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Number of light and heavy neutrinos in Seesaw type I

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Seesaw mass matrix and diagonalization



Mass matrix

$$\mathcal{L}^{D+M} = \frac{1}{2} \tilde{N}^T C^\dagger M \tilde{N} + H.c.$$

$$M = \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix}$$

$$\tilde{N} = \begin{pmatrix} \nu_L \\ \nu_R^C \end{pmatrix}$$

two mass scales

$$M_R \gg M_D$$

Block diagonalization

$$M = \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix}$$

$$\lambda(M_R) \gg |M_D|$$

$$W^T M W \simeq \begin{pmatrix} M_{light} & 0 \\ 0 & M_{heavy} \end{pmatrix}$$

$$\left. \begin{aligned} M_{light} &\simeq -M_D^T M_R^{-1} M_D \\ M_{heavy} &\simeq M_R \end{aligned} \right\} \text{Seesaw}$$

e.g.

$$M_R \sim 10^{15} [\text{GeV}], \quad M_D \sim 10^2 [\text{GeV}] \Rightarrow M_{light} \sim 10^{-2} [\text{eV}]$$

$$\text{LHC [wishful thinking?]} \quad M_R \sim 10^3 [\text{GeV}], \quad M_D \sim 10^{-4} [\text{GeV}] \Rightarrow M_{light} \sim 10^{-2} [\text{eV}]$$



Why $\lambda(M_R) \gg |M_D|$?

Example

$$M_R = \begin{pmatrix} 100 & 100 & 100 \\ 100 & 100 & 100 \\ 100 & 100 & 100 \end{pmatrix} \Rightarrow \text{rank one matrix} \Rightarrow \lambda(A) = (300, 0, 0)$$

$$M = \begin{pmatrix} 0 & 0 & 0 & 1 & 1 & 1 \\ 0 & 0 & 0 & 1 & 1 & 1 \\ 0 & 0 & 0 & 1 & 1 & 1 \\ 1 & 1 & 1 & 100 & 100 & 100 \\ 1 & 1 & 1 & 100 & 100 & 100 \\ 1 & 1 & 1 & 100 & 100 & 100 \end{pmatrix} \Rightarrow$$

M_R, M_D rank one
 M has four **zero**
eigenvalues!!!

Czakon, Gluza and Zralek [2] gave
a formal proof that for $M_R \gg M_D$
we can not get a fourth light sterile
neutrino.



In general a full mass matrix M is complex and symmetric

Not Hermitian, not even normal!!!

Consequences:

1. A complex symmetric matrix may not even be diagonalizable

$$A = \begin{pmatrix} 2i & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow \begin{matrix} A^\dagger \neq A, A^\dagger A \neq AA^\dagger \\ \text{algebraic multiplicity} \neq \text{geometric multiplicity} \Rightarrow \text{not diagonalizable} \end{matrix}$$

2. Even if it is, we have $M_{heavy} \simeq M_R \Rightarrow$ small differences can dramatically change eigenvalues

$$A = \begin{pmatrix} \mu & 1 & 0 \\ 0 & \mu & 1 \\ 0 & 0 & \mu \end{pmatrix}, \quad E = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ \epsilon & 0 & 0 \end{pmatrix}.$$

If μ — eigenvalue of A , λ — (any) eigenvalue of $A + E$, then

$$|\lambda - \mu| = \epsilon^{1/3}$$

If $\epsilon = 10^{-15}$, then $|\lambda - \mu| = 10^{-5}$ — spectrum shifted by a number 10 orders larger than the parameter ϵ !



Number of neutrinos in CP invariant case

In the CP invariant seesaw scenario with two mass scales $M_R \gg M_D$, $M_D \in M_{3 \times n}$, $M_R \in M_{n \times n}$ and $\lambda(M_R) \gg |m_D|$ exactly 3 light neutrinos are present.

Matrix Decomposition

$$M = \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & M_R \end{pmatrix} + \begin{pmatrix} 0 & M_D \\ M_D^T & 0 \end{pmatrix} \equiv \hat{M}_R + \hat{M}_D$$



From Weyl's inequalities

$$\alpha_j + \beta_k \leq \gamma_{j+k-1} \quad \text{and} \quad \gamma_{n-j-k} \leq \alpha_{n-j+1} + \beta_{n-k+1}$$

$$|\lambda_i(M) - \lambda_i(\hat{M}_R)| \leq \rho(\hat{M}_D) \leq \|\hat{M}_D\|_F = \sqrt{\sum_{i,j} m_{ij}^2}, \quad \rho(\hat{M}_D) = \max\{|\lambda(\hat{M}_D)|\}$$

But

$$\sqrt{\sum_{i,j} m_{ij}^2} \sim |m_D| \Rightarrow \rho(\hat{M}_D) \leq |m_D|, \quad m_D \in M_D$$

On the other hand matrix \hat{M}_R has at least 3 eigenvalues equal 0.

Thus $|\lambda_i(M) - 0| \leq \rho(\hat{M}_D)$ for $\lambda_i(\hat{M}_R) = 0$

Hence, three eigenvalues of M must be smaller than $|m_D|$.



Conclusion (I): At least 3 light neutrinos exist.

Similar steps with assumption that $\lambda(M_R) \gg |m_D|$ gives
Conclusion (II): Remaining masses must be heavy.

Conclusions (I) and (II) implies Conclusion (III):
Heavy masses $N_{1,\dots,n}$ are maximally shifted by $|m_D|$ from eigenvalues of M_R .

e.g.

$$\lambda_1(M_R) = 100, |m_D| \sim 1 \Rightarrow \lambda_1(M) \simeq 100 \pm 1$$



- [1] Rejandra Bhatia. *Matrix Analysis*. Springer, 1996.
- [2] M. Czakon, J. Gluza, M. Zralek. *Seesaw mechanism and four light neutrino states*. Phys. Rev., D64: 117304, 2001