

SIMP, Hierarchy Problem, Quintessence

Hitoshi Murayama (Berkeley, Kavli IPMU)
Niels Bohr Institute, Aug 13, 2018

+Eric Kuflik, Yonit Hochberg
Katelin Schutz, Robert McGehee
Chien-I Chang



BERKELEY LAB



KAVLI
IPMU

10th

*Anniversary Symposium
with Ceremonial Session*

Oct 16 (Mon) - 18 (Wed), 2017

Kashiwanoha Conference Center

Kavli IPMU



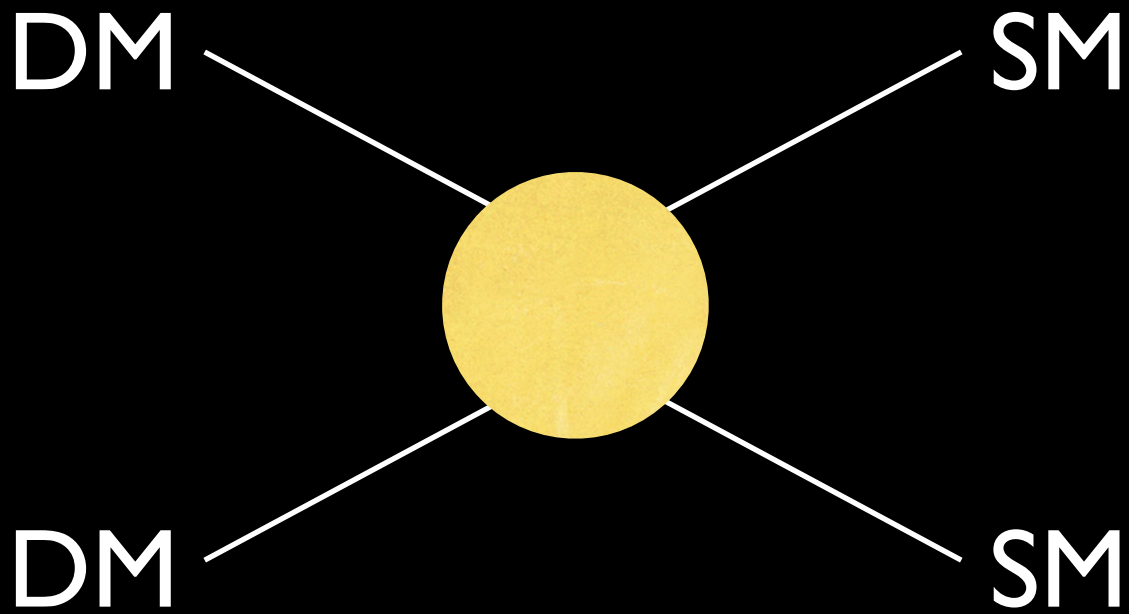
Kavli IPMU is there to stay!

Wither WIMP?



$$\frac{n_{\text{DM}}}{s} = 4.4 \times 10^{-10} \frac{\text{GeV}}{m_{\text{DM}}}$$

WIMP Miracle



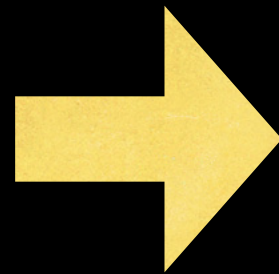
$$\langle \sigma_{2 \rightarrow 2\nu} \rangle \approx \frac{\alpha^2}{m^2}$$

$$\alpha \approx 10^{-2}$$

$$m \approx 300 \text{ GeV}$$

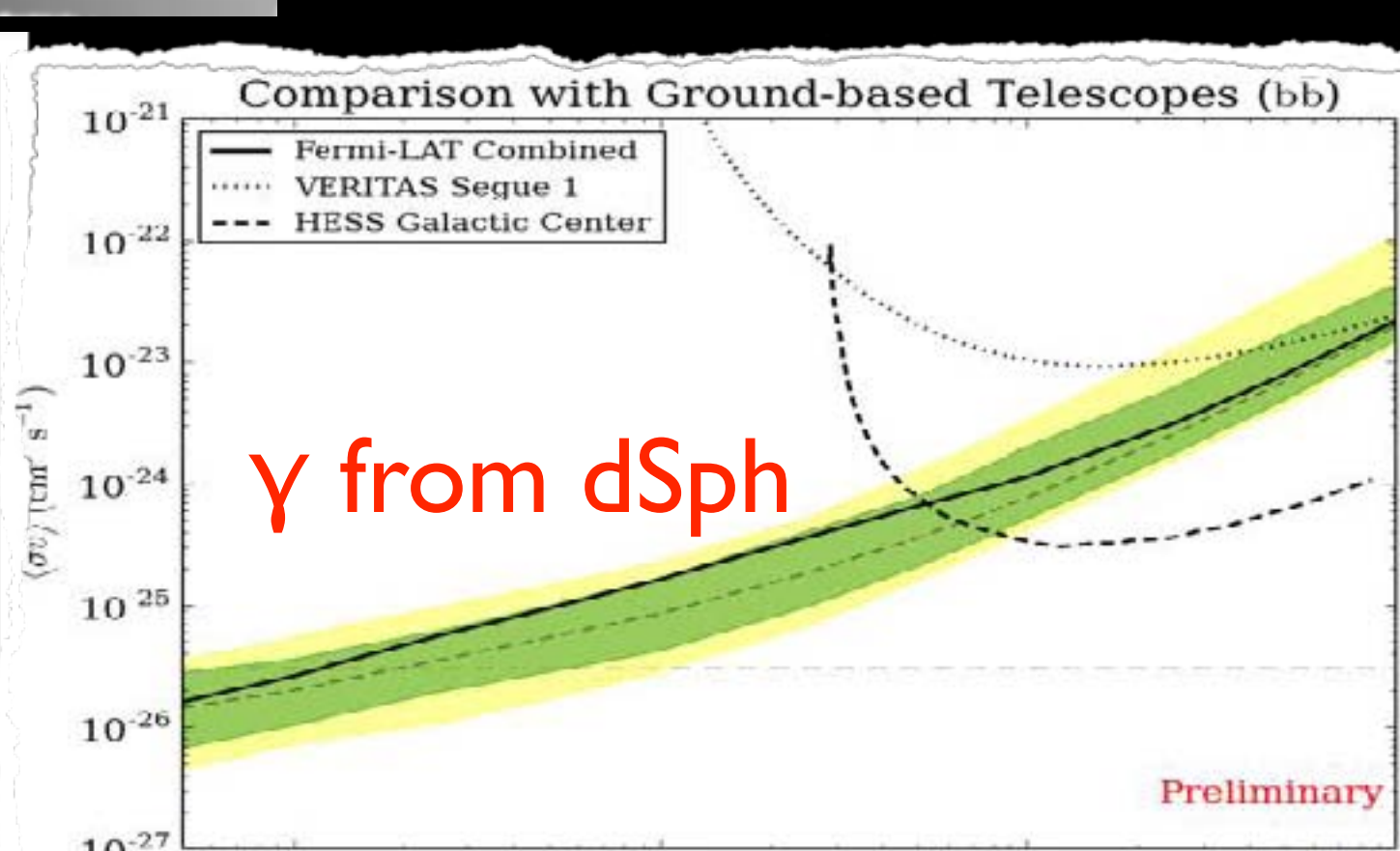
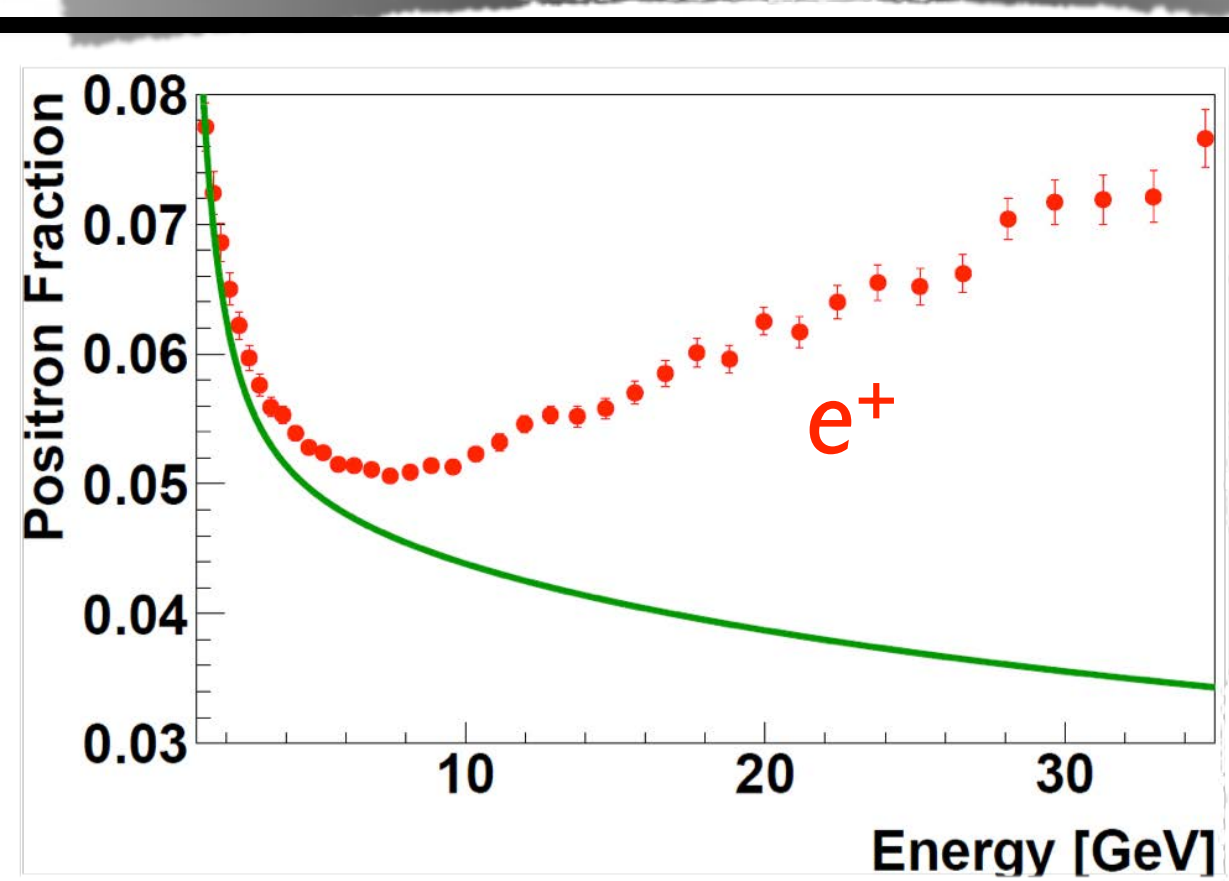
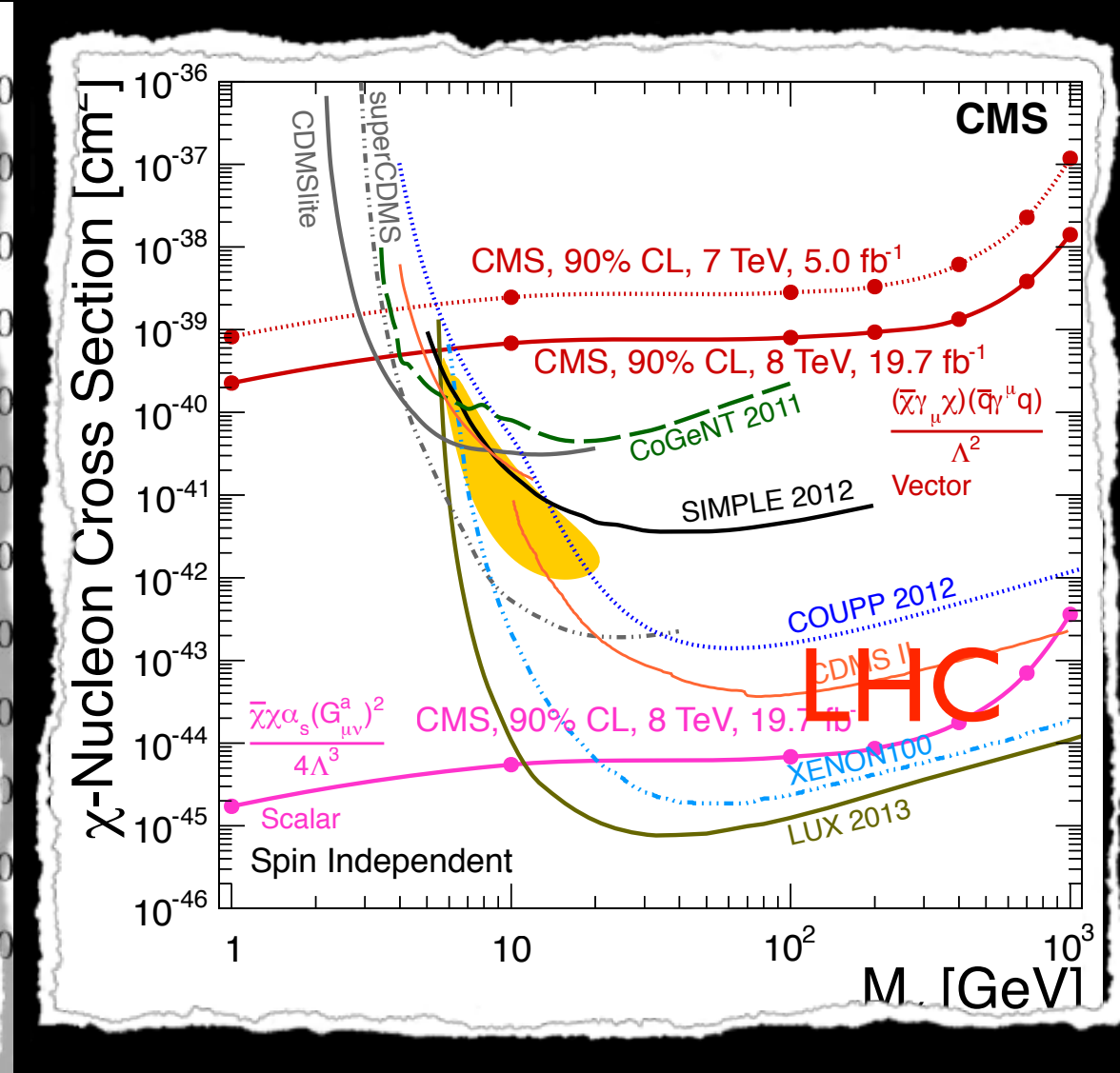
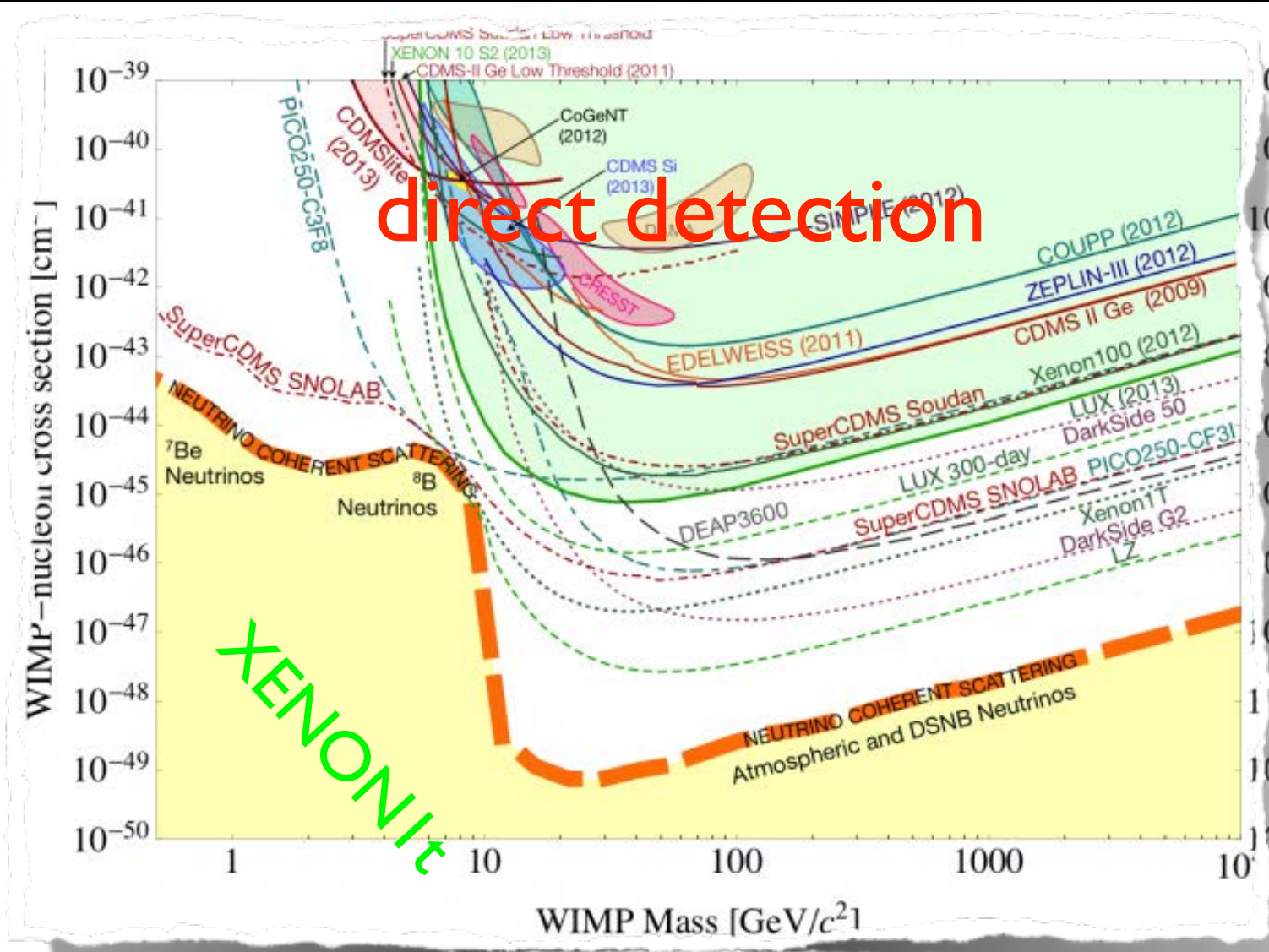
Relevant to
hierarchy problem?

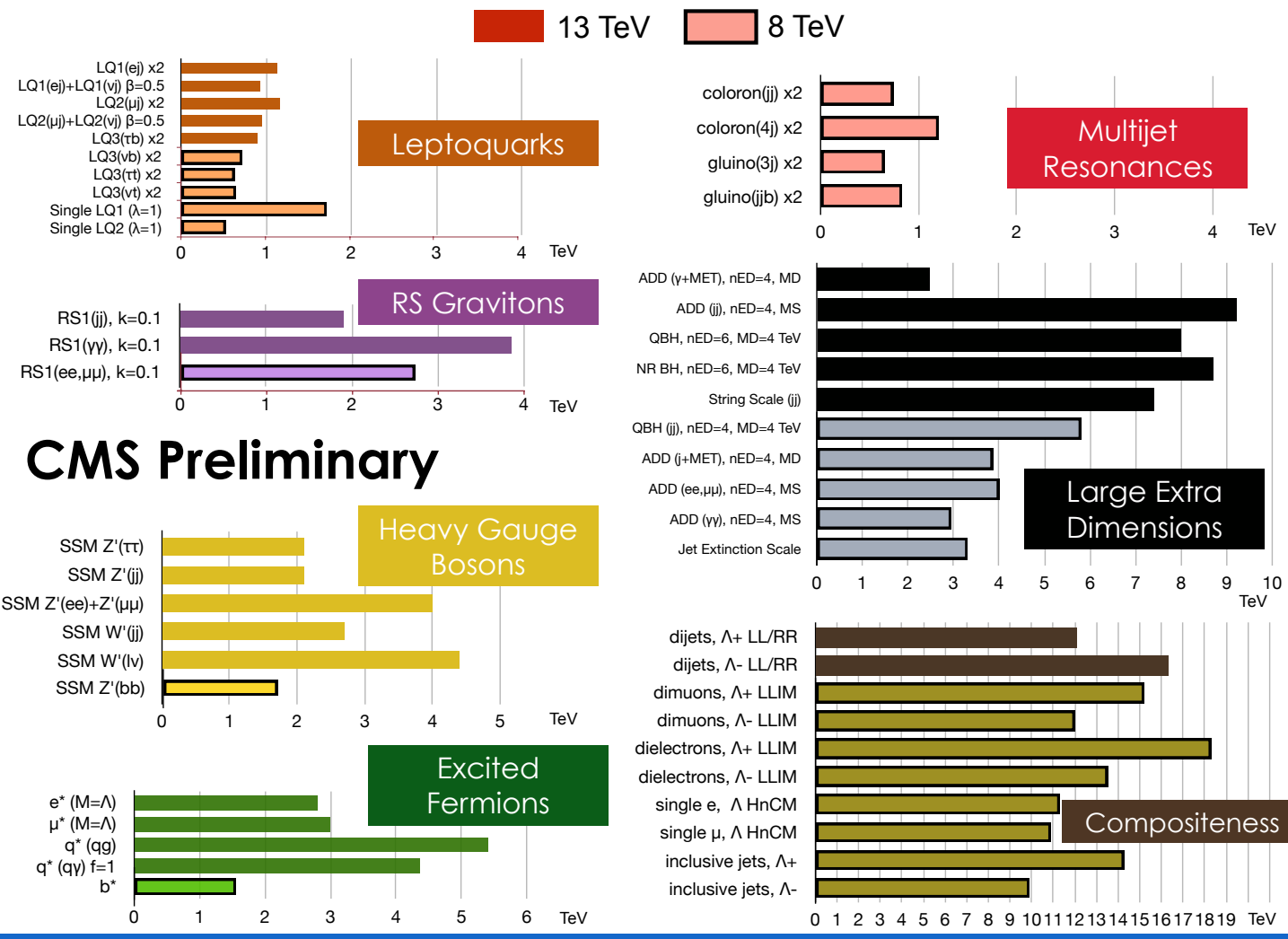
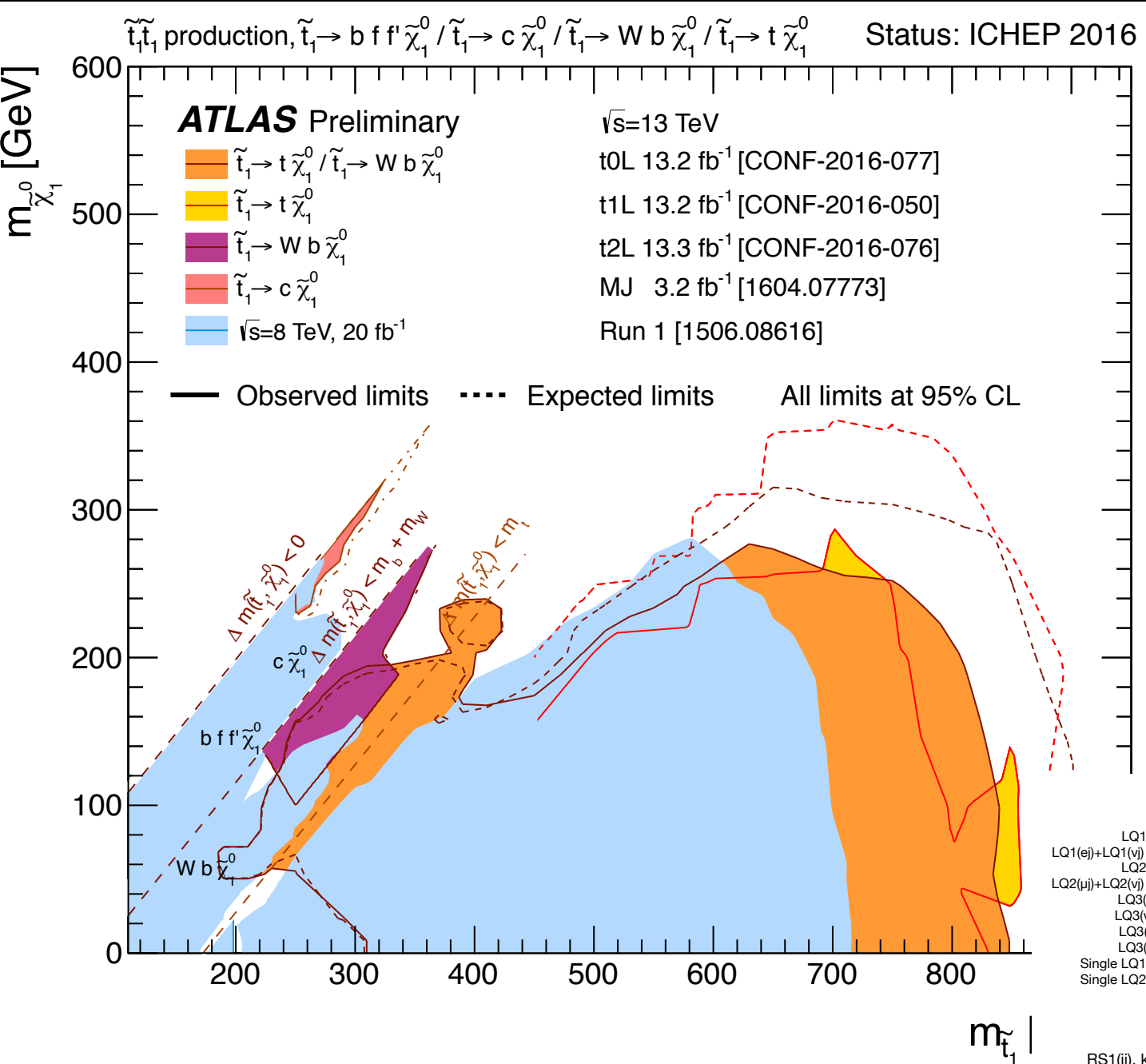
“weak” coupling
“weak” mass scale



correct abundance

Miracle²





no sign of new physics at TeV scale

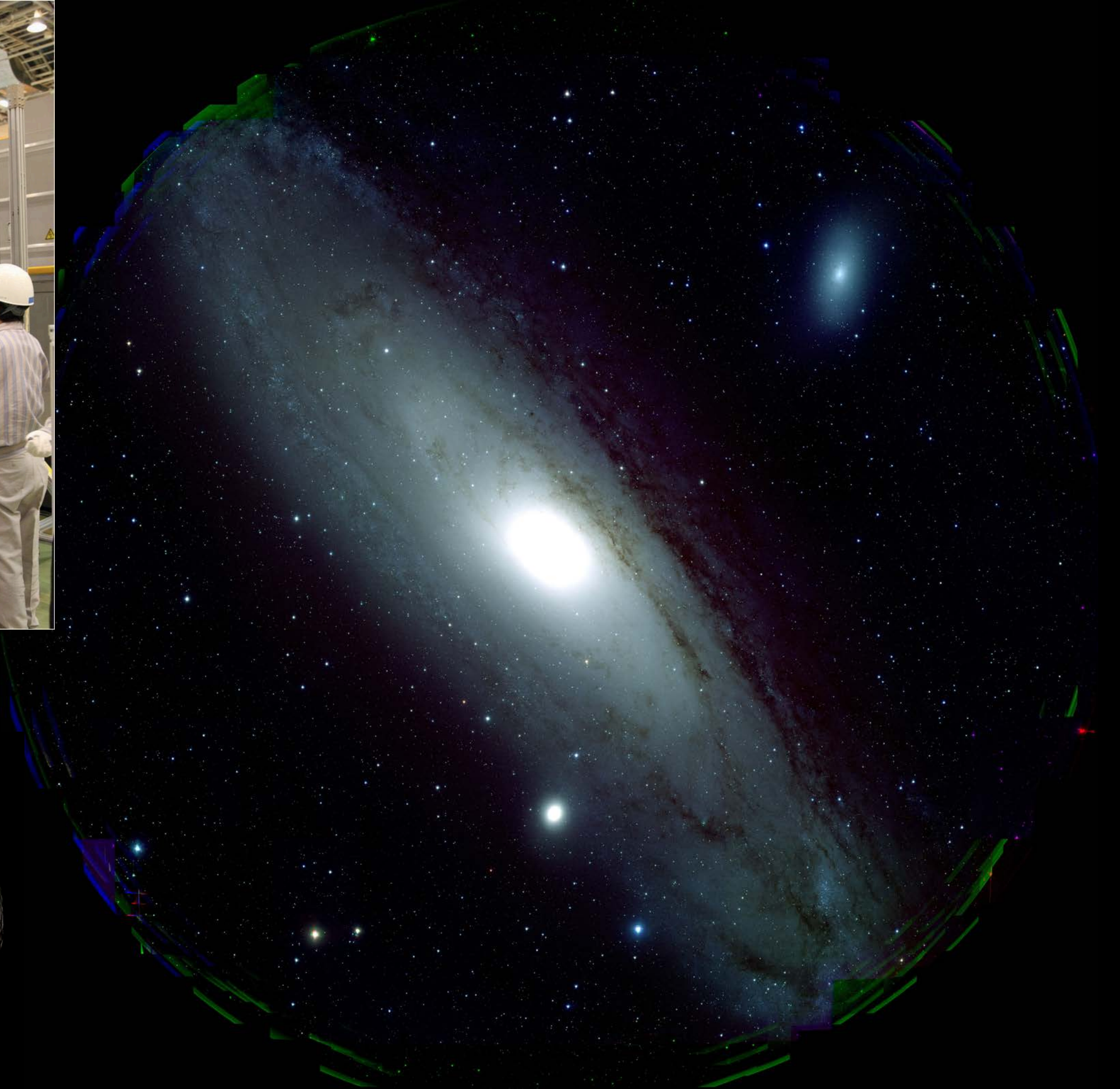
recent thinking

- dark matter definitely exists
 - hierarchy problem may be optional?
- need to explain dark matter on its own
- perhaps we should decouple these two
- do we really need big ideas like SUSY?
- perhaps we can solve it with ideas more familiar to us?

What do we know
about Dark Matter?



1.8 sq. deg.
870M pixels

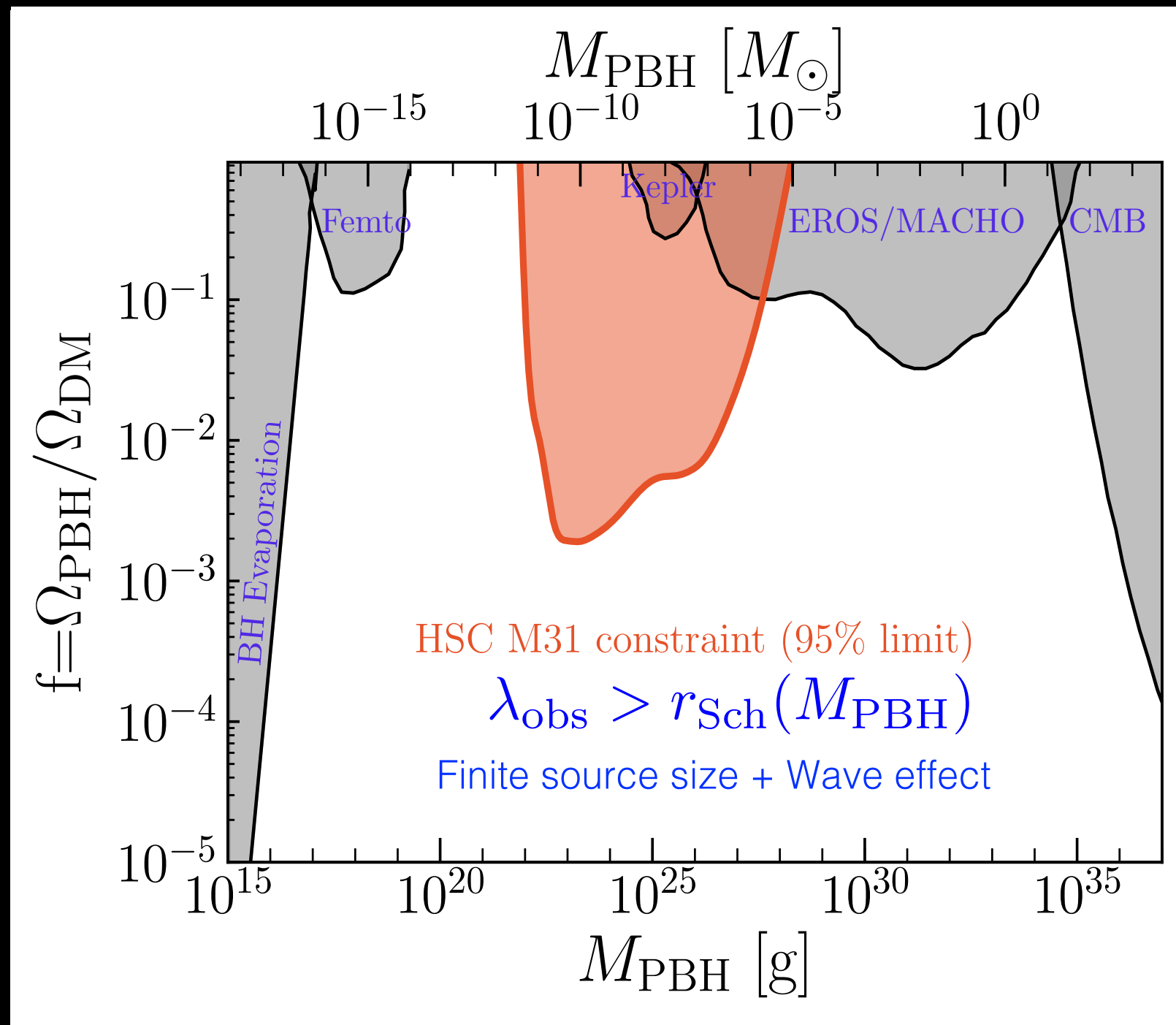


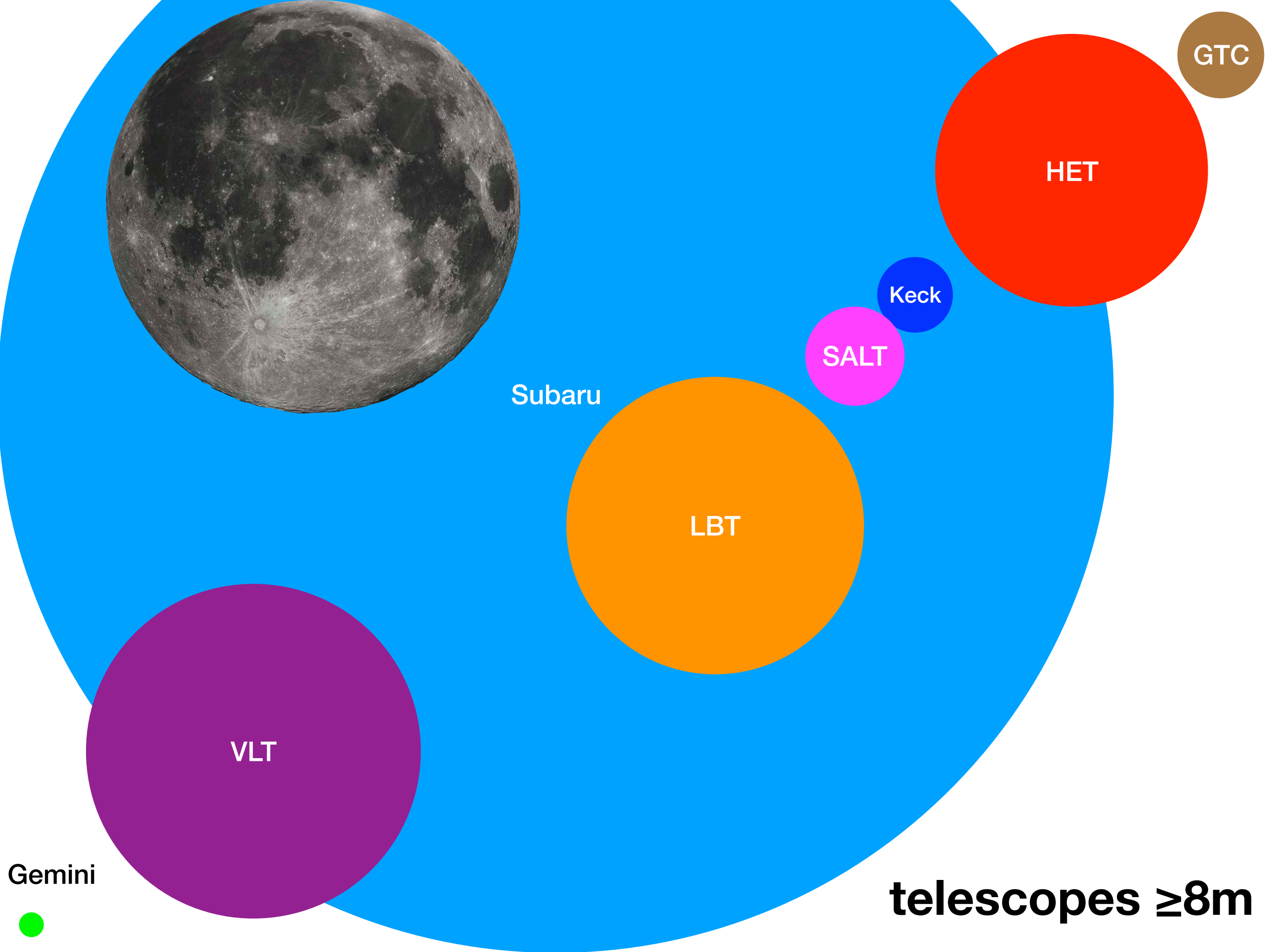
Best limit on Black Hole dark matter



Niikura, Takada et al. arXiv:1701.02151

Reading out every 2 minutes of HSC obs. of M31 to look for microlensing (just *one* night in Nov, 2015)





GTC

HET

Keck

SALT

Subaru

LBT

VLT

Gemini



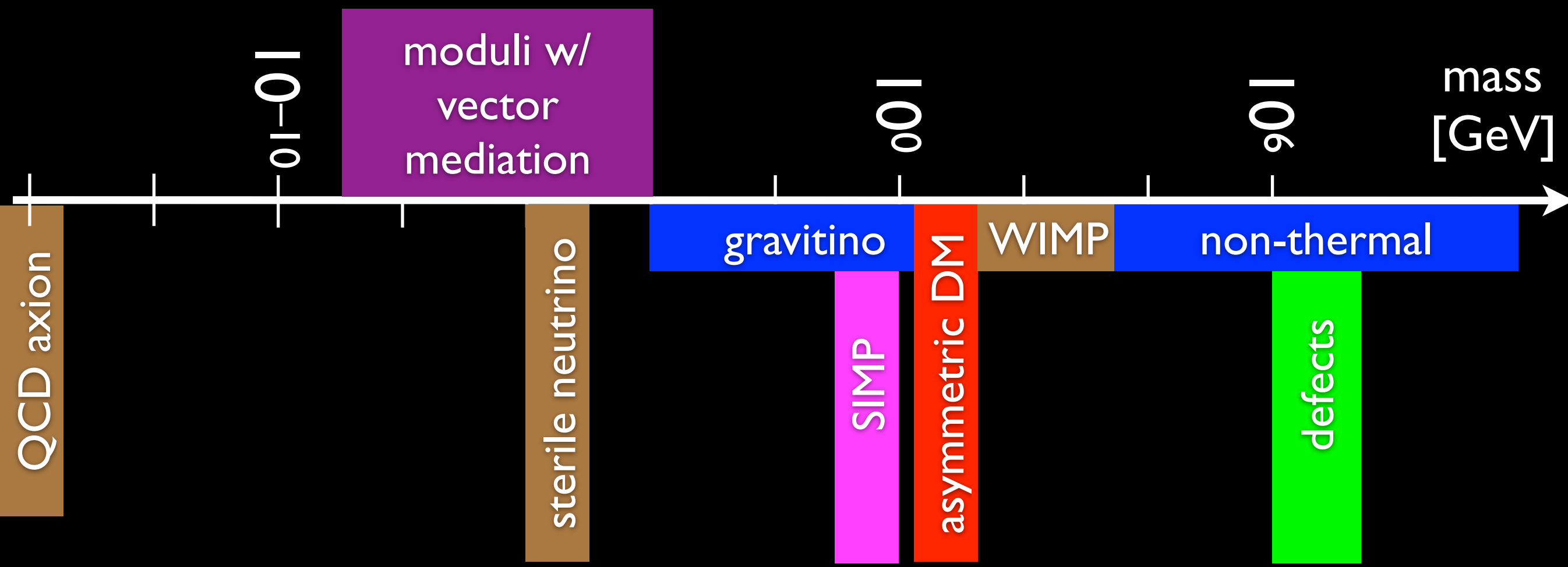
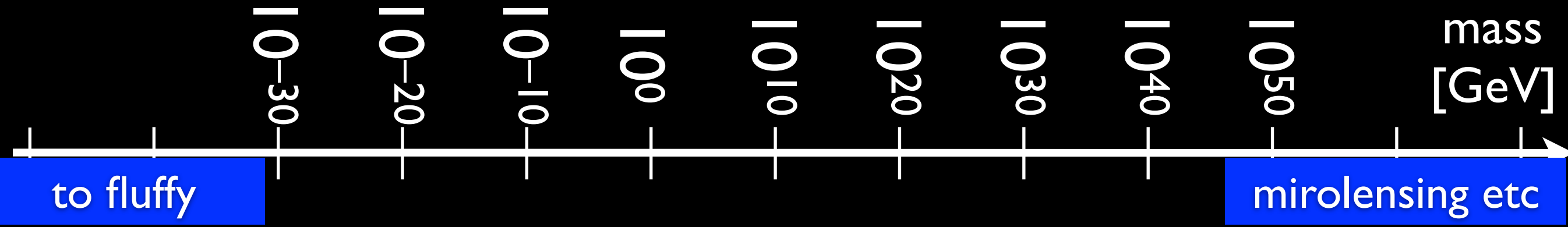
telescopes $\geq 8\text{m}$



Mass Limits

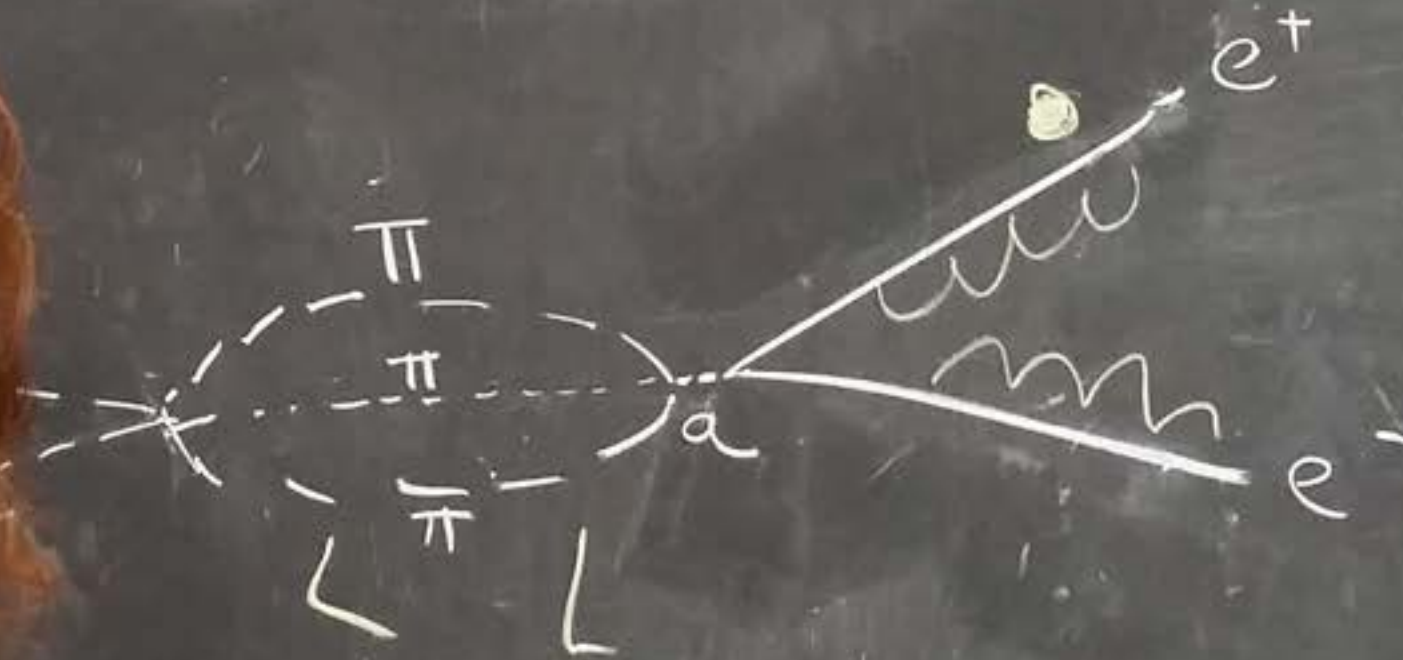
“Uncertainty Principle”

- Clumps to form structure
- imagine $V = G_N \frac{Mm}{r}$
- “Bohr radius”: $r_B = \frac{\hbar^2}{G_N M m^2}$
- too small $m \Rightarrow$ won’t “fit” in a galaxy!
- $m > 10^{-22}$ eV “uncertainty principle” bound
(modified from Hu, Barkana, Gruzinov, astro-ph/0003365)



Seminar in Berkeley

Strongly Interacting Massive Particle (SIMP)

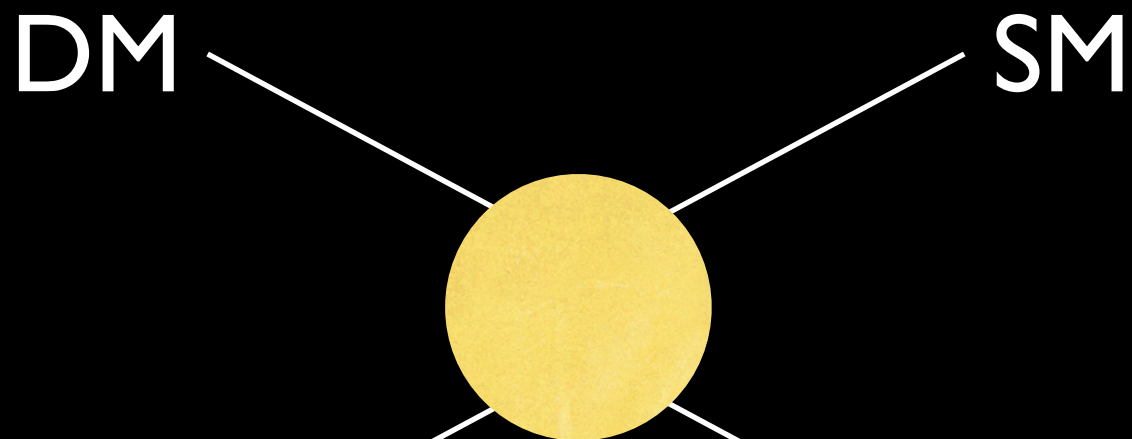


Yonit Hochberg



$$\frac{n_{\text{DM}}}{s} = 4.4 \times 10^{-10}$$

Miracles

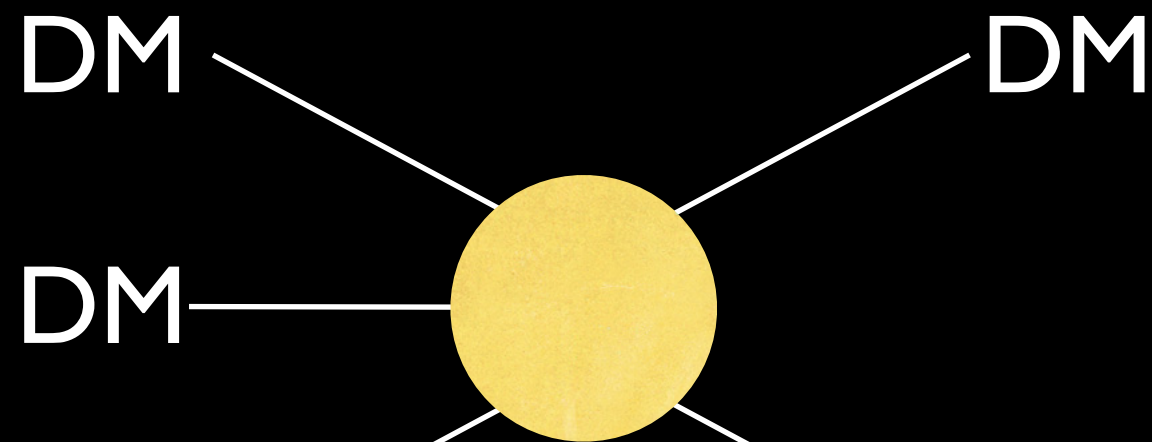
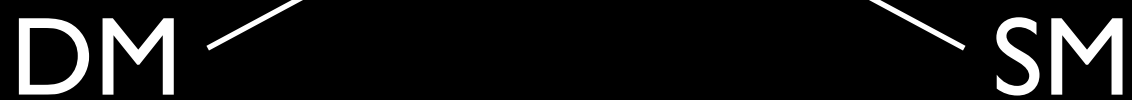


$$\langle \sigma_{2 \rightarrow 2\nu} \rangle \approx \frac{\alpha^2}{m^2}$$

$$\alpha \approx 10^{-2}$$

$$m \approx 300 \text{ GeV}$$

WIMP miracle!



$$\langle \sigma_{3 \rightarrow 2\nu^2} \rangle \approx \frac{\alpha^3}{m^5}$$

$$\alpha \approx 4\pi$$

Hochberg, Kuflik,
Volansky, Wacker

$$m \approx 300 \text{ MeV}$$

arXiv:1402.5143

SIMP miracle!

3 → 2

~~LEE WEINBERG~~ FREEZE-OUT

Back of the envelope calculation

$$\Gamma_{\text{ann}} \simeq H|_{\text{freezeout}}$$

$$\Gamma_{\text{ann}} \simeq n_{\text{dm}}^2 \langle \sigma v^2 \rangle_{3 \rightarrow 2}$$

$$H \simeq \frac{T^2}{M_{\text{pl}}}$$

$$m_{\text{dm}} n_{\text{dm}} \sim m_p n_b$$

$$\langle \sigma v^2 \rangle_{3 \rightarrow 2} \simeq \frac{\alpha^3}{m_{\text{dm}}^2}$$

$$n_b \sim \eta_b s$$

Eric Kuflik

$$\eta_b \simeq T_{\text{eq}}/m_p$$

$$s \simeq T^3$$

$$\Gamma_{\text{ann}} \simeq \frac{T_{\text{eq}}^2 \alpha^3}{x_F^3 \times m_{\text{dm}}} =$$

$$H \simeq \frac{m_{\text{dm}}^2}{M_{\text{pl}} x_F^2}$$

THE SIMP MIRACLE

A coincidence of scales

$$m_{\text{dm}} \simeq \alpha \left(T_{\text{eq}}^2 M_{\text{pl}} / x_F^4 \right)^{\frac{1}{3}}$$

$$m_{\text{dm}} \simeq \alpha \times 100 \text{ MeV}$$

- If $\alpha \sim 1$, the strong scale emerges ($x_F \sim 20$)
- Like the WIMP, no input of scales or particle physics Eric Kuflik

Strongly interacting sub-GeV dark matter

But not UV complete; coupling blows up very quickly

Dark Pions



$$\frac{n_{DM}}{s} = 4.4 \times 10^{-10}$$

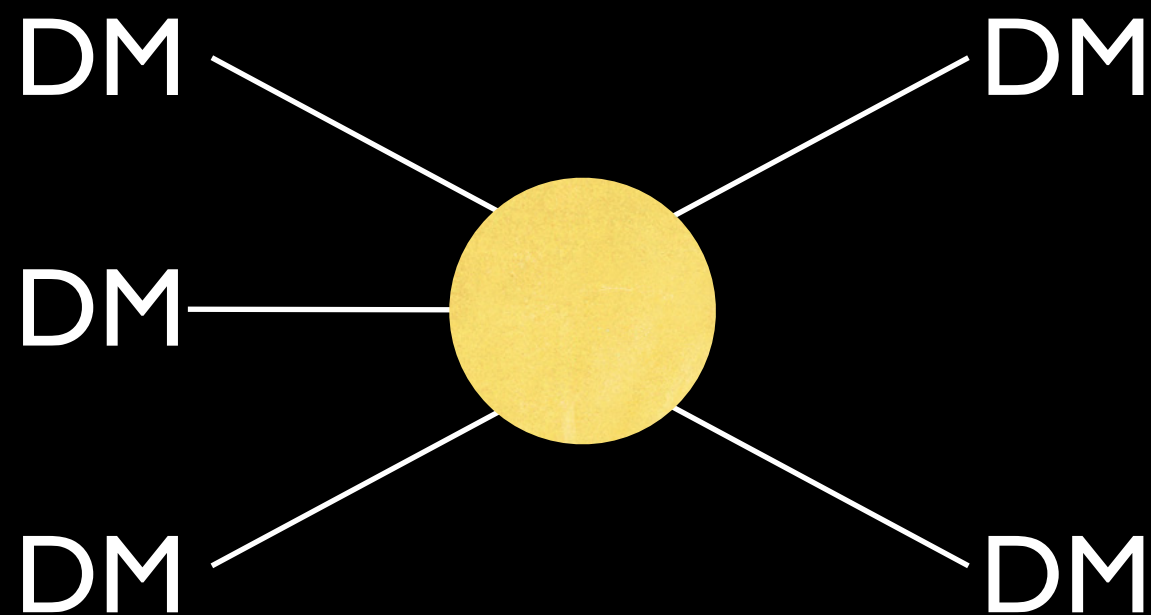
SIMPLest Miracle

- Not only the mass scale is similar to QCD
- dynamics itself can be QCD! Miracle³
- DM = pions = $q\bar{q}$
- e.g. $SU(4)/Sp(4) = S^5$

$$\mathcal{L}_{\text{chiral}} = \frac{1}{16f_\pi^2} \text{Tr} \partial^\mu U^\dagger \partial_\mu U$$

$$\mathcal{L}_{\text{WZW}} = \frac{8N_c}{15\pi^2 f_\pi^5} \epsilon_{abcde} \epsilon^{\mu\nu\rho\sigma} \pi^a \partial_\mu \pi^b \partial_\nu \pi^c \partial_\rho \pi^d \partial_\sigma \pi^e + O(\pi^7)$$

$$\pi_5(G/H) \neq 0$$

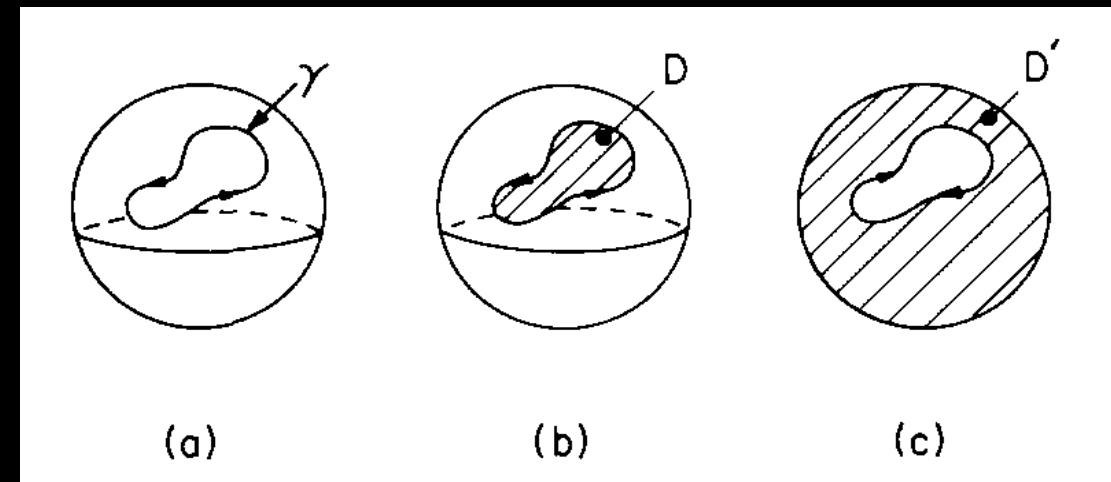


+HM

arXiv:1411.3727

Wess-Zumino term

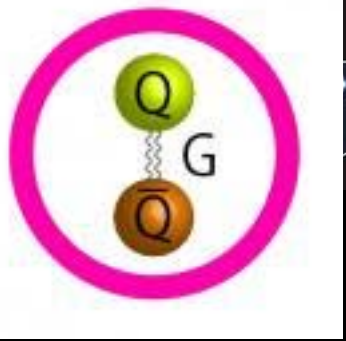
- $\mathcal{L}_{WZ} \sim \epsilon_{abcde} \epsilon^{\mu\nu\rho\sigma} \pi^a \partial_\mu \pi^b \partial_\nu \pi^c \partial_\rho \pi^d \partial_\sigma \pi^e$
- $SU(N_c)$ gauge theory
 - $\pi_5(SU(N_f)) = \mathbb{Z}$ ($N_f \geq 3$)
- $Sp(N_c)$ gauge theory
 - $\pi_5(SU(2N_f)/Sp(N_f)) = \mathbb{Z}$ ($N_f \geq 2$)
- $SO(N_c)$ gauge theory
 - $\pi_5(SU(N_f)/SO(N_f)) = \mathbb{Z}$ ($N_f \geq 3$)



Witten

also, non-abelian vector bosons (vector SIMP)

+S-M Choi, HM Lee, Y. Mambrini, M. Pierre, arXiv:1707.01434



SIMPLest Miracle

$$\begin{pmatrix} u_1 \\ d_1 \end{pmatrix} \begin{pmatrix} u_2 \\ d_2 \end{pmatrix} \begin{pmatrix} u_3 \\ d_3 \end{pmatrix} \begin{pmatrix} u_4 \\ d_4 \end{pmatrix}$$

- SU(2) gauge theory with four doublets
- SU(4)=SO(6) flavor symmetry
- $\langle q^i q^j \rangle \neq 0$ breaks it to Sp(2)=SO(5)
- coset space SO(6)/SO(5)=S⁵, 5 pions
- $\pi_5(S^5)=\mathbb{Z} \Rightarrow$ Wess-Zumino term
- $\mathcal{L}_{WZ} = \epsilon_{abcde} \epsilon^{\mu\nu\rho\sigma} \pi^a \partial_\mu \pi^b \partial_\nu \pi^c \partial_\rho \pi^d \partial_\sigma \pi^e$

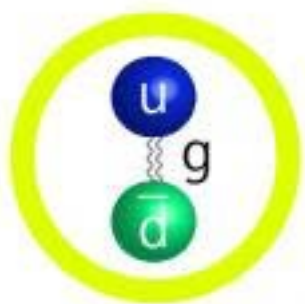
LAGRANGIANS

Quark theory

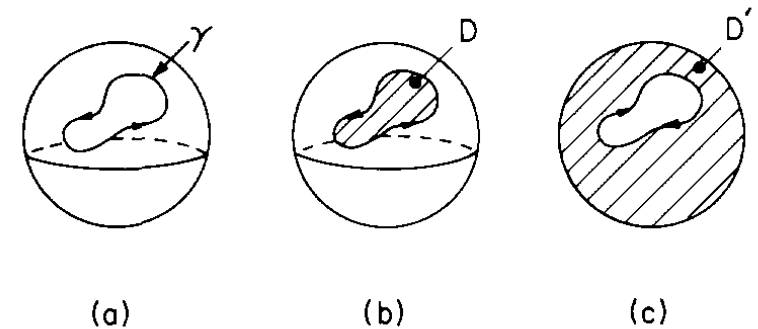
$$\mathcal{L}_{\text{quark}} = -\frac{1}{4} F_{\mu\nu}^a F^{\mu\nu a} + \bar{q}_i i \not{D} q_i - \frac{1}{2} m_Q J^{ij} q_i q_j + h.c.$$

Sigma theory

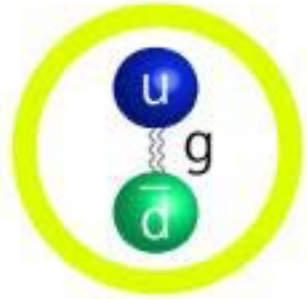
$$\mathcal{L}_{\text{Sigma}} = \frac{f_\pi^2}{16} \text{Tr} \partial_\mu \Sigma \partial^\mu \Sigma^\dagger - \frac{1}{2} m_Q \mu^3 \text{Tr} J \Sigma + h.c. - \frac{i N_c}{240 \pi^2} \int \text{Tr} (\Sigma^\dagger d\Sigma)^5$$



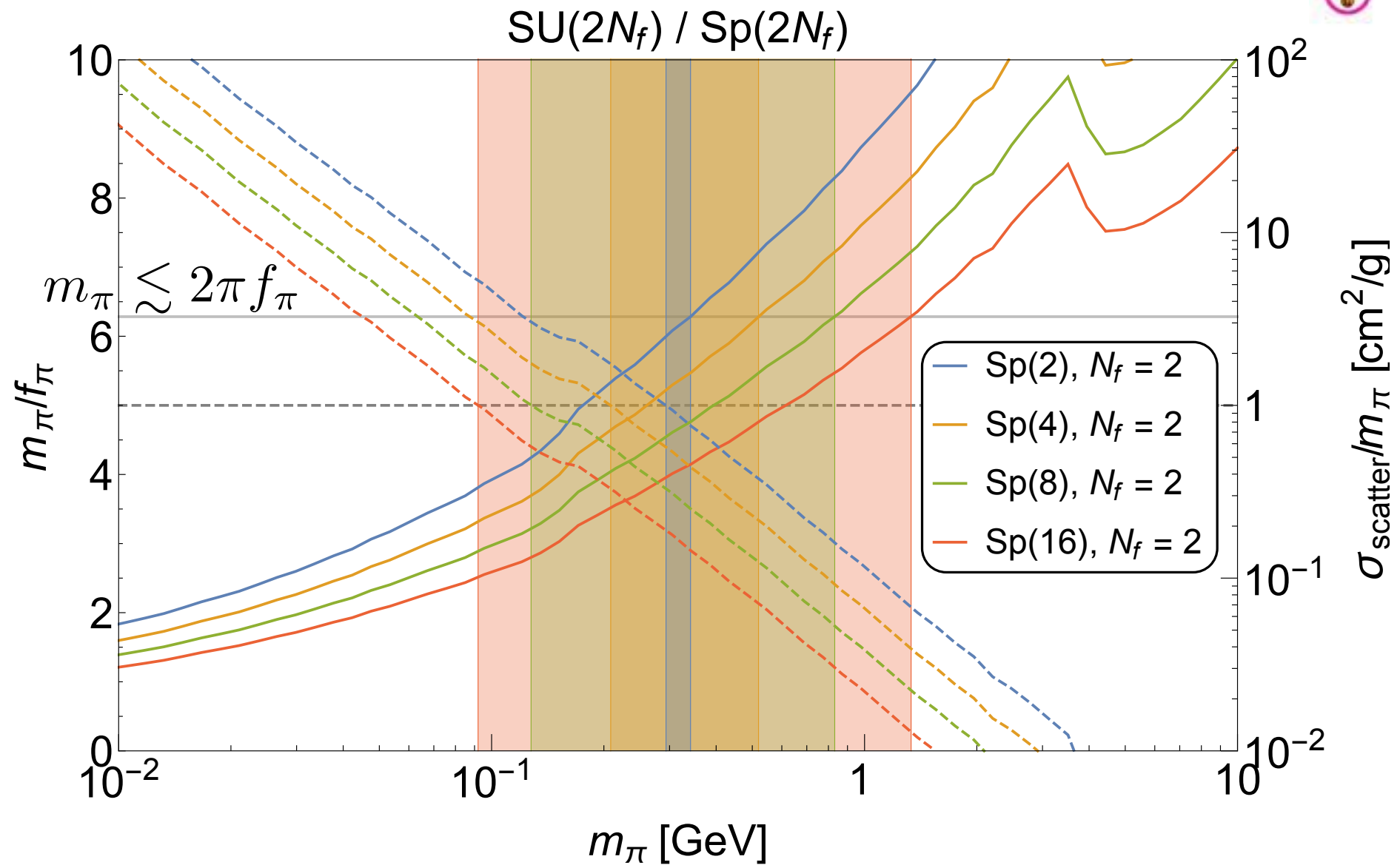
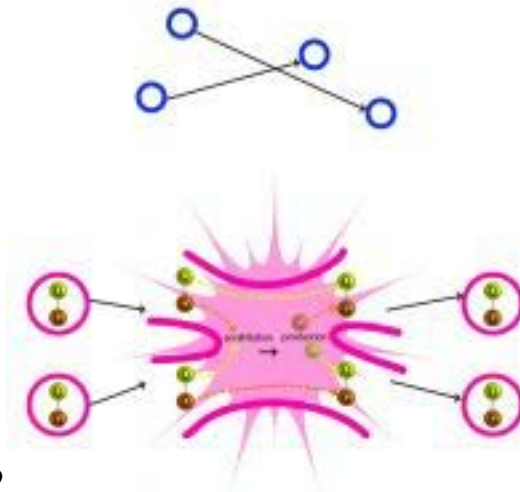
Pion theory



$$\mathcal{L}_{\text{pion}} = \frac{1}{4} \text{Tr} \partial_\mu \pi \partial^\mu \pi - \frac{m_\pi^2}{4} \text{Tr} \pi^2 + \frac{m_\pi^2}{12 f_\pi^2} \text{Tr} \pi^4 - \frac{1}{6 f_\pi^2} \text{Tr} (\pi^2 \partial^\mu \pi \partial_\mu \pi - \pi \partial^\mu \pi \pi \partial_\mu \pi) + \frac{2 N_c}{15 \pi^2 f_\pi^5} \epsilon^{\mu\nu\rho\sigma} \text{Tr} [\pi \partial_\mu \pi \partial_\nu \pi \partial_\rho \pi \partial_\sigma \pi] + \mathcal{O}(\pi^6)$$



The Results

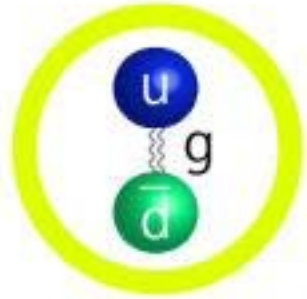


Solid curves: solution to Boltzmann eq.

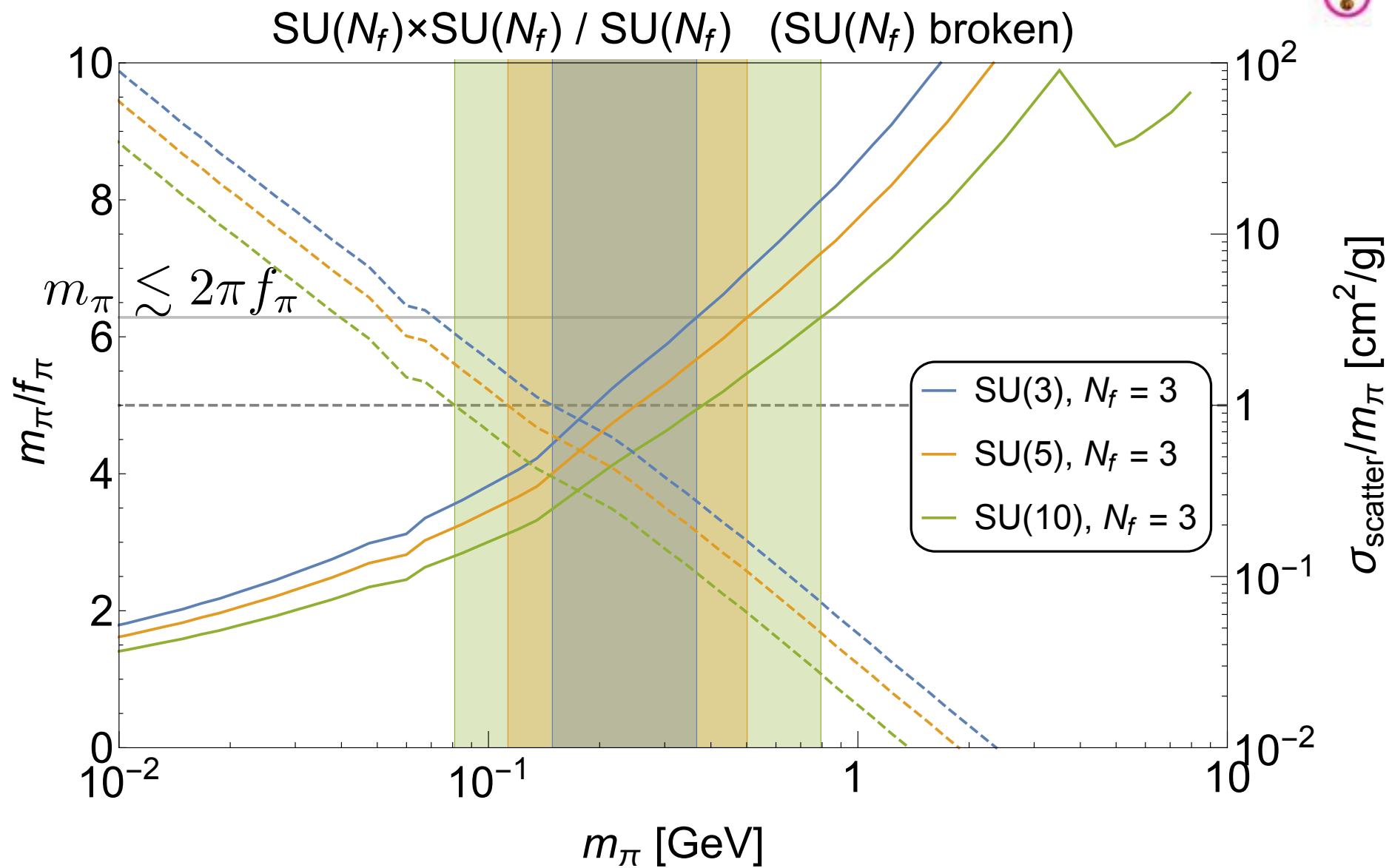
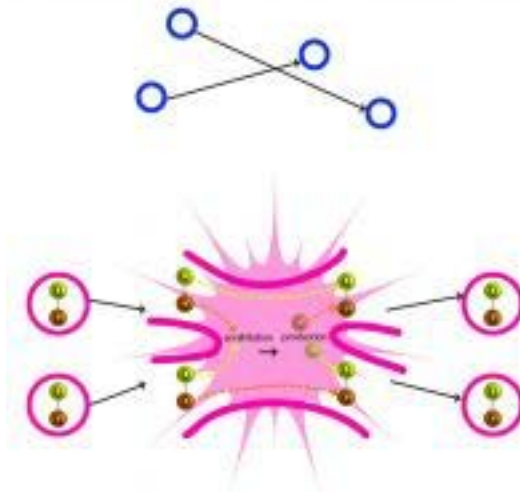
Dashed curves: along that solution

$$\frac{m_\pi}{f_\pi} \propto m_\pi^{3/10}$$

$$\frac{\sigma_{scatter}}{m_\pi} \propto m_\pi^{-9/5}$$



The Results

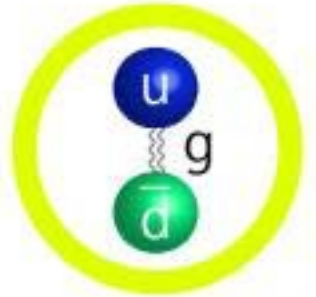


Solid curves: solution to Boltzmann eq.

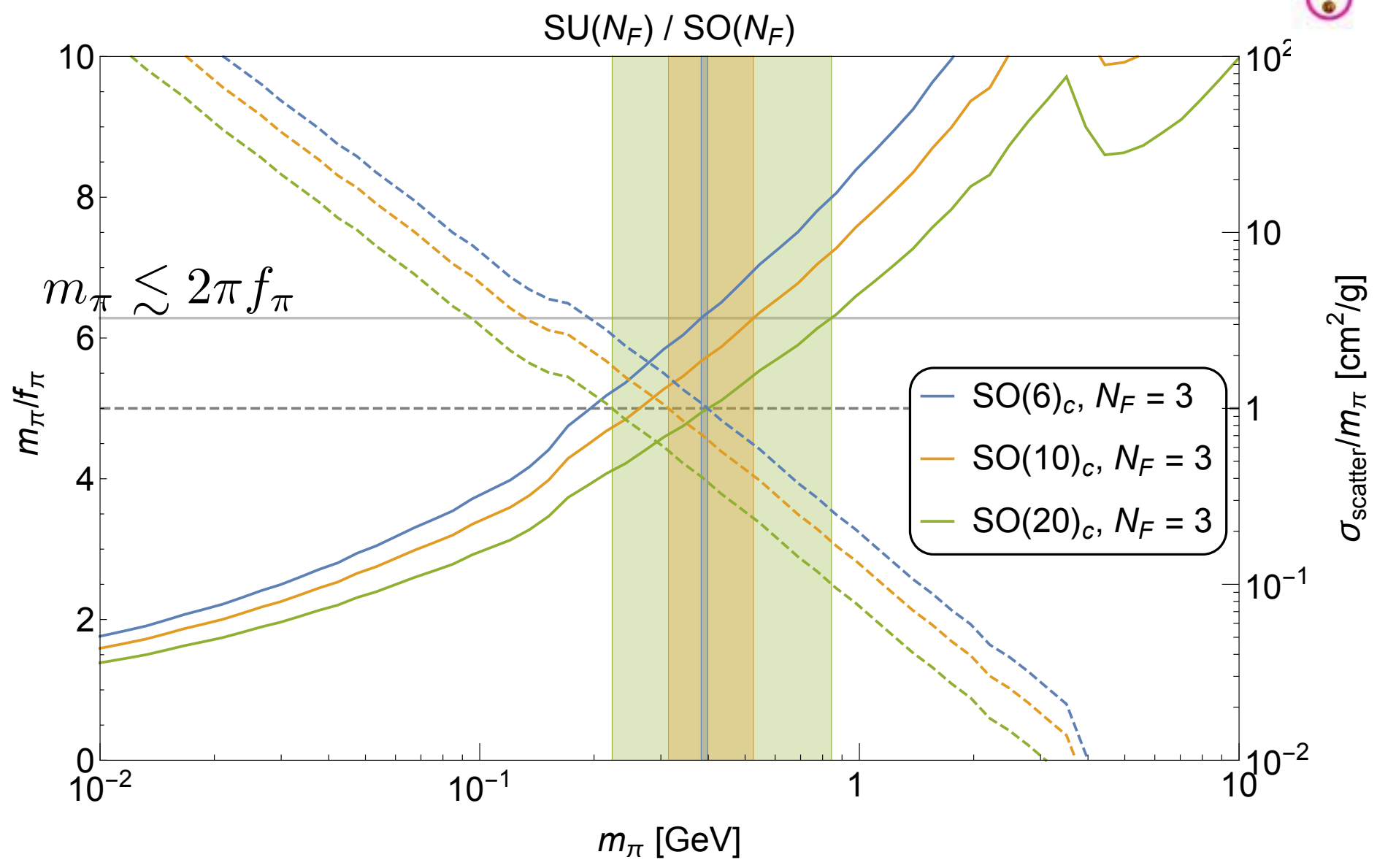
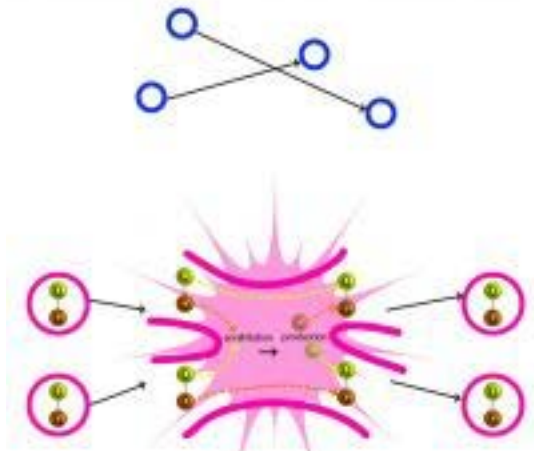
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The Results



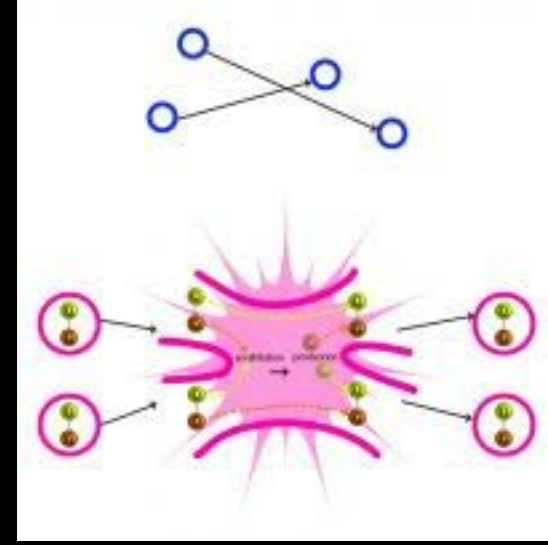
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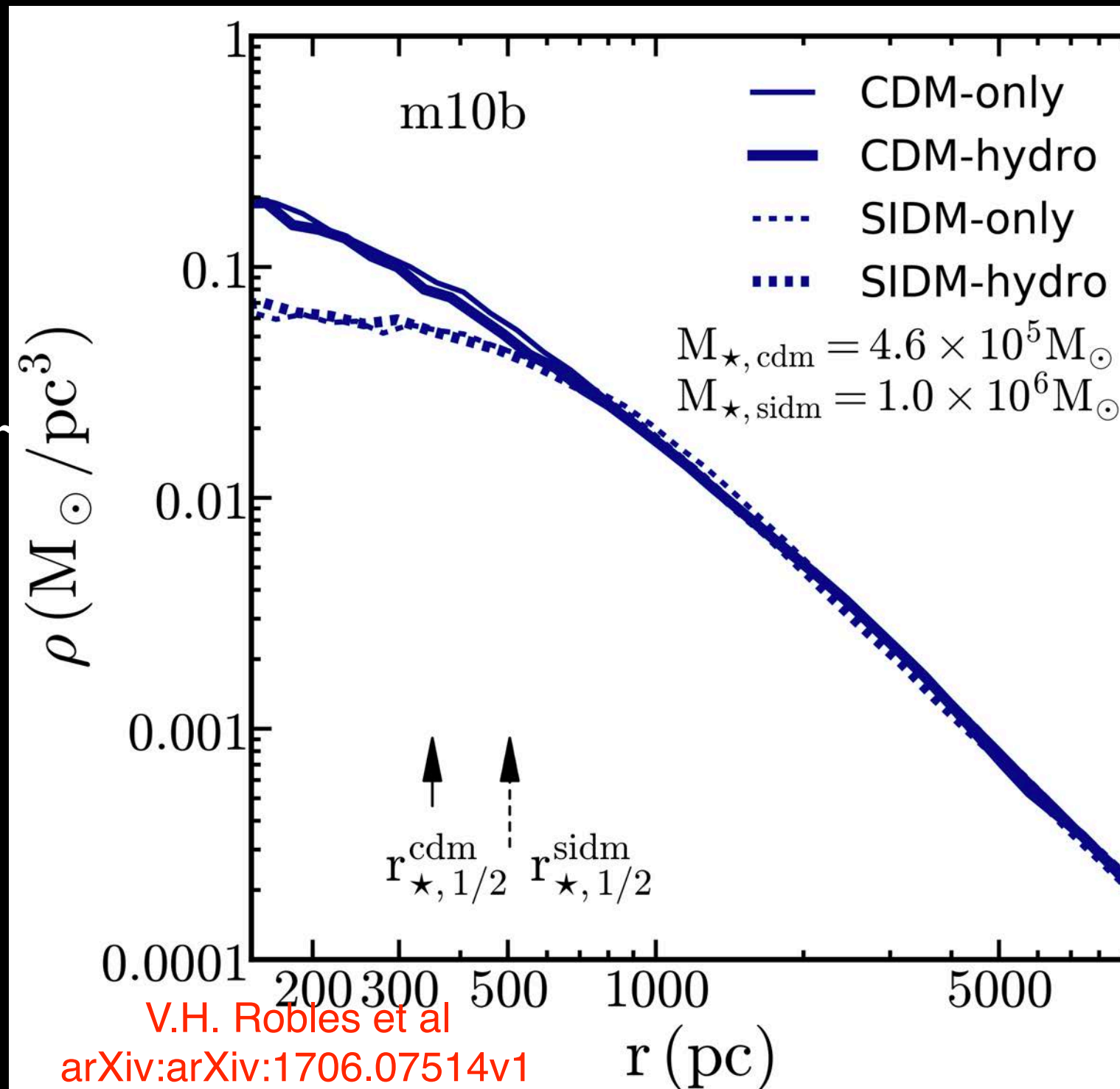
$$\frac{\sigma_{\text{scatter}}}{m_\pi} \propto m_\pi^{-9/5}$$

self interaction

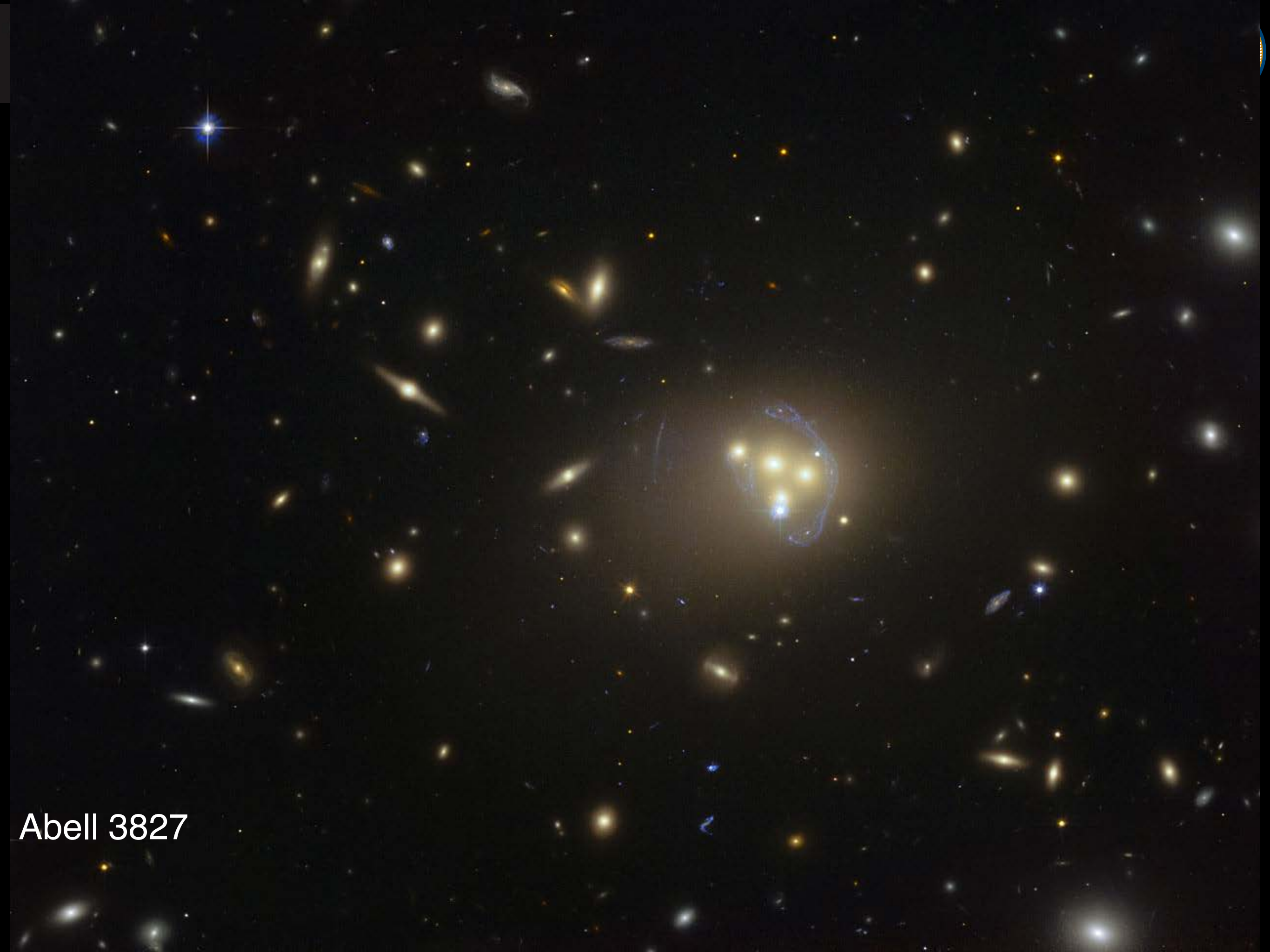


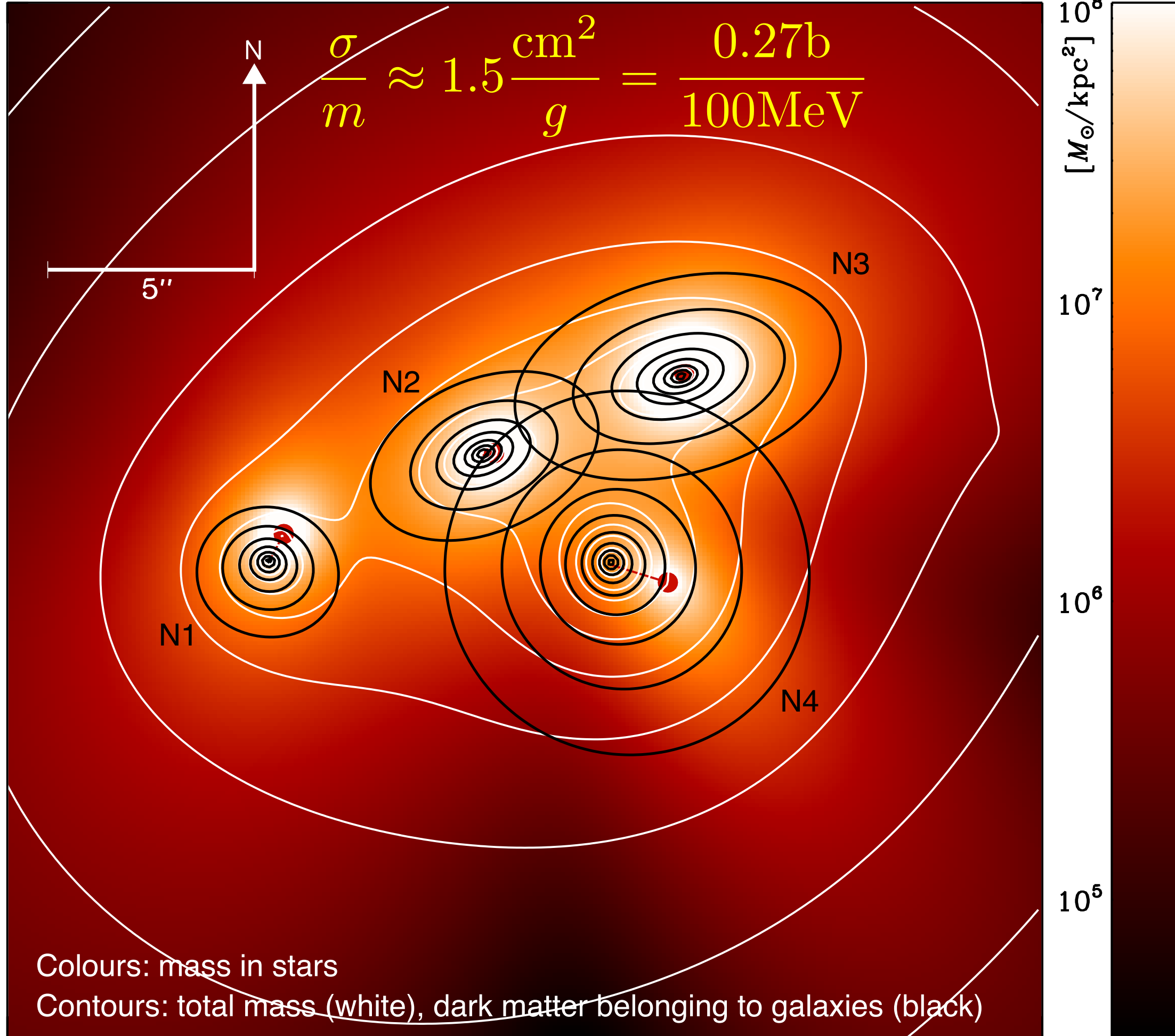
- $\sigma/m \sim \text{cm}^2/g$
 $\sim 10^{-24} \text{cm}^2 / 300 \text{MeV}$
- flattens the cusps in NFW profile
- suppresses substructure
- actually desirable for dwarf galaxies?

SIDM
Spergel & Steinhardt
(2000)
now complete theory



Abell 3827



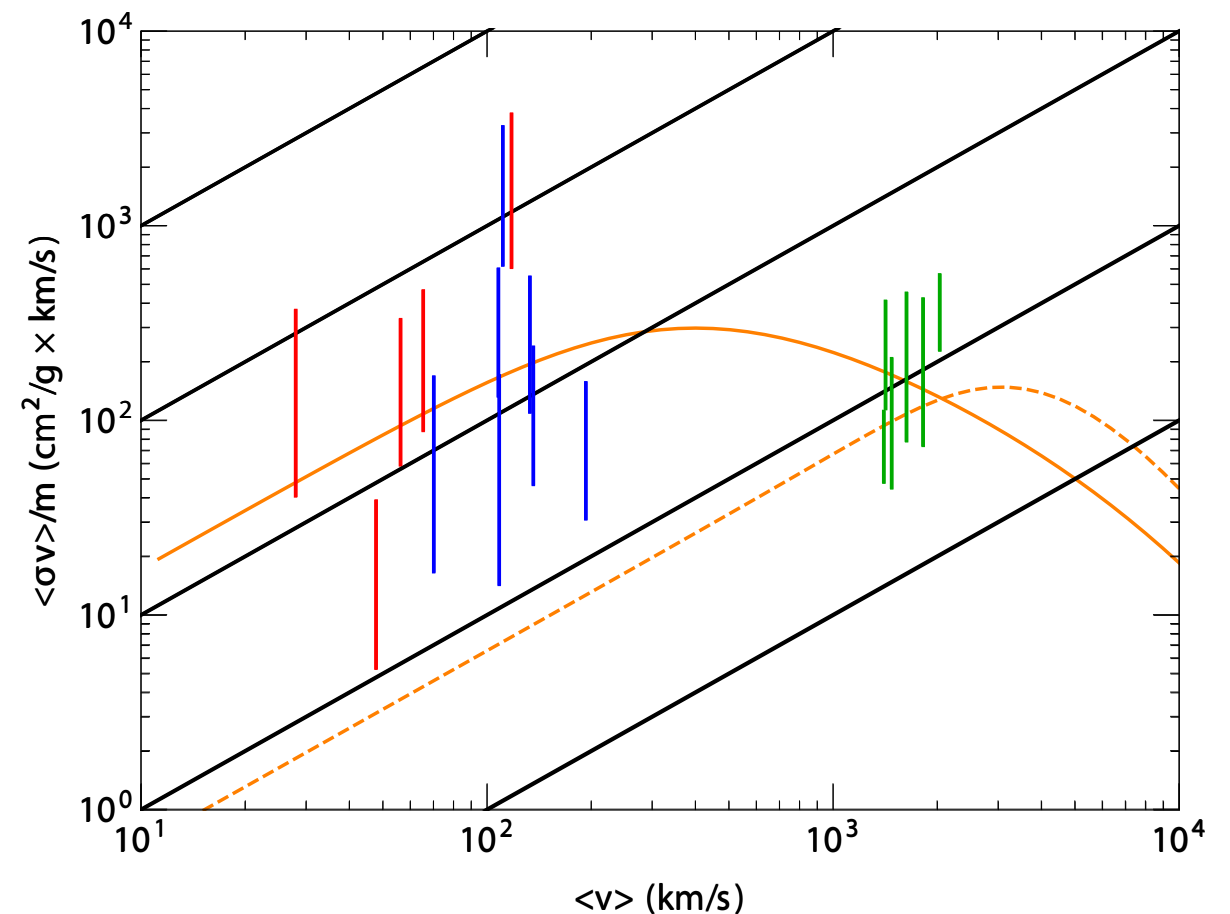


velocity dependence?

- cluster data prefer smaller σ ?
- near constant $\langle\sigma v\rangle$?
- Sommerfeld effect (S. Tulin, H.-B. Yu, and K.M. Zurek, arXiv:1302.3898)
 - requires light mediator
- near-threshold resonance can “fit” the data
- *i.e.*, $\pi\pi \rightarrow \sigma \rightarrow \pi\pi$
 - (Xiaoyong Chu, Camilo Garcia-Cely, Yonit Hochberg, Eric Kuik, HM)

$$\mathcal{L} = m_R g RDM^2 .$$

n	a	b	γ_0	v_R (km/s)	m_R (GeV)	$\chi^2/\text{d.o.f.}$
0	24	32	$10^{-4.3}$	1829	26	2.1



M. Kaplinghat, S. Tulin, and H.-B. Yu,
arXiv:1508.03339.

PFS subsystems distribution

Software system

Spectrograph system (SpS)

... in Prime focus unit "POpt2" with Wide Field Corrector "WFC".

Calibration system

Fiber connectors

Fiber cable

Prime Focus Instrument

Wide-field corrector



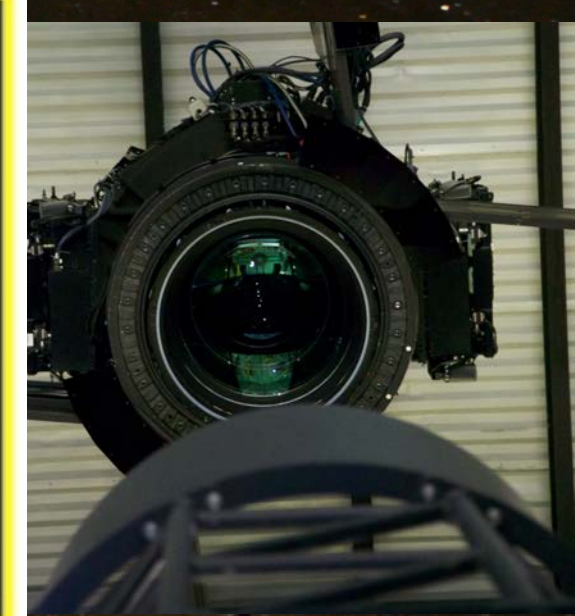
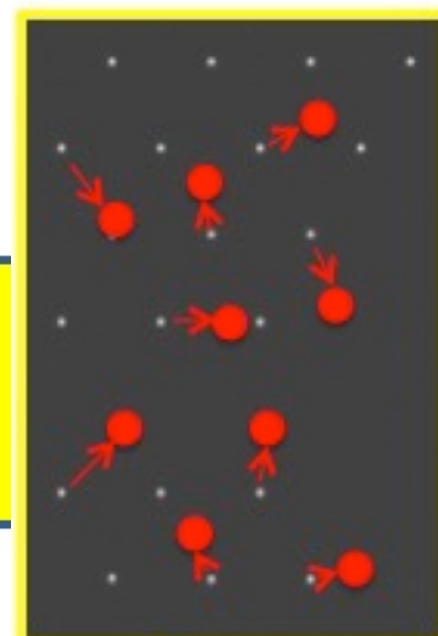
2400 fibers steered by positioners

On the TUE floor (IR side)

4 spectrographs

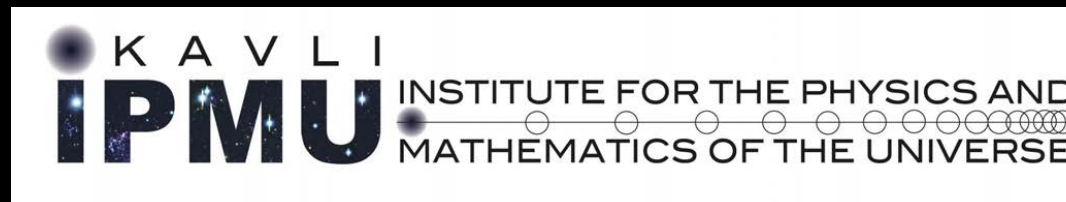
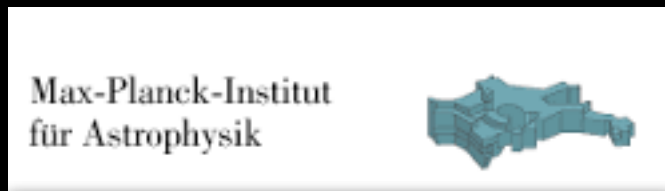
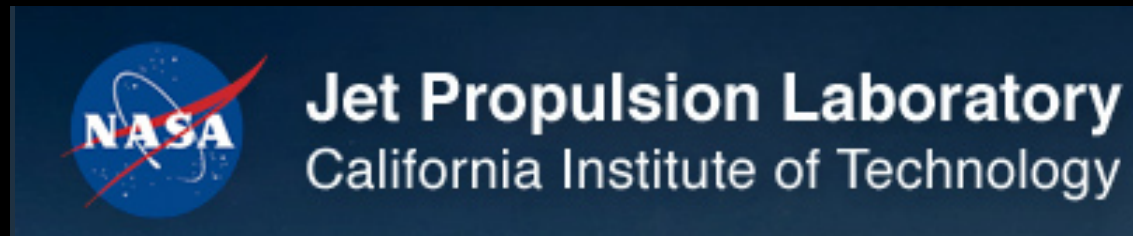
Fiber cable

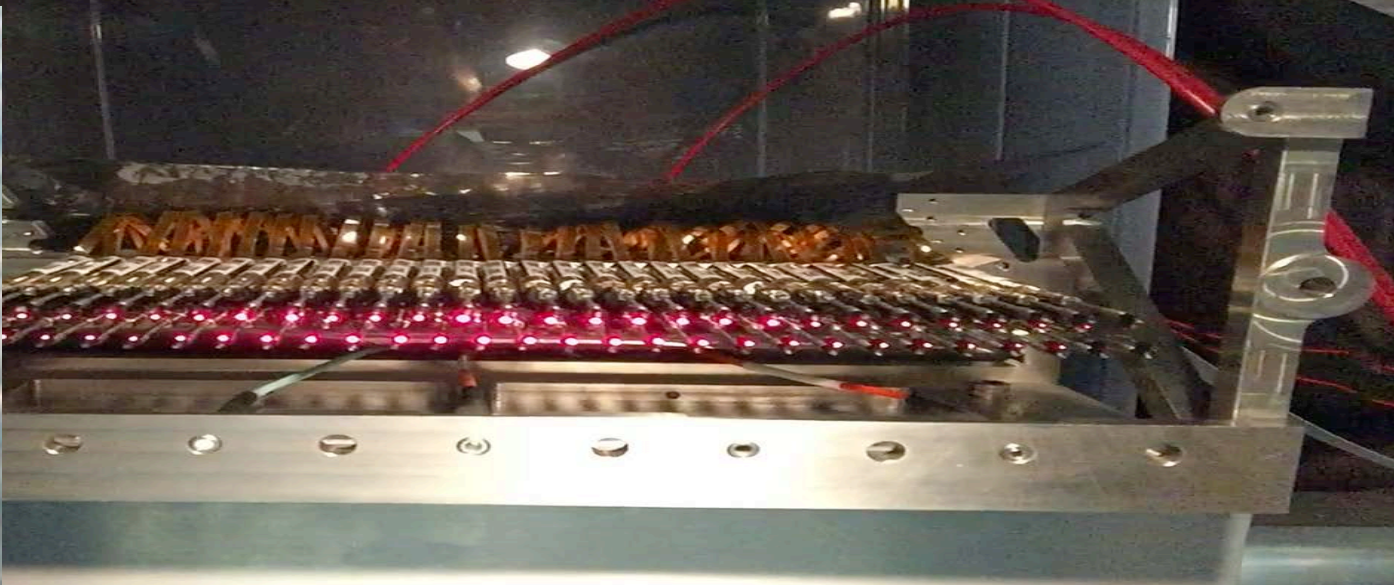
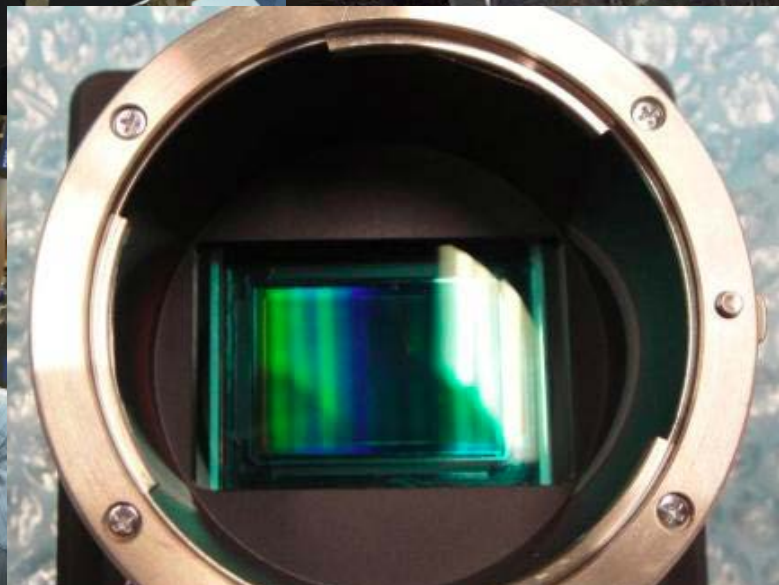
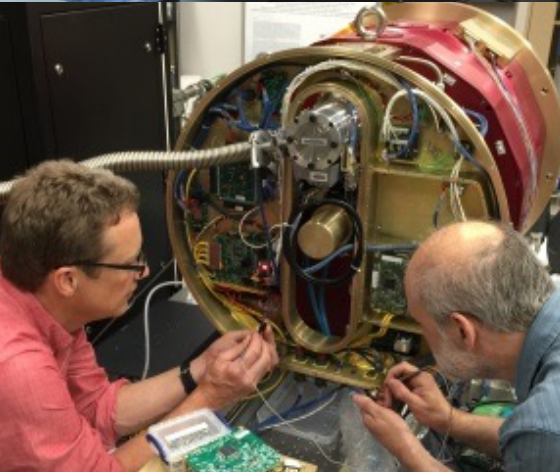
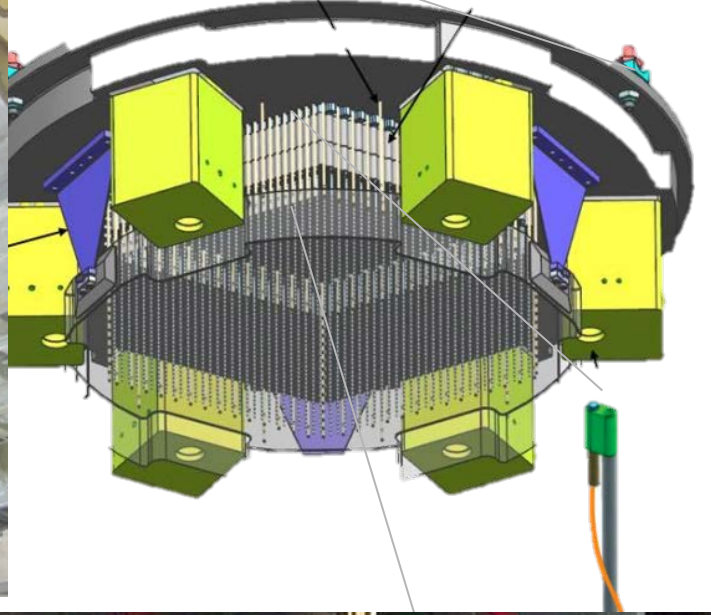
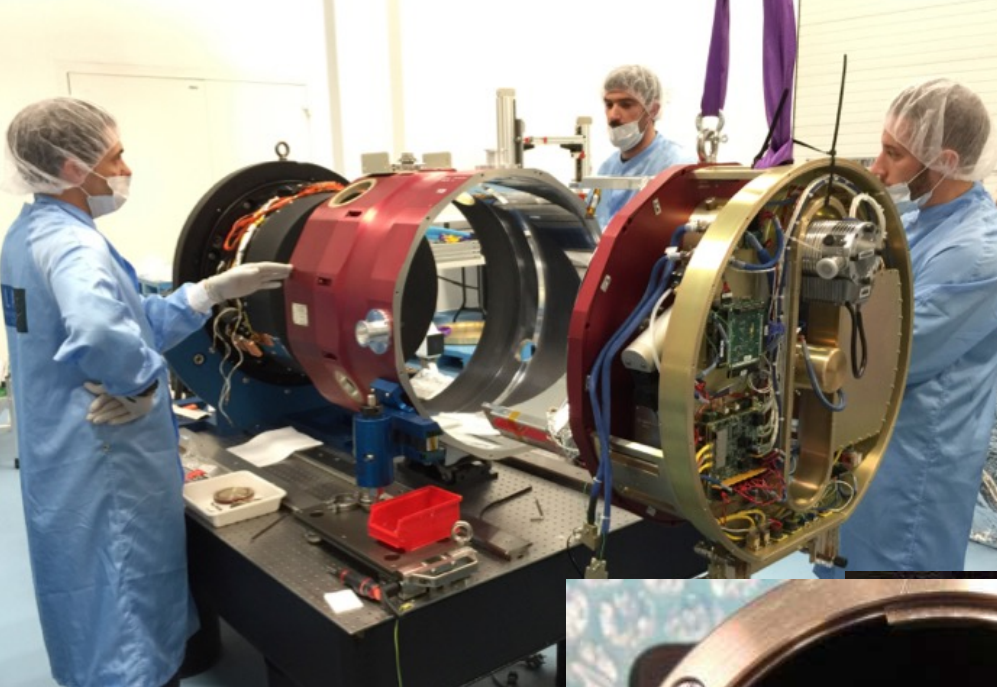
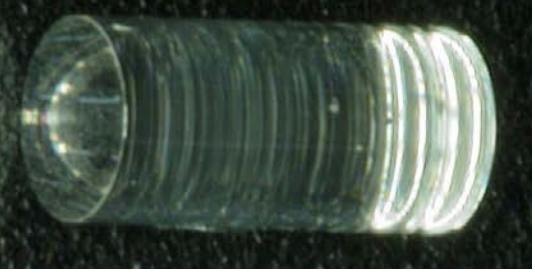
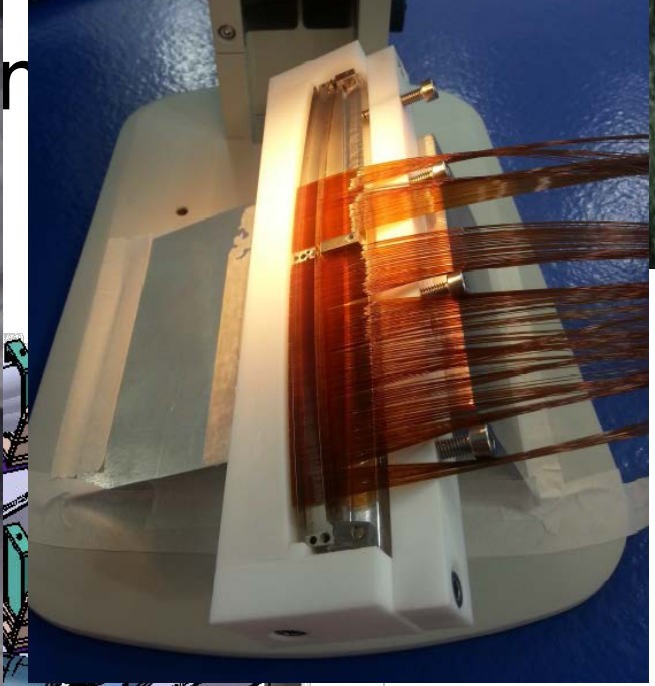
Metrology camera as a Cassegrain instrument



Subaru Telescope

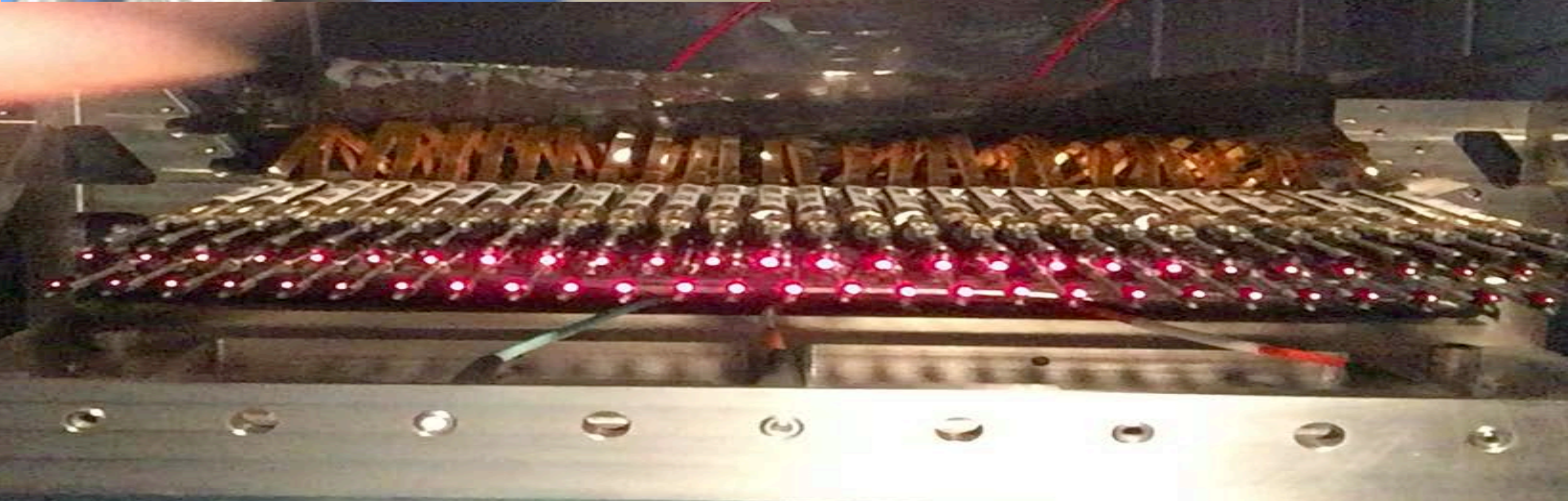
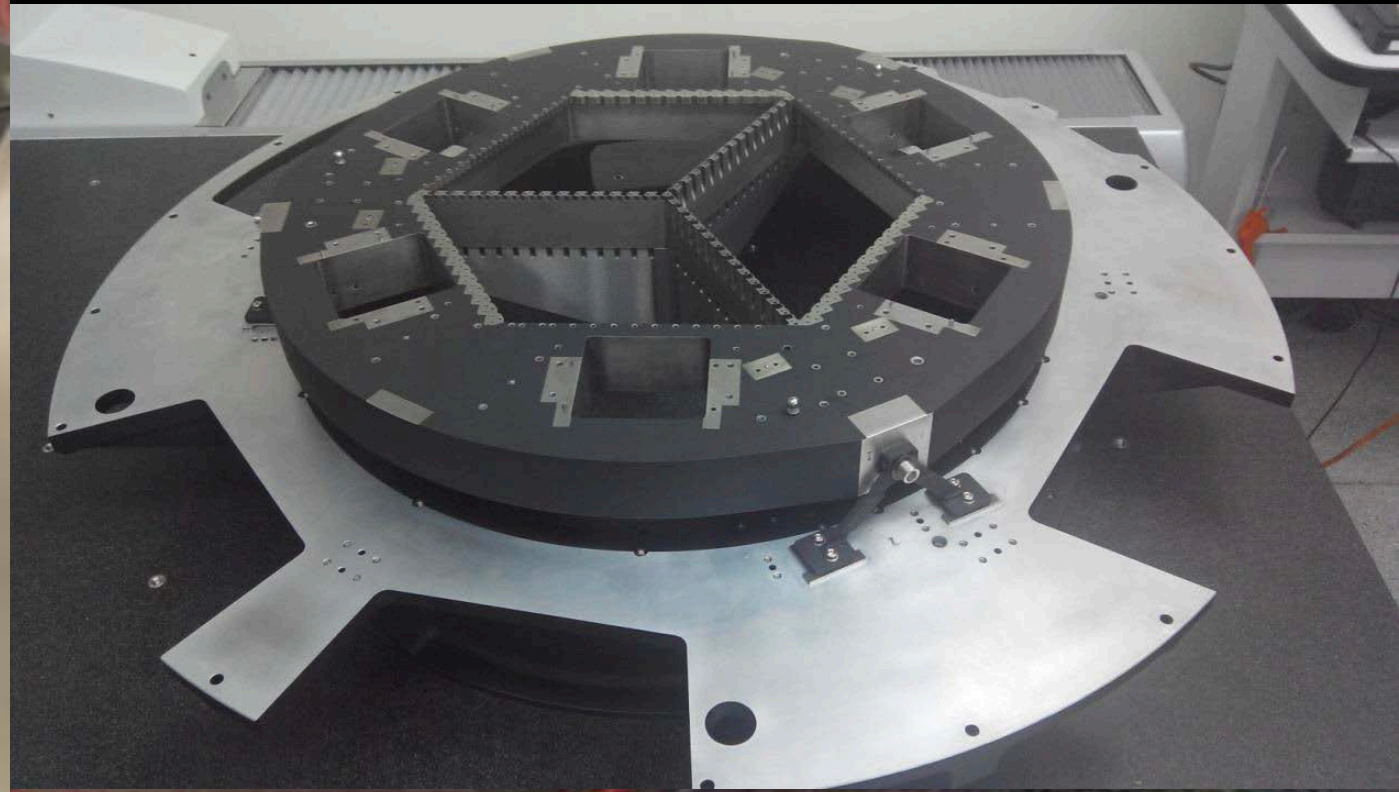
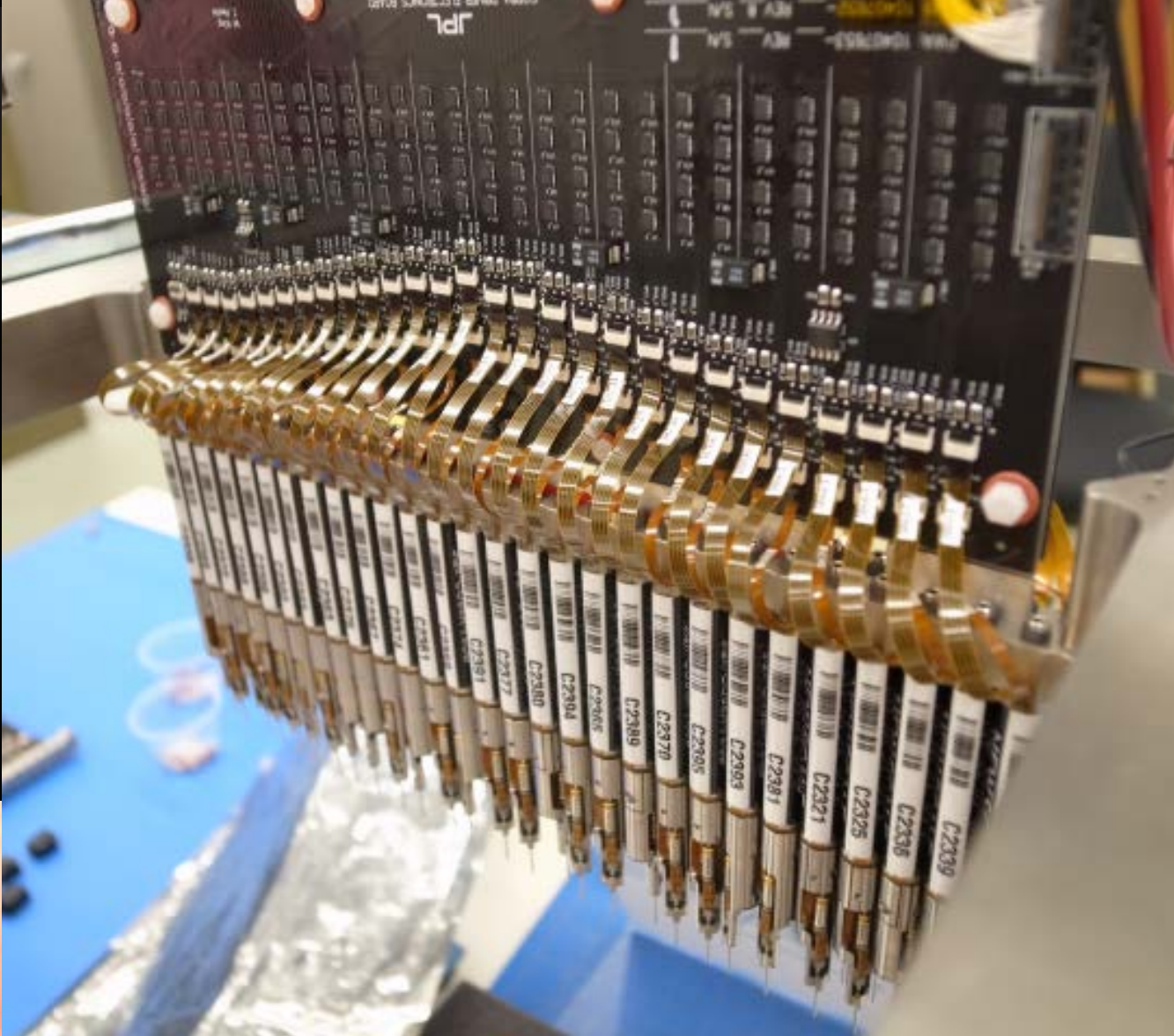
PFS collaboration







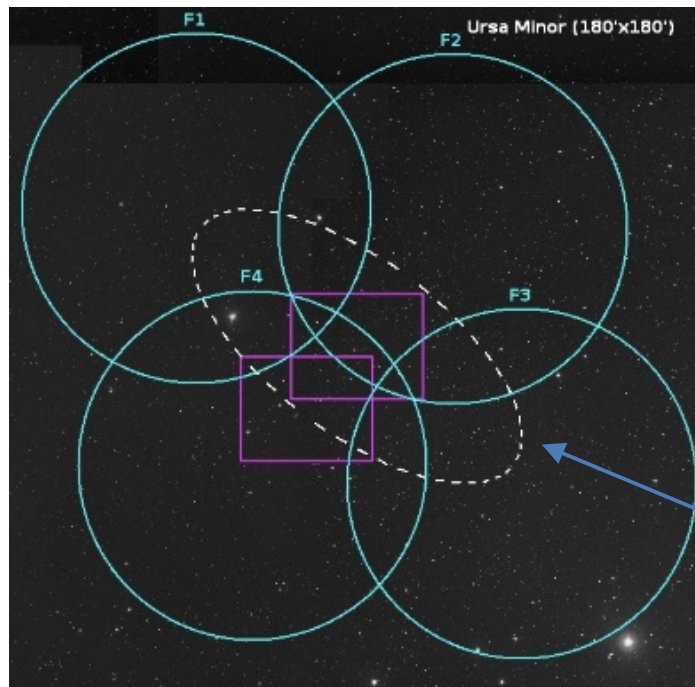
2396 Cobras



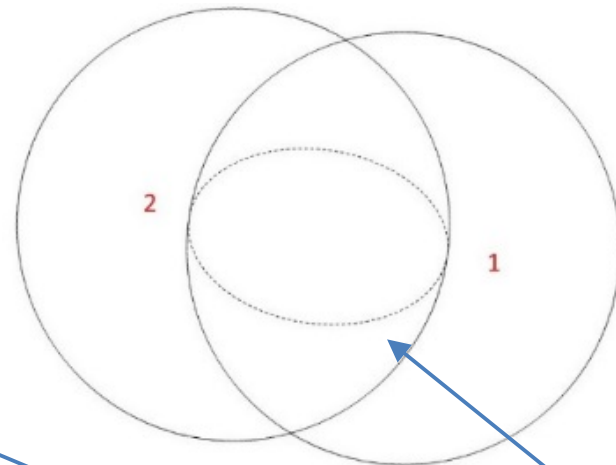
PFS pointings for MW satellites

~ HSC imaging data are available for all samples ~

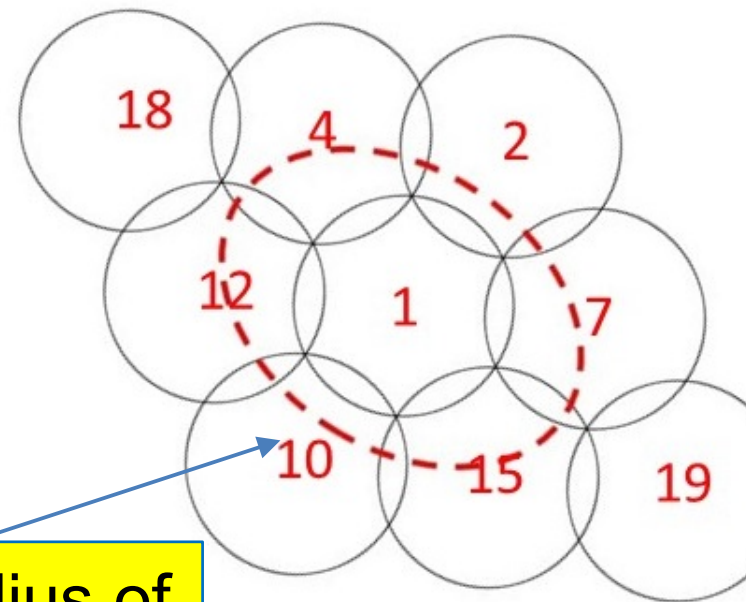
Ursa Minor



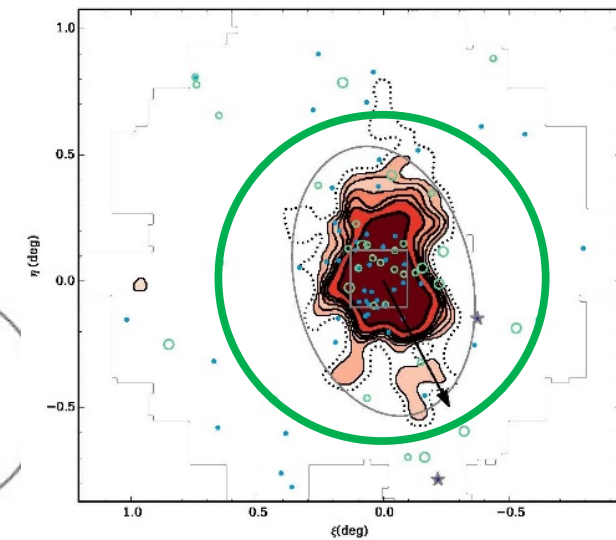
Draco



Sextans

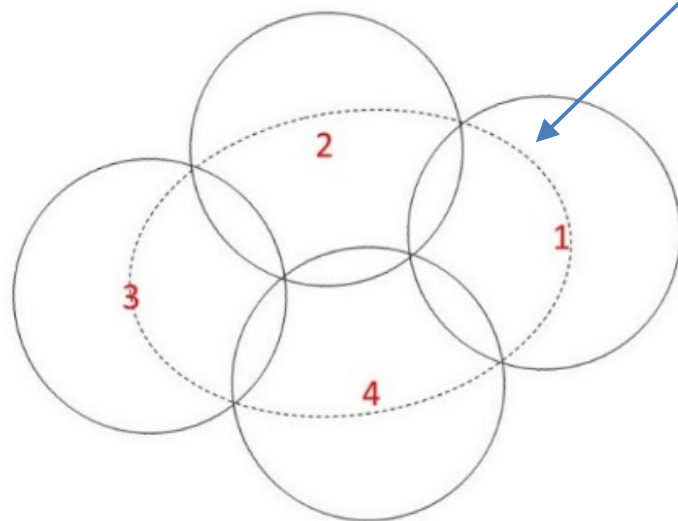


Bootes I

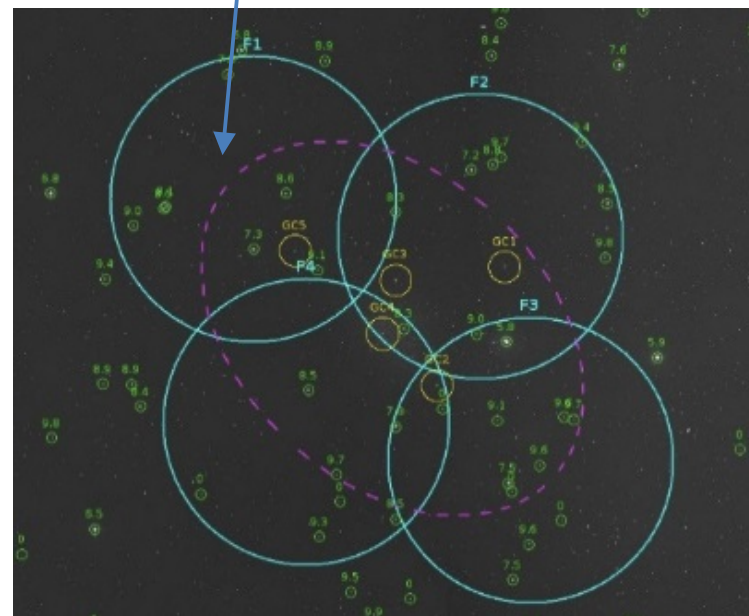


tidal radius of stellar comp.

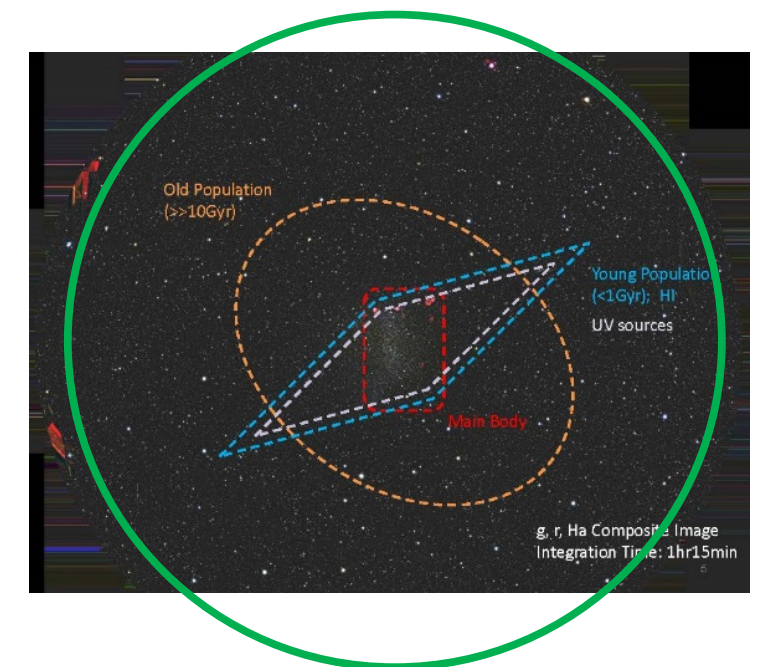
Sculptor



Fornax

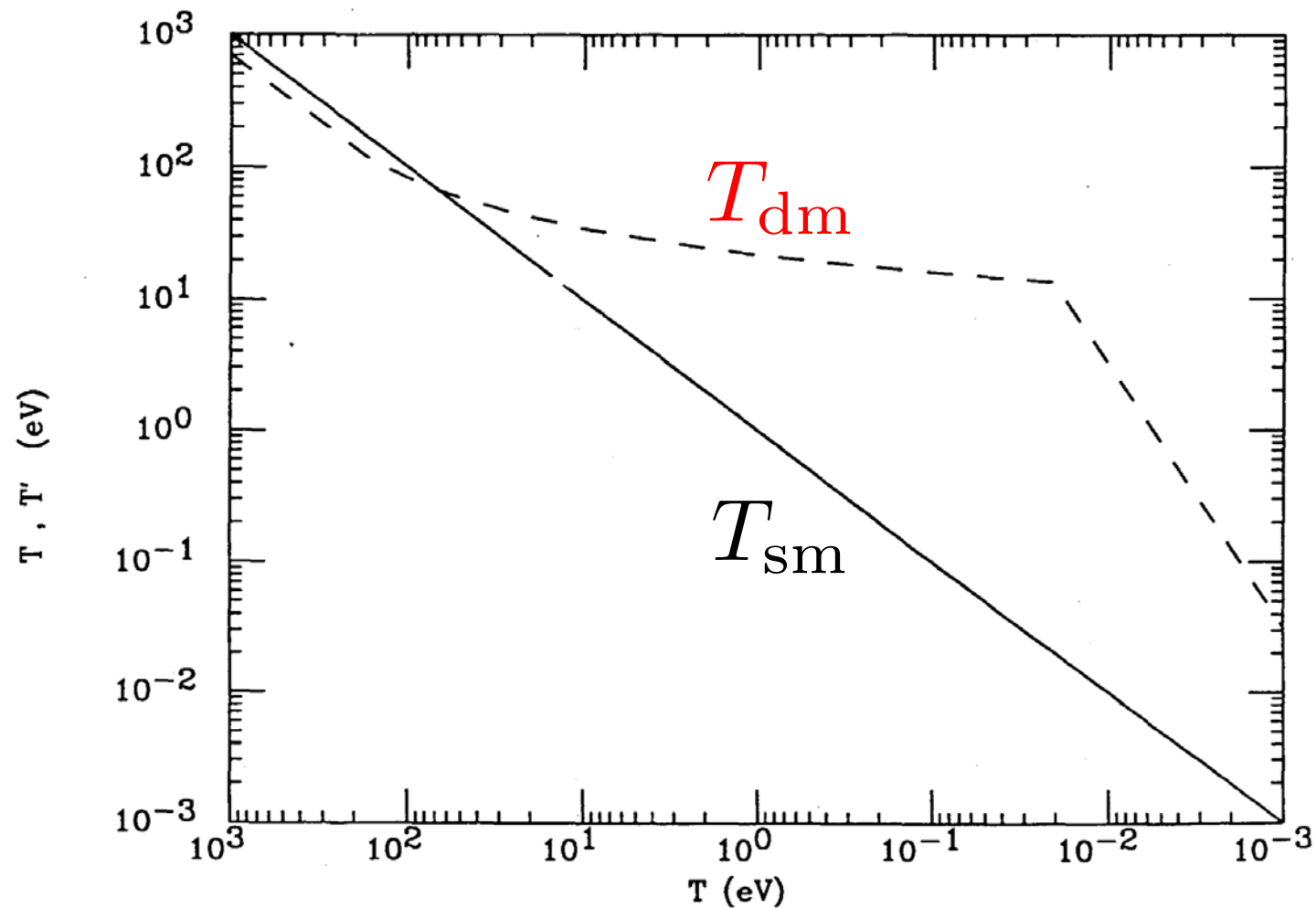


NGC6822



Interactions with SM

if totally decoupled



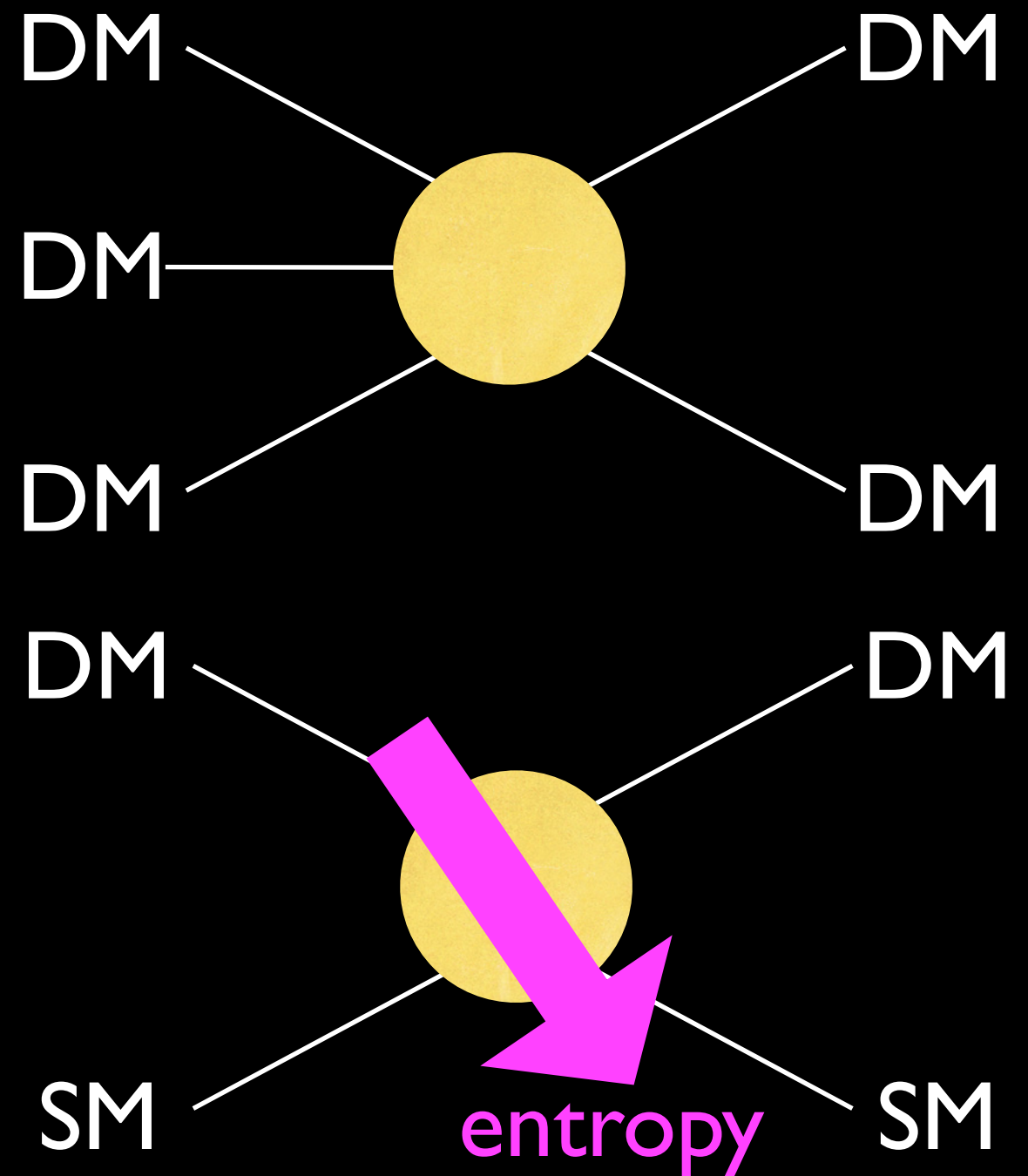
Carlson, Hall and Machacek,
Astrophys. J. 398, 43 (1992)

- $3 \rightarrow 2$ annihilations without heat exchange is excluded by structure formation, [de Laix, Scherrer and Schaefer, Astrophys. J. 452, 495 (1995)]



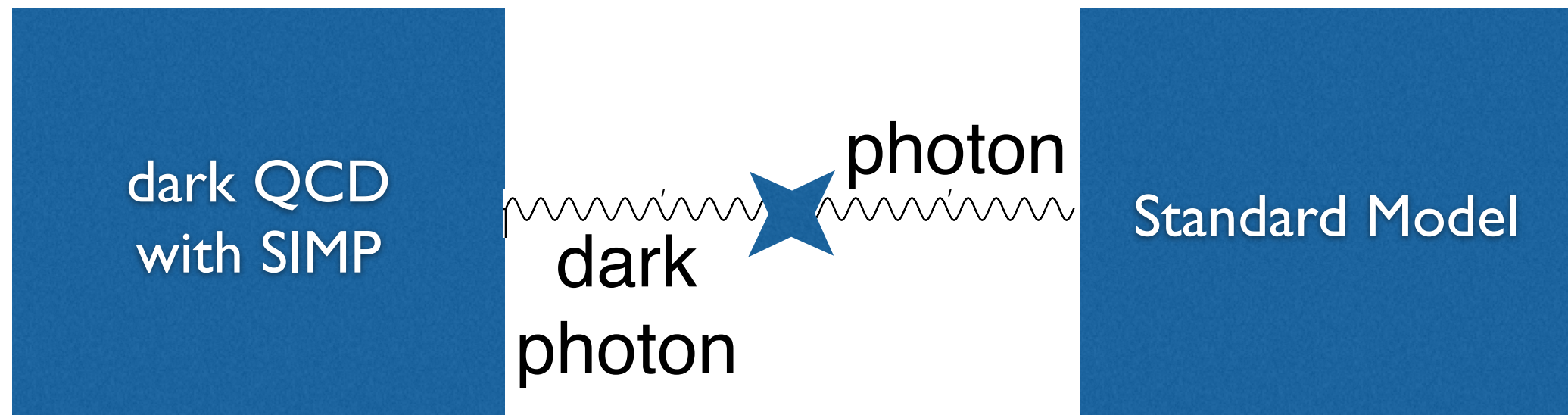
communication

- 3 to 2 annihilation
- excess entropy *must* be transferred to e^\pm, γ
- need communication at some level
- leads to experimental signal



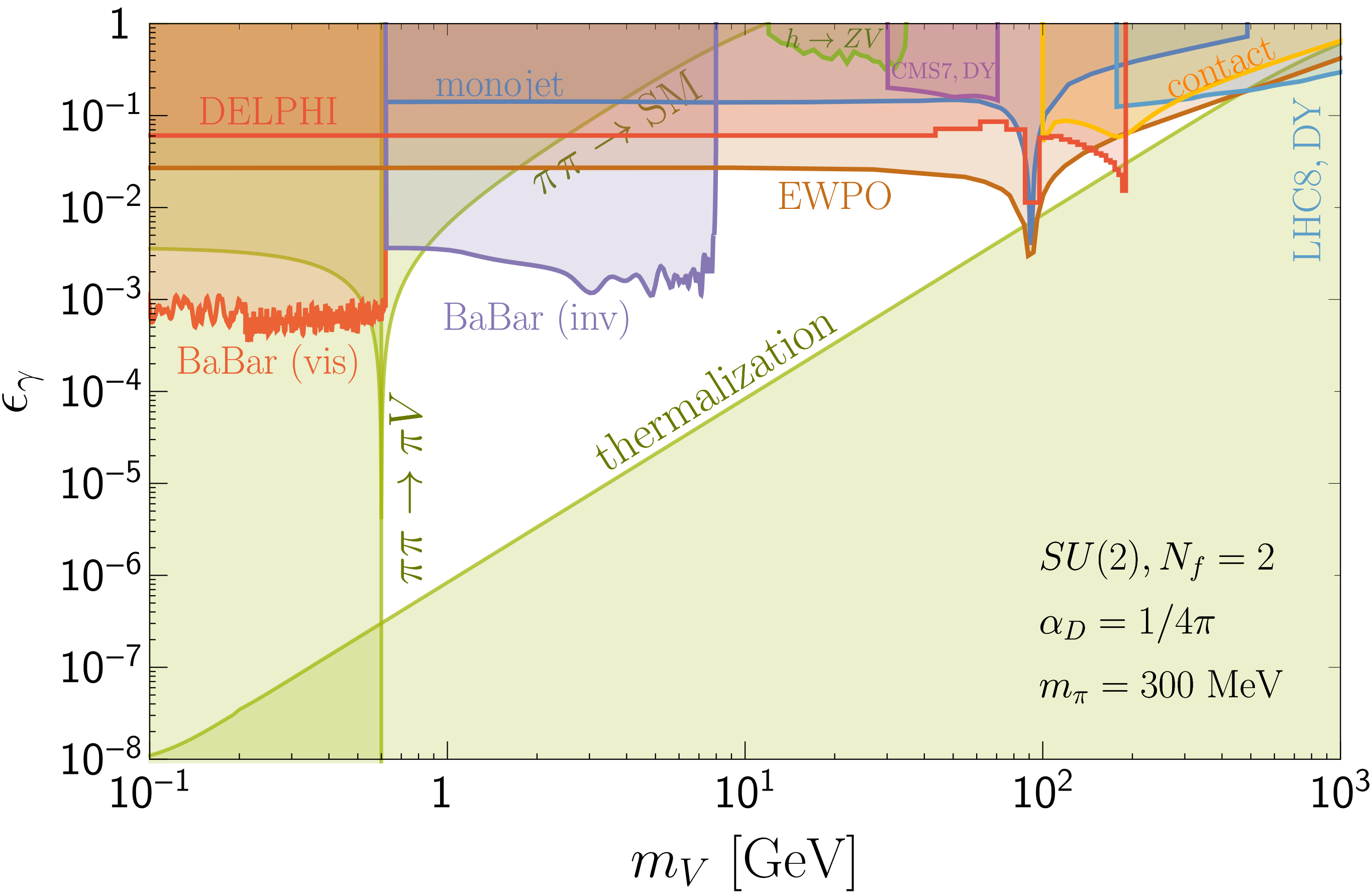
vector portal

arXiv:1512.07917

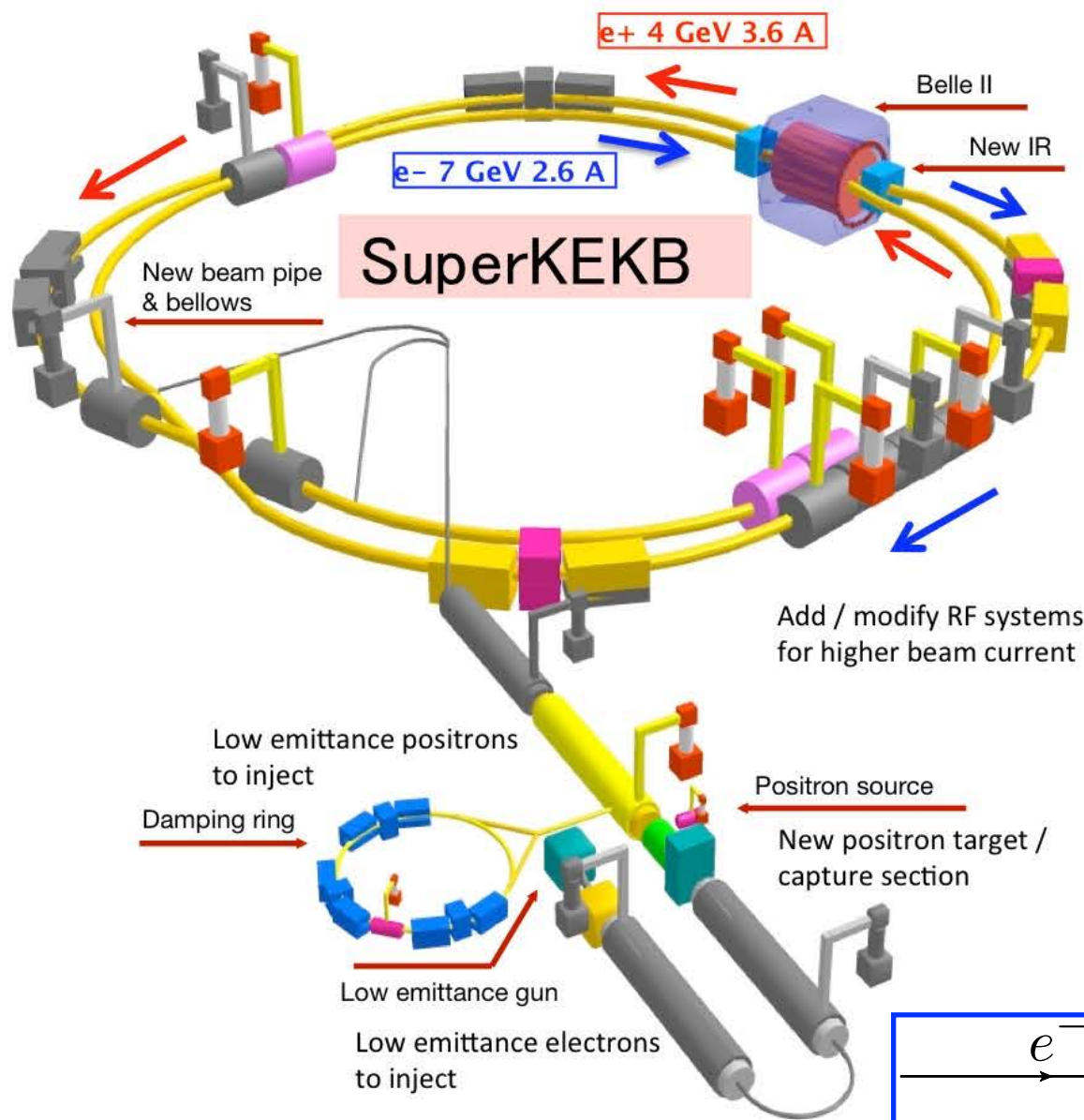


$$\frac{\epsilon_\gamma}{2} F_{\mu\nu} F_D^{\mu\nu}$$

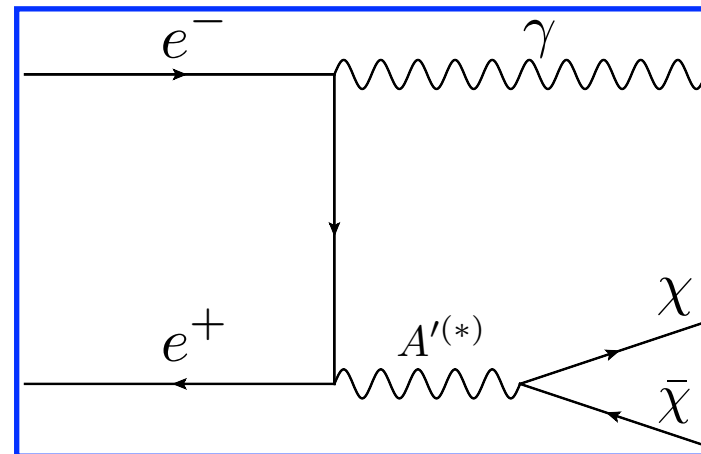
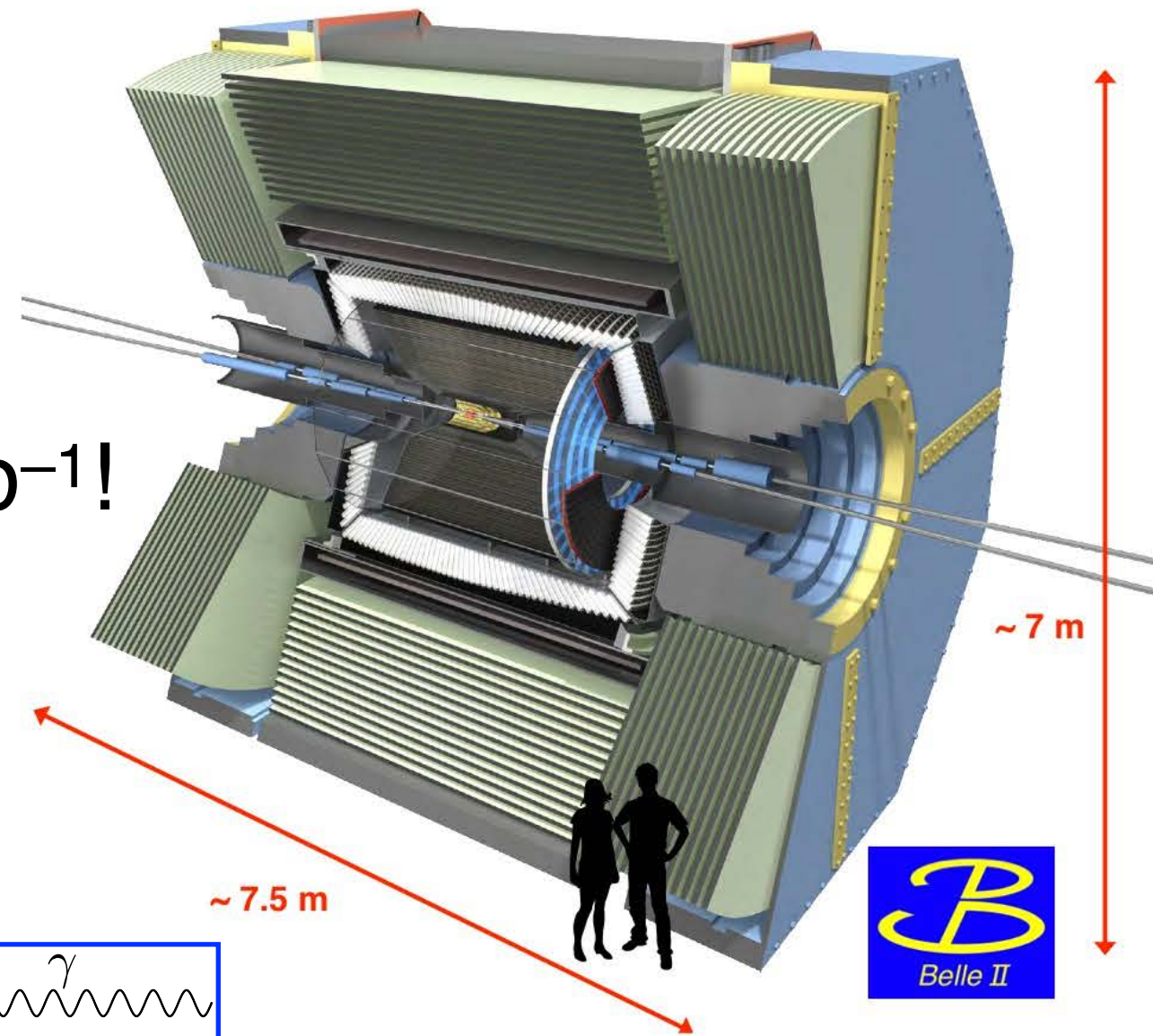
also axion portal: +Katelin Schutz, Robert McGehee, arXiv:1806.10139



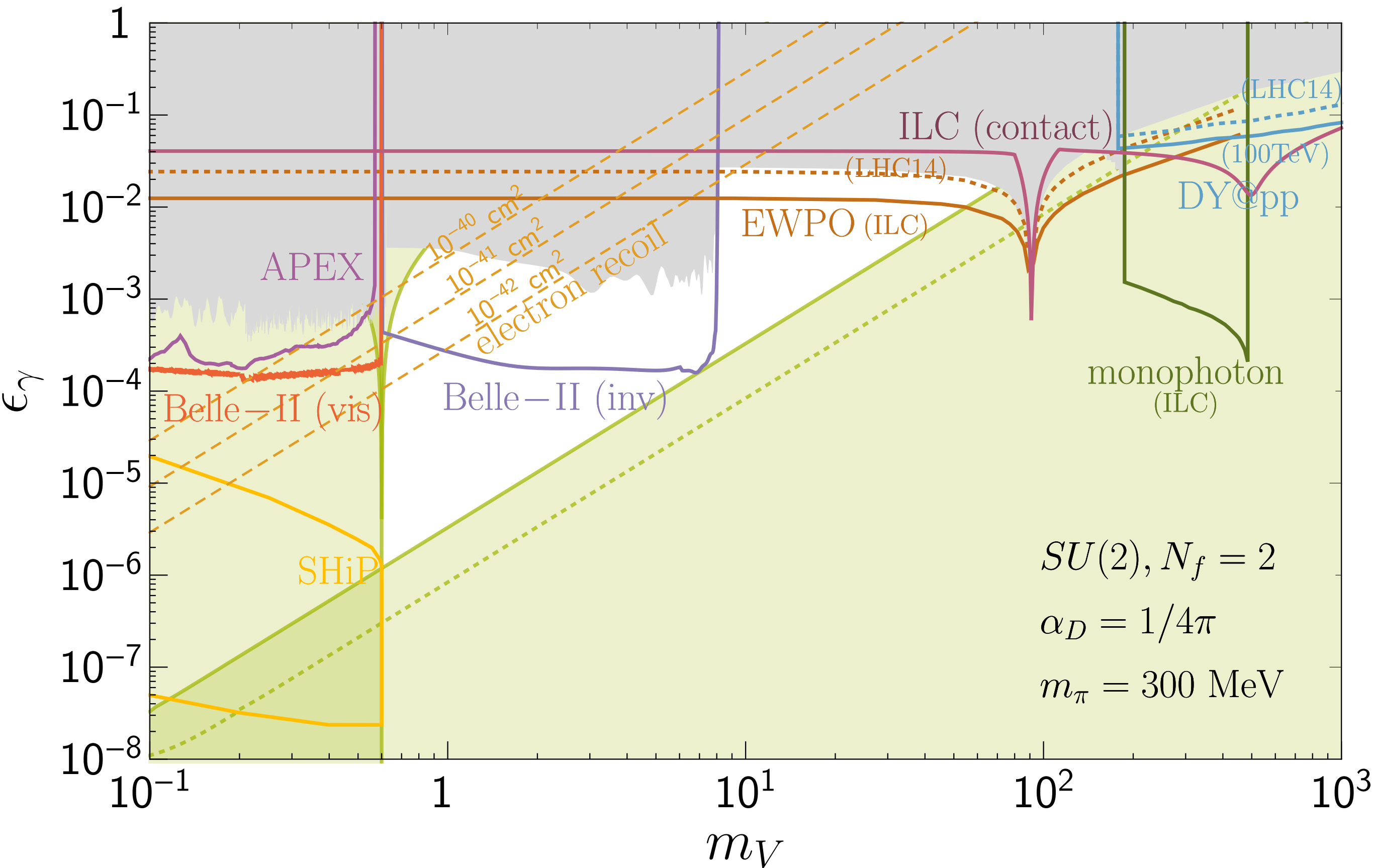
$e^+ e^-$ colliders

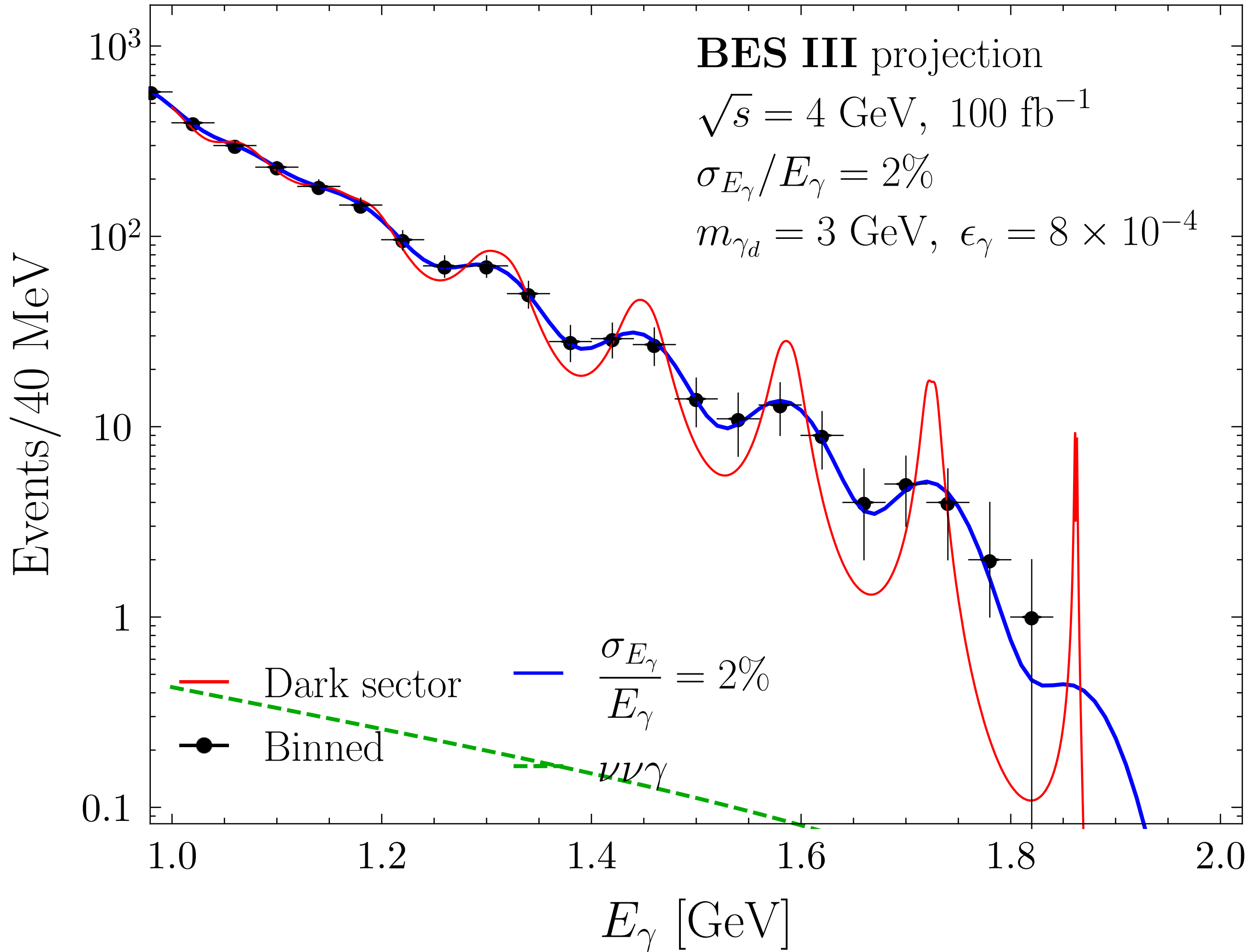


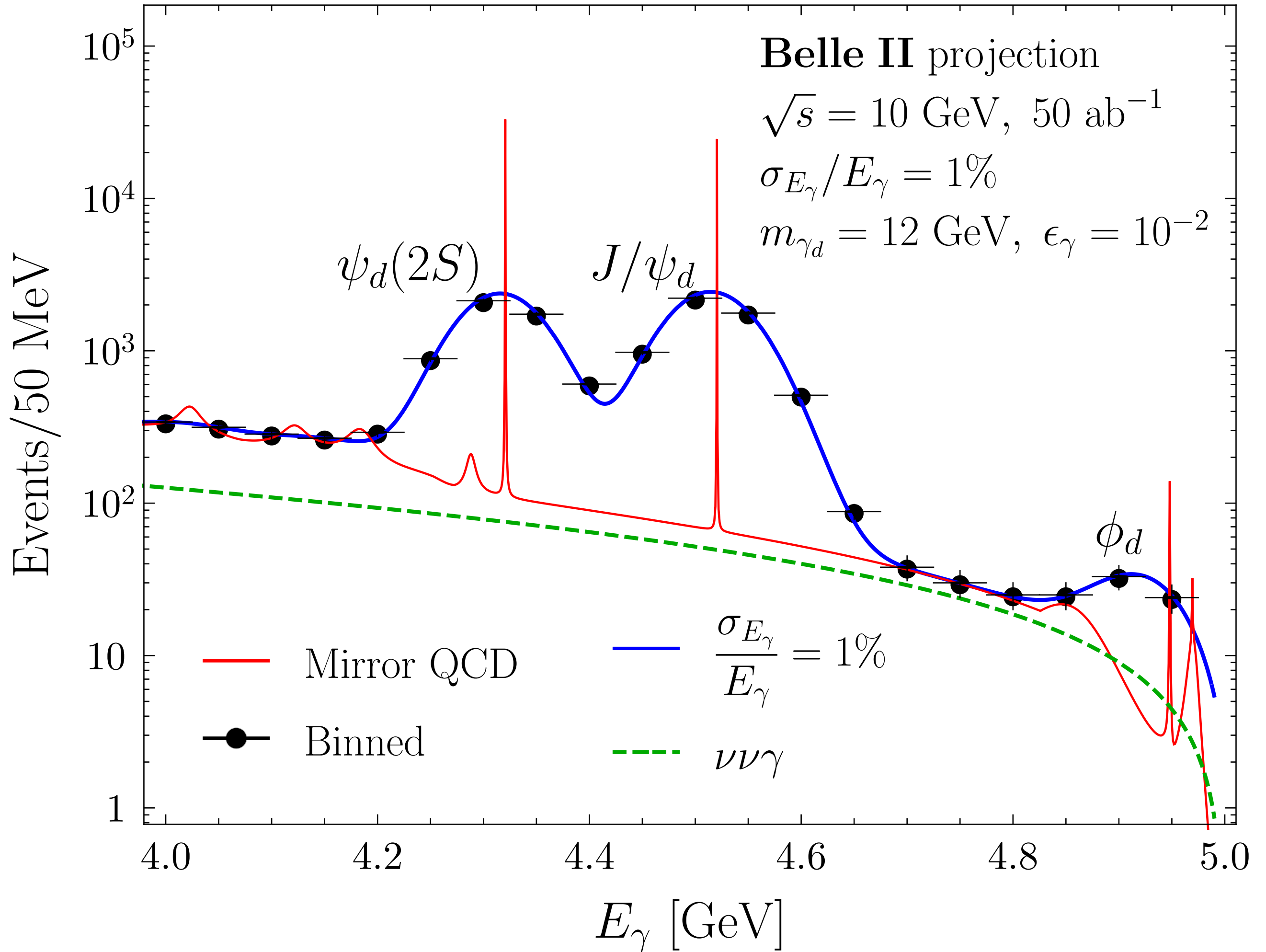
50 ab^{-1} !



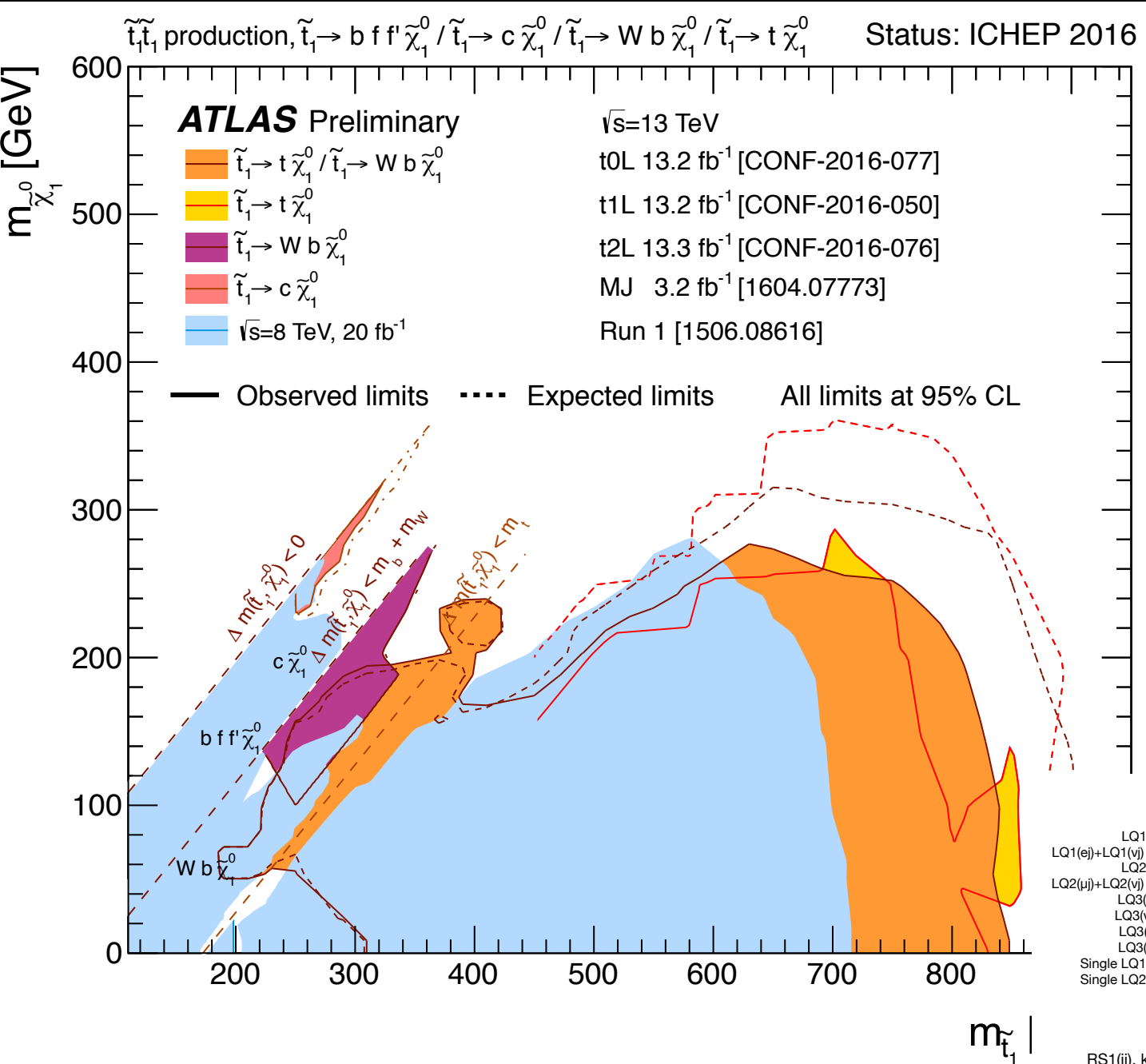
$$E_\gamma = \frac{\sqrt{s}}{2} \left(1 - \frac{M_{\text{inv}}^2}{s} \right)$$



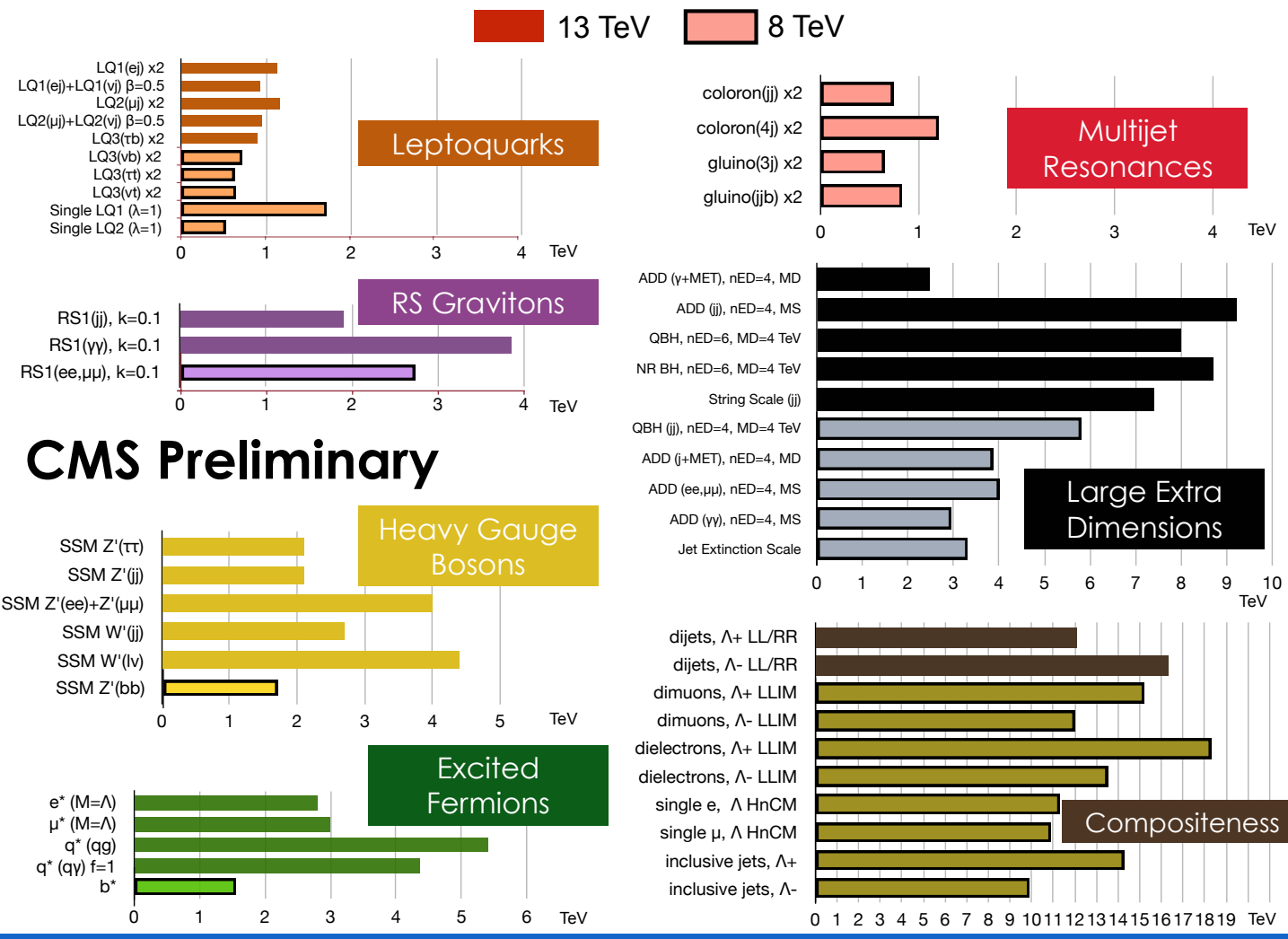




Back to Hierarchy Problem



LHC excludes mostly colored particles



no sign of new physics that explains hierarchy problem!

Twin Higgs

- Take two mirror copies of the SM:

$$(SM_A) \times (SM_B)$$



Z_2

An exchange symmetry. $A \leftrightarrow B$.

- Assume Higgs potential has an $SU(4)$ or $SO(8)$ global symmetry in the UV.

$$V = -\mu^2(|H_{SM}|^2 + |H_{twin}|^2) + \lambda(|H_{SM}|^2 + |H_{twin}|^2)^2 + \kappa|H_{SM}|^2|H_{twin}|^2$$

- Take a small hierarchy of Higgs vevs:

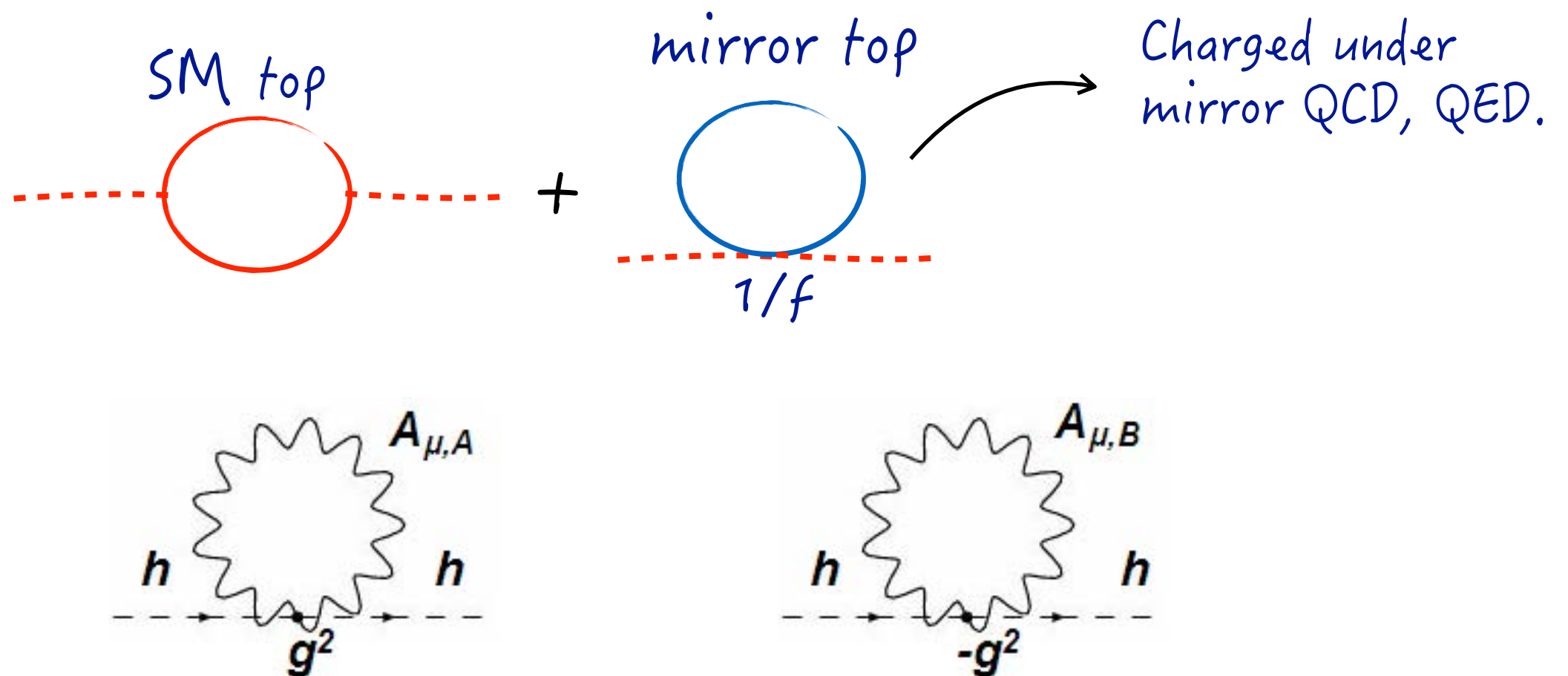
$$\langle H_A \rangle = v \quad \langle H_B \rangle = f \quad \text{with } v < f.$$

$SO(8)/SO(7)$: 7 NGBs

3: eaten by $SU(2)_{twin}$

Twin Higgs

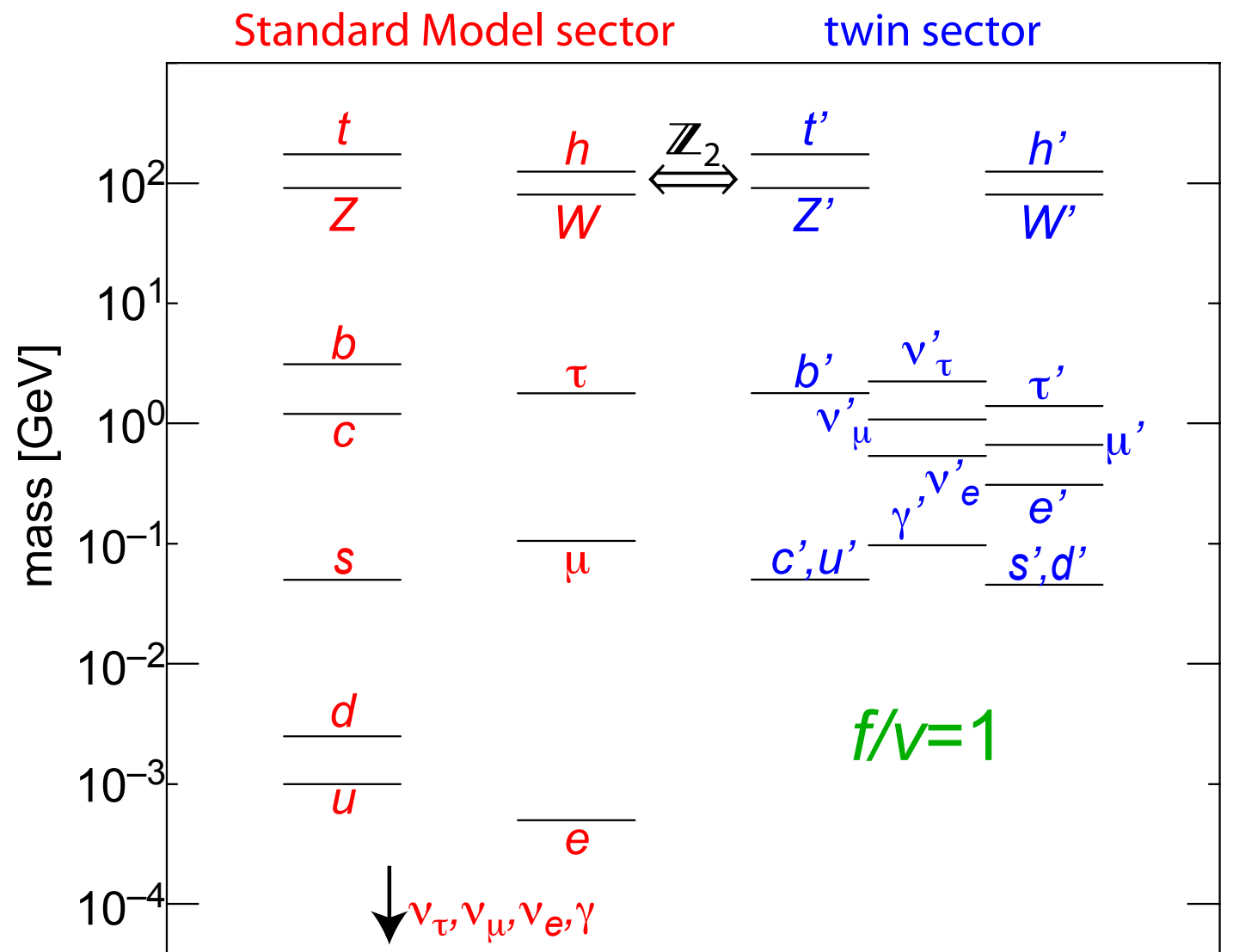
- All NP within LHC reach is SM neutral.
- Pseudo Nambu-Goldstone Higgs, cancellation ...



Twin spectrum

arXiv:1805.09345

- light fermion spectrum not important to hierarchy because of small Yukawas
- Planck: $\Delta N_{eff} < 0.29$
 - assume leptons heavy
 - no N_{eff} problem
- assume exact SU(2)
 - (u,d) vs (c,s)
- lightest pions are triplet in $SU(2)_{d,s}$
 - stable



mesons

- assume exact SU(2)
 - (u,d) vs (c,s)
- approximate SU(4)
- $m_{u,c} = m_{d,s}(1 + \Delta)$
- 15 mesons
- lightest pions are triplet in SU(2)_{d,s}
 - stable
- triplet in SU(2)_{u,c} also practically stable
- 8 charged ones stable
- only one unstable η

meson M	particle content	$m_M^2 \propto$	m_M
$\theta^0(\mathbf{3}, \mathbf{1})$	$u'\bar{c}', c'\bar{u}', \frac{1}{\sqrt{2}}(u'\bar{u}' - c'\bar{c}')$	$2m_{u'}$	$m_\pi(1 + \Delta)$
$D^+(\mathbf{2}, \mathbf{2})$	$u'\bar{d}', c'\bar{d}', u'\bar{s}', c'\bar{s}'$	$m_{u'} + m_{d'}$	$m_\pi(1 + \frac{\Delta}{2})$
$D^-(\mathbf{2}, \mathbf{2})$	$d'\bar{u}', s'\bar{u}', d'\bar{c}', s'\bar{c}'$	$m_{u'} + m_{d'}$	$m_\pi(1 + \frac{\Delta}{2})$
$\eta^0(\mathbf{1}, \mathbf{1})$	$\frac{1}{2}(d'\bar{d}' + s'\bar{s}' - u'\bar{u}' - c'\bar{c}')$	$m_{u'} + m_{d'}$	$m_\pi(1 + \frac{\Delta}{2})$
$\pi^0(\mathbf{1}, \mathbf{3})$	$d'\bar{s}', s'\bar{d}', \frac{1}{\sqrt{2}}(d'\bar{d}' - s'\bar{s}')$	$2m_{d'}$	m_π

TABLE I. Decomposition of the meson $SU(4)_f$ 15-plet under $SU(2)_U \times SU(2)_D \times U(1)_{EM}$. The third column shows the linear combination of quark masses that determines the meson masses-squared. From top to bottom, the meson masses go from heaviest to lightest, assuming $m_{d'} = m_{s'} < m_{u'} = m_{c'} = m_{d',s'}(1 + \Delta)$.

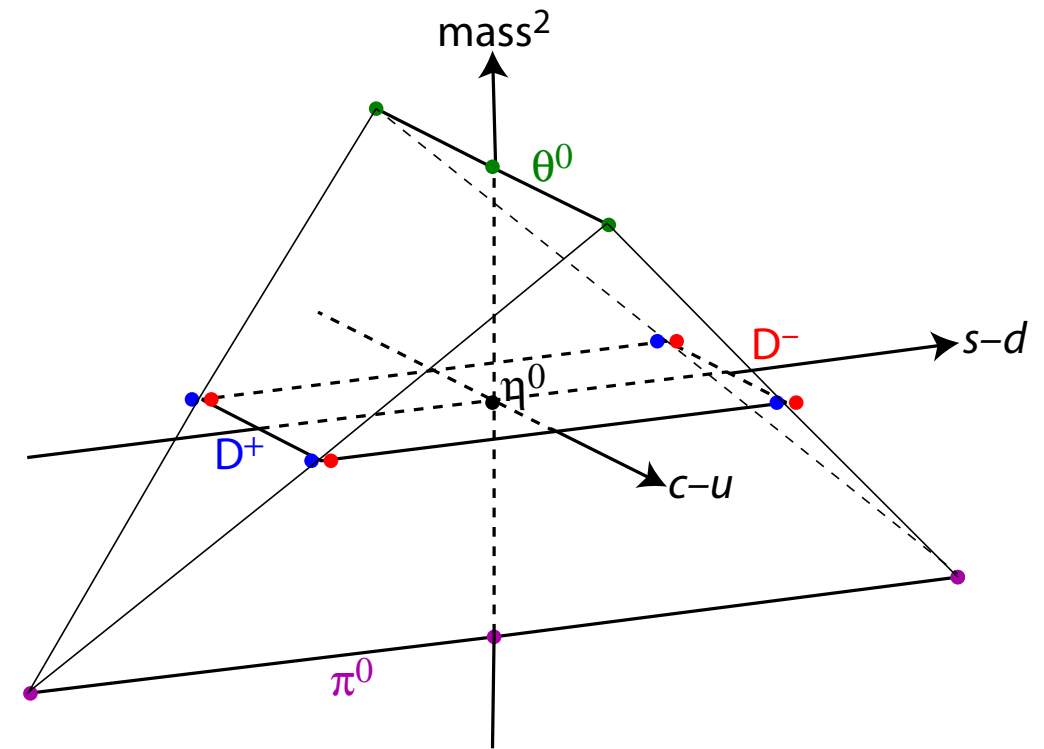
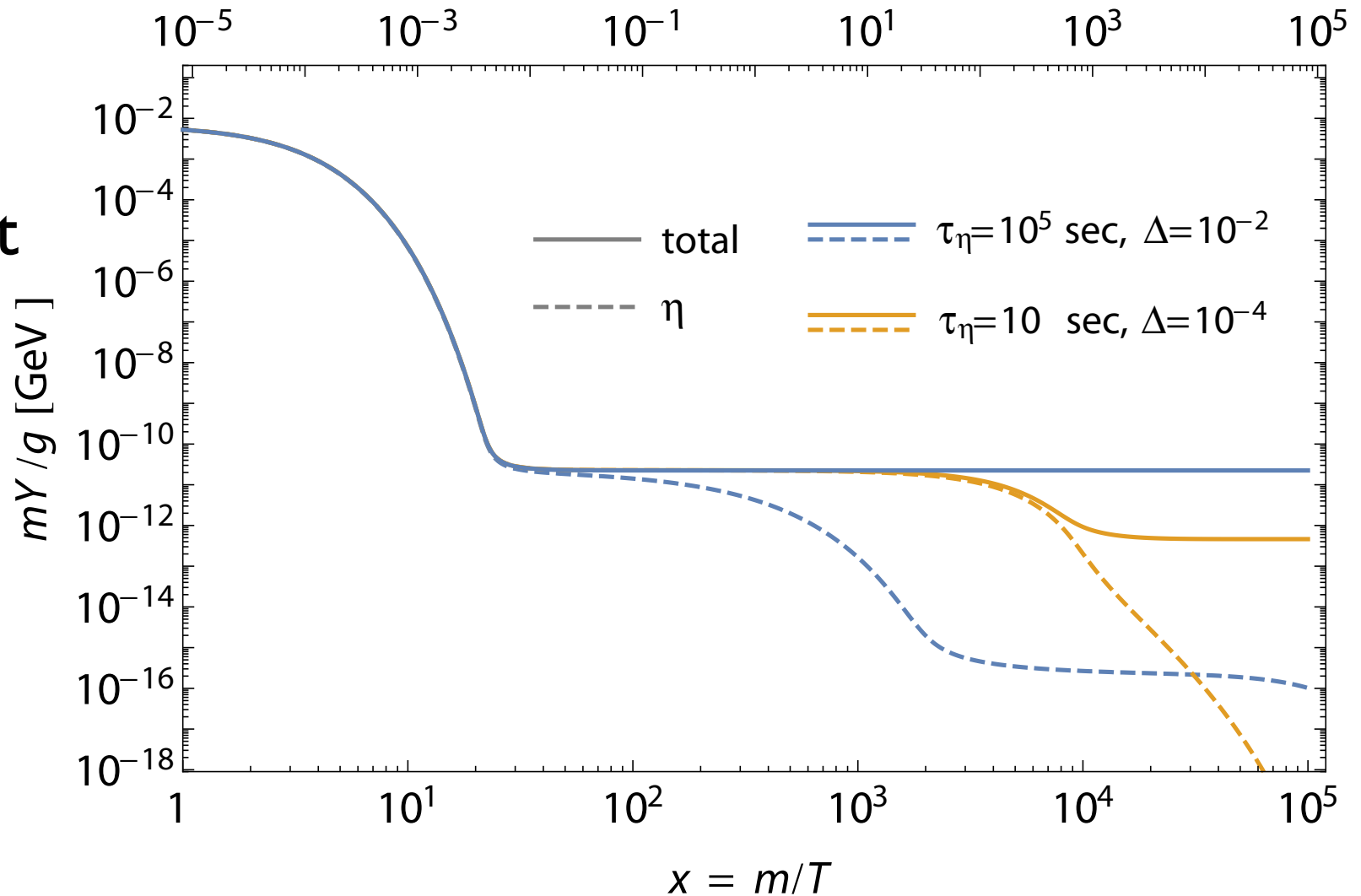
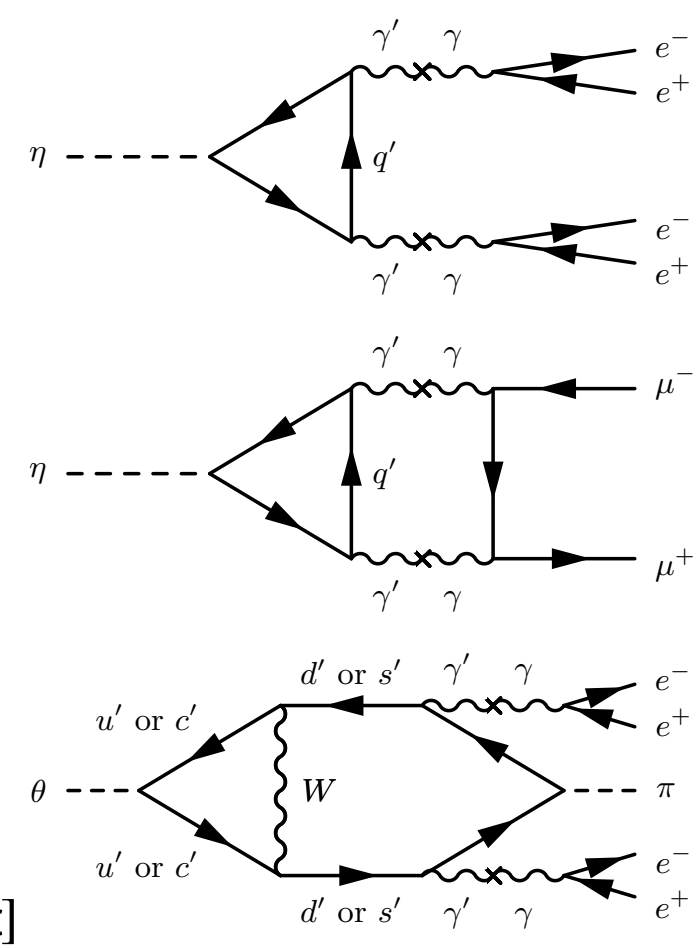


FIG. 2. A visual representation of the meson spectrum.

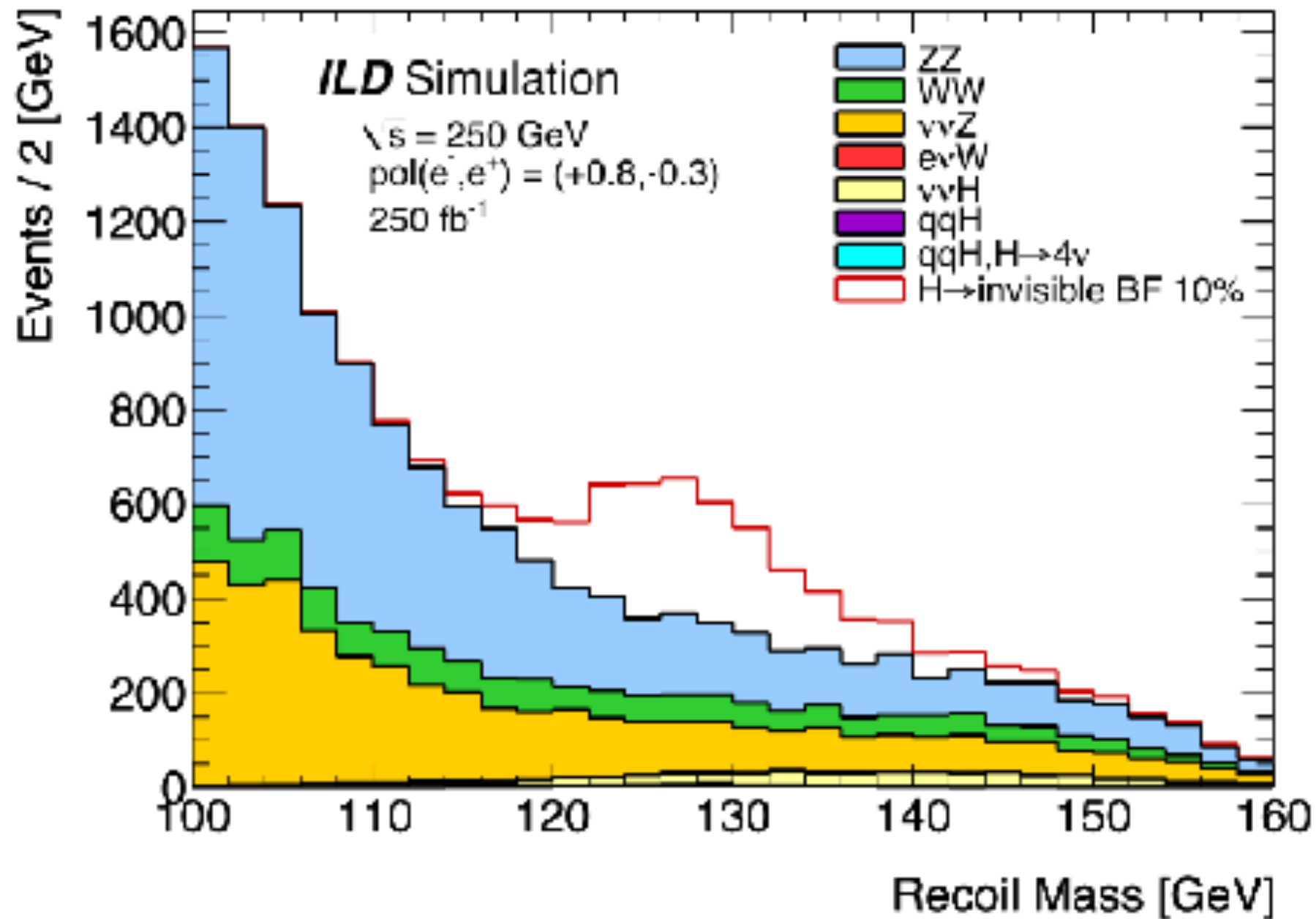
$\eta\eta \rightarrow \pi\pi$

- $\eta \rightarrow \mu^+\mu^-$ or $2(e^+e^-)$
- $\tau_\eta \sim 10^5$ sec
- 1. $\pi\pi \rightarrow \eta\eta$ kinematically not possible
- 2. lose η by $\eta\eta \rightarrow \pi\pi$
- 3. η decays before BBN limit
- $\pi\pi \rightarrow \eta\eta \rightarrow \text{SM}$ shouldn't happen in halo today
- can easily be avoided $\Delta > v_{\text{halo}}^2 \sim 10^{-6}$



invisible Higgs @ ILC

Keisuke Fujii



$\text{BR}(H \rightarrow \text{invisible}) \sim (v/f)^2$ down to 0.3%

Quintessence

Swampland

EFT

String Landscape

$$|\nabla V| > cV$$

(meta)-stable
positive vacuum energy

Swampland

$$w = -1 + \frac{2c^2}{3 + c^2}$$

“Quintessence”

- Accelerated Expansion happened before
 - Inflation
- Is current acceleration also by scalar field?
 - very difficult to keep flat potential for Q
 - SUSY broken $> \text{TeV}$
 - at the least, $m_{3/2} > (\text{TeV}^2/M_{Pl}) \sim \text{eV}$
 - *i.e.*, vector mediation (Hook, HM)
 - more typically $m_{3/2} > \text{TeV}$
 - but we need $m_Q \sim H_0 \sim 10^{-33} \text{ eV}$

shift symmetry

- incorporate into supergravity
- shift symmetry (monodromy) in Kähler
 - $Q \rightarrow Q + i\alpha$
 - $K(Q, Q^*) = K(Q + Q^*) \sim (Q + Q^*)^2 / 2$

$$V = e^K \left((K_i W + W_i)^* K_{\bar{i}}^{-1j} (K_j W + W_j) - 3|W|^2 \right)$$

$$= |W_Q|^2 - 3m_{3/2}(W(Q) + W^*(Q))$$

- need $m_{3/2}W(Q) \sim m_{3/2}\Lambda^3 \sim H_0^2$
- *any* potential can be lifted to supergravity
- also radiatively stable $\delta K \sim m_{3/2}^2 \Lambda^6$
- no fifth force

Chien-I Chiang, HM, arXiv:1808.02279

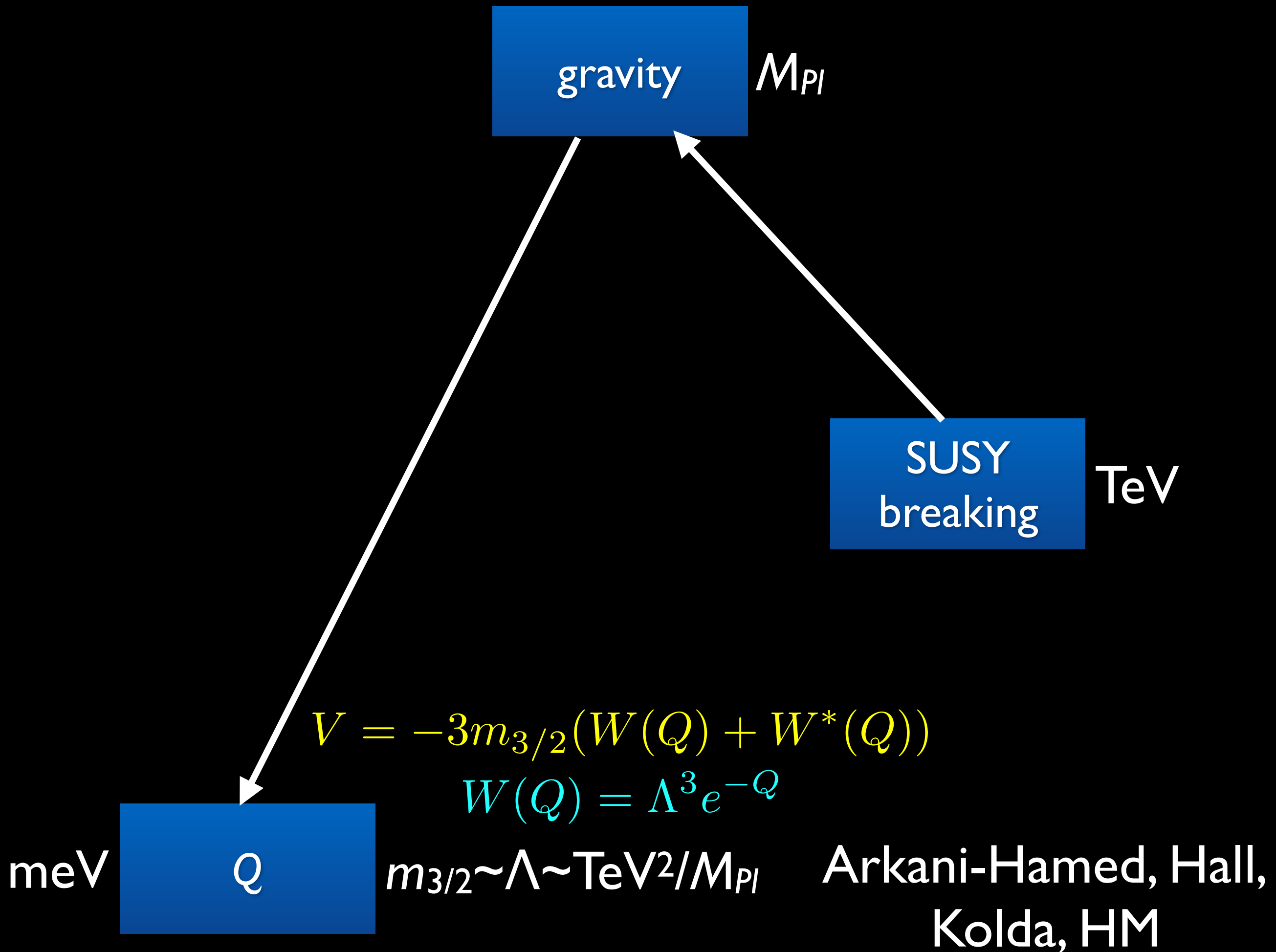
axion-like example

- incorporate into supergravity
- shift symmetry (monodromy) in Kähler, W
 - $Q \rightarrow Q + i\alpha$
 - $K(Q, Q^*) = K(Q + Q^*) \sim (Q + Q^*)^2 / 2$

$$V = e^K \left((K_i W + W_i)^* K_{\bar{i}}^{-1j} (K_j W + W_j) - 3|W|^2 \right)$$

$$= |W_Q|^2 - 3m_{3/2}(W(Q) + W^*(Q))$$

- shift symmetry = $U(1)_R$ symmetry
- $W(Q) = \Lambda^3 e^{-Q}$, $V = 3m_{3/2} \Lambda^3 \cos(\text{Im } Q)$
- symmetry spontaneously broken by $m_{3/2}$



shift symmetry

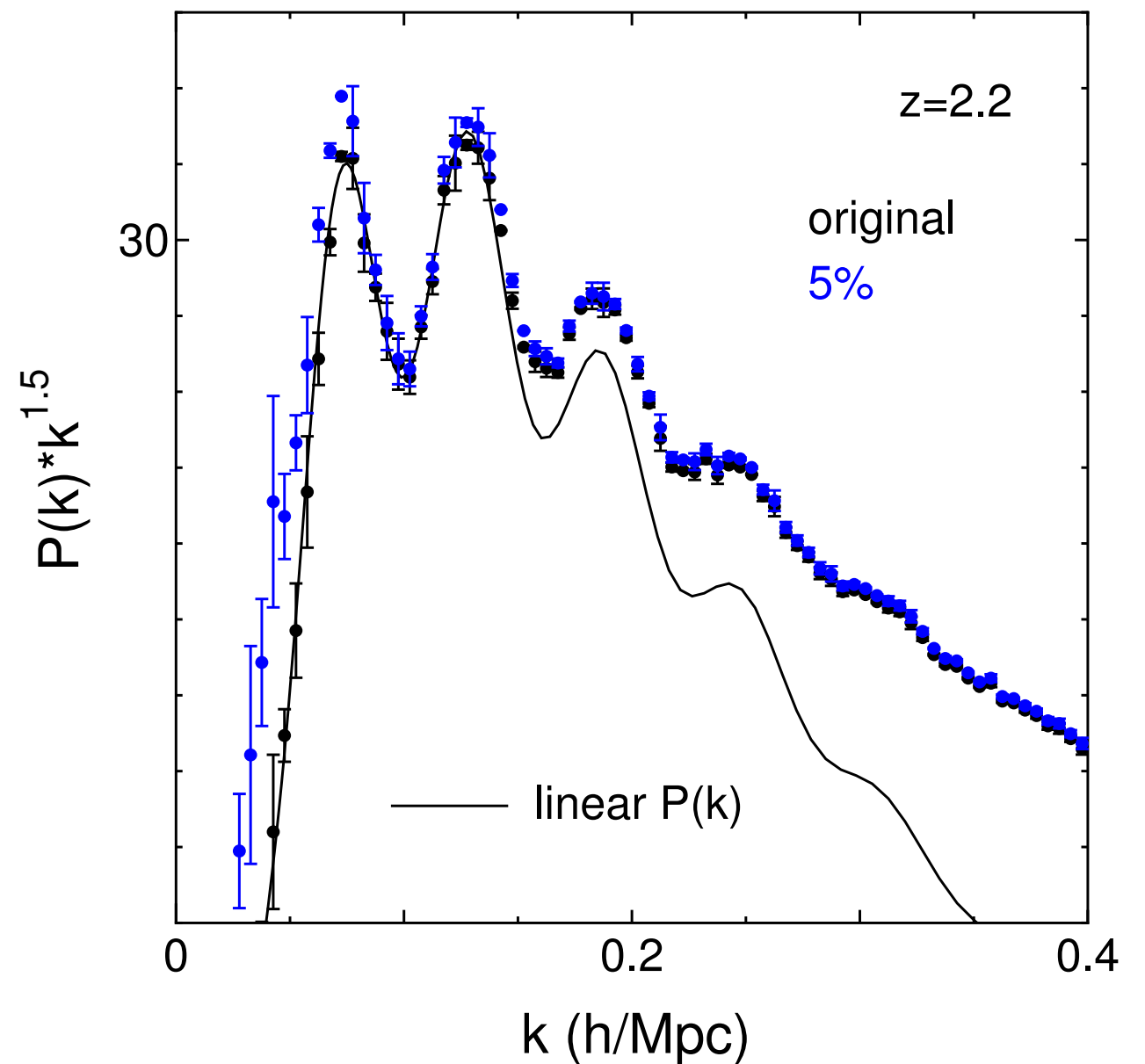
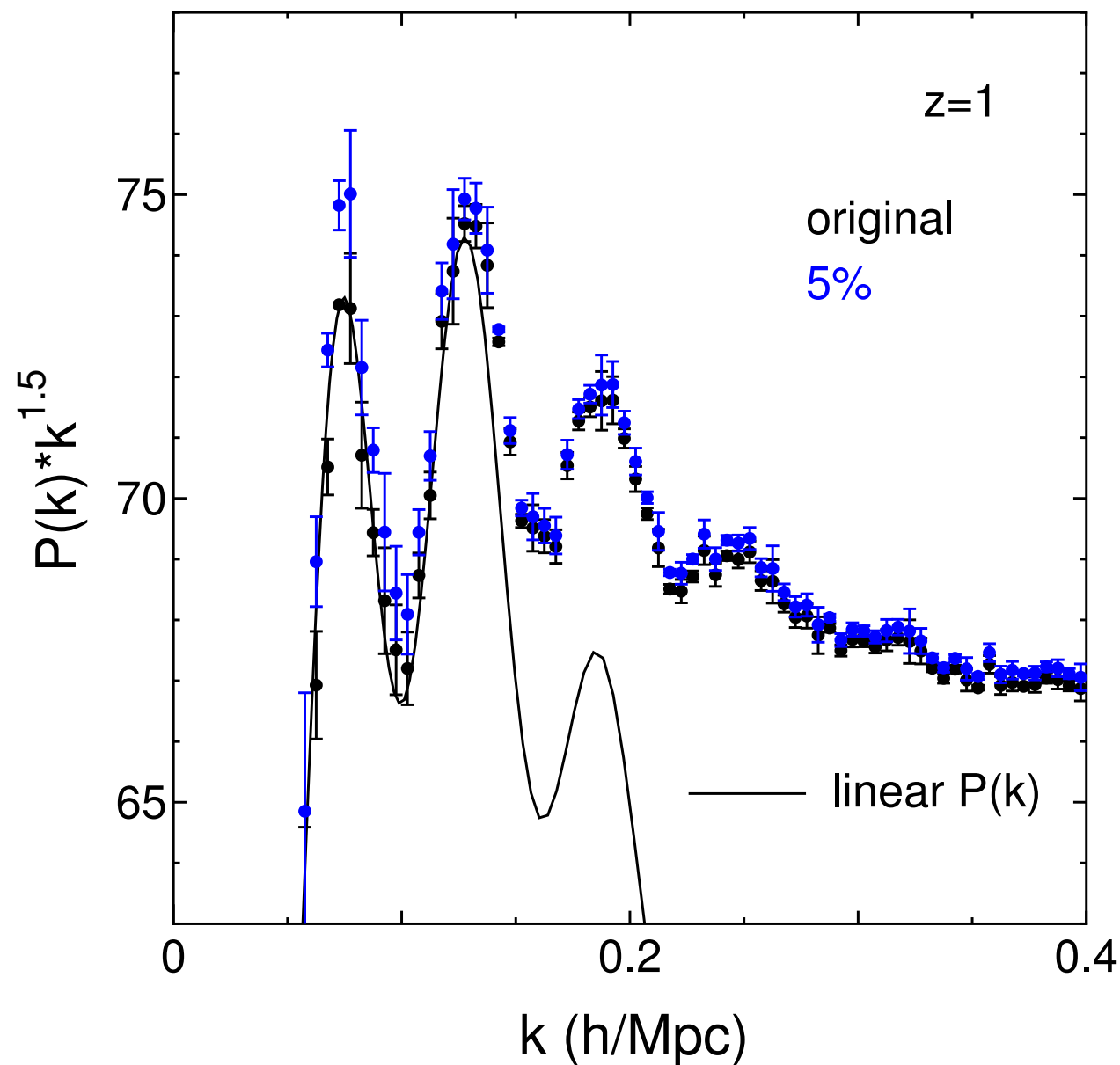
- incorporate into supergravity
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 - $Q \rightarrow Q + i\alpha$
 - $K(Q, Q^*) = K(Q + Q^*) \sim (Q + Q^*)^2 / 2$

$$V = e^K \left((K_i W + W_i)^* K_{\bar{i}}^{-1j} (K_j W + W_j) - 3|W|^2 \right)$$

$$= |W_Q|^2 - 3m_{3/2}(W(Q) + W^*(Q))$$

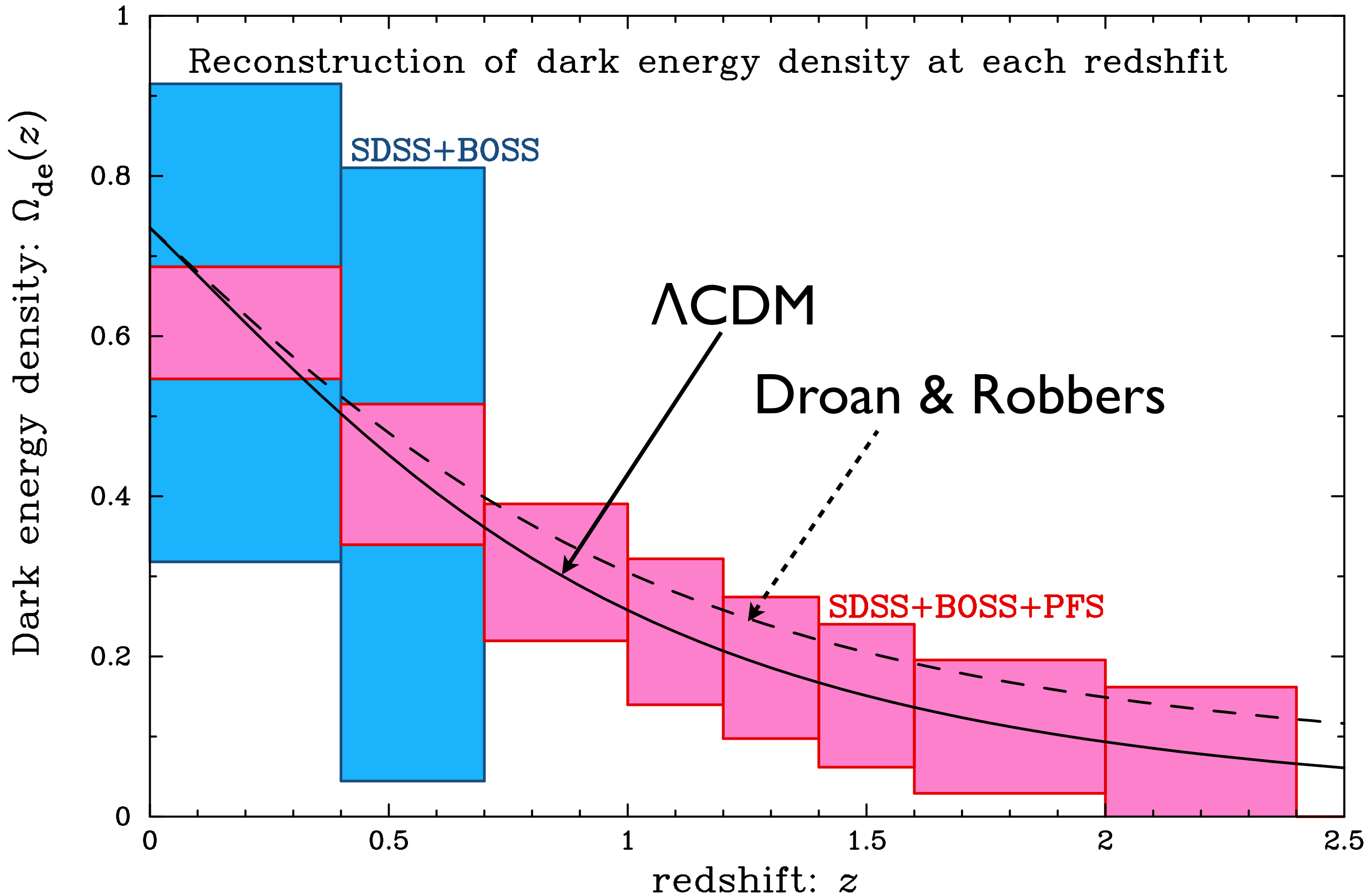
- need $m_{3/2}W(Q) \sim m_{3/2}\Lambda^3 \sim H_0^2$
- **any potential can be lifted to supergravity**
- also radiatively stable $\delta m_Q^2 \sim m_{3/2}^4 \Lambda^6$
- no fifth force

Baryon Acoustic Oscillations



ELGs [OII] $> 8.5\sigma$, 15 min exposure

cosmology $H(z)$ with BAO





PFS Rocks!

Subaru Prime Focus Spectrograph

<http://pfs.ipmu.jp/intro.html>

Conclusion

- We still have *no idea* what dark matter is
- WIMP still the main paradigm
 - but no sign of it
- SIMP can solve problems with DM profile
 - core/cusp, missing satellites
 - very rich phenomenology
 - Exciting dark spectroscopy
 - Can also address the hierarchy problem
- cosmological constant in swampland?
 - Can lift any quintessence potential to supergravity
 - renewed interest in observation