

# MIXED HARMONIC FLOW CORRELATIONS

## RECENT RESULTS ON SYMMETRIC CUMULANTS

Matthew Luzum

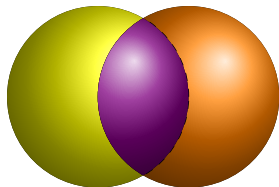
F. Gardim, F. Grassi, ML, J. Noronha-Hostler;  
arXiv:1608.02982

Universidade de São Paulo

*NBI Mini Workshop*  
19 September, 2016

# EVOLUTION OF HEAVY-ION COLLISION

- Hydro picture: system thermalizes and expands as fluid
- Particles emitted independently at end of evolution



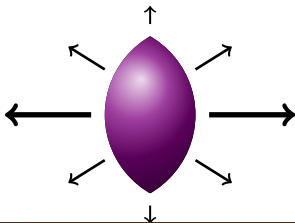
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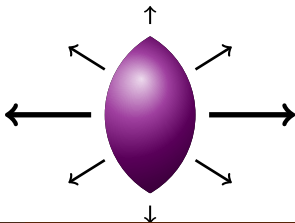
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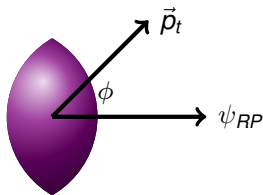
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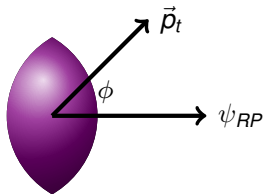


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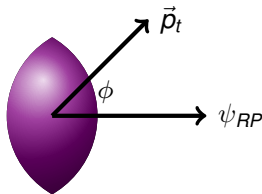
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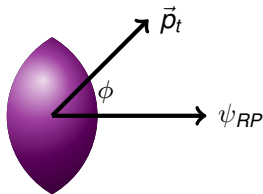
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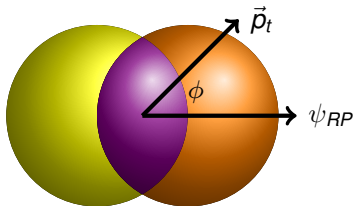
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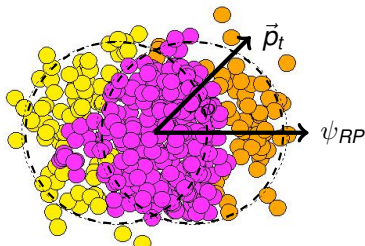
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- No symmetry  $\implies$  more harmonics, flow coefficients are *vectors*
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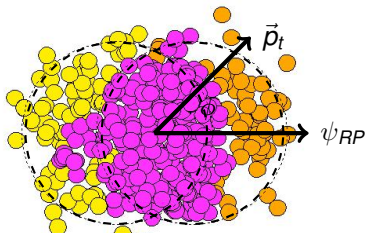
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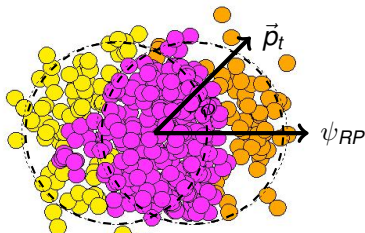
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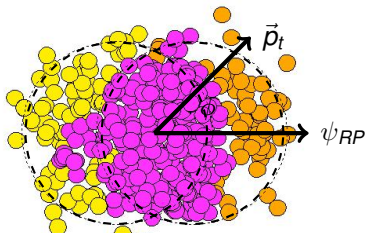
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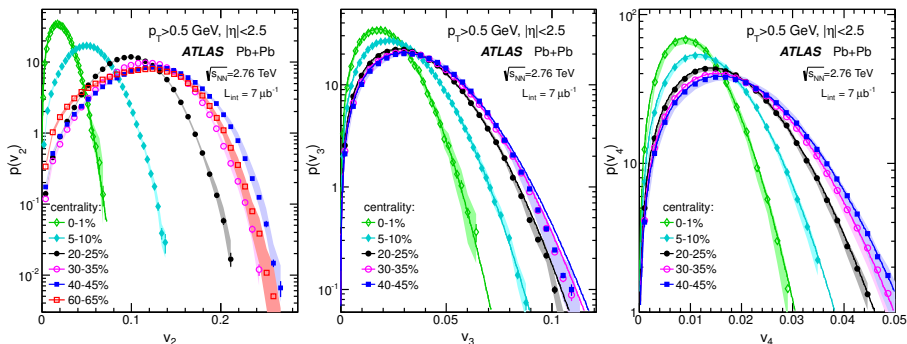


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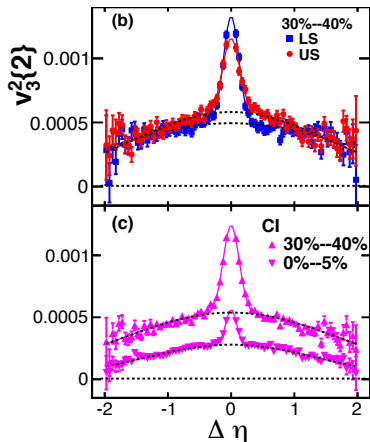
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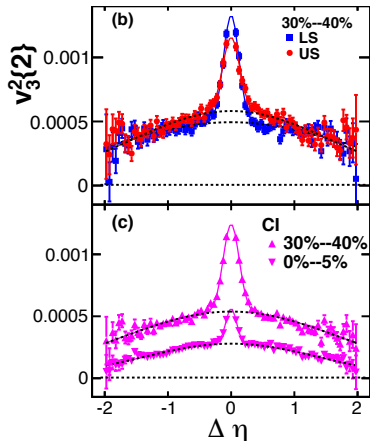
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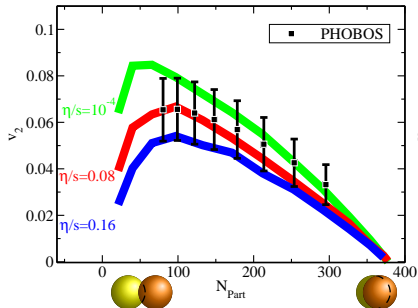


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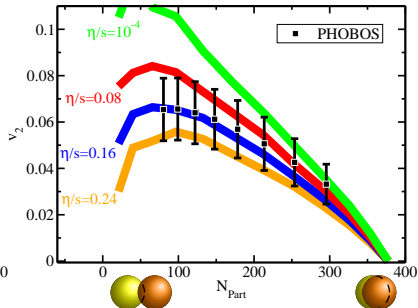
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- **Different harmonics can also have non-trivial correlations**

## ELLIPTIC FLOW ALONE

"Glauber" initial conditions

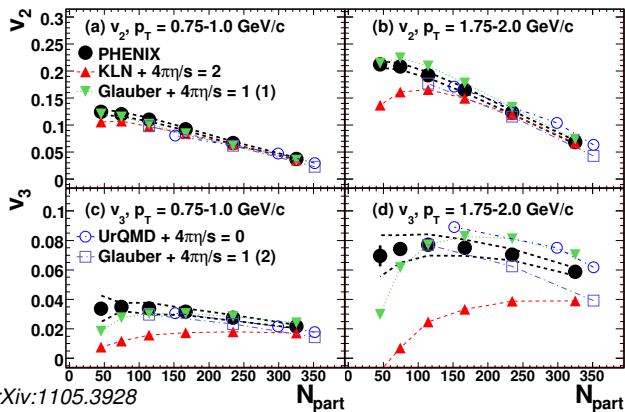


"CGC" initial conditions

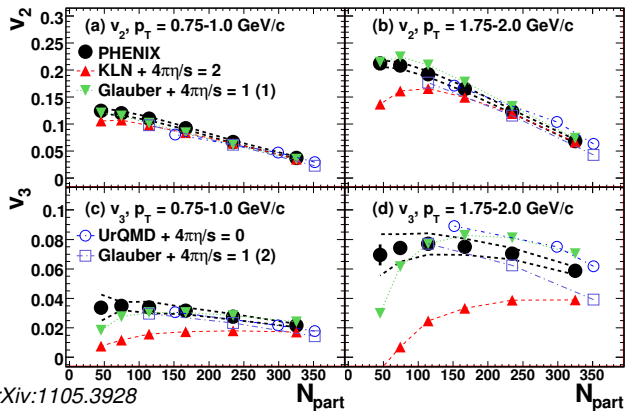


(ML & Romatschke, *Phys.Rev. C78 (2008) 034915*)

- Can fit  $v_2$  with different combinations of initial cond. + viscosity
- $\implies$  Neither initial conditions nor viscosity individually constrained

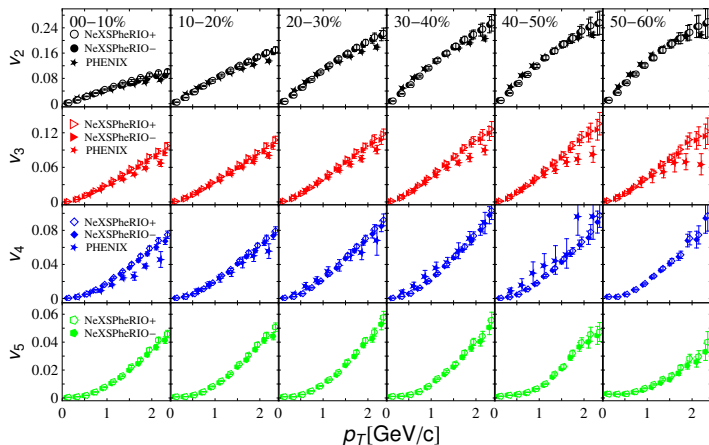
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- Combining  $v_2$  and  $v_3$  can rule out models of initial conditions
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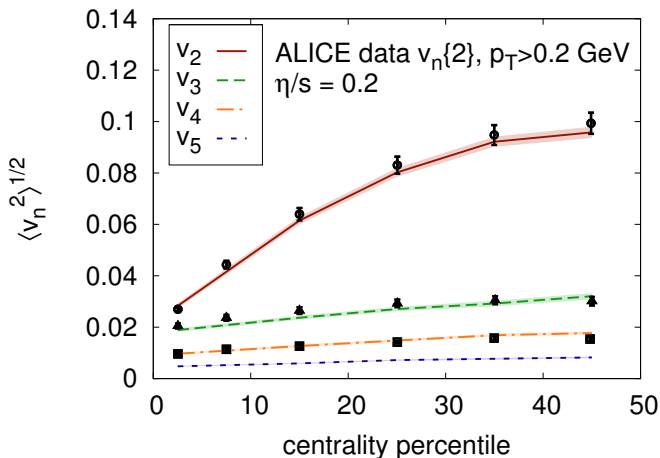
# CALCULATIONS VS RHIC DATA



(F. Gardim, F. Grassi, ML, J.-Y. Ollitrault, *Phys.Rev.Lett.* 109 (2012) 202302)

- NeXus initial conditions +  $\eta/s = 0$  at RHIC

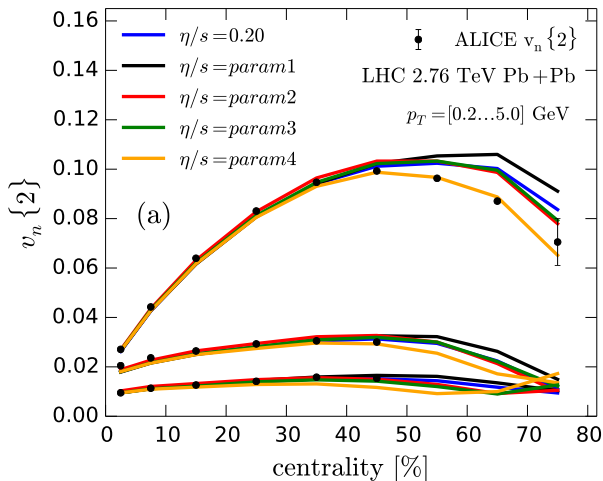
# CALCULATIONS VS LHC1 DATA



(C. Gale, S. Jeon, B. Schenke, P. Tribedy, R. Venugopalan, Phys.Rev.Lett. 110 (2013) 012302)

- IP-Glasma initial conditions +  $\eta/s = 0.2$  at LHC run 1

# CALCULATIONS VS LHC1 DATA



(H. Niemi, K.J. Eskola, R. Paatelainen, *Phys.Rev. C93* (2016) 024907)

- EKRT initial conditions + various  $\eta/s(T)$  at LHC run 1

# FLOW MEASUREMENTS

Basic building block is  $n$ -particle correlation:

$$\langle m \rangle_{n_1, n_2, \dots, n_m} \equiv \left\langle \left\langle \cos(n_1 \phi + n_2 \phi \dots + n_m \phi) \right\rangle_{m \text{ particles}} \right\rangle$$

$$\stackrel{(\text{flow})}{=} \langle V_{n_1} V_{n_2} \dots V_{n_m} \rangle$$

$$\sum n_i = 0$$

$$\langle \dots \rangle_{2 \text{ particles}} \equiv \frac{1}{N_{\text{pairs}}} \sum_{i=1}^M \sum_{j \neq i} \dots$$

$$\langle \dots \rangle \equiv \frac{\sum_{\text{events}} W \dots}{\sum_{\text{events}} W}$$

Typical weights:  $W_{\langle 2 \rangle} = M(M-1)$

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$$\stackrel{\text{(flow)}}{=} \sqrt{\langle v_n^2 \rangle},$$

Enforce gap in rapidity between pair  $\implies$  suppressed non-flow

$$-v_n\{4\}^4 \equiv \langle 4 \rangle_{n,n,-n,-n} - 2\langle 2 \rangle_{n,-n}^2$$

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# SYMMETRIC CUMULANTS

Can also define a mixed-harmonic cumulant:

$$SC(n, m) \equiv \langle 4 \rangle_{n, m, -n, -m} - \langle 2 \rangle_{n, -n} \langle 2 \rangle_{m, -m} \\ \stackrel{\text{(flow)}}{=} \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle.$$

Independent information better encoded in normalized quantity:

$$NSC(n, m) \equiv \frac{SC(n, m)}{\langle 2 \rangle_{n, -n} \langle 2 \rangle_{m, -m}} \\ \stackrel{\text{(flow)}}{=} \frac{\langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle}{\langle v_n^2 \rangle \langle v_m^2 \rangle}.$$

(Bilandzic, Christensen, Gulbrandsen, Hansen, Zhou, *Phys. Rev. C* 89, no. 6, 064904 (2014))

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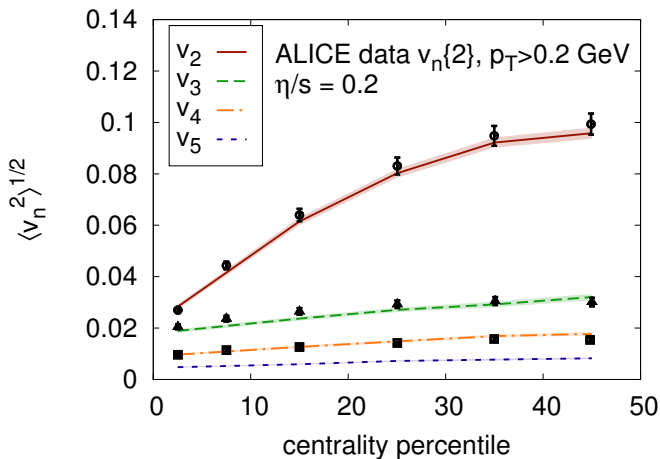
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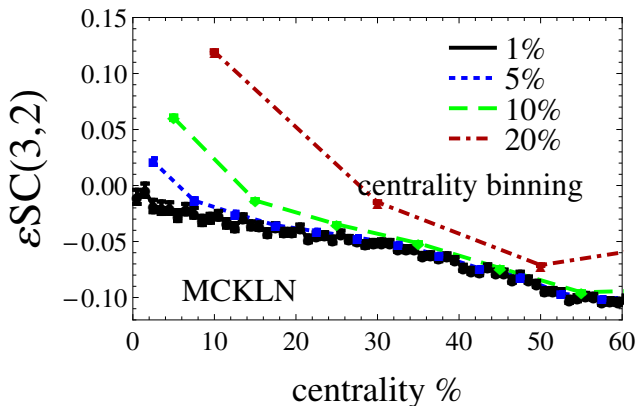
(Bilandzic, Christensen, Gulbrandsen, Hansen, Zhou, *Phys. Rev. C* 89, no. 6, 064904 (2014))

## CENTRALITY BINNING



- The magnitude of all  $v_n$  increase from central  $\rightarrow$  peripheral
- $\implies$  impact parameter fluctuations increase NSC( $n, m$ )

## CENTRALITY BINNING



- Large centrality bins produce a bias
- Biggest effect in central collisions, where  $v_n$  vs. centrality steepest

## EVENT WEIGHTING

$$\text{NSC}(n, m) \stackrel{(\text{flow})}{=} \frac{\langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle}{\langle v_n^2 \rangle \langle v_m^2 \rangle}$$

- Typically, 4-particle correlations have  $\sim M^4$  event weights, while 2-particle correlations have  $\sim M^2$  weighting
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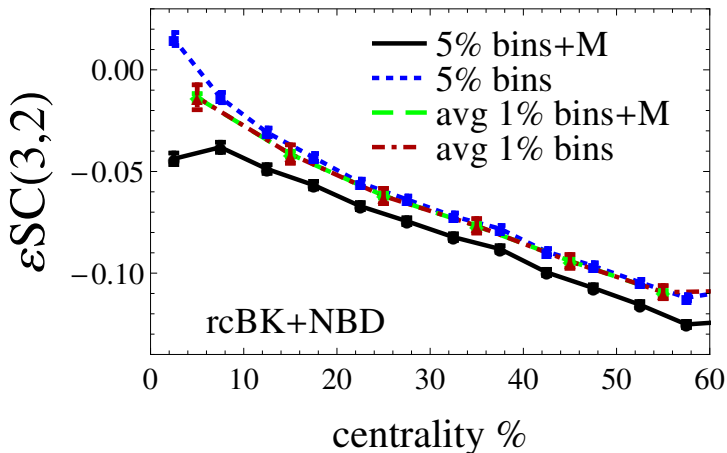
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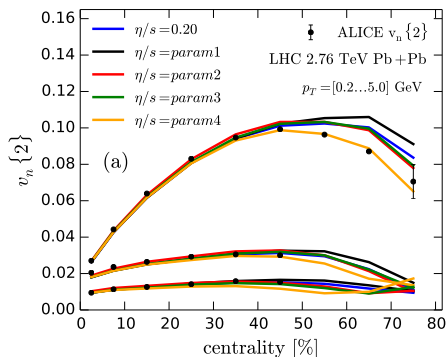
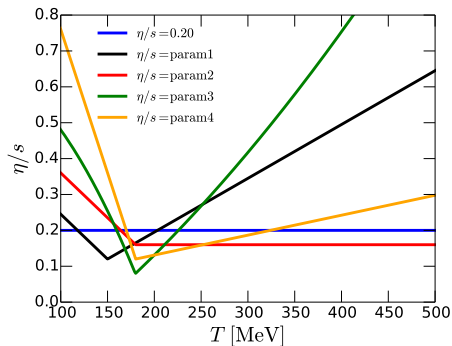


## EVENT WEIGHTING



- $\implies$  *negative* bias with large centrality bins
- No bias with small centrality bins

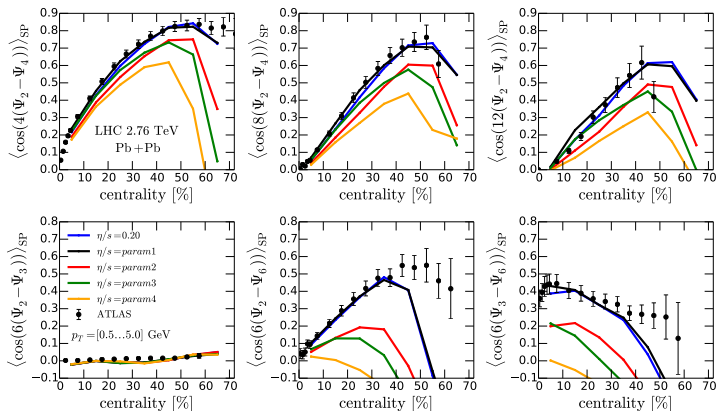
## FIRST CALCULATION



(H. Niemi, K.J. Eskola, R. Paatelainen, *Phys.Rev. C93* (2016) 024907)

- Various  $\eta/s(T)$  fit single harmonic data at RHIC + LHC.
- Other mixed harmonic measurements ruled out some

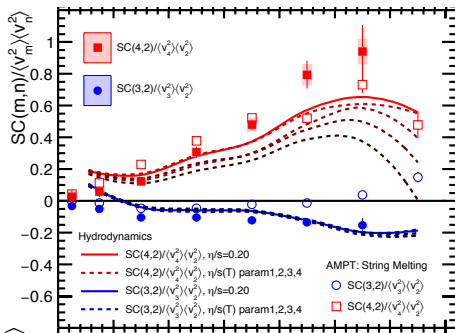
## FIRST CALCULATION



(H. Niemi, K.J. Eskola, R. Paatelainen, *Phys.Rev. C93 (2016) 024907*)

- Various  $\eta/s(T)$  fit single harmonic data at RHIC + LHC.
- Other mixed harmonic measurements ruled out some

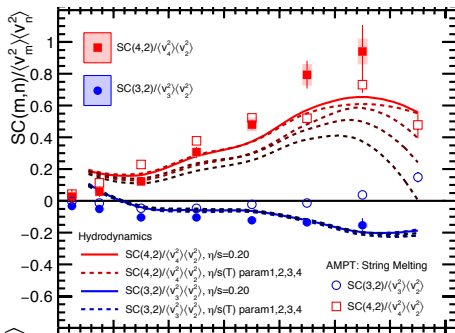
## FIRST CALCULATION / FIRST DATA



ALICE, arXiv:1604.07663; Niemi, Eskola, Paatelainen, 1505.02677

- Doesn't seem to fit NSC(3,2) and NSC(4,2)
- (may be affected by centrality binning effects)
- NSC(3,2) insensitive to  $\eta/s(T)$ , NSC(4,2) is sensitive

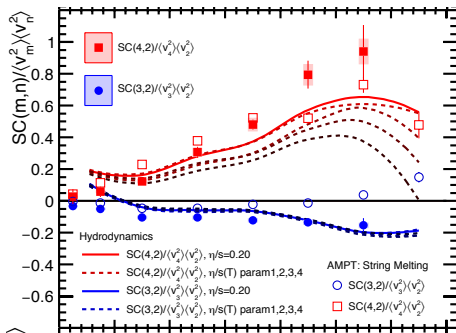
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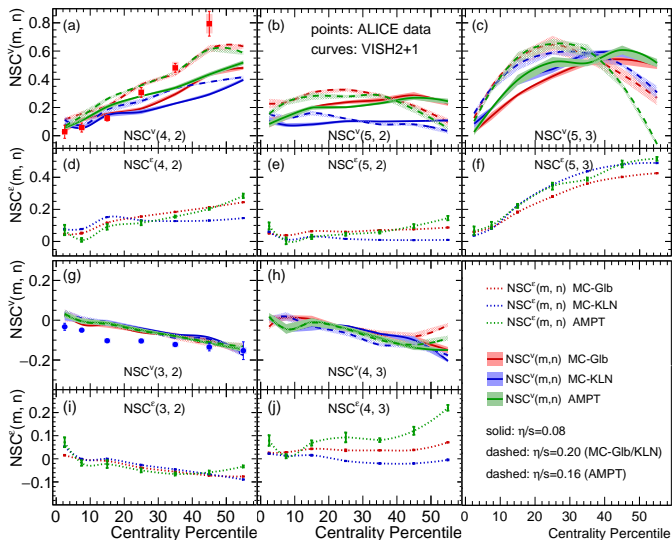
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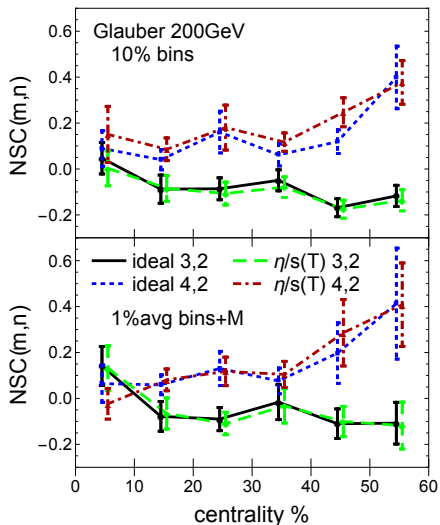
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Zhu, Zhou, Xu, Song; arXiv:1608.05305

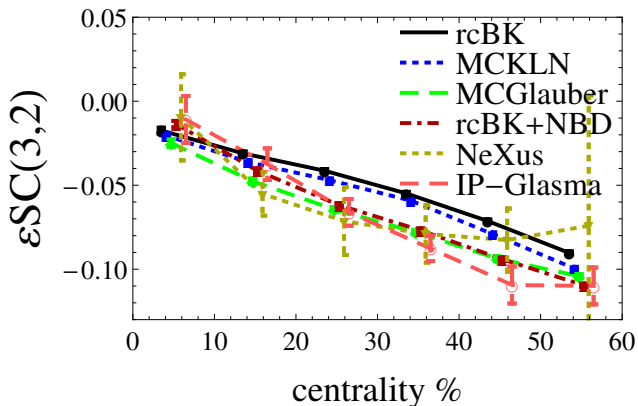
$NSC(3,2)$  also insensitive to magnitude of  $\eta/s$  ( $NSC(4,2)$  sensitive)



Even that observation can depend on centrality binning!

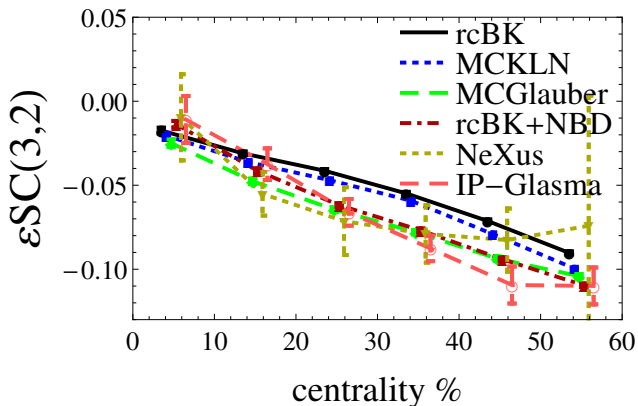


## NSC(3,2) VS. INITIAL CONDITIONS



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# SUMMARY/CONCLUSIONS

- Mixed harmonic correlation measurements shed new light on heavy-ion collisions
- Symmetric cumulants measure correlation between magnitudes of flow vectors of different harmonics, with suppression of non-flow correlations
- Different sensitivity to initial conditions and viscosity than previous observables.

# EXTRA SLIDES

# CUMULANT EXPANSION

Idea: (Teaney, Yan; *Phys.Rev. C83 (2011) 064904*)

- Characterize density by moments (or cumulants) of 2-D Fourier transform

$$\rho(\mathbf{k}) = \int d^2x \rho(\mathbf{x}) e^{i\mathbf{k}\cdot\mathbf{x}}$$

- (Small  $m$  = small power of  $k$  in Taylor series)
- Can write complete hydro response as set of functions

$$v_n e^{in\psi_n} = f(W_{m,n})$$

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# CUMULANT EXPANSION: TAYLOR SERIES IN $\varepsilon_{m,n}$

- Must make quantities with correct symmetries out of  $W_{m,n} \sim \langle r^m e^{in\phi} \rangle$
- If anisotropies are small, can arrange in Taylor series. E.g., to first order:

$$V_n \equiv v_n e^{in\psi_n} = \sum_{p=0}^{\infty} C_{n+2p,n} W_{n+2,p}$$

- If hydro is sensitive to large scale structure, first terms are most important (smaller powers of  $k$  in Fourier transform).
- If single term sufficient:

$$v_n e^{in\psi_n} = -C \frac{\langle r^n e^{in\phi} \rangle}{\langle r^n \rangle} \equiv C \varepsilon_n e^{in\phi_n}$$

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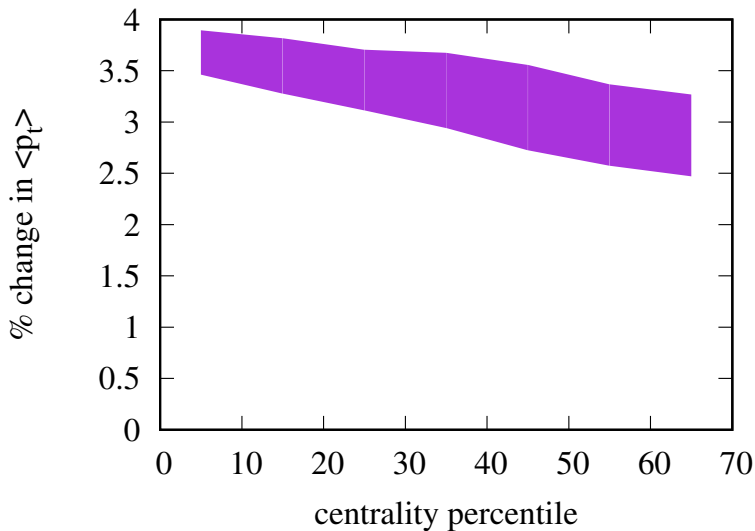
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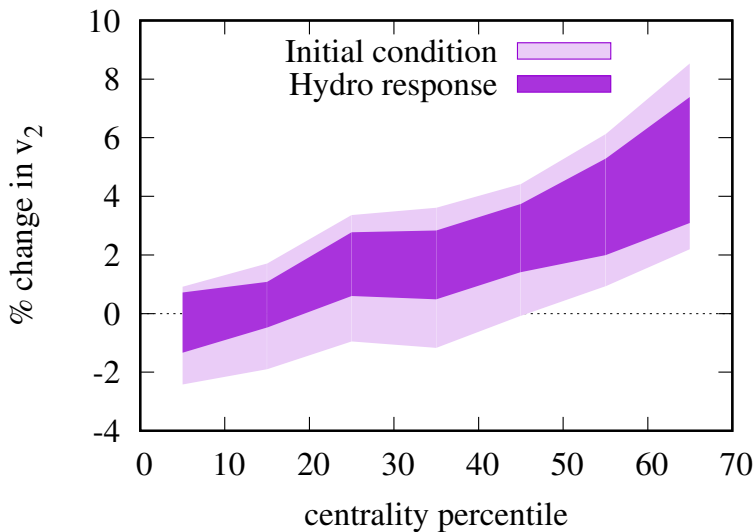
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## FINAL PREDICTION

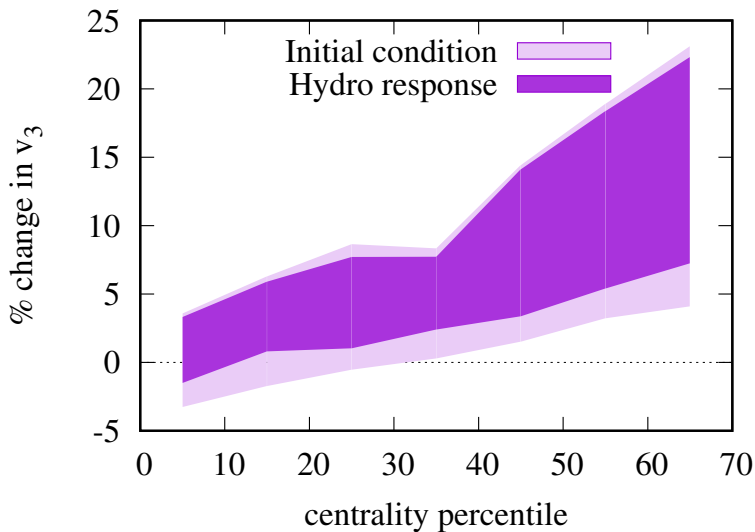




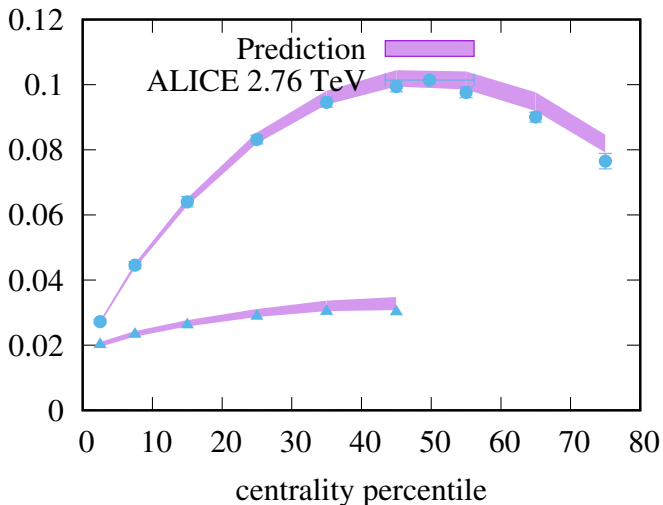
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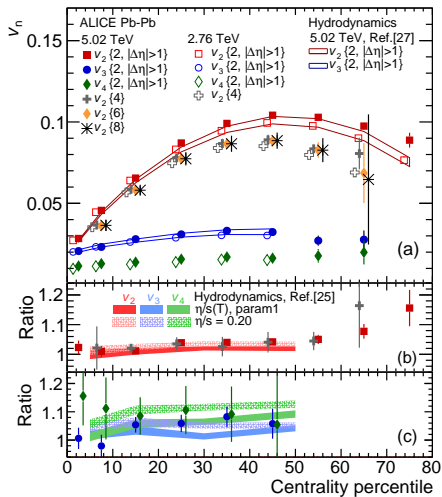
## FINAL PREDICTION



# PREDICTION PREVIEW: ABSOLUTE $v_2\{2\}$ AND $v_3\{2\}$

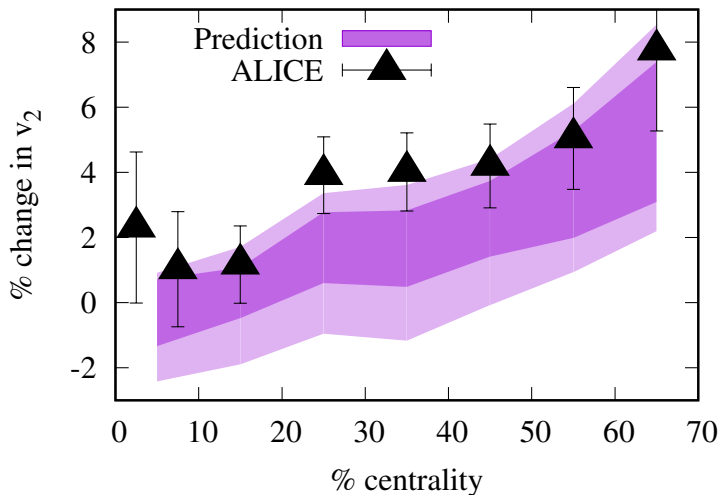


## FIRST DATA FROM RUN 2



(ALICE Collaboration, Phys.Rev.Lett. 116 (2016) 132302)

## COMPARISON WITH DATA



## COMPARISON WITH DATA

