

N=2* Holography and Localization

Konstantin Zarembo
(Nordita, Stockholm)

X. Chen-Lin, A. Dekel, K.Z. 1512.06420

X. Chen-Lin, K.Z. 1502.01942

K.Z. 1410.6114

X. Chen-Lin, J. Gordon, K.Z. 1408.6040

A. Buchel, J. Russo, K.Z. 1301.1597

Current Themes in Holography, Copenhagen, 29.04.16

N=2* theory

vector multiplet

A_μ

ψ

Φ, Φ'

mass = 0

hypermultiplets

Z

χ

$\tilde{\chi}$

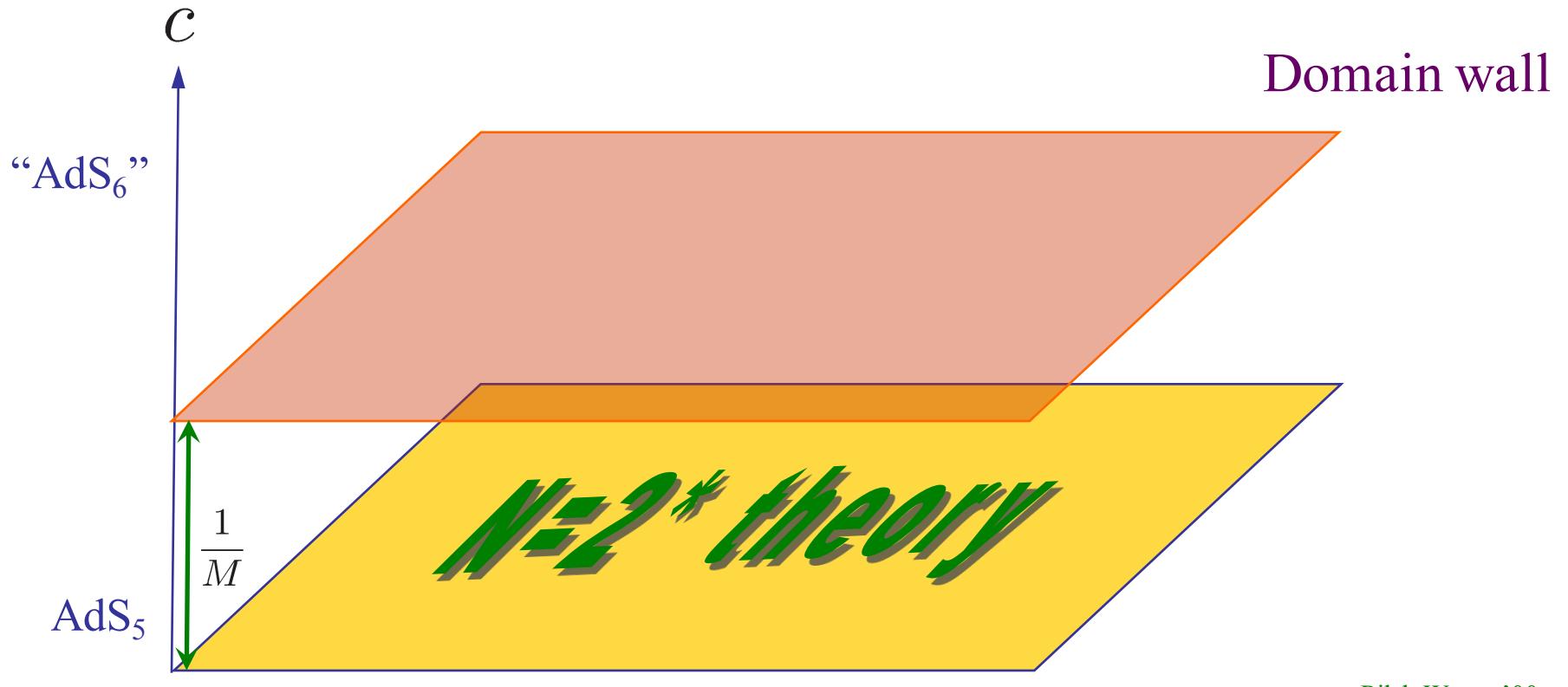
+ conj.

\tilde{Z}

mass = $\pm M$

- relevant perturbation of N=4 super-Yang-Mills

Holographic dual



Pilch, Warner'00

$$ds^2 = \frac{A}{c^2 - 1} M^2 dx_\mu^2 + \frac{1}{A (c^2 - 1)^2} dc^2$$

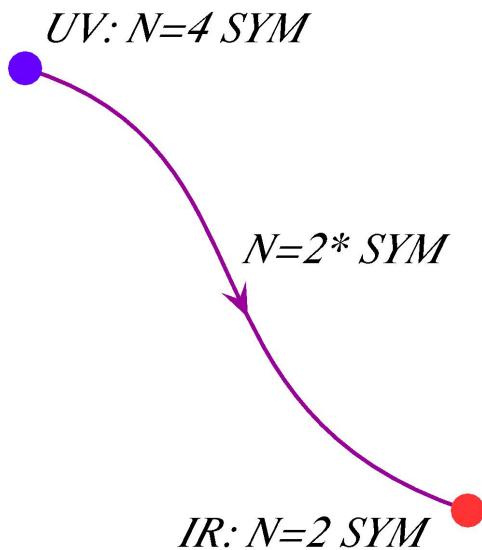
$$A = c + \frac{c^2 - 1}{2} \ln \frac{c - 1}{c + 1}$$

$$c = 1 + \frac{M^2 z^2}{2} + \dots$$

- UV regularization of pure N=2 SYM

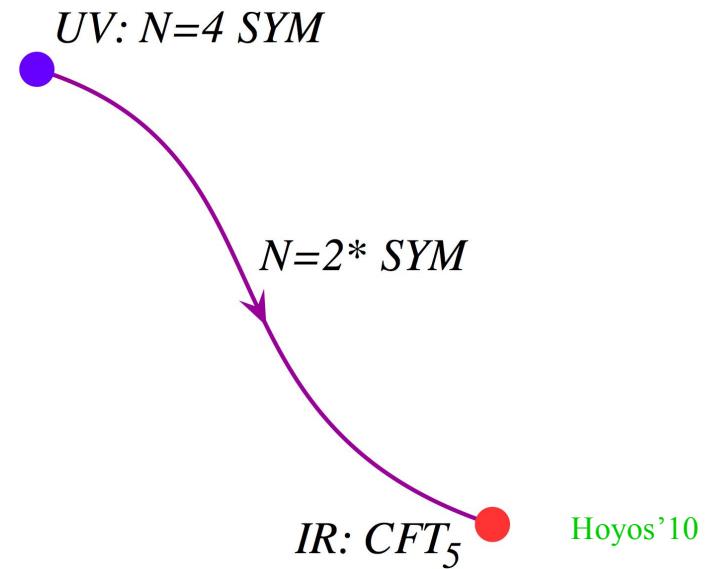
In this talk:
 $N=\infty$
 $\lambda = g_{\text{YM}}^2 N$

Weak coupling



$$\Lambda = M e^{-\frac{4\pi^2}{\lambda}}$$

Strong coupling

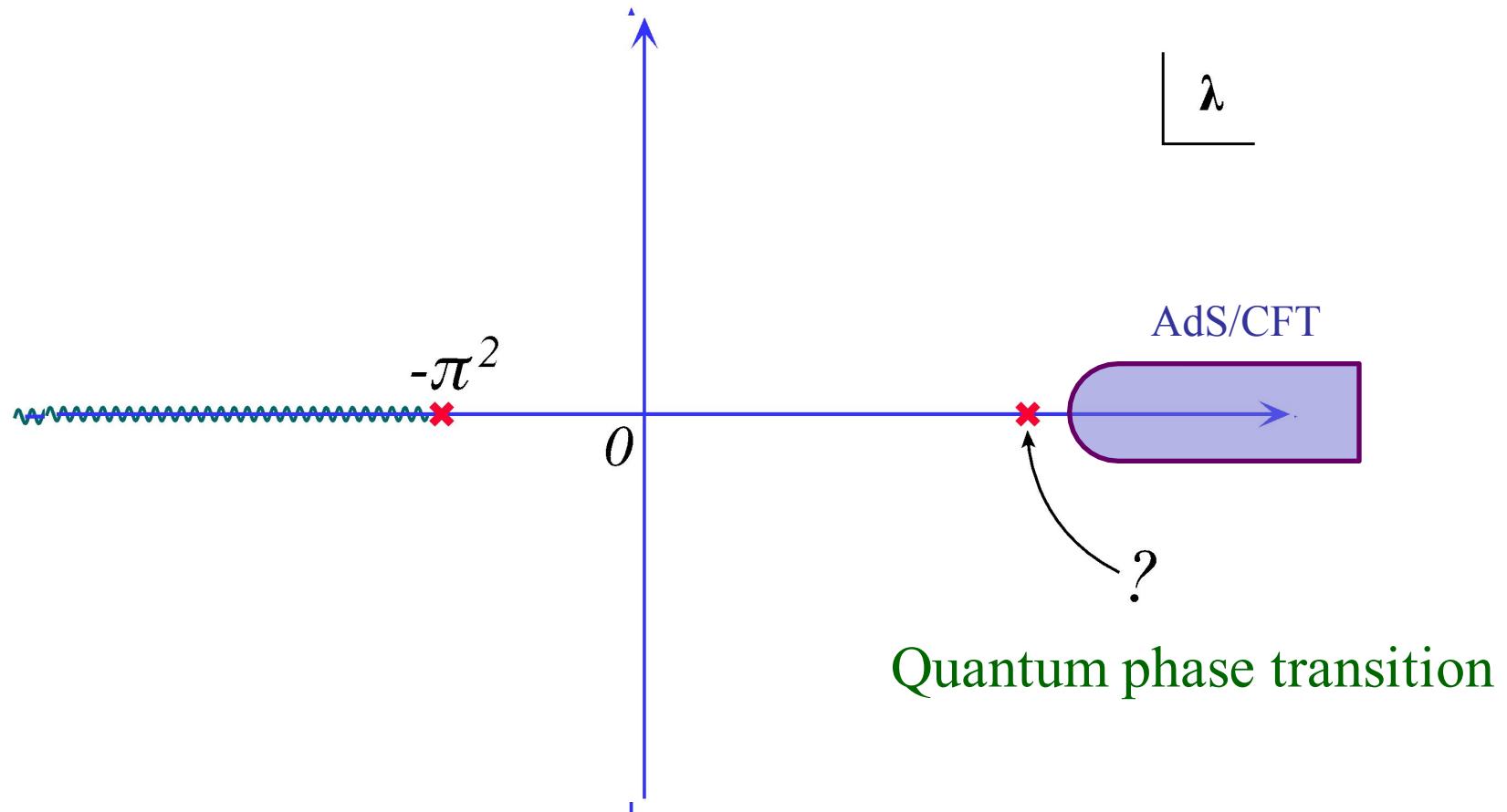


Dimensional crossover
by Eguchi-Kawai mechanism

Young, Z.'14

Phase transitions

In N=4 super-Yang-Mills:



Master field

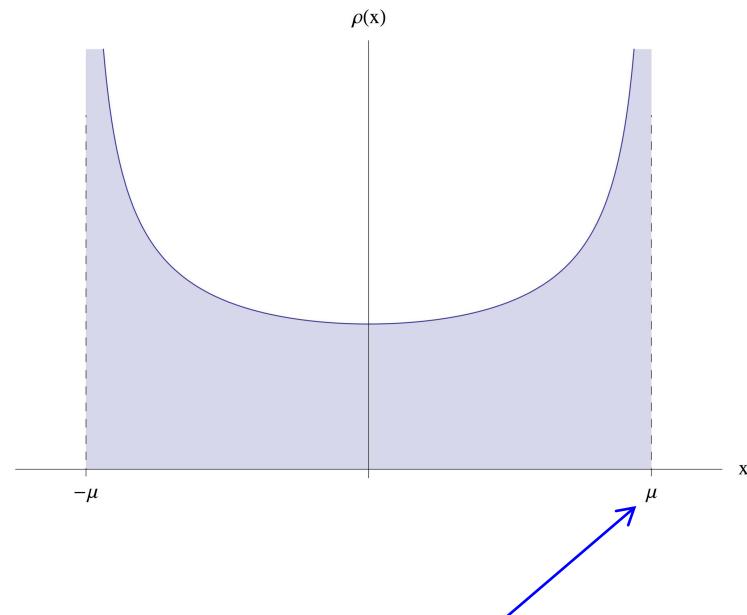
Symmetry breaking: $SU(N) \rightarrow U(1)^{N-1}$

$$\langle \Phi \rangle = \text{diag} (a_1, \dots, a_N)$$

Large-N master field:

$$\rho(x) = \left\langle \frac{1}{N} \text{tr} \delta (x - \Phi) \right\rangle$$

$$= \frac{1}{N} \sum_{i=1}^N \delta (x - a_i)$$



typical scale of symmetry breaking

Exact solution

Localization

Pestun'07

Russo,Z.'13

SW theory

Seiberg,Witten'94

Billo,Frau,Fucito,Lerda,Morales,Poghossian,Ricci Pacifici'14
Hollowood,Prem Kumar'15

$$\int_{-\mu}^{\mu} dy \rho(y) \left(\frac{2}{x-y} - \frac{1}{x-y+M} - \frac{1}{x-y-M} \right) = 0$$

$$\int_{-\mu}^{\mu} dx \rho(x) \ln \frac{|M^2 - x^2|}{x^2} = \frac{8\pi^2}{\lambda}$$

$$\int_{-\mu}^{\mu} dx \rho(x) = 1$$

RG log

Weak Coupling

$$\lambda \ll 1 \quad \xrightarrow{\hspace{1cm}} \quad \mu \ll M$$

$$\int_{-\mu}^{\mu} dy \rho(y) \left(\frac{2}{x-y} - \frac{1}{x-y+M} - \frac{1}{x-y-M} \right) = 0$$

$$\rho(x) = \frac{1}{\pi \sqrt{\mu^2 - x^2}} \quad \text{Douglas,Shenker'95}$$

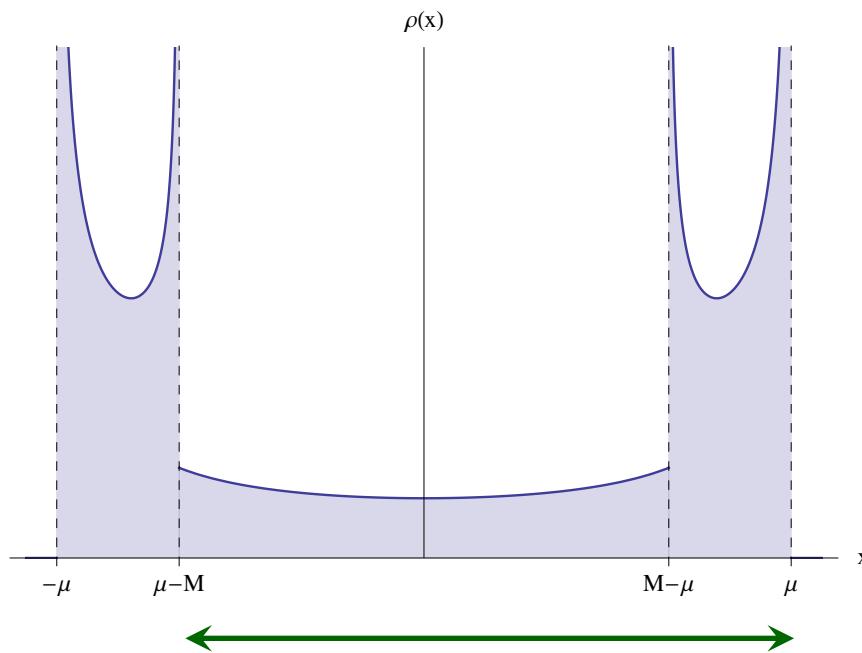
$$\frac{\mu}{M} = 2 e^{-\frac{4\pi^2}{\lambda}} - 4 e^{-\frac{12\pi^2}{\lambda}} - 16 e^{-\frac{20\pi^2}{\lambda}} - 58 e^{-\frac{28\pi^2}{\lambda}} - 324 e^{-\frac{36\pi^2}{\lambda}} - 1856 e^{-\frac{44\pi^2}{\lambda}} + \dots$$

OPE in condensates

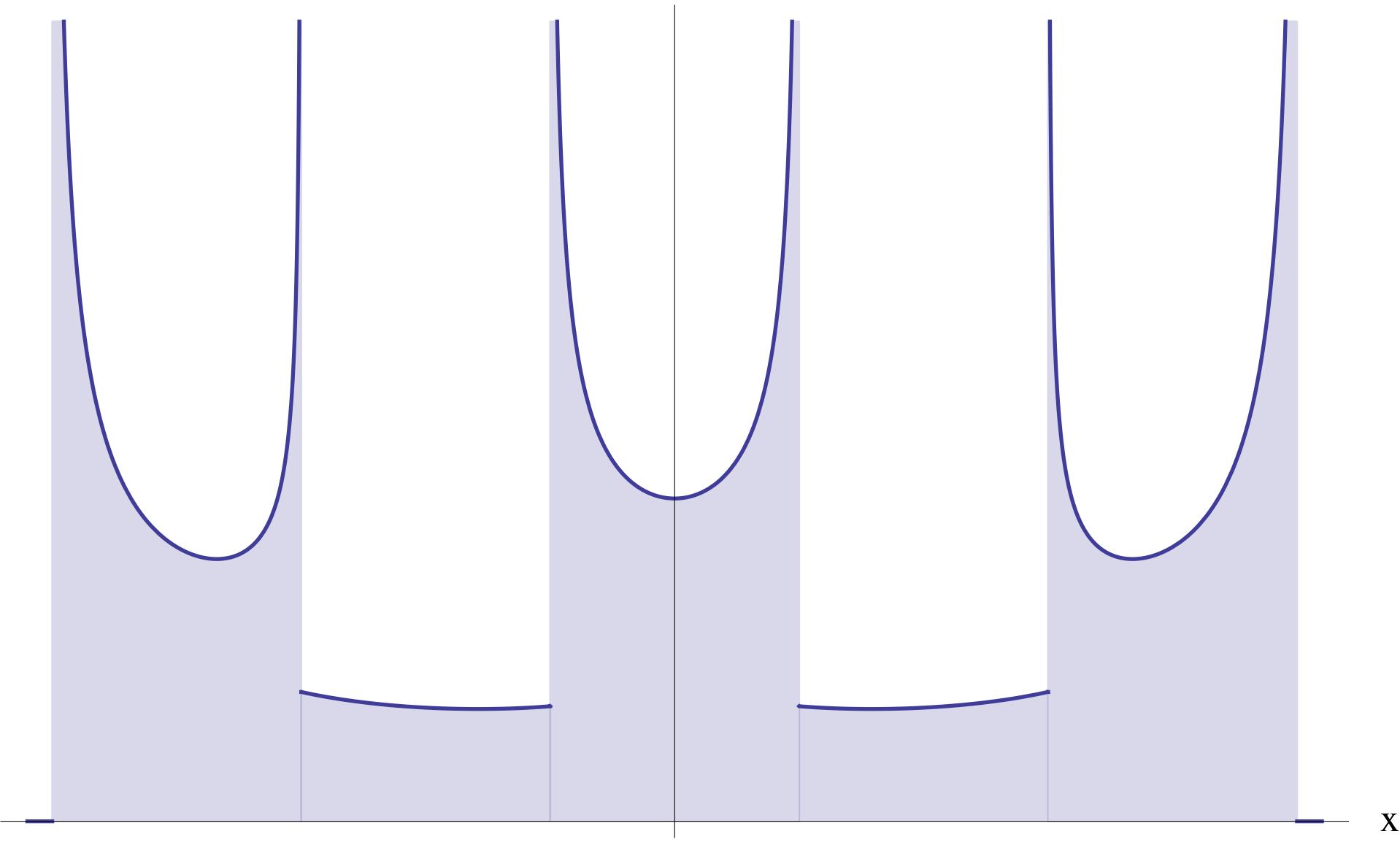
β-function of N=2 SYM!

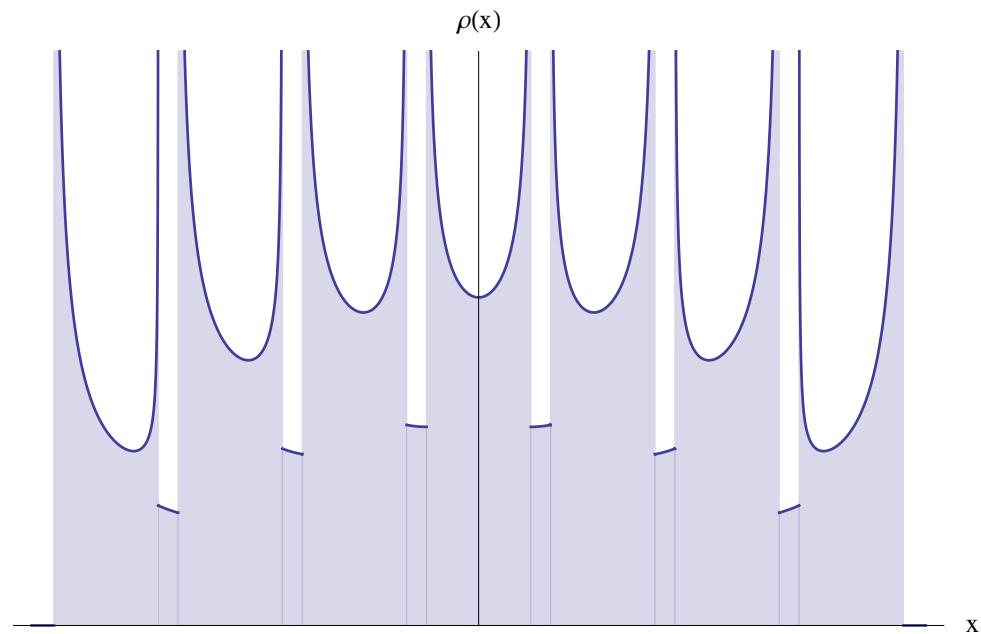
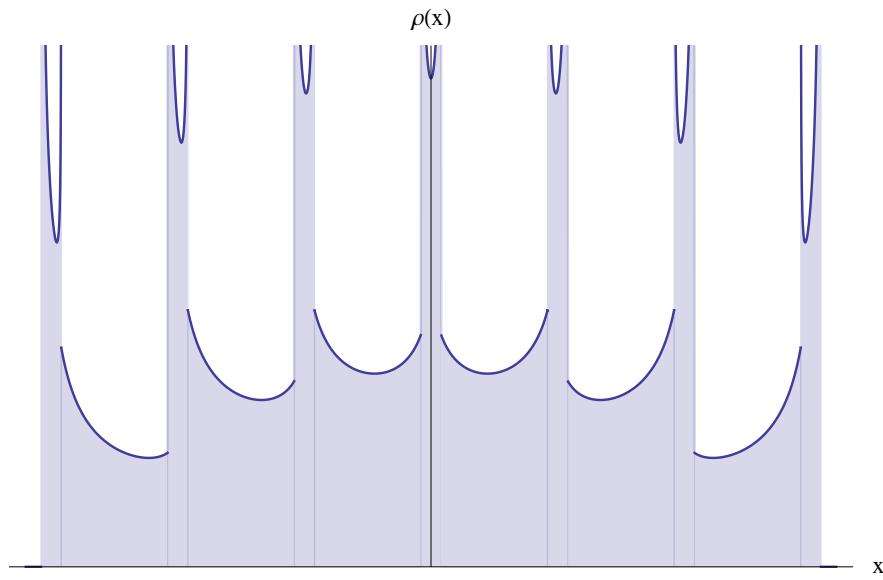
Phase transition

Weak-coupling solution is valid up to $\mu_c = \frac{M}{2} \iff \lambda < \lambda_c = 35.42\dots$

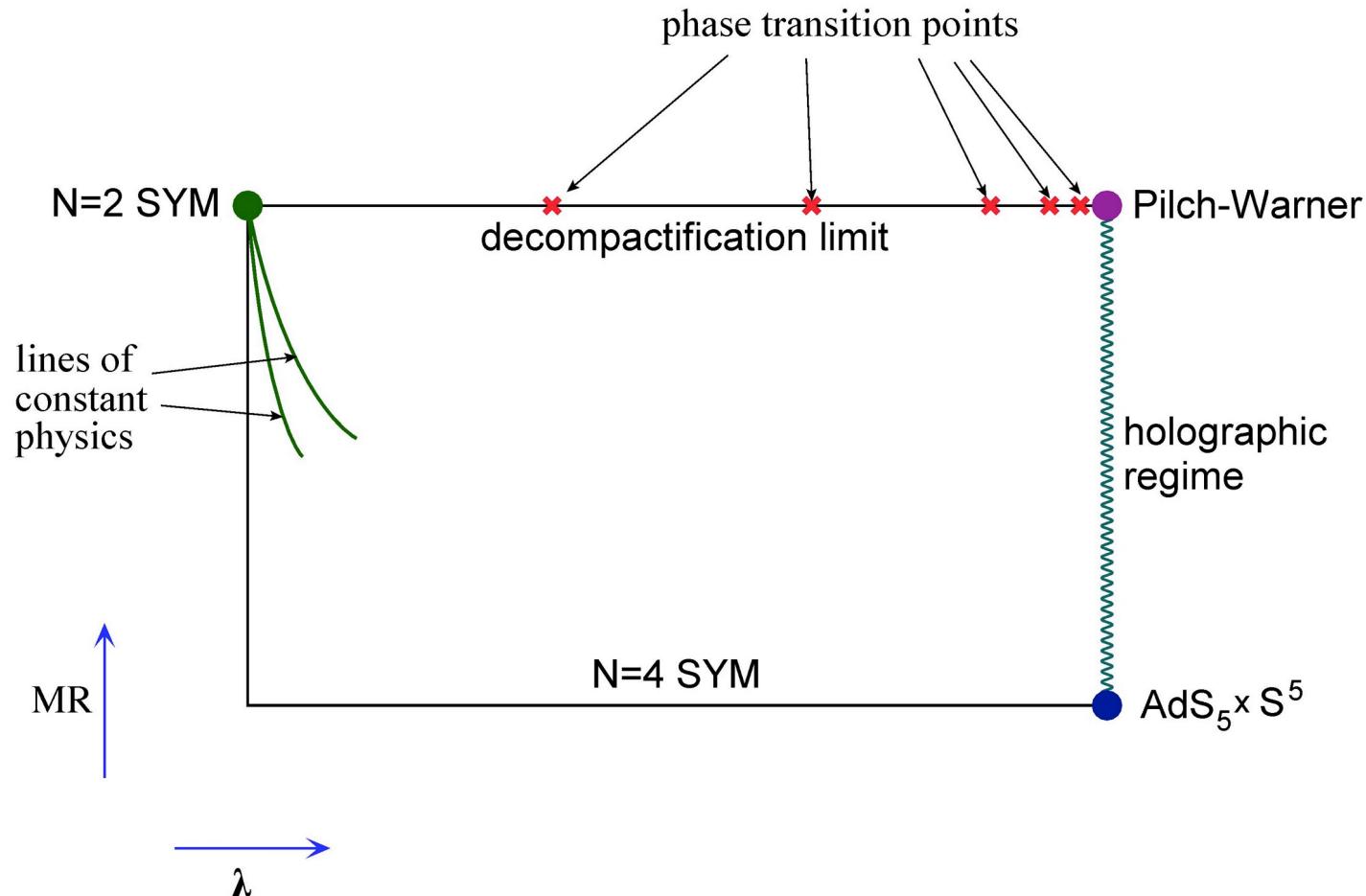


resonance on massless hyper: $m_{ij}^h = |a_i - a_j \pm M|$

$\rho(x)$ 

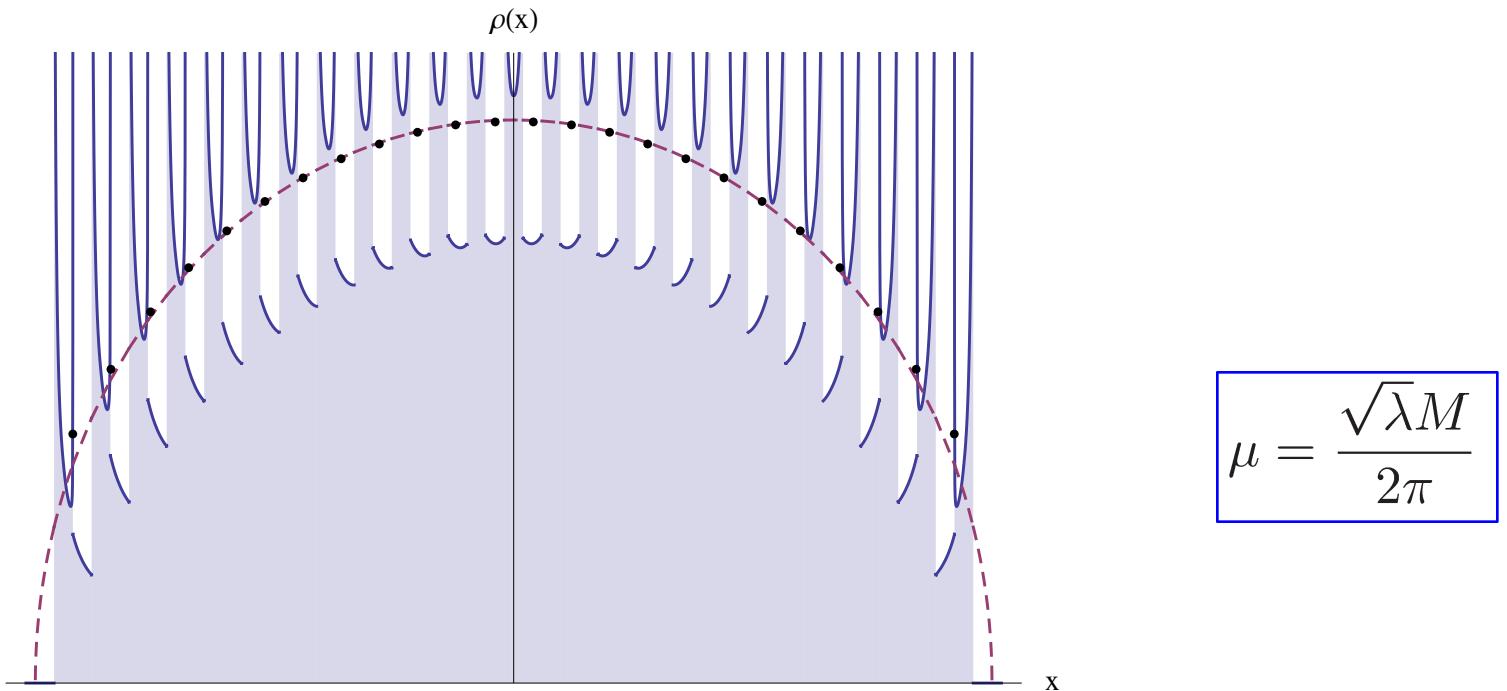


Phase diagram



R: radius of S^4

Strong coupling



$$\rho(x) \simeq \frac{2}{\pi\mu^2} \sqrt{\mu^2 - x^2} + \frac{1}{\pi} \sqrt{\frac{M}{2\mu^5}} \left[(\mu - x) \zeta \left(\frac{1}{2}, \left\{ \frac{\mu + x}{M} \right\} \right) + (\mu + x) \zeta \left(\frac{1}{2}, \left\{ \frac{\mu - x}{M} \right\} \right) \right],$$

$\mathcal{O}\left(\frac{1}{\mu}\right)$ ←
 $\mathcal{O}\left(\frac{M^{1/2}}{\mu^{3/2}}\right) = \mathcal{O}\left(\frac{\lambda^{-1/4}}{\mu}\right)$ ←

Perimeter law

$$W(C) = \left\langle \frac{1}{N} \operatorname{tr} P \exp \oint_C ds (i\dot{x}^\mu A_\mu + |\dot{x}| \Phi) \right\rangle$$

 substitute $\langle \Phi \rangle$

$$W(C) = \frac{1}{N} \sum_i e^{La_i} \xrightarrow{L \rightarrow \infty} e^{L\mu} \quad (\text{perimeter law})$$

$$\ln W(C) \simeq \frac{\sqrt{\lambda}ML}{2\pi} \left(1 - \frac{\pi}{\sqrt{\lambda}} + \dots \right)$$

Buchel,Russo,Z.'13

Chen-Lin,Gordon,Z.'14

Z.'14

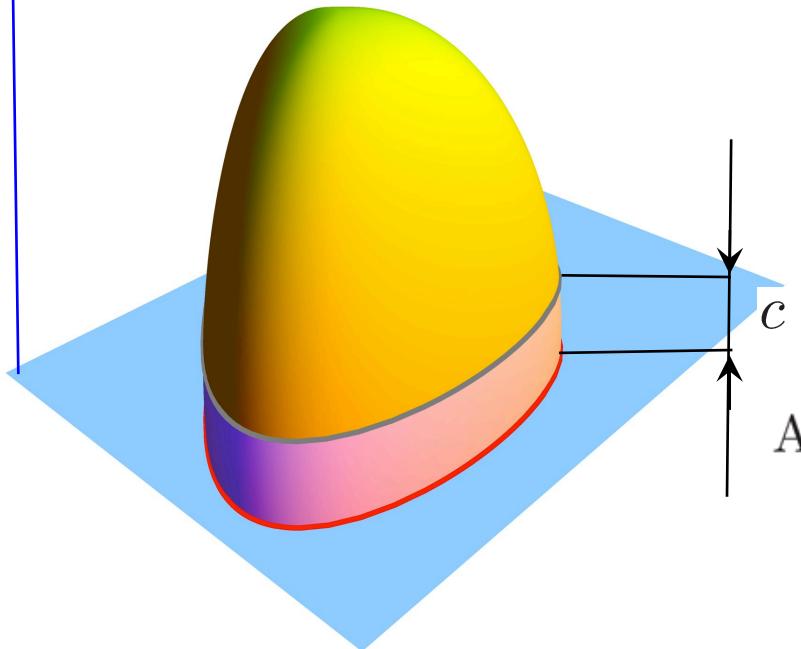
c

Minimal area

Metric of the gravity dual:

$$ds^2 = \frac{\rho^6}{c^2 - 1} M^2 dl^2 + \frac{1}{\rho^6 (c^2 - 1)^2} dc^2$$

Pilch, Warner '00



$c \sim 1$

$$\text{Area} = ML \int_{1+\frac{\varepsilon^2 M^2}{2}}^{\infty} \frac{dc}{(c^2 - 1)^{\frac{3}{2}}} = \frac{L}{\varepsilon} - ML$$



renormalized away

$$\ln W(C) \simeq \frac{\sqrt{\lambda}ML}{2\pi} \quad (\lambda \rightarrow \infty, ML \gg 1)$$

agrees!

- free energy also agrees

Bobev, Elvang, Freedman, Pufu '13

Higher representations

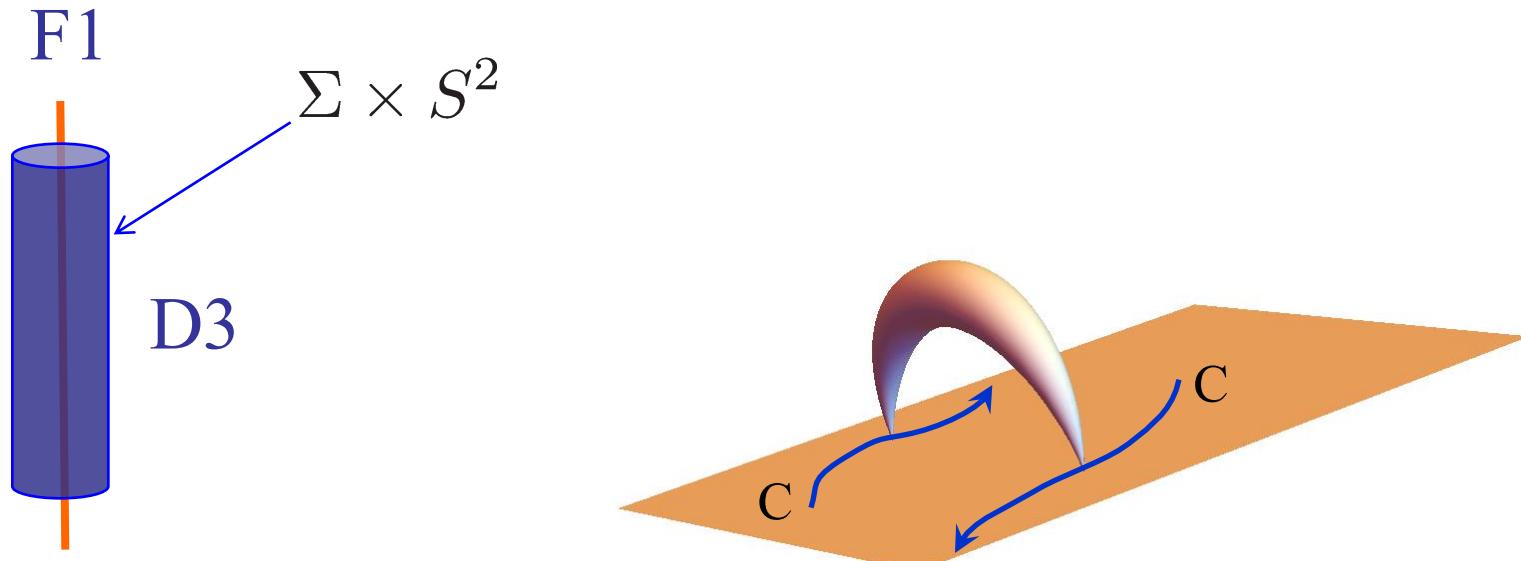
$$W_{\mathcal{R}}(C) = \left\langle \frac{1}{N} \operatorname{tr}_{\mathcal{R}} \text{P exp} \oint_C ds (i\dot{x}^\mu A_\mu + |\dot{x}| \Phi) \right\rangle$$

$$\mathcal{R}_k = \overbrace{\boxed{} \boxed{} \boxed{} \boxed{} \boxed{}}^k$$

$$W_{\mathcal{R}_k} \simeq e^{\frac{NM^2L^2}{2\pi^2} f\left(\frac{\pi\sqrt{\lambda}k}{2NML}\right)} \underset{ML \gg 1}{\simeq} e^{\frac{\sqrt{\lambda}kML}{2\pi}}$$

Chen-Lin,Z.'15

$$f(\kappa) = \kappa\sqrt{1+\kappa^2} + \operatorname{arcsinh} \kappa$$



D3-brane embedding (for straight Wilson line):

$$r = \frac{\sqrt{\lambda}k}{N} \sqrt{c^2 - 1}, \quad F = -\frac{i \text{Vol}(S^2)}{(c^2 - 1)^{\frac{3}{2}}}$$

$$W_{\mathcal{R}_k} = e^{-S_{D3}} = e^{\frac{\sqrt{\lambda}k M L}{2\pi}}$$
Chen-Lin,Dekel,Z.'15

agrees w. QFT

Conclusions

- Basic agreement between first-principles QFT calculations and **non-conformal** holography:
Strong coupling \longleftrightarrow Classical gravity
- Quantum corrections? $1/\sqrt{\lambda}$
- Strong-weak coupling phase transitions do not obstruct holography
- What is holographic interpretation of these phase transitions?