

REMEDIOS

- LHC precision tests (di/tri-boson): what is being searched for? -

$L=27\text{ km}$

$E_{\text{beam}}=362\text{ MJ}$

Precision=??



AlpTransit

$L=35\text{ km}$

$E_{\text{train}}=362\text{ MJ}$

Precision= 10^{-6}



Francesco Riva
(CERN)

In collaboration with
Liu, Pomarol, Rattazzi 1603.03064,
Azatov, Contino, Machado 1607.05236

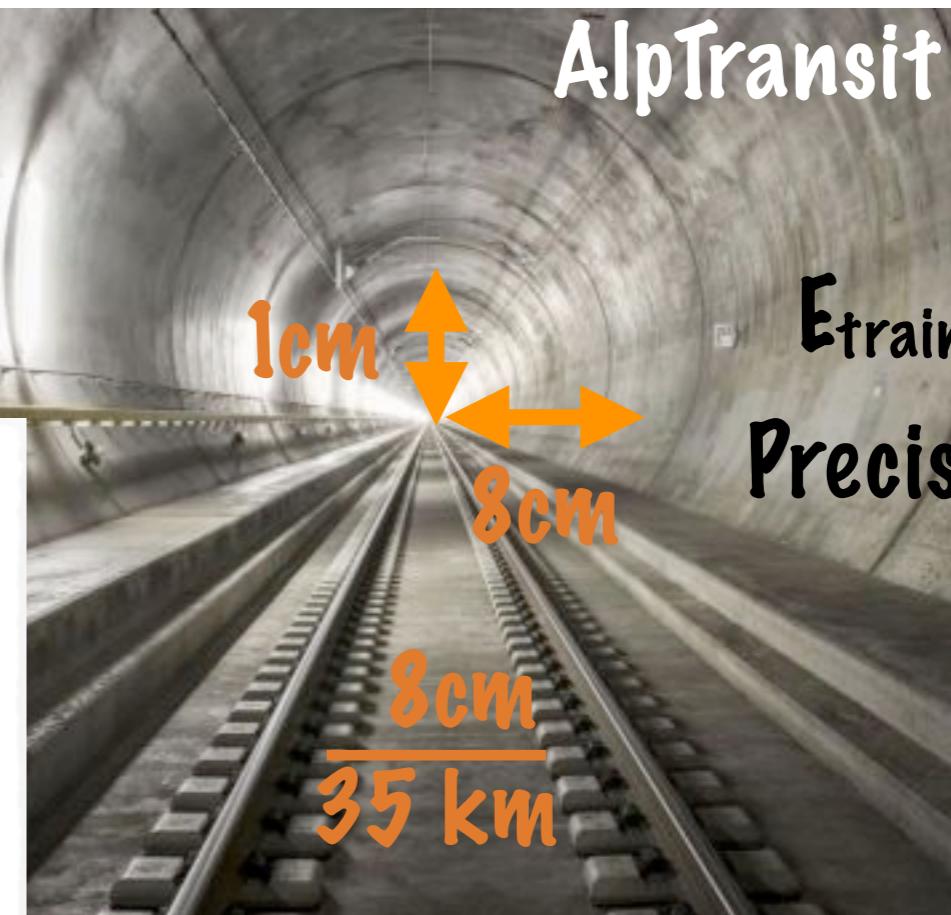
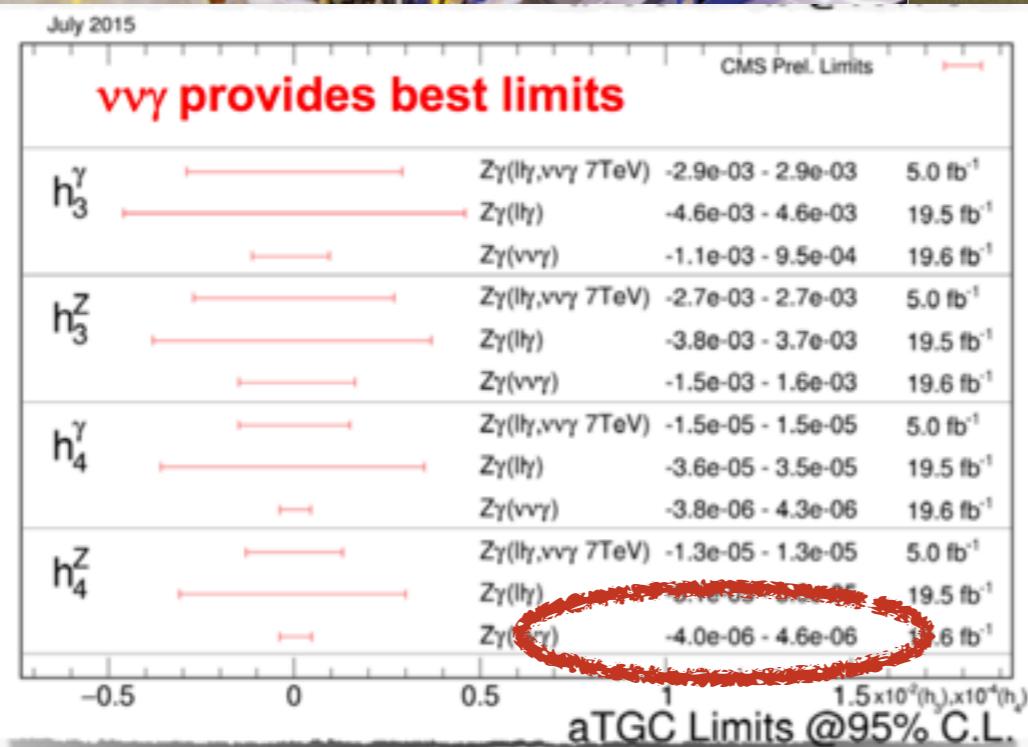
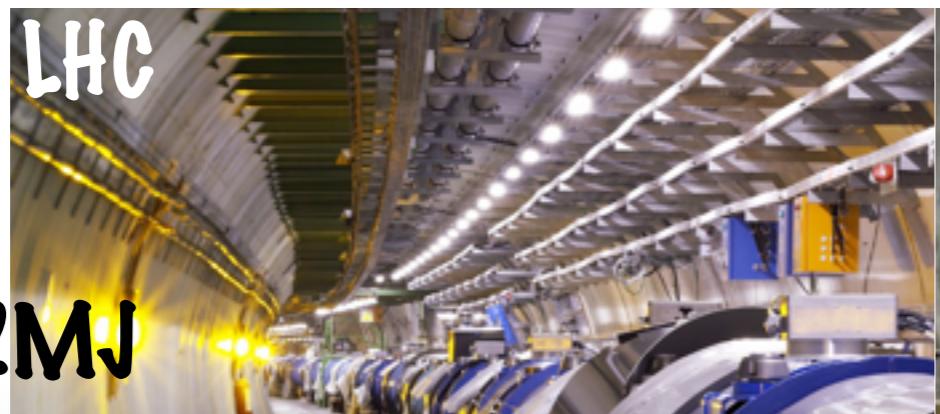
REMEDIOS

- LHC precision tests (di/tri-boson): what is being searched for? -

$L=27\text{ km}$

$E_{\text{beam}}=362\text{ MJ}$

Precision =



$L=35\text{ km}$

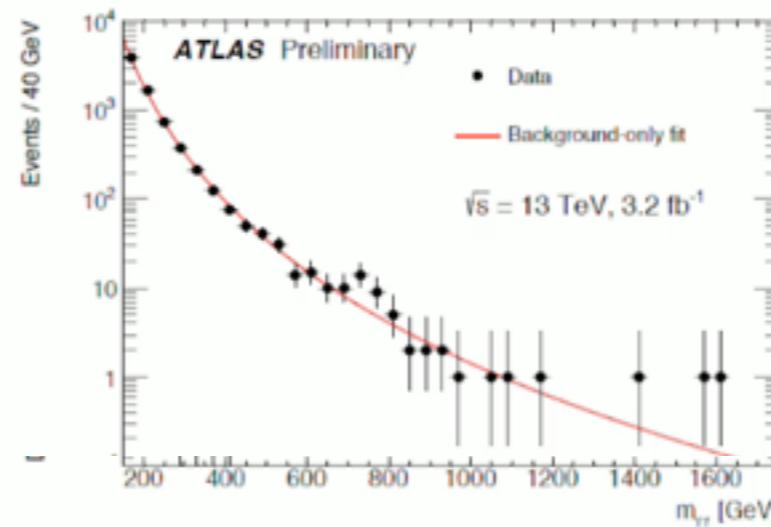
$E_{\text{train}}=362\text{ MJ}$
Precision = 10^{-6}

Francesco Riva
(CERN)

In collaboration with
Liu, Pomarol, Rattazzi 1603.03064,
Azatov, Contino, Machado 1607.05236

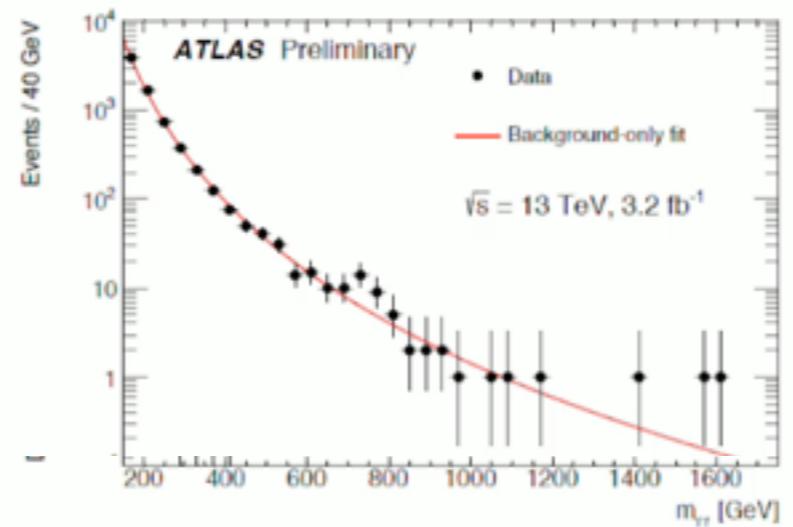
Two modes of exploration at LHC:

A) Direct Searches:



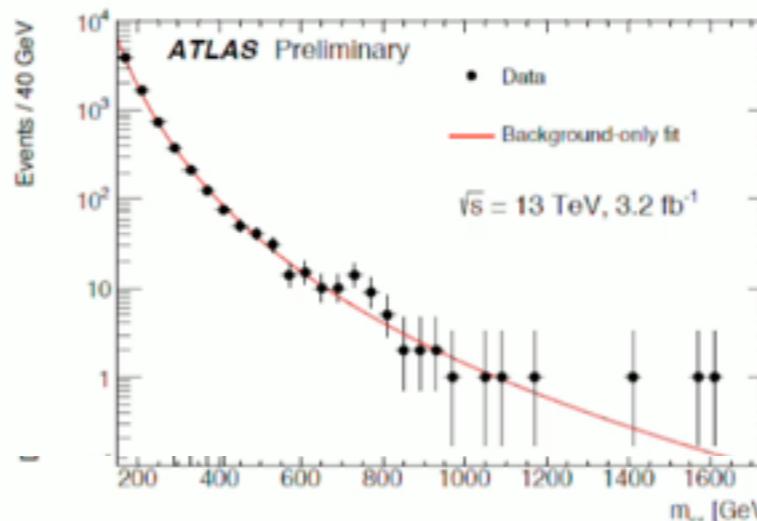
Two modes of exploration at LHC:

A) Direct Searches:

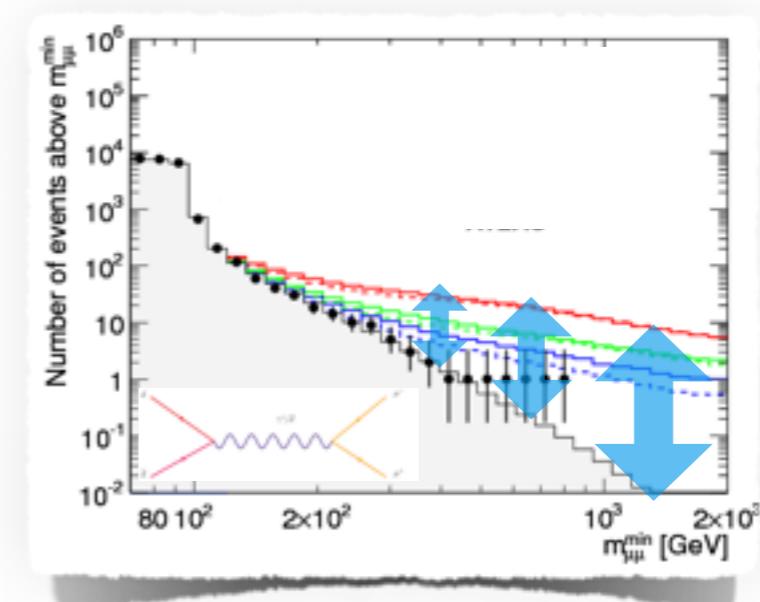
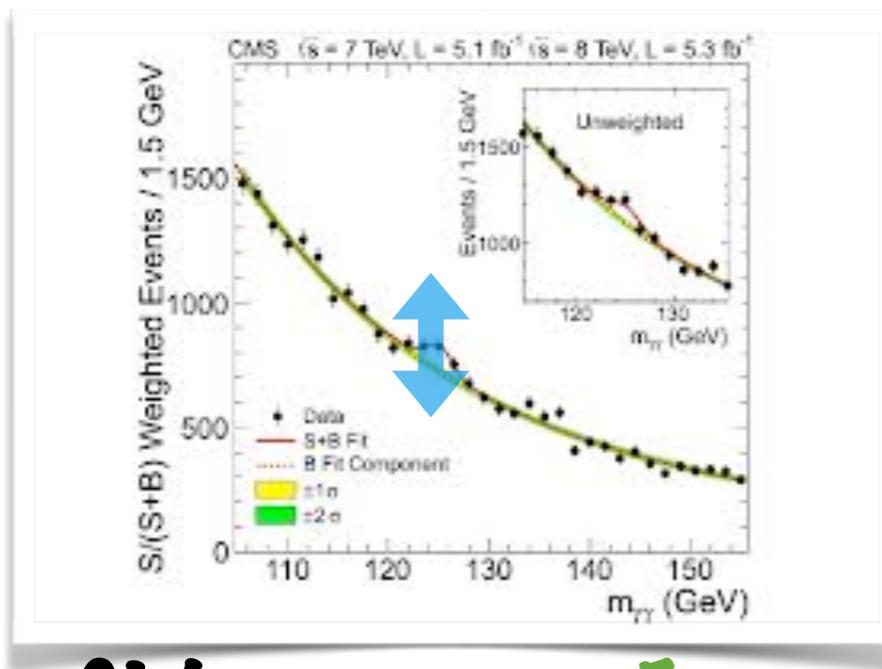


Two modes of exploration at LHC:

A) Direct Searches:



B) Indirect Searches:

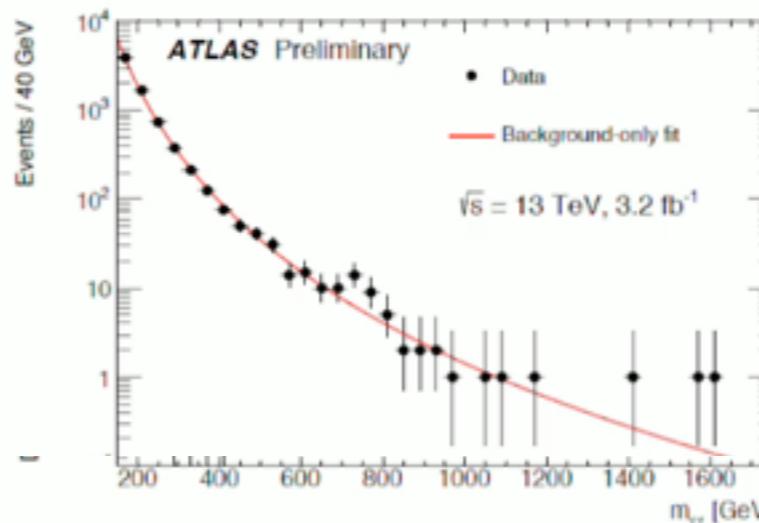


1) On SM resonance $E=m_Z, m_h$

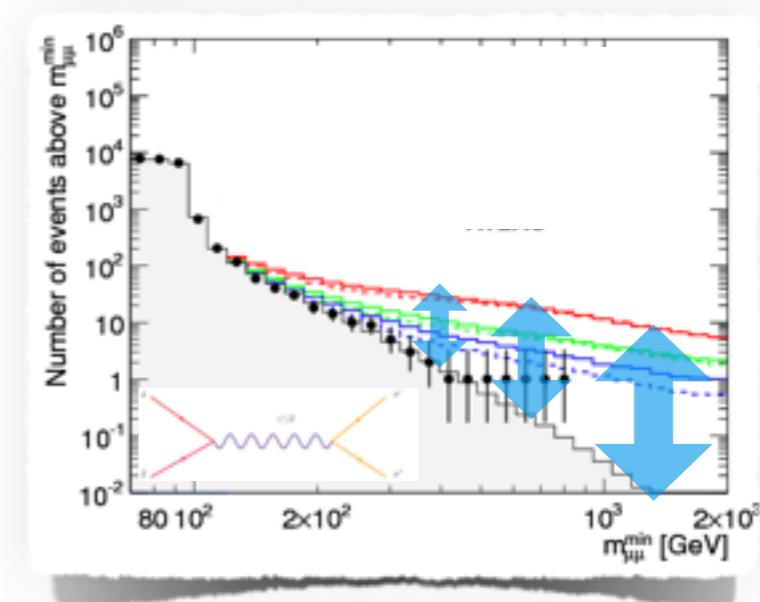
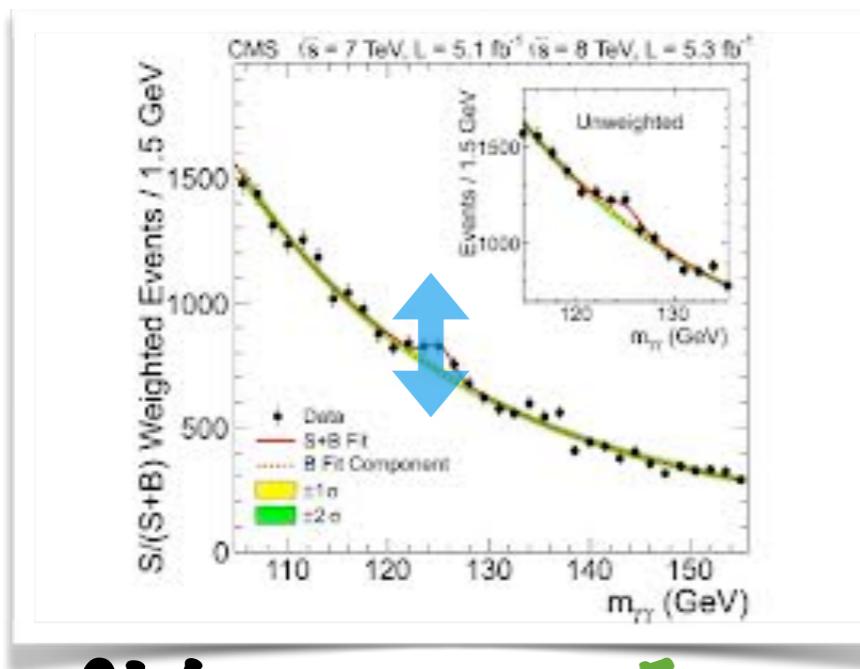
2) Off SM resonance $E \gg m_Z, m_h$

Two modes of exploration at LHC:

A) Direct Searches:



B) Indirect Searches:

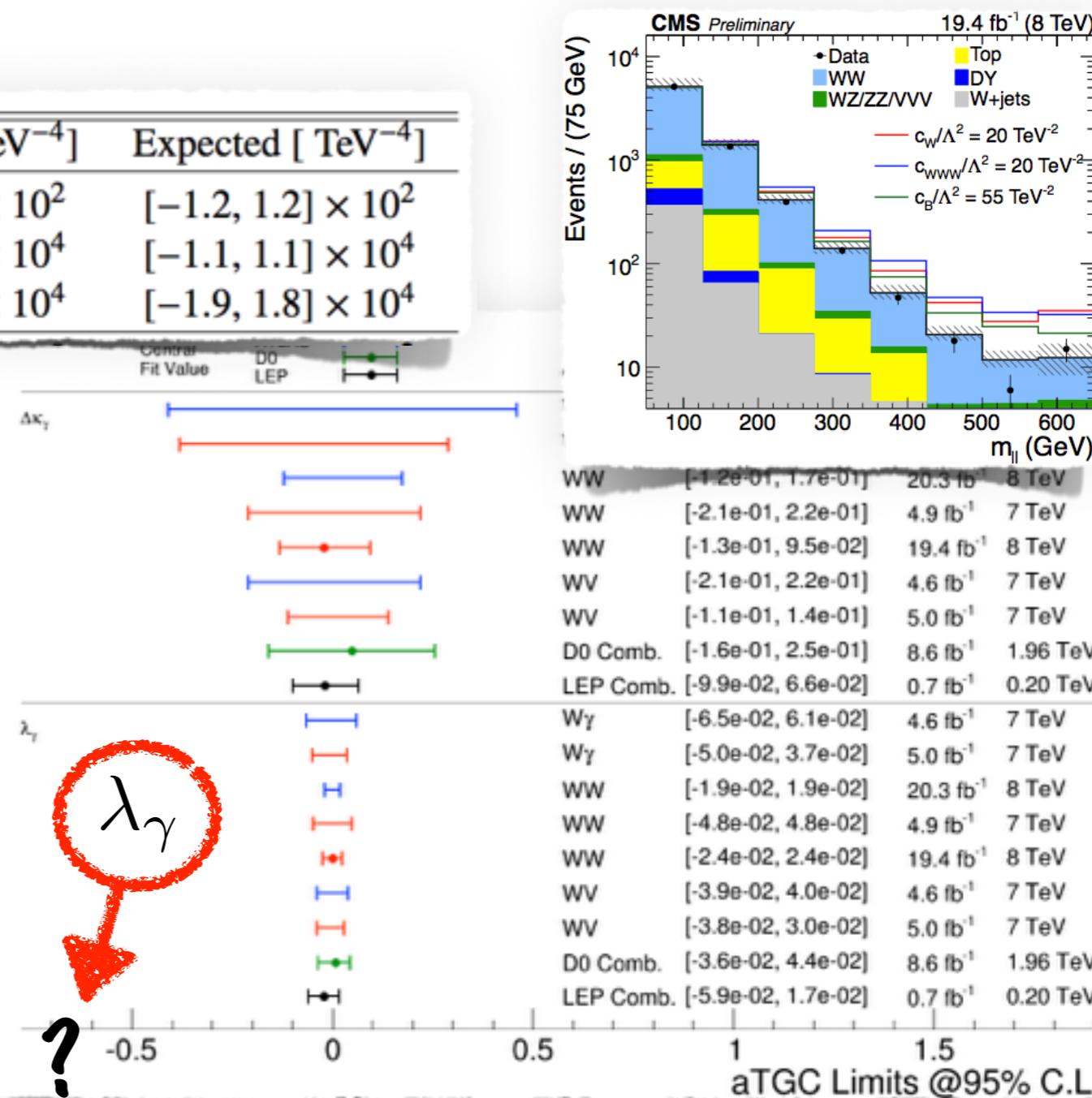
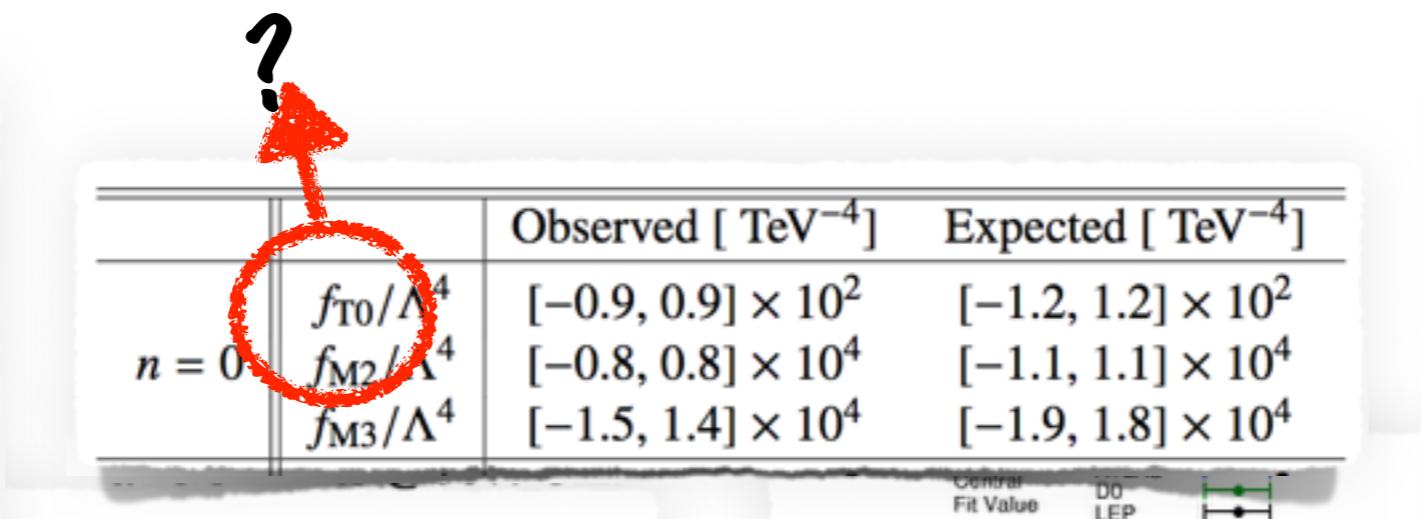
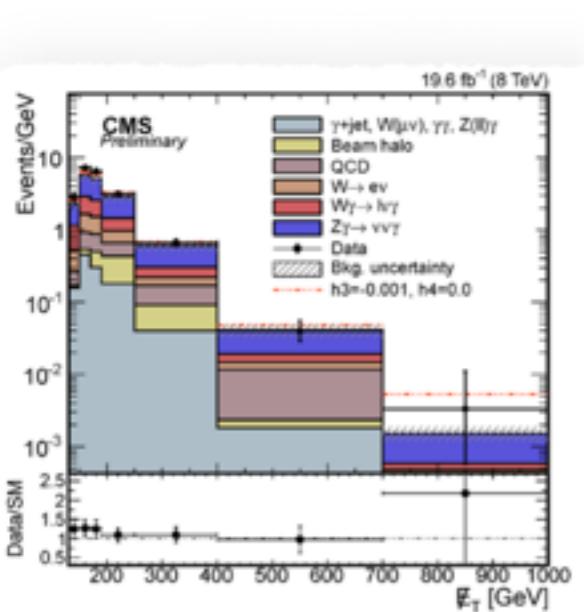


1) On SM resonance $E=m_Z, m_h$

2) Off SM resonance $E \gg m_Z, m_h$

From LHC

Plenty of Data in multiboson processes:



Why nobody cares?

Part 1

why nobody cares

1 Smallness

These parameters \leftrightarrow EFT coefficients:

$$\lambda_\gamma \leftrightarrow$$

dim-6

$$\epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$$

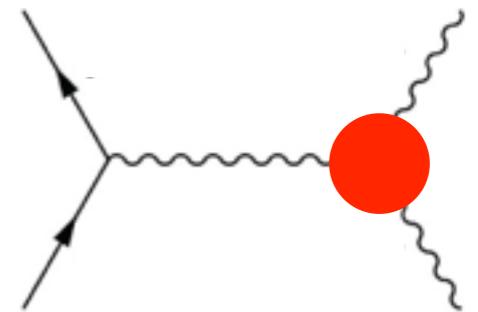
$$h_3^Z \leftrightarrow$$

dim-8

$$H^\dagger D_\mu H D_\rho B_\nu^\rho \tilde{B}^{\mu\nu}$$

$$f_{T,0} \leftrightarrow$$

$$W_{\mu\nu}^4$$



Involve the **transverse polarizations** of vectors

1 Smallness

These parameters \leftrightarrow EFT coefficients:

$$\lambda_\gamma \leftrightarrow$$

dim-6

$$\epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$$

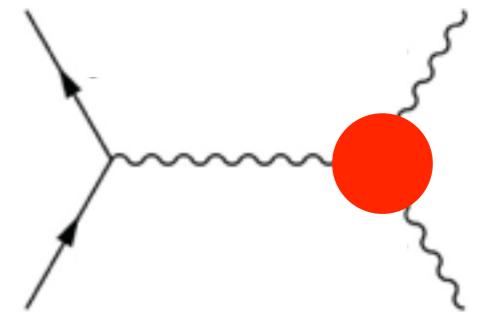
$$h_3^Z \leftrightarrow$$

dim-8

$$H^\dagger D_\mu H D_\rho B_\nu^\rho \tilde{B}^{\mu\nu}$$

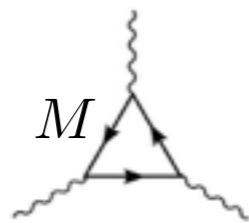
$$f_{T,0} \leftrightarrow$$

$$W_{\mu\nu}^4$$



Involve the **transverse polarizations of vectors**

In popular models (e.g. SUSY, CH), V_T elementary, these are tiny:



$$g\lambda_\gamma \sim \frac{g^3}{16\pi^2} \frac{m_W^2}{M^2}$$

SILH: Giudice, Grojean, Pomarol, Rattazzi'2007

2 Non-Interference for BSM_6 amplitudes

Azatov,Contino,Machado,FR'16

There is in fact one more obstruction to test transverse vectors:

For $E \gg m_W$ states have well defined helicity h

Amplitudes for $2 \rightarrow 2$ with different h_{tot} don't interfere

2 Non-Interference for BSM_6 amplitudes

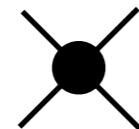
Azatov,Contino,Machado,FR'16

There is in fact one more obstruction to test transverse vectors:

For $E \gg m_W$ states have well defined helicity h

Amplitudes for $2 \rightarrow 2$ with different h_{tot} don't interfere

Selection Rule:



A_4	$ h(A_4^{\text{SM}}) $	$ h(A_4^{\text{BSM}}) $
$VVVV$	0	4,2
$VV\phi\phi$	0	2
$VV\psi\psi$	0	2
$V\psi\psi\phi$	0	2
$\psi\psi\psi\psi$	2,0	2,0
$\psi\psi\phi\phi$	0	0
$\phi\phi\phi\phi$	0	0

Any BSM dim-6 operator

tree level + at least one transverse vector ➤ No interference

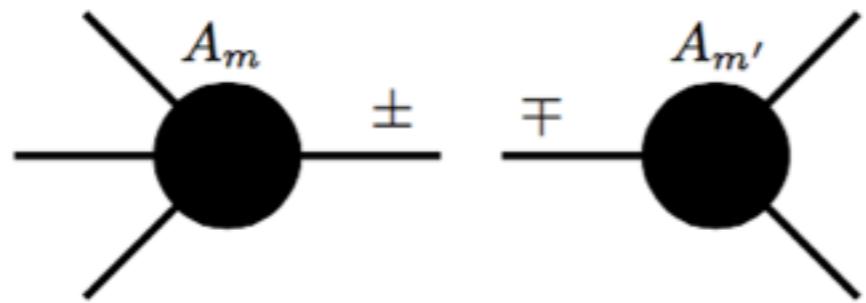
2 Non-Interference for BSM₆ amplitudes

Azatov,Contino,Machado,FR'16

How?**i) Helicity sums:**

$$h(A_n) = h(A_m) + h(A_{m'})$$

\nwarrow
 $n=m+m'-2 \text{ legs}$



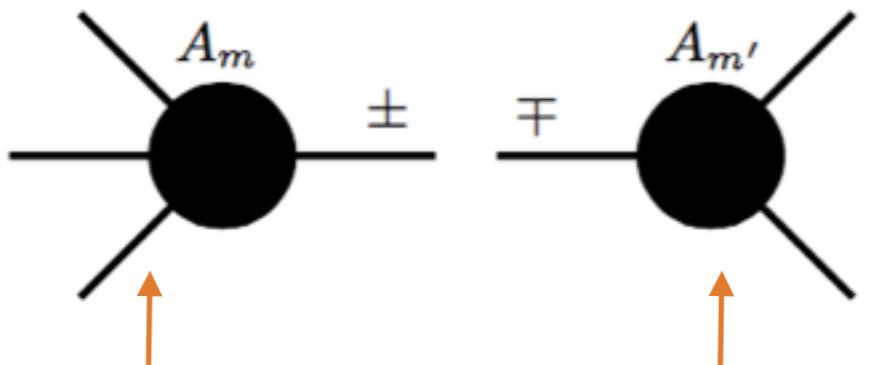
2 Non-Interference for BSM_6 amplitudes

Azatov,Contino,Machado,FR'16

How?**i) Helicity sums:**

$$h(A_n) = h(A_m) + h(A_{m'})$$

$n=m+m'-2$ legs



$p \in \mathbb{C}$ so that on-shell condition $p^2 = 0$ satisfied also for A_3

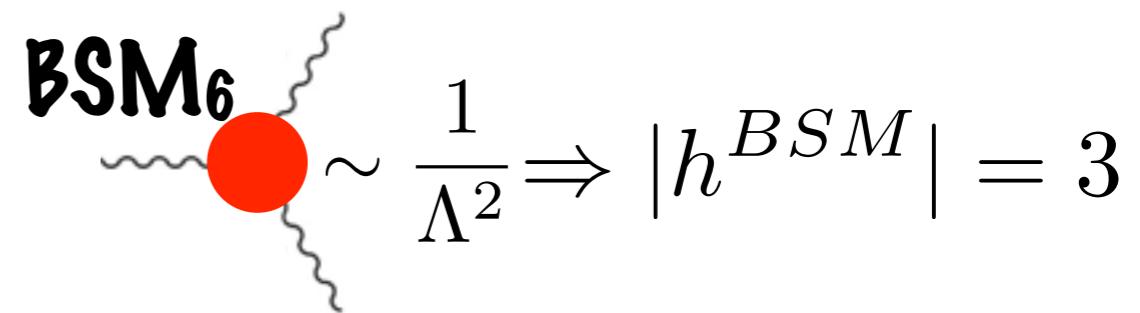
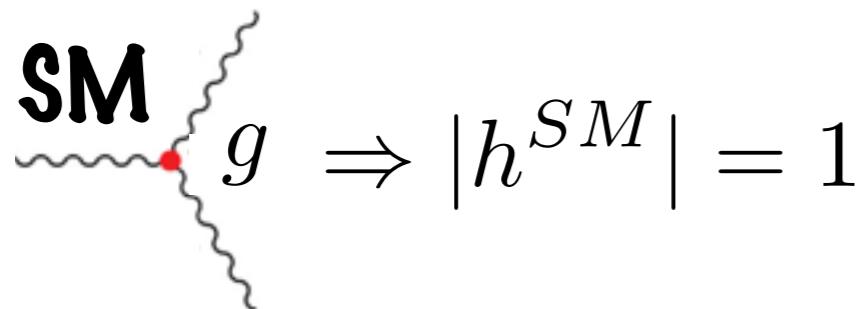
SM or BSM_6

2 Non-Interference for BSM_6 amplitudes

Azatov,Contino,Machado,FR'16

How?**ii) Helicity of 3-point \leftrightarrow coupling dimension:**

$$|h(A_3)| = 1 - [g]$$



From: Little group scaling + dimensional analysis

(see e.g. Elvang,Huang'13)

2 Non-Interference for BSM₆ amplitudes

Azatov,Contino,Machado,FR'16

How?

iii) SUSY* Ward Identities: $|h(A_4^{SM})| < 2$ (except $\psi^+ \psi^+ \psi^+ \psi^+$)
(aka Ultra Helicity Violation)

2 Non-Interference for BSM_6 amplitudes

Azatov,Contino,Machado,FR'16

How?

iii) SUSY* Ward Identities:
 (aka Ultra Helicity Violation)

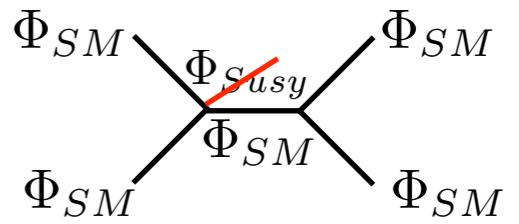
$$|h(A_4^{SM})| < 2$$

(except $\psi^+ \psi^+ \psi^+ \psi^+$)

*in the limit of either $y_u=0$ or $y_d=y_l=0$:

SM upliftable to SUSY+R-parity (with 1 Higgs doublet)

Grisaru,Pendleton,vanNieuwenhuizen'77



2 Non-Interference for BSM_6 amplitudes

Azatov,Contino,Machado,FR'16

How?

iii) SUSY* Ward Identities:
 (aka Ultra Helicity Violation)

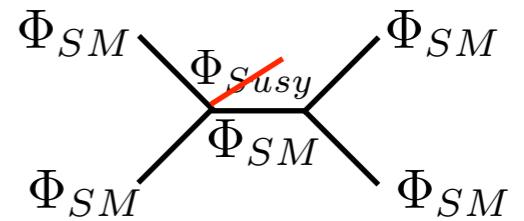
$$|h(A_4^{SM})| < 2$$

(except $\psi^+ \psi^+ \psi^+ \psi^+$)

*in the limit of either $y_u=0$ or $y_d=y_l=0$:

SM upliftable to SUSY+R-parity (with 1 Higgs doublet)

Grisaru,Pendleton,vanNieuwenhuizen'77



→ e.g. $0 = \langle 0 | [Q, \Psi^+ V^+ V^+ V^+] | 0 \rangle = \sum_i \langle 0 | \Psi^+ \dots [Q, V^+] \dots V^+ | 0 \rangle \propto \langle 0 | V^+ V^+ V^+ V^+ | 0 \rangle$

$[Q, \psi] \sim V \quad [Q, V] \sim \psi$

$\psi^+ \psi^+ \dots = 0$

$SM : A(V^+ V^+ V^+ V^-) = A(V^+ V^+ \psi^+ \psi^-) = A(V^+ V^+ \phi \phi) = A(V^+ \psi^+ \psi^+ \phi) = 0$

2 Non-Interference for BSM_6 amplitudes

Azatov,Contino,Machado,FR'16

How?

iii) SUSY* Ward Identities:
 (aka Ultra Helicity Violation)

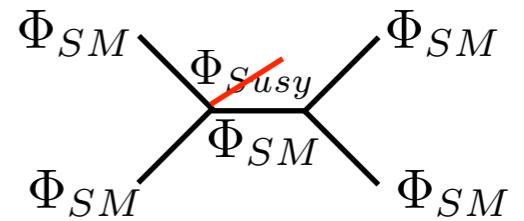
$$|h(A_4^{SM})| < 2$$

(except $\psi^+ \psi^+ \psi^+ \psi^+$)

*in the limit of either $y_u=0$ or $y_d=y_l=0$:

SM upliftable to SUSY+R-parity (with 1 Higgs doublet)

Grisaru,Pendleton,vanNieuwenhuizen'77



→ e.g. $0 = \langle 0 | [Q, \Psi^+ V^+ V^+ V^+] | 0 \rangle = \sum_i \langle 0 | \Psi^+ ... [Q, V^+] ... V^+ | 0 \rangle \propto \langle 0 | V^+ V^+ V^+ V^+ | 0 \rangle$

$[Q, \psi] \sim V$ $[Q, V] \sim \psi$

$\psi^+ \psi^+ \dots = 0$

$$SM : A(V^+ V^+ V^+ V^-) = A(V^+ V^+ \psi^+ \psi^-) = A(V^+ V^+ \phi \phi) = A(V^+ \psi^+ \psi^+ \phi) = 0$$

BSM : Operators with transverse V not supersymmetrizable

Elias-Miro,Espinosa,Pomarol'14

2 Non-Interference for BSM₆ amplitudes

Azatov,Contino,Machado,FR'16

How?

i) Helicity sums $h(A_n) = h(A_m) + h(A_{m'})$

ii) Helicity of 3-point \leftrightarrow coupling dimension $|h(A_3)| = 1 - [g]$

iii) SUSY* Ward Identities $|h(A_4^{SM})| < 2$

2 Non-Interference for BSM_6 amplitudes

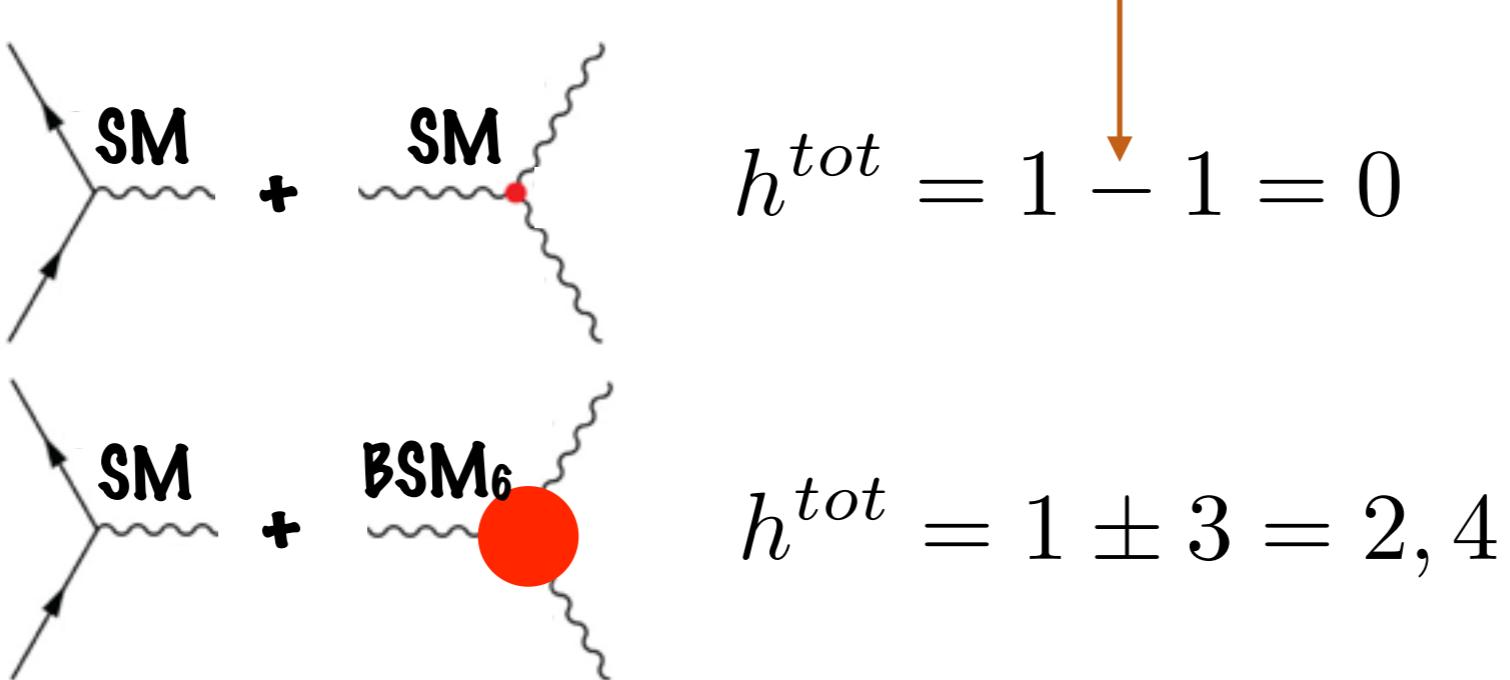
Azatov,Contino,Machado,FR'16

How?

i) Helicity sums $h(A_n) = h(A_m) + h(A_{m'})$

ii) Helicity of 3-point \leftrightarrow coupling dimension $|h(A_3)| = 1 - [g]$

iii) SUSY* Ward Identities $|h(A_4^{SM})| < 2$



2 Non-Interference for BSM_6 amplitudes

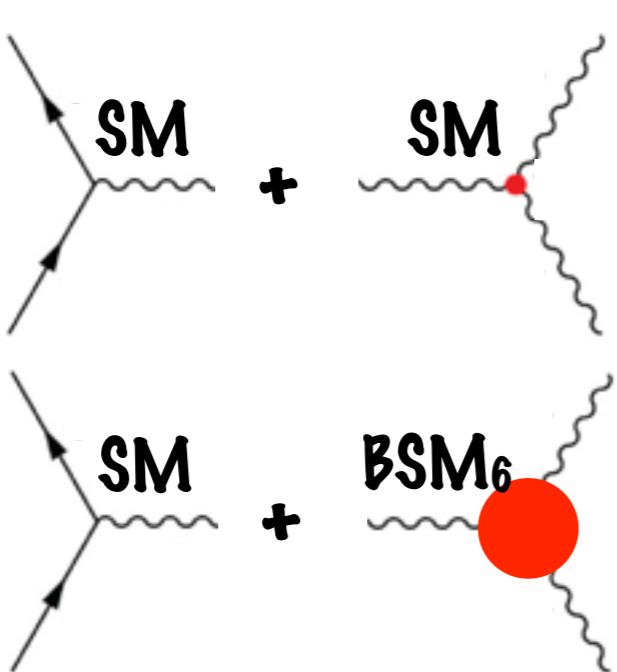
Azatov,Contino,Machado,FR'16

How?

i) Helicity sums $h(A_n) = h(A_m) + h(A_{m'})$

ii) Helicity of 3-point \leftrightarrow coupling dimension $|h(A_3)| = 1 - [g]$

iii) SUSY* Ward Identities $|h(A_4^{SM})| < 2$



$$h^{tot} = 1 - 1 = 0$$

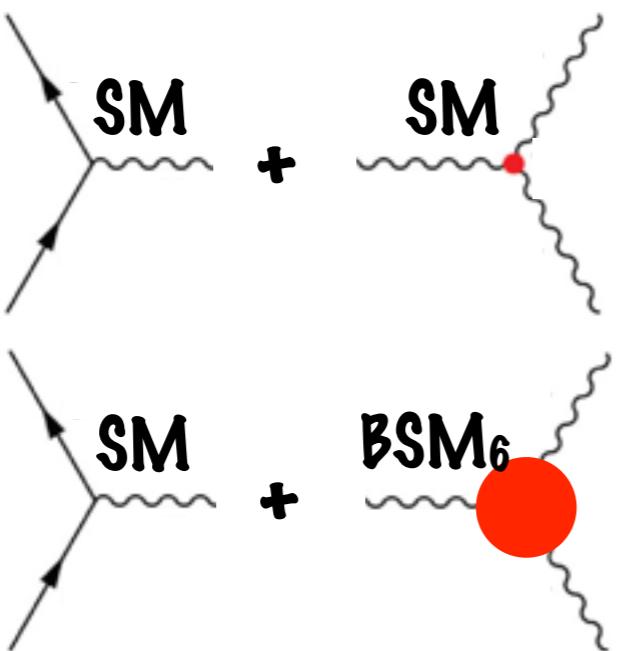
$$h^{tot} = 1 \pm 3 = 2, 4$$

No Interference
(dim-6, 4-point)

2 Non-Interference for BSM₆ amplitudes

Azatov,Contino,Machado,FR'16

How?

i) Helicity sums $h(A_n) = h(A_m) + h(A_{m'})$ ii) Helicity of 3-point \leftrightarrow coupling dimension $|h(A_3)| = 1 - [g]$ iii) SUSY* Ward Identities $|h(A_4^{SM})| < 2$ 

$$h^{tot} = 1 - 1 = 0$$

No Interference
(dim-6, 4-point)

$$h^{tot} = 1 \pm 3 = 2, 4$$

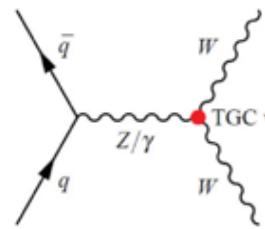
interference vanishes

$$\blacktriangleright \sigma \propto \sum |Amp|^2 \simeq SM^2 \left(1 + c_i \frac{E^2}{M^2} + c_i^2 \frac{E^4}{M^4} \right)$$

Small effects, are even smaller!

Non-Interference+Smallness

Present constraints:



LEP2: $E = 130 - 209 \text{ GeV}$
 $\lambda\gamma \in [-0.059, 0.017] \sim \frac{g^2}{16\pi^2} \frac{m_W^2}{M^2} \Rightarrow M \gtrsim 30 \text{ GeV}$
95% C.L. (LEP EW: 1302.3415)

LHC: $E \gtrsim 1 \text{ TeV}$...
 $\lambda_\gamma \in [-0.019, +0.019]$

...so small that it doesn't even make sense as an EFT.

why nobody cares?

Because there is no structured scenario
where these searches are self-consistent
(need strong coupling)

Part 2

why caring (Remedios)

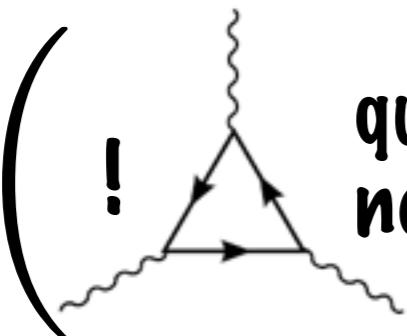
Strongly Coupled BSM?

How can SM be light and **weakly** coupled at $E < m_W$ and
strongly coupled at $E \gg m_W$?



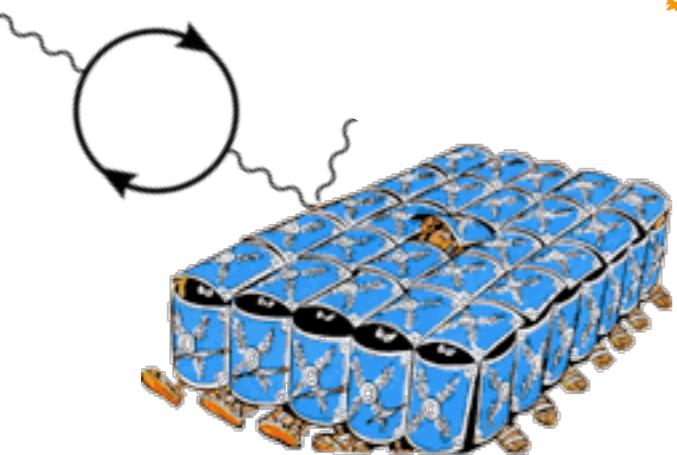
Strongly Coupled BSM?

How can SM be light and **weakly** coupled at $E < m_W$ and
strongly coupled at $E \gg m_W$?

(!  quantum effects generally propagate new couplings to the whole SM !)

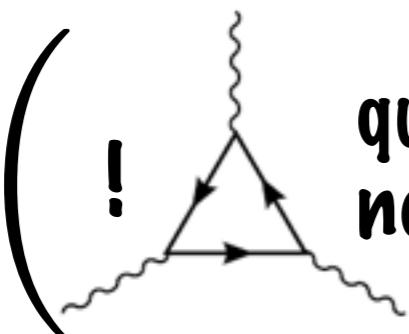
Need shielding:

Approximate Symmetries broken in SM



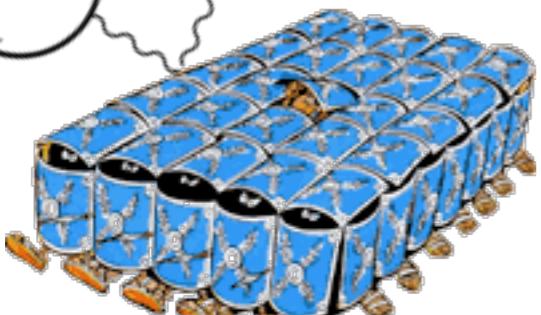
Strongly Coupled BSM?

How can SM be light and **weakly** coupled at $E < m_W$ and
strongly coupled at $E \gg m_W$?

(!  quantum effects generally propagate new couplings to the whole SM !)

Need shielding:

Approximate Symmetries broken in SM



gsm weak
SM Lagrangian
(dim-4)

Marginal

$$A \simeq g_{SM}^2 \left(1 + \frac{g_*^2}{g_{SM}^2} \frac{E^2}{M^2} \right) \equiv g^2(E)$$



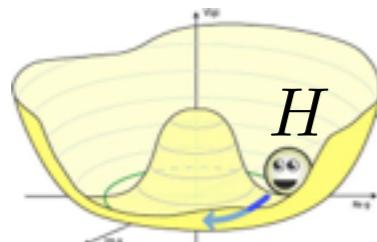
g* strong
Higher-dim
operators
(dim-6, dim-8...)
Irrelevant

Strongly Coupled BSM and Approximate Symmetries

Scalars: Composite Higgs

Higgs is a Pseudo Goldstone boson from new strong sector (symm=SO(5)/SO(4))

Georgi,Kaplan'84; Agashe,Contino,Nomura,Pomarol'04; Giudice,Grojean,Pomarol,Rattazzi'07;...



Shift Symm: $H \rightarrow H + c_{+n.l.}$

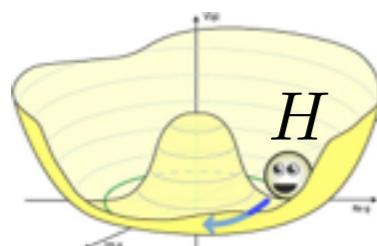
$$g_* \partial_\mu H_{+n.l.} \checkmark \quad \epsilon H \times \text{symmetry}$$

Strongly Coupled BSM and Approximate Symmetries

Scalars: Composite Higgs

Higgs is a Pseudo Goldstone boson from new strong sector (symm=SO(5)/SO(4))

Georgi,Kaplan'84; Agashe,Contino,Nomura,Pomarol'04; Giudice,Grojean,Pomarol,Rattazzi'07;...



Shift Symm: $H \rightarrow H + c_{+n.l.}$

$g_* \partial_\mu H_{+n.l.} \checkmark$ $\epsilon H \times$
symmetry

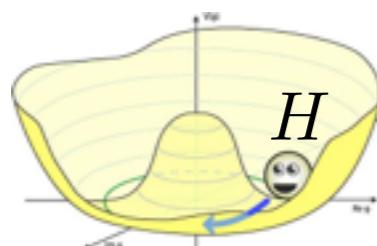
- ▶ $\frac{g_*^2}{M^2} (\partial_\mu |H|^2)^2$ large
- ▶ $\lambda |H|^4$ small

Strongly Coupled BSM and Approximate Symmetries

Scalars: Composite Higgs

Higgs is a Pseudo Goldstone boson from new strong sector (symm=SO(5)/SO(4))

Georgi,Kaplan'84; Agashe,Contino,Nomura,Pomarol'04; Giudice,Grojean,Pomarol,Rattazzi'07;...



Shift Symm: $H \rightarrow H + c_{+n.l.}$

$$g_* \partial_\mu H_{+n.l.} \checkmark \quad \epsilon H \times \text{symmetry}$$

- ▶ $\frac{g_*^2}{M^2} (\partial_\mu |H|^2)^2$ large
- ▶ $\lambda |H|^4$ small
- ▶ Large effects $V_L V_L \rightarrow V_L V_L$ at high- E :

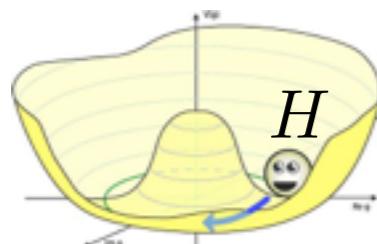
$$A \sim \lambda \left(1 + \frac{g_*^2}{\lambda} \frac{E^2}{M^2} \right)$$

Strongly Coupled BSM and Approximate Symmetries

Scalars: Composite Higgs

Higgs is a Pseudo Goldstone boson from new strong sector (symm=SO(5)/SO(4))

Georgi,Kaplan'84; Agashe,Contino,Nomura,Pomarol'04; Giudice,Grojean,Pomarol,Rattazzi'07;...



Shift Symm: $H \rightarrow H + c_{+n.l.}$

$$g_* \partial_\mu H_{+n.l.} \checkmark \quad \epsilon H \times \text{symmetry}$$

- ▶ $\frac{g_*^2}{M^2} (\partial_\mu |H|^2)^2$ large
- ▶ $\lambda |H|^4$ small
- ▶ Large effects $V_L V_L \rightarrow V_L V_L$ at high- E :

$$A \sim \lambda \left(1 + \frac{g_*^2}{\lambda} \frac{E^2}{M^2} \right)$$

Vectors: ?

Strong transverse vectors at high- E

Problem: Naively, contrary to scalars, gauge bosons couple universally

Gauge bosons: **weak** SM coupling ($\partial_\mu + igA_\mu$)

Strong transverse vectors at high-E

Problem: Naively, contrary to scalars, gauge bosons couple universally
Gauge bosons: weak SM coupling ($\partial_\mu + igA_\mu$)

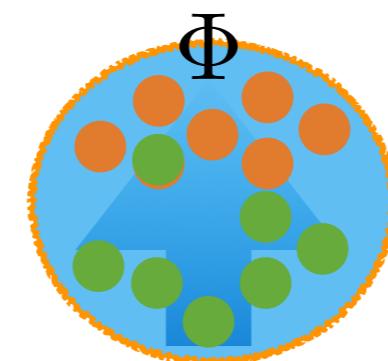
However:

$$g = 0$$

- ▶ A_μ composite, g^*
- ▶ No light charged d.o.f
(but large dipoles possible)

$$\mathcal{L} = L \left(\frac{\partial_\mu}{M}, g_* \frac{\hat{F}_{\mu\nu}}{M^2}, \Phi \right)$$

(Euler-Heisenberg)



- ▶ Strong higher-d interactions

Strong transverse vectors at high-E

Problem: Naively, contrary to scalars, gauge bosons couple universally
Gauge bosons: weak SM coupling ($\partial_\mu + igA_\mu$)

However:

$$g = 0$$



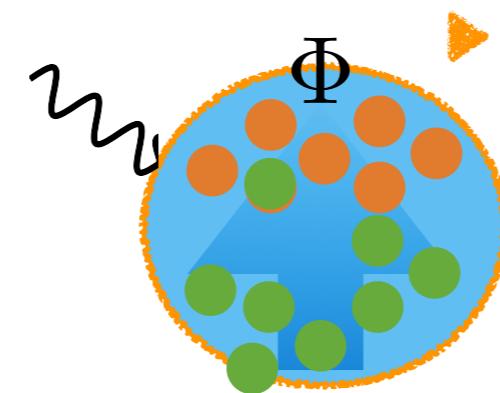
$$g \neq 0$$

► A_μ composite, g^*

► No light charged d.o.f
(but large dipoles possible)

$$\mathcal{L} = L \left(\frac{\partial_\mu}{M}, g_* \frac{\hat{F}_{\mu\nu}}{M^2}, \Phi \right)$$

(Euler-Heisenberg)



$$q = \frac{g}{g_*}$$

► Light d.o.f. with charge $\frac{g}{g_*} \ll 1$

$$\mathcal{L} = L \left(\frac{\partial_\mu + igA_\mu}{M}, g_* \frac{F_{\mu\nu}}{M^2}, \Phi \right)$$

► Weak marginal interactions

► Strong higher-d interactions

Strong transverse vectors at high-E

What symmetry is broken by g and not by g^* ? none...

Multipole interactions
(field strengths)

$$g_* F_{\mu\nu}^a$$

$$U(1)_{local}^3 \times SU(2)_{global}$$

Non-abelian

Monopole interactions
(covariant derivative)

$$\partial_\mu + ig A_\mu^a T^a$$

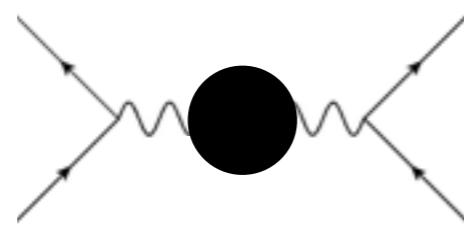


Inonu-Wigner Contraction ('53)
#generators invariant
(Like Poincaré->Galilei)

- Symmetry deformation: a new selection rule to build EFTs
(different from spurions)

Strong transverse vectors: Implications

► $\frac{1}{M^2} (D_\rho W_\mu^{a,\nu})^2$



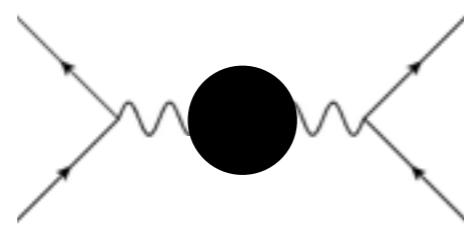
$$W, Y \simeq \frac{m_W^2}{M^2} \lesssim 10^{-3}$$

Barbieri, Pomarol, Rattazzi, Strumia'04

→ $M \gtrsim 2 \text{ TeV}$

Strong transverse vectors: Implications

► $\frac{1}{M^2} (D_\rho W_\mu^{a,\nu})^2$

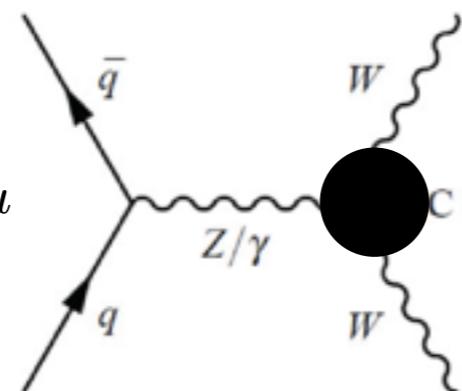


$$W, Y \simeq \frac{m_W^2}{M^2} \lesssim 10^{-3}$$

$$\rightarrow M \gtrsim 2 \text{ TeV}$$

Barbieri, Pomarol, Rattazzi, Strumia'04

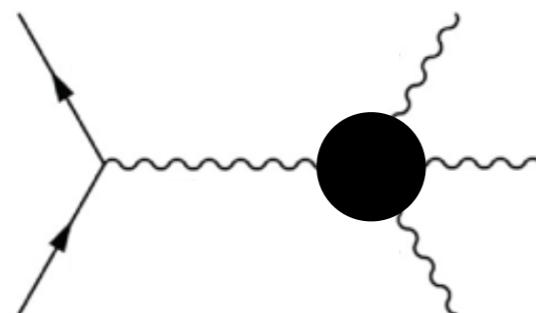
► $\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^b W^{c\rho\mu}$



$$\lambda_\gamma \approx g_* \frac{m_W^2}{M^2}$$

$$M \gtrsim \sqrt{\frac{g_*}{4\pi}} 2.2 \text{ TeV}$$

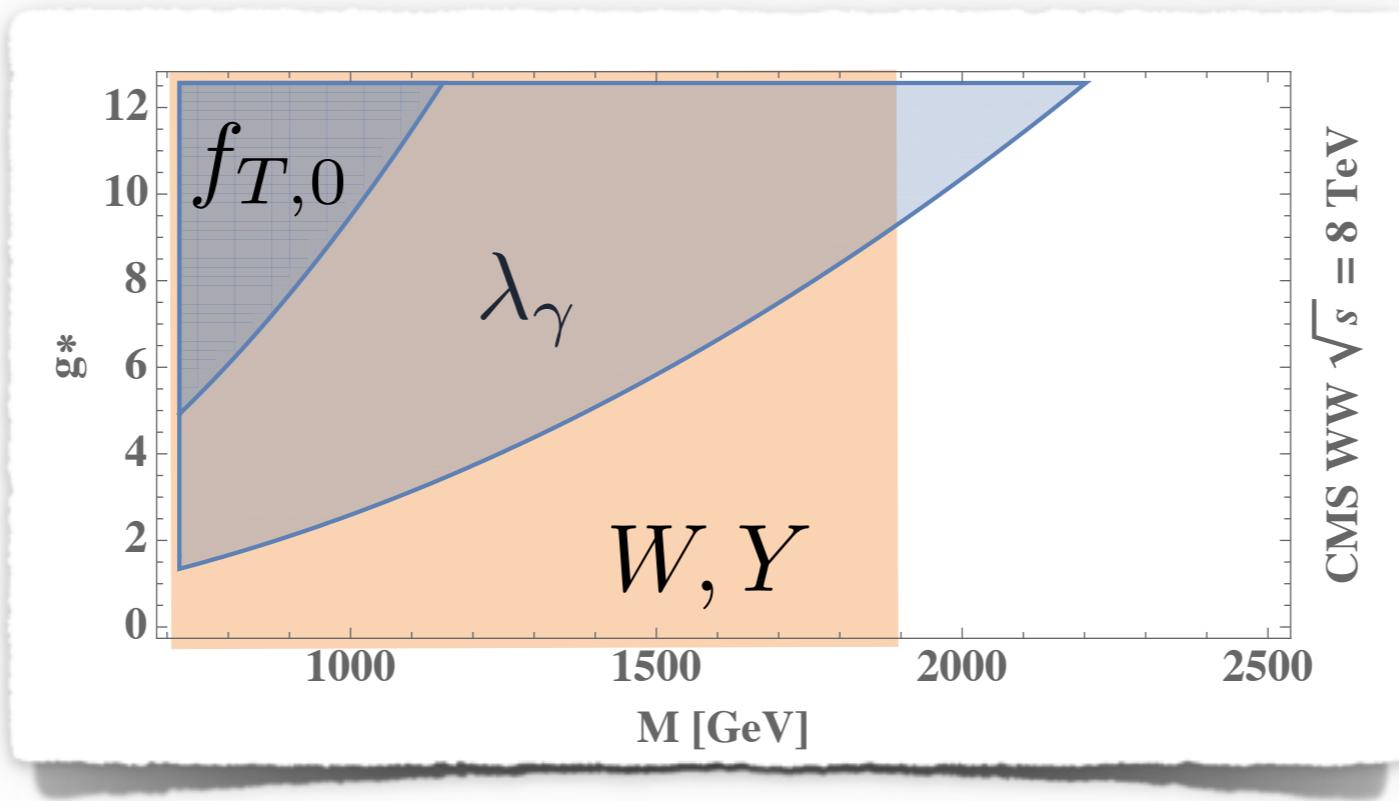
► $\frac{g_*^2}{M^4} W_{\mu\nu}^4$



$$f_{T,0} \approx \frac{g_*^2}{M^4}$$

$$M \gtrsim \sqrt{\frac{g_*}{4\pi}} 1.1 \text{ TeV}$$

Strong transverse vectors: Implications



So far the only models that motivate these searches consistently

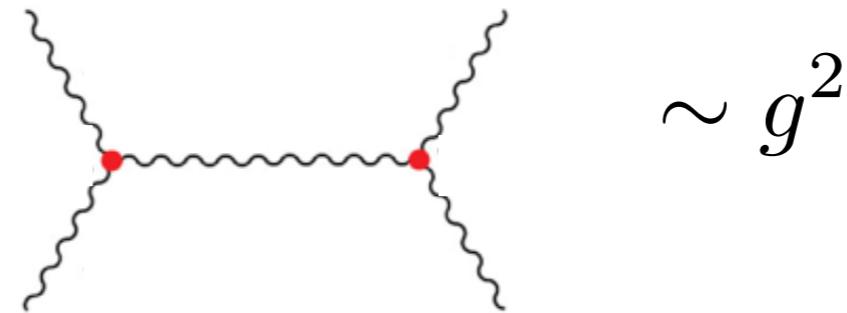
Remedios

(Remedios the Beauty was not a creature of this world-
Gabriel Garcia Marquez)

Strong transverse vectors: Implications

W_T scattering

SM

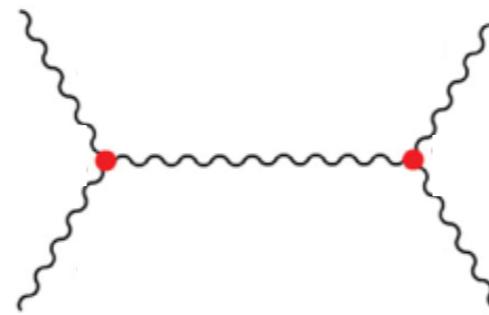


$$\sim g^2$$

Strong transverse vectors: Implications

W_T scattering

SM

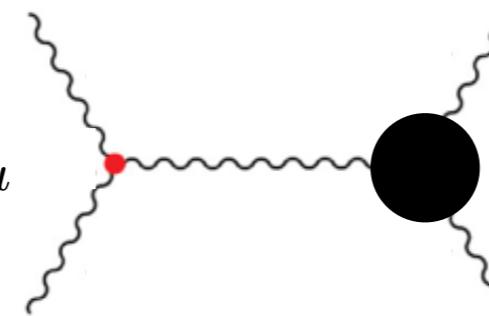


$$\sim g^2$$



Strong vectors:

D6 $\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$

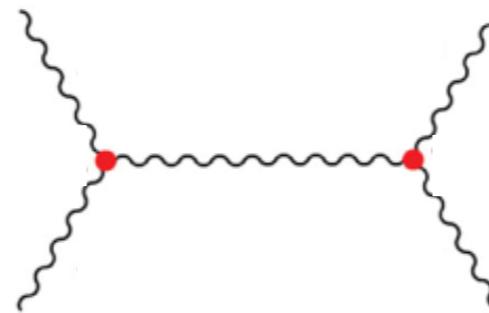


$$\sim g g_* \frac{E^2}{M^2}$$

Strong transverse vectors: Implications

W_T scattering

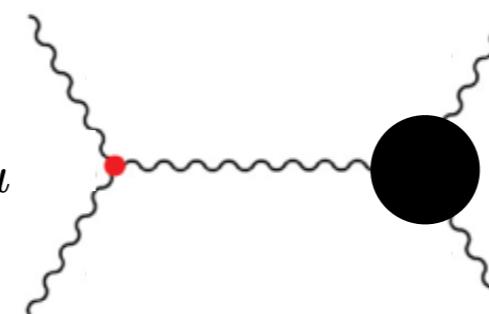
SM



$$\sim g^2$$

Strong vectors:

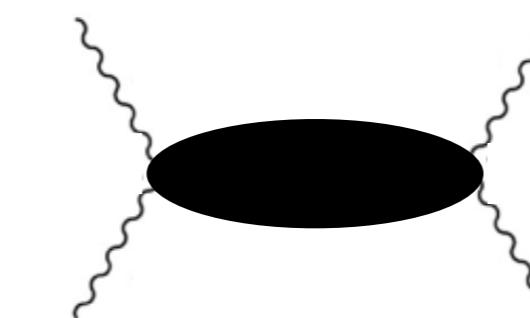
D6 $\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$



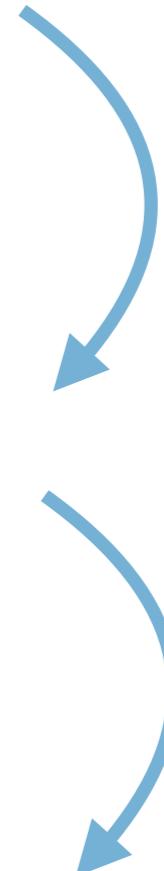
$$\sim gg_* \frac{E^2}{M^2}$$

D8 $\frac{g_*^2}{M^4} W_{\mu\nu}^4$

(NDA: 4-point vertex = coupling²)



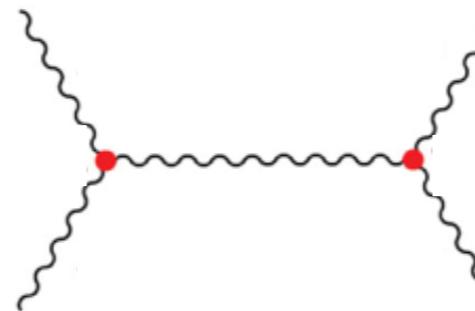
$$\sim g_*^2 \frac{E^4}{M^4}$$



Strong transverse vectors: Implications

W_T scattering

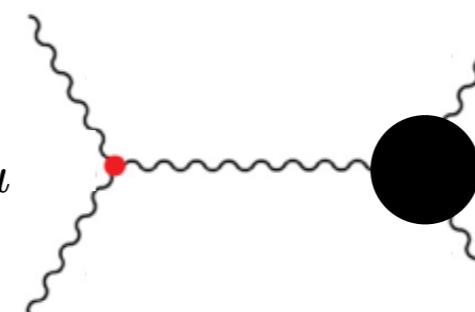
SM



$$\sim g^2$$

Strong vectors:

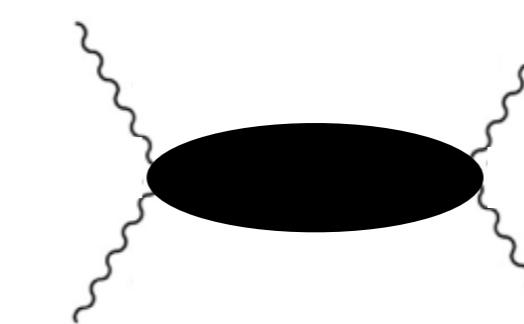
D6 $\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$



$$\sim gg_* \frac{E^2}{M^2}$$

D8 $\frac{g_*^2}{M^4} W_{\mu\nu}^4$

(NDA: 4-point vertex = coupling²)



$$\sim g_*^2 \frac{E^4}{M^4}$$

$$\frac{g_*}{g} \frac{E^2}{M^2} \lesssim 1$$

$$\frac{g_*}{g} \frac{E^2}{M^2} \lesssim 1$$

► $\frac{\delta \mathcal{A}_{BSM}}{\mathcal{A}_{SM}} \lesssim 1$ dimension-6 analysis ok

Strong transverse vectors: Implications

W_T scattering

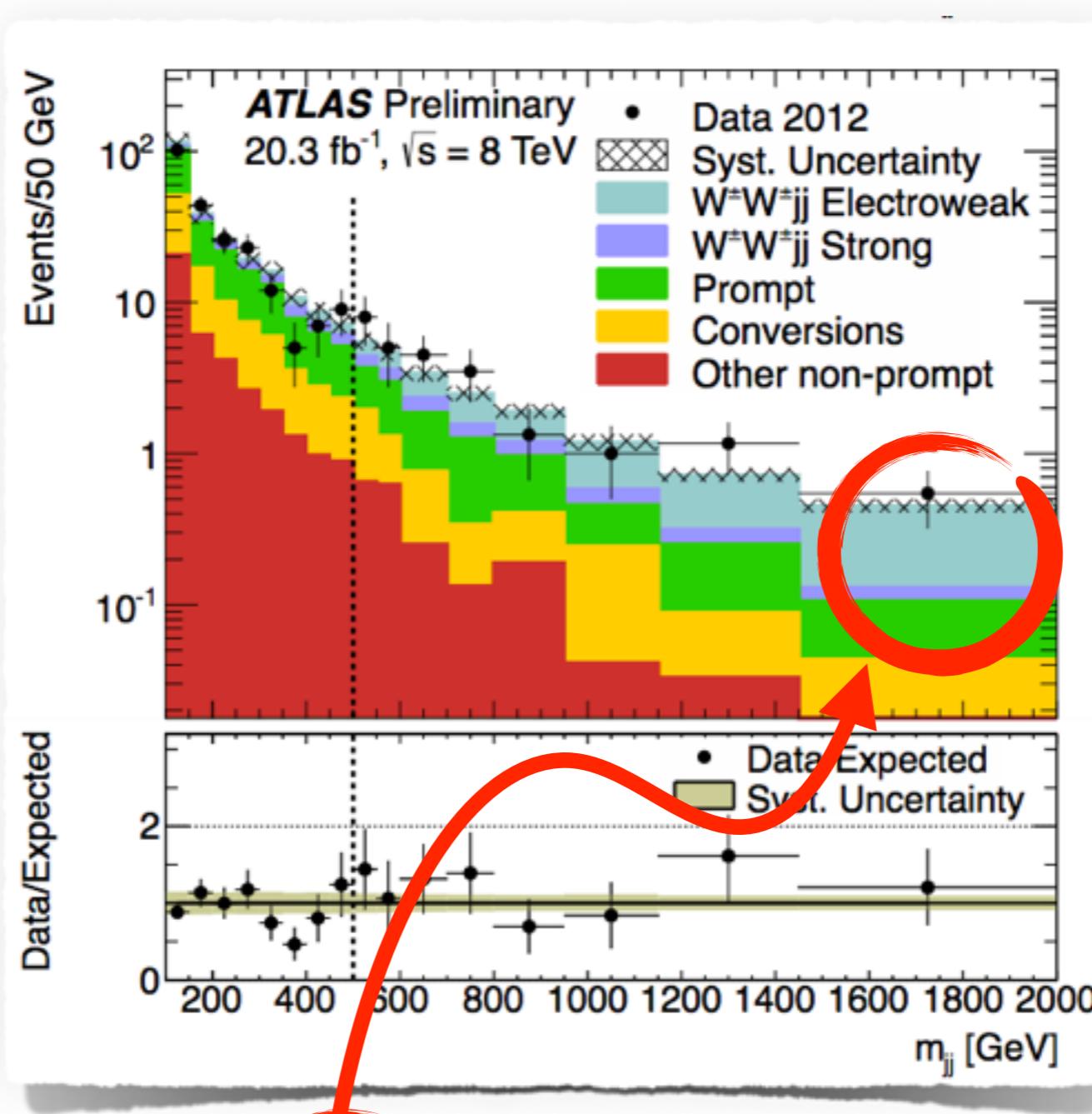
SM

Strong vectors:

$$D6 \quad \frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\ell}^b$$

$$D8 \quad \frac{g_*^2}{M^4} W_{\mu\nu}^4$$

(NDA: 4-point vertex \approx col)



$$\frac{g_*}{g} \frac{E^2}{M^2} \lesssim 1$$

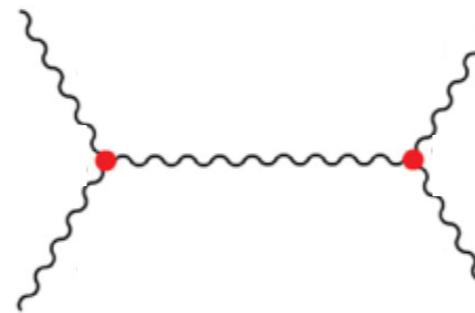
$$\frac{g_*}{g} \frac{E^2}{M^2} \lesssim 1$$

► $\frac{\delta \mathcal{A}_{BSM}}{\mathcal{A}_{SM}} \lesssim 1$ dimension-6 analysis ok

Strong transverse vectors: Implications

W_T scattering

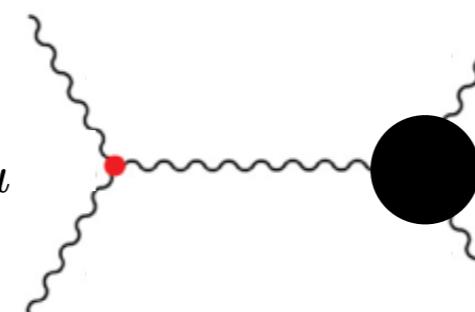
SM



$$\sim g^2$$

Strong vectors:

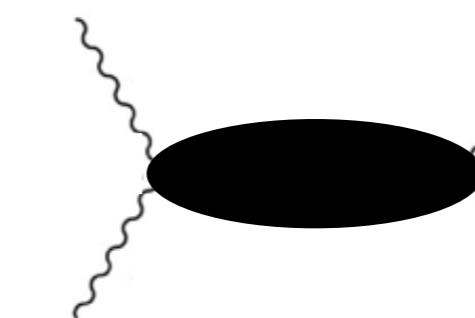
D6 $\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$



$$\sim gg_* \frac{E^2}{M^2}$$

D8 $\frac{g_*^2}{M^4} W_{\mu\nu}^4$

(NDA: 4-point vertex = coupling²)



$$\sim g_*^2 \frac{E^4}{M^4}$$

$$\frac{g_*}{g} \frac{E^2}{M^2} \lesssim 1$$

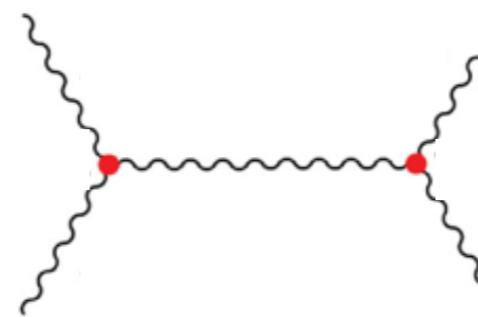
$$\frac{g_*}{g} \frac{E^2}{M^2} \lesssim 1$$

► $\frac{\delta \mathcal{A}_{BSM}}{\mathcal{A}_{SM}} \lesssim 1$ dimension-6 analysis ok

Strong transverse vectors: Implications

W_T scattering

SM



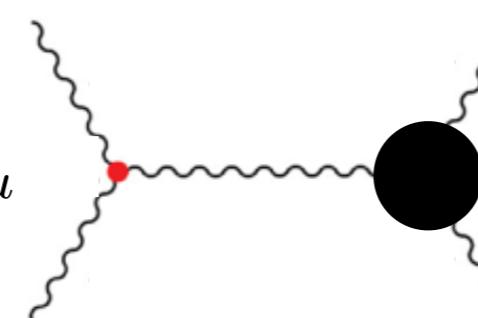
$$\sim g^2$$



$$\frac{g_*}{g} \frac{E^2}{M^2}$$

Strong vectors:

D6 $\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^{b} W^{c\rho\mu}$



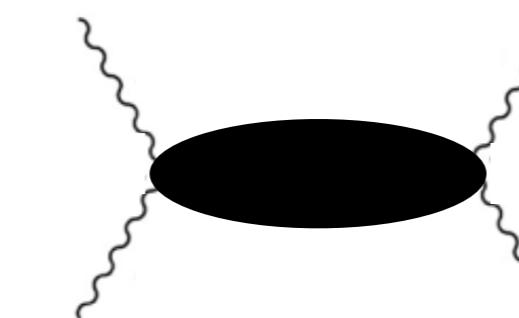
$$\sim gg_* \frac{E^2}{M^2}$$



$$\frac{g_*}{g} \frac{E^2}{M^2}$$

D8 $\frac{g_*^2}{M^4} W_{\mu\nu}^4$

(NDA: 4-point vertex = coupling²)



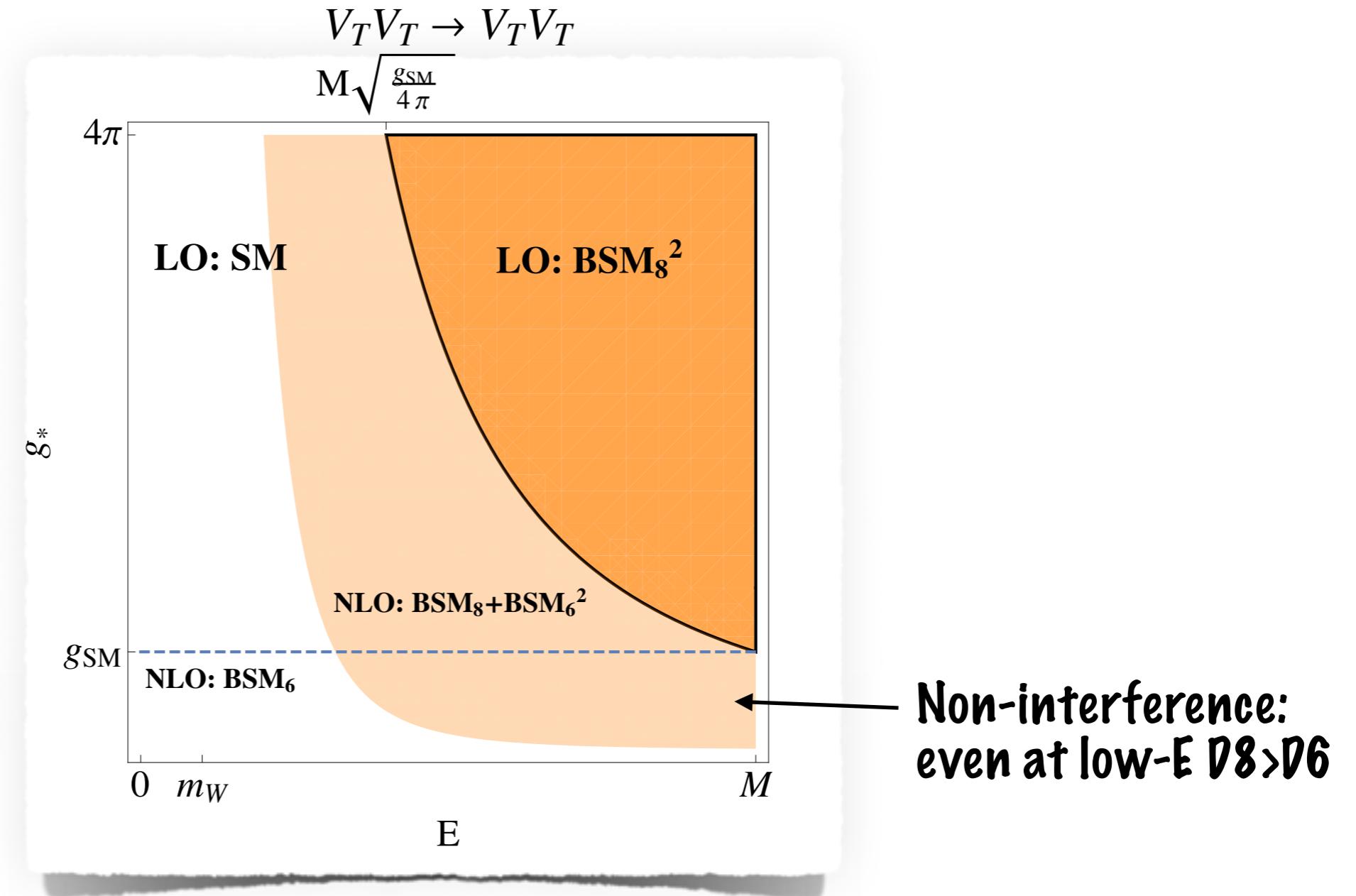
$$\sim g_*^2 \frac{E^4}{M^4}$$

► $\frac{\delta \mathcal{A}_{BSM}}{\mathcal{A}_{SM}} \gtrsim 1$ dim-6 not ok! dim-8 dominate!

(EFT E/M expansion still valid: dimension-10 small)

Strong transverse vectors: Implications

W_T scattering

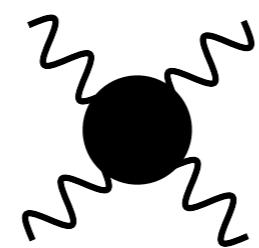


dim-8 always large for WW scattering

(EFT E/M expansion still valid: dimension-10 small)

Approximate (Non-linear) symmetries go beyond shielding the SM...

► Remedios Model:



Dim-8

$$\frac{g_*^2}{M^4} (W_{\mu\nu}^a)^4$$

Dim-6

$$\frac{gg_*}{M^2} W_\mu^{a\nu} W_{\nu\rho}^b \cancel{W^{b\mu}} \cancel{W^{a\rho}}$$

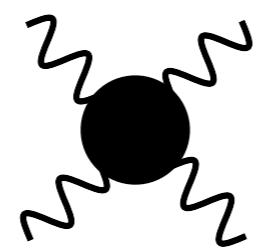
$g \ll g_*$

(dimension-10 smaller by E^2/M^2)

Dimension-8 Liu,Pomarol,Rattazzi,FR'16

Approximate (Non-linear) symmetries go beyond shielding the SM...

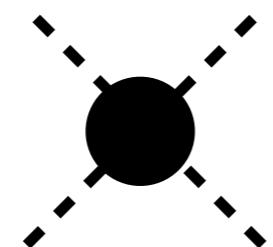
► Remedios Model:



Dim-8

$$\frac{g_*^2}{M^4} (W_{\mu\nu}^a)^4$$

► Composite Higgs:
(flat coset)



$$H \in ISO(4)/SO(4) \xrightarrow{SO(4) \rtimes T_4} H \rightarrow H + c, \\ H \rightarrow RH,$$

Proportional to
coset curvature

$$\frac{g_*^2}{M^4} (D_\mu H^\dagger D_\nu H)^2$$

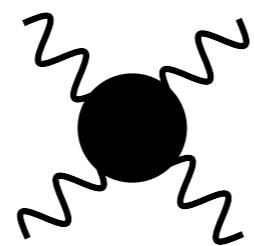
$$\epsilon_c \frac{g_*^2}{M^2} (H^\dagger D_\mu H)^2$$

$g \ll g_*$

(dimension-10 smaller by E^2/M^2)

Approximate (Non-linear) symmetries go beyond shielding the SM...

► Remedios Model:



Dim-8

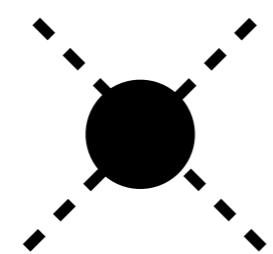
$$\frac{g_*^2}{M^4} (W_{\mu\nu}^a)^4$$

Dim-6

$$\frac{gg_*}{M^2} W_\mu^{a\nu} W_{\nu\rho}^b \cancel{W^{\rho\mu}} \cancel{W^{a\rho}}$$

$g \ll g_*$

► Composite Higgs:
(flat coset)



$$H \in ISO(4)/SO(4) \xrightarrow{SO(4) \rtimes T_4}$$

$$H \rightarrow H + c, \\ H \rightarrow RH,$$

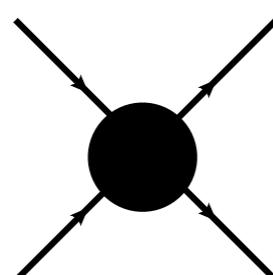
Proportional to
coset curvature

$$\frac{g_*^2}{M^4} (D_\mu H^\dagger D_\nu H)^2$$

$$\epsilon_c \frac{g_*^2}{M^2} (H^\dagger D_\mu H)^2$$

► Fermions as Pseudo-Goldstinos

Bardeen, Visnjic'82; Bellazzini,FR'soon



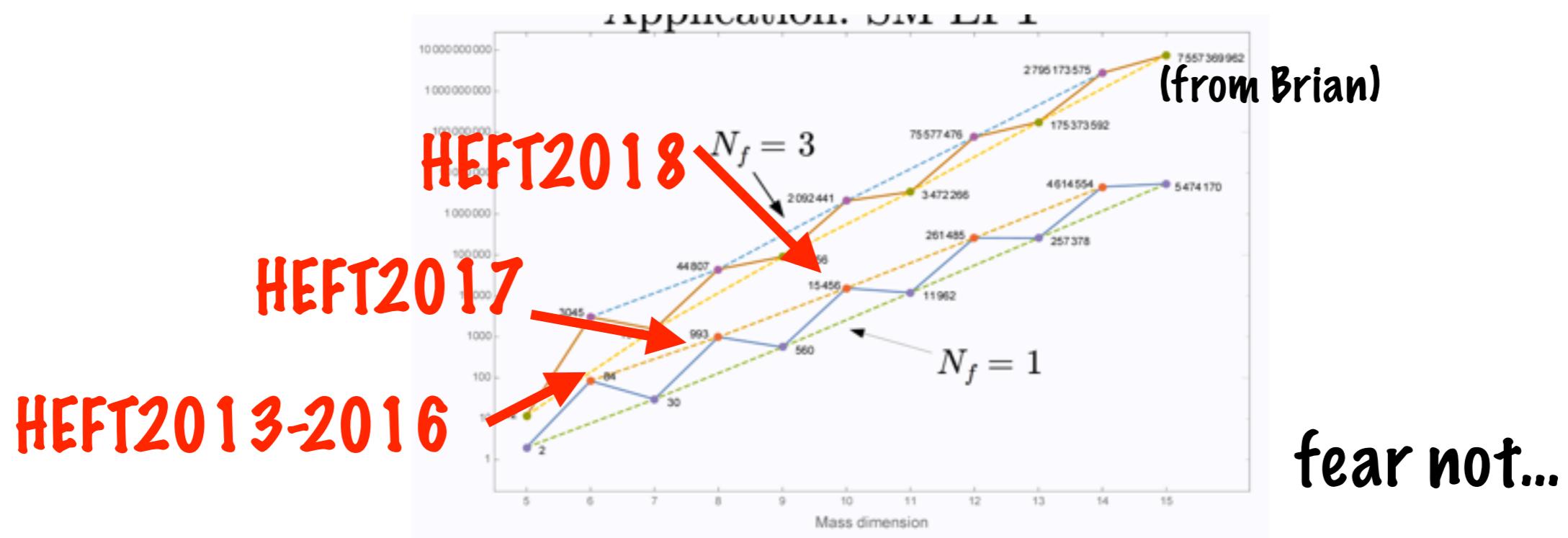
$$\delta\psi = \xi + \frac{i}{2F^2} \partial_\mu \psi (\bar{\psi} \gamma^\mu \xi - \bar{\xi} \gamma^\mu \psi)$$

$$\frac{g_*^2}{M^4} (\partial_\mu \bar{\psi} \gamma^\mu \psi)^2$$

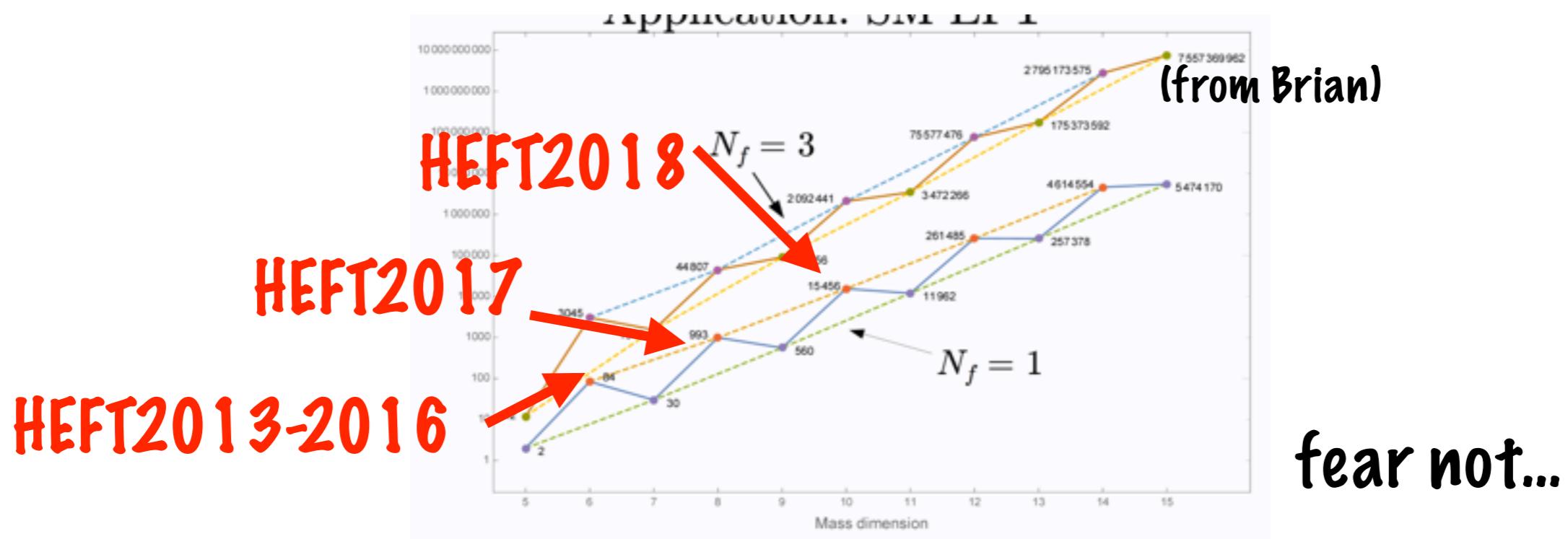
~~$$\epsilon_{\text{SUSY}} \frac{g_*^2}{M^2} (\bar{\psi} \gamma^\mu \psi)^2$$~~

(dimension-10 smaller by E^2/M^2)

HEFT 993, 15456,...

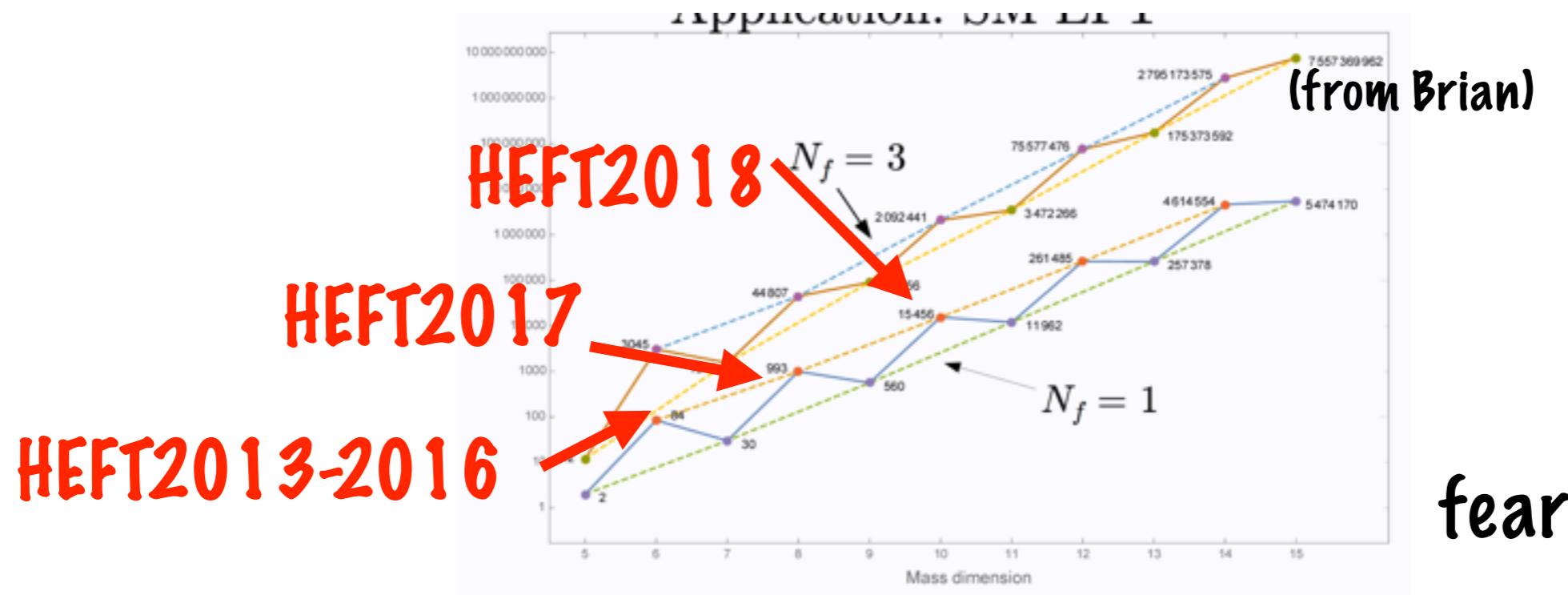


HEFT 993, 15456,...



Can there be symmetries that protect dimension-10 in $2 \rightarrow 2$ processes?

HEFT 993, 15456,...



Can there be symmetries that protect dimension-10 in $2 \rightarrow 2$ processes?

NO! Positivity constraints from analyticity, unitarity, crossing symm.:

Adams, Arkani-hamed, Dubovsky, Nicolis, Rattazzi '06;

Bellazzini '16; ...

$$c_8 > 0$$

$$\mathcal{L} = \partial_\mu \pi \partial^\mu \pi + \frac{c_8}{M^4} (\partial_\mu \pi \partial^\mu \pi)^2 + \frac{c_{10}}{M^6} (\partial_\mu \pi \partial_\nu \partial^\mu \pi)^2$$

► A symmetry that protects dim-10 can never be considered exact

Conclusions

LHC/LEP: many €€resources€€ to test transverse vectors

- ▶ Not BSM motivated (non-interference, smallness in ordinary models)

Henning,Lu,Melia,Murayama'15; Lehman,Martin'15

Conclusions

LHC/LEP: many €€resources€€ to test transverse vectors

- ▶ Not BSM motivated (non-interference, smallness in ordinary models)

Remedios: structurally robust scenario for strong coupling

- large multipoles/small monopoles
- deformed symmetry

Henning,Lu,Melia,Murayama'15; Lehman,Martin'15

Conclusions

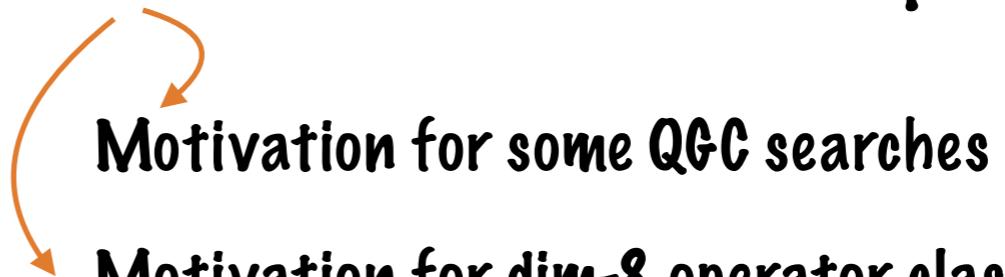
LHC/LEP: many €€resources€€ to test transverse vectors

- ▶ Not BSM motivated (non-interference, smallness in ordinary models)

Remedios: structurally robust scenario for strong coupling

- large multipoles/small monopoles
- deformed symmetry

Approximate symmetries can lead to dim-8 domination in $2 \rightarrow 2$ processes



Henning,Lu,Melia,Murayama'15; Lehman,Martin'15

- ▶ New patterns for deviations from the SM

Conclusions

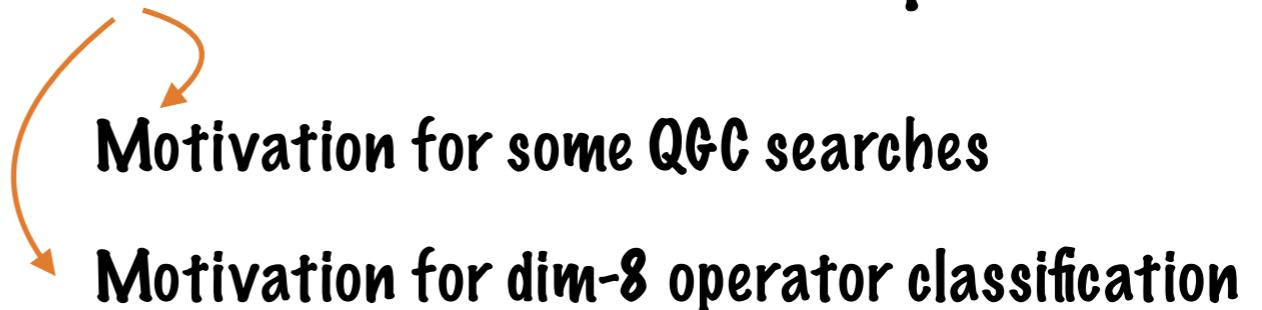
LHC/LEP: many €€resources€€ to test transverse vectors

- ▶ Not BSM motivated (non-interference, smallness in ordinary models)

Remedios: structurally robust scenario for strong coupling

- large multipoles/small monopoles
- deformed symmetry

Approximate symmetries can lead to dim-8 domination in $2 \rightarrow 2$ processes



Henning,Lu,Melia,Murayama'15; Lehman,Martin'15

- ▶ New patterns for deviations from the SM

So far the hierarchy problem has shaped our picture of TeV-physics.

Here “Data-driven”: what can we learn from well-measured quantities?

Conclusions

Precision Tests at LHC



Strong transverse vectors: UV?

Through Partial Compositeness?

$$\mathcal{L}_{mix} = \epsilon_A A_\mu J^\mu + \epsilon_F F_{\mu\nu} \mathcal{O}^{\mu\nu}$$

Marginal $\dim[J] = 3$

Relevant if $\dim[\mathcal{O}] < 2$

Unitarity $\dim[\mathcal{O}] \geq 2$

Ferrara,Gatto,Grillo'74,Mack'79

- ▶ Either free or irrelevant
- ▶ No CFT or Warped Extra-D model exists

Strong transverse vectors: UV?

Through P

Strong transverse vectors=Remedios

Remedy to motivate (some) LHC searches

No warped UV model

(Remedios the Beauty was not a creature of this world-
Gabriel Garcia Marquez)

► No C

...