evolution of electrons

Electron density > 1 MeV [1/m³]



Temporal evolution of electrons

Electron density > 1 MeV [1/m³]



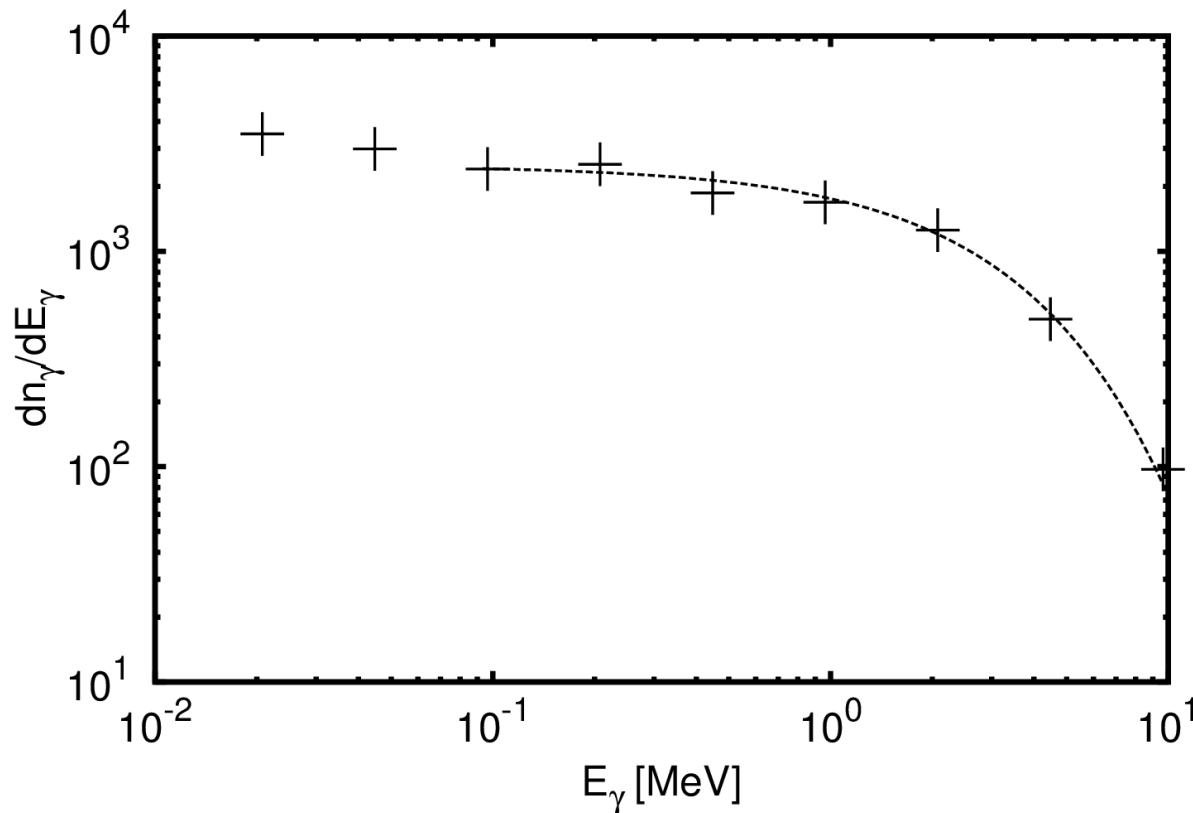
[C. Köhn et al., 2014. The importance of electron-electron Bremsstrahlung for terrestrial gamma-ray flashes, electron beams and electron-positron beams. *J. Phys. D.: Appl. Phys.*, vol. 47, 252001]

From photons to other species

Starting with a spectrum

$$\frac{dn_\gamma}{dE_\gamma} \sim e^{-E_\gamma/3\text{ MeV}}$$

with energies ≤ 40 MeV from
stepped lightning leader



Motion of photons

3D Monte Carlo code

Free motion between collisions

Motion of photons

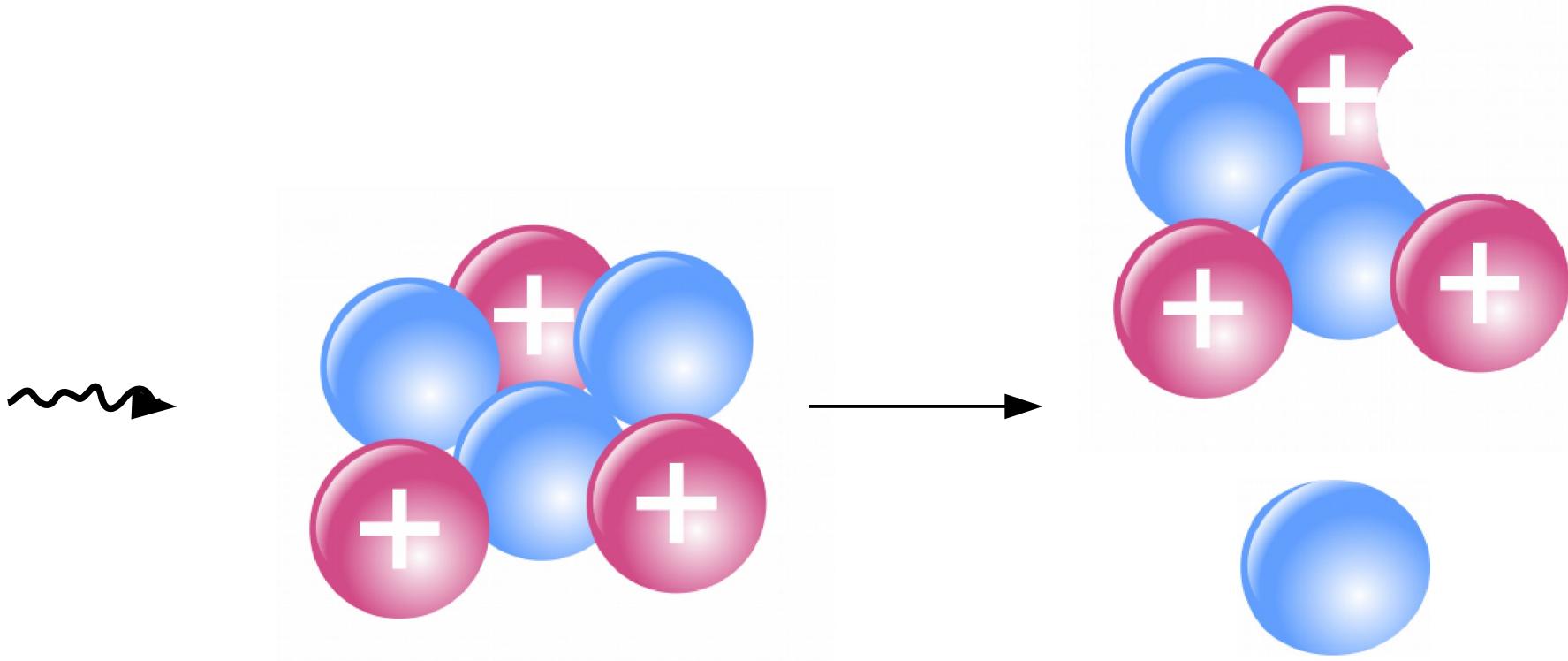
3D Monte Carlo code

Free motion between collisions

Differential cross sections to model collisions:

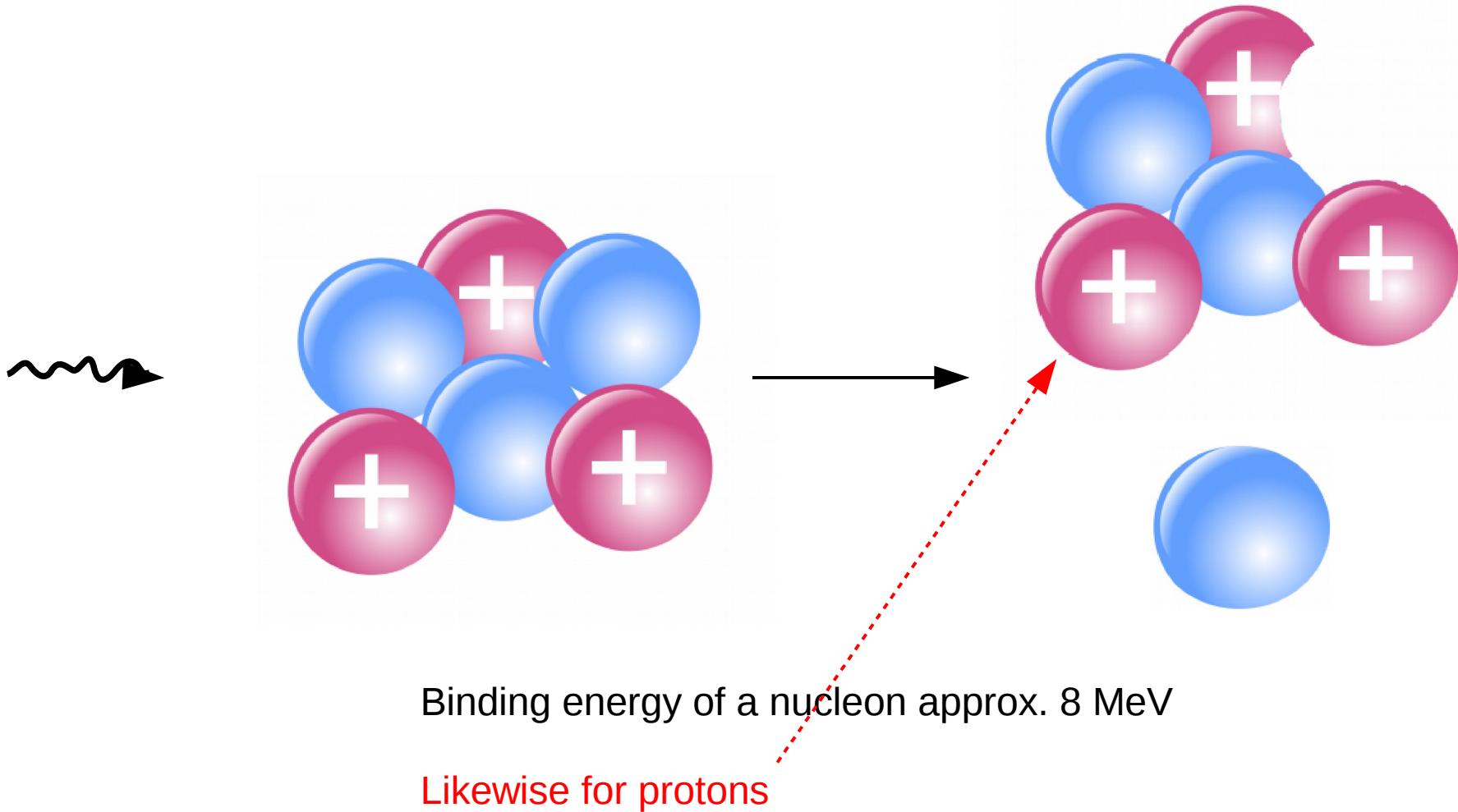
- Photo ionization (**Photons disappear**)
- Compton scattering
- Pair production (**Photons disappear**)
- Neutron production (**Photons disappear**)
- Proton production (**Photons disappear**)

Photoproduction of neutrons

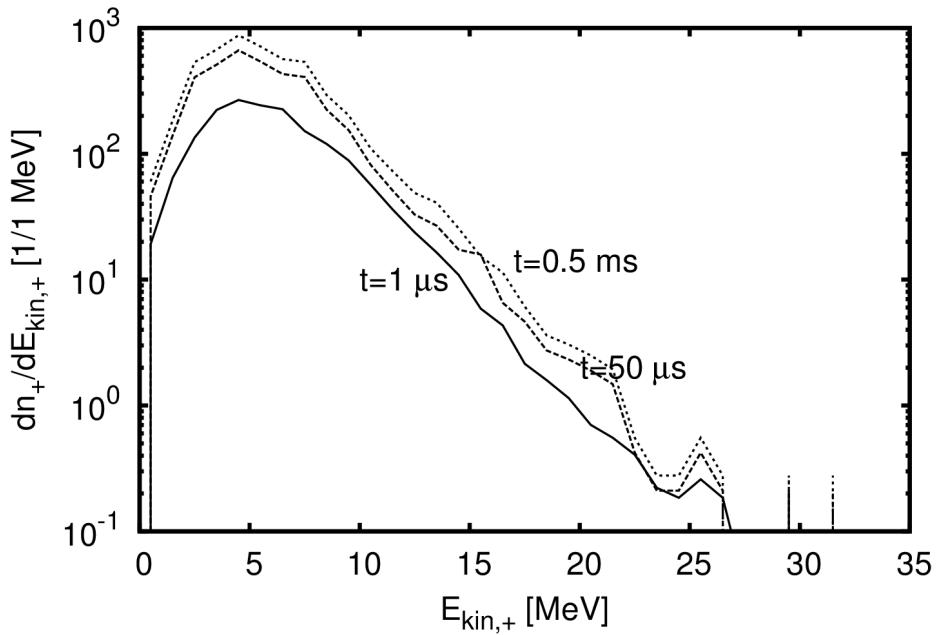


Binding energy of a nucleon approx. 8 MeV

Photoproduction of neutrons



Generation and propagation of positrons

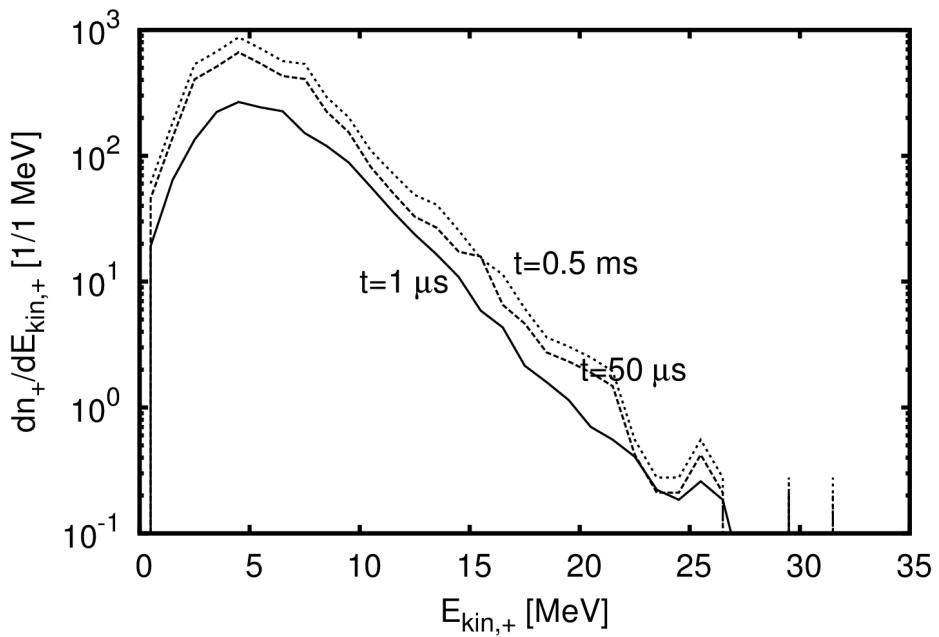


Maximum of
energy distribution
at approx 7 MeV

Energies of up to 30 MeV

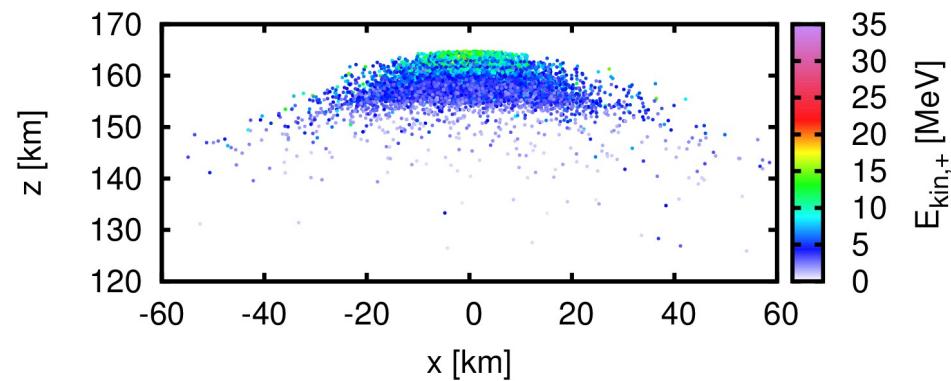
[C. Köhn and U. Ebert., 2015. Calculation of beams of positrons, neutrons and protons associated with terrestrial gamma-ray flashes. *J. Geophys. Res.*, vol. 120, doi:10.1002/2014JD022229]

Generation and propagation of positrons



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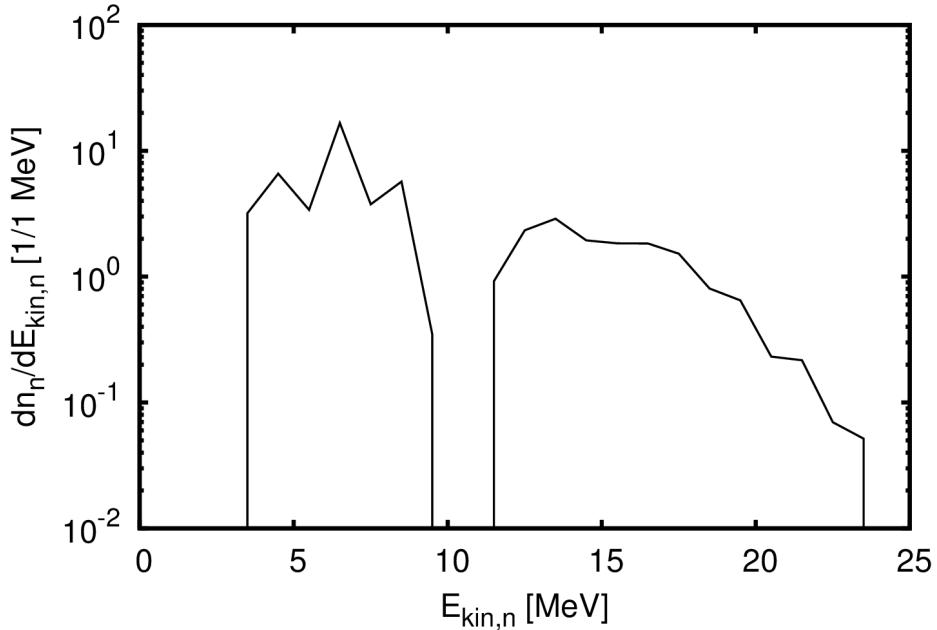
Energies of up to 30 MeV



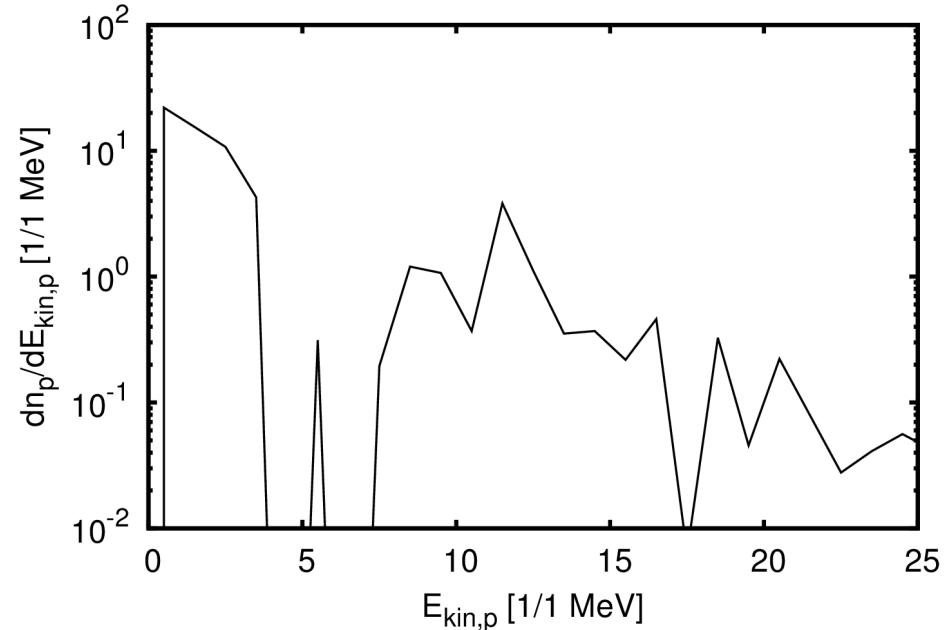
Positron beam limited
by relativistic positrons

Hadron energy distributions

Neutrons:



Protons:



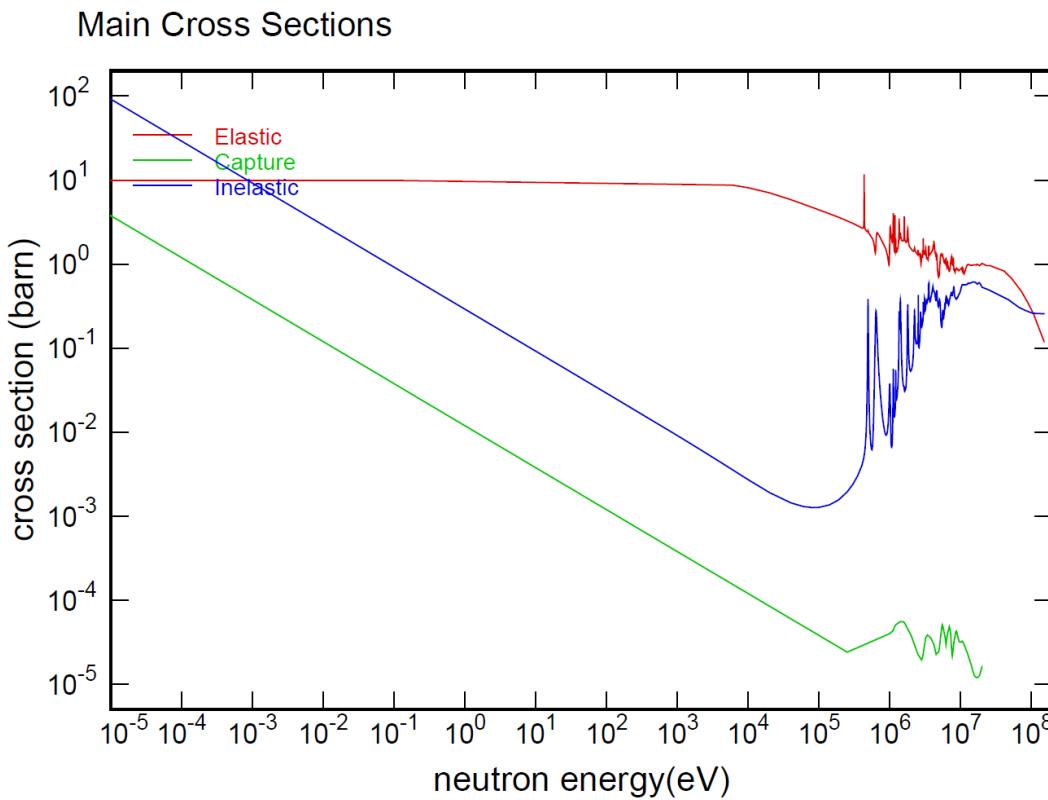
Energies of up to tens of MeV
Resonances because of cross sections

Propagation of hadrons through the atmosphere?

[C. Köhn and U. Ebert., 2015. Calculation of beams of positrons, neutrons and protons associated with terrestrial gamma-ray flashes. J. Geophys. Res., vol. 120, doi:10.1002/2014JD022229]

Neutron cross sections

for nitrogen

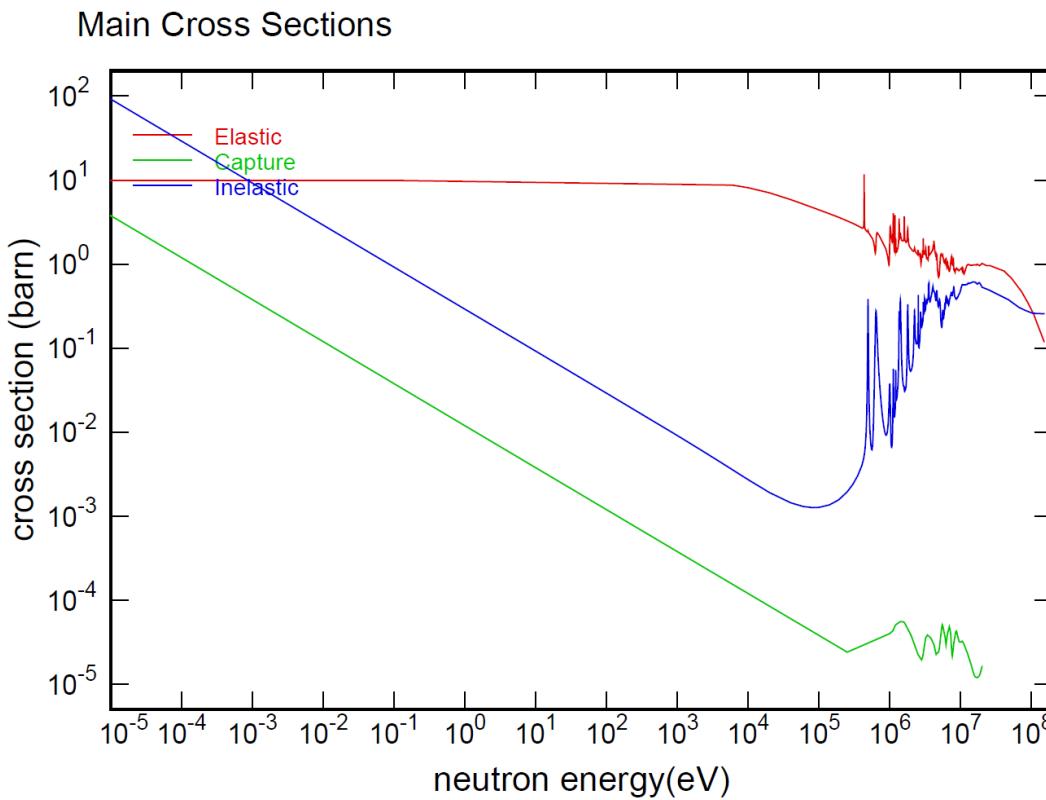


Elastic scattering:

- Energy transfer to air nuclei (recoil)
- Change of direction

Neutron cross sections

for nitrogen



Elastic scattering:

- Energy transfer to air nuclei (recoil)
- Change of direction

Inelastic scattering:

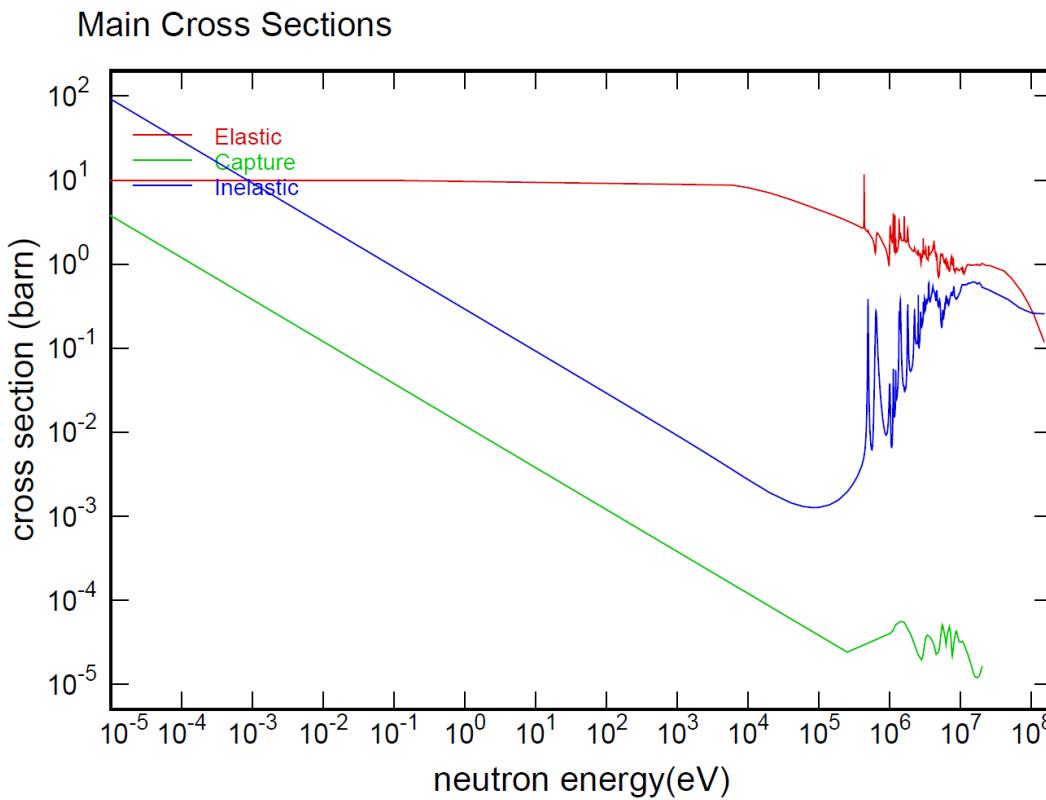
- Energy transfer to nuclei (excitations) in MeV range
- Change of direction
- Loss of neutrons, e.g.



- Changes constitution of atmosphere

Neutron cross sections

for nitrogen



Elastic scattering:

- Energy transfer to air nuclei (recoil)
- Change of direction

Inelastic scattering:

- Energy transfer to nuclei (excitations) in MeV range
- Change of direction
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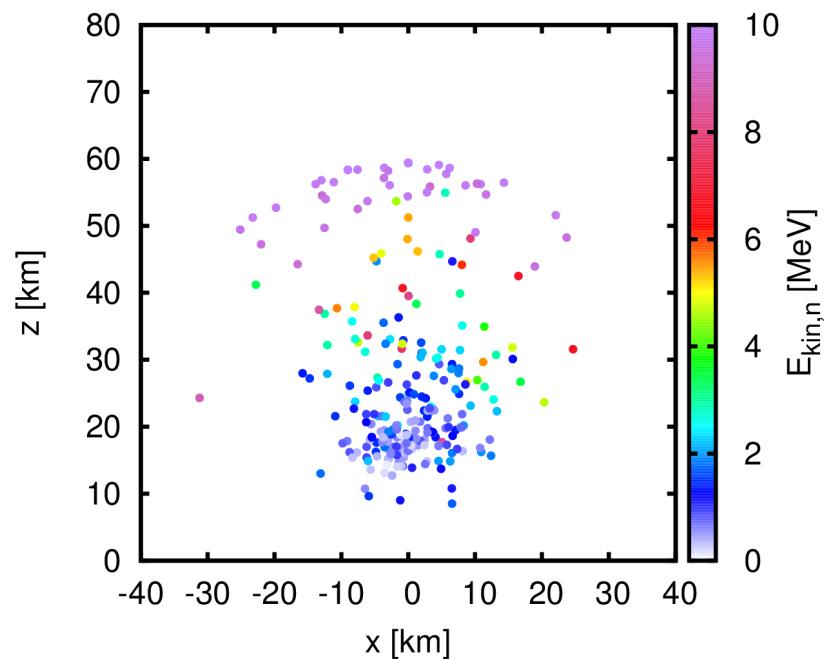
- Changes constitution of atmosphere

Capture:

- Loss of neutrons

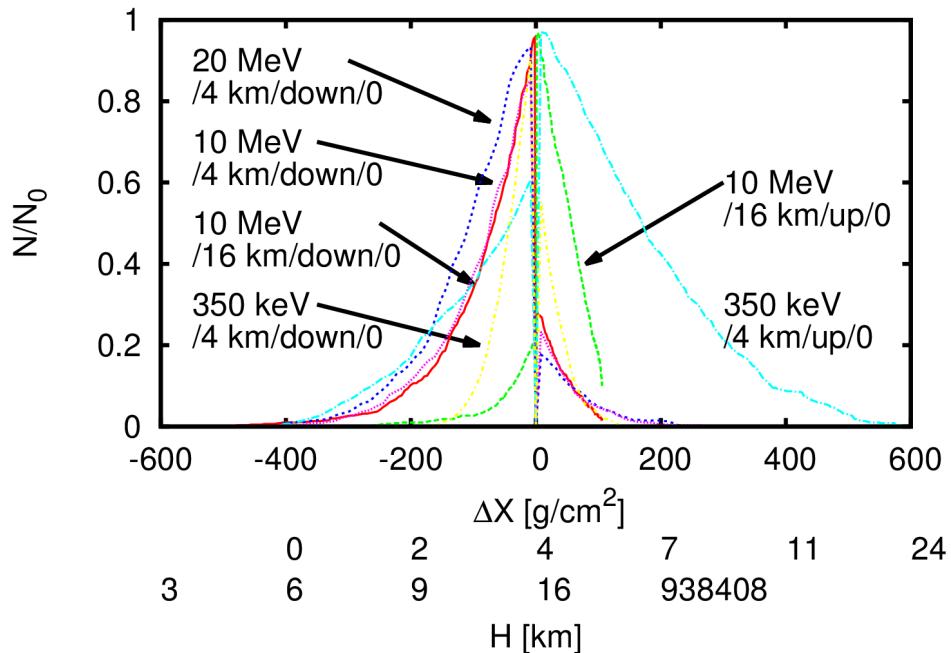
Temporal evolution of neutron beam

Beam of monoenergetic test neutrons



Neutron characteristics

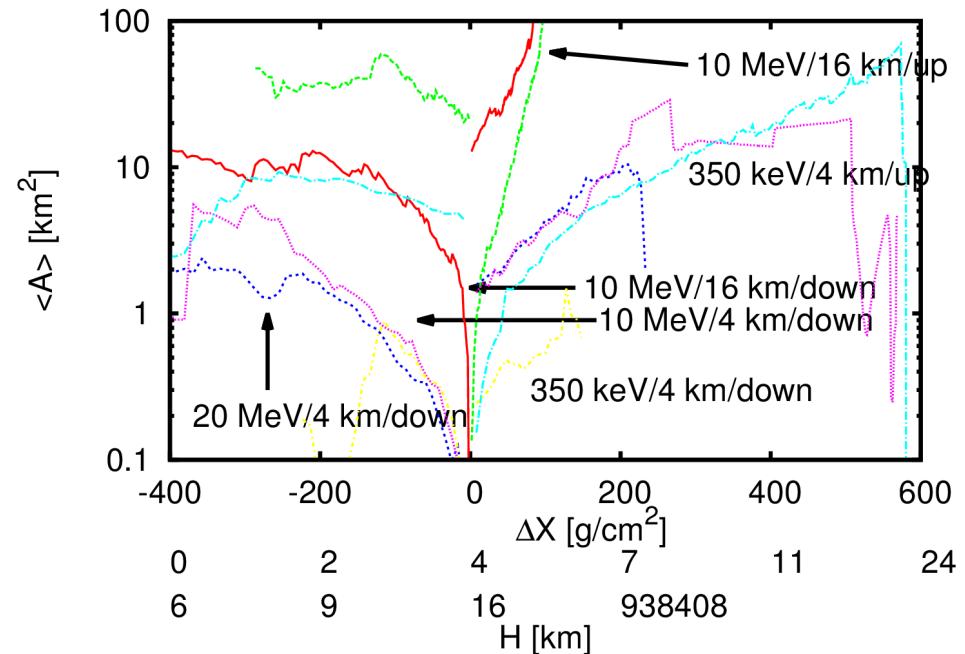
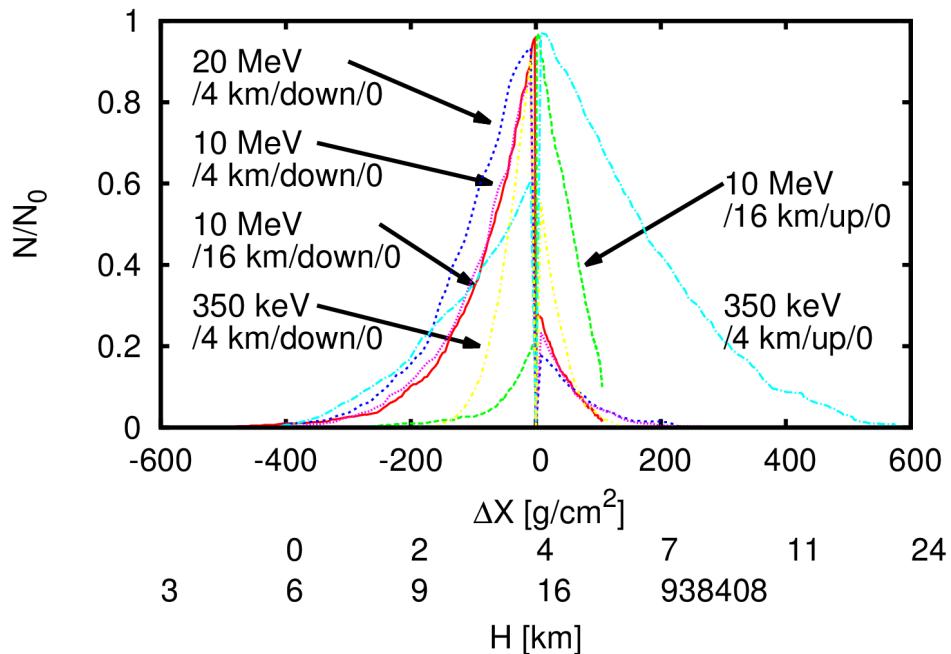
Start with different monoenergetic neutron beams at different altitudes



- Small probability of detecting neutrons at ground for all simulated cases
- For upward beams satellite detection possible

Neutron characteristics

Start with different monoenergetic neutron beams at different altitudes



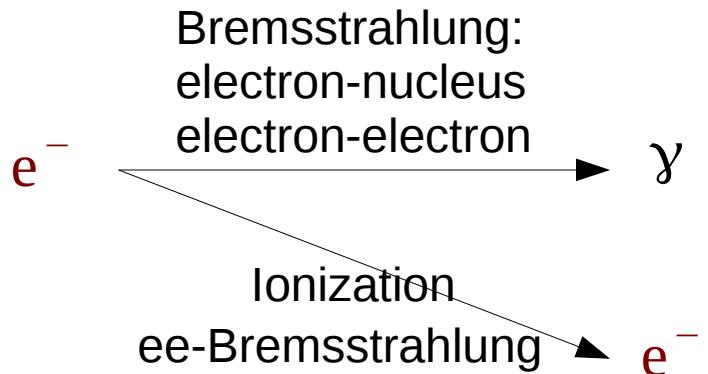
- Small probability of detecting neutrons at ground for all simulated cases
- For upward beams satellite detection possible

However: Mean cross sections of areas of tens of km^2

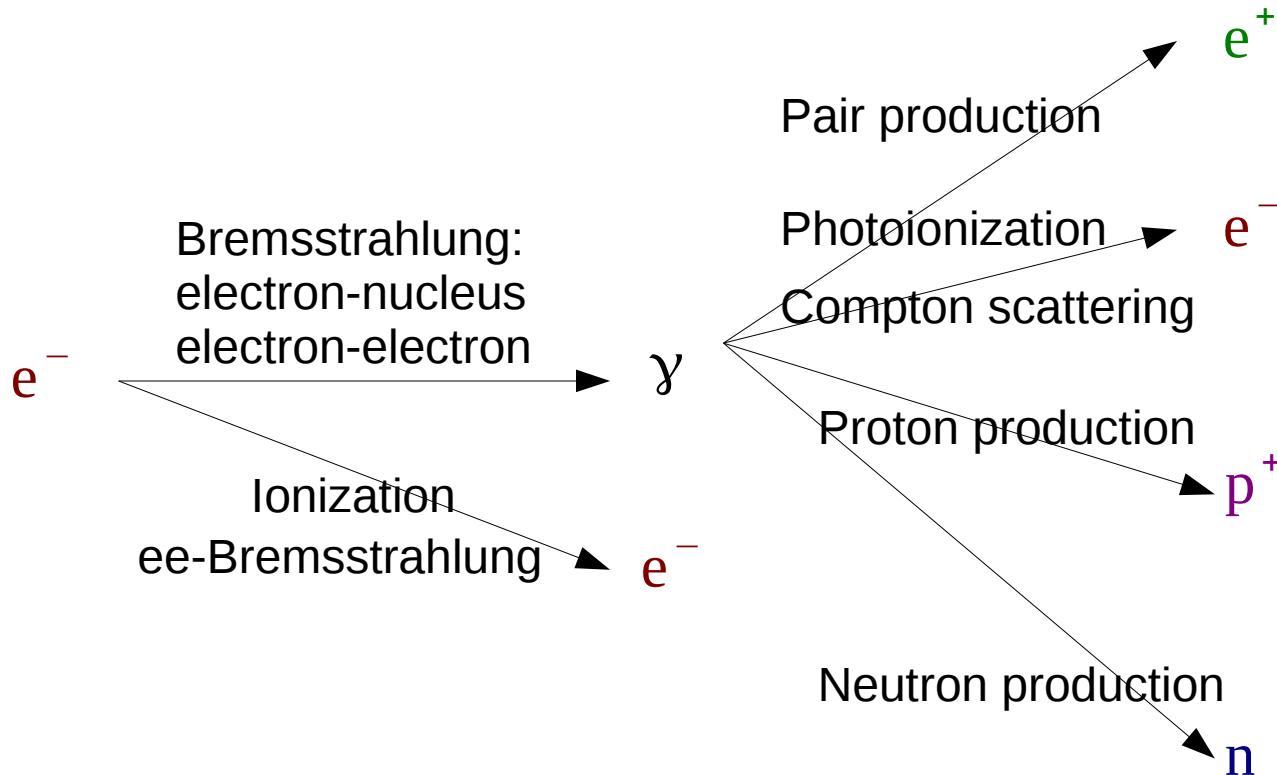
Towards a full feedback model

e⁻

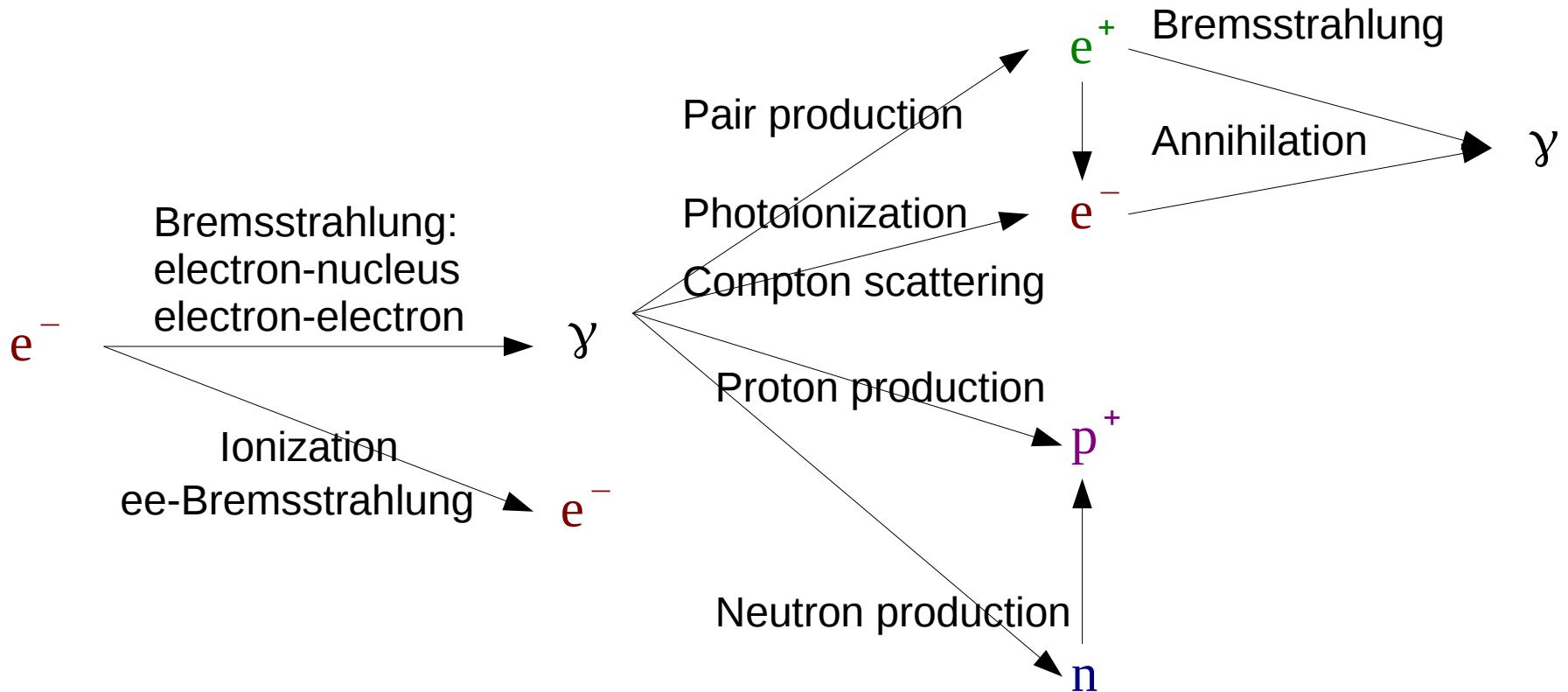
Conversion from electrons to different species



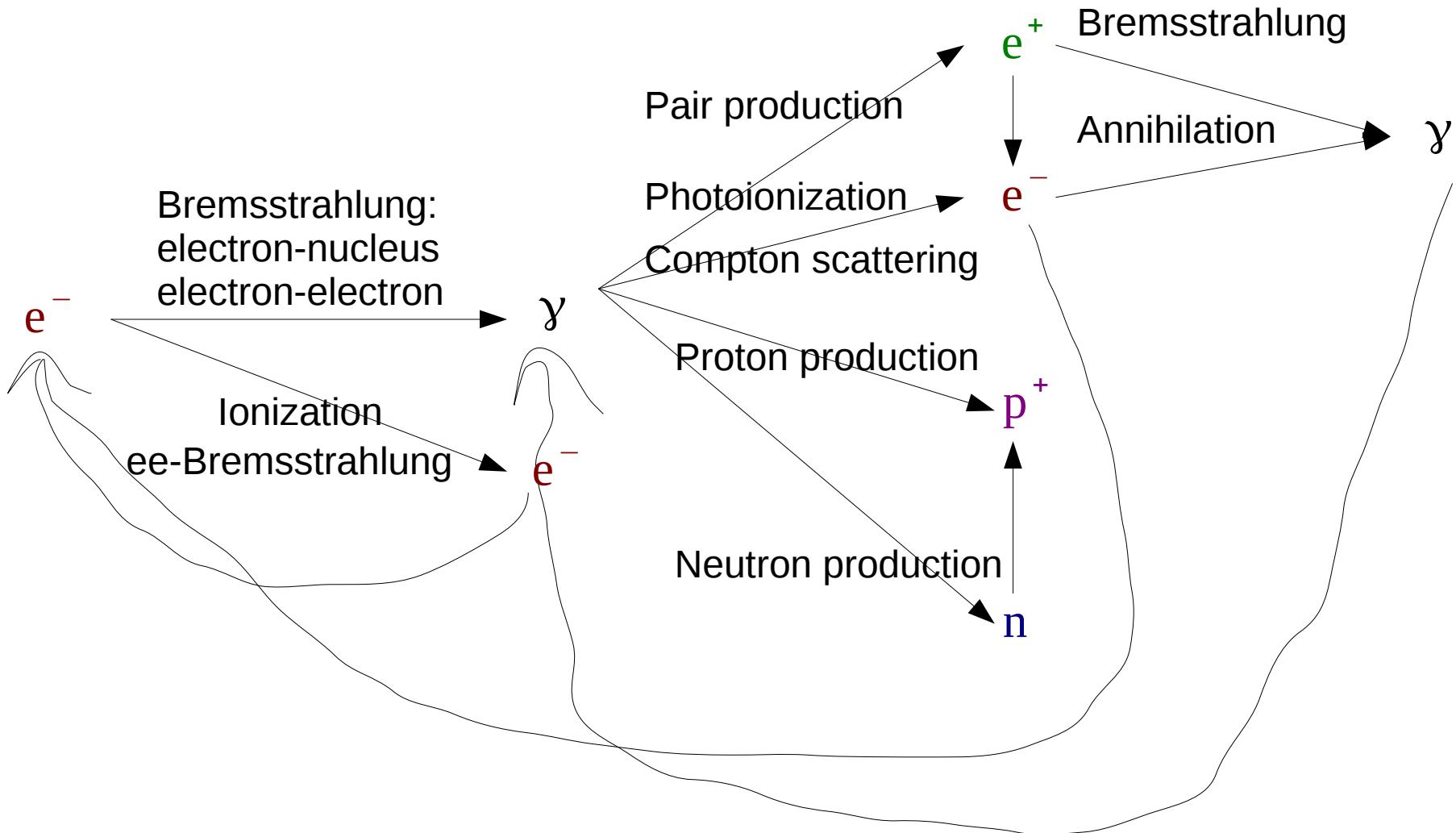
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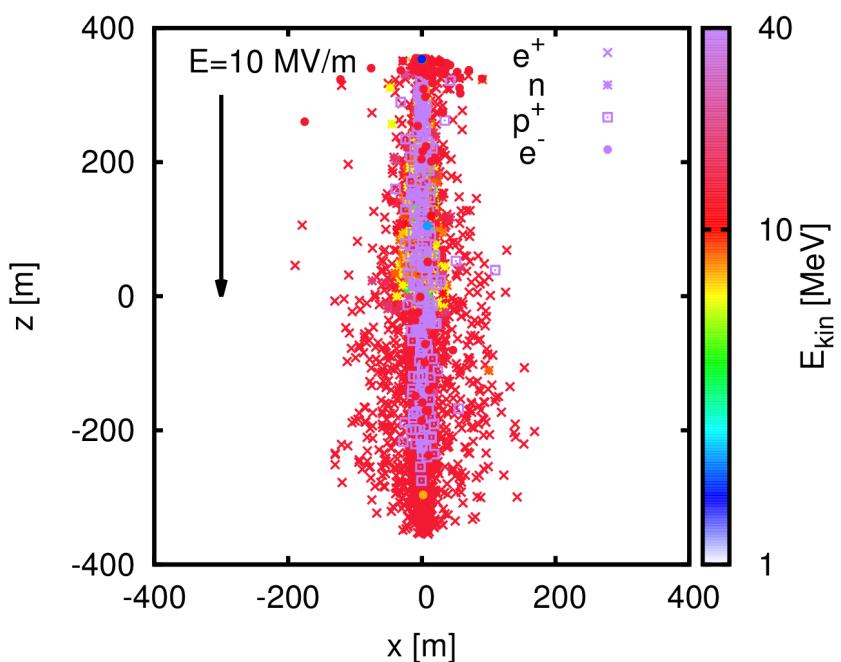


Example: Motion in a constant field

Start with upwards
photon beam $\frac{dn_\gamma}{dE_\gamma} \sim e^{-E_\gamma/3\text{ MeV}}$

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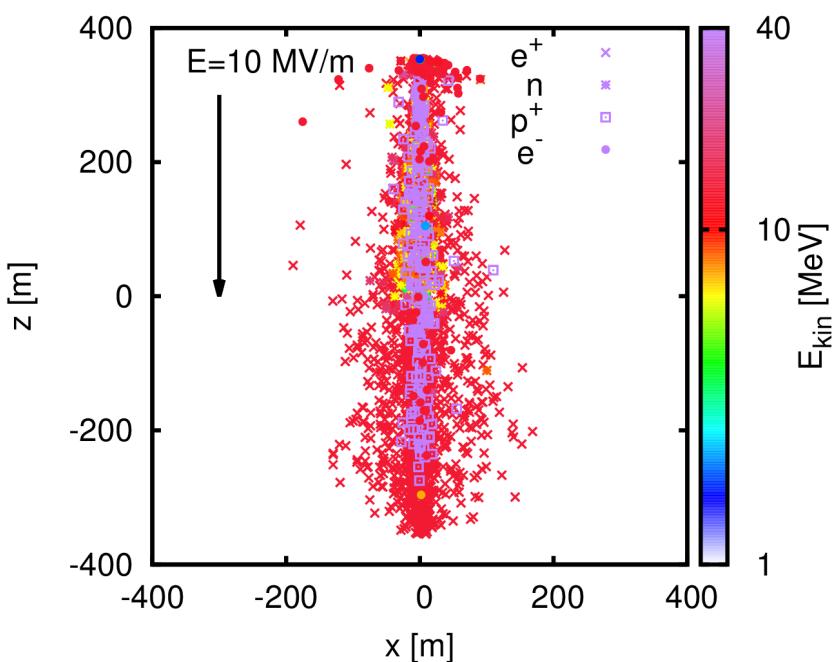


i) photon produce
leptons and hadrons
along their path

Position and energy of
different species after 1 μs

Example: Motion in a constant field

Start with upwards photon beam $\frac{dn_\gamma}{dE_\gamma} \sim e^{-E_\gamma/3\text{ MeV}}$

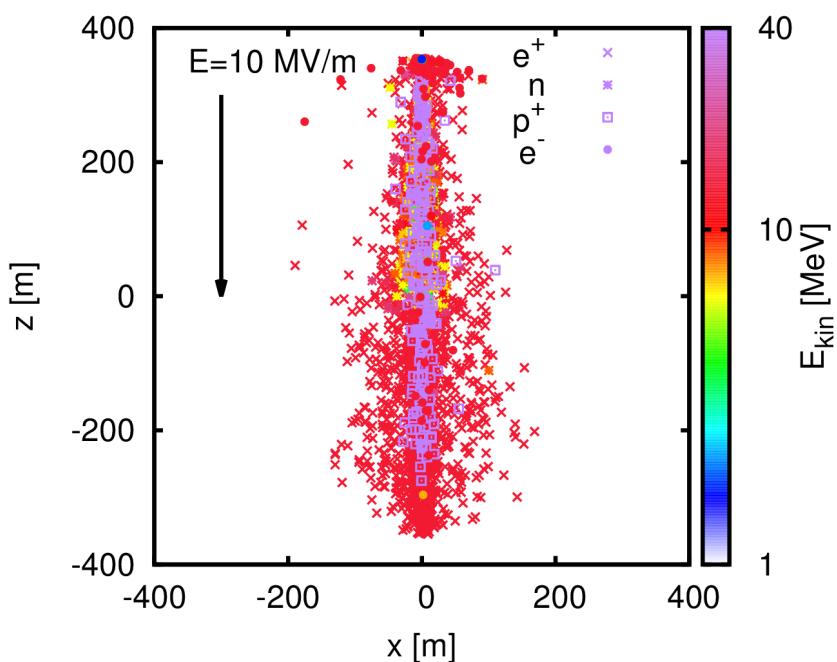


- i) photon produce leptons and hadrons along their path
- ii) electrons move against the electric field

Position and energy of different species after 1 μ s

Example: Motion in a constant field

Start with upwards photon beam $\frac{dn_\gamma}{dE_\gamma} \sim e^{-E_\gamma/3\text{ MeV}}$

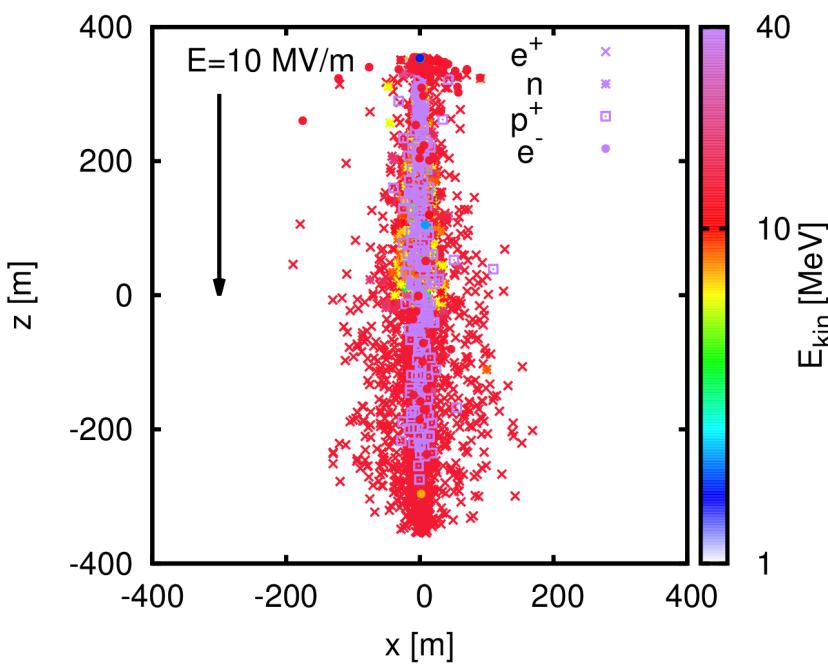


- i) photon produce leptons and hadrons along their path
- ii) electrons move against the electric field
- iii) positrons in the opposite direction

Position and energy of different species after 1 μ s

Example: Motion in a constant field

Start with upwards photon beam $\frac{dn_\gamma}{dE_\gamma} \sim e^{-E_\gamma/3\text{ MeV}}$



Position and energy of
different species after 1 μs

- i) photon produce leptons and hadrons along their path
- ii) electrons move against the electric field
- iii) positrons in the opposite direction
- iv) protons and neutrons much more slower than leptons

Conclusion

