

MC tools for EFT studies at the LHC

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Overview

- Guided tour of current EFT tools market - mainly Higgs
- Status of theory predictions for LHC processes
 - Higher order calculations
 - MC event generation
- EFT at the LHC
 - Contrast with anomalous coupling description from MC implementation & LHC analysis point of view
- Overview of available tools & what they can do
 - Specific production processes
 - More general model/event generation frameworks
 - Higgs decays, global fits & EFT basis translation

The goal

- Develop the required technology to make accurate & precise predictions for EFT at the LHC
- Include all processes sensitive to the Wilson coefficients of our operator expansion
- Combine run 1 & 2 data with low energy experiments into global fits
- EFT interpretation may guide is in searching for and constraining new physics at the TeV scale

How to get there

- Theory
 - Increasing precision (loops) in predictions including EFT operators
 - QFT: theory uncertainty (PDF/scale) under control
 - EFT: theory uncertainty more complex, coupled to the question of the validity of the expansion in addition to scale dependence & mixing of operators.
 - Can we be truly model independent?
- Practical
 - Tools for higher order calculations, event generation
 - Fixed order & matching/merging with parton shower
- Experimental
 - Reporting & analysis of results compatible with EFT interpretation

State of the art

- Continuing progress in higher order calculations for LHC processes
- NLO QCD is the new LO
- NNLO (N^3 LO) in QCD and/or N(N)LO EW corrections are the cutting edge
- MC tools are following suit
 - Wide array of integrator/event generation software to choose from
 - Essential components of the precision LHC programme
 - Automation at NLO in QCD: from Lagrangian to LHC events
 - Also some EW corrections

EFT at the LHC

- The LHC will continue to provide much complementary information to previous constraints
 - Probing blind directions in fits
 - Higher order predictions are needed to keep up with data & improve sensitivity
- Theory
 - Basis definitions & anomalous coupling (AC) Lagrangians (+ translations)
 - We know how to perform loop calculations & renormalise EFT

[R. Alonso, E. Jenkins, A. Manohar & M. Trott;
JHEP 1310 (2013) 087, JHEP 1401 (2014) 035 & JHEP 1404 (2014) 159*]*
- Tools
 - Increasing numbers of predictions including anomalous coupling/EFT exist
 - A great deal of experimental interest in expanding these

EFT & AC: implementation

- Anomalous coupling Lagrangians start out in the mass eigenbasis of SM fields e.g. after EWSB
 - Physically intuitive & can directly calculate/implement MC generators
 - $SU(3) \times U(1)_Q$ gauge symmetry only
 - Expand in canonical dimension, derivatives,...
- EFT expansions are the mapping from UV theories
 - Linear: Higgs is a doublet, $SU(2)_L$ symmetry also present
 - Non-linear: Higgs as a singlet + chiral EW Lagrangian
- Must first break EW symmetry and map to the mass eigenbasis in order to make predictions for the LHC
 - Generally a one-way map $EFT \rightarrow AC$ & parameters will be (cor)related

EFT & AC: map

$$\begin{aligned}\mathcal{O}_H &= \frac{\bar{c}_H}{2} \partial_\mu (\Phi^\dagger \Phi) \partial^\mu (\Phi^\dagger \Phi) \\ &= \frac{\bar{c}_H}{\Lambda^2} \frac{v^2}{2} \partial_\mu h \partial^\mu h + \mathcal{O}(h^3, h^2) \\ h &\rightarrow h(1 + \delta h), \quad \delta h = -\frac{\bar{c}_H}{\Lambda^2} \frac{v^2}{4}\end{aligned}$$

$$\begin{aligned}\mathcal{O}_W|_{\Phi=\langle\Phi\rangle} &= \frac{ig}{2} \bar{c}_W \left[\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi \right] D^\nu W_{\mu\nu}^k|_{\Phi=\langle\Phi\rangle} \\ &= \frac{gv^2}{16} \bar{c}_W \left[2gW_+^{\mu\nu} W_{\mu\nu}^- + g(W_3^{\mu\nu} - g' B^{\mu\nu}) W_{\mu\nu}^3 \right] + \text{aGC} \\ W_\pm^\mu &\rightarrow W_\pm^\mu [1 + \delta W] \\ B^\mu &\rightarrow B^\mu [1 + \delta B] + yW_3^\mu \\ W_3^\mu &\rightarrow W_3^\mu [1 + \delta W] + zB^\mu\end{aligned}$$

$$y + z \equiv -\frac{v^2}{\Lambda^2} \frac{gg'}{4} \left(\bar{c}_B + \frac{\bar{c}_W}{2} \right)$$

- Canonical mass eigenbasis Lagrangian generically different from the SM case
 - Non canonical kinetic terms & kinetic mixing
 - Infinite freedom in how one redefines fields (equivalent up to $1/\Lambda^2$)
 - EFT contributions to masses and couplings
- Some extra book-keeping required...

EFT & AC: map

- ‘SM-like’ EW sector is modified
 - New relations between SM inputs and EW parameters
 - $v(G_F, c_i)$, $\sin\theta_W(m_W, m_Z, \alpha_{EM}, c_i)$, ...
 - Higgs couplings rescaled
 - Z couplings to fermions
 - SM gauge bosons self interactions
 - Many of these effects are constrained pre-LHC (i.e. by EWPO)
- Additional $1/\Lambda^2$ pieces on top of direct effects from new operators/Lorentz structures
 - Depends on SM input scheme used
 - Can contribute to SM backgrounds!

[Trott & Passarino; LHCHXSWG-DRAFT-2016-005]

[Falkowski et.al; LHCHXSWG-INT-2015-001]

[Williams, KM & Sanz; JHEP 1608 (2016) 039]

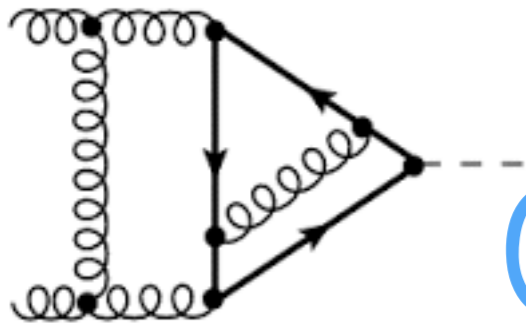
[Ge, He & Xiao; JHEP 1610 (2016) 007]

[Degrande, Fuks, Mawatari, KM, Sanz; 1609.04833]

EFT & AC: interpretation

- Tree-level effects
 - Important to ensure it is consistent before moving to higher orders in PT
- AC implementations do not need to consider such things
 - Potentially complicates the ‘one-way map’
 - There is a way to at least preserve SM input relations (Higgs Basis)
[A. Falkowski et.al; LHCHSWG-INT-2015-001]
- If AC are used to simulate signal & interpret data
 - From a pure Lorentz structure basis, the EFT parameter space is a subset of the general AC description
 - Mapping constraints from a multi parameter AC fit to an EFT interpretation may be tricky & relevant parameter space is larger than it originally seems
 - e.g. $H \rightarrow WW$: might also need to include rescalings of the WW background
 - Assess how big these effects can be

Tools for LHC processes



Gluon-gluon fusion

$$\mathcal{L}_{\text{SILH}} = y_u \frac{g_s^2 \bar{c}_g}{m_W^2} |\Phi|^2 - y_u \frac{\bar{c}_u}{v^2} |\Phi|^2 \bar{Q}_L \tilde{\Phi} u_R + \frac{\bar{c}_H}{2v^2} \partial^\mu |\Phi|^2 \partial_\mu |\Phi|^2 - \lambda \frac{\bar{c}_6}{v^2} |\Phi|^6$$

$$\mathcal{L}_{\text{eff.}} = -m_t \bar{t}t \left(c_t \frac{h}{v} + c_{tt} \frac{h^2}{2v^2} \right) - \frac{m_h^2}{2v} c_3 h^3 + \frac{\alpha_S}{\pi} G^{\mu\nu} G_{\mu\nu} \left(c_g \frac{h}{v} + c_{gg} \frac{h^2}{2v^2} \right)$$

- Much effort had been invested into this process
- EFT: Modified Higgs, top & gluon interactions
 - Contribution from chromomagnetic operator often not included
 - Arises at one-loop in perturbative UV completions
 - First calculated in: *[C. Degrande et al.; JHEP 1207 (2012) 036]*
 - Included in implementation of:

[F. Maltoni, E. Vryonidou, C. Zhang, arXiv:1607.05330]

Single Higgs

- HiGlu

[M. Spira; arXiv:hep-ph/9510347]

<http://tiger.web.psi.ch/higlu>

- QCD: NNLO for SM
- Quark mass effects at NLO
- NLO EW & NNLO EW/QCD in the SM
- EFT: rescaling of top yukawa (NNLO QCD) & c_g (NNLL scale dep. only)

- SusHi

[R. Harlander, S. Liebler, H. Mantler; arXiv:1605.03190]

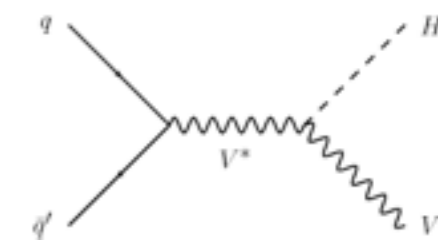
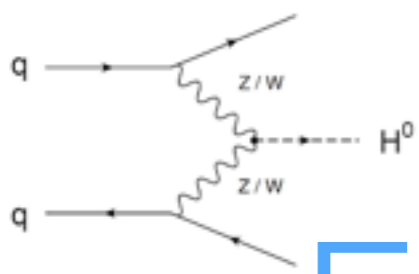
<http://sushi.hepforge.org>

- QCD: NNLO, N³LO in heavy top limit + $1/m_t$ expansion
- Recently included Dim-5 gg operator c_g (no $1/m_t$ expansion)
- EFT: now same as HiGlu
- Event generation: aMCSusHi - MadGraph5_aMC@NLO interface (no c_g)

Double Higgs

- HPAIR *[S. Dawson, M. Dittmaier, M. Spira; Phys. Rev. D58:115012]*
<http://tiger.web.psi.ch/hpair>
 - NLO in QCD in either SM or MSSM (heavy top limit)
 - Events: HiggsPair - HERWIG++ <https://herwig.hepforge.org>
- EFT contribution
 - 5-parameter anomalous coupling Lagrangian
 - Alternative ‘SILH’ basis input maps to the anomalous couplings
 - No dipole operator
 - Also LO events: HiggsPair in HERWIG 7.0 *[F. Goertz et al.; JHEP 1504 (2015) 167]*

$$c_t = 1 - \frac{\bar{c}_H}{2} - \bar{c}_u, \quad c_{tt} = -\frac{1}{2}(\bar{c}_H - 3\bar{c}_u), \quad c_3 = 1 - \frac{3\bar{c}_H}{2} + \bar{c}_6, \quad c_g = c_{gg} = \frac{4\pi}{\alpha_2}\bar{c}_g$$



EW production: VH, VBF

$$\mathcal{L}_\Phi = \frac{\bar{c}_H}{\Lambda^2} \partial^\mu |\Phi|^2 \partial_\mu |\Phi|^2 + \frac{\bar{c}_T}{\Lambda^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] - \lambda \frac{\bar{c}_6}{\Lambda^2} |\Phi|^6$$

$$\mathcal{L}_G = \frac{g'^2}{4\Lambda^2} \bar{c}_{BB} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{ig}{2\Lambda^2} \bar{c}_W [\Phi^\dagger T_{2k} \overleftrightarrow{D}_\mu \Phi] D_\nu W^{k,\mu\nu} + \frac{ig'}{2\Lambda^2} \bar{c}_B [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] \partial_\nu B^{\mu\nu}$$

$$+ \frac{ig}{\Lambda^2} \bar{c}_{HW} [D_\mu \Phi^\dagger T_{2k} D_\nu \Phi] W^{k,\mu\nu} + \frac{ig'}{\Lambda^2} \bar{c}_{HB} [D_\mu \Phi^\dagger D_\nu \Phi] B^{\mu\nu}$$

$$\mathcal{L}_C = \frac{i\bar{c}_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu Q_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\bar{c}'_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu T_{2k} Q_L] [\Phi^\dagger T^{2k} \overleftrightarrow{D}_\mu \Phi]$$

$$+ \frac{i\bar{c}_{Hu}}{v^2} [\bar{u}_R \gamma^\mu u_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{i\bar{c}_{Hd}}{v^2} [\bar{d}_R \gamma^\mu Q_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] - \left[\frac{i\bar{c}_{Hud}}{v^2} [\bar{u}_R \gamma^\mu Q_R] [\tilde{\Phi}^\dagger \overleftrightarrow{D}_\mu \Phi] + \text{h.c.} \right]$$

- Many more operators can contribute
 - Operators also affect non Higgs rates i.e. backgrounds
 - “Classic” new V-V-h Lorentz structure
 - V-h-f-f Contact interactions

EW production: VH, VBF

- The usually considered bunch that maps to the familiar “L3” parametrisation of anomalous Higgs couplings to gauge bosons
 - Deviations in the tails of distributions probing large momentum flow in the production vertex
 - “Easy” for QCD corrections as they factorise to the initial state quark lines

$$\mathcal{L}_G = \frac{g'^2}{4\Lambda^2} \bar{c}_{BB} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{ig}{2\Lambda^2} \bar{c}_W [\Phi^\dagger T_{2k} \overleftrightarrow{D}_\mu \Phi] D_\nu W^{k,\mu\nu} + \frac{ig'}{2\Lambda^2} \bar{c}_B [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] \partial_\nu B^{\mu\nu}$$

$$+ \frac{ig}{\Lambda^2} \bar{c}_{HW} [D_\mu \Phi^\dagger T_{2k} D_\nu \Phi] W^{k,\mu\nu} + \frac{ig'}{\Lambda^2} \bar{c}_{HB} [D_\mu \Phi^\dagger D_\nu \Phi] B^{\mu\nu} + \text{CP violating}$$

$$\mathcal{L}_{HAC} = -\frac{1}{4} g_{hzz}^{(1)} Z_{\mu\nu} Z^{\mu\nu} h - g_{hzz}^{(2)} Z_\nu \partial_\mu Z^{\mu\nu} h + \frac{1}{2} g_{hzz}^{(3)} Z_\mu Z^\mu h - \frac{1}{4} \tilde{g}_{hzz} Z_{\mu\nu} \tilde{Z}^{\mu\nu} h$$

$$- \frac{1}{2} g_{hww}^{(1)} W^{\mu\nu} W_{\mu\nu}^\dagger h - \left[g_{hww}^{(2)} W^\nu \partial^\mu W_{\mu\nu}^\dagger h + \text{h.c.} \right] + g_{hww}^{(3)} W_\mu W^{\dagger\mu} h - \frac{1}{2} \tilde{g}_{hww} W^{\mu\nu} \tilde{W}_{\mu\nu}^\dagger h$$

$$- \frac{1}{2} g_{haz}^{(1)} Z_{\mu\nu} F^{\mu\nu} h - g_{haz}^{(2)} Z_\nu \partial_\mu F^{\mu\nu} h - \frac{1}{2} \tilde{g}_{haz} Z_{\mu\nu} \tilde{F}^{\mu\nu} h$$

EW production: VH, VBF

[A. Denner et al.; JHEP 1203 (2012) 075]

- HAWK

<http://omnibus.uni-freiburg.de/~sd565/programs/hawk/hawk>

- VH & VBF
- NLO in EW/QCD for the SM
- Anomalous coupling NLO in QCD, 2 parameter subset given others constrained by LEP

[J. Baglio et al.; arXiv:1404.3940]

- VBFNLO

<https://www.itp.kit.edu/vbfnlo>

- General purpose tool for Higgs and weak boson production
- NLO QCD (LO event generation)
- Focus on VBF topologies, many processes available:
- H+X, V+X, HH+X, VV+X, VVV+X (also WH)
- Full L3 AC parametrisation

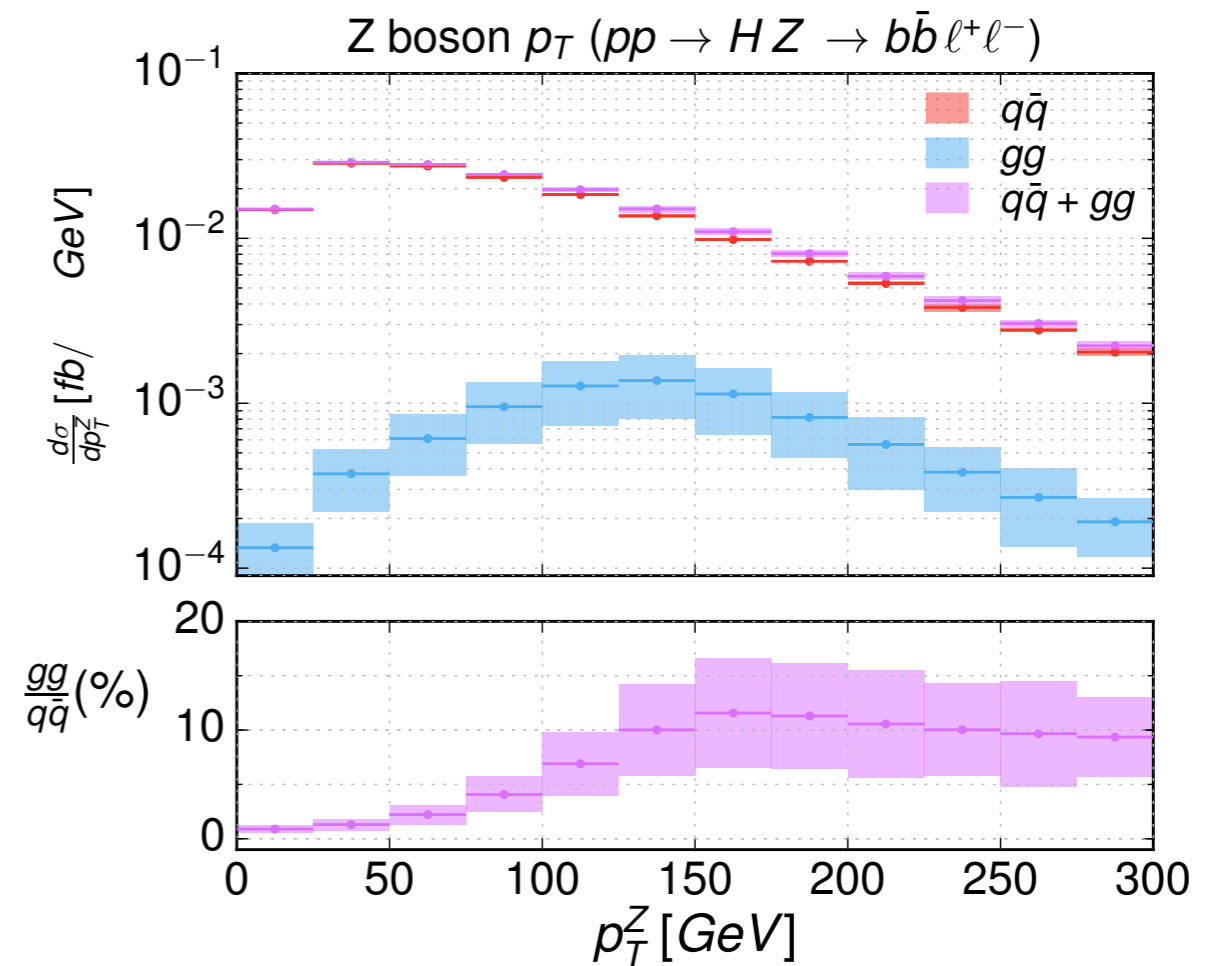
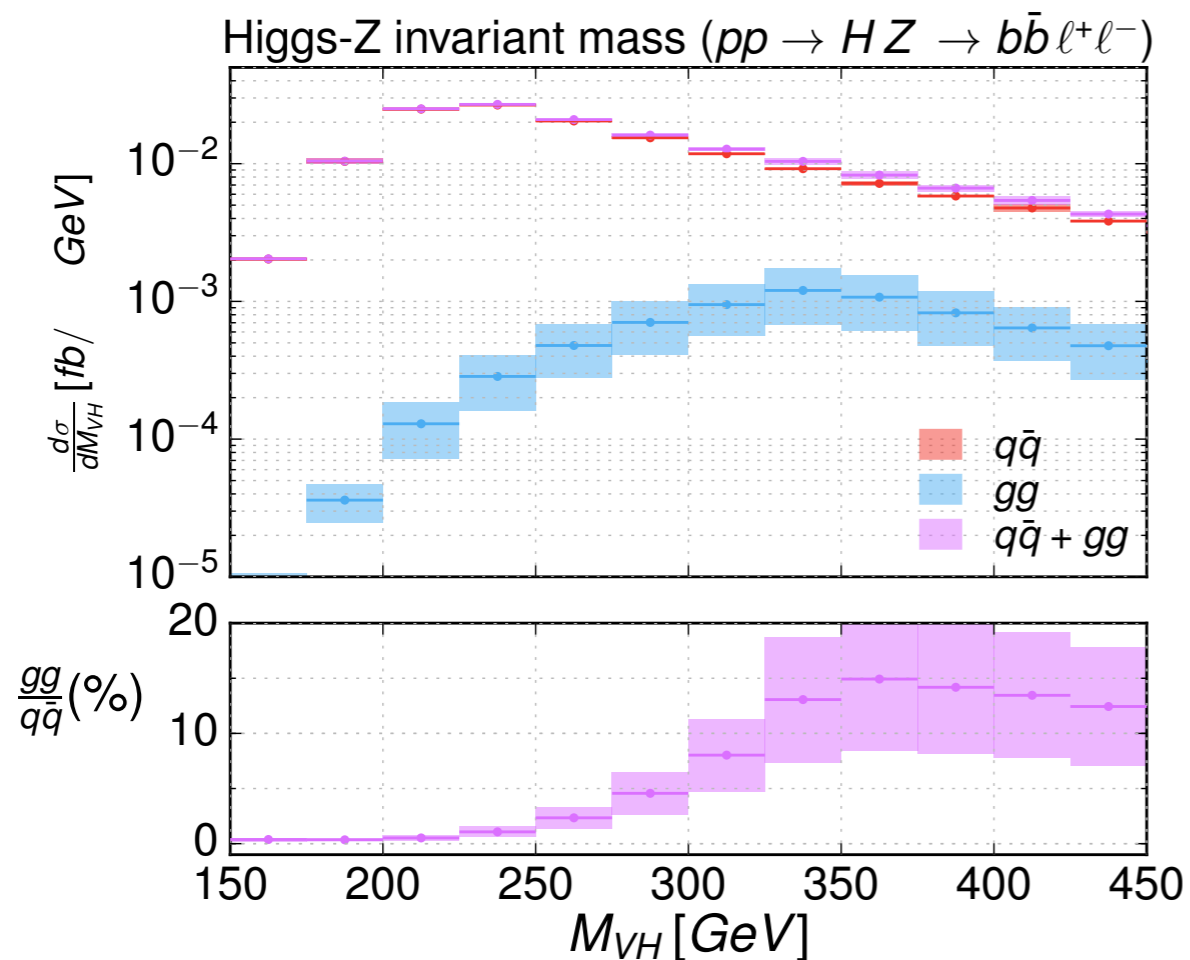
POWHEG-BOX/MCFM

- Implementation of WH/ZH in POWHEG-BOX framework
 - NLO in QCD and matched to parton shower
 - Built from MCFM matrix elements for the SM processes
 - $gg \rightarrow ZH$ contribution should be added separately, widely available
- EFT contribution
 - Map from “gauge” part of SILH to anomalous couplings (5 operators)
 - Also includes CP-violating operators
 - Relevant modifications to EW parameter relations taken into account
- In the paper
 - gg -contribution, K-factors
 - Distributions for benchmarks based on current constraints & UV models

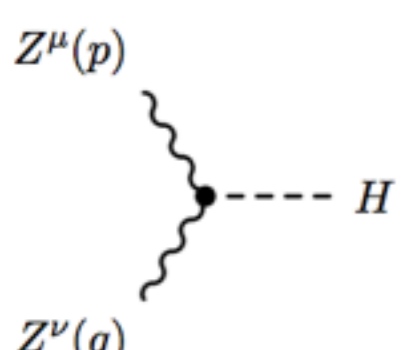
[Gorbahn, No & Sanz; JHEP 1510 (2015) 036]

$gg \rightarrow ZH \rightarrow l^+l^- bb$

- gg initiated process (formally NNLO)
 - Gluon PDF plus kinematics of EFT searches warrant its inclusion
 - Well known to ‘mimic’ EFT effects if not properly taken into account



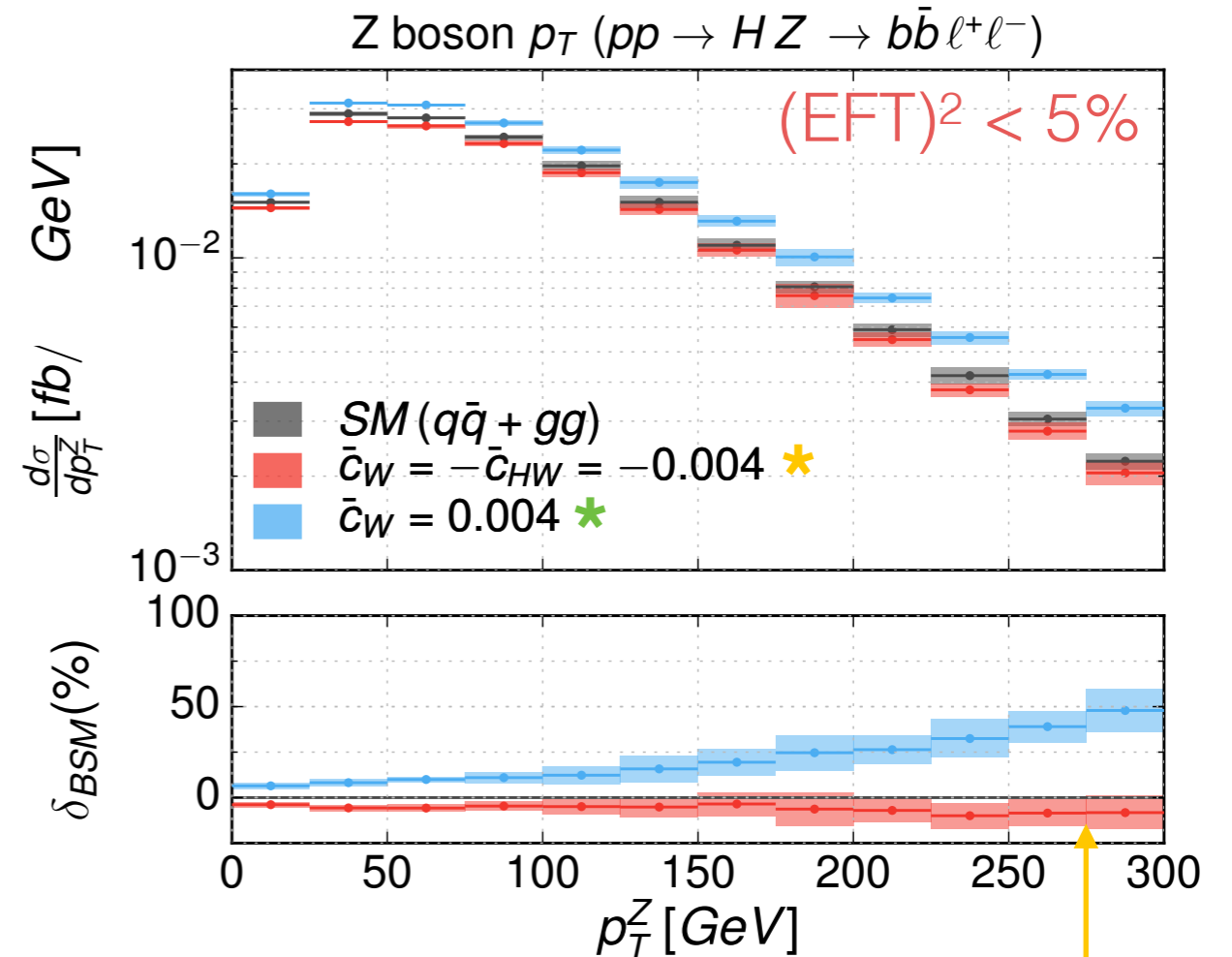
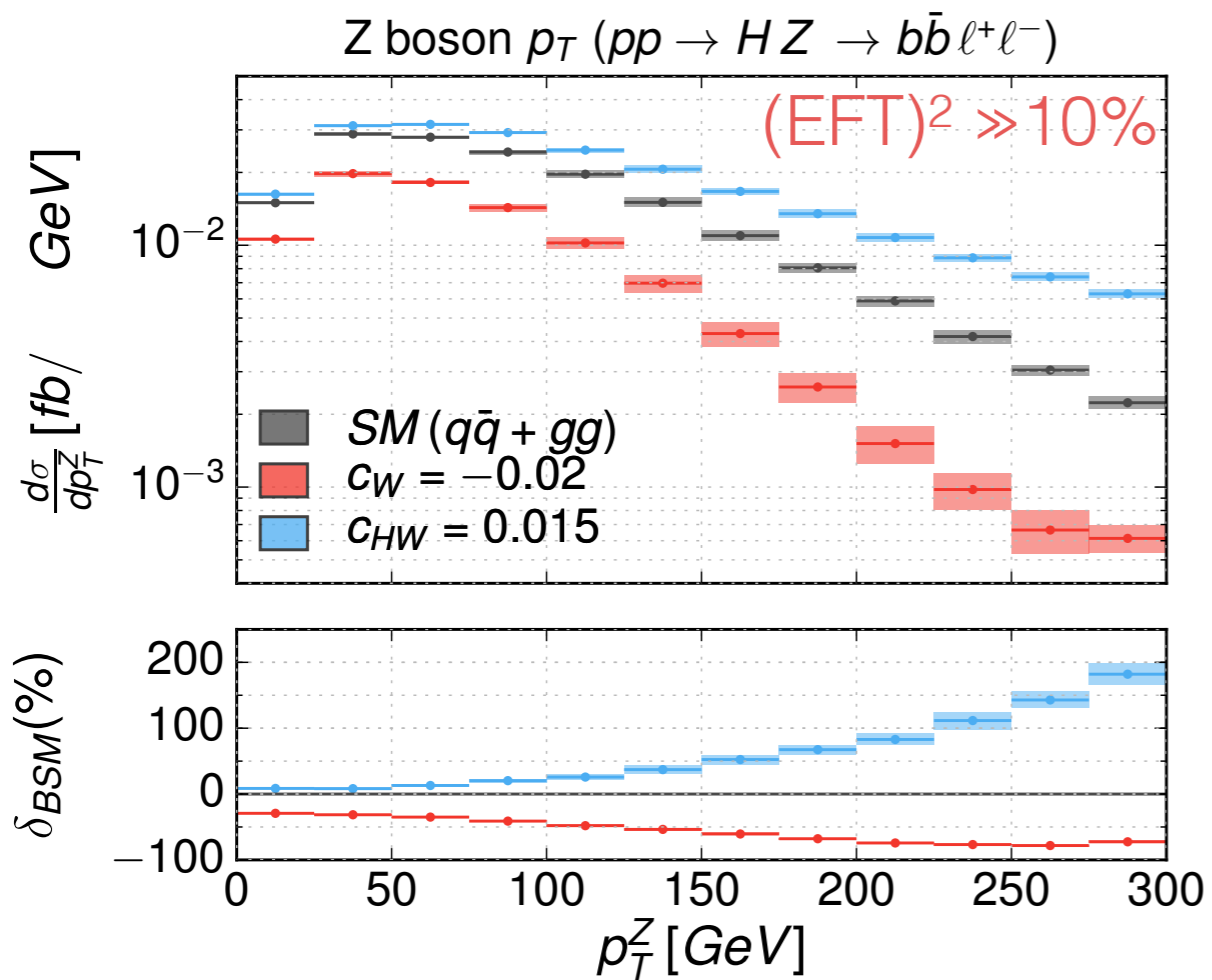
$pp \rightarrow ZH \rightarrow l^+l^- bb$

$Z^\mu(p)$

 $Z^\nu(q)$

$$i \left[\frac{g}{\cos \theta_W} M_Z + g_{hzz}^{(1)} (\eta^{\mu\nu} p \cdot q - q^\mu p^\nu) + g_{hzz}^{(2)} ((p^2 + q^2)\eta^{\mu\nu} - p^\mu p^\nu + q^\mu q^\nu) \right]$$

$g_{hzz}^{(1)} \propto \bar{c}_{HW}, \quad g_{hzz}^{(2)} \propto (\bar{c}_{HW} + \bar{c}_W)$

“BM B”
“BM A”



* Benchmark B does not show “EFT-like” features

EW production: VH, VBF

- Not commonly considered contact interactions
 - Also affect V-f-f vertices: constrained by previous experiments e.g. LEP
 - Not all of them much more than LHC constraints on $(c_W, c_B, c_{HW}, c_{HB})$
 - Flavour matrices: FCNC, ... interplay between LHC and other experiments
 - Coloured particles: non-trivial QCD corrections
 - Curious to see how much they can impact VH, VBF at the LHC given existing constraints or conversely, whether LHC can be complementary

$$\mathcal{L}_C = \frac{i\bar{c}_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu Q_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\bar{c}'_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu T_{2k} Q_L] [\Phi^\dagger T^{2k} \overleftrightarrow{D}_\mu \Phi]$$

$$+ \frac{i\bar{c}_{Hu}}{v^2} [\bar{u}_R \gamma^\mu u_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{i\bar{c}_{Hd}}{v^2} [\bar{d}_R \gamma^\mu Q_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] - \left[\frac{i\bar{c}_{Hud}}{v^2} [\bar{u}_R \gamma^\mu Q_R] [\tilde{\Phi}^\dagger \overleftrightarrow{D}_\mu \Phi] + \text{h.c.} \right]$$

+ lepton currents (some of which affect SM inputs)

Only HiggsPO includes these so far: more work required!

MC event generation frameworks

Frameworks

- WHIZARD *[W. Kilian, T. Ohl, J.Reuter; Eur.Phys.J. C71, 1742]*
<http://whizard.hepforge.org>
 - General purpose LO event generator w/ internal parton shower
 - Possibility to interface with loop-level matrix elements
 - EFT operators involving tops/higgs/gauge bosons

- JHU generator & MELA <http://spin.pha.jhu.edu>
 - Dedicated LO event generator for Higgs production/characterisation
 - $pp \rightarrow X \rightarrow VV, VBF, X+JJ, pp \rightarrow VX, ee \rightarrow VX$
 - Matrix Element Likelihood Analysis package for sophisticated signal optimisation techniques

- VBFNLO could have been in this section too

Frameworks

- FeynRules/Madgraph5_aMC@NLO suite
 - Most popular event generation framework <https://launchpad.net/mg5amcnlo>
 - NLO in QCD + PS, loop-induced processes
- Interfaces to many useful particle physics software
 - FeynRules + NLOCT for “automated” generation of NLOQCD-ready model file from Lagrangian: compatible with many MC event generators
 - PYTHIA/HERWIG for parton shower
 - Delphes for fast detector simulation
 - SysCalc, MadDM, Madspin,...
- Brings NLO QCD to the masses!
 - EFT community has been making great use of this suite
 - NLO calculations/implementations owe much to the EFT-RGE calculation

Higgs Characterisation

- Anomalous coupling Lagrangian for Higgs couplings
 - General Lorentz structures in EW gauge-Higgs sector
 - Gluon-gluon effective vertex (LO)
 - Currently used by experimental collaborations for some Higgs analyses
- UFO (FeynRules) model at NLO in QCD
 - VBF, VH + CP-violating
- Related model: BSM Characterisation (LO)
 - Extends to a larger set of operators: all except four fermion
 - Based on collection of Lorentz structure in Higgs Basis definition
 - Output for ROSETTA: basis translator program which implements maps from popular EFT basis choices to the mass eigenbasis

HiggsPO

- Alternative but related framework for Higgs couplings
 - Lorentz structures weighted by general form factors
 - Defined at the amplitude level
 - Momentum expansion around poles of gauge bosons
- UFO model at NLO in QCD
 - Higgs decays to 4 leptons,..
 - EW Higgs production
 - Including $hVff$ contact terms

*[M. Gonzalez-Alonso et al.; Eur.Phys.J. C75 (2015) 128
& Eur.Phys.J. C75 (2015) 341]*

[M. Bordone et al.; Eur.Phys.J. C75 (2015) 8, 385]

[A. Greljo et al.; Eur.Phys.J. C76 (2016) 3, 158]

Higgs Effective Lagrangian

- LO FeynRules model for the ‘SILH’ basis of dim-6 EFT
 - 39 operators included (no four fermion operator)
 - Effect of field redefinitions and SM inputs included
 - Largest collection of dim-6 EFT operators in one place

[A. Alloul, B. Fuks & V. Sanz; JHEP 1404 (2014) 110]

<http://feynrules.irmp.ucl.ac.be/wiki/HEL>

- Upgrade to NLO QCD has begun
 - First results for EW Higgs production: VH & VBF
 - So far, only 5 operators that map to the L3 parametrisation
 - These require no new counterterms for QCD corrections
 - Cross checked VH with POWHEG-BOX/MCFM implementation

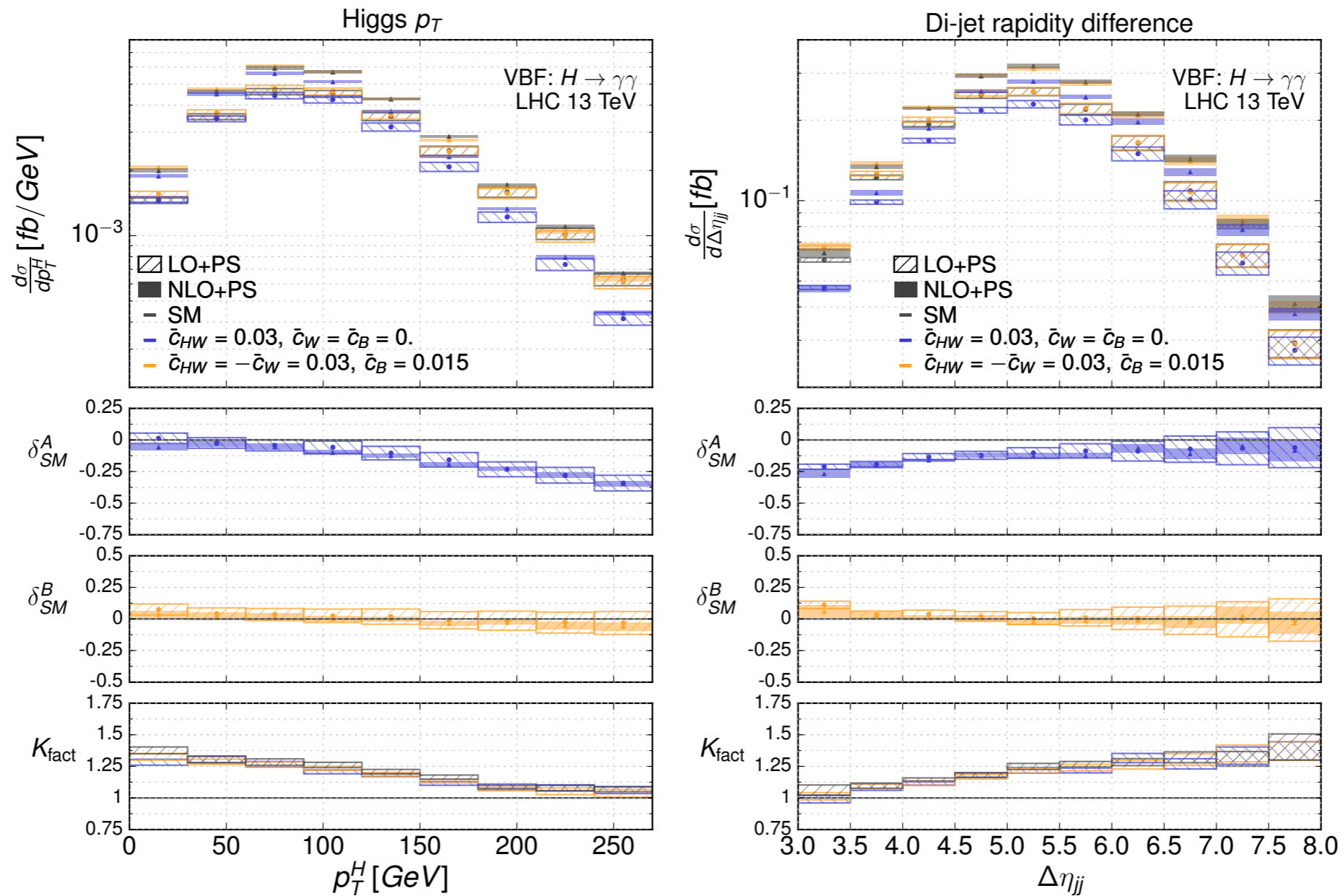
[C. Degrande, B. Fuks, K. Mawatari, KM & V. Sanz; arXiv:1609.04833]

<http://feynrules.irmp.ucl.ac.be/wiki/HELatNLO>

$pp \rightarrow H jj \rightarrow \gamma\gamma jj$

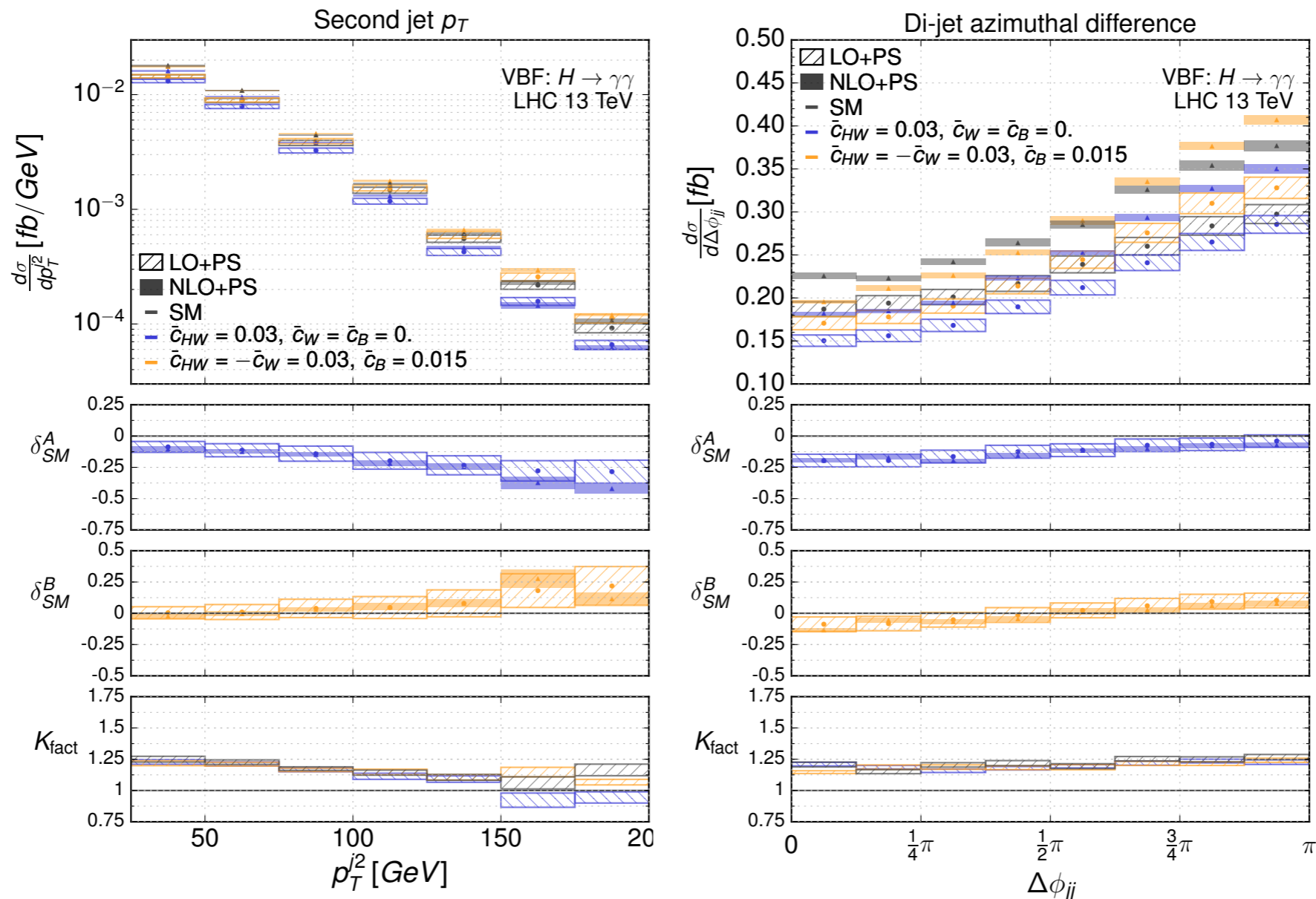
generate p p > h j j \$\$ w+ w- z a QCD=0 [QCD]

No ggF but VBF cuts, fixed scale of m_W



$$pp \rightarrow H jj \rightarrow \gamma\gamma jj$$

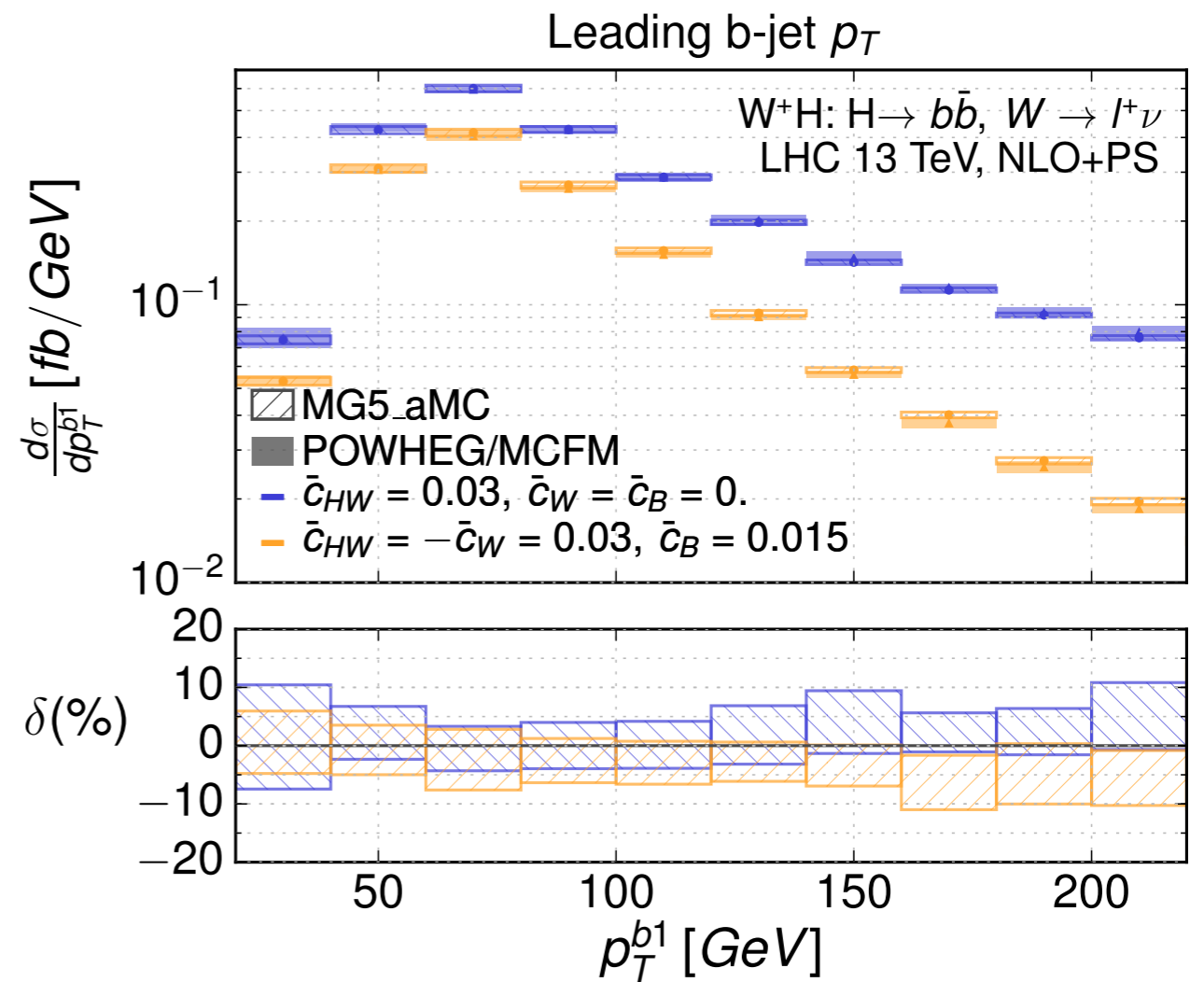
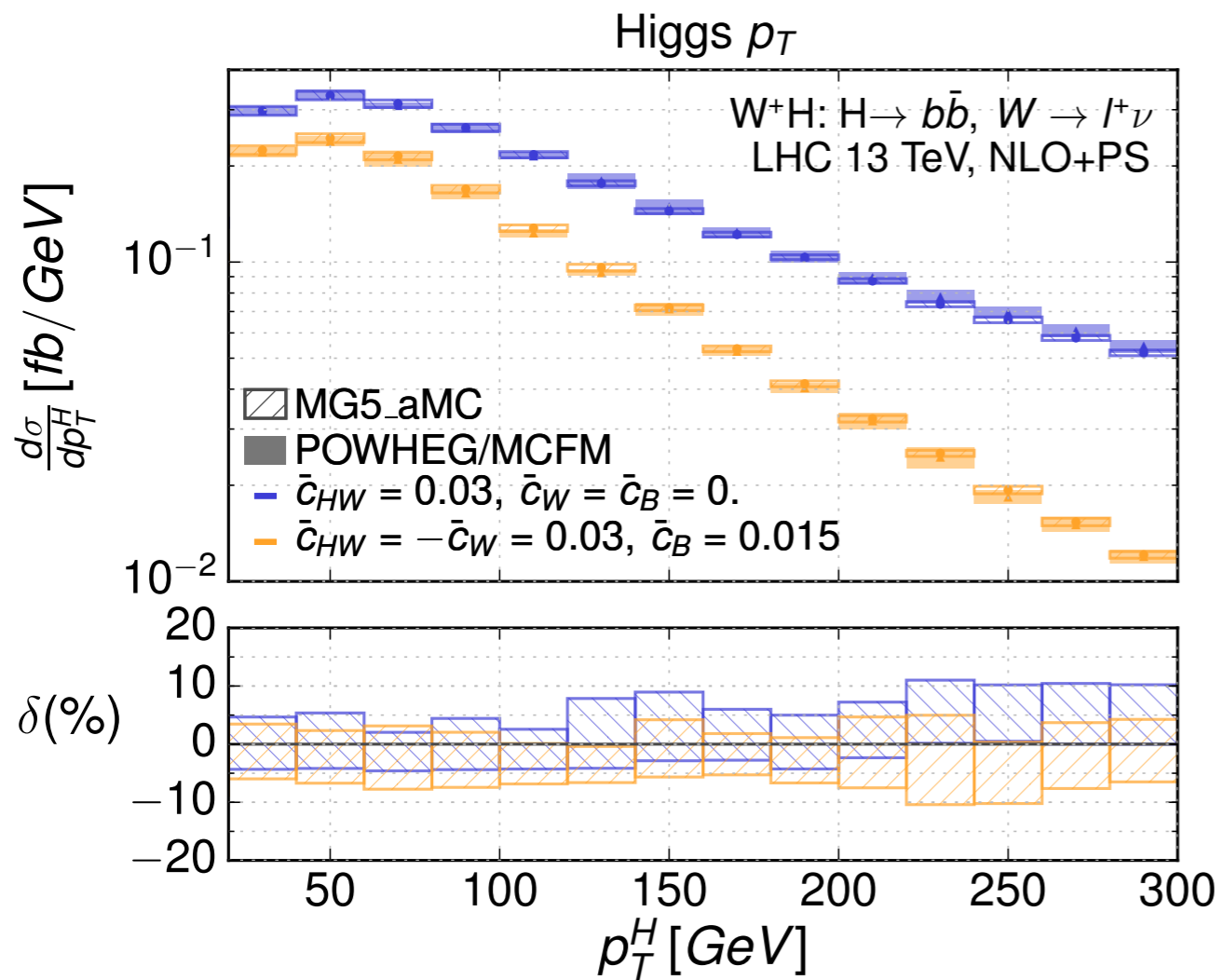
Generally effects of order 25-50% present, mild sensitivity to “benchmark B”. Correlating VH & VBF may help disentangle this coupling structure.



Reduced scale uncert. improves prospects of distinguishing ‘UV motivated’ benchmark B

Validation

Reasonable agreement between POWHEG/MCFM & Madgraph5_aMC@NLO implementations



Top/Higgs EFT@NLO

- Cen's talk yesterday (talk to him to request models?)
- NLO model implementations of dim-6 operators relevant for:
 - tt (e^+e^- , pp), single top, ttH, ttV
 - [C. Degrande et al.; PRD 91 (2015) 034024]*
 - [G. Durieux et al.; PRD 91 (2015) no.7, 074017]*
 - [O. Bylund et al.; JHEP 1605 (2016) 052]*
 - [C. Zhang; PRL 116 (2016) no.16, 162002]*
 - [F. Maltoni et al.; arXiv:1607.05330]*
 - H, HH, Hj
 - top FCNC
- Coloured operators: non trivial QCD corrections
 - Compared RGE-improved vs. full NLO
 - Incorporate a new theory error contribution based on EFT scale dependence
 - Doesn't cancel in, e.g., cross section ratios as opposed to usual PDF, scale uncertainties
- Big step towards global EFT fit at NLO QCD accuracy

Others

Higgs decays: eHDECAY

<https://www.itp.kit.edu/~maggie/eHDECAY/>

- Extension of HDECAY
*[A. Djouadi, J. Kalinowski, M. Spira, *Comp. Phys. Comm.* 108 (1998) 56]*
- QCD & EW corrections to Higgs BRs
- EFT contribution truncated at $(1/\Lambda)^2$
 - Anomalous coupling (“non-linear”) Lagrangian
 - Alternative ‘SILH’ basis input maps to the anomalous couplings
 - SILH input also includes some NLO EW corrections
 - Unclear if modifications to EW parameters from SM inputs included
- Only publicly available tool
 - ROSETTA interface to this program exists for other basis choices
 - See Darren’s talk for improved calculation in some final states

Fitting

- There are many active EFT fitting collaborations
- LHC Higgs, EWPO, top, Flavour & combinations thereof
- So far, mostly LO EFT predictions used
 - The next generation will at least employ NLO QCD
 - Allow for study of theory uncertainties
- Mostly private theorist codes & experimental analyses
- One public tool: HEPfit
 - Statistical code for combination of direct/indirect constraints on new physics from Flavour, EWPO, Higgs signal strengths,...
 - Allows for definition of general model & effects on observables of interest
 - Implements EW chiral Lagrangian <http://hepfit.roma1.infn.it>

Rosetta

<http://rosetta.hepforge.org>

Version 2.0 just out!
Original definitions of
Warsaw & SILH bases

- Operator basis translation tool for SMEFT
 - Encodes relations between EFT “bases”
 - SILH, Warsaw, Higgs, HISZ,...
 - Special BSMC Lagrangian output for FeynRules model
- Opens up user base for basis specific codes
- Modular + extendable: user defined bases
- Existing interfaces: eHDECAY, Lilith (Higgs rates), EWPO
- Future
 - “Basis independent” implementation of RGEs
 - Output for HC (caveats..) and HELatNLO models

Outlook

- Constant progress in higher order calculations & MC tools
 - Anomalous couplings descriptions abound
 - First theoretical results at NLO in EFT
 - EFT implementations via automation of NLO QCD
- We are heading towards a complete basis at NLO in QCD
 - Top sector: almost there
 - Higgs sector: a little way to go
- Next generation of global fits will hopefully incorporate there

Thoughts

- When experiments employ anomalous coupling description, careful when mapping to EFT interpretation
 - At this point its a practical matter for experiments to use AC descriptions
- We should continue discussion on definitions of EFT theory errors & effect on constraints
 - Missing higher orders, amplitudes squared, EFT scale
 - EFT “validity” might be considered as a part of the theory uncertainty
 - Blows up as EFT is approaches its limits for a given observable?
 - Can MC tools incorporate these effects?

End of the tour.
Thank you

Backup

HiGlu

- Single Higgs production cross section via gluon fusion
 - NNLO in QCD in either SM or MSSM
 - Includes quark mass effects at NLO
 - CP-even or -odd scalar
 - Some differential distributions
 - SM NLO EW & NNLO QCD/EW contributions
- EFT contribution
 - Rescaling of top Yukawa, c_t
 - Non-zero dim-5 gluon-gluon operator, c_g

$$\mathcal{L}_{\text{eff.}} = \frac{\alpha_S}{\pi} \left(\frac{c_t}{12} (1 + \delta) + c_g \right) G^{\mu\nu} G_{\mu\nu} \frac{h}{v}, \quad (\text{or CP-odd})$$

NNLO in QCD

NNLL scale dependence only

SusHi

- Single Higgs production cross section via gluon fusion
 - Recently included dim-5 gluon-gluon operator
 - Quark mass effects also at NLO
 - NNLO, N³LO in heavy top limit + 1/m_t expansion
 - CP-even or -odd scalar
- EFT contribution
 - Same as HiGlu
 - 1/m_t expansion not available when including c_g
- aMCSusHi interface for MadGraph5_aMC@NLO (no c_g)

$$\mathcal{L}_{\text{eff.}} = \frac{\alpha_S}{\pi} \left(\frac{c_t}{12} (1 + \delta) + c_g \right) G^{\mu\nu} G_{\mu\nu} \frac{h}{v}, \quad (\text{or CP-odd})$$

HPAIR

- Higgs pair production cross section via gluon fusion
 - NLO in QCD in either SM or MSSM
 - Heavy top limit
- EFT contribution
 - 5-parameter anomalous coupling Lagrangian
 - Alternative ‘SILH’ basis input maps to the anomalous couplings
 - Dipole operator?

$$\mathcal{L}_{\text{SILH}} = y_u \frac{g_s^2 \bar{c}_g}{m_W^2} |\Phi|^2 - y_u \frac{\bar{c}_u}{v^2} |\Phi|^2 \bar{Q}_L \tilde{\Phi} u_R + \frac{\bar{c}_H}{2v^2} \partial^\mu |\Phi|^2 \partial_\mu |\Phi|^2 - \lambda \frac{\bar{c}_6}{v^2} |\Phi|^6$$

$$\mathcal{L}_{\text{eff.}} = -m_t \bar{t} t \left(c_t \frac{h}{v} + c_{tt} \frac{h^2}{2v^2} \right) - \frac{m_h^2}{2v} c_3 h^3 + \frac{\alpha_S}{\pi} G^{\mu\nu} G_{\mu\nu} \left(c_g \frac{h}{v} + c_{gg} \frac{h^2}{2v^2} \right)$$

$$c_t = 1 - \frac{\bar{c}_H}{2} - \bar{c}_u, \quad c_{tt} = -\frac{1}{2}(\bar{c}_H - 3\bar{c}_u), \quad c_3 = 1 - \frac{3\bar{c}_H}{2} + \bar{c}_6, \quad c_g = c_{gg} = \frac{4\pi}{\alpha_2} \bar{c}_g$$

Also HiggsPair: LO event generation with HERWIG 7.0

SILH operators

- Higgs-EW gauge boson operators in SILH basis

$$\begin{aligned} \mathcal{L}_{D6} = & \frac{1}{\Lambda^2} \left[\frac{g'^2}{4} \bar{c}_{BB} \Phi^\dagger \Phi B^{\mu\nu} B_{\mu\nu} + \frac{ig}{2} \bar{c}_W [\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig'}{2} \bar{c}_B [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \right. \\ & + ig \bar{c}_{HW} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] W_{\mu\nu}^k + ig' \bar{c}_{HB} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} \\ & \left. + \frac{g'^2}{4} \tilde{c}_{BB} \Phi^\dagger \Phi B^{\mu\nu} \tilde{B}_{\mu\nu} + ig \tilde{c}_{HW} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] \tilde{W}_{\mu\nu}^k + ig' \tilde{c}_{HB} [D^\mu \Phi^\dagger D^\nu \Phi] \tilde{B}_{\mu\nu} \right] \end{aligned}$$

$$\Phi^\dagger \overleftrightarrow{D}^\mu \Phi \equiv (D^\mu \Phi^\dagger) \Phi - \Phi^\dagger (D^\mu \Phi)$$

MCFM/POWHEG-box implementation:
operator definition has a factor 1/2 for (c_W, c_{HW})

- New Lorentz structures (1) & (2):

$$\begin{aligned} \mathcal{L}_{HAC} = & -\frac{1}{4} g_{hzz}^{(1)} Z_{\mu\nu} Z^{\mu\nu} h - g_{hzz}^{(2)} Z_\nu \partial_\mu Z^{\mu\nu} h + \frac{1}{2} g_{hzz}^{(3)} Z_\mu Z^\mu h - \frac{1}{4} \tilde{g}_{hzz} Z_{\mu\nu} \tilde{Z}^{\mu\nu} h \\ & - \frac{1}{2} g_{hww}^{(1)} W^{\mu\nu} W_{\mu\nu}^\dagger h - \left[g_{hww}^{(2)} W^\nu \partial^\mu W_{\mu\nu}^\dagger h + \text{h.c.} \right] + g_{hww}^{(3)} W_\mu W^{\dagger\mu} h - \frac{1}{2} \tilde{g}_{hww} W^{\mu\nu} \tilde{W}_{\mu\nu}^\dagger h \\ & - \frac{1}{2} g_{haz}^{(1)} Z_{\mu\nu} F^{\mu\nu} h - g_{haz}^{(2)} Z_\nu \partial_\mu F^{\mu\nu} h - \frac{1}{2} \tilde{g}_{haz} Z_{\mu\nu} \tilde{F}^{\mu\nu} h \end{aligned}$$

Feynman Rules

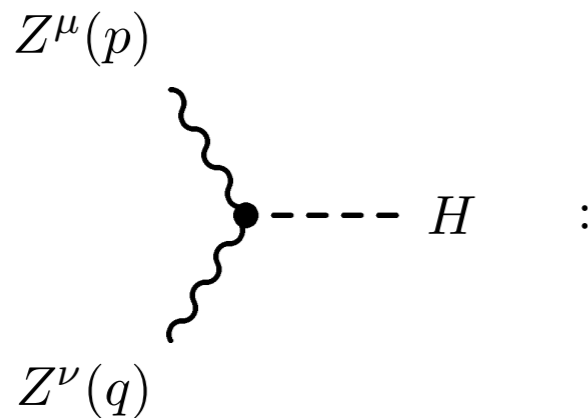


Diagram showing a Z boson (wavy line) with momentum p and index μ entering from the top, and another Z boson (wavy line) with momentum q and index ν entering from the bottom. A Higgs boson (dashed line) with momentum H exits to the right from a central vertex.

$$: i \left[\eta^{\mu\nu} \left(\frac{g}{\cos \theta_W} M_Z + g_{hzz}^{(1)} p \cdot q + g_{hzz}^{(2)} (p^2 + q^2) \right) - g_{hzz}^{(1)} q^\mu p^\nu - \tilde{g}_{hzz} \epsilon^{\mu\nu\rho\sigma} q_\rho p_\sigma - g_{hzz}^{(2)} (p^\mu p^\nu + q^\mu q^\nu) \right]$$

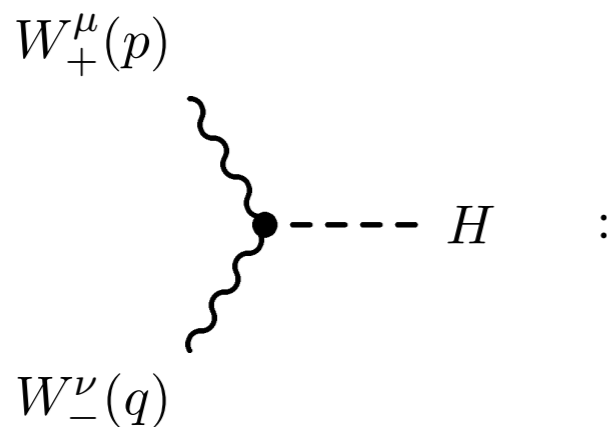


Diagram showing a W_+ boson (wavy line) with momentum p and index μ entering from the top, and a W_- boson (wavy line) with momentum q and index ν entering from the bottom. A Higgs boson (dashed line) with momentum H exits to the right from a central vertex.

$$: i \left[\eta^{\mu\nu} \left(g M_W + g_{hww}^{(1)} p \cdot q + g_{hww}^{(2)} (p^2 + q^2) \right) - g_{hww}^{(1)} q^\mu p^\nu - \tilde{g}_{hww} \epsilon^{\mu\nu\rho\sigma} q_\rho p_\sigma - g_{hww}^{(2)} (p^\mu p^\nu + q^\mu q^\nu) \right]$$

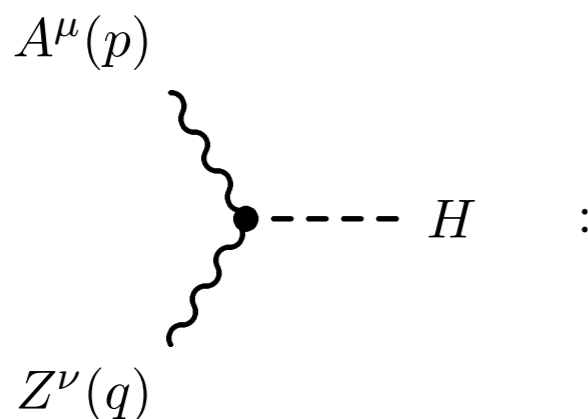


Diagram showing an Axion (wavy line) with momentum p and index μ entering from the top, and a Z boson (wavy line) with momentum q and index ν entering from the bottom. A Higgs boson (dashed line) with momentum H exits to the right from a central vertex.

$$: i \left[\eta^{\mu\nu} \left(g_{haz}^{(1)} p \cdot q + g_{haz}^{(2)} p^2 \right) - g_{haz}^{(1)} q^\mu p^\nu - \tilde{g}_{haz} \epsilon^{\mu\nu\rho\sigma} q_\rho p_\sigma - g_{haz}^{(2)} p^\mu p^\nu \right]$$

BSM \rightarrow

Mapping to AC/(i.e. HC)

| Coupling | HEL@NLO |
|-----------------|---|
| $g_{hzz}^{(1)}$ | $\frac{e^2 v}{2\hat{c}_W^2 \hat{s}_W^2} \frac{1}{\Lambda^2} [\hat{c}_W^2 \bar{c}_{HW} + 2\hat{s}_W^2 \bar{c}_{HB} - 2\hat{s}_W^4 \bar{c}_{BB}]$ |
| $g_{hzz}^{(2)}$ | $\frac{e^2 v}{4\hat{s}_W^2 \hat{c}_W^2 \Lambda^2} [\hat{c}_W^2 (\bar{c}_{HW} + \bar{c}_W) + 2\hat{s}_W^2 (\bar{c}_B + \bar{c}_{HB})]$ |
| $g_{hzz}^{(3)}$ | $\frac{g^2 v}{2\hat{c}_W^2} + \frac{e^4 v^3}{8\hat{c}_W^4 \hat{s}_W^2 \Lambda^2} [\hat{c}_W^2 \bar{c}_W + 2\bar{c}_B]$ |
| $g_{haz}^{(1)}$ | $\frac{e^2 v}{4\hat{s}_W \hat{c}_W \Lambda^2} [\bar{c}_{HW} - 2\bar{c}_{HB} + 4\hat{s}_W^2 \bar{c}_{BB}]$ |
| $g_{haz}^{(2)}$ | $\frac{e^2 v}{4\hat{s}_W \hat{c}_W \Lambda^2} [\bar{c}_{HW} + \bar{c}_W - 2(\bar{c}_B + \bar{c}_{BB})]$ |
| $g_{hww}^{(1)}$ | $\frac{e^2 v}{2\hat{s}_W^2 \Lambda^2} \bar{c}_{HW}$ |
| $g_{hww}^{(2)}$ | $\frac{ve^2}{4\Lambda^2 \hat{s}_W^2} [\bar{c}_W + \bar{c}_{HW}]$ |
| $g_{hww}^{(3)}$ | $\frac{g^2 v}{2}$ |

Anomalous couplings (AC) equivalent to
Higgs Characterisation (HC)

Limits from global fits

- A number of global fits to data deriving constraints on EFT Wilson coefficients have been performed
 - LHC, LEP & other low-energy experiments
- Marginalised constraints from EWPO + LHC Run 1 data on coefficients of interest

[Sanz et al.; JHEP 1503 (2015) 157]

| Operator | Coefficient | Constraints |
|--|---|-------------------|
| $\mathcal{O}_W = \frac{ig}{2} \left(H^\dagger T_{2k} \overleftrightarrow{D}^\mu H \right) D^\nu W_{\mu\nu}^k$ | $\frac{m_W^2}{\Lambda^2} (\bar{c}_W - \bar{c}_B)$ | (-0.035, 0.005) |
| $\mathcal{O}_B = \frac{ig'}{2} \left(H^\dagger \overleftrightarrow{D}^\mu H \right) \partial^\nu B_{\mu\nu}$ | $\frac{m_W^2}{\Lambda^2} (\bar{c}_W + \bar{c}_B)$ | (-0.0033, 0.0018) |
| $\mathcal{O}_{HW} = ig(D^\mu H)^\dagger T_{2k}(D^\nu H)W_{\mu\nu}^k$ | $\frac{m_W^2}{\Lambda^2} \bar{c}_{HW}$ | (-0.07, 0.03) |
| $\mathcal{O}_{HB} = ig'(D^\mu H)^\dagger (D^\nu H)B_{\mu\nu}$ | $\frac{m_W^2}{\Lambda^2} \bar{c}_{HB}$ | (-0.045, 0.075) |

stronger & weaker directions...

See also: [Falkowski & Riva; JHEP 1502 (2015) 039], [Berthier & Trott; JHEP 1505 (2015) 024], [Corbett et al.; JHEP 1508 (2015) 156], [Englert et al.; EPJC 76 (2016) 7, 393]

EFT Benchmarks

- To showcase the usage of both implementations, we select points in c_W, c_{HW} parameter space that:
 - Approximately saturate these limits
 - Select particular Lorentz structures in the new vertices
 - Are also motivated from a BSM point of view
- Tightly constrained direction in (c_B, c_W) forces $c_B \sim -c_W/2$

$$\mathcal{L}_{\text{new}} = -\frac{1}{4} g_{hvv}^{(1)} V_{\mu\nu} V^{\mu\nu} h - g_{hvv}^{(2)} V_\nu \partial_\mu V^{\mu\nu} h$$

- We then pick benchmark points that single out:

I) $V_\nu \partial_\mu V^{\mu\nu} h : g_{hvv}^{(1)} = 0, g_{hvv}^{(2)} \neq 0 \rightarrow \bar{c}_{HW} = 0, \bar{c}_W \neq 0$

II) $V_{\mu\nu} V^{\mu\nu} h : g_{hvv}^{(2)} = 0, g_{hvv}^{(1)} \neq 0 \rightarrow \bar{c}_W = -\bar{c}_{HW}$

EFT Benchmarks

- Pattern II) is a feature of matching conditions that arise in a large class of UV completions, e.g. 2HDM

[Gorbahn, No & Sanz; JHEP 1510 (2015) 036]

- Constraints then become tighter:

$$c_{HW} = -\bar{c}_W = (0.0008, 0.04)$$

- Summary of benchmarks used, roughly compatible with current limits

| POWHEG/MCFM | \bar{c}_{HW} | \bar{c}_W | \bar{c}_B | $g_{hvv}^{(1)}$ | $g_{hvv}^{(2)}$ |
|-------------|----------------|-------------|-------------|-----------------|-----------------|
| I | 0 | 0.008 | 0 | X | ✓ |
| II | 0.008 | -0.008 | 0 | ✓ | X |
| MG5_aMC | | | | | |
| A | 0.03 | 0 | 0 | ✓ | ✓ |
| B | 0.03 | -0.03 | 0.015 | ✓ | X |

Selection of results

- ZH in POWHEG-BOX/MCFM, including SM $gg \rightarrow ZH$
- WH, VBF in FR-NLOCT/Madgraph5_aMC@NLO
- Used PYTHIA8 for Higgs decay, PS and Hadronisation
 - Rescaled rates by eHDECAY BRs to capture EFT contributions
- Events were reconstructed using Fastjet thanks to MadAnalysis5 “reco” mode and analysed according to some realistic event selection procedure also in MA5
- Theoretical uncertainties due to scale variation were quantified but not PDF uncertainties

POWHEG-BOX/MCFM

- Higgs associated production with a leptonically decaying W or Z at NLO in QCD matched to parton shower
 - Include EFT effects via a mapping to AC/HC (also CP violating)
- At NLO, the initial state current factorises from the final state, even when the Higgs decays to b 's
 - Drell-Yan-like NLO corrections which are well known
- Builds upon previous work in the SM matched to parton shower in the same framework as well as fixed order predictions including anomalous couplings
- Matrix elements based on MCFM code interfaced with POWHEG-BOX for which the SM process was already implemented

Simulation

- For definiteness we specified that the Higgs decay to bb , allowing PYTHIA to perform the decay but scaling the rates by the BR predicted by eHDECAY
- Used CTEQ10 PDFs for NLO predictions and CTEQ6L1 PDFs for LO comparisons
- Modification of EW parameters taken into account in the (m_Z, m_W, G_F) input scheme
- Scale uncertainty determined by varying μ_R, μ_F together around a central scale of $\mu_0 = m_{VH}$
 - Envelope of $\mu_0/2$ and $2 \mu_0$

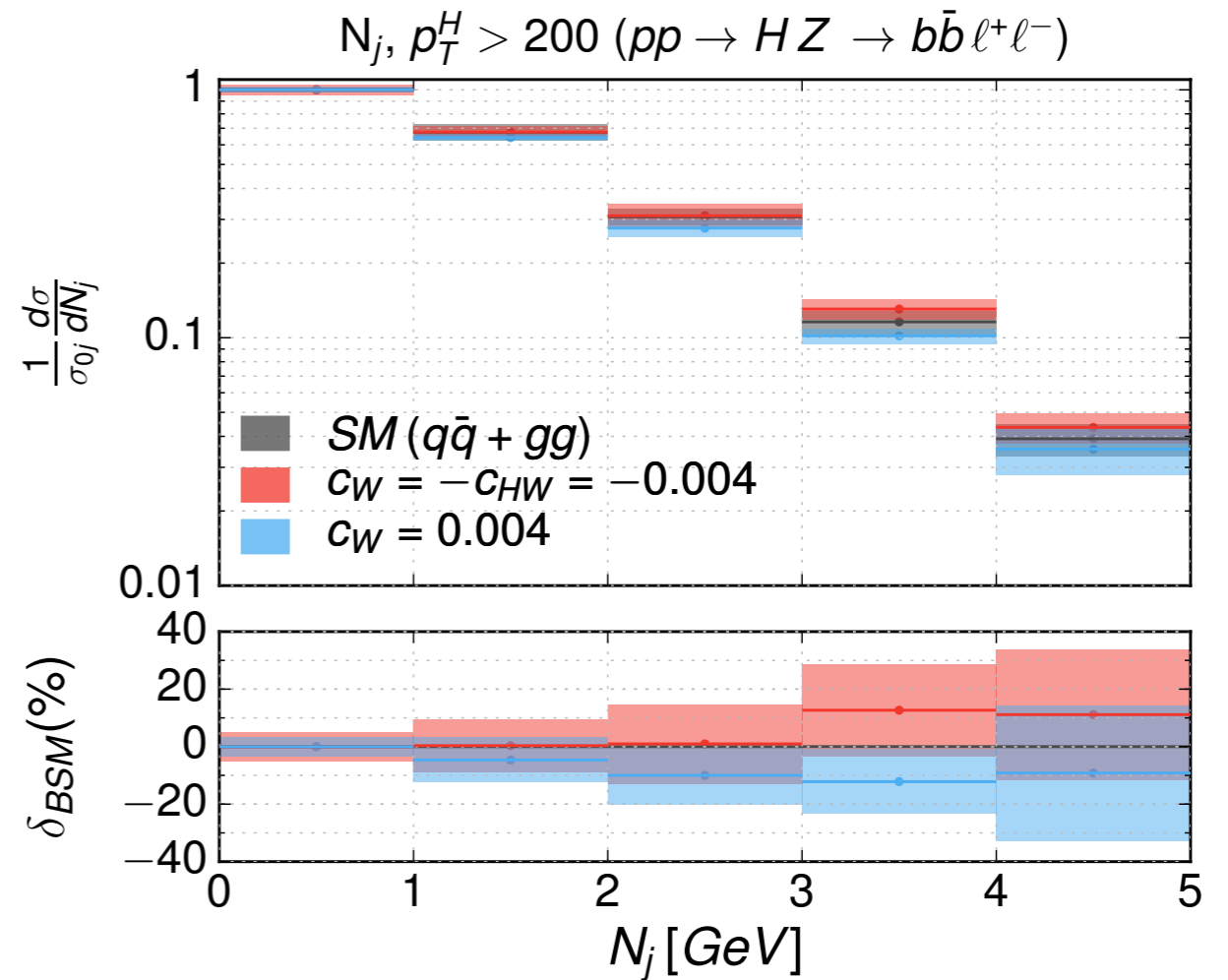
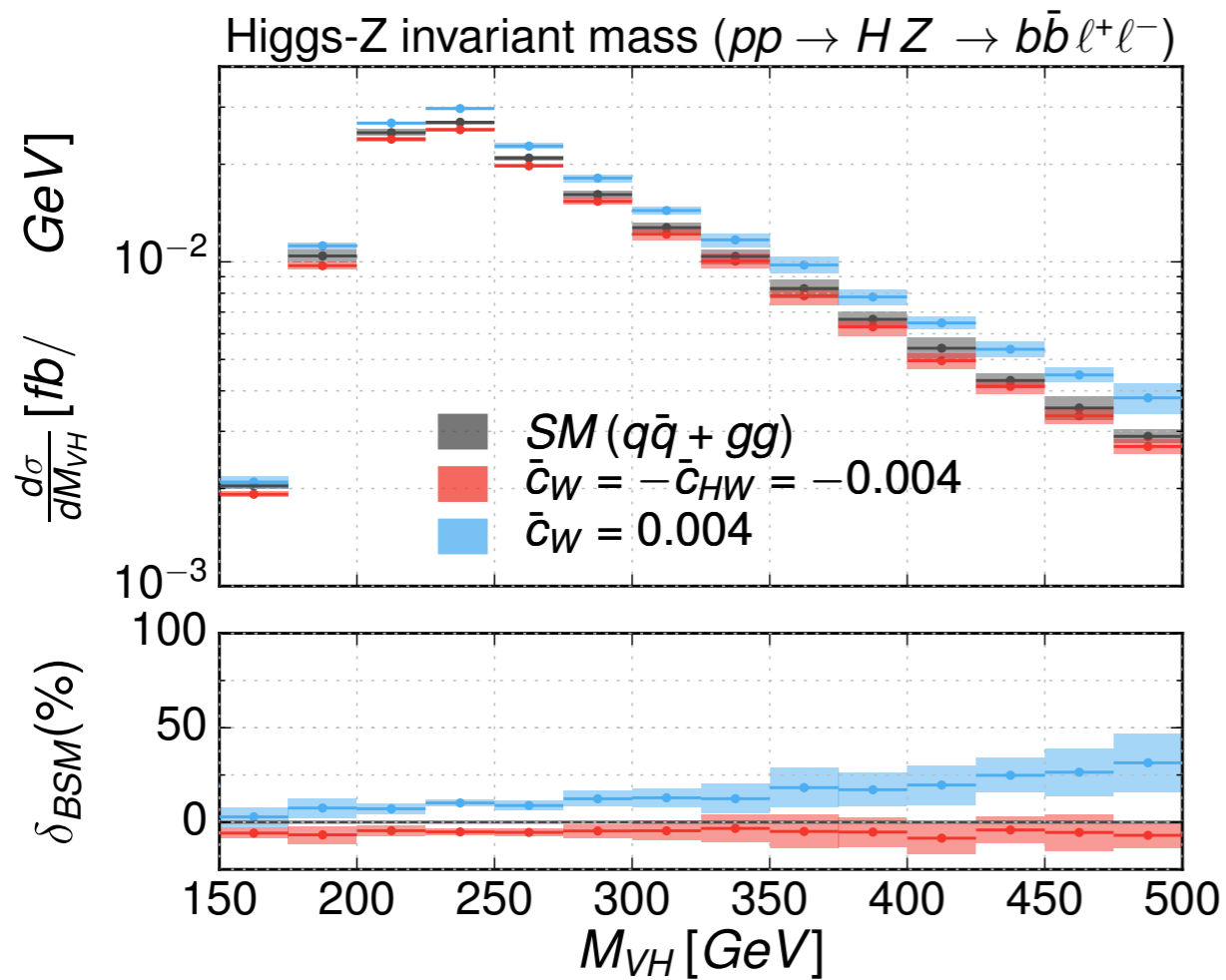
Selection

| Process | |
|---|--|
| $H Z \rightarrow b\bar{b} \ell^+ \ell^-$ | $H W \rightarrow b\bar{b} \ell \nu$ |
| Jets | |
| k_T algorithm: $\Delta R=0.4, p_T > 25 \text{ GeV} \ \& \ \eta_b < 2.5$ | |
| Cuts | |
| 2 b -jets, $p_T > 25 \text{ GeV}, \eta_b < 2.5$ | |
| 1 lepton, ℓ^\pm (e or μ) | 2 leptons, ℓ^+, ℓ^- (e or μ) |
| $p_T^\ell < 25 \text{ GeV}, \eta_\ell < 2.5$ | |

MA5 performs b -jet identification based on truth level jet information (presence of b -hadrons in jet)

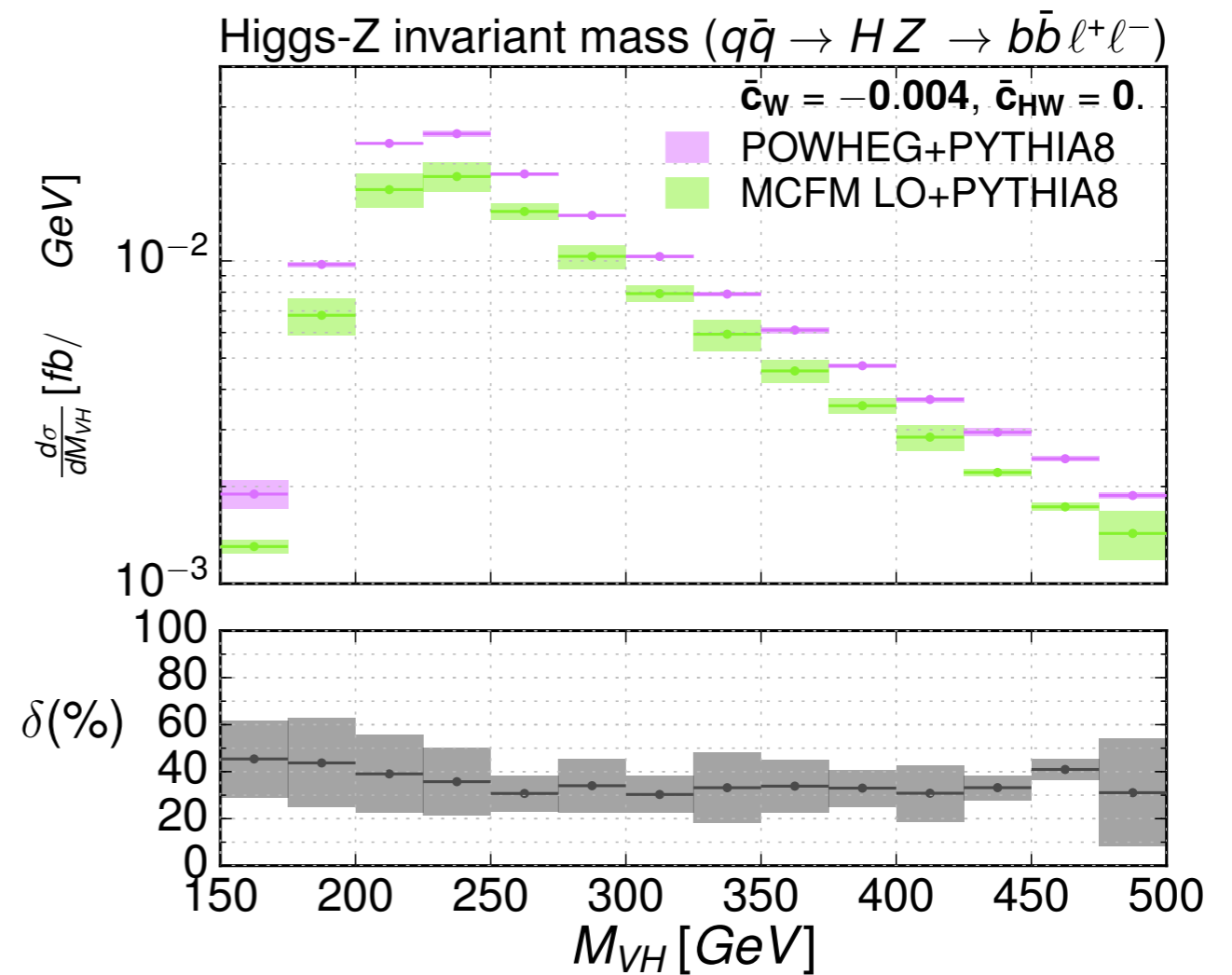
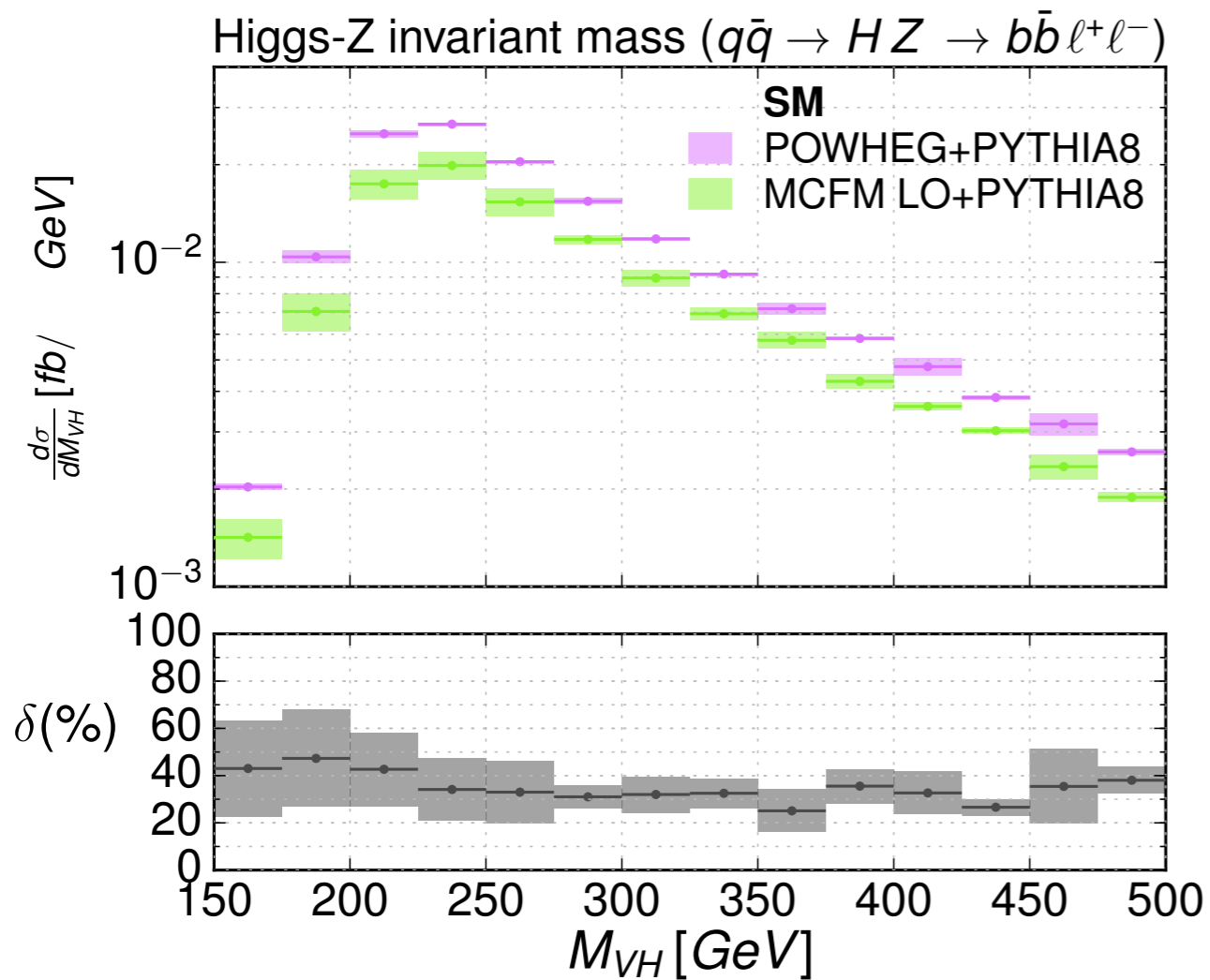
$$pp \rightarrow ZH \rightarrow l^+l^- bb$$

N_j exhibits some difference but stats too low to distinguish



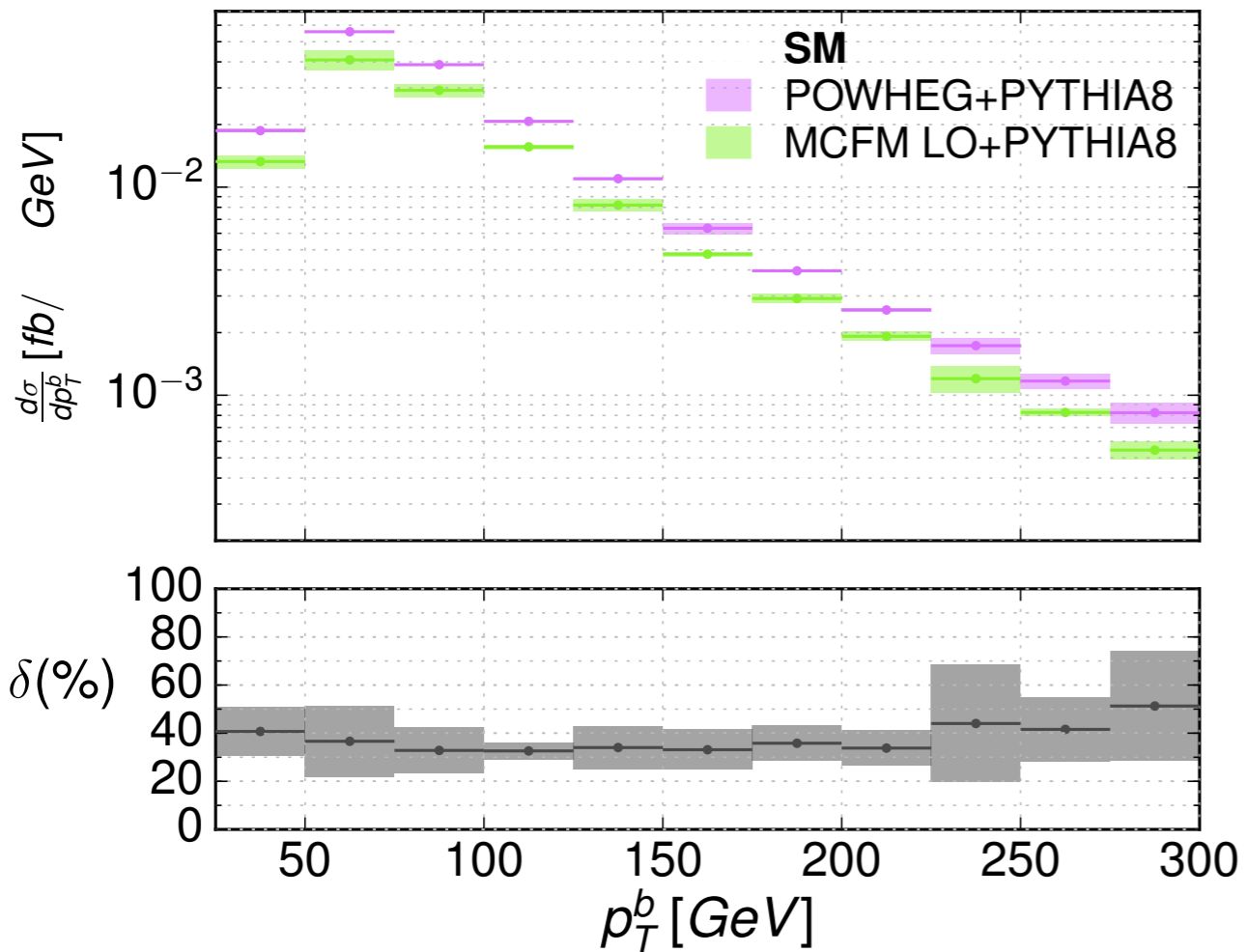
K-factors

No significant difference between SM & EFT
Relatively flat

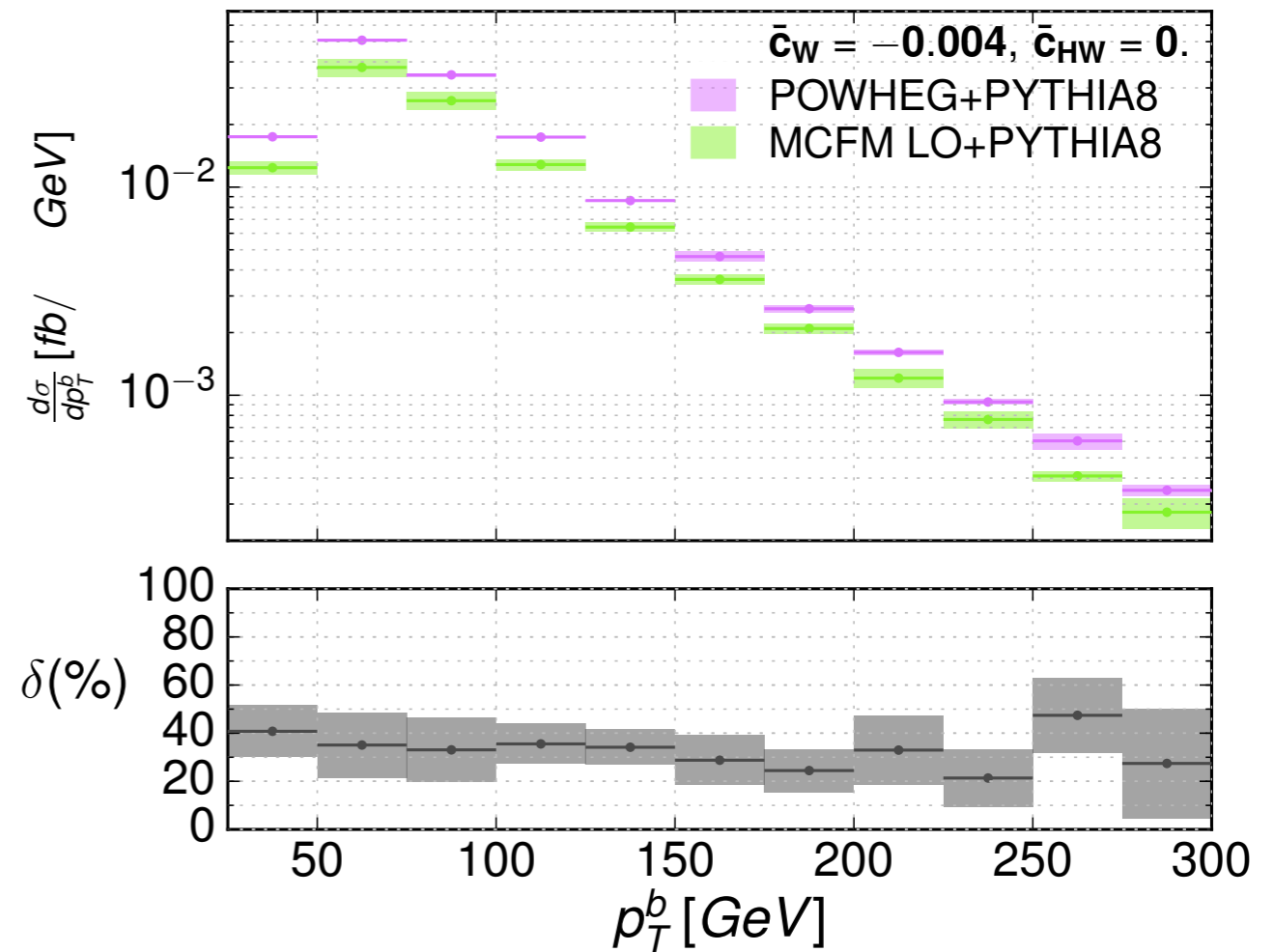


K-factors

Leading b -jet p_T ($q\bar{q} \rightarrow HZ \rightarrow b\bar{b}l^+l^-$)



Leading b -jet p_T ($q\bar{q} \rightarrow HZ \rightarrow b\bar{b}l^+l^-$)



FR+NLOCT/MG5_aMC@NLO

<http://feynrules.irmp.ucl.ac.be/wiki/HELatNLO>

- HEL@NLO: Alternative & independent implementation within the FeynRules + NLOCT framework
 - Generate NLO ready UFO file
 - Simulation performed with MadGraph5_aMC@NLO ~ any process!
[Alloul, Fuks & Sanz; JHEP 1404 (2014) 110]
- Builds upon previous LO implementation of full SILH basis
- Modification of EW parameters fully taken into account in the (m_Z, α_S, G_F) input scheme
- Scale uncertainty determined by varying (reweighting) μ_R, μ_F independently around a central scale of μ_0
 - Envelope of 9 combinations of $(1/2, 2) \times \mu_0$

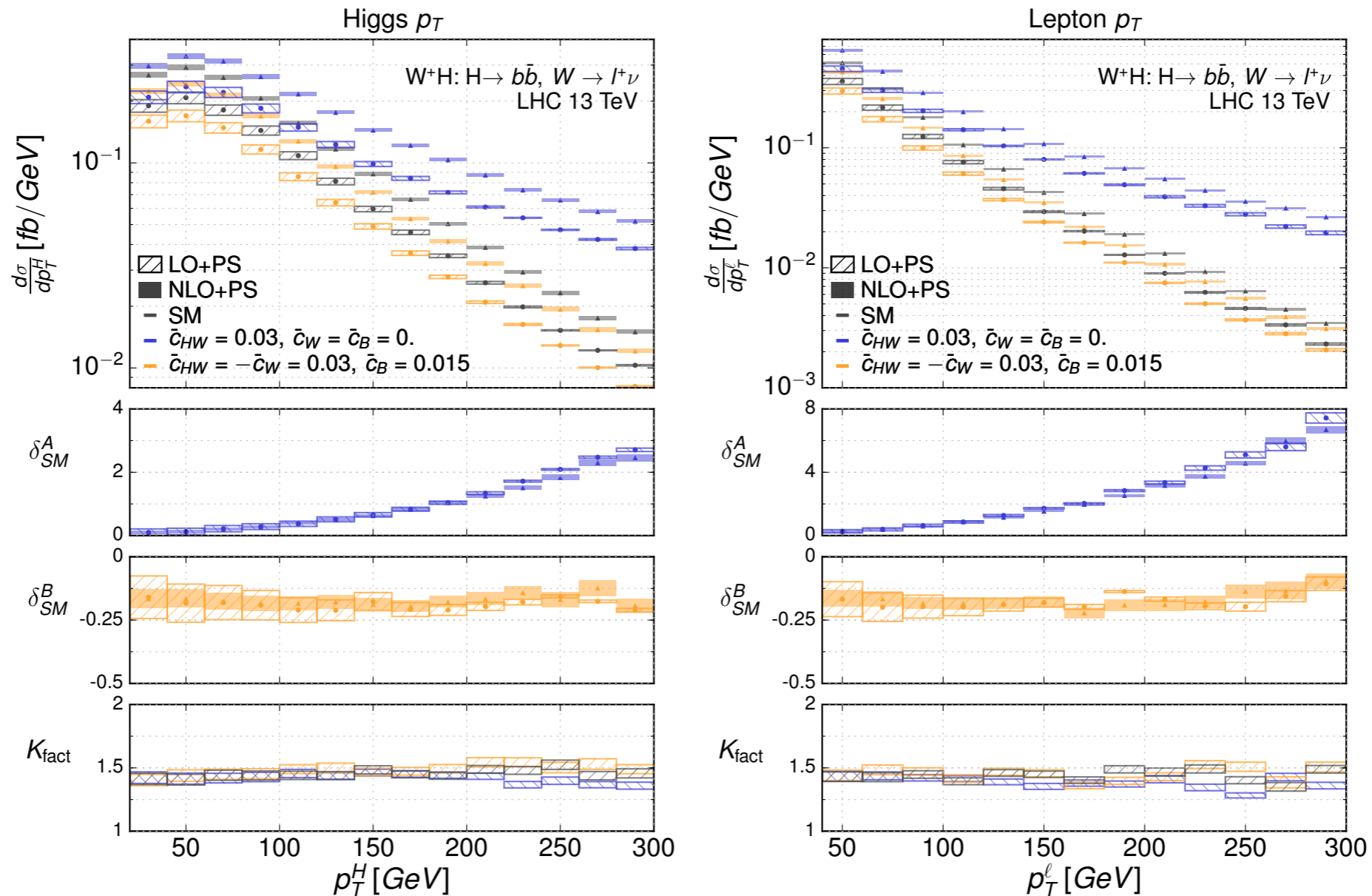
Simulation

- For WH we specified that the Higgs decay to bb, as in previous example while for VBF, decay Higgs to photons
- Used NNPDF23_nlo_as_0118_qed PDFs
- Made use of recent MG5 feature to select only interference terms for comparison
 - Specify coupling order squared , e.g., “NP²≤2” to get interference
- Validated results against POWHEG-BOX implementation

```
generate p p > h ve e+ [QCD]
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$$pp \rightarrow W^+ H \rightarrow l^+ \nu bb$$

Benchmarks correspond to ‘large’ values of Wilson coefficients as described in previous analysis since they saturate current limits

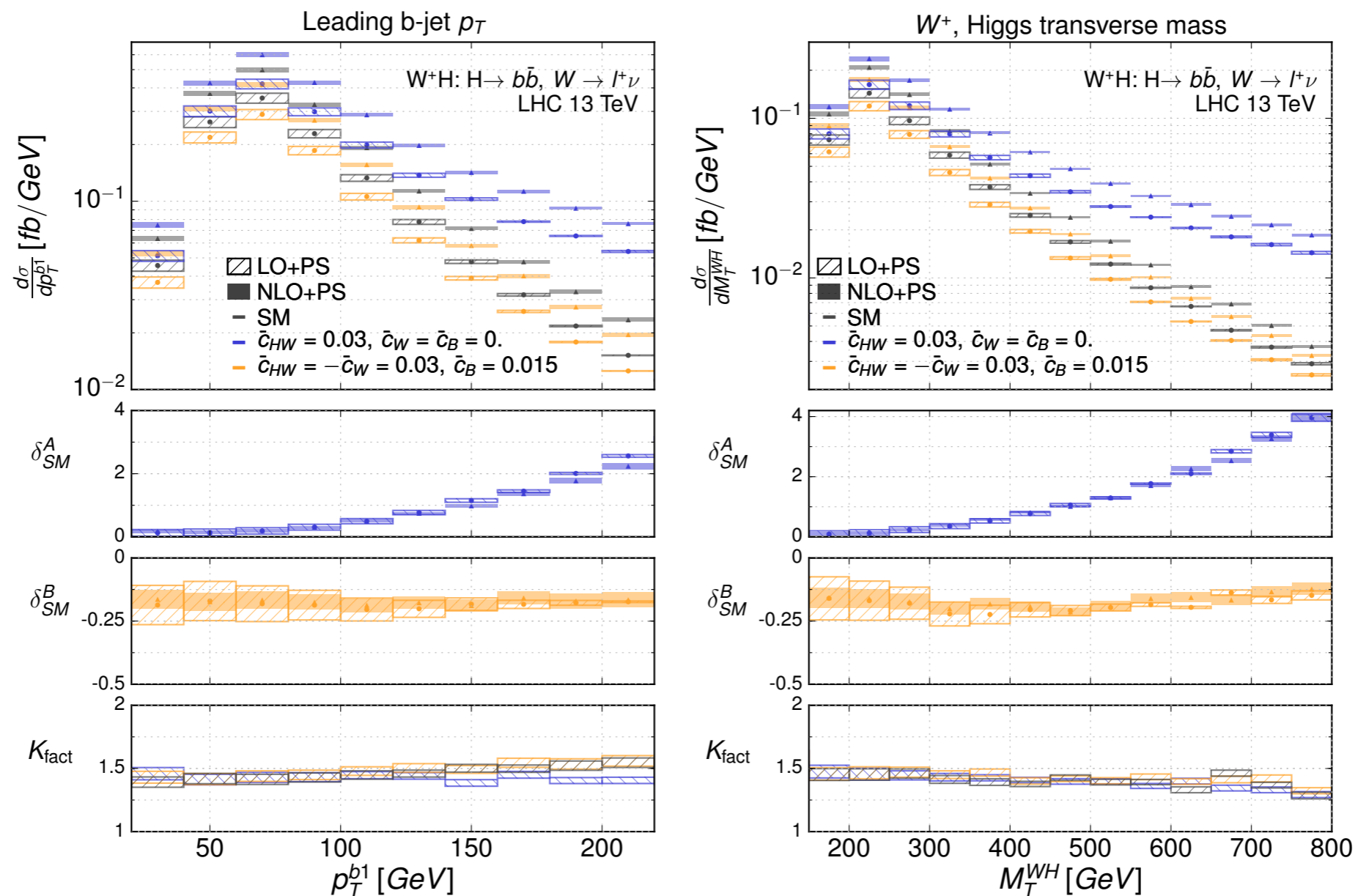


$$\delta = \text{TOT}/\text{SM} - 1$$

NLO/LO

$$pp \rightarrow W^+ H \rightarrow l^+ \nu bb$$

Again, benchmark **B)** does not exhibit strong “EFT” features
 The $g_{h\nu\nu}^{(2)}$ Lorentz structure is responsible for these



Validation

